

ON A THEOREM OF W. GUSTIN

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1. W. Gustin⁽¹⁾ has recently shown that any pair of functions harmonic in respective domains of a euclidean space satisfies a certain bilinear integral identity, and then applied it obtaining systematically new proofs of a few fundamental theorems in classical harmonic function theory. His principal theorem may be stated as follows:

"Let ϕ_1 be harmonic in a domain D_1 containing a point (expressed by a vector) q_1 and ϕ_2 be harmonic in D_2 containing q_2 . Then the bilinear integral expression

$$\int_{\Omega} \phi_1(q_1 + p_1 x) \phi_2(q_2 + p_2 x) d\omega_x$$

depends only on the product $p_1 p_2$, provided the closed sphere with radius p_1 about q_1 and the closed sphere with radius p_2 about q_2 are contained in D_1 and D_2 , respectively. Here the integral is taken such that the unit vector x extends over the periphery Ω of the unit sphere with surface element $d\omega_x$, the dimension of the space being arbitrary."

Gustin has given two proofs of the theorem; the first being based on Poisson integral formula and the second on Green's bilinear integral identity. In this Note we shall give a brief proof of which will furthermore clarify the essential nature of the theorem.

2. Now, we may suppose, without loss of generality, that q_1 and q_2 both coincide with the origin, since the harmonicity remains invariant by any translation. As well known⁽²⁾, any function $\phi(p x)$ harmonic in a closed sphere $0 \leq p \leq p^*$ can be expanded in a uniformly convergent series of the form

$$\phi(p x) = \sum_{n=0}^{\infty} p^n Y_n(x) \quad (0 \leq p \leq p^*),$$

$Y_n(x)$ for each n , denoting a spherical surface harmonic of order n . (As to spherical surface harmonics, cf. Remark 2 at the end of the present Note.) Hence we may put

$$\phi_1(p_1 x) = \sum_{n=0}^{\infty} p_1^n Y_n^{(1)}(x)$$

and

$$\phi_2(p_2 x) = \sum_{n=0}^{\infty} p_2^n Y_n^{(2)}(x),$$

where $Y_n^{(1)}(x)$ and $Y_n^{(2)}(x)$ are spherical surface harmonics of order n . Remembering the orthogonality character of spherical surface harmonics

$$\int_{\Omega} Y_m^{(1)}(x) Y_n^{(2)}(x) d\omega_x = 0 \quad (m \neq n),$$

we deduce immediately the relation

$$\int_{\Omega} \phi_1(p_1 x) \phi_2(p_2 x) d\omega_x = \sum_{n=0}^{\infty} (p_1 p_2)^n \int_{\Omega} Y_n^{(1)}(x) Y_n^{(2)}(x) d\omega_x,$$

yielding the desired result.

3. Remark 1. In Gustin's paper the dimension of basic space is assumed to be not less than two. But, if the space is one-dimensional, the bilinear integral expression may be considered to degenerate into the sum

$$\phi_1(q_1 + p_1) \phi_2(q_2 + p_2) + \phi_1(q_1 - p_1) \phi_2(q_2 - p_2).$$

On the other hand, the only harmonic functions in one-dimensional space are linear functions, i.e., of the form

$$\phi(p x) = a p x + b, \quad p \geq 0, x = \pm 1,$$

a and b being constants. It is quite easy to see that the above expression depends on the aggregate $p_1 p_2$ alone for any pair of such linear ϕ_1 and ϕ_2 .

Remark 2. In an N -dimensional euclidean space, the rectangular cartesian and polar coordinates, (ξ_1, \dots, ξ_N) and $(\rho, \vartheta_1, \dots, \vartheta_{N-1})$ are connected in the following manner:

$$\xi_j = \rho \left(\prod_{k=1}^{j-1} \sin \vartheta_k \right) \cos \vartheta_j \quad (1 \leq j \leq N-1),$$

$$\xi_N = \rho \prod_{k=1}^{N-1} \sin \vartheta_k;$$

$$\rho \geq 0, \quad 0 \leq \vartheta_j \leq \pi \quad (1 \leq j \leq N-2), \quad 0 \leq \vartheta_{N-1} \leq 2\pi;$$

the empty product being understood, in the usual way, to denote unity. The square of line element is given by

$$d\sigma^2 = \sum_{j=1}^N d\xi_j^2 = d\rho^2 + \rho^2 \sum_{j=1}^{N-1} \left(\prod_{k=1}^{j-1} \sin^2 \vartheta_k \right) d\vartheta_j^2.$$

On the other hand, by introducing general orthogonal curvilinear coordinates $(\sigma_1, \dots, \sigma_N)$ with $ds^2 = \sum_{j=1}^N g_j d\sigma_j^2$, the Laplacian operator

$$\Delta \equiv \sum_{j=1}^N \frac{\partial^2}{\partial \xi_j^2}$$

is transformed into (*)

$$\Delta = \frac{1}{Jg} \sum_{j=1}^N \frac{\partial}{\partial \sigma_j} \left(\frac{\sqrt{g}}{g_j} \frac{\partial}{\partial \sigma_j} \right), \quad g \equiv \prod_{j=1}^N g_j,$$

which reduces, in our case of polar coordinates, to

$$\Delta = \frac{1}{\rho^N} \frac{\partial}{\partial \rho} \left(\rho^{N-1} \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \Delta^* \equiv \frac{\partial^2}{\partial \rho^2} + \frac{N-1}{\rho} \frac{\partial}{\partial \rho} + \frac{1}{\rho^2} \Delta^*$$

with

$$\Delta^* = \sum_{j=1}^{N-1} \left(\prod_{k=1}^{j-1} \cot^2 \theta_k \right) \cdot \left(\frac{\partial^2}{\partial \theta_j^2} + (N-j+1) \cot \theta_j \frac{\partial}{\partial \theta_j} \right).$$

Hence, for any solid harmonics of the form $\rho^n Y_n(\theta_1, \dots, \theta_{N-1})$, we have

$$0 = \Delta(\rho^n Y_n) = \rho^{n-2} (\mathfrak{n}(\mathfrak{n}+N-2) Y_n + \Delta^* Y_n),$$

i.e.,

$$\Delta^* Y_n + \mathfrak{n}(\mathfrak{n}+N-2) Y_n = 0.$$

The last relation is the self-adjoint partial differential equation for spherical surface harmonics Y_n of order \mathfrak{n} , which belong to the eigen-value $\mathfrak{n}(\mathfrak{n}+N-2)$ (5). In case of N variables, a general homogeneous function (polynomial) of order \mathfrak{n} (with respect to cartesian coordinates) possesses

$\binom{\mathfrak{n}+N-1}{N-1}$ coefficients. Hence, the maximal number of linearly independent Y_n is, in general, equal to

$$\binom{\mathfrak{n}+N-1}{N-1} - \binom{\mathfrak{n}+N-3}{N-1} = 2 \binom{\mathfrak{n}+N-3}{N-2} + \binom{\mathfrak{n}+N-3}{N-3}.$$

(*) Received September 1, 1950.

(1) William Gustin, A bilinear integral identity for harmonic functions. Amer. Journ. of Math. 70(1948), 212-220.

(2) Cf., e.g., R. Courant u. D. Hilbert, Methoden der mathematischen Physik, I. Berlin (1931), p.443, where the completeness of the system is shown for three-dimensional case.

(3) A. Dinghas, Geometrische Anwendungen der Kugelfunktionen. Göttinger Nachr. Neue Folge 1, No.18 (1940), 213-235.

(4) See, for instance, H. Cartan, Leçons sur la géométrie des espaces de Riemann. Paris (1928), pp.48-49; for particular case of $N=3$ see also Courant-Hilbert, loc. cit., pp.194-195.

(5) Courant-Hilbert, loc. cit., pp. 270 and 441.

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