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# A Liouville hyperbolic souvlaki* 

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#### Abstract

We construct a transient bounded-degree graph no transient subgraph of which embeds in any surface of finite genus.

Moreover, we construct a transient, Liouville, bounded-degree, Gromov-hyperbolic graph with trivial hyperbolic boundary that has no transient subtree. This answers a question of Benjamini. This graph also yields a (further) counterexample to a conjecture of Benjamini and Schramm. In an appendix by Gábor Pete and Gourab Ray, our construction is extended to yield a unimodular graph with the above properties.


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## 1 Introduction

A well-known result of Benjamini \& Schramm [6] states that every non-amenable graph contains a non-ambenable tree. This naturally motivates seeking for other properties that imply a subtree with the same property. However, there is a simple example of a transient graph that does not contain a transient tree [6] (such a graph had previously also been obtained by McGuinness [23]). We improve this by constructing -in Section 7a transient bounded-degree graph no transient subgraph of which embeds in any surface of finite genus (even worse, every transient subgraph has the complete graph $K^{r}$ as a minor for every $r$ ). This answers a question of I. Benjamini (private communication).

Given these examples, it is natural to ask for conditions on a transient graph that would imply a transient subtree. In this spirit, Benjamini [4, Open Problem 1.62] asks whether hyperbolicity is such a condition. We answer this in the negative by constructing -in Section 6- a transient hyperbolic (bounded-degree) graph that has no transient subtree. While preparing this manuscript, T. Hutchcroft and A. Nachmias (private communication) provided a simpler example with these properties, which we sketch in

[^0]Section 6.1. In an appendix by Gábor Pete and Gourab Ray, our construction is extended to yield a unimodular graph with the above properties.

A related result of Thomassen states that if a graph satisfies a certain isoperimetric inequality, then it must have a transient subtree [27].

The starting point for this paper was the following problem of Benjamini and Schramm
Conjecture 1.1 ([7, 1.11. Conjecture]). Let $M$ be a connected, transient, Gromovhyperbolic, Riemannian manifold with bounded local geometry, with the property that the union of all bi-infinite geodesics meets every ball of sufficiently large radius. Then $M$ admits non constant bounded harmonic functions. Similarly, a Gromov-hyperbolic bounded valence, transient graph, with $C$-dense bi-infinite geodesics has non constant bounded harmonic functions.

The term $C$-dense here means that every vertex of the graph is at distance at most some constant $C$ from a bi-infinite geodesic. We remark that in order to disprove - the second assertion of - this, it suffices to find a transient, Gromov-hyperbolic bounded valence (aka. degree) graph with the Liouville propertyinfinite geodesics (rather than bi-infinite ones) having the remaining properties ; for given such a graph $G$, one can attach a disjoint 1-way infinite path to each vertex of $G$, to obtain a graph having 1-dense bi-infinite geodesics while preserving all other properties. As pointed out by I. Benjamini (private communication), it is not hard to prove that any 'lattice' in a horoball in 4dimensional hyperbolic space has these properties. We prove that our example also has these properties, thus providing a further counterexample to Conjecture 1.1. A perhaps surpising aspect of our example is that all of its geodesics eventually coincide despite its transience; see Section 2.

In Section 2.1 we provide a sketch of this construction, from which the expert reader might be able to deduce the details.

Although we do not formally provide a counterexample to the first assertion of Conjecture 1.1, we believe it is easy to obtain one by blowing up the edges of our graph into tubes.

## 2 The hyperbolic Souvlaki

In this section we construct a bounded-degree graph $\Psi$ with the following properties

1. it is hyperbolic, and its hyperbolic boundary consists of a single point;
2. for every vertex $x$ of $\Psi$, there is a unique infinite geodesic starting at $x$, and any two 1-way infinite geodesics of $\Psi$ eventually coincide;
3. it is transient;
4. every subtree of $\Psi$ is recurrent;
5. it has the Liouville property.

This graph thus yields a counterexample to [4, Open Problem 1.62] and Conjecture 1.1 as mentioned in the Introduction.

### 2.1 Sketch of construction

Let us sketch the construction of this graph $\Psi$, and outline the reasons why it has the above properties. It consists of an 1-way infinite path $S=s_{0} s_{1} \ldots$, on which we glue a sequence $M_{i}$ of finite increasing subgraphs of an infinite '3-dimensional' hyperbolic graph $H_{3}$. For example, $H_{3}$ could be the 1-skeleton of a regular tiling of 3-dimensional hyperbolic space, and the $M_{i}$ could be taken to be copies of balls of increasing radii around some origin in $H_{3}$, although it was more convenient for our proofs to construct different $H_{3}$ and $M_{i}$.

In order to glue $M_{i}$ on $S$, we identify the subpath $s_{2^{i}} \ldots s_{2^{i+2}-1}$ with a geodesic of the same length in $M_{i}$. Thus $M_{i}$ intersects $M_{i-1}$ and $M_{i+1}$ but no other $M_{j}$, and this intersection is a subpath of $S$; see Figure 4. (Our graph can be quasi-isometrically embedded in $\mathbb{H}^{5}$, but probably not in $\mathbb{H}^{4}$.) We call this graph a hyperbolic souvlaki, with skewer $S$ and meatballs $M_{i}$. We detail its construction in Section 2.

To prove that this graph is transient, we construct a flow of finite energy from $s_{0}$ to infinity (Section 4). This flow carries a current of strength $2^{-i}$ inside $M_{i}$ out of each vertex in $s_{2^{i}} \ldots s_{2^{i+1}-1}$, and distributes it evenly to the vertices in $s_{2^{i+1}} \ldots s_{2^{i+2}}$ for every $i$. These currents can be thought of as flowing on spheres of varying radii inside $M_{i}$, avoiding each other, and it was important to have at least three dimensions for this to be possible while keeping the energy dissipated under control.

To prove that our graph has the Liouville property, we observe that random walk has to visit $S$ infinitely often, and has enough time to 'mix' inside the $M_{i}$ between subsequent visits to $S$ (Section 5).

### 2.2 Formal construction

We now explain our precise construction, which is similar but not identical to the above sketch. We start by constructing a hyperbolic graph $H_{3}$ which we will use as a model for the 'meatballs' $M_{i}$; more precisely, the $M_{i}$ will be chosen to be increasing subgraphs of $H_{3}$.

Let $T_{3}$ denote the infinite tree with one vertex $r$, which we call the root, of degree 3 and all other vertices of degree 4 . For $n=1,2, \ldots$, we put a cycle -of length $3^{n}$ - on the vertices of $T_{3}$ that are at distance $n$ from $r$ in such a way that the resulting graph is planar ${ }^{1}$; see Figure 1. We denote this graph by $H_{2}$. It is not hard to see that $H_{2}$ is hyperbolic, for instance by checking that any two infinite geodesics starting at $r$ either stay at bounded distance or diverge exponentially, and using [26, Section 2.20].


Figure 1: The ball of radius 3 around the root of $\mathrm{H}_{2}$.

Recall that a ray is a 1-way infinite path. We will now turn $H_{2}$ into a '3-dimensional' hyperbolic graph $H_{3}$, in such a way that each ray inside $T_{3}$ (or $H_{2}$ ) starting at $r$ gives rise to a subgraph of $H_{3}$ isomorphic to the graph $W$ of Figure 2, which is a subgraph of the Cayley graph of the Baumslag-Solitar group $B S(1,2)$. Formally, we construct $W$ from

[^1]infinitely many vertex disjoint double rays ${ }^{2} D_{0}, D_{1}, D_{2}, .$. , where $D_{i}=\ldots r_{i}^{-2} r_{i}^{-1} r_{i}^{0} r_{i}^{1} r_{i}^{2} \ldots$ Then we add all edges of the form $r_{i}^{k} r_{i+1}^{2 k}$.


Figure 2: The graph $W$ : a subgraph of the standard Cayley graph of the Baumslag-Solitar group $B S(1,2)$. It is a plane hyperbolic graph.

To define $H_{3}$, we let the height $h(t)$ of a vertex $t \in V\left(H_{2}\right)$ be its distance $d(r, t)$ from the root $r$. For a vertex $w$ of $W$, we say that its height $h(w)$ is $n$ if $w$ lies in $D_{n}$, the $n$th horizontal double ray in Figure 2.

We define the vertex set of $H_{3}$ to consist of all ordered pairs $(t, w)$ where $t$ is a vertex of $H_{2}$ and $w$ is a vertex of $W$ and $h(w)=h(t)$. The edge set of $H_{3}$ consists of all pairs of pairs $(t, w)\left(t^{\prime}, w^{\prime}\right)$ such that either

- $t t^{\prime} \in E\left(H_{2}\right)$ and $w w^{\prime} \in E(W)$, or
- $t t^{\prime} \in E\left(H_{2}\right)$ and $w=w^{\prime}$, or
- $t=t^{\prime}$ and $w w^{\prime} \in E(W)$.


Figure 3: A subgraph of $H_{3}$. Edges of the form $(t, w)\left(t^{\prime}, w^{\prime}\right)$ with $t=t^{\prime}$ and $w w^{\prime} \in E(W)$ are missing from the figure: these are all the edges joining corresponding vertices in consecutive components of the figure.

Thus every vertex $t$ of $H_{2}$ gives rise to a double ray in $H_{3}$, which consists of those vertices of $H_{3}$ that have $t$ as their first coordinate. Similarly, every vertex $w$ of $W$ gives rise to a cycle in $H_{3}$, the length of which depends on $h(w)$. We call the vertices on any such cycle cocircular. Every ray of $T_{2}$ starting at $r$ gives rise to a copy of $W$, and if two such paths share their first $k$ vertices, then the corresponding copies of $W$ share their

[^2]first $k$ levels of $h$. It is not hard to prove that $H_{3}$ is a hyperbolic graph, but we will omit the proof as we will not use this fact.

We next construct $\Psi$ by glueing a sequence of finite subgraphs $M_{n}$ of $H_{3}$ along a ray $S$. We could choose the subgraph $M_{n}$ to be a ball in $H_{3}$, but we found it more convenient to work with somewhat different subgraphs of $H_{3}$ : we let $M_{n}$ be the finite subgraph of $H_{3}$ spanned by those vertices $(t, w)$ such that $w$ lies in a certain box $B_{n} \subseteq W$ of $W$ defined as follows. Consider a subpath $P_{n}$ of the bottom double-ray of $W$ of length $3 \cdot 2^{n}$, and let $B_{n}$ consist of those vertices $w$ that lie in or above $P_{n}$ (as drawn in Figure 2) and satisfy $h(w) \leq 2^{n+1}$.

This completes the definition of $M_{n}$. We let $S_{n}$ denote the vertices of $M_{n}$ corresponding to $P_{n}$, and we index the vertices of $S_{n}$ as $\left\{r(x), 0 \leq x \leq 3 \cdot 2^{n}\right\}$. Note that $S_{n}$ is a geodesic of $M_{n}$. We subdivide $S_{n}$ into three parts: $L_{n}:=\left\{r(x), 0 \leq x<2^{n}\right\}, m_{n}:=r\left(2^{n}\right)$ and $R_{n}:=\left\{r(x), 2^{n}<x \leq 3 \cdot 2^{n}\right\}$. We define the ceiling $F_{n}$ of $M_{n}$ to be its vertices of maximum height, i.e. the vertices $(t, w) \in V\left(M_{n}\right)$ with $h(w)=2^{n+1}$.

Finally, it remains to describe how to glue the $M_{n}$ together to form $\Psi$. We start with a ray $S$, the first vertex of which we denote by $o$ and call the root of $\Psi$. We glue $M_{1}$ on $S$ by identifying $S_{1}$ with the initial subpath of $S$ of length $\left|S_{1}\right|$. Then, for $n=2,3, \ldots$, we glue $M_{n}$ on $S$ in such a way that $L_{n}$ is identified with $R_{n-1}$ (where we used the fact that $\left|L_{n}\right|=\left|R_{n-1}\right|=2^{n}$ by construction), $m_{n}$ is identified with the following vertex of $S$, and $R_{n}$ is identified with the subpath of $S$ following that vertex and having length $\left|R_{n}\right|=2^{n+1}$. Of course, we perform this identification in such a way that the linear orderings of $L_{n}$ and $R_{n}$ are given by the induced linear ordering of $S$. We let $\Psi$ denote the resulting graph. We think of $M_{n}$ as a subgraph of $\Psi$.

### 2.3 Properties of $\Psi$

By construction, for $j>i$ we have $M_{i} \cap M_{j}=\emptyset$ unless $j=i+1$, in which case $M_{i} \cap M_{j}=R_{i}=L_{j} \subset S$. The following fact is easy to see.

For every $n, R_{n}$ separates $L_{n}$ (and $o$ ) from infinity.
The following assertion will be important for the proof of the Liouville property.
There is a uniform lower bound $p>0$ for the probability $\mathbb{P}_{v}\left[\tau_{F_{n}}<\tau_{S_{n}}\right]$ that random walk in $\Psi$ from any vertex of $L_{n}$ will visit the ceiling $F_{n}$ before returning to $S_{n}$.
Indeed, we can let $p$ be the probability for random walk on $H_{2}$ starting at the root $o$ to never visit $o$ again; this is positive because $H_{2}$ is transient. Then (2.2) holds because in a random walk from $S_{n}$ on $M_{n}$, any steps inside the copies of $H_{2}$ behave like random walk on $H_{2}$ until hitting $F_{n}$, and the steps 'parallel' to $S_{n}$ do not have any influence.

## 3 Hyperbolicity

In this section we prove that $\Psi$ is hyperbolic in the sense of Gromov [16].
Lemma 3.1. The graph $\Psi$ is hyperbolic, and has a one-point hyperbolic boundary.
Proof. We claim that for every vertex $x \in V(\Psi)$, there is a unique 1-way infinite geodesic starting at $x$. Indeed, this geodesic $x_{0} x_{1} \ldots$, takes a step from $x_{i}$ towards the root of $T_{3}$ inside the copy of $H_{2}$ corresponding $x_{i}$ whenever such an edge exists in $\Psi$, and it takes a horizontal step in the direction of infinity whenever such an edge does not exist. To see that $\gamma$ is the unique infinite geodesic starting at $x$, suppose there is a second such geodesic $\delta$. Clearly, $\delta$ has infinitely many vertices on the skewer $S$ as all components of $\Psi \backslash S$ are finite. In fact, it is not hard to see that $\delta$ eventually coincides with $S$ as the latter contains the unique geodesic between any two of its vertices. Thus $\gamma$ and $\delta$ meet,
and we can let $y$ be their first common vertex. Now consider their subpaths $x \gamma y$ and $x \delta y$ from $x$ to $y$. Note that $\Psi$ has two types of edges: those that lie in a copy of $H_{2}$, and horizontal ones. It is easy to see that any $x-y$ path must have at least as many edges of each type as $x \gamma y$. Moreover, by considering the first edge $e$ at which $x \delta y$ deviates from $x \gamma y$, it is not hard to check that $x \delta y$ has more edges of the same type as $e$ as $x \gamma y$, which leads to a contradiction.

The hyperbolicity of $\Psi$ now follows from a well-known fact saying that a space is hyperbolic if and only if any two geodesics with a common starting point are either at bounded distance or diverge exponentially in a certain sense; see [26, Section 2.20]. We skip the details here as in our case the condition is trivially satisfied due to the above claim-namely, any two geodesics from a given point are at bounded distance since they coincide.

As all infinite geodesics eventually coincide with $S$, we also immediately have that the hyperbolic boundary of $G$ consists of just one point.

## 4 Transience

In this section we prove that $\Psi$ is transient. We do so by displaying a flow from $o$ to infinity having finite Dirichlet energy; transience then follows from Lyons' criterion:
Theorem 4.1 (T. Lyons' criterion (see [20] or [21])). A graph $G$ is transient, if and only if $G$ admits a flow of finite energy from a vertex to infinity.

We refer the reader to [21] or [15] for the basics of electrical networks needed to understand this theorem.


Figure 4: The structure of the graph $\Psi$, with the 'balls' intersecting along the ray and the flow inside the ball.

To construct this flow $f$, we start with the flow $t$ on the tree $T_{3} \subset H_{2}$ which sends the amount $3^{-n}$ through each directed edge of $T_{3}$ from a vertex of distance $n-1$ from the root to a vertex of distance $n$ from the root. Note that $t$ has finite Dirichlet energy.

Our flow $f$ will be as described in the introduction, that is, it is composed of flows $g(n)$ in $M_{n}$. These flows flow from $L_{n}$ to $R_{n}$. The flow $g(n)$ in turn is composed of 'atomic'
flows, one for each $v \in L_{n}$. Roughly, these atomic flows imitate $t$ from above for some levels, then use the edges parallel to $S_{n}$ to bring it 'above' $R_{n}$, and then collect it back to (two vertices of) $S_{n}$ imitating $t$ in the inverse direction. A key idea here is that although the energy dissipated along the long paths parallel to $S_{n}$ is proportional to their length, by going up enough levels with the $t$-part of these flows, we can ensure that the flow $i$ carried by each such path is very small compared to its length $\ell$. Thus its contribution $i^{2} \ell$ to the Dirichlet energy can be controlled: although going up one level doubles $\ell$, and triples the number of long paths we have, each of them now carries $1 / 3$ of the flow, and so its contribution to the energy is multiplied by a factor of $1 / 9$. Thus all in all, we save a factor of $6 / 9$ by going up one more level - and we have made the $M_{i}$ high enough that we can go up enough levels.

We now describe $g(n)$ precisely. For every $n \in \mathbb{N}$, let us first enumerate the vertices of $L_{n}$ as $l^{j}=l_{n}^{j}$, with $j$ ranging from 1 to $\left|L_{n}\right|=2^{n}$, in the order they appear on $S_{n}$ as we move from the midpoint $m_{n}$ towards the root $o$. Likewise, we enumerate the vertices of $R_{n}$ as $r^{j}=r_{n}^{j}$, with $j$ ranging from 1 to $\left|R_{n}\right|=2\left|L_{n}\right|$, in the order they appear on $S_{n}$ as we move from the midpoint $m_{n}$ towards infinity. Thus $r^{1}, l^{1}$ are the two neighbours of $m_{n}$ on $S$. We will let $g(n)$ be the union of $\left|L_{n}\right|$ subflows $g^{j}=g_{n}^{j}$, where $g^{j}$ flows from $l^{j}$ into $r^{2 j}$ and $r^{2 j-1}$. More precisely, $g^{j}$ sends $1 /\left|L_{n}\right|=2^{-n}$ units of current out of $l^{j}$, and half as many units of current into each of $r^{2 j}$ and $r^{2 j-1}$.

We define $g^{j}$ as follows. In the copy of $H_{2}$ containing the source $l^{j}$ of $g^{j}$, we multiply the flow $t$ from above by the factor $2^{-n}$, and truncate it after $j$ layers; we call this the out-part of $g^{j}$. Then, from each endpoint $x$ of that flow, we send the amount of flow that $x$ receives from $l^{j}$, which equals $2^{-n} 3^{-j}$, along the horizontal path $P_{x}$ joining $x$ to the copy $C_{1}$ of $H_{2}$ containing $r^{2 j-1}$. We let half of that flow continue horizontally to reach the copy $C_{2}$ of $H_{2}$ containing $r^{2 j}$; call this the middle-part of $g^{j}$. Finally, inside each of $C_{1}, C_{2}$, we put a copy of the out-part of $g^{j}$ multiplied by $1 / 2$ and with directions inverted; this is called the in-part of $g^{j}$. Note that the union of these three parts is a flow of intensity $2^{-n}$ from $l^{j}$ to $r^{2 j}$ and $r^{2 j-1}$, each of the latter receiving $2^{-n-1}$ units of current.

Let us calculate the energy $E\left(g^{j}\right)$. The contribution to $E\left(g^{j}\right)$ by its out-part is bounded above by $2^{-2 n} E(t)$ because that part is contained in the flow $2^{-n} t$. Similarly, the contribution of the in-part is half of the contribution of the out-part. The contribution of the middle-part is $3^{j} \cdot(2 j+1) 2^{j} \cdot\left(2^{-n} 3^{-j}\right)^{2}$ : the factor $3^{j}$ counts the number of horizontal paths used by the flow, each of which has length $(2 j+1) 2^{j}$, and carries $2^{-n} 3^{-j}$ units of current (except for its last $2^{j}$ edges, from $C_{1}$ to $C_{2}$, which carry half as much, but we can afford to be generous). Note that this expression equals $2^{-2 n}(2 j+1)(6 / 9)^{j}$, which is upper bounded by $k 2^{-2 n}$ for some constant $k$.

Adding up these contributions, we see that $E\left(g^{j}\right) \leq K 2^{-2 n}$ for some constant $K$ (which depends on neither $n$ nor $j$ ).

Now let $g(n)$ be the union of the $2^{n}$ flows $g^{j}$. Note that $g^{j}, g^{i}$ are disjoint for $i \neq j$, and therefore the energy $E(g(n))$ of $g$ is just the sum $\sum_{j<2^{n}} E\left(g^{j}\right)$. By the above bound, this yields $E(g(n)) \leq K 2^{-n}$.

Now let $f=\bigcup_{n \in \mathbb{N}} g(n)$ be the union of all the flows $g(n)$. Then $g(n), g(m)$ are disjoint for $n \neq m$, because they are in different $M_{i}^{\prime} s$. Thus $E(f)=\sum_{n} E(g(n)) \leq K$ is finite. Since $g(n)$ removes as much current from each vertex of $L_{n}$ as $g(n-1)$ inputs, $f$ is a flow from $o$ to infinity. Hence $\Psi$ is transient by Lyons' criterion (Theorem 4.1).

## 5 Liouville property

In this section we prove that $\Psi$ is Liouville, i.e. it admits no bounded non-constant harmonic functions.

We remark that a well-known theorem of Ancona [2] states that in any non-amenable hyperbolic graph the hyperbolic boundary coincides with the Martin boundary. We cannot apply this fact to our case in order to deduce the Liouville property from the fact that our hyperbolic boundary is trivial, because our graph turns out to be amenable.

We will use some elementary facts about harmonic functions that can be found e.g. in [14].

Let $h$ be a bounded non-constant harmonic functions on a graph $G$. We may assume that the range of $h$ is contained in $[0,1]$. Recall that, by the bounded martingale convergence theorem, if $\left(X_{n}\right)_{n \in \mathbb{N}}$ is a simple random walk on $G$, then $h\left(X_{n}\right)$ converges almost surely. We call such a function $h$ sharp, if this $\operatorname{limit}^{\lim _{n}} h\left(X_{n}\right)$ is either 0 or 1 almost surely. It is well-known that if a graph admits a bounded non-constant harmonic function, then it admits a sharp harmonic function, see [14, Section 4].

So let us assume from now on that $h: V(\Psi) \rightarrow[0,1]$ is a sharp bounded harmonic function on $\Psi$.

We first recall some basic facts from [14, Section 7]; we repeat some of the proofs for the convenience of the reader.
Lemma 5.1. If $h$ is a sharp harmonic function, then $h(z)=\mathbb{P}_{z}\left[\lim h\left(Z_{n}\right)=1\right]$ for every vertex $z$, where $Z_{n}$ denotes a random walk from $z$.


Figure 5: The path $P_{a}$ in the proof of the Liouville property.

Lemma 5.2. If $h$ is a sharp harmonic function that is not constant, then for every $\epsilon>0$ there are $a, z \in V$ with $h(a)<\epsilon$ and $h(z)>1-\epsilon$.

Let $\mathcal{A}$ be a shift-invariant event of our random walk, i.e. an event not depending on the first $n$ steps for every $n$. (The probability space we work with here is the space of 1 -way infinite walks, endowed with the natural probability measure induced by simple
random walk. The only kind of event we will later consider is the event $1^{s}$ that $s\left(Z_{n}\right)$ converges to 1 , where $s$ is our fixed sharp harmonic function.) By an event here we mean a measurable subset of the space of 1-way infinite walks in our fixed graph; the starting vertex of the walks is not fixed, it can be an arbitrary vertex of our graph. A walk here is a sequence $v_{0}, v_{1}, \ldots$ of vertices of the graph such that $v_{i}$ is joined with an edge to $v_{i+1}$ for every $i \geq 0$. As usual, we consider the $\sigma$-algebra generated by sets of walks that start with a fixed finite sequence of steps and are arbitrary after those steps.

For $r \in(0,1 / 2]$, let

$$
\begin{aligned}
& A_{r}:=\left\{v \in V \mid \mathbb{P}_{v}[\mathcal{A}]>1-r\right\} \text { and } \\
& Z_{r}:=\left\{v \in V \mid \mathbb{P}_{v}[\mathcal{A}]<r\right\},
\end{aligned}
$$

where $\mathbb{P}_{v}[\cdot]$ denotes the law of random walk from a vertex $v$. Note that $A_{r} \cap Z_{r}=\emptyset$ for every such $r$.

By Lemma 5.1, if we let $\mathcal{A}:=1^{s}$ then we have $A_{r}=\{v \in V \mid s(v)>1-r\}$ and $Z_{r}=\{v \in V \mid s(v)<r\}$.
Lemma 5.3. For every $\epsilon, \delta \in(0,1 / 2]$, and every $v \in A_{\epsilon}$, we have
$\mathbb{P}_{v}\left[\right.$ visit $\left.V \backslash A_{\delta}\right]<\epsilon / \delta$. Similarly, for every $v \in Z_{\epsilon}$, we have
$\mathbb{P}_{v}\left[\right.$ visit $\left.V \backslash Z_{\delta}\right]<\epsilon / \delta$.
Proof. Start a random walk $\left(Z_{n}\right)$ at $v$, and consider a stopping time $\tau$ at the first visit to $V \backslash A_{\delta}$. If $\tau$ is finite, let $z=Z_{\tau}$ be the first vertex of random walk outside $A_{\delta}$. Since $z \notin A_{\delta}$, the probability that $s\left(X_{n}\right)$ converges to 0 for a random walk $\left(X_{n}\right)$ starting from $z$ is at least $\delta$ by the definition of $A_{\delta}$. Thus, conditioning on ever visiting $V \backslash A_{\delta}$, the event $\mathcal{A}$ fails with probability at least $\delta$ since $\mathcal{A}$ is a shift-invariant event and our random walk has the Markov property. But $\mathcal{A}$ fails with probability less than $\epsilon$ because $v \in A_{\epsilon}$, and so $\mathbb{P}_{v}\left[\right.$ visit $\left.V \backslash A_{\delta}\right]<\epsilon / \delta$ as claimed.

The second assertion follows by the same arguments applied to the complement of $\mathcal{A}$.

Corollary 5.4. If a random walk from $v \in A_{\epsilon}$ (respectively, $v \in Z_{\epsilon}$ ) visits a set $W \subset V$ with probability at least $\kappa$, then there is a $v-W$ path all vertices of which lie in $A_{\epsilon / \kappa}$ (resp. $Z_{\epsilon / \kappa}$ ).

Proof. Apply Lemma 5.3 with $\delta=\epsilon / \kappa$. Then the probability that random walk always stays within $A_{\epsilon / \kappa}$ is larger than $1-\kappa$. Hence there is a nonzero probability that random walk meets $W$ and along its trace only has vertices from $A_{\epsilon / \kappa}$.

Easily, $h$ is uniquely determined by its values on the skewer $S$. Indeed, for every other vertex $v$, note that random walk $X_{n}$ from $v$ visits $S$ almost surely, and so $h(v)=\mathbb{E} h\left(X_{\tau(R)}\right)$, where $\tau(S)$ denotes the first hitting time of $S$ by $X_{n}$. The same argument implies that
$h$ is radially symmetric, i.e. for every two cocircular vertices $v, w$, we have
$h(v)=h(w)$.
Indeed, this follows from the fact that cocircular vertices have the same hitting distribution to $S$, which is easy to see (for any vertex on a circle, random walk has the same probability to move to some other circle).

We claim that, given an arbitrarily small $\epsilon>0$, all but finitely many of the $L_{n}$ contain a vertex in $A_{\epsilon}$.

Indeed if not, then since random walk from $o$ has to visit all $L_{n}$ by transience and (2.1) (where we use the fact that $L_{n}=R_{n-1}$ ), we would have $\mathbb{P}\left[\lim h\left(X_{n}\right)=1\right]=0$ for random walk $X_{n}$ from $o$ by Lemma 5.1, because if our random walk visits infinitely many vertices $y$ such that $h(y)<1-\epsilon$ then $h\left(X_{n}\right)$ cannot converge to 1 . But that probability is
equal to $h(o)$ by Lemma 5.1, and if it is zero, then using Lemma 5.1 again easily implies that $h$ is identically zero, contrary to our assumption that it is not constant.

Similarly, all but finitely many of the $L_{n}$ contain a vertex in $Z_{\epsilon}$, because as $h$ is sharp, $h\left(X_{n}\right)$ must converge to either 0 or 1 . Thus we can find a late enough $M_{n}$ such that $L_{n}$ contains a vertex $a \in A_{\epsilon}$ as well as a vertex $z \in Z_{\epsilon}$. We assume that $a$ and $z$ are the last vertices of $L_{n}$ (in the ordering of $L_{n}$ induced by the well-ordering of $S$ ) that are in $A_{\epsilon}$ and $Z_{\epsilon}$ respectively. Assume without loss of generality that $a$ appears before $z$ in the ordering of $L_{n}$.

Note that, since $R_{n}$ separates $a$ from infinity (2.1), random walk from $a$ visits $R_{n}$ almost surely. Thus we can apply Corollary 5.4 with $W:=R_{n}$ and $\kappa=1$ to obtain an $a-R_{n}$ path $P_{a}$ with all its vertices in $A_{\epsilon}$. We may assume that $P_{a} \subset M_{n}$ by taking a subpath contained in $M_{n}$ if needed. Indeed, $P_{a}$ can meet $L_{n}$ only in vertices that are not past $a$ in the linear ordering of $L_{n}$.

Let $O_{a}$ denote the set of vertices $\left\{x=(t, w) \in M_{n} \mid\right.$ there is $\left(t^{\prime}, w^{\prime}\right) \in V\left(P_{a}\right)$ with $w^{\prime}=$ $w\}$ obtained by 'rotating' $P_{a}$ around $S$. By (5.1), we have $O_{a} \subset A_{\epsilon}$ since $P_{a} \subset A_{\epsilon}$. Note that $O_{a}$ separates $z$ from the ceiling $F_{n}$ of $M_{n}$. But as random walk from $z \in Z_{\epsilon}$ visits $F_{n}$ before returning to $S$ with probability uniformly bounded below by (2.2), we obtain a contradiction to Lemma 5.3 with $\delta=1 / 2$ for $\epsilon$ small enough compared to that bound.

## 6 A transient hyperbolic graph with no transient subtree

In this section we explain how our souvlaki construction can be slightly modified so that it does not contain any transient subtrees but remains transient and hyperbolic (and Liouville). This answers a question of I. Benjamini (private communication). The question is motivated by the fact that it is not too easy to come up with transient graphs that do not have transient subtrees [6].

We start with a very fast growing function $f: \mathbb{N} \rightarrow \mathbb{N}$, whose precise definition we reveal at the end of the proof. Roughly speaking, we will attach a sequence of finite graphs $\left(M_{f(n)}\right)_{n \in \mathbb{N}}$ similar to the 'meatballs' from above to a ray $S$ (the 'skewer') in such a way that most of the intersection of $S$ with a fixed meatball is not contained in any other meatball. Formally, we let $P_{m}$ be the 'bottom path' of $M_{m}$ as defined in Section 2, and we tripartition $P_{f(n)}$ as follows: Let $L_{n}$ consist of the first $2^{n}$ vertices on $P_{f(n)}$, and $R_{n}$ consist of its last $2^{n+1}$ vertices. The set of the remaining vertices of $P_{f(n)}$ we denote by $Z_{n}$, which by our choice of $f$ will be much larger than $R_{n}$. As before, we glue the $M_{f(n)}$ on $S$ by identifying $P_{f(n)}$ with a subpath of $S$. We start by glueing $M_{f(1)}$ on the initial segment of $S$ of the appropriate length. Then we recursively glue the other $M_{f(n)}$ in such a way that $L_{n}$ is identified with $R_{n-1}$. We call the resulting graph $\bar{\Psi}$.
Theorem 6.1. $\bar{\Psi}$ is a bounded degree transient gromov-hyperbolic graph that does not contain a transient subtree.

Proof. The hyperbolicity of $\bar{\Psi}$ can be proved by the arguments we used for the original souvlaki $\Psi$. Also $\bar{\Psi}$ is transient by an argument analogue given to the one for $\Psi$ : the obvious analogue of the flow $f$ described in Section 4 is in $M_{f(n)}$ a flow of intensity one from $L_{n}$ to $R_{n}$ of energy at most constant times $2^{-n}$. The computation is analogue to the one given above. ${ }^{3}$ So it remains to show that $\bar{\Psi}$ does not have a transient subtree.

Let $T$ be any subtree of $\bar{\Psi}$. We want to prove that $T$ is not transient. Easily, we may assume that $T$ does not have any degree 1 vertices. We will show that the following quotient $Q$ of $T$ is not transient: for each $n$, we identify all vertices in $L_{n}$ to a new vertex $v_{n}$.

[^3]Note that the vertices $v_{n}$ and $v_{n+1}$ are cut-vertices of $Q$; let $Q_{n}$ be the union of those components of $Q-v_{n}-v_{n+1}$ that send edges to both vertices $v_{n}$ and $v_{n+1}$. We will show that in $Q_{n}$ the effective resistance from $v_{n}$ to $v_{n+1}$ is bounded away from 0 , from which the recurrence of $T$ will follow using Lyons' criterion.

Let $d=\left|L_{n+1}\right|$. We claim that there is some constant $c=c(d)$ only depending on $d$ such that there are at most $c$ vertices of $Q_{n}$ with a degree greater than 2: indeed, $Q_{n} \backslash\left\{v_{n}, v_{n+1}\right\}$ is a forest with at most $d\left(v_{n}\right)+d\left(v_{n+1}\right)$ leaves. Since these degrees are bounded also the number of leaves is bounded. Hence all but boundedly many vertices of $Q_{n}$ have degree two.

Next, we observe that $Q_{n}$ has maximum degree at most $d$. Furthermore, the distance between $v_{n}$ and $v_{n+1}$ in $Q_{n}$ is at least $Z_{n}$, which -by the choice of $f$ - is huge compared to $d$ and so also compared to $c$. Hence it remains to prove the following:
Lemma 6.2. For every constant $C$ and every $m$ there is some $s=s(m, C)$, such that for every finite graph $K$ with maximum degree at most $C$ and at most $C$ vertices of degree greater than 2, and for any two vertices $x$, $y$ of $K$ with distance at least $s$, the effective resistance between $x$ and $y$ in $K$ is at least $m$.

Proof. We start with a large natural number $R$ the value of which we reveal later, and set $s=R \cdot C$.

Let $K^{\prime}$ be the graph obtained from $K$ by suppressing all vertices of degree 2; suppressing a vertex $x$ of degree 2 means replacing $x$ and its two incident edges with a single edge between the neighbours of $x$. The length of an edge of $K^{\prime}$ is the number of times it is subdivided in $K$. Let $N^{\prime}$ be the electrical network with underlying graph $K^{\prime}$, where the resistance of an edge of $K^{\prime}$ is its length. Clearly, the effective resistance between $x$ and $y$ in the graph $K$ is equal to the effective resistance between $x$ and $y$ in the network $K^{\prime}$. Hence it suffices to show that the effective resistance between $x$ and $y$ in $K^{\prime}$ is at least $m$.

We colour an edge of $K^{\prime}$ black if it has length at least $R$. Note that $K^{\prime}$ has at most $C$ vertices. Thus every $x$ - $y$-path in $K^{\prime}$ has length at most $C$, but in $K$ any such path has length at least $s$. Therefore each $x-y$-path in $K^{\prime}$ contains a black edge. Hence in $K^{\prime}$ there is an $x$ - $y$-cut consisting of black edges only. This cut has at most $C^{2}$ edges. Thus by Rayleigh's monotonicity law [21] the effective resistance in $K^{\prime}$ between $x$ and $y$ is at least the one of that cut, which is as large as we want: indeed, we can pick $R$ so large that the latter resistance exceeds $m$.

Now we reveal how large we have picked $f(n)$ : recall that $d=2^{n+1}$ and that $\left|Z_{n}\right|=f(n)-3 \cdot 2^{n}$. We pick $f(n)$ large enough that $\left|Z_{n}\right| \geq s(1, \max (c(d), d))$, where $s$ is as given by the last lemma. With these choices the effective resistance between $v_{n}$ and $v_{n+1}$ in $Q_{n}$ is at least 1. So $Q$ cannot be transient by Lyons' criterion (Theorem 4.1) as the $Q_{n}$ are disjoint and any flow to infinity has to traverse all but finitely many of them with a constant intensity. By Rayleigh's monotonicity law [21], $T$ is recurrent too.

### 6.1 Another transient hyperbolic graph with no transient subtree

We now sketch another construction of a transient hyperbolic graph with no transient subtree, provided by Tom Hutchcroft and Asaf Nachmias (private communication).

Let $[0,1]^{3}$ be the unit cube. For each $n \geq 0$, let $D_{n}$ be the set of closed dyadic subcubes of length $2^{-n}$. For each $n \geq 0$, let $G_{n}$ be the graph with vertex set $\bigcup_{i=0}^{n} D_{i}$, and where two cubes $x$ and $y$ are adjacent if and only if

- $x \supset y, x \in D_{i}$ and $y \in D_{i+1}$ for some $i \in\{0, \ldots n-1\}$,
- $y \supset x, y \in D_{i}$ and $x \in D_{i+1}$ for some $i \in\{0, \ldots n-1\}$, or
- $x, y \in D_{i}$ for some $i \in\{0, \ldots, n\}$ and $x \cap y$ is a square.

Then the graphs $G_{n}$ are uniformly Gromov hyperbolic and, since the subgraph of $G_{n}$ induced by $D_{n}$ is a cube in $\mathbb{Z}^{3}$ (of size $4^{n}$ ), the effective resistance between two corners this cube are bounded above uniformly in $n$. Moreover, the distance between these two points in $G_{n}$ is at least $n$.

Let $T$ be a binary tree, and let $G$ be the graph formed by replacing each edge of $T$ at height $k$ from the root with a copy of $G_{3^{k}}$, so that the endpoints of each edge of $T$ are identified with opposite corners in the corresponding copy of $D_{3^{k}}$. Since the graphs $G_{n}$ are uniformly hyperbolic and $T$ is a tree, it is easily verified that $G$ is also hyperbolic. The effective resistance from the root to infinity in $G$ is at most a constant multiple of the effective resistance to infinity of the root in $T$, so that $G$ is transient. However, $G$ does not contain a transient tree, since every tree contained in $G$ is isomorphic to a binary tree in which each edge at height $k$ from the root has been stretched by at least $3^{k}$, plus some finite bushes.

## 7 A transient graph with no embeddable transient subgraph

We say that a graph $H$ has a graph $K$ as a minor, if $K$ can be obtained from $H$ by deleting vertices and edges and by contracting edges. Let $K^{r}$ denote the complete graph on $r$ vertices.
Proposition 7.1. There is a transient bounded degree graph $G$ such that every transient subgraph of $G$ has a $K^{r}$ minor for every $r \in \mathbb{N}$.

In particular, $G$ has no transient subgraph that embeds in any surface of finite genus.
We now construct this graph $G$. We will start with the infinite binary tree with root $o$, and replace each edge at distance $r$ from $o$ with a gadget $D_{2^{r}}$ which we now define. Given $n\left(=2^{r}\right)$, the vertices of $D_{n}$ are organized in $2 n+1$ levels numbered $-n, \ldots,-1,0,1, \ldots, n$. Each level $i$ has $2^{n-|i|}$ vertices, and two levels $i, j$ form a complete bipartite graph whenever $|i-j|=1$; otherwise there is no edge between levels $i, j$. Any edge of $D_{n}$ from level $i \geq 0$ to level $i+1$ or from level $-i$ to level $-(i+1)$ is given a resistance equal to $2^{n-|i|}$ (we will later subdivide such edges into paths of that many edges each having resistance 1). With this choice, the effective resistance $R_{i}$ between levels $i$ and $i+1$ of $D_{n}$ is $2^{n-|i|}$ divided by the number of edges between those two levels, that is, $R_{i}=\frac{2^{n-|i|}}{2^{n-|i|} 2^{n-|i|-1}}=2^{-n+|i|+1}$, and so the effective resistance in $D_{n}$ between its two vertices at levels $n$ and $-n$ is $O(1)$

Let $G^{\prime}$ be the graph obtained from the infinite binary tree with root o by replacing each edge $e$ at distance $n$ from $o$ with a disjoint copy of $D_{n}$, attaching the two vertices at levels $n$ and $-n$ of $D_{n}$ to the two end-vertices of $e$. We will later modify $G^{\prime}$ to obtain a bounded degree $G$ with similar properties satisfying Proposition 7.1.

Note that as $D_{n}$ has effective resistance $O(1)$, the graph $G^{\prime}$ is transient by Lyons' criterion.

We are claiming that if $H$ is a transient subgraph of $G^{\prime}$, then $H$ has a $K^{r}$ minor for every $r \in \mathbb{N}$.

This will follow from the following basic fact of finite extremal graph theory [22, 18, 12]
Theorem 7.2. For every $r \in \mathbb{N}$ there is a constant $c_{r}$ such that every graph of average degree at least $c_{r}$ has a $K^{r}$ minor.
Lemma 7.3. If $H$ is a transient subgraph of $G^{\prime}$, then $H$ has a $K^{r}$ minor for every $r \in \mathbb{N}$.
Proof. Suppose that $H$ has no $K^{r}$ minor for some $r$, and fix any $m \in \mathbb{N}$. For every copy $C$ of the gadget $D_{n}$ in $G^{\prime}$ where $n>m$, consider the bipartite subgraph $G_{m}=G_{m}(C)$ of $H$ spanned by levels $m$ and $m+1$ of $C \cap H$. By Theorem 7.2, the average degree of $G_{m}$ is at most $c_{r}$. Thus, if we identify each of the partition classes of $G_{m}$ into one vertex, we
obtain a graph with 2 vertices and at most $\frac{3}{2} 2^{n-m} c_{r}$ parallel edges, each of resistance $2^{n-m}$, so that the effective resistance of the contracted graph is at least $\frac{2}{3 c_{r}}=: C_{r}$.

Now repeating this argument for $m+1, m+2, \ldots$, we see that the effective resistance between the two partition classes of $G_{m+k}$ (which is edge-disjoint to $G_{m}$ ) is also at least the same constant $C_{r}$. This easily implies that the effective resistance between the two endvertices of $C \cap H$ for any copy $C$ of $D_{n}$ is $\Omega(n)$. Since $G^{\prime}$ has $2^{r}$ copies of $D_{2^{r}}$ at each 'level' $r$, we obtain that the effective resistance from $o$ (which we may assume without loss of generality to be contained in $H$ ) to infinity in $H$ is $\Omega\left(\sum_{r} 2^{r} / 2^{r}\right)=\infty$.

Thus $H$ can have no electrical flow from a vertex to infinity, and by Lyons' criterion (Theorem 4.1) it is not transient.

Recall that the edges of $G^{\prime}$ had resistances greater than 1. By replacing each edge of resistance $k$ by a path of length $k$ with edges having resistance 1 , we do not affect the transience of $G^{\prime}$. We now modify $G^{\prime}$ further into a graph $G$ of bounded degree, which will retain the desired property.

Let $x$ be a vertex of some copy $C$ of $D_{n}$, at some level $j \neq n,-n$ of $C$. Then $x$ sends edges to the two neighbouring levels $j \pm 1$. Each of those levels $L, L^{\prime}$, sends $2^{k \pm 1}$ edges to $x$ for some $k$. Now disconnect all the edges from $L$ to $x$, attach a binary tree $T_{L}$ of depth $k \pm 1$ to $x$, and then reconnect those edges, one at each leaf of $T_{L}$.

Do the same for the other level $L^{\prime}$, attaching a new tree $T_{L^{\prime}}$ of appropriate depth to $x$.
Note that this operation affects the edges incident with $x$ only, and every other vertex of $G^{\prime}$, even those adjacent with $x$, retains its vertex degree. Thus we can perform this operation on every such vertex $x$ simultaneously, with the understanding that if $e=x x^{\prime}$ is an edge of $G^{\prime}$, and both $x, x^{\prime}$ are replaced by trees $T, T^{\prime}$ respectively by the above operation, then $e$ becomes an edge joining a leaf of $T$ to a leaf of $T^{\prime}$; see Figure 6. There are many ways to match the leaves of the trees coming from vertices in one layer of $D_{n}$ to the leaves of the trees coming from vertices in a subsequent layers, and so we have not uniquely identified the resulting graph, but what matters is that such a matching is possible because we have the same number of leaves on each side.


Figure 6: The tree $T_{L}$ we replaced $x$ with in order to turn $G^{\prime}$ into a bounded degree graph $G$, and a few similar trees for other vertices in the level of $x$ and the level $L$ above.

Let $G$ denote a graph obtained by performing this operation to every vertex $x$ as above. Note that $G$ has maximum degree 6 (we did not need to modify the vertices at levels $n,-n$ in $C$, as they already had degree 6).

Now let's check that $G$ is still transient, by considering the obvious flow to infinity: we start from the canonical flow $f$ of strength 1 from $o$ to infinity in $G^{\prime}$. Recall that every edge $e=x x^{\prime}$ of $G^{\prime}$ of resistance $k$ was subdivided into a path $P_{e}$ of length $k$ consisting of edges of resistance 1 , then $x, x^{\prime}$ were replaced by trees $T, T^{\prime}$, and now $P_{e}$ joins a leaf of
$T$ to a leaf of $T^{\prime}$ in $G$. Note that there is a unique path $Q_{e} \supset P_{e}$ in $T \cup P_{e} \cup T^{\prime}$ from the root of $T$ to the root of $T^{\prime}$. For each edge $e$ of $G^{\prime}$, we send a flow of intensity $f(e)$ along that path $Q_{e}$; easily, this induces a flow $j$ on $G$ from $o$ to infinity.

We claim that the energy of $j$ is finite, which means that $G$ is transient by Lyons' criterion. Indeed, the contribution of the path $P_{e}$ to the energy of $j$ coincides with the contribution of $e$ to the energy of $f$, and so their total contribution is finite. Let us now bound the contributions of the trees we introduced when defining $G$ from $G^{\prime}$. For this, we will use the following basic observation about flows on binary trees

Let $T$ be a binary tree of depth $k$, and let $j$ be a flow from the root of $T$ to its leaves such that every two edges at the same layer carry the same flow. Then the energy dissipated by $j$ in all of $T$ equals $\left(2^{k+1}-1\right)$ times the energy
dissipated by $j$ in the last layer of $T$.
Indeed, it is straightforward to check that the energy dissipated in each layer equals twice the energy dissipated in the next layer, and so the energy dissipated by $j$ in all of $T$ eqals $\left(1+2+\ldots+2^{k}\right)$ times the energy dissipated by $j$ in the last layer.

Consider now two consecutive levels $L, M$ in a copy of some gadget $D_{n}$ in $G^{\prime}$, and suppose $L$ has $2^{k}$ vertices and $M$ has $2^{k+1}$ vertices. Recall that each $L-M$ edge had resistance $2^{k}$ in $G^{\prime}$. Furthermore, the $f$ value is the same for all these edges; let $b$ denote that common value. Thus, letting $E$ denote the number of $L-M$ edges, the total energy dissipated by $f$ on $L-M$ edges is $E 2^{k} b^{2}$.

Note that for each tree $T$ we introduced in the definition of $G$, each leaf of $T$ was joined with exactly one edge of $G^{\prime}$. It follows that for each such tree $T$ between the layers $L$ and $M$, the value of $j$ at any edge in the last layer of $T$ is $b$. Since each $L-M$ edge of $G^{\prime}$ gave rise to exactly two such last-layer edges, namely one in the tree substituting each of its end-vertices, the total energy dissipated by $j$ in all last-layer edges of $G$ between the layers $L$ and $M$ is $2 E b^{2}$. By (7.1), the total energy dissipated by $j$ in all layers of all trees we introduced between layers $L$ and $M$, equals that amount multiplied by a constant smaller than $2^{k+1}$. Recalling that the total energy dissipated by $f$ on $L-M$ edges was $E 2^{k} b^{2}$, we see that the energy dissipated by $j$ between layers $L$ and $M$ is less than 5 times that dissipated by $f$. Since this holds for each copy of each $D_{n}$, we deduce that $j$ has finite energy since $f$ does, proving that $G$ is transient too.

Note that $G^{\prime}$ can be obtained from $G$ by contracting edges. Thus any transient $H \subseteq G$ has a transient minor $H^{\prime} \subseteq G^{\prime}$, because contracting edges preserves transience by Lyons' criterion. As we have proved that $H^{\prime}$ has a $K^{r}$ minor (Lemma 7.3), so does $H$ as any minor of $H^{\prime}$ is a minor of $H$.

Despite Proposition 7.1, the following remains open
Question 7.4 (I. Benjamini (private communication)). Does every bounded-degree transient graph have a transient subgraph which is sphere-packable in $R^{3}$ ?

## 8 Problems

It is not hard to see that our hyperbolic souvlaki $\Psi$ is amenable, that is, we have $\inf _{\emptyset \neq S \subset \Psi \text { finite }} \frac{|\partial S|}{|S|}=0$, where $\partial S=\{v \in V(\Psi) \backslash S \mid$ there exists $w \in S$ adjacent to $v\}$. We do not know if this is an essential feature:
Problem 8.1. Is there a non-amenable counterexample to Conjecture 1.1?
Similarly, one can ask
Problem 8.2. Is there a non-amenable, hyperbolic graph with bounded-degrees, $C$ dense infinite geodesics, and the Liouville property, the hyperbolic boundary of which consists of a single point?

## A Liouville hyperbolic souvlaki

Here we did not ask for transience as it is implied by non-amenability [6].
We conclude with further questions asked by I. Benjamini (private communication)
Problem 8.3. Is there a uniformly transient counterexample to Conjecture 1.1? Is there an 1-ended counterexample?

Here uniformly transient means that there is an upper bound on the effective resistance between any vertex of the graph and infinity.

Our last problem, also by I. Benjamini (private communication), is motivated by our construction in Section 6 and the Appendix, where unimodular random graphs are defined.

Problem 8.4. Is there a bounded degree unimodular random graph that is non-Liouville but contains no transient subtree?

## A Appendix: a unimodular Liouville hyperbolic souvlaki

authored by

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Dedicated to Findus, who has always wanted to grow a tree out of meatballs [24]

## A. 1 Motivation and the result

Graphs with strange properties as constructed in the main paper cannot be transitive. On the one hand, a transitive transient graph has at least 3-dimensional volume growth (as follows from an extension of Gromov's polynomial growth theorem by Trofimov [28] and Losert [19]; see also [29, Theorem 5.11]), hence by the Coulhon-Saloff-Coste isoperimetric inequality (see [11] or [21, Theorem 6.29]) it has at least 3-dimensional isoperimetry, and hence by Thomassen's result [27] it contains a transient subtree. On the other hand, a transitive transient hyperbolic graph must be non-amenable [10], and non-amenable transitive graphs are non-Liouville [17] and contain non-amenable subtrees [6].

This raises the question whether such exceptional graphs can possess any sort of homogeneity. Beyond transitivity, a very natural class of graphs, especially when random walks are considered, is the class of unimodular random graphs, introduced in [8], studied in depth in [1] and in many works since; see [25, Chapter 14] for an overview. Here are the main definitions.

Let $\mathcal{G}_{\star}$ be the space of isomorphism classes of locally finite labeled rooted graphs, and let $\mathcal{G}_{\star \star}$ be the space of isomorphism classes of locally finite labeled graphs with an ordered pair of distinguished vertices, each equipped with the natural local topology: two (doubly) rooted graphs are "close" if they agree in "large" neighborhoods of the root(s).

[^4]Definition A.1. We say that a Borel measure $\mu$ on $\mathcal{G}_{\star}$ is unimodular if it obeys the Mass Transport Principle:

$$
\int_{\mathcal{G}_{\star}} \sum_{x \in V(G)} f(G, o, x) d \mu(G, o)=\int_{\mathcal{G}_{\star}} \sum_{x \in V(G)} f(G, x, o) d \mu(G, o)
$$

for any Borel function $f: \mathcal{G}_{\star \star} \rightarrow[0, \infty]$.
There are several other equivalent definitions; see [25, Definition 14.1]. Probably the nicest one, which works in most situations (e.g., bounded degree non-deterministic graphs), is that the Markov chain on $\mathcal{G}_{\star}$ generated by continuous time random walk on $G$ with rate 1 exponential clocks on the edges is reversible.

An important class of unimodular graphs consists of Cayley graphs of finitely generated groups and of invariant random subgraphs of a Cayley graph. Another one is the class of sofic measures: the closure of the set of finite graphs with a uniform random root under local weak convergence, which is just weak convergence of measures in the space $\mathcal{G}_{\star}$.

Since many results on random walks and harmonic functions on transitive graphs generalize to unimodular or, more generally, stationary random graphs [5, 13, 3], it is natural to ask what the situation is in the present case. Here is our answer:
Theorem A.2. There exists a bounded degree unimodular random graph that is a.s. transient and hyperbolic, but Liouville and has no transient subtree.

Our construction will be a splice between the one in the main paper and the so-called $d$-regular canopy tree, which is the local weak limit of larger and larger balls in the $d$-regular tree. It is partly motivated by [9], where similar counterexamples for Bernoulli percolation are constructed based on the canopy tree. However, making the splice is not entirely straightforward here, since we have to put the meatballs on the canopy tree in a way that the graph remains unimodular. For this, the exponentially growing meatballs of the original construction would not work.

## A. 2 The construction

Take a $d$-ary tree $T_{n}$ of height $n$ (i.e., the root has $d$ children, each of which has $d$ children, and so on, stopping with the $n$th descendent generation). We are going to replace each edge of $T_{n}$ by a modification of the graphs $M_{k}$ in the Souvlaki construction of the main paper. Recall what $W, H_{2}, H_{3}$ are. Consider a subpath of the bottom double ray $P_{k}$ of $W$ of length $(k-1)^{2}+k^{2}+k^{4}$. Define $M_{k}$ to be the subgraph of $H_{3}$ induced by vertices of the form $(t, w)$ such that $w$ lies at or above $P_{k}$ and has height $h(w)$ at most $k$. We call $M_{k}$ the meatballs.

We now "replace" each edge of $T_{n}$ at height $n-k+1$ (edges such that the vertex closer to the root has height $n-k$ ) by $M_{k}$ for $k=1, \ldots, n$. The word replace is within quotes because we have to specify the way we glue adjacent meatballs. We divide $P_{k}=L_{k} \cup R_{k} \cup A_{k}$, where $L_{k}$ is the segment of the leftmost $k^{2}$ vertices, $R_{k}$ is the segment of the rightmost $(k-1)^{2}$ vertices, and $A_{k}$ is the middle $k^{4}$ vertices. Now let $M_{k}^{L}$ denote the set of vertices $(t, w)$ so that $w$ lies on or above $L_{k} \cup A_{k}$ and has height at most $k$. Define $M_{k}^{R}$ to be the graph induced by the rest of the vertices. Note that $M_{k}^{R}$ and $M_{k}^{L}$ are joined together by a set of edges. Let $B_{k}$ denote the endpoint of these edges that lie in $M_{k}^{L}$.

Since we have the tree $T_{n}$ instead of just a line, we need to modify $M_{k}$ a bit, so that it branches into $d$ copies of $M_{k}^{R}$ for the identifications. That is, we take one copy of $M_{k}^{L}$ and $d$ copies of $M_{k}^{R}$, then glue each of the latter with $M_{k}^{L}$ along $B_{k}$. Call this new gadget $M_{k}^{\prime}$. The height function $h$ extends to $M_{k}^{\prime}$, with values between 0 and $k$.

Now take an edge $e$ at height $n-k+1$, for $k \geq 1$. Remove it and replace it with $M_{k}^{\prime}$ so that $e$ corresponds to the segment $L_{k} \cup A_{k}$, while its $d$ children $e_{1}, \ldots e_{d}$ at height $n-k+2$ correspond to the $d$ copies of $R_{k}$ for $e$. Since each $R_{k}$ contains $(k-1)^{2}$ vertices, we can identify them with the copies of $L_{k-1}$ for the edges $e_{i}$. This completes the gluing procedure. For an edge $e$ at height $n$, replace $e$ just by $M_{1}^{L}$, without branching into copies of $M_{1}^{R}$. Call the new graph so obtained $T_{n}^{\prime}$.

Now pick a uniform random vertex $\rho_{n}$ from $T_{n}^{\prime}$ and take a weak limit. Call the limit $(T, \rho)$. Clearly this graph is unimodular, Gromov hyperbolic and bounded degree, from arguments in the main paper.

## A. 3 Proofs

Root height Pick any integer $d>6$. At height $i$ of $M_{k}^{L}$, the number of vertices is $3^{i} 2^{i}\left(k^{2}+k^{4}\right)$, with the factor $3^{i}$ coming from $H_{2}$, and the factor $2^{i}$ coming from $W$. Thus, the volume of $M_{k}^{L}$ is

$$
v_{k}:=\frac{6^{k+1}-1}{5}\left(k^{4}+k^{2}\right)
$$

and the probability that the uniform root in $T_{n}$ is a vertex in one of the $M_{k}^{L} \mathbf{s}$ is

$$
p_{k, n}:=\frac{v_{k} d^{n-k+1}}{\sum_{j=1}^{n} v_{j} d^{n-j+1}}=\frac{v_{k} d^{-k}}{\sum_{j=1}^{n} v_{j} d^{-j}}
$$

Since $d>6$ implies that $v_{k} d^{-k}$ is summable, the limit $p_{k}:=\lim _{n \rightarrow \infty} p_{k, n}$ is a proper probability distribution for $k=1,2, \ldots$. Therefore, the root in the limit $(T, \rho)$ is almost surely at a level corresponding to a finite $k$, at a finite distance from the leaves (of the underlying canopy tree that is the local weak limit of the original trees $T_{n}$ ). In other words, we can think of the limit $(T, \rho)$ as a souvlaki with a canopy tree skewer, with a random root somewhere.

Constructing a good flow One can think of the canopy tree as an infinite spine with finite bushes hanging off of it. Similarly, our canopy tree souvlaki has an infinite spine, a "traditional" infinite souvlaki. It is of course enough to show that this infinite spine is transient. We will construct for each $k \geq 1$ a unit flow $g$ from $R_{k+1}$ to $L_{k+1}$, with an energy that is summable in $k$. Concatenating these flows yields a flow along the spine to infinity, with finite energy, hence the spine turns out to be transient by Terry Lyons' criterion. (Unfortunately, the roles of $R \mathrm{~s}$ and $L \mathrm{~s}$ are now swapped compared to the main paper, due to the way that the infinite limit is constructed.)

Note that $R_{k+1}$ has $k^{2}$ vertices and $L_{k+1}$ has $(k+1)^{2}$ vertices. We name the vertices in $R_{k+1}$ as $r_{1}, r_{2}, \ldots, r_{k^{2}}$ and the vertices in $L_{k+1}$ as $l_{1}, \ldots, l_{(k+1)^{2}}$. We will construct a flow $g_{j}$ from $r_{j}$ to $L_{k+1}$ and $g$ will be the sum the flows $g_{j}$. The flow $g$ will have outflow $1 / k^{2}$ from each vertex in $R_{k+1}$ and inflow $1 /(k+1)^{2}$ for each vertex in $L_{k+1}$.

In the main paper, since the meatballs had exponentially growing lengths, there was a straightforward division of the total flow from each vertex into two vertices. Since here everything is polynomial, the division is slightly more complicated, but this is just a technicality. We deal with this as follows.

Since $H_{2}$ is transient, there exists a natural unit flow $t$ (equally branching off at each vertex) with finite energy. Then, the flow $g_{j}$ is just $\frac{1}{k^{2}} t$ up to height $k$, for each $j=1, \ldots, k^{2}$. After this, $g_{j}$ takes $k^{2} /(k+1)^{2}$ fraction of the total incoming flow at height $k$, flows along the horizontal edges at height $k$ until reaching above $l_{j}$, then flows along the tree proportionally with $t$ (in reverse direction) to reach $l_{j}$. So, for $j=1, \ldots, k^{2}$, the total flow into $l_{j}$ is already $1 /(k+1)^{2}$.

There is still $\left(1-k^{2} /(k+1)^{2}\right)$ fraction of the outflow from $R_{k+1}$ left at level $k$. In fact, there is $\left(k^{-2}-(k+1)^{-2}\right) 3^{-k}$ amount left in each vertex at height $k$, in each $g_{j}$, $1 \leq j \leq k^{2}$. Flow this amount horizontally along level $k$ until we reach above the vertex $l_{k^{2}+1}$. The total flow here (summed over $j$ ) is now $\left((k+1)^{2}-k^{2}\right) /(k+1)^{2}$. Now we take $1 /(k+1)^{2}$ out of this, and flow it along the tree above $l_{k^{2}+1}$ proportionally with $t$, in reverse direction. We flow the remaining amount of flow at level $k$ horizontally to reach above $l_{k^{2}+2}$, and again drop $1 /(k+1)^{2}$ amount along the tree with $t$. We continue like this until all the flow is exhausted. Note that we input a flow $1 /(k+1)^{2}$ for each $l_{j}$, $k^{2}<j \leq(k+1)^{2}$. Thus all in all, this is a unit flow from $R_{k+1}$ to $L_{k+1}$.

Energy computation The total energy cost for going up the tree to height $k$ for each $r_{j}$ is $E(t) / k^{4}$. The flow received at each vertex is $3^{-k} / k^{2}$. On each horizontal edge at height $k$, there are at most $k^{2}$ flows that we are summing up, hence total flow through is at most $3^{-k}$. The total length of horizontal paths is $O\left(k^{4}\right) 3^{k} 2^{k}$, hence the total energy along the horizontal edges is $O\left(k^{4}\right) 6^{k} 3^{-2 k}$. The energy for going down each tree above $l_{j}$ is again at most $E(t) / k^{4}$. Therefore, the total energy dissipation is

$$
O\left(k^{4}\right) 6^{k} 3^{-2 k}+O\left(k^{2}\right) E(t) / k^{4}=O\left(k^{4}\right)(2 / 3)^{k}+O\left(1 / k^{2}\right)=O\left(1 / k^{2}\right)
$$

which is summable in $k$. This concludes the proof of transience.

Liouville property Removing the infinite spine, the canopy tree souvlaki falls apart into finite pieces. Thus, random walk started anywhere in the graph will almost surely hit the spine. This and the Optional Stopping Theorem for bounded martingales imply that any bounded harmonic function is determined by its restriction to the spine. Also, the radial symmetry of the harmonic function in the meatballs along the spine is preserved, as in the main paper. Thus the graph is Liouville by the same argument as in the main paper.

Transient subtree We first claim that any subtree inside the infinite spine must be recurrent. This can be shown by the same argument as in the main paper. Namely, in the proof of Theorem 6.1, the degrees $d\left(v_{k}\right)$ and $d\left(v_{k+1}\right)$ are $O\left(k^{2}\right)$, hence following the proof of Lemma 6.2 shows that taking $s=k^{4}$ is enough for the effective resistance between $v_{k}$ and $v_{k+1}$ to be uniformly positive. This is exactly the choice we made in defining $A_{k}$, thus the effective resistance of any subtree to infinity is infinite, hence it is recurrent.

Now, for a general subtree, when we do the contraction into the vertices $v_{k}$, then the portions of the subtree inside the finite bushes off the spine get contracted into finite pieces, each attached to the contracted graph from the spine at a single vertex. These finite pieces do not influence transience of the contracted graph, hence the original subtree is also recurrent.

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    ${ }^{\ddagger}$ University of Warwick

[^1]:    ${ }^{1}$ Formally, we pick a cyclic ordering on the neighbours of $r$ and a linear ordering on the outer neighbours of every other vertex of $T_{3}$. Given a cyclic ordering on the vertices at level $n$ of $T_{3}$, we get a cyclic ordering at level $n+1$ by replacing each vertex by the linear ordering on its outer neighbours. Now we add edges between any two vertices that are adjacent in any of these cyclic orderings.

[^2]:    ${ }^{2}$ A double ray is a 2-way infinite path.

[^3]:    ${ }^{3}$ Although $Z_{n}$ gets larger if $f(n)$ increases, the flow $f$ then branches more before 'traversing' $Z_{n}$. Since the increase of $Z_{n}$ has an additive effect on the energy while the branching has a multiplicative effect, the effect due to branching dominates, hence the energy remains bounded.

[^4]:    ${ }^{4}$ Partially supported by the Hungarian National Research, Development and Innovation Office, NKFIH grant K109684, and by the MTA Rényi Institute "Lendület" Limits of Structures Research Group.
    ${ }^{5}$ Supported by the Engineering and Physical Sciences Research Council under grant EP/103372X/1.

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