# CHARACTERIZATIONS OF ASYMMETRIC TRUNCATED TOEPLITZ OPERATORS 

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#### Abstract

The aim of this paper is to investigate asymmetric truncated Toeplitz operators with $L^{2}$-symbols between two different model spaces given by inner functions such that one divides the other. The class of symbols corresponding to the zero operator is described. Asymmetric truncated Toeplitz operators are characterized in terms of operators of rank at most 2, and the relations with the corresponding symbols are studied.


## 1. Introduction

Toeplitz operators on the Hardy space $H^{2}$, which are compositions of a multiplication operator and the orthogonal projection from $L^{2}$ onto $H^{2}$, constitute a classical topic in operator theory. In his important paper [19], Sarason explored truncated Toeplitz operators, thus generating huge interest in this class of operators (see, e.g., [4], [6], [7], [9]-[12], and [15]). Instead of the Hardy space $H^{2}$, they act on a model space $K_{\theta}^{2}=H^{2} \ominus \theta H^{2}$ associated with a given nonconstant inner function $\theta$, and a multiplication operator is composed with the orthogonal projection from $H^{2}$ onto $K_{\theta}^{2}$.

This work was inspired, on the one hand, by the work of Sarason, and, on the other hand, by [8], where asymmetric truncated Toeplitz operators were introduced (in the context of the Hardy space $H^{p}$ of the half-plane, with $1<p<\infty$ )

[^0]and studied in the case of bounded symbols. Asymmetric truncated Toeplitz operators involve the composition of a multiplication operator with two projections from $H^{2}$ onto a model space, associated with (possibly different) nonconstant inner functions $\alpha$ and $\theta$. They are a natural generalization of rectangular Toeplitz matrices, which appear in various contexts, such as in the study of finite-time convolution equations, signal processing, control theory, probability, approximation theory, and diffraction problems (see, e.g., [2]-[4], [16], [17], [1], and [20]).

Here we consider bounded asymmetric truncated Toeplitz operators with $L^{2}$-symbols, defined between two model spaces $K_{\theta}^{2}$ and $K_{\alpha}^{2}$, where $\alpha$ divides $\theta$ $(\alpha \leq \theta)$. We study various properties of these operators and their relations with the corresponding symbols, and we obtain a necessary and sufficient condition for a bounded operator between two model spaces to be an asymmetric truncated Toeplitz operator in terms of rank 2 operators, thus generalizing a corresponding result by Sarason for the case where $\alpha=\theta$. In the asymmetric case, however, a more complex connection between the operators and their symbols is revealed, which is not apparent when the two model spaces involved are the same.

Our article is organized as follows. In Section 2, we present some auxiliary results on model spaces, their direct sum decompositions, and the associated projections and conjugations. In Section 3, asymmetric truncated Toeplitz operators with $L^{2}$-symbols are introduced and some of their basic properties are presented. In Section 4, the class of all possible symbols for a given asymmetric truncated Toeplitz operator is described, and a necessary and sufficient condition for the operator to be zero is obtained. In Sections 5 and 6, we generalize Sarason's characterization of truncated Toeplitz operators in terms of operators of rank 2 at most, highlighting the similarities and differences between the symmetric and the asymmetric cases. In Section 7, we obtain conditions under which that characterization can be done in terms of a rank 1 operator. In Section 8, we address the inverse problem of determining a symbol for a given asymmetric truncated Toeplitz operator in terms of the action of $A$ and its adjoint on certain reproducing kernel functions. Moreover, we give the formulas for a symbol of the operator satisfying the characterizations in terms of rank 2 operators presented in Sections 5 and 6.

## 2. Model spaces, COnjugations, and decompositions

Let $L^{2}$ denote the space $L^{2}(\mathbb{T}, m)$, where $\mathbb{T}$ is the unit circle and $m$ is the normalized Lebesgue measure on $\mathbb{T}$, and let $H^{2}$ be the Hardy space on the unit disk $\mathbb{D}$, identified as usual with a subspace of $L^{2}$. Similarly, $L^{\infty}=L^{\infty}(\mathbb{T}, m)$, and we denote by $H^{\infty}$ the space of all analytic and bounded functions on $\mathbb{D}$. Denoting by $H_{0}^{2}$ the subspace consisting of all functions in $H^{2}$ which vanish at 0 , we have $L^{2} \ominus H^{2}=\overline{H_{0}^{2}}$, and we denote by $P$ and $P^{-}$the orthogonal projections from $L^{2}$ onto $H^{2}$ and $\overline{H_{0}^{2}}$, respectively.

With any given inner function $\theta$ we associate the so-called model space $K_{\theta}^{2}$, defined by $K_{\theta}^{2}=H^{2} \ominus \theta H^{2}$. We also have $K_{\theta}^{2}=H^{2} \cap \theta \overline{H_{0}^{2}}$; thus,

$$
f \in K_{\theta}^{2} \quad \text { if and only if } \quad \bar{\theta} f \in \overline{H_{0}^{2}} \quad \text { and } \quad f \in H^{2} .
$$

In particular, if $f \in K_{\theta}^{2}$, then $\theta \bar{f} \in H_{0}^{2}$. Let $P_{\theta}$ be the orthogonal projection $P_{\theta}: L^{2} \rightarrow K_{\theta}^{2}$. We then have the following.

Proposition 2.1. Let $\theta$ be a nonconstant inner function, and let $K_{\theta}^{2}=H^{2} \ominus \theta H^{2}$ be the associated model space. Then
(1) $P_{\theta}=\theta P^{-} \bar{\theta} P=\theta P^{-} \bar{\theta}-P^{-}$,
(2) $P_{\theta} f=\theta P^{-} \bar{\theta} f=f-\theta P \bar{\theta} f$ for all $f \in H^{2}$, and
(3) $P_{\theta} \bar{f}=P_{\theta} P \bar{f}=\overline{f(0)} P_{\theta} 1=\overline{f(0)}(1-\overline{\theta(0)} \theta)$ for all $f \in H^{2}$.

Model spaces are also equipped with conjugations (antilinear isometric involutions), which are important tools in the study of model spaces and truncated Toeplitz operators (see, e.g., [13], [14], and [18]). For a given inner function $\theta$, the conjugation $C_{\theta}$ is defined by $C_{\theta}: L^{2} \rightarrow L^{2}$,

$$
C_{\theta} f(z)=\theta \overline{z f(z)}
$$

It is worth noting that $C_{\theta}$ preserves the space $K_{\theta}^{2}$ and maps $\theta H^{2}$ onto $L^{2} \ominus H^{2}$.
Recall that, for $\lambda \in \mathbb{D}$, the kernel function in $H^{2}$ denoted by $k_{\lambda}$ is given by $k_{\lambda}(z)=\frac{1}{1-\lambda z}$. Similarly, for an inner function $\theta$, in $K_{\theta}^{2}$ the kernel function $k_{\lambda}^{\theta}$ is given by $k_{\lambda}^{\theta}=P_{\theta} k_{\lambda}$; that is, $k_{\lambda}^{\theta}=k_{\lambda}(1-\overline{\theta(\lambda)} \theta)$. The set $\left\{k_{\lambda}^{\theta}: \lambda \in \mathbb{D}\right\}$ is linearly dense in $K_{\theta}^{2}$. Since $k_{\lambda}^{\theta} \in K_{\theta}^{\infty}$, where $K_{\theta}^{\infty}$ denotes the subspace $H^{\infty} \cap K_{\theta}^{2}$, the space $K_{\theta}^{\infty}$ is dense in $K_{\theta}^{2}$ (see [19, p. 494]).

Defining $\tilde{k}_{\lambda}^{\theta}=C_{\theta} k_{\lambda}^{\theta}$, we have in particular

$$
k_{0}^{\theta}(z)=1-\overline{\theta(0)} \theta(z), \quad \tilde{k}_{0}^{\theta}(z)=\bar{z}(\theta(z)-\theta(0))
$$

It is easy to see that, for all $f \in K_{\theta}^{2}$,

$$
\left\langle f, k_{0}^{\theta}\right\rangle=f(0), \quad\left\langle f, \tilde{k}_{0}^{\theta}\right\rangle=\overline{\left(C_{\theta} f\right)(0)} .
$$

Now let us consider two nonconstant inner functions $\alpha$ and $\theta$. If $\bar{\alpha} \theta$ is an inner function, then we say that $\alpha$ divides $\theta$, and we write $\alpha \leq \theta$.

Proposition 2.2. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. Then the following hold:
(1) $K_{\theta}^{2}=K_{\alpha}^{2} \oplus \alpha K_{\underline{\theta}}^{2}$,
(2) $P_{\theta}=P_{\alpha}+\alpha P_{\frac{\theta}{\alpha}} \frac{\alpha}{\bar{\alpha}}$,
(3) $k_{0}^{\theta}=k_{0}^{\alpha}+\overline{\alpha(0)} \alpha k_{0}^{\frac{\theta}{\alpha}}$,
(4) $\tilde{k}_{0}^{\theta}=\frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}+\alpha \tilde{k}_{0}^{\frac{\theta}{\alpha}}$, and
(5) $P_{\alpha} k_{0}^{\theta} \stackrel{ }{=} k_{0}^{\alpha}, P_{\alpha} \tilde{k}_{0}^{\theta}=\frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}$.

Proof. Items (1), (2), and (3) were proved in [12, p. 97]. We need to prove only (4), and (5) follows immediately from (3) and (4). In fact, we have

$$
\tilde{k}_{0}^{\theta}=\bar{z}(\theta-\theta(0))=\alpha \bar{z} \frac{\theta}{\alpha}-\bar{z} \alpha(0) \frac{\theta}{\alpha}(0)=\alpha \bar{z}\left(\frac{\theta}{\alpha}-\frac{\theta}{\alpha}(0)\right)+\frac{\theta}{\alpha}(0) \bar{z}(\alpha-\alpha(0)),
$$

where $\bar{z}\left(\frac{\theta}{\alpha}-\frac{\theta}{\alpha}(0)\right)=\tilde{k}_{0}^{\frac{\theta}{\alpha}}, \bar{z}(\alpha-\alpha(0))=\tilde{k}_{0}^{\alpha}$.

The following proposition describes some relations between decompositions and conjugations. Note that, if $\alpha \leq \theta$, then any $f \in K_{\theta}^{2}$ can be uniquely decomposed as $f=f_{1}+\alpha f_{2}$ for some $f_{1} \in K_{\alpha}^{2}$ and some $f_{2} \in K_{\frac{\theta}{\alpha}}^{2}$, or as $f=f_{2}+\frac{\theta}{\alpha} f_{1}$ for some $f_{1} \in K_{\alpha}^{2}$ and some $f_{2} \in K_{\underline{\theta}}^{2}$. Then the conjugation $C_{\theta}$ can be seen as $C_{\theta}: K_{\theta}^{2}=$ $K_{\alpha}^{2} \oplus \alpha K_{\frac{\theta}{\alpha}}^{2} \rightarrow K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \stackrel{\alpha}{\alpha} K_{\alpha}^{2}$, or as $C_{\theta}: K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \frac{\theta}{\alpha} K_{\alpha}^{2} \rightarrow K_{\theta}^{2}=K_{\alpha}^{2} \oplus \alpha K_{\frac{\theta}{\alpha}}^{2}$. Now we have the following.

Proposition 2.3. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. If $f_{1} \in K_{\alpha}^{2}$ and $f_{2} \in K_{\frac{\theta}{\alpha}}^{2}$, then
(1) $C_{\theta}\left(f_{1}+\alpha f_{2}\right)=C_{\frac{\theta}{\alpha}} f_{2}+\frac{\theta}{\alpha} C_{\alpha} f_{1}$, and
(2) $C_{\theta}\left(f_{2}+\frac{\theta}{\alpha} f_{1}\right)=C_{\alpha} f_{1}+\alpha C_{\frac{\theta}{\alpha}} f_{2}$.

Proof. If we let $f_{1} \in K_{\alpha}^{2}, f_{2} \in K_{\frac{\theta}{\alpha}}^{2}$, then $C_{\theta}\left(f_{1}+\alpha f_{2}\right)=\theta \bar{z} \bar{f}_{1}+\theta \bar{z} \bar{\alpha} \bar{f}_{2}=\frac{\theta}{\alpha} \alpha \bar{z} \bar{f}_{1}+$ $\frac{\theta}{\alpha} \bar{z}_{2}=C_{\frac{\theta}{\alpha}}\left(f_{2}\right)+\frac{\theta}{\alpha} C_{\alpha}\left(f_{1}\right)$, and $C_{\theta}\left(f_{2}+\frac{\theta}{\alpha} f_{1}\right)=\theta \bar{z} \bar{f}_{2}+\theta \bar{z} \overline{\bar{\alpha}} \overline{\bar{\alpha}} \bar{f}_{1}=\alpha \bar{z} \bar{f}_{1}+\alpha \frac{\theta}{\alpha} \bar{z} \bar{f}_{2}=$ $C_{\alpha}\left(f_{1}\right)+\alpha_{\alpha}^{\alpha} C_{\frac{\theta}{\alpha}}\left(f_{2}\right)$.
Corollary 2.4. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. We have, for $f \in K_{\theta}^{2}$,
(1) $P_{\alpha} C_{\theta} f=C_{\alpha} P_{\alpha}\left(\frac{\bar{\theta}}{\bar{\alpha}} f\right)$, and
(2) $P_{\frac{\theta}{\alpha}} C_{\theta} f=C_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}} \bar{\alpha} f$.

Let $S$ be the unilateral shift on the Hardy space $H^{2}$, and, for a nonconstant inner function $\theta$, let $S_{\theta}=P_{\theta} S_{\mid K_{\theta}^{2}}$ be the compression of $S$ to $K_{\theta}^{2}$. The space $K_{\theta}^{2}$ is invariant for $S^{*}$; thus, $\left(S_{\theta}\right)^{*}=S^{*}{ }_{\mid K_{\theta}^{2}}$. Note that, for any $f \in K_{\theta}^{2}$,

$$
\begin{align*}
& S_{\theta} f=z f-\overline{\left(C_{\theta} f\right)(0)} \theta=S f-\left\langle f, \tilde{k}_{0}^{\theta}\right\rangle \theta,  \tag{2.1}\\
& S_{\theta}^{*} f=\bar{z}(f-f(0)) \tag{2.2}
\end{align*}
$$

In particular,

$$
\begin{equation*}
S_{\theta}^{*} k_{0}^{\theta}=-\overline{\theta(0)} \tilde{k}_{0}^{\theta}, \quad S_{\theta} \tilde{k}_{0}^{\theta}=-\theta(0) k_{0}^{\theta} . \tag{2.3}
\end{equation*}
$$

The function $k_{0}^{\theta}$ is a cyclic vector for $S_{\theta}$, and $\tilde{k}_{0}^{\theta}$ is a cyclic vector for $S_{\theta}^{*}$ (see [19, Lemma 2.3]). In what follows we will use the defect operators $I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}=k_{0}^{\theta} \otimes k_{0}^{\theta}$ and $I_{K_{\theta}^{2}}-S_{\theta}^{*} S_{\theta}=\tilde{k}_{0}^{\theta} \otimes \tilde{k}_{0}^{\theta}$, using the notation $(x \otimes y) z=\langle z, y\rangle x$ for any $x, y, z$ in a Hilbert space $H$ (see [19, Lemma 2.4]). We list below some simple properties that will be used later.

Proposition 2.5. Let $\alpha$ and $\theta$ be nonconstant inner functions such that $\alpha \leq \theta$. Then
(1) $P_{\alpha} S_{\theta}=S_{\alpha} P_{\alpha}$ on $K_{\theta}^{2}$,
(2) $S_{\theta}^{*} P_{\alpha}=S_{\alpha}^{*}$ on $K_{\alpha}^{2}$,
(3) $S_{\theta}^{n} k_{0}^{\theta} \in K_{\theta}^{\infty}$ for $n \geq 0$, and
(4) $\left(S_{\theta}^{n}\right) * \tilde{k}_{0}^{\theta} \in K_{\theta}^{\infty}$ for $n \geq 0$.

Proof. The proof of (1) and (2) follows immediately. By [12, Theorem 9.2.2] we know that $S_{\theta}^{n} f=P_{\theta} S^{n} f$ for any $f \in K_{\theta}^{2}, n \geq 0$; hence,

$$
S_{\theta}^{n} k_{0}^{\theta}=P_{\theta} S^{n} k_{0}^{\theta}=P_{\theta}\left(z^{n}(1-\overline{\theta(0)} \theta)\right)=P_{\theta} z^{n}=z^{n}-\theta P\left(\bar{\theta} z^{n}\right)
$$

Since $P\left(\bar{\theta} z^{n}\right)$ is a polynomial, then we get $S_{\theta}^{n} k_{0}^{\theta} \in K_{\theta}^{\infty}$, which proves (3). On the other hand,

$$
\left(S_{\theta}^{*}\right)^{n} \tilde{k}_{0}^{\theta}=P_{\theta} \bar{z}^{n+1}(\theta-\theta(0))=\theta \bar{z}^{n+1}-P^{-} \theta \bar{z}^{n+1}
$$

since $P^{-}\left(\theta \bar{z}^{n}\right)$ is a polynomial in $\bar{z}$, we get (4).

## 3. Asymmetric truncated Toeplitz operators

Let $\alpha, \theta$ be nonconstant inner functions. For $\varphi \in L^{2}$, we define an operator $A_{\varphi}^{\theta, \alpha}: \mathcal{D} \subset K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ as $A_{\varphi}^{\theta, \alpha} f=P_{\alpha}(\varphi f)$ having domain $\mathcal{D}=\mathcal{D}\left(A_{\varphi}^{\theta, \alpha}\right)=\{f \in$ $\left.K_{\theta}^{2}: \varphi f \in L^{2}\right\}$. The operator $A_{\varphi}^{\theta, \alpha}$ is closed and densely defined in $K_{\theta}^{2}$. Note that $K_{\theta}^{\infty} \subset \mathcal{D}\left(A_{\varphi}^{\theta, \alpha}\right)$. The operator $A_{\varphi}^{\theta, \alpha}$ will be called an asymmetric truncated Toeplitz operator. If this operator is bounded, then it admits a unique bounded extension to $K_{\theta}^{2}, A_{\varphi}^{\theta, \alpha}: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$. By $\mathcal{T}(\theta, \alpha)$ we denote the space of all bounded asymmetric truncated Toeplitz operators. For $\alpha=\theta$ we will write $A_{\varphi}^{\theta}$ instead of $A_{\varphi}^{\theta, \theta}$. (Such operators are called truncated Toeplitz operators; they were studied by Sarason in [19].) In addition we will write $\mathcal{T}(\theta)$ instead of $\mathcal{T}(\theta, \theta)$. It is easy to see that the following holds.

Proposition 3.1. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. Then

$$
A_{k_{0}^{\alpha}}^{\theta, \alpha}=P_{\alpha}=A \frac{\theta, \alpha}{k_{0}^{\theta}} .
$$

If $\varphi \in L^{2}$, then we have, for all $f \in K_{\alpha}^{\infty}$,

$$
A_{\varphi}^{\theta, \alpha} f=A_{\varphi}^{\alpha} f
$$

Proposition 3.2. Let $\alpha, \theta$ be any inner functions, and let $\varphi \in L^{2}$. Then

$$
\left\langle A_{\varphi}^{\theta, \alpha} f, g\right\rangle=\left\langle f, A_{\bar{\varphi}}^{\alpha, \theta} g\right\rangle \quad \text { for all } f \in \mathcal{D}\left(A_{\varphi}^{\theta, \alpha}\right), g \in \mathcal{D}\left(A_{\bar{\varphi}}^{\alpha, \theta}\right)
$$

Moreover, $\mathcal{D}\left(A_{\bar{\varphi}}^{\alpha, \theta}\right)=\mathcal{D}\left(\left(A_{\varphi}^{\theta, \alpha}\right)^{*}\right)$, and $\left(A_{\varphi}^{\theta, \alpha}\right)^{*}=A_{\bar{\varphi}}^{\alpha, \theta}$.
Proof. A straightforward calculation shows that

$$
\begin{aligned}
\left\langle A_{\varphi}^{\theta, \alpha} f, g\right\rangle & =\left\langle P_{\alpha}(\varphi f), g\right\rangle=\langle\varphi f, g\rangle=\int f \varphi \bar{g} d m \\
& =\langle f, \bar{\varphi} g\rangle=\left\langle f, P_{\theta}(\bar{\varphi} g)\right\rangle=\left\langle f, A_{\bar{\varphi}}^{\alpha, \theta} g\right\rangle
\end{aligned}
$$

Note also that $\mathcal{D}\left(A_{\bar{\varphi}}^{\alpha, \theta}\right)=\left\{g \in K_{\alpha}^{2}: \bar{\varphi} g \in L^{2}\right\}$. Hence, if $g \in \mathcal{D}\left(A_{\bar{\varphi}}^{\alpha, \theta}\right)$, then there is $h \in K_{\theta}^{2}$ such that

$$
\left\langle A_{\varphi}^{\theta, \alpha} f, g\right\rangle=\langle f, h\rangle
$$

for all $f \in \mathcal{D}\left(A_{\varphi}^{\theta, \alpha}\right)$, that is, such that $\varphi f \in L^{2}$; therefore, $g \in \mathcal{D}\left(\left(A_{\varphi}^{\theta, \alpha}\right)^{*}\right)$. In fact, taking $h=P_{\theta}(\bar{\varphi} f)$, we have $\left\langle A_{\varphi}^{\theta, \alpha} f, g\right\rangle=\left\langle P_{\alpha}(\varphi f), g\right\rangle=\langle\varphi f, g\rangle=\langle f, \bar{\varphi} g\rangle=$ $\left\langle f, P_{\theta}(\bar{\varphi} g)\right\rangle=\langle f, h\rangle$. The converse is similarly true.

Proposition 3.3. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. Let $A_{\psi}^{\theta, \alpha}$ be an asymmetric truncated Toeplitz operator with $\psi \in H^{2}$. Then

$$
S_{\alpha} A_{\psi}^{\theta, \alpha} f=A_{\psi}^{\theta, \alpha} S_{\theta} f \quad \text { for all } f \in K_{\theta}^{\infty} .
$$

Proof. Let $f \in K_{\theta}^{\infty}$. Then by Proposition 2.1,

$$
S_{\alpha} A_{\psi}^{\theta, \alpha} f=S_{\alpha} P_{\alpha}(\psi f)=P_{\alpha} z(\psi f-\alpha P(\bar{\alpha} \psi f))=P_{\alpha}(z \psi f)
$$

since $P_{\alpha}(\alpha z P(\bar{\alpha} \psi f))=0$. On the other hand, taking (2.1) into account, for all $f \in K_{\theta}^{\infty}$, we have

$$
A_{\psi}^{\theta, \alpha} S_{\theta} f=P_{\alpha} \psi S_{\theta} f=P_{\alpha} \psi\left(z f-\left\langle f, \tilde{k}_{0}^{\theta}\right\rangle \theta\right)=P_{\alpha}(z \psi f)
$$

since $P_{\alpha}(\psi \theta)=0$.
Remark 3.4. Theorem 3.1.16 of [5] implies that, for nonconstant inner functions $\alpha, \theta$ such that $\alpha \leq \theta$, if a bounded operator $A: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ intertwines $S_{\alpha}, S_{\theta}$ (i.e., $S_{\alpha} A=A S_{\theta}$ ), then $A=A_{\psi}^{\theta, \alpha}$ for some $\psi \in H^{\infty}$.

Example 3.5. One can ask whether a similar result as in Proposition 3.3 can be obtained for $A_{\psi}^{\alpha, \theta}$ with $\alpha \leq \theta$ and $\psi \in H^{2}$; the answer is no. For example, let $\alpha=z^{2}, \theta=z^{n}$, let $n>5, \psi=z^{3}$, and let $f=z$. Then $S_{\theta} A_{\psi}^{\alpha, \theta} f=z^{5}$, but $A_{\psi}^{\alpha, \theta} S_{\alpha} f=0$.

In the next proposition, we describe the action of some (not necessarily bounded) asymmetric truncated Toeplitz operators on some particular functions. These properties will be used later on.

Proposition 3.6. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$ with $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$, where $\chi_{\frac{\theta}{\alpha}} \in K_{\frac{\theta}{\alpha}}^{2}$ and $\chi_{\alpha} \in K_{\alpha}^{2}$. Then
(1) $A_{\psi}^{\theta, \alpha} k_{0}^{\theta}=\psi$,
(2) $A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta}=\overline{\chi(0)} k_{0}^{\alpha}-\overline{\theta(0)} \alpha\left(\bar{\chi}_{\alpha}-\overline{\chi_{\alpha}(0)}\right)$,
(3) $A_{\psi}^{\theta, \alpha} \tilde{k}_{0}^{\theta}=\frac{\theta}{\alpha}(0) \psi(0) \tilde{k}_{0}^{\alpha}-\theta(0) S_{\alpha}^{*} \psi=\bar{z}\left(\frac{\theta}{\alpha}(0) \psi(0) \alpha-\theta(0) \psi\right)$,
(4) $A_{\bar{\chi}}^{\theta, \alpha} \tilde{k}_{0}^{\theta}=P_{\alpha} C_{\theta} \chi=C_{\alpha} P_{\alpha}(\overline{\bar{\theta}} \chi)=C_{\alpha} \chi_{\alpha}$,
(5) $A_{\bar{\psi}}^{\alpha, \theta} k_{0}^{\alpha}=\overline{\psi(0)} k_{0}^{\theta}-\overline{\alpha(0)}\left(\alpha \bar{\psi}-\overline{\frac{\theta}{\alpha}(0) \psi(0)} \theta\right)$,
(6) $A_{\chi}^{\alpha, \theta} k_{0}^{\alpha}=\chi-\overline{\alpha(0)} \alpha \chi_{\frac{\theta}{\alpha}}$,
(7) $A_{\bar{\psi}}^{\alpha, \theta} \tilde{k}_{0}^{\alpha}=C_{\alpha} \psi$, and
(8) $A_{\chi}^{\alpha, \theta} \tilde{k}_{0}^{\alpha}=P_{\theta} C_{\theta}\left(\bar{\chi} \frac{\theta}{\alpha}\right)-\alpha(0) S_{\theta}^{*} \chi=S^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)+\left\langle\chi_{\alpha}, k_{0}^{\alpha}\right\rangle \tilde{k}_{0}^{\theta}-\alpha(0) S_{\theta}^{*} \chi$.

Proof. Item (1) follows from

$$
A_{\psi}^{\theta, \alpha} k_{0}^{\theta}=P_{\alpha}(\psi-\overline{\theta(0)} \theta \psi)=\psi .
$$

To prove item (2), we first calculate

$$
P_{\alpha}(\theta \bar{\chi})=P_{\alpha}\left(\alpha \frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}\right)+P_{\alpha}\left(\alpha \bar{\chi}_{\alpha}\right)=P_{\alpha}\left(\alpha \bar{\chi}_{\alpha}\right)=\alpha\left(\bar{\chi}_{\alpha}-\overline{\chi_{\alpha}(0)}\right) .
$$

Hence, by Proposition 2.1(3), we have

$$
A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta}=P_{\alpha} \bar{\chi}-\overline{\theta(0)} P_{\alpha}(\theta \bar{\chi})=\overline{\chi(0)} k_{0}^{\alpha}-\overline{\theta(0)} \alpha\left(\bar{\chi}_{\alpha}-\overline{\chi_{\alpha}(0)}\right)
$$

To prove item (3), note that, by Proposition 2.2,

$$
\begin{aligned}
A_{\psi}^{\theta, \alpha} \tilde{k}_{0}^{\theta} & =P_{\alpha}\left(\psi P_{\theta} \tilde{k}_{0}^{\theta}\right)=P_{\alpha}\left(\psi P_{\alpha} \tilde{k}_{0}^{\theta}\right)+P_{\alpha}\left(\psi \alpha P_{\frac{\theta}{\alpha}} \bar{\alpha} \tilde{k}_{0}^{\theta}\right)=P_{\alpha}\left(\psi P_{\alpha} \tilde{k}_{0}^{\theta}\right) \\
& =P_{\alpha}\left(\psi \frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}\right)=\frac{\theta}{\alpha}(0) P_{\alpha}(\psi \bar{z} \alpha)-\theta(0) P_{\alpha}(\psi \bar{z}) \\
& =\frac{\theta}{\alpha}(0) \psi(0) \tilde{k}_{0}^{\alpha}-\theta(0) S_{\alpha}^{*} \psi
\end{aligned}
$$

since

$$
\begin{aligned}
P_{\alpha} \psi \bar{z} \alpha & =\left(\alpha P^{-} \bar{\alpha}-P^{-}\right)(\psi \alpha \bar{z})=\alpha P^{-}(\psi \bar{z})-\psi(0) \alpha(0) \bar{z} \\
& =\alpha \psi(0) \bar{z}-\psi(0) \alpha(0) \bar{z}=\psi(0) \tilde{k}_{0}^{\alpha}
\end{aligned}
$$

To show item (4), we calculate

$$
A_{\bar{\chi}}^{\theta, \alpha} \tilde{k}_{0}^{\theta}=P_{\alpha}(\bar{\chi} \bar{z}(\theta-\theta(0)))=P_{\alpha}(\theta \bar{\chi} \bar{z})-\theta(0) P_{\alpha}(\bar{\chi} \bar{z})=P_{\alpha} C_{\theta} \chi
$$

and by Corollary 2.4,

$$
P_{\alpha}\left(C_{\theta} \chi\right)=C_{\alpha} P_{\alpha} \frac{\bar{\theta}}{\bar{\alpha}} \chi=C_{\alpha} \chi_{\alpha} .
$$

The equality in item (5) follows from

$$
A_{\bar{\psi}}^{\alpha, \theta} k_{0}^{\alpha}=P_{\theta}(\bar{\psi}(1-\overline{\alpha(0)} \alpha))=\overline{\psi(0)} k_{0}^{\theta}-\overline{\alpha(0)} P_{\theta}(\alpha \bar{\psi})
$$

and

$$
P_{\theta}(\alpha \bar{\psi})=\alpha \bar{\psi}-\theta P(\bar{\theta} \alpha \bar{\psi})=\alpha \bar{\psi}-\theta P \frac{\bar{\theta}}{\bar{\alpha}} \bar{\psi}=\alpha \bar{\psi}-\theta \frac{\bar{\theta}}{\alpha}(0) \psi(0)
$$

The equality in item (6) follows from

$$
A_{\chi}^{\alpha, \theta} k_{0}^{\alpha}=P_{\theta}(\chi(1-\overline{\alpha(0)} \alpha))=\chi-\overline{\alpha(0)} P_{\theta}(\alpha \chi)=\chi-\overline{\alpha(0)} \alpha \chi_{\frac{\theta}{\alpha}}
$$

and item (7) follows from

$$
A_{\bar{\psi}}^{\alpha, \theta} \tilde{k}_{0}^{\alpha}=P_{\theta}(\bar{\psi} \bar{z}(\alpha-\alpha(0)))=P_{\theta}(\alpha \bar{z} \bar{\psi})-\alpha(0) P_{\theta}(\bar{z} \bar{\psi})=P_{\theta} C_{\alpha} \psi=C_{\alpha} \psi
$$

Finally, item (8) follows from

$$
A_{\chi}^{\alpha, \theta} \tilde{k}_{0}^{\alpha}=P_{\theta}\left(\chi \tilde{k}_{0}^{\alpha}\right)=P_{\theta}(\alpha \chi \bar{z})-\alpha(0) P_{\theta}(\bar{z} \chi)=P_{\theta} C_{\theta}\left(\bar{\chi} \frac{\theta}{\alpha}\right)-\alpha(0) S_{\theta}^{*} \chi
$$

Since $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$, we have $\frac{\theta}{\alpha} \bar{\chi}=\frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}+\bar{\chi}_{\alpha}$; hence

$$
\begin{aligned}
P_{\theta}\left(\frac{\theta}{\alpha} \bar{\chi}\right) & =P_{\theta}\left(\frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}\right)+P_{\theta} \bar{\chi}_{\alpha}=\theta P^{-} \bar{\theta} \frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}+P_{\theta} \overline{\chi_{\alpha}(0)} \\
& =\theta\left(\bar{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}-\bar{\alpha}(0) \bar{\chi}_{\frac{\theta}{\alpha}}(0)\right)+\overline{\chi_{\alpha}(0)} k_{0}^{\theta} .
\end{aligned}
$$

Thus

$$
\begin{aligned}
P_{\theta} C_{\theta}\left(\frac{\theta}{\alpha} \bar{\chi}\right) & =C_{\theta} P_{\theta}\left(\frac{\theta}{\alpha} \bar{\chi}\right)=\bar{z}\left(\alpha \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right)+\chi_{\alpha}(0) C_{\theta} k_{0}^{\theta} \\
& =S^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)+\left\langle\chi_{\alpha}, k_{0}^{\alpha}\right\rangle \tilde{k}_{0}^{\theta} .
\end{aligned}
$$

## 4. The symbols of zero in $\mathcal{T}(\theta, \alpha)$

In what follows, for $f \in H^{2}$, let $f_{i}, f_{o}$ denote the inner and outer factors, respectively, in an inner-outer factorization $f=f_{i} f_{o}$. Recall that $f_{i}$ and $f_{o}$ are defined up to multiplication by a constant. For $\alpha$ and $\theta$ nonconstant inner functions, let $G C D(\theta, \alpha)$ denote the greatest common divisor of $\theta$ and $\alpha$, which is also defined up to a constant.

In this section, we study the symbols for which the corresponding asymmetric truncated Toeplitz operator is the zero operator. This is equivalent to the kernel of the operator being equal to $K_{\theta}^{2}$. We start by characterizing the kernels of some asymmetric truncated Toeplitz operators. The following result generalizes the disk versions of Theorem 7.2 in [8] and Theorem 3.2 in [7]. Note also that in Theorem 4.1, Lemma 4.2, and Corollary 4.3, we do not assume that $\alpha$ divides $\theta$.

Theorem 4.1. Let $\alpha$ and $\theta$ be nonconstant inner functions, and let $\varphi \in H^{\infty}$, $\varphi \neq 0$ with inner factor $\varphi_{i}$. Then
(1) $\operatorname{ker} A_{\varphi}^{\theta, \alpha}=K_{\theta}^{2} \cap \frac{\alpha}{\psi} H^{2}$ with $\psi=G C D\left(\alpha, \varphi_{i}\right)$; in particular, if $\frac{\alpha}{\psi} \leq \theta$, then $\operatorname{ker} A_{\varphi}^{\theta, \alpha}=\frac{\alpha}{\psi} K_{\frac{\theta \psi}{\alpha}}^{2} ;$ and
(2) if $\theta \leq \alpha$, then $\operatorname{ker} A_{\stackrel{\varphi}{\varphi}}^{\theta, \alpha}=\operatorname{ker}\left(A_{\varphi}^{\alpha, \theta}\right)^{*}=K_{G C D(\theta, \psi)}^{2}$ with $\psi=G C D\left(\alpha, \varphi_{i}\right)$.

Proof. (1) Let $f \in K_{\theta}^{2}$. Note that $f \in \operatorname{ker} A_{\varphi}^{\theta, \alpha}$ if and only if $\varphi f \in \alpha H^{2}$, which is equivalent to $\alpha$ dividing $\varphi_{i} f_{i}$, where $f_{i}$ is the inner factor of $f$. Since $\frac{\alpha}{\psi}$ and $\frac{\varphi_{i}}{\psi}$ are relatively prime, then $f_{i}$ is divisible by $\frac{\alpha}{\psi}$; hence $\operatorname{ker} A_{\varphi}^{\theta, \alpha}=\left\{f \in K_{\theta}^{2}: f \in\right.$ $\left.\frac{\alpha}{\psi} H^{2}\right\}=K_{\theta}^{2} \cap \frac{\alpha}{\psi} H^{2}$. If $\frac{\alpha}{\psi} \leq \theta$, then the result follows from the decomposition $K_{\theta}^{2}=K_{\frac{\alpha}{\psi}}^{2} \oplus \frac{\alpha}{\psi} K_{\frac{\theta \psi}{\alpha}}^{2}$.
(2) If $\theta \leq \alpha$, in which case $K_{\theta}^{2} \subset K_{\alpha}^{2}$, then we have

$$
\begin{aligned}
\operatorname{ker} A_{\bar{\varphi}}^{\theta, \alpha} & =\left\{f \in K_{\theta}^{2}: P_{\alpha}(\bar{\varphi} f)=0\right\}=\left\{f \in K_{\theta}^{2}: P(\bar{\varphi} f)=0\right\} \\
& =\left\{f \in K_{\theta}^{2}: \bar{\varphi} f \in \overline{H_{0}^{2}}\right\}=\left\{f \in K_{\theta}^{2}: \varphi \alpha \bar{z} \bar{f} \in \alpha H^{2}\right\} \\
& =\left\{f \in K_{\theta}^{2}: \varphi C_{\alpha} f \in \alpha H^{2}\right\}=\left\{f \in K_{\theta}^{2}: C_{\alpha} f \in \operatorname{ker} A_{\varphi}^{\alpha}=K_{\alpha}^{2} \cap \frac{\alpha}{\psi} H^{2}\right\}
\end{aligned}
$$

by (1). Now $C_{\alpha} f \in K_{\alpha}^{2} \cap \frac{\alpha}{\psi} H^{2}$ if and only if $f \in K_{\alpha}^{2} \cap \psi \overline{H_{0}^{2}}$, which is equivalent to $f \in K_{\alpha}^{2} \cap K_{\psi}^{2}=K_{\psi}^{2}$; hence $f \in K_{\theta}^{2} \cap K_{\psi}^{2}=K_{G C D(\theta, \psi)}^{2}$.

Lemma 4.2. Let $\alpha$ and $\theta$ be nonconstant inner functions, and let $\varphi \in H^{2}$. If we assume that $A_{\varphi}^{\alpha, \theta}$ is an asymmetric truncated Toeplitz operator, then $A_{\varphi}^{\alpha, \theta}=0$ if and only if $\varphi \in \theta H^{2}$.

Proof. If $\varphi \in \theta H^{2}$ and $f \in K_{\alpha}^{\infty}$, then $A_{\varphi}^{\alpha, \theta} f=P_{\theta}(\varphi f)=P_{\theta}\left(\theta \varphi_{1} f\right)=0$, where $\varphi=\theta \varphi_{1}$ and $\varphi_{1} \in H^{2}$. Since $K_{\alpha}^{\infty}$ is dense in $K_{\alpha}^{2}$, we have $A_{\varphi}^{\alpha, \theta}=0$.

For the converse implication, let $\psi=G C D\left(\theta, \varphi_{i}\right)$. If $A_{\varphi}^{\alpha, \theta}=0$ and $f \in K_{\alpha}^{\infty}$, then $\varphi f \in \theta H^{2}$; hence, $\theta$ divides $\varphi_{i} f_{i}$, and as a result $\frac{\varphi_{i}}{\psi} f_{i} \in \frac{\theta}{\psi} H^{2}$. Since $\frac{\varphi_{i}}{\psi}$ and $\frac{\theta}{\psi}$ are relatively prime, we have $f_{i} \in \frac{\theta}{\psi} H^{2}$. By the density of $K_{\alpha}^{\infty}$ in $K_{\alpha}^{2}$, we have $K_{\alpha}^{2} \subset \frac{\theta}{\psi} H^{2}$. Since $K_{\alpha}^{2}$ is the kernel of the Toeplitz operator $T_{\bar{\alpha}}$, then it is near invariant with respect to functions from $\overline{H^{\infty}}$; that is, if $g \in \operatorname{ker} T_{\bar{\alpha}}$ and $h \in H^{\infty}$, then $\bar{h} g \in H^{2}$ implies $\bar{h} g \in \operatorname{ker} T_{\bar{\alpha}}$. Following the proof of [6, Proposition 2.4], $K_{\alpha}^{2} \subset \frac{\theta}{\psi} H^{2}$ implies that, for every $g_{0} \in K_{\alpha}^{2}$, we have $\frac{\bar{\theta}}{\psi} g_{0}=g_{1} \in H^{2}$ so that $g_{1} \in K_{\alpha}^{2}$. Repeating this reasoning for $g_{1}$ we conclude that $g_{0}$ can be indefinitely divided by $\frac{\theta}{\psi}$, which is possible only if $\frac{\theta}{\psi}$ is a constant.

Corollary 4.3. If $\varphi \in H^{\infty}$, then $A_{\varphi}^{\theta, \alpha}=0$ if and only if $\varphi \in \alpha H^{\infty}$, and $A_{\bar{\varphi}}^{\theta, \alpha}=0$ if and only if $\varphi \in \theta H^{\infty}$.
Proof. If $A_{\varphi}^{\theta, \alpha}=0$, then, by the previous lemma, $\bar{\alpha} \varphi \in H^{2} \cap L^{\infty}=H^{\infty}$; thus $\varphi \in \alpha H^{\infty}$. The converse is obvious.

The next theorem establishes a necessary and sufficient condition for a bounded asymmetric truncated Toeplitz operator to be the zero operator in terms of its symbol.
Theorem 4.4. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$. If we let $A_{\varphi}^{\theta, \alpha}: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ be a bounded asymmetric truncated Toeplitz operator with $\varphi \in L^{2}$, then $A_{\varphi}^{\theta, \alpha}=0$ if and only if $\varphi \in \alpha H^{2}+\overline{\theta H^{2}}$.
Proof. If we assume that $\varphi=\alpha h_{1}+\bar{\theta} \overline{h_{2}}$ for some $h_{1}, h_{2} \in H^{2}$, and we take $f \in K_{\theta}^{\infty}$, then $\varphi f=\alpha h_{1} f+\bar{\theta} \overline{h_{2}} f$ and $P_{\alpha}\left(\alpha h_{1} f\right)=0$. Since $\bar{\theta} f \in \bar{\theta} K_{\theta}^{2} \subset \overline{z H^{2}}$, we have $\bar{\theta} f \bar{h}_{2} \in \overline{z H^{2}}$ and $P_{\alpha}\left(\bar{\theta} f \bar{h}_{2}\right)=0$; hence $A_{\varphi}^{\theta, \alpha}=0$ on a dense subset of $K_{\theta}^{2}$, which implies that $A_{\varphi}^{\theta, \alpha}=0$.

For the converse implication, let us take $\varphi=\psi+\bar{\chi} \in L^{2}$ with $\psi, \chi \in H^{2}$ such that $A_{\varphi}^{\theta, \alpha}=0$. Note that $A_{\psi}^{\theta, \alpha}=-A_{\bar{\chi}}^{\theta, \alpha}$; thus $A_{\psi}^{\alpha} f=A_{\psi}^{\theta, \alpha} f=-A_{\bar{\chi}}^{\theta, \alpha} f=-A_{\bar{\chi}}^{\alpha} f$ for all $f \in K_{\alpha}^{\infty}$ by Proposition 3.1. Hence $A_{\psi}^{\alpha}$ commutes on $K_{\alpha}^{\infty}$ not only with $S_{\alpha}$ but also with $S_{\alpha}^{*}$ (see [19] and Proposition 3.3). Consequently, since $k_{0}^{\alpha} \in K_{\alpha}^{\infty}$, $S_{\alpha}^{*} k_{0}^{\alpha}=-\overline{\alpha(0)} \tilde{k}_{0}^{\alpha} \in K_{\alpha}^{\infty}$ (cf. (2.3)), and $\left(I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}\right) k_{0}^{\alpha} \in K_{\alpha}^{\infty}$, it follows that

$$
\begin{aligned}
A_{\psi}^{\theta, \alpha}\left(I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}\right) k_{0}^{\alpha} & =\left(A_{\psi}^{\alpha}-S_{\alpha} A_{\psi}^{\alpha} S_{\alpha}^{*}\right) k_{0}^{\alpha}=\left(A_{\psi}^{\alpha}-S_{\alpha} S_{\alpha}^{*} A_{\psi}^{\alpha}\right) k_{0}^{\alpha} \\
& =\left(I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}\right) A_{\psi}^{\alpha} k_{0}^{\alpha}=\left(k_{0}^{\alpha} \otimes k_{0}^{\alpha}\right) A_{\psi}^{\alpha} k_{0}^{\alpha}=\left\langle A_{\psi}^{\alpha} k_{0}^{\alpha}, k_{0}^{\alpha}\right\rangle k_{0}^{\alpha}
\end{aligned}
$$

On the other hand, by Lemma 2.4 in [19],

$$
A_{\psi}^{\theta, \alpha}\left(I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}\right) k_{0}^{\alpha}=A_{\psi}^{\alpha}\left(k_{0}^{\alpha} \otimes k_{0}^{\alpha}\right) k_{0}^{\alpha}=\left(\left(A_{\psi}^{\alpha} k_{0}^{\alpha}\right) \otimes k_{0}^{\alpha}\right) k_{0}^{\alpha}=\left\langle k_{0}^{\alpha}, k_{0}^{\alpha}\right\rangle A_{\psi}^{\alpha} k_{0}^{\alpha}
$$

hence,

$$
\left\langle k_{0}^{\alpha}, k_{0}^{\alpha}\right\rangle A_{\psi}^{\alpha} k_{0}^{\alpha}=\left\langle A_{\psi}^{\alpha} k_{0}^{\alpha}, k_{0}^{\alpha}\right\rangle k_{0}^{\alpha},
$$

and it follows that there is $c \in \mathbb{C}$ such that

$$
\begin{equation*}
A_{\psi}^{\alpha} k_{0}^{\alpha}=c k_{0}^{\alpha} . \tag{4.1}
\end{equation*}
$$

Accordingly,

$$
0=\left(A_{\psi}^{\alpha}-c I_{K_{\alpha}^{2}}\right) k_{0}^{\alpha}=P_{\alpha}((\psi-c)(1-\overline{\alpha(0)} \alpha))=P_{\alpha}(\psi-c),
$$

which implies that $\psi-c \in \alpha H^{2}$.
Let us now consider $A_{\bar{\chi}}^{\theta, \alpha}$. We have $A_{\bar{\chi}}^{\theta, \alpha}=-A_{\psi}^{\theta, \alpha}$, and $A_{\psi}^{\theta, \alpha}$ intertwines $S_{\alpha}$ and $S_{\theta}$ on $K_{\theta}^{\infty}$ by Proposition 3.3. Thus $A_{\bar{\chi}}^{\theta, \alpha} S_{\theta} f=S_{\alpha} A_{\tilde{\chi}}^{\theta, \alpha} f$ for $f \in K_{\theta}^{\infty}$. In addition, $A_{\bar{\chi}}^{\theta, \alpha} S_{\theta}^{*} P_{\alpha}=A_{\bar{\chi}}^{\alpha} S_{\alpha}^{*} P_{\alpha}$ since $K_{\alpha}^{2}$ is invariant for both $S_{\theta}^{*}$ and $A_{\bar{\chi}}^{\theta, \alpha}$. Since $P_{\alpha} k_{0}^{\theta}=k_{0}^{\alpha} \in K_{\alpha}^{\infty} \subset K_{\theta}^{\infty}$, and, as above, $S_{\theta}^{*} P_{\alpha} k_{0}^{\theta} \in K_{\theta}^{\infty}$, we have

$$
\begin{aligned}
A_{\bar{\chi}}^{\theta, \alpha}\left(I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}\right) P_{\alpha} k_{0}^{\theta} & =\left(A_{\bar{\chi}}^{\theta, \alpha} P_{\alpha}-S_{\alpha} A_{\bar{\chi}}^{\theta, \alpha} S_{\theta}^{*} P_{\alpha}\right) k_{0}^{\theta} \\
& =\left(A_{\bar{\chi}}^{\theta, \alpha} P_{\alpha}-S_{\alpha} A_{\bar{\chi}}^{\alpha} S_{\alpha}^{*} P_{\alpha}\right) k_{0}^{\theta}=\left(A_{\bar{\chi}}^{\theta, \alpha} P_{\alpha}-S_{\alpha} S_{\alpha}^{*} A_{\bar{\chi}}^{\alpha} P_{\alpha}\right) k_{0}^{\theta} \\
& =\left(I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}\right)\left(A_{\bar{\chi}}^{\theta, \alpha}\right)_{\mid K_{\alpha}^{2}} P_{\alpha} k_{0}^{\theta}=\left(k_{0}^{\alpha} \otimes k_{0}^{\alpha}\right) A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\alpha} \\
& =\left\langle A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\alpha}, k_{0}^{\alpha}\right\rangle k_{0}^{\alpha}
\end{aligned}
$$

and

$$
\begin{aligned}
A_{\bar{\chi}}^{\theta, \alpha}\left(I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}\right) P_{\alpha} k_{0}^{\theta} & =A_{\bar{\chi}}^{\theta, \alpha}\left(k_{0}^{\theta} \otimes k_{0}^{\theta}\right) P_{\alpha} k_{0}^{\theta} \\
& =\left(A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta} \otimes k_{0}^{\theta}\right) k_{0}^{\alpha}=\left\langle k_{0}^{\alpha}, k_{0}^{\theta}\right\rangle A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta} .
\end{aligned}
$$

It then follows that $A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta}=c_{1} k_{0}^{\alpha}$.
Since $A_{\vec{\chi}}^{\theta, \alpha}=-A_{\psi}^{\theta, \alpha}$, by Proposition 3.3 the operator $A_{\widetilde{\chi}}^{\theta, \alpha}$ intertwines $S_{\theta}$ and $S_{\alpha}$ on $K_{\theta}^{\infty}$. By Proposition 2.5, we know that $S_{\theta}^{n} k_{0}^{\theta} \in K_{\theta}^{\infty}$ for $n \geq 0$; thus we have

$$
A_{\bar{\chi}}^{\theta, \alpha} S_{\theta}^{n} k_{0}^{\theta}=S_{\alpha}^{n} A_{\bar{\chi}}^{\theta, \alpha} k_{0}^{\theta}=c_{1} S_{\alpha}^{n} P_{\alpha} k_{0}^{\theta}=c_{1} P_{\alpha} S_{\theta}^{n} k_{0}^{\theta} .
$$

From the linear density of $S_{\theta}^{n} k_{0}^{\theta}, n \geq 0$ in $K_{\theta}^{2}$, we have

$$
\begin{equation*}
A_{\bar{\chi}}^{\theta, \alpha}=c_{1} P_{\alpha} \in \mathcal{T}(\theta, \alpha) . \tag{4.2}
\end{equation*}
$$

Since $A_{\bar{\chi}}^{\theta, \alpha}=-A_{\psi}^{\theta, \alpha}$, by (4.1) and (4.2) we then have that $c_{1}=-c$. Note that $\left(A_{\bar{\chi}+c}^{\theta, \alpha}\right)^{*}=A_{\chi+\bar{c}}^{\alpha, \theta}$, by Proposition 3.2. Using Lemma 4.2, we obtain that $A_{\chi+\bar{c}}^{\alpha, \theta}=0$ if and only if $\bar{\chi}+c \in \overline{\theta H^{2}}$; therefore $\varphi=(\psi-c)+(\bar{\chi}+c)$ is in $\alpha H^{2}+\overline{\theta H^{2}}$.
Corollary 4.5. Let $\alpha \leq \theta$ be nonconstant inner functions, and let $A_{\varphi}^{\theta, \alpha} \in \mathcal{T}(\theta, \alpha)$. For $\varphi \in L^{2}$ there are functions $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$ such that $A_{\varphi}^{\theta, \alpha}=A_{\psi+\bar{\chi}}^{\theta, \alpha}$. Moreover, $A_{\psi+\bar{\chi}}^{\theta, \alpha}=A_{\psi_{1}+\bar{\chi}_{1}}^{\theta, \alpha}$ if and only if $\psi_{1}=\psi+c k_{0}^{\alpha}$, $\chi_{1}=\chi-\bar{c} k_{0}^{\theta}$ for some constant $c$.

Proof. Let $\varphi=\varphi_{+}+\varphi_{-}$with $\varphi_{+} \in H^{2}$ and $\varphi_{-} \in \overline{H^{2}}$. If we put $\psi=P_{\alpha} \varphi_{+}$and $\chi=P_{\theta} \bar{\varphi}_{-}$, then $\varphi-\psi-\bar{\chi} \in \alpha H^{2}+\overline{\theta H^{2}}$; hence $A_{\varphi}^{\theta, \alpha}=A_{\psi+\bar{\chi}}^{\theta, \alpha}$ by Theorem 4.4.

The proof of the second statement is similar to the proof of the corollary on page 499 in [19]. Note first that, by Proposition 3.1, $A_{\varphi}^{\theta, \alpha}=A_{\psi+\bar{\chi}}^{\theta, \alpha}=A_{\psi_{1}+\bar{\chi}_{1}}^{\theta, \alpha}$. On the other hand, if $A_{\psi_{1}+\bar{\chi}_{1}}^{\theta, \alpha}=A_{\varphi}^{\theta, \alpha}$, then $\psi-\psi_{1}+\bar{\chi}-\bar{\chi}_{1} \in \alpha H^{2}+\overline{\theta H^{2}}$; hence $\psi-\psi_{1}+\bar{\chi}-\bar{\chi}_{1}=\alpha h_{1}+\bar{\theta} \bar{h}_{2}$ for some $h_{1}, h_{2} \in H^{2}$. Applying $P_{\alpha}$ and taking into account that $P_{\alpha} 1=k_{0}^{\alpha}$, we get $\psi-\psi_{1}=-P_{\alpha}\left(\bar{\chi}-\bar{\chi}_{1}\right)+P_{\alpha} \bar{\theta} \bar{h}_{2}=-c_{1} P_{\alpha} 1=-c_{1} k_{0}^{\alpha}$ for some constant $c_{1}$ since $P_{\alpha} \overline{H^{2}}$ contains only constant functions. Similarly, we
also have that $\bar{\psi}-\bar{\psi}_{1}+\chi-\chi_{1}=\bar{\alpha} \bar{h}_{1}+\theta h_{2}$, and, applying $P_{\theta}$, we obtain analogously that $\chi-\chi_{1}=c_{2} k_{0}^{\theta}$. Since we must have $A_{\psi+\bar{\chi}}^{\theta, \alpha}-A_{\psi_{1}+\bar{\chi}_{1}}^{\theta, \alpha}=A_{-c_{1} k_{0}^{\alpha}+\bar{c}_{2} \bar{k}_{0}^{\theta}}^{\theta, \alpha}=0$, by Proposition 3.1, $-c_{1} P_{\alpha}+\bar{c}_{2} P_{\alpha}=0$, which implies that $\bar{c}_{2}=c_{1}$.

The following properties can be immediately obtained from the previous results by taking adjoints.
Corollary 4.6. If we let $A_{\varphi}^{\alpha, \theta}: K_{\alpha}^{2} \rightarrow K_{\theta}^{2}, A_{\varphi}^{\alpha, \theta} \in \mathcal{T}(\alpha, \theta), \alpha \leq \theta, \varphi \in L^{2}$, then $A_{\varphi}^{\alpha, \theta}=0$ if and only if $\varphi \in \theta H^{2}+\overline{\alpha H^{2}}$.

Corollary 4.7. If we let $A_{\varphi}^{\alpha, \theta}: K_{\alpha}^{2} \rightarrow K_{\theta}^{2}, A_{\varphi}^{\alpha, \theta} \in \mathcal{T}(\alpha, \theta), \alpha \leq \theta, \varphi \in L^{2}$, then there are functions $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$ such that $A_{\varphi}^{\alpha, \theta}=A_{\bar{\psi}+\chi}^{\alpha, \theta}$.

## 5. First characterization in terms of rank 2 operators

In [19, Theorem 4.1], a characterization of truncated Toeplitz operators in $\mathcal{T}(\theta)$ was given by using certain rank 2 operators defined in terms of the kernel function $k_{0}^{\theta}$. Here we obtain an analogous result for asymmetric truncated Toeplitz operators $\mathcal{T}(\theta, \alpha)$ using the kernel functions $k_{0}^{\alpha}$ and $k_{0}^{\theta}$.

Theorem 5.1. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $A: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ be a bounded operator. Then $A \in \mathcal{T}(\theta, \alpha)$ if and only if there are $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$ such that

$$
\begin{equation*}
A-S_{\alpha} A S_{\theta}^{*}=\psi \otimes k_{0}^{\theta}+k_{0}^{\alpha} \otimes \chi \tag{5.1}
\end{equation*}
$$

Proof. If we assume that $A \in \mathcal{T}(\theta, \alpha)$, then $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$ for some $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$. Note that if $f \in K_{\theta}^{\infty}$, then $S_{\theta}^{*} f \in K_{\theta}^{\infty}$ (cf. (2.2)). Hence by Proposition 3.3,

$$
\begin{aligned}
\left(A_{\psi}^{\theta, \alpha}-S_{\alpha} A_{\psi}^{\theta, \alpha} S_{\theta}^{*}\right) f & =\left(A_{\psi}^{\theta, \alpha}-A_{\psi}^{\theta, \alpha} S_{\theta} S_{\theta}^{*}\right) f \\
& =\left(A_{\psi}^{\theta, \alpha}\left(I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}\right)\right) f=\left(A_{\psi}^{\theta, \alpha} k_{0}^{\theta} \otimes k_{0}^{\theta}\right) f \\
& =\left(P_{\alpha}(\psi(1-\overline{\theta(0)} \theta)) \otimes k_{0}^{\theta}\right) f=\left(\psi \otimes k_{0}^{\theta}\right) f .
\end{aligned}
$$

Moreover, since $A_{\bar{\chi}}^{\theta}$ commutes with $S_{\theta}^{*}$ on $K_{\theta}^{\infty}$ (see [19, p. 498]) and by Proposition 2.5, we have

$$
\begin{aligned}
\left(A_{\bar{\chi}}^{\theta, \alpha}-S_{\alpha} A_{\bar{\chi}}^{\theta, \alpha} S_{\theta}^{*}\right) f & =\left(A_{\bar{\chi}}^{\theta, \alpha}-S_{\alpha} P_{\alpha} A_{\bar{\chi}}^{\theta} S_{\theta}^{*}\right) f \\
& =\left(A_{\bar{\chi}}^{\theta, \alpha}-S_{\alpha} P_{\alpha} S_{\theta}^{*} A_{\bar{\chi}}^{\theta}\right) f=\left(P_{\alpha} A_{\bar{\chi}}^{\theta}-P_{\alpha} S_{\theta} S_{\theta}^{*} A_{\bar{\chi}}^{\theta}\right) f \\
& =P_{\alpha}\left(I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}\right) A_{\bar{\chi}}^{\theta} f=P_{\alpha}\left(k_{0}^{\theta} \otimes k_{0}^{\theta}\right) A_{\bar{\chi}}^{\theta} f \\
& =\left(\left(P_{\alpha} k_{0}^{\theta}\right) \otimes\left(A_{\chi}^{\theta} k_{0}^{\theta}\right)\right) f=\left(k_{0}^{\alpha} \otimes \chi\right) f .
\end{aligned}
$$

Since $K_{\theta}^{\infty}$ is dense in $K_{\theta}^{2}$, we obtain (5.1).
For the converse implication, note that, for $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$ and for all $f \in K_{\theta}^{\infty}$, $g \in K_{\alpha}^{\infty}$, we have

$$
\begin{equation*}
\left\langle A_{\psi+\bar{\chi}}^{\theta, \alpha} f, g\right\rangle=\sum_{n=0}^{\infty}\left(\left\langle f, S_{\theta}^{n} k_{0}^{\theta}\right\rangle\left\langle S_{\alpha}^{n} \psi, g\right\rangle+\left\langle f, S_{\theta}^{n} \chi\right\rangle\left\langle S_{\alpha}^{n} k_{0}^{\alpha}, g\right\rangle\right) . \tag{5.2}
\end{equation*}
$$

Indeed, we obtain (5.2) as in the proof of [19, Lemma 4.2] by taking $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$ instead of $\psi, \chi \in K_{u}^{2}$ and $S_{\alpha}, S_{\theta}^{*}$ instead of $S_{u}, S_{u}^{*}$, respectively. The convergence $\left\langle A_{\psi+\bar{\chi}}^{\theta, \alpha}\left(S_{\theta}^{*}\right)^{N+1} f, S_{\alpha}^{* N+1} g\right\rangle \rightarrow 0$ as $N \rightarrow \infty$ follows from

$$
\begin{aligned}
\left\langle S_{\alpha}^{N} A_{\psi+\bar{\chi}}^{\theta, \alpha}\left(S_{\theta}^{*}\right)^{N} f, g\right\rangle & =\left\langle S_{\alpha}^{N} A_{\psi}^{\theta, \alpha}\left(S_{\theta}^{*}\right)^{N} f, g\right\rangle+\left\langle S_{\alpha}^{N} A_{\bar{\chi}}^{\theta, \alpha}\left(S_{\theta}^{*}\right)^{N} f, g\right\rangle \\
& =\left\langle\left(S_{\theta}^{*}\right)^{N} f, A_{\bar{\psi}}^{\alpha, \theta}\left(S_{\alpha}^{*}\right)^{N} g\right\rangle+\left\langle S_{\alpha}^{N} P_{\alpha} A_{\bar{\chi}}^{\theta}\left(S_{\theta}^{*}\right)^{N} f, g\right\rangle \\
& =\left\langle\left(S_{\theta}^{*}\right)^{N} f,\left(S_{\theta}^{*}\right)^{N} A_{\bar{\psi}}^{\alpha, \theta} g\right\rangle+\left\langle\left(S_{\theta}^{*}\right)^{N} A_{\bar{\chi}}^{\theta} f, P_{\alpha}\left(S_{\alpha}^{*}\right)^{N} g\right\rangle \\
& =\left\langle\left(S_{\theta}^{*}\right)^{N} f,\left(S_{\theta}^{*}\right)^{N} A_{\bar{\psi}}^{\alpha, \theta} g\right\rangle+\left\langle\left(S_{\theta}^{*}\right)^{N} A_{\bar{\chi}}^{\theta} f,\left(S_{\alpha}^{*}\right)^{N} g\right\rangle,
\end{aligned}
$$

where the last expression tends to zero, when $N \rightarrow \infty$, by the strong convergence $\left(S^{*}\right)^{N} \rightarrow 0$.

Now assume that a bounded operator $A: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ satisfies (5.1). For any positive integer $n$ we have

$$
S_{\alpha}^{n} A S_{\theta}^{* n}-S_{\alpha}^{n+1} A S_{\theta}^{*(n+1)}=\left(S_{\alpha}^{n} \psi \otimes S_{\theta}^{n} k_{0}^{\theta}\right)+\left(S_{\alpha}^{n} k_{0}^{\alpha} \otimes S_{\theta}^{n} \chi\right)
$$

Adding for $n=0,1, \ldots, N$, we obtain

$$
A=\sum_{n=0}^{N}\left(S_{\alpha}^{n} \psi \otimes S_{\theta}^{n} k_{0}^{\theta}+S_{\alpha}^{n} k_{0}^{\alpha} \otimes S_{\theta}^{n} \chi\right)+S_{\alpha}^{N+1} A S_{\theta}^{*(N+1)}
$$

Taking into account that $S_{\theta}^{* n} \rightarrow 0$ in the strong operator topology, we get

$$
\begin{equation*}
A=\sum_{n=0}^{\infty}\left(S_{\alpha}^{n} \psi \otimes S_{\theta}^{n} k_{0}^{\theta}+S_{\alpha}^{n} k_{0}^{\alpha} \otimes S_{\theta}^{n} \chi\right) \tag{5.3}
\end{equation*}
$$

hence, by comparing (5.2) and (5.3), we conclude that the right-hand side of (5.3) is equal to $A_{\psi+\bar{\chi}}^{\theta, \alpha}$.

We can obtain a similar characterization for operators from $\mathcal{T}(\alpha, \theta)$ by taking adjoints in (5.1). Namely, we have the following.

Corollary 5.2. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $A: K_{\alpha}^{2} \rightarrow K_{\theta}^{2}$ be a bounded operator. Then $A \in \mathcal{T}(\alpha, \theta)$ if and only if there are $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$ such that

$$
A-S_{\theta} A S_{\alpha}^{*}=k_{0}^{\theta} \otimes \psi+\chi \otimes k_{0}^{\alpha}
$$

## 6. SECOND CHARACTERIZATION IN TERMS OF RANK 2 OPERATORS

Sarason also obtained a characterization for truncated Toeplitz operators belonging to $\mathcal{T}(\theta)$ involving the function $\tilde{k}_{0}^{\theta}=C_{\theta} k_{0}^{\theta}$ instead of $k_{0}^{\theta}$ by a simple application of the conjugation $C_{\theta}$ to the result of Theorem 5.1 in the case $\alpha=\theta$. Here we will show that an analogous result holds for operators belonging to $\mathcal{T}(\theta, \alpha)$, $\alpha \leq \theta$. We cannot, however, use the same reasoning for $\alpha \neq \theta$ because, as we will see, the case of asymmetric truncated Toeplitz operators is more complex. The relation between a symbol of an asymmetric truncated Toeplitz operator and a rank 2 operator appearing in (6.1) is more involved.

Theorem 6.1. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $A: K_{\theta}^{2} \rightarrow K_{\alpha}^{2}$ be a bounded operator. Then $A \in \mathcal{T}(\theta, \alpha)$ if and only if there are $\mu \in K_{\alpha}^{2}$ and $\nu \in K_{\theta}^{2}$ such that

$$
\begin{equation*}
A-S_{\alpha}^{*} A S_{\theta}=\mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu \tag{6.1}
\end{equation*}
$$

Moreover, if $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$ with $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$, then $A$ satisfies (6.1) with

$$
\begin{equation*}
\mu=C_{\alpha} P_{\alpha}\left(\frac{\bar{\theta}}{\bar{\alpha}} \chi\right), \quad \nu=C_{\alpha} \psi+S^{*}\left(\alpha P_{\frac{\theta}{\alpha}} \chi\right) \tag{6.2}
\end{equation*}
$$

Proof. Assume first that $A \in \mathcal{T}(\theta, \alpha)$. By Theorem 5.1, there are $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$ such that

$$
A-S_{\alpha} A S_{\theta}^{*}=\psi \otimes k_{0}^{\theta}+k_{0}^{\alpha} \otimes \chi
$$

hence

$$
S_{\alpha}^{*}\left(A-S_{\alpha} A S_{\theta}^{*}\right) S_{\theta}=S_{\alpha}^{*} \psi \otimes S_{\theta}^{*} k_{0}^{\theta}+S_{\alpha}^{*} k_{0}^{\alpha} \otimes S_{\theta}^{*} \chi
$$

Using $I_{K_{\alpha}^{2}}-S_{\alpha}^{*} S_{\alpha}=\tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\alpha}$ and $I_{K_{\theta}^{2}}-S_{\theta}^{*} S_{\theta}=\tilde{k}_{0}^{\theta} \otimes \tilde{k}_{0}^{\theta}$, after some calculations we get

$$
\begin{aligned}
A-S_{\alpha}^{*} A S_{\theta}= & \tilde{k}_{0}^{\alpha} \otimes A^{*} \tilde{k}_{0}^{\alpha}+A \tilde{k}_{0}^{\theta} \otimes \tilde{k}_{0}^{\theta}-\left\langle A \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\theta} \\
& -S_{\alpha}^{*} \psi \otimes S_{\theta}^{*} k_{0}^{\theta}-S_{\alpha}^{*} k_{0}^{\alpha} \otimes S_{\theta}^{*} \chi
\end{aligned}
$$

Thus, by using (2.3), we obtain

$$
\begin{align*}
A-S_{\alpha}^{*} A S_{\theta}= & \left(A \tilde{k}_{0}^{\theta}-\left\langle A \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle \tilde{k}_{0}^{\alpha}+\theta(0) S_{\alpha}^{*} \psi\right) \otimes \tilde{k}_{0}^{\theta} \\
& +\tilde{k}_{0}^{\alpha} \otimes\left(A^{*} \tilde{k}_{0}^{\alpha}+\alpha(0) S_{\theta}^{*} \chi\right) \tag{6.3}
\end{align*}
$$

For the reverse implication, if we assume that a bounded operator $A$ satisfies the equality (6.1) for $\mu \in K_{\alpha}^{2}, \nu \in K_{\theta}^{2}$, then

$$
S_{\alpha} A S_{\theta}^{*}-S_{\alpha} S_{\alpha}^{*} A S_{\theta} S_{\theta}^{*}=S_{\alpha} \mu \otimes S_{\theta} \tilde{k}_{0}^{\theta}+S_{\alpha} \tilde{k}_{0}^{\alpha} \otimes S_{\theta} \nu
$$

By using $I_{K_{\alpha}^{2}}-S_{\alpha} S_{\alpha}^{*}=k_{0}^{\alpha} \otimes k_{0}^{\alpha}, I_{K_{\theta}^{2}}-S_{\theta} S_{\theta}^{*}=k_{0}^{\theta} \otimes k_{0}^{\theta}$ and (2.3), we obtain

$$
\begin{aligned}
A-S_{\alpha} A S_{\theta}^{*}= & A k_{0}^{\theta} \otimes k_{0}^{\theta}+k_{0}^{\alpha} \otimes A^{*} k_{0}^{\alpha}-\left\langle A k_{0}^{\theta}, k_{0}^{\alpha}\right\rangle k_{0}^{\alpha} \otimes k_{0}^{\theta} \\
& +\overline{\theta(0)} S_{\alpha} \mu \otimes k_{0}^{\theta}+\alpha(0) k_{0}^{\alpha} \otimes S_{\theta} \nu
\end{aligned}
$$

hence

$$
A-S_{\alpha} A S_{\theta}^{*}=\left(A k_{0}^{\theta}-\left\langle A k_{0}^{\theta}, k_{0}^{\alpha}\right\rangle k_{0}^{\alpha}+\overline{\theta(0)} S_{\alpha} \mu\right) \otimes k_{0}^{\theta}+k_{0}^{\alpha} \otimes\left(A^{*} k_{0}^{\alpha}+\overline{\alpha(0)} S_{\theta} \nu\right)
$$

which, by Theorem 5.1, implies that $A \in \mathcal{T}(\theta, \alpha)$.
If $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$ with $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$, then we can obtain formulas for $\mu$ and $\nu$ in terms of $\psi$ and $\chi$ by using Proposition 3.6 and equality (6.3). Namely, we can take

$$
\begin{aligned}
\mu & =A_{\psi+\bar{\chi}}^{\theta, \alpha} \tilde{k}_{0}^{\theta}-\left\langle A_{\psi}^{\theta, \alpha} \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle \tilde{k}_{0}^{\alpha}+\theta(0) S_{\alpha}^{*} \psi, \quad \text { and } \\
\nu & =A_{\bar{\psi}+\chi}^{\alpha, \theta} \tilde{k}_{0}^{\alpha}-\overline{\left\langle A_{\bar{\chi}}^{\theta, \alpha} \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle} \tilde{k}_{0}^{\theta}+\alpha(0) S_{\theta}^{*} \chi
\end{aligned}
$$

Since $\psi \in K_{\alpha}^{2}$, we have $\alpha \bar{\psi} \in H_{0}^{2}$; thus $(\alpha \bar{\psi})(0)=0$, and, since $C_{\alpha} S_{\alpha}^{*}=S_{\alpha} C_{\alpha}$ (see [19, Lemma 2.1]), we then get

$$
\begin{aligned}
\left\langle S_{\alpha}^{*} \psi, \tilde{k}_{0}^{\alpha}\right\rangle & =\left\langle k_{0}^{\alpha}, C_{\alpha} S_{\alpha}^{*} \psi\right\rangle=\left\langle k_{0}^{\alpha}, S_{\alpha} C_{\alpha} \psi\right\rangle=\left\langle k_{0}^{\alpha}, S_{\alpha} \alpha \bar{z} \bar{\psi}\right\rangle \\
& =\left\langle k_{0}^{\alpha}, \alpha \bar{\psi}\right\rangle=\langle 1-\overline{\alpha(0)} \alpha, \alpha \bar{\psi}\rangle=\langle 1, \alpha \bar{\psi}\rangle-\overline{\alpha(0)}\langle 1, \bar{\psi}\rangle=-\overline{\alpha(0)} \psi(0)
\end{aligned}
$$

Thus by (3) and (4) in Proposition 3.6 we have

$$
\begin{aligned}
\left\langle A_{\psi}^{\theta, \alpha} \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle & =\left\langle\frac{\theta}{\alpha}(0) \psi(0) \tilde{k}_{0}^{\alpha}-\theta(0) S_{\alpha}^{*} \psi, \tilde{k}_{0}^{\alpha}\right\rangle=\frac{\theta}{\alpha}(0) \psi(0)\left\|\tilde{k}_{0}^{\alpha}\right\|^{2}+\theta(0) \overline{\alpha(0)} \psi(0) \\
& =\frac{\theta}{\alpha}(0) \psi(0)\left(1-|\alpha(0)|^{2}\right)+\theta(0) \overline{\alpha(0)} \psi(0)=\frac{\theta}{\alpha}(0) \psi(0)
\end{aligned}
$$

Taking into account that $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$, where $\chi_{\frac{\theta}{\alpha}} \in K_{\frac{\theta}{\alpha}}^{2}$ and $\chi_{\alpha} \in K_{\alpha}^{2}$, we have

$$
\left\langle A_{\bar{\chi}}^{\theta, \alpha} \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle=\left\langle C_{\alpha} P_{\alpha}\left(\frac{\bar{\theta}}{\bar{\alpha}} \chi\right), \tilde{k}_{0}^{\alpha}\right\rangle=\left\langle k_{0}^{\alpha}, \chi_{\alpha}\right\rangle,
$$

and it then follows that

$$
\mu=C_{\alpha} \chi_{\alpha} \in K_{\alpha}^{2}
$$

Moreover, taking into account (7) and (8) in Proposition 3.6, it follows that

$$
\nu=C_{\alpha} \psi+S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right) \in K_{\theta}^{2}
$$

By taking adjoints in (6.1) we obtain the following similar characterization for operators from $\mathcal{T}(\alpha, \theta)$.

Corollary 6.2. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $A: K_{\alpha}^{2} \rightarrow K_{\theta}^{2}$ be a bounded operator. Then $A \in \mathcal{T}(\alpha, \theta)$ if and only if there are $\mu \in K_{\alpha}^{2}, \nu \in K_{\theta}^{2}$ such that

$$
A-S_{\theta}^{*} A S_{\alpha}=\tilde{k}_{0}^{\theta} \otimes \mu+\nu \otimes \tilde{k}_{0}^{\alpha} .
$$

It is clear that, if an asymmetric truncated Toeplitz operator $A$ satisfies equation (6.1) with some $\mu, \nu$, then that equation is also satisfied if $\mu, \nu$ are replaced by

$$
\begin{equation*}
\mu^{\prime}=\mu+\bar{b} \tilde{k}_{0}^{\alpha}, \quad \nu^{\prime}=\nu-b \tilde{k}_{0}^{\theta} \tag{6.4}
\end{equation*}
$$

respectively, for any $b \in \mathbb{C}$. On the other hand, it is also true that the symbol of $A=A_{\psi+\bar{\chi}}^{\theta, \alpha} \in \mathcal{T}(\theta, \alpha)$ is not unique, and by Corollary 4.5 we can replace $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$ with

$$
\begin{equation*}
\psi^{\prime}=\psi+c k_{0}^{\alpha} \in K_{\alpha}^{2}, \quad \chi^{\prime}=\chi-\bar{c} k_{0}^{\theta} \in K_{\theta}^{2} \tag{6.5}
\end{equation*}
$$

respectively, for any $c \in \mathbb{C}$. Using (6.2), it is easy to see that the following relation between the freedom of choice holds: of $\mu, \nu$ on the one hand and of $\psi, \chi$ on the other.

Corollary 6.3. Let $\mu \in K_{\alpha}^{2}$ and $\nu \in K_{\theta}^{2}$ be defined by (6.2) for given $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$, and let $\mu^{\prime} \in K_{\alpha}^{2}$ and $\nu^{\prime} \in K_{\theta}^{2}$ be defined analogously for $\psi^{\prime} \in K_{\alpha}^{2}$ and $\chi^{\prime} \in K_{\theta}^{2}$. If (6.5) holds, then

$$
\mu^{\prime}=\mu-c \frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}, \quad \nu^{\prime}=\nu+\bar{c} \overline{\frac{\theta}{\alpha}(0)} \tilde{k}_{0}^{\theta} .
$$

The examples below illustrate the result of Theorem 6.1 in the case of Toeplitz matrices.
Example 6.4. Let us consider $\alpha=z^{2}, \theta=z^{5}$, and a Toeplitz operator $A=A_{\psi+\bar{\chi}}^{z^{5}}$. If we assume that $\psi=a_{0}+a_{1} z$ and that $\chi=\bar{b}_{0}+\bar{b}_{-1} z+\bar{b}_{-2} z^{2}+\bar{b}_{-3} z^{3}+\bar{b}_{-4} z^{4}=$ $\left(\bar{b}_{0}+\bar{b}_{-1} z+\bar{b}_{-2} z^{2}\right)+z^{3}\left(\bar{b}_{-3}+\bar{b}_{-4} z\right)$, then $C_{z^{2}} \psi=\bar{a}_{1}+\bar{a}_{0} z, C_{z^{2}} P_{z^{2}} \bar{z}^{3} \chi=b_{-4}+b_{-3} z$, and $S^{*}\left(z^{2}\left(\bar{b}_{0}+\bar{b}_{-1} z+\bar{b}_{-2} z^{2}\right)\right)=\bar{b}_{0} z+\bar{b}_{-1} z^{2}+\bar{b}_{-2} z^{3}$. Note that $A-S_{z^{2}}^{*} A S_{z^{5}}$ has a matrix representation

$$
\left(\begin{array}{ccccc}
0 & 0 & 0 & 0 & b_{-4} \\
a_{1} & a_{0}+b_{0} & b_{-1} & b_{-2} & b_{-3}
\end{array}\right)
$$

which can be expressed as

$$
\left(b_{-4}+b_{-3} z\right) \otimes z^{4}+z \otimes\left(\bar{a}_{1}+\left(\bar{a}_{0}+\bar{b}_{0}\right) z+\bar{b}_{-1} z^{2}+\bar{b}_{-2} z^{3}\right) .
$$

On the other hand, let $A-S_{z^{2}}^{*} A S_{z^{5}}$ have a matrix representation

$$
\left(\begin{array}{ccccc}
0 & 0 & 0 & 0 & b_{0} \\
a_{0} & a_{1} & a_{2} & a_{3} & a_{4}+b_{1}
\end{array}\right)
$$

which can be expressed as

$$
\mu \otimes z^{4}+z \otimes \nu=\left(b_{0}+b_{1} z\right) \otimes z^{4}+z \otimes\left(\bar{a}_{0}+\bar{a}_{1} z+\bar{a}_{2} z^{2}+\bar{a}_{3} z^{3}+\bar{a}_{4} z^{4}\right)
$$

Note that if $\nu=\nu_{z^{2}}+z^{2} \nu_{z^{3}}=\left(\bar{a}_{0}+\bar{a}_{1} z\right)+z^{2}\left(\bar{a}_{2}+\bar{a}_{3} z+\bar{a}_{4} z^{2}\right)$, then $\psi=C_{z^{2}} P_{z^{2}} \nu=$ $a_{1}+a_{0} z$, and $\chi=\bar{a}_{2} z+\bar{a}_{3} z^{2}+\left(\bar{b}_{1}+\bar{a}_{4}\right) z^{3}+\bar{b}_{0} z^{4}$. Hence by Theorem 5.1 we have

$$
\begin{equation*}
A-S_{z^{2}} A S_{z^{5}}^{*}=\left(a_{1}+a_{0} z\right) \otimes 1+1 \otimes\left(\bar{a}_{2} z+\bar{a}_{3} z^{2}+\left(\bar{b}_{1}+\bar{a}_{4}\right) z^{3}+\bar{b}_{0} z^{4}\right) \tag{6.6}
\end{equation*}
$$

Requiring that $\nu_{z^{3}}$ be orthogonal to $z^{2}$ (see the proof of Theorem 8.5) determines that $a_{4}=0$.

On the other hand, we have some freedom in defining $\psi$ and $\chi$, namely, that $\psi_{1}=s+a_{0} z$ and $\chi_{1}=\bar{t}+\bar{a}_{2} z+\bar{a}_{3} z^{2}+\left(\bar{b}_{1}+\bar{a}_{4}\right) z^{3}+\bar{b}_{0} z^{4}$ also satisfy (6.6) if we assume that $t+s=a_{1}$.

Example 6.5. Let us now take $\alpha=z^{3}, \theta=z^{3}((\lambda-z) /(1-\bar{\lambda} z))^{2}, \lambda \in \mathbb{D}$, and let us consider the operator $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$, where $\psi=a_{0}+a_{1} z+a_{2} z^{2} \in K_{\alpha}^{2}$, and $\chi=\left(\bar{b}_{0}+\bar{b}_{1} z+\bar{b}_{2} z^{2}+\bar{b}_{3} z^{3}+\bar{b}_{4} z^{4}\right)(1-\bar{\lambda} z)^{-2} \in K_{\theta}^{2}($ see [12, Corollary 5.7.3]). Then by Theorem 6.1

$$
A-S_{\alpha}^{*} A S_{\theta}=\mu \otimes\left(\lambda^{2} z^{2}-2 \lambda z^{3}+z^{4}\right)(1-\bar{\lambda} z)^{-2}+z^{2} \otimes \nu
$$

where $\mu=b_{4}+\left(b_{3}+2 \bar{\lambda} b_{4}\right) z+\left(b_{2}+3 \bar{\lambda}^{2} b_{4}+\bar{\lambda} b_{3}\right) z^{2}$ and $\nu=\left(\bar{a}_{2}+\left(\bar{a}_{1}-2 \bar{\lambda} \bar{a}_{2}\right) z+\right.$ $\left.\left(\bar{b}_{0}+\bar{a}_{0}-2 \bar{\lambda} \bar{a}_{1}+\bar{\lambda} \bar{a}_{2}\right) z^{2}+\left(\bar{b}_{1}+\bar{\lambda}^{2} \bar{a}_{1}-2 \bar{\lambda} \bar{a}_{0}\right) z^{3}+\bar{\lambda}^{2} \bar{a}_{0} z^{4}\right)(1-\bar{\lambda} z)^{-2}$.

## 7. Characterizations in terms of Rank 1 operators

Our aim now is to describe the classes of symbols of an operator $A \in \mathcal{T}(\theta, \alpha)$ for which the right-hand side of (6.1) is a rank 1 operator. The corresponding question regarding the equation (5.1) is trivial by Corollary 4.5, since the right-hand side of (5.1) is a rank 1 operator if and only if $\psi=c \cdot k_{0}^{\alpha}$ or $\chi=c \cdot k_{0}^{\theta}$ with $c \in \mathbb{C}$. In the case $\alpha=\theta$ the question regarding the equality (6.1) also has an easy answer, since the relation between the symbols in (5.1) and (6.1) is $\psi=C_{\theta} \nu$ and $\chi=C_{\theta} \mu$. For $\alpha \neq \theta$ we need the following lemma.

Lemma 7.1. Let $\alpha \leq \theta$ be nonconstant inner functions, and let $A_{\psi+\bar{\chi}}^{\theta, \alpha} \in \mathcal{T}(\theta, \alpha)$. Assume that $\psi \in K_{\alpha}^{2}$ and that $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha} \in K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \frac{\theta}{\alpha} K_{\alpha}^{2}$. Then
(1) $P_{\widetilde{C} \tilde{k}_{0}^{\theta}} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)=-\overline{\frac{\theta}{\alpha}(0)} \chi_{\frac{\theta}{\alpha}}(0)\left\|\tilde{k}_{0}^{\theta}\right\|^{-2} \tilde{k}_{0}^{\theta}$,
(2) $P_{\widetilde{C} \tilde{k}_{0}^{\theta}} C_{\alpha} \psi=\overline{\frac{\theta}{\alpha}(0) \psi(0)}\left\|\tilde{k}_{0}^{\theta}\right\|^{-2} \tilde{k}_{0}^{\theta}$,
(3) $P_{\alpha} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)=P_{\widetilde{C} \tilde{k}_{0}^{\alpha}} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)=\chi_{\frac{\theta}{\alpha}}(0) \tilde{k}_{0}^{\alpha}$,
(4) $\left(P_{\theta}-P_{\alpha}\right) S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)=\alpha S^{*} \chi_{\frac{\theta}{\alpha}}$,
(5) $P_{\widetilde{C} \tilde{k}_{0}^{\alpha}} C_{\alpha} P_{\alpha}(\chi \overline{\bar{\alpha}})=\overline{\chi_{\alpha}(0)}\left\|\tilde{k}_{0}^{\alpha}\right\|^{-2} \tilde{k}_{0}^{\alpha}$, and
(6) $P_{\widetilde{\mathbb{E}} \tilde{k}_{0}^{\alpha}} C_{\alpha} \psi=\overline{\psi(0)}\left\|\tilde{k}_{0}^{\alpha}\right\|^{-2} \tilde{k}_{0}^{\alpha}$.

Proof. To prove (1), it is enough to calculate

$$
\begin{aligned}
\left\langle S^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right), \tilde{k}_{0}^{\theta}\right\rangle & =\left\langle\bar{z}\left(\alpha \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right), \bar{z}(\theta-\theta(0))\right\rangle \\
& =\left\langle\chi_{\frac{\theta}{\alpha}}, \frac{\theta}{\alpha}\right\rangle-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\langle 1, \theta\rangle-\overline{\theta(0)}\left\langle\alpha \chi_{\frac{\theta}{\alpha}}, 1\right\rangle+\alpha(0) \overline{\theta(0)} \chi_{\frac{\theta}{\alpha}}(0) \\
& =-\frac{\theta}{\alpha}(0) \chi_{\frac{\theta}{\alpha}}(0)
\end{aligned}
$$

To show (2), note that by Proposition 2.2

$$
\begin{aligned}
\left\langle C_{\alpha} \psi, \tilde{k}_{0}^{\theta}\right\rangle & =\left\langle C_{\alpha} \psi, P_{\alpha} \tilde{k}_{0}^{\theta}\right\rangle=\left\langle C_{\alpha} \psi, \frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}\right\rangle \\
& =\left\langle C_{\alpha} \frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}, \psi\right\rangle=\overline{\frac{\theta}{\alpha}(0)}\left\langle k_{0}^{\alpha}, \psi\right\rangle=\overline{\frac{\theta}{\alpha}(0)} \overline{\psi(0)}
\end{aligned}
$$

The equalities in (3) follow from

$$
\begin{aligned}
P_{\alpha} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right) & =P_{\alpha}\left(\bar{z} \alpha \chi_{\frac{\theta}{\alpha}}\right)=\alpha P^{-} \bar{\alpha} \bar{z} \alpha \chi_{\frac{\theta}{\alpha}}-P^{-} \bar{z} \alpha \chi_{\frac{\theta}{\alpha}}=\alpha P^{-} \bar{z} \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0) \bar{z} \\
& =\alpha \chi_{\frac{\theta}{\alpha}}(0) \bar{z}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0) \bar{z}=\chi_{\frac{\theta}{\alpha}}(0) \tilde{k}_{0}^{\alpha}=P_{\mathbb{C} \tilde{k}_{0}^{\alpha}} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}} .\right.
\end{aligned}
$$

To see (4) it is enough to calculate

$$
\begin{aligned}
\left(P_{\theta}-P_{\alpha}\right) S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right) & =\alpha P_{\frac{\theta}{\alpha}} \bar{\alpha} S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)=\alpha P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \bar{z}\left(\alpha \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right)\right) \\
& =\alpha P_{\frac{\theta}{\alpha}}\left(\bar{z} \chi_{\frac{\theta}{\alpha}}\right)=\alpha S_{\frac{\theta}{\alpha}}^{*} \chi_{\frac{\theta}{\alpha}}=\alpha S^{*} \chi_{\frac{\theta}{\alpha}} .
\end{aligned}
$$

The proof of (5) follows from $\left\langle C_{\alpha} P_{\alpha}(\chi \overline{\bar{\alpha}}), \tilde{k}_{0}^{\alpha}\right\rangle=\left\langle k_{0}^{\alpha}, \chi_{\alpha}\right\rangle=\overline{\chi_{\alpha}(0)}$, and, to show that (6) holds, note that $\left\langle C_{\alpha} \psi, \tilde{k}_{0}^{\alpha}\right\rangle=\left\langle k_{0}^{\alpha}, \psi\right\rangle=\overline{\psi(0)}$.

Theorem 7.2. Let $\alpha \leq \theta$ be nonconstant inner functions, and let $A_{\psi+\bar{\chi}}^{\theta, \alpha} \in$ $\mathcal{T}(\theta, \alpha)$, where $\psi \in K_{\alpha}^{2}$, and $\chi \in K_{\theta}^{2}$. Then
(1) $A_{\psi+\bar{\chi}}^{\theta, \alpha}-S_{\alpha}^{*} A_{\psi+\bar{\chi}}^{\theta, \alpha} S_{\theta}=\mu \otimes \tilde{k}_{0}^{\theta}$ for $\mu \in K_{\alpha}^{2}$ if and only if there is $s \in \mathbb{C}$ such

$$
\text { that } \psi=s k_{0}^{\alpha}, P_{\frac{\theta}{\alpha}} \chi=-\bar{s} k_{0}^{\frac{\theta}{\alpha}}, \text { and }
$$

(2) $A_{\psi+\bar{\chi}}^{\theta, \alpha}-S_{\alpha}^{*} A_{\psi+\bar{\chi}}^{\theta, \alpha} S_{\theta}=\tilde{k}_{0}^{\alpha} \otimes \nu$ for $\nu \in K_{\theta}^{2}$ if and only if $P_{\alpha}\left(\chi \frac{\bar{\theta}}{\bar{\alpha}}\right)=$ const $\cdot k_{0}^{\alpha}$. Proof. Assume first that $A_{\psi+\bar{\chi}}^{\theta, \alpha}-S_{\alpha}^{*} A_{\psi+\bar{\chi}}^{\theta, \alpha} S_{\theta}=\mu \otimes \tilde{k}_{0}^{\theta}$ for $\mu \in K_{\alpha}^{2}$. Now, the right-hand side of the equation (6.1) reduces to $\mu \otimes \tilde{k}_{0}^{\theta}$ if and only if $\nu=c \cdot \tilde{k}_{0}^{\theta}$ with $c \in \mathbb{C}$, which is equivalent to $\nu-P_{\mathbb{C} \tilde{k}_{0}^{\theta}} \nu=0$. Let $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha} \in K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \frac{\theta}{\alpha} K_{\alpha}^{2}$. By formulas (6.2) and Lemma 7.1 we thus have that

$$
\begin{align*}
0= & C_{\alpha} \psi+S_{\theta}^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)+\overline{\frac{\theta}{\alpha}(0)}\left(\chi_{\frac{\theta}{\alpha}}(0)-\overline{\psi(0)}\right)\left\|k_{0}^{\theta}\right\|^{-2} \tilde{k}_{0}^{\theta} \\
= & \left(C_{\alpha} \psi+\left(\chi_{\frac{\theta}{\alpha}}(0)+\left|\frac{\theta}{\alpha}(0)\right|^{2}\left(\chi_{\frac{\theta}{\alpha}}(0)-\overline{\psi(0)}\right)\left\|k_{0}^{\theta}\right\|^{-2}\right) \tilde{k}_{0}^{\alpha}\right) \\
& +\alpha\left(S_{\theta}^{*} \chi_{\frac{\theta}{\alpha}}+\overline{\frac{\theta}{\alpha}(0)}\left(\chi_{\frac{\theta}{\alpha}}(0)-\overline{\psi(0)}\right)\left\|k_{0}^{\theta}\right\|^{-2} \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right) . \tag{7.1}
\end{align*}
$$

Since the right-hand side of (7.1) is an orthogonal sum, each term must be zero. Thus, for $\left.s=-\overline{\chi_{\frac{\theta}{\alpha}}(0)}-\left|\frac{\theta}{\alpha}(0)\right|^{2} \overline{\overline{\chi_{\frac{\theta}{\alpha}}(0)}}-\psi(0)\right)\left\|k_{0}^{\theta}\right\|^{-2}$, we get $\psi=s k_{0}^{\alpha}$. Moreover, $S_{\theta}^{*} \chi_{\frac{\theta}{\alpha}}=-\overline{\frac{\theta}{\alpha}(0)}\left(\chi_{\frac{\theta}{\alpha}}(0)-\overline{\psi(0)}\right)\left\|k_{0}^{\theta}\right\|^{-2} \tilde{k}_{0}^{\frac{\theta}{\alpha}}$; hence, proceeding as in (8.7), we get

$$
\chi_{\frac{\theta}{\alpha}}=\chi_{\frac{\theta}{\alpha}}(0) k_{0}^{\frac{\theta}{\alpha}}-\frac{\bar{\theta}}{\alpha}(0)\left(\chi_{\frac{\theta}{\alpha}}(0)-\overline{\psi(0)}\right)\left\|k_{0}^{\theta}\right\|^{-2} S_{\frac{\theta}{\alpha}} \tilde{k}_{0}^{\frac{\theta}{\alpha}}=-\bar{s} k_{0}^{\frac{\theta}{\alpha}}
$$

by Lemma 2.3.
The converse is immediate from (6.2). The proof of (2) is analogous.
Remark 7.3. When the right-hand side of the characterization (5.1) reduces to a rank 1 operator const $\cdot k_{0}^{\alpha} \otimes k_{0}^{\theta}$, this operator can immediately be expressed in terms of the symbol $\psi+\bar{\chi}$ as

$$
\begin{aligned}
\text { const } \cdot k_{0}^{\alpha} \otimes k_{0}^{\theta} & =P_{\mathbb{C} k_{0}^{\alpha}} \psi \otimes k_{0}^{\theta}+k_{0}^{\alpha} \otimes P_{\mathbb{C} k_{0}^{\theta}} \chi \\
& =\left(\psi(0)\left\|\tilde{k}_{0}^{\alpha}\right\|^{-2}+\overline{\chi(0)}\left\|\tilde{k}_{0}^{\theta}\right\|^{-2}\right) k_{0}^{\alpha} \otimes k_{0}^{\theta}
\end{aligned}
$$

It might be of independent interest to consider the case when the right-hand side in the equation (6.1) reduces to a rank 1 operator const $\cdot \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\theta}$. By equations (6.2) and Lemma 7.1(1)-(2),(5), this operator can be expressed in terms of the symbol $\psi+\bar{\chi}$ as

$$
\begin{aligned}
\text { const } \cdot \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\theta} & =P_{\widetilde{C} \tilde{k}_{0}^{\alpha}} \mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes P_{\widetilde{\mathbb{C}} \tilde{k}_{0}^{\theta}} \\
& =\left(\overline{\chi_{\alpha}(0)}\left\|\tilde{k}_{0}^{\alpha}\right\|^{-2}+\frac{\theta}{\alpha}(0)\left(\psi(0)-\overline{\chi_{\frac{\theta}{\alpha}}(0)}\left\|\tilde{k}_{0}^{\theta}\right\|^{-2}\right)\right) \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\theta} .
\end{aligned}
$$

A similar question can be asked regarding the case when the right-hand side of the equation (6.1) reduces to a rank 1 operator const $\cdot \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\alpha}$. By equations (6.2),

Proposition 2.2(4), and Lemma 7.1(3),(5)-(6), we have

$$
\text { const } \begin{aligned}
\tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\alpha} & =P_{\mathbb{C} \tilde{k}_{0}^{\alpha}} \mu \otimes \tilde{k}_{0}^{\alpha}+\tilde{k}_{0}^{\alpha} \otimes P_{\mathbb{C} \tilde{k}_{0}^{\alpha}} \nu \\
& =\left(\overline{\chi(0)}-\overline{\chi_{\frac{\theta}{\alpha}}(0)}|\alpha(0)|^{2}+\psi(0)\right)\left\|k_{0}^{\alpha}\right\|^{-2} \tilde{k}_{0}^{\alpha} \otimes \tilde{k}_{0}^{\alpha}
\end{aligned}
$$

## 8. An inverse problem: From the operator to the symbol

In the case of a classical Toeplitz operator $T_{\varphi}$ on $H^{2}$, the (unique) symbol $\varphi$ can be obtained from the operator by the formula $\lim _{n \rightarrow \infty} \bar{z}^{n} T_{\varphi} z^{n}$. In the case of a truncated Toeplitz operator, that is, of the form $A_{\varphi}^{\alpha, \theta}$ with $\alpha=\theta$, one can obtain a symbol belonging to $H^{2}+\overline{H^{2}}$ from the action of $A_{\varphi}^{\theta}$ on $k_{0}^{\theta}$ and $\tilde{k}_{0}^{\theta}$ (see [4]). A similar result can be obtained for an asymmetric truncated Toeplitz operator $A \in \mathcal{T}(\theta, \alpha)$ by considering the action of the operator $A$ and its adjoint on reproducing kernel functions of the same kind.

Note first that, if $A$ has a symbol $\psi+\bar{\chi}$ with $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$, writing $\chi=$ $\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$, where $\chi_{\frac{\theta}{\alpha}} \in K_{\frac{\theta}{\alpha}}^{2}$, and $\chi_{\alpha} \in K_{\alpha}^{2}$, then by Proposition 3.6 we have the following equations on $L^{2}$ :

$$
\left\{\begin{array}{l}
\psi-\overline{\theta(0)} \alpha \bar{\chi}_{\alpha}=A k_{0}^{\theta}-\overline{\chi(0)} k_{0}^{\alpha}+\overline{\theta(0) \chi_{\alpha}(0)} \alpha  \tag{8.1}\\
-\theta(0) \bar{z} \psi+\bar{z} \alpha \bar{\chi}_{\alpha}=A \tilde{k}_{0}^{\theta}-\bar{z} \frac{\theta}{\alpha}(0) \psi(0) \alpha .
\end{array}\right.
$$

Taking the scalar product with $k_{0}^{\alpha}$ and $\tilde{k}_{0}^{\alpha}$, respectively, we obtain

$$
\left\{\begin{array}{l}
\psi(0)+\overline{\theta(0)} \alpha(0) \overline{\chi_{\alpha}(0)}+\left\|k_{0}^{\alpha}\right\|^{2} \overline{\chi(0)}=\left\langle A k_{0}^{\theta}, k_{0}^{\alpha}\right\rangle  \tag{8.2}\\
\frac{\theta}{\alpha}(0) \psi(0)+\overline{\chi_{\alpha}(0)}=\left\langle A \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle,
\end{array}\right.
$$

and, by using the notation $a=\psi(0), b=\overline{\chi_{\alpha}(0)}, c=\overline{\chi(0)}$ in (8.2), we obtain the following system of equations for the unknowns $a, b, c$ :

$$
\left\{\begin{array}{l}
a+\overline{\theta(0)} \alpha(0) b+\left\|k_{0}^{\alpha}\right\|^{2} c=\left\langle A k_{0}^{\theta}, k_{0}^{\alpha}\right\rangle  \tag{8.3}\\
-\frac{\theta}{\alpha}(0) a+b=\left\langle A \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle .
\end{array}\right.
$$

Hence we have proven the following.
Lemma 8.1. Let $A \in \mathcal{T}(\theta, \alpha), A=A_{\psi+\bar{\chi}}$, where $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$, and $\chi=$ $\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$, with $\chi_{\frac{\theta}{\alpha}} \in K_{\frac{\theta}{\alpha}}^{2}$ and $\chi_{\alpha} \in K_{\alpha}^{2}$. Then the values $\psi(0), \overline{\chi_{\alpha}(0)}, \overline{\chi_{\frac{\theta}{\alpha}}(0)}$, and $\overline{\chi(0)}=\overline{\chi_{\frac{\theta}{\alpha}}(0)}+\overline{\frac{\theta}{\alpha}(0)} \overline{\chi_{\alpha}(0)}$ can be determined from $\left\langle A k_{0}^{\theta}, k_{0}^{\alpha}\right\rangle,\left\langle A \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\alpha}\right\rangle$ as a solution to the system (8.3). The solution is unique if we fix a concrete value of $\psi(0)$ or $\chi(0)$.

Note that by Corollary 4.5 we can always choose a concrete value of either $\psi(0)$ or $\chi(0)$.

From equations (8.1) we cannot obtain $\chi_{\frac{\theta}{9}}$; hence we need another equality, which we formulate by using Proposition $3.6(5-6)$ and applying $C_{\theta}$. We get

$$
C_{\theta}\left(A_{\bar{\psi}}^{\alpha, \theta} k_{0}^{\alpha}\right)=\psi(0) \tilde{k}_{0}^{\theta}-\alpha(0)\left(\frac{\theta}{\alpha} \bar{z} \psi-\frac{\theta}{\alpha}(0) \psi(0) \bar{z}\right),
$$

and by Proposition 2.3

$$
C_{\theta}\left(A_{\chi}^{\alpha, \theta} k_{0}^{\alpha}\right)=\alpha \bar{z} \bar{\chi}_{\alpha}+\theta \bar{z} \bar{\chi}_{\frac{\theta}{\alpha}}-\alpha(0) \frac{\theta}{\alpha} \bar{z} \bar{\chi}_{\frac{\theta}{\alpha}} .
$$

The operator $A$ must satisfy the condition

$$
\begin{equation*}
-\alpha(0) \frac{\theta}{\alpha} \bar{z} \psi+\alpha \bar{z} \bar{\chi}_{\alpha}+(\alpha-\alpha(0)) \bar{z} \frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}=C_{\theta} A^{*} k_{0}^{\alpha}-\psi(0) \bar{z} \theta . \tag{8.4}
\end{equation*}
$$

To obtain a symbol for the operator $A$, we take the analytic functions $\psi=X$, $\alpha \bar{\chi}_{\alpha}=Y$, and $\frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}=Z$ as unknowns, and we substitute $a=\psi(0), b=\overline{\chi_{\alpha}(0)}$, $c=\overline{\chi(0)}$. Then from (8.1) and (8.4) we get the system of equations

$$
\left\{\begin{array}{l}
X-\overline{\theta(0)} Y=A k_{0}^{\theta}-c k_{0}^{\alpha}+b \overline{\theta(0)} \alpha  \tag{8.5}\\
-\theta(0) X+Y=z A \tilde{k}_{0}^{\theta}-\frac{\theta}{\alpha}(0) a \alpha \\
-\alpha(0) \frac{\theta}{\alpha} X+Y+(\alpha-\alpha(0)) Z=z C_{\theta} A^{*} k_{0}^{\alpha}-a \theta
\end{array}\right.
$$

Note that the determinant of the matrix of the coefficients of the system (8.5) is $\left\|k_{0}^{\theta}\right\|^{2}(\alpha-\alpha(0)) \neq 0$; hence the solution of the system (8.5) is unique.

We have thus the following.
Theorem 8.2. Let $\alpha \leq \theta$ be nonconstant inner functions, and let $A \in \mathcal{T}(\theta, \alpha)$. Suppose that a certain value is given for either $\psi(0)$ or $\chi(0)$, and let $a, b, c$ be the corresponding solutions to (8.3). If $X, Y, Z$ satisfy the linear system (8.5), then $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$, where $\psi=X$, and $\chi=\frac{\theta}{\alpha} \bar{Z}+\theta \bar{Y}$.
Proof. If $A$ has the symbol $\psi+\bar{\chi}$ with fixed value $\psi(0)$ or $\chi(0)$, then $\psi$ and $\chi$ are uniquely determined. Since, as shown above, $X=\psi, Y=\alpha \bar{\chi}_{\alpha}, Z=\frac{\theta}{\alpha} \bar{\chi}_{\frac{\theta}{\alpha}}$ satisfy the system (8.5), the result follows from the uniqueness of the solution to that system.

The characterizations of asymmetric truncated Toeplitz operators in terms of operators of rank 2 at most, as obtained in previous sections, allow us also to obtain a symbol for the operator.

In fact, regarding the first characterization, it follows from the proof of Theorem 5.1 that, if $A$ is a bounded operator and satisfies the equality (5.1), then $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$. Note that by Corollary 4.5 we know that $\psi$ and $\chi$ are not unique, and we can adjust the value of either $\psi$ or $\chi$ at the origin.

For $\alpha=\theta$ the characterization (6.1) of truncated Toeplitz operators in Theorem 6.1 reduces to that of Sarason (see [19, Remark, p. 501]). In that case the relation between $\psi, \chi$ in the symbol of $A_{\psi+\bar{\chi}}^{\theta}$ and $\mu, \nu$ is given by the conjugation $C_{\theta}$, namely, $\mu=C_{\theta} \chi$ and $\nu=C_{\theta} \psi$. Thus one can also immediately associate a symbol of the form $\psi+\bar{\chi}$ to a truncated Toeplitz operator satisfying that equality. In the asymmetric case, however, Theorem 6.1 unveils a more complex connection between the rank 2 operator on the right-hand side of (6.1) and the symbols of $A_{\psi+\bar{\chi}}^{\theta, \alpha}$. Finding a symbol in terms of $\mu$ and $\nu$ for an operator $A$ satisfying equality (6.1) is more difficult.

To solve that problem in the case of asymmetric truncated Toeplitz operators, we start with two auxiliary results.

Lemma 8.3. Let $\psi \in K_{\alpha}^{2}, \chi \in K_{\theta}^{2}$. Assume that $\chi=\chi_{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha} \chi_{\alpha}$ according to the decomposition $K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \frac{\theta}{\alpha} K_{\alpha}^{2}$. If

$$
\mu=C_{\alpha} P_{\alpha}\left(\frac{\bar{\theta}}{\bar{\alpha}} \chi\right)+\bar{b} \tilde{k}_{0}^{\alpha}, \quad \nu=C_{\alpha} \psi+S^{*}\left(\alpha P_{\frac{\theta}{\alpha}} \chi\right)-b \tilde{k}_{0}^{\theta}
$$

for fixed $b \in \mathbb{C}$, then

$$
\begin{aligned}
\psi & =C_{\alpha} \nu_{\alpha}-\left(\overline{\chi_{\frac{\theta}{\alpha}}(0)}-\bar{b} \frac{\theta}{\alpha}(0)\right) k_{0}^{\alpha} \\
\chi_{\alpha} & =C_{\alpha} \mu-b k_{0}^{\alpha}, \quad \chi_{\frac{\theta}{\alpha}}=S_{\frac{\theta}{\alpha}} \nu_{\frac{\theta}{\alpha}}+\left(\chi_{\frac{\theta}{\alpha}}(0)-b \frac{\theta}{\alpha}(0)\right) k_{0}^{\frac{\theta}{\alpha}}
\end{aligned}
$$

where $\nu=\nu_{\alpha}+\alpha \nu_{\frac{\theta}{\alpha}}$ according to the decomposition $K_{\theta}^{2}=K_{\alpha}^{2} \oplus \alpha K_{\frac{\theta}{\alpha}}^{2}$.
Proof. Let us calculate using Proposition 2.1:

$$
\begin{aligned}
P_{\alpha}\left(S^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)\right) & =P_{\alpha}\left(\bar{z}\left(\alpha \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right)\right)=P_{\alpha}\left(\bar{z} \alpha \chi_{\frac{\theta}{\alpha}}\right) \\
& =\alpha P^{-} \bar{\alpha}\left(\bar{z} \alpha \chi_{\frac{\theta}{\alpha}}\right)-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0) \bar{z} \\
& =\alpha P^{-}\left(\bar{z} \chi_{\frac{\theta}{\alpha}}\right)-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0) \bar{z}=\left(\alpha \chi_{\frac{\theta}{\alpha}}(0)-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right) \bar{z} \\
& =\chi_{\frac{\theta}{\alpha}}(0) \tilde{k}_{0}^{\alpha} .
\end{aligned}
$$

Hence, since $C_{\alpha} \psi \in K_{\alpha}^{2}$ and $\tilde{k}_{0}^{\theta}=\frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}+\alpha \tilde{k}_{0}^{\frac{\theta}{\alpha}}$ (Proposition 2.2), we have

$$
P_{\alpha} \nu=C_{\alpha} \psi+\left(\chi_{\frac{\theta}{\alpha}}(0)-b \frac{\theta}{\alpha}(0)\right) \tilde{k}_{0}^{\alpha}
$$

thus

$$
\psi=C_{\alpha} P_{\alpha} \nu-C_{\alpha}\left(\left(\chi_{\frac{\theta}{\alpha}}(0)-b \frac{\theta}{\alpha}(0)\right) \tilde{k}_{0}^{\alpha}\right)=C_{\alpha} \nu_{\alpha}-\left(\overline{\chi_{\frac{\theta}{\alpha}}(0)}-\bar{b} \frac{\theta}{\alpha}(0)\right) k_{0}^{\alpha}
$$

Now let us consider $\chi$. Since $\mu=P_{\alpha} C_{\alpha}\left(\chi_{\overline{\bar{\alpha}}}^{\overline{\bar{\theta}}}\right)+\bar{b} \tilde{k}_{0}^{\alpha}=C_{\alpha} \chi_{\alpha}+\bar{b} \tilde{k}_{0}^{\alpha}$, we have

$$
\chi_{\alpha}=C_{\alpha}\left(\mu-\bar{b} \tilde{k}_{0}^{\alpha}\right)=C_{\alpha} \mu-b k_{0}^{\alpha} .
$$

On the other hand, by Proposition 2.1,

$$
\begin{align*}
\nu_{\frac{\theta}{\alpha}} & =P_{\frac{\theta}{\alpha}}(\bar{\alpha} \nu)=P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} S^{*}\left(\alpha \chi_{\frac{\theta}{\alpha}}\right)-b \bar{\alpha} \tilde{k}_{0}^{\theta}\right)=P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \bar{z}\left(\alpha \chi_{\frac{\theta}{\alpha}}-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0)\right)\right)-b \tilde{k}_{0}^{\frac{\theta}{\alpha}} \\
& =P_{\frac{\theta}{\alpha}}\left(\bar{z} \chi_{\frac{\theta}{\alpha}}\right)-\alpha(0) \chi_{\frac{\theta}{\alpha}}(0) P_{\frac{\theta}{\alpha}}(\bar{\alpha} \bar{z})-b \tilde{k}_{0}^{\alpha}=P_{\frac{\theta}{\alpha}}\left(\bar{z} \chi_{\frac{\theta}{\alpha}}\right)-b \tilde{k}_{0}^{\frac{\theta}{\alpha}} \\
& =\frac{\theta}{\alpha} P^{-}\left(\frac{\bar{\theta}}{\bar{\alpha}} \bar{z} \chi_{\frac{\theta}{\alpha}}\right)-P^{-}\left(\bar{z} \chi_{\frac{\theta}{\alpha}}\right)-b \tilde{k}_{0}^{\frac{\theta}{\alpha}}=\bar{z} \chi_{\frac{\theta}{\alpha}}-\bar{z} \chi_{\frac{\theta}{\alpha}}(0)-b \tilde{k}_{0}^{\frac{\theta}{\alpha}} \\
& =S_{\frac{\theta}{\alpha}}^{*} \chi_{\frac{\theta}{\alpha}}-b \tilde{k}_{0}^{\frac{\theta}{\alpha}} \tag{8.6}
\end{align*}
$$

since $\frac{\bar{\theta}}{\bar{\alpha}} \chi_{\frac{\theta}{\alpha}} \perp H^{2}$. Hence, using a proper defect operator and (2.3), we obtain

$$
\begin{align*}
\chi_{\frac{\theta}{\alpha}} & =\left(S_{\frac{\theta}{\alpha}} S_{\frac{\theta}{\alpha}}^{*}+k_{0}^{\frac{\theta}{\alpha}} \otimes k_{0}^{\frac{\theta}{\alpha}}\right) \chi_{\frac{\theta}{\alpha}} \\
& =S_{\frac{\theta}{\alpha}} S_{\frac{\theta}{\alpha}}^{*} \chi_{\frac{\theta}{\alpha}}+\left(k_{0}^{\frac{\theta}{\alpha}} \otimes k_{0}^{\frac{\theta}{\alpha}}\right) \chi_{\frac{\theta}{\alpha}}  \tag{8.7}\\
S_{\frac{\theta}{\alpha}} \nu_{\frac{\theta}{\alpha}}+b S_{\frac{\theta}{\alpha}} \tilde{k}_{0}^{\frac{\theta}{\alpha}}+\chi_{\frac{\theta}{\alpha}}(0) k_{0}^{\frac{\theta}{\alpha}} & =\left(\chi_{\frac{\theta}{\alpha}}(0)-b \frac{\theta}{\alpha}(0)\right) k_{0}^{\frac{\theta}{\alpha}}+S_{\frac{\theta}{\alpha}} \nu_{\frac{\theta}{\alpha}} .
\end{align*}
$$

Lemma 8.4. Let $A \in \mathcal{T}(\theta, \alpha)$ satisfy the equation

$$
A-S_{\alpha}^{*} A S_{\theta}=\mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu
$$

for $\mu \in K_{\alpha}^{2}, \nu \in K_{\theta}^{2}$. Then $\mu$ and $\nu$ can be chosen such that $P_{\frac{\theta}{\alpha}}(\bar{\alpha} \nu)$ is orthogonal to $\tilde{k}_{0}^{\frac{\theta}{\alpha}}$. In this case, $\mu$ and $\nu$ are uniquely determined.

Proof. According to the decomposition $K_{\theta}^{2}=K_{\alpha}^{2} \oplus \alpha K_{\frac{\theta}{\alpha}}^{2}$ we can write $\nu=P_{\alpha} \nu+$ $\alpha P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu=\nu_{\alpha}+\alpha \nu_{\frac{\theta}{\alpha}}$. By Proposition 2.2 recall that $\tilde{k}_{0}^{\theta}=\frac{\theta}{\alpha}(0) \tilde{k}_{0}^{\alpha}+\alpha \tilde{k}_{0}^{\frac{\theta}{\alpha}}$. Let us define

$$
\mu^{\prime}=\mu+\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} \overline{\left\langle\nu_{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle \tilde{k}_{0}^{\alpha}, \quad \nu^{\prime}=\nu-\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2}\left\langle\nu_{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle \tilde{k}_{0}^{\theta} . .{ }^{\theta} .}
$$

(see (6.4)). We have

$$
A-S_{\alpha}^{*} A S_{\theta}=\mu^{\prime} \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu^{\prime}, \quad \mu^{\prime} \in K_{\alpha}^{2}, \quad \nu^{\prime} \in K_{\theta}^{2}
$$

Moreover, $\nu_{\frac{\theta}{\alpha}}^{\prime}$ is orthogonal to $\tilde{k}_{0}^{\frac{\theta}{\alpha}}$ since

$$
\left\langle\nu_{\frac{\theta}{\alpha}}^{\prime}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle=\left\langle\nu_{\frac{\theta}{\alpha}}-\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2}\left\langle\nu_{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle \tilde{k}_{0}^{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle=0
$$

To show uniqueness it is enough to note that $\mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu=0$ if and only if $\mu=0$ and $\nu=0$. Assume then that $\mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu=0$, which is equivalent to $\mu=c_{1} \tilde{k}_{0}^{\alpha}, \nu=c_{2} \tilde{k}_{0}^{\theta}$, and $c_{1}+\bar{c}_{2}=0$ for $c_{1}, c_{2} \in \mathbb{C}$. Then by Proposition 2.2(4),

$$
P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu=c_{2} P_{\frac{\theta}{\alpha}} \bar{\alpha} \tilde{k}_{0}^{\theta}=c_{2} \tilde{k}_{0}^{\frac{\theta}{\alpha}} .
$$

Since $P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu \perp \tilde{k}_{0}^{\frac{\theta}{\alpha}}$, we then obtain that $c_{2}=0$, which implies that $c_{1}=0$; hence $\mu=0$ and $\nu=0$.

When investigating symbols of the asymmetric truncated Toeplitz operator, it is worth keeping in mind Corollary 4.5, which states that it is enough to find one of them.

Theorem 8.5. Let $\alpha, \theta$ be nonconstant inner functions such that $\alpha \leq \theta$, and let $A$ be a bounded operator satisfying

$$
\begin{equation*}
A-S_{\alpha}^{*} A S_{\theta}=\mu \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu \tag{8.8}
\end{equation*}
$$

for $\mu \in K_{\alpha}^{2}, \nu \in K_{\theta}^{2}$. Then $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$, where

$$
\begin{aligned}
\psi & =C_{\alpha} P_{\alpha}\left(\nu-c \tilde{k}_{0}^{\theta}\right)=C_{\alpha} P_{\alpha} \nu-\bar{c} \frac{\theta}{\alpha}(0) k_{0}^{\alpha} \in K_{\alpha}^{2} \quad \text { and } \\
\chi & =S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}} \bar{\alpha}\left(\nu-c \tilde{k}_{0}^{\theta}\right)+\frac{\theta}{\alpha} C_{\alpha}\left(\mu+\bar{c} \tilde{k}_{0}^{\alpha}\right) \\
& =\left(S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu+c \frac{\theta}{\alpha}(0) k_{0}^{\frac{\theta}{\alpha}}\right)+\frac{\theta}{\alpha}\left(C_{\alpha} \mu+c k_{0}^{\alpha}\right) \in K_{\theta}^{2}=K_{\frac{\theta}{\alpha}}^{2} \oplus \frac{\theta}{\alpha} K_{\alpha}^{2} \quad \text { with } \\
c & =\left\langle P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} .
\end{aligned}
$$

Proof. Let $\mu_{0}=\mu+\bar{c} \tilde{k}_{0}^{\alpha}, \nu_{0}=\nu-c \tilde{k}_{0}^{\theta}$ with $c=\left\langle P_{\frac{\theta}{\alpha}} \bar{\alpha} \nu, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2}$. We have $P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \nu_{0}\right) \perp \tilde{k}_{0}^{\frac{\theta}{\alpha}}$, and (8.8) is also satisfied; that is,

$$
A-S_{\alpha}^{*} A S_{\theta}=\mu_{0} \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu_{0}
$$

On the other hand, $A \in \mathcal{T}(\theta, \alpha)$, so $A=A_{\psi+\bar{\chi}}^{\theta, \alpha}$ with $\psi \in K_{\alpha}^{2}$ and $\chi \in K_{\theta}^{2}$ and, by (6.2) and (6.4), there are also $\mu^{\prime}=C_{\alpha} P_{\alpha}(\chi \overline{\bar{\alpha}})+\bar{b} \tilde{b}_{0}^{\alpha} \in K_{\alpha}^{2}$ and $\nu^{\prime}=$ $C_{\alpha} \psi+S_{\theta}^{*}\left(\alpha P_{\frac{\theta}{\alpha}} \chi\right)-b \tilde{k}_{0}^{\theta} \in K_{\theta}^{2}(b \in \mathbb{C})$ such that

$$
A-S_{\alpha}^{*} A S_{\theta}=\mu^{\prime} \otimes \tilde{k}_{0}^{\theta}+\tilde{k}_{0}^{\alpha} \otimes \nu^{\prime}
$$

Take

$$
b=-\chi_{\frac{\theta}{\alpha}}(0) \overline{\frac{\theta}{\alpha}(0)}\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-26},
$$

where $\chi_{\frac{\theta}{\alpha}}=P_{\frac{\theta}{\alpha}} \chi$. Using calculations as in (8.6), we have $P_{\frac{\theta}{\alpha}}\left(\bar{\alpha}\left(\nu^{\prime}-b \tilde{k}_{0}^{\theta}\right)\right) \perp \tilde{k}_{0}^{\frac{\theta}{\alpha}}$ because

$$
\begin{aligned}
\left\langle P_{\frac{\theta}{\alpha}} \bar{\alpha}\left(\nu^{\prime}-b \tilde{k}_{0}^{\theta}\right), \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle & =\left\langle S^{*} \chi_{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle-b\left\langle P_{\frac{\theta}{\alpha}} \bar{\alpha} \tilde{k}_{0}^{\theta}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle \\
& =\left\langle\bar{z}\left(\chi_{\frac{\theta}{\alpha}}-\chi_{\frac{\theta}{\alpha}}(0)\right), \bar{z}\left(\frac{\theta}{\alpha}-\frac{\theta}{\alpha}(0)\right)\right\rangle-b\left\langle\tilde{k}_{0}^{\frac{\theta}{\alpha}}, \tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\rangle \\
& =-\chi_{\frac{\theta}{\alpha}}(0) \frac{\theta}{\alpha}(0)-b\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{2}=0
\end{aligned}
$$

hence, by Lemma 8.4, $\mu_{0}=\mu^{\prime}$ and $\nu_{0}=\nu^{\prime}$. On the other hand, by Lemma 8.3,

$$
\psi=C_{\alpha} P_{\alpha} \nu_{0}-\left(\overline{\chi_{\frac{\theta}{\alpha}}(0)}+\overline{\chi_{\frac{\theta}{\alpha}}(0)}\left|\frac{\theta}{\alpha}(0)\right|^{2}\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2}\right) k_{0}^{\alpha}=C_{\alpha} P_{\alpha} \nu_{0}-\overline{\chi_{\frac{\theta}{\alpha}}(0)}\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} k_{0}^{\alpha}
$$

and

$$
\begin{aligned}
\chi & =S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \nu_{0}\right)+\chi_{\frac{\theta}{\alpha}}(0)\left(1+\frac{\left|\frac{\theta}{\alpha}(0)\right|^{2}}{\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{2}}\right) k_{0}^{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha}\left(C_{\alpha} \mu_{0}+\chi_{\frac{\theta}{\alpha}}(0) \frac{\overline{\frac{\theta}{\alpha}(0)}}{\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{2}} k_{0}^{\alpha}\right) \\
& =S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \nu_{0}\right)+\chi_{\frac{\theta}{\alpha}}(0)\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} k_{0}^{\frac{\theta}{\alpha}}+\frac{\theta}{\alpha}\left(C_{\alpha} \mu_{0}+\chi_{\frac{\theta}{\alpha}}(0) \frac{\theta}{\alpha}(0)\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} k_{0}^{\alpha}\right) \\
& =S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \nu_{0}\right)+\frac{\theta}{\alpha} C_{\alpha} \mu_{0}+\chi_{\frac{\theta}{\alpha}}(0)\left\|\tilde{k}_{0}^{\frac{\theta}{\alpha}}\right\|^{-2} k_{0}^{\theta} .
\end{aligned}
$$

Thus, taking Corollary 4.5 into account, we conclude that $A=A_{\psi^{\prime}+\overline{\chi^{\prime}}}^{\theta, \alpha}$ with $\psi^{\prime}=$ $C_{\alpha} P_{\alpha} \nu_{0}=C_{\alpha} P_{\alpha}\left(\nu-c \tilde{k}_{0}^{\theta}\right)$ and that $\chi^{\prime}=S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}}\left(\bar{\alpha} \nu_{0}\right)+\frac{\theta}{\alpha} C_{\alpha} \mu_{0}=S_{\frac{\theta}{\alpha}} P_{\frac{\theta}{\alpha}}(\bar{\alpha}(\nu-$ $\left.\left.c \tilde{k}_{0}^{\theta}\right)\right)+\frac{\theta}{\alpha} C_{\alpha}\left(\mu+\bar{c} \tilde{k}_{0}^{\alpha}\right)$.

Acknowledgments. Câmara's work was partially supported by Fundação para a Ciência e a Tecnologia (FCT/Portugal), through Project PEst-OE/EEI/ LA0009/2013. The work of Kliś-Garlicka and Ptak was partially supported by the Ministry of Science and Higher Education of the Republic of Poland.

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[^0]:    Copyright 2017 by the Tusi Mathematical Research Group.
    Received Sep. 15, 2016; Accepted Dec. 15, 2016.
    First published online Aug. 30, 2017.
    *Corresponding author.
    2010 Mathematics Subject Classification. Primary 47B35; Secondary 30H10, 47A15.
    Keywords. model space, truncated Toeplitz operator, kernel functions, conjugation.

