# OBLIQUELY REFLECTED BROWNIAN MOTION IN NONSMOOTH PLANAR DOMAINS 

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We construct obliquely reflected Brownian motions in all bounded simply connected planar domains, including nonsmooth domains, with general reflection vector fields on the boundary. Conformal mappings and excursion theory are our main technical tools. A key intermediate step, which may be of independent interest, is an alternative characterization of reflected Brownian motions in smooth bounded planar domains with a given field of angles of oblique reflection on the boundary in terms of a pair of quantities, namely an integrable positive harmonic function, which represents the stationary distribution of the process, and a real number that represents, in a suitable sense, the asymptotic rate of rotation of the process around a reference point in the domain. Furthermore, we also show that any obliquely reflected Brownian motion in a simply connected Jordan domain can be obtained as a suitable limit of obliquely reflected Brownian motions in smooth domains.

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1. Introduction. Obliquely reflected Brownian motion (ORBM) arises naturally in some applied probabilistic models, for example, in queuing theory; see Ramanan (2006), Williams (1998) and the references therein. This part of the theory of ORBMs is mostly concerned with processes confined to the positive orthant of the Euclidean space with constant reflection direction on each face. ORBMs in nonsmooth (fractal) domains serve as a toy model for some biological phenomena [see Hołyst et al. (2000)]. In this paper, we will construct and investigate ORBMs in bounded simply connected planar domains, including nonsmooth domains, with variable and possibly nonsmooth reflection directions. Conformal mappings will be our main technical tool. The construction of ORBM in a general nonsmooth domain is difficult because the process (if it exists) is nonsymmetric and, therefore, the (symmetric) Dirichlet form approach [see Chen (1993), Fukushima (1967) and the references therein], very successful in the case of normally reflected Brownian motion, is not applicable to ORBM with general nonsmooth reflection directions.

A conceptual problem with obliquely reflected Brownian motion is that the oblique reflection represents, in heuristic terms, a slight push away from the boundary accompanied by a proportional push along the boundary. In fractal domains, the concepts of "normal" direction at a boundary point and moving "along" the boundary do not have a meaning according to classical definitions. Hence, describing and classifying ORBMs in such nonsmooth domains requires a new approach. The key to our study is the observation that ORBMs in smooth domains can be fully and uniquely classified using two "parameters"-an integrable positive harmonic function $h$ and a real number $\mu_{0}$. The harmonic function $h$ represents the density of the stationary distribution of the process and the real number $\mu_{0}$ represents, in an appropriate sense, the asymptotic rate of rotation around a reference point in the domain. This alternative characterization of ORBM will allow us to construct and investigate ORBM in nonsmooth planar domains with general reflection on the boundary. More specifically, we will first show in Theorems 3.2 and 3.5 that $h$ and $\mu_{0}$ provide a parametrization of ORBMs in the unit disc alternative to the reflection vector field on the boundary. Then we will show in Theorems 3.173.19 how ORBMs in nonsmooth domains can be constructed and classified.

Yet another "parametrization" of ORBMs in simply connected domains is given by "rotation rates" $\mu(z)$ of the process around points $z$ in the domain. Every function $\mu(z)$ representing rotation rates is harmonic but not every harmonic function $\mu(z)$ represents rotation rates for an ORBM.

We will also discuss some ORBMs with degenerate ("tangential") "reflection" along the boundary. The infinitely strong tangential push generates jumps along the boundary, a feature not normally associated with models labeled "Brownian." We will show that ORBMs with "degenerate" boundary behavior are processes that recently appeared in the probabilistic literature in a different context.

The present paper can be viewed as a first step in a much more ambitious project to define ORBMs in $d$-dimensional nonsmooth domains with $d \geq 2$. In the two-dimensional case, especially in simply connected domains, one can give a
meaning to the "angle of reflection" even in domains with fractal boundary by approximating the boundary with continuous curves, defining the angle of reflection on these curves, then defining the corresponding ORBMs and finally passing to the limit (see Theorem 3.19 below). The same program is questionable in higherdimensional domains. It is not clear how to define the direction of reflection on a fractal boundary or how to define the direction of reflection on a sequence of approximating smooth surfaces in a "consistent" way. We believe that our approach via the stationary density [see Kang and Ramanan (2014) for a characterization of stationary distributions of ORBMs in $d$-dimenesional piecewise smooth domains] and appropriate "rotations about $(d-2)$-dimensional sets" may be the right approach to the high-dimensional version of the problem but we leave it for a future project.

There are two classes of domains to which some of our results should extend in a fairly straightforward way: unbounded simply connected planar domains and finitely connected bounded planar domains. These generalizations are also left for a future article.

Some results for ORBM in multidimensional domains were obtained in Dupuis and Ishii (1993, 2008), Ramanan (2006), Williams (1998) under rather restrictive assumptions about smoothness of the boundary of the domain and/or the direction of reflection. The theory of nonsymmetric Dirichlet forms was used to construct families of ORBMs in Duarte (2012), Kim, Kim and Yun (1998) under fairly strong assumptions. A fairly explicit formula for the stationary distribution for ORBM in a smooth planar domain was derived in Harrison, Landau and Shepp (1985). Some results on convergence of ORBMs have been recently obtained in Sarantsev (2016a, 2016b) but the setting of those papers is considerably different from ours.

The article is organized as follows. Section 2 contains a review of some basic probabilistic and analytic facts used in the article. It also contains a theorem relating reflection vector fields on the boundary of a domain and harmonic functions inside the domain; this theorem is the fundamental analytic ingredient of our arguments. Our main results are stated in Section 3. Their proofs are given in Section 4. Our proofs are based in part on ideas developed in Burdzy and Marshall (1993).

## 2. Preliminaries.

2.1. Reflected Brownian motion. We will identify $\mathbb{C}$ and $\mathbb{R}^{2}$. Let $\mathcal{B}(x, r)=$ $\left\{z \in \mathbb{R}^{2}:|x-z|<r\right\}$ and $D_{*}=\mathcal{B}(0,1)$. Suppose that $D \subset \mathbb{C}$ is a bounded open set with smooth boundary and $\theta: \partial D \rightarrow(-\pi / 2, \pi / 2)$ is a Borel measurable function satisfying $\sup _{x \in \partial D}|\theta(x)|<\pi / 2$. Let $\mathbf{n}(x)$ denote the unit inward normal vector at $x \in \partial D$ and let $\mathbf{t}(x)=e^{-i \pi / 2} \mathbf{n}(x)$ be the unit vector tangent to $\partial D$ at $x$.

Let $\mathbf{v}_{\theta}(x)=\mathbf{n}(x)+\tan \theta(x) \mathbf{t}(x)$, let $B$ be standard two-dimensional Brownian motion and consider the following Skorokhod equation:

$$
\begin{equation*}
X_{t}=x_{0}+B_{t}+\int_{0}^{t} \mathbf{v}_{\theta}\left(X_{s}\right) d L_{s} \quad \text { for } t \geq 0 \tag{2.1}
\end{equation*}
$$

Here, $x_{0} \in \bar{D}$ and $L$ is the local time of $X$ on $\partial D$. In other words, $L$ is a nondecreasing continuous process that does not increase when $X$ is in $D$, that is, $\int_{0}^{\infty} \mathbf{1}_{D}\left(X_{t}\right) d L_{t}=0$, almost surely. If $\theta$ is $C^{2}$ then equation (2.1) has a unique pathwise solution $(X, L)$ such that $X_{t} \in \bar{D}$ for all $t \geq 0$, by Dupuis and Ishii (1993), Corollary 5.2 [see also Dupuis and Ishii (2008)]. The process $X$ is a continuous strong Markov process on $\bar{D}_{*}$, and is called obliquely reflected Brownian motion in $D$ with reflecting vector field $\mathbf{v}_{\theta}$. When $\theta \equiv 0$, that is, when $\mathbf{v}_{\theta}=\mathbf{n}, X$ is called normally reflected Brownian motion in $D$. The goal of this paper is to construct and characterize obliquely reflected Brownian motions when $\theta$ is nonsmooth and can possibly take values in $[-\pi / 2, \pi / 2]$, and when $\partial D$ is also possibly nonsmooth.

Consider the case when $D=D_{*}$ and recall that we are assuming that $\theta$ is measurable and $\|\theta\|_{\infty}<\pi / 2$. Then one can show that (2.1) has a unique pathwise solution using the decomposition of the process in $D_{*}$ into the radial and angular parts, and an argument similar to that in [Lions and Sznitman (1984), Remark 4.2(ii)]. In both cases discussed above, the ORBM $X$ is a strong Markov process. Since $X$ does not visit the origin as it behaves like a Brownian motion inside the disk $D_{*}$, applying Itô's formula to $Y_{t}=f\left(X_{t}\right)$ with $f(x)=|x|$, we obtain

$$
\begin{equation*}
d Y_{t}=d W_{t}+\frac{1}{Y_{t}} d t-d L_{t} \tag{2.2}
\end{equation*}
$$

where $W_{t}=\int_{0}^{t} \frac{X_{s}}{\left|X_{s}\right|} \cdot d B_{s}$ is a one-dimensional Brownian motion. Note that $L_{t}$ increases only when $Y_{t}=1$. Thus, $Y_{t}$ is a 2-dimensional Bessel process in (0,1] reflected at 1. It is known [see Bass and Chen (2005)] that the one-dimensional SDE (2.2) has a unique strong solution and all its weak solutions have the same distribution. It follows that the distribution of $(|X|, L)$ is independent of the reflection angle $\theta$. Theorem 3.5 proved below implies that this property continues to hold for ORBMs in $D_{*}$ with nonsmooth reflection angles $\theta$ including those that could be tangential in some subset of the boundary $\partial D_{*}$.

It is known that [see Theorem 3.1(ii) below] the submartingale problem formulation of ORBM is equivalent to the one given above. Let $\mathcal{C}$ be the family of all real functions $f \in C^{2}(\bar{D})$ such that

$$
\frac{\partial}{\partial \mathbf{n}} f(x)+\tan \theta(x) \frac{\partial}{\partial \mathbf{t}} f(x) \geq 0, \quad x \in \partial D
$$

We will say that $\left\{\mathbb{P}_{z}: z \in \bar{D}\right\}$ is a solution of the submartingale problem defining an ORBM with the angle of reflection $\theta$ if $\mathbb{P}_{z}\left(X_{0}=z\right)=1$ for every $z \in \bar{D}$, and

$$
\begin{equation*}
f\left(X_{t}\right)-\frac{1}{2} \int_{0}^{t} \Delta f\left(X_{s}\right) d s, \quad t \geq 0 \tag{2.3}
\end{equation*}
$$

is a submartingale under $\mathbb{P}_{z}$ for every $z \in \bar{D}$ and $f \in \mathcal{C}$.
2.2. Review of excursion theory. We will use excursion theory of Brownian motion in our characterization of obliquely reflected Brownian motion. This section contains a brief review of the excursion theory needed in this paper. See, for example, Maisonneuve (1975) for the foundations of the theory in the abstract setting and Burdzy (1987) for the special case of excursions of Brownian motion. Although Burdzy (1987) does not discuss reflected Brownian motion, all of the results we will use from that book readily apply in the present context.

Let $\mathbb{P}_{x}$ denote the distribution of the process $X$ with $X_{0}=x$, defined by (2.1) or (2.3), and let $\mathbb{E}_{x}$ be the corresponding expectation. Let $\mathbb{P}_{x}^{D}$ denote the distribution of Brownian motion starting from $x \in D$ and killed upon exiting $D$.

An "exit system" for excursions of an ORBM $X$ from $\partial D$ is a pair $\left(L_{t}^{*}, H^{x}\right)$ consisting of a positive continuous additive functional $L_{t}^{*}$ of $X$ and a family of "excursion laws" $\left\{H^{x}\right\}_{x \in \partial D}$. Let $\boldsymbol{\Delta}$ denote the "cemetery" point outside $\bar{D}$ and let $\mathcal{C}$ be the space of all functions $f:[0, \infty) \rightarrow \bar{D} \cup\{\boldsymbol{\Delta}\}$ that are continuous and take values in $\bar{D}$ on some interval $[0, \zeta)$, and are equal to $\Delta$ on $[\zeta, \infty)$. For $x \in \partial D$, the excursion law $H^{x}$ is a $\sigma$-finite (positive) measure on $\mathcal{C}$, such that the canonical process is strong Markov on $\left(t_{0}, \infty\right)$, for every $t_{0}>0$, with transition probabilities $\mathbb{P}^{D}$. Moreover, $H^{x}$ gives zero mass to paths that do not start from $x$. We will be concerned only with the "standard" excursion laws; see Definition 3.2 of Burdzy (1987). For every $x \in \partial D$, there exists a unique standard excursion law $H^{x}$ in $D$, up to a multiplicative constant.

Excursions of $X$ from $\partial D$ will be denoted e or $\mathrm{e}_{s}$, that is, if $s<u, X_{s}, X_{u} \in \partial D$, and $X_{t} \notin \partial D$ for $t \in(s, u)$ then $\mathrm{e}_{s}=\left\{\mathrm{e}_{s}(t)=X_{t+s}, t \in[0, u-s)\right\}, \zeta\left(\mathrm{e}_{s}\right)=u-s$ and $\mathrm{e}_{s}(t)=\boldsymbol{\Delta}$ for $t \geq \zeta$. By convention, $\mathrm{e}_{t} \equiv \boldsymbol{\Delta}$ if $\inf \left\{s>t: X_{s} \in \partial D\right\}=t$.

Let $\sigma_{t}=\inf \left\{s \geq 0: L_{s}^{*}>t\right\}$ and $\mathcal{E}_{u}=\left\{\mathrm{e}_{s}: s<\sigma_{u}\right\}$. Let $I$ be the set of left endpoints of all connected components of $(0, \infty) \backslash\left\{t \geq 0: X_{t} \in \partial D\right\}$. The following is a special case of the exit system formula of Maisonneuve (1975). For every $x \in \bar{D}$, every bounded predictable process $V_{t}$ and every universally measurable function $f: \mathcal{C} \rightarrow[0, \infty)$ that vanishes on excursions $\mathrm{e}_{t}$ identically equal to $\boldsymbol{\Delta}$, we have

$$
\begin{align*}
\mathbb{E}_{x}\left[\sum_{t \in I} V_{t} \cdot f\left(\mathrm{e}_{t}\right)\right] & =\mathbb{E}_{x}\left[\int_{0}^{\infty} V_{\sigma_{s}} H^{X\left(\sigma_{s}\right)}(f) d s\right] \\
& =\mathbb{E}_{x}\left[\int_{0}^{\infty} V_{t} H^{X_{t}}(f) d L_{t}^{*}\right] \tag{2.4}
\end{align*}
$$

Here and elsewhere, $H^{x}(f)=\int_{\mathcal{C}} f d H^{x}$. Informally speaking, (2.4) says that the right continuous version $\mathcal{E}_{t+}$ of the process of excursions is a Poisson point process on the local time scale with variable intensity $H^{\cdot}(f)$.

The normalization of the exit system is somewhat arbitrary, for example, if $\left(L_{t}^{*}, H^{x}\right)$ is an exit system and $c \in(0, \infty)$ is a constant then $\left(c L_{t}^{*},(1 / c) H^{x}\right)$ is also an exit system. One can even make $c$ dependent on $x \in \partial D$. Theorem 7.2 of Burdzy (1987) shows how to choose a "canonical" exit system; that theorem is stated for the usual planar Brownian motion but it is easy to check that both the
statement and the proof apply to normally reflected Brownian motion (i.e., ORBM with $\theta \equiv 0$ ). According to that result, if $D$ is Lipschitz then we can take $L_{t}^{*}$ to be the continuous additive functional $L^{X}$ whose Revuz measure is a constant multiple of the surface area measure $d x$ on $\partial D$ and $H^{x}$, to be standard excursion laws normalized so that

$$
\begin{equation*}
H^{x}(A)=\lim _{\delta \downarrow 0} \frac{1}{\delta} \mathbb{P}_{x+\delta \mathbf{n}(x)}^{D}(A), \tag{2.5}
\end{equation*}
$$

for any event $A$ in a $\sigma$-field generated by the process on an interval $\left[t_{0}, \infty\right)$, for any $t_{0}>0$. The Revuz measure of $L^{X}$ is the measure $d x /(2|D|)$ on $\partial D$, that is, if the initial distribution of $X$ is the uniform probability measure $\mu$ on $D$, then

$$
\begin{equation*}
\mathbb{E}_{\mu}\left[\int_{0}^{1} \mathbf{1}_{A}\left(X_{s}\right) d L_{s}^{X}\right]=\int_{A} \frac{d x}{2|D|}, \tag{2.6}
\end{equation*}
$$

for any Borel set $A \subset \partial D$. It has been shown in Burdzy, Chen and Jones (2006) that $L_{t}^{*}=L_{t}^{X}$.

Let $K_{x}(\cdot)$ denote the Poisson kernel for $D_{*}$, that is, $K_{x}(\cdot)$ vanishes continuously on $\partial D_{*} \backslash\{x\}$ and is harmonic and strictly positive in $D_{*}$. We normalize $K_{x}$ so that $K_{x}(0)=1$ for all $x$. It is easy to see that the following equality holds up to a multiplicative constant:

$$
\begin{equation*}
\int_{A} K_{x}(y) d y=\lim _{\delta \downarrow 0} \frac{1}{\delta} \mathbb{E}_{x+\delta \mathbf{n}(x)}^{D_{*}}\left[\int_{0}^{\infty} \mathbf{1}_{A}\left(X_{s}\right) d s\right], \quad A \subset D_{*} \tag{2.7}
\end{equation*}
$$

In view of (2.5), this means that $K_{x}(\cdot)$ is (a constant multiple of) the density of the expected occupation measure for the excursion law $H^{x}$, that is,

$$
\begin{equation*}
\int_{A} K_{x}(y) d y=H^{x}\left(\int_{0}^{\infty} \mathbf{1}_{A}\left(X_{s}\right) d s\right), \quad A \subset D_{*} \tag{2.8}
\end{equation*}
$$

We omitted the multiplicative constant in (2.7) and (2.8) because it is equal to 1 ; see the proof of Theorem 3.12(ii).
2.3. Analytic preliminaries. Recall that $\mathcal{B}(x, r)=\left\{z \in \mathbb{R}^{2}:|x-z|<r\right\}$ and $D_{*}:=\mathcal{B}(0,1)$. Let $\theta: \partial D_{*} \rightarrow[-\pi / 2, \pi / 2]$ be a Borel measurable function. Typically, $|d x|$ will refer to the arc length measure on $\partial D_{*}$ and $d z$ will refer to the two-dimensional Lebesgue measure on $D_{*}$. The notation $|A|$ will represent either the arc length measure of $A \subset \partial D_{*}$ or the two-dimensional Lebesgue measure of $A \subset D_{*} ;$ the meaning should be clear from the context. Let $\|\cdot\|_{L^{1}(D)}$ denote the $L^{1}$ norm for real functions on an open bounded set $D$ with respect to two-dimensional Lebesgue measure $d z$ on $D$ and let $L^{1}(D)$ be the family of real functions in $D$ with finite $L^{1}$ norm. We will abbreviate $\|\cdot\|_{L^{1}\left(D_{*}\right)}$ as $\|\cdot\|_{1}$. Similar conventions will apply to $L^{\infty}=L^{\infty}\left(\partial D_{*}\right)$ with respect to the measure $|d x|$ on $\partial D_{*}$. As usual, we identify functions that are equal to each other a.e. $|d x|$ on $\partial D_{*}$.

For a function $f$ and constant $c$, the notation $f \not \equiv c$ will mean that $f$ is not identically equal to $c$. If $f$ is harmonic and nonnegative in $D_{*}$ then

$$
\|f\|_{1}=\int_{0}^{1} \int_{0}^{2 \pi} f\left(r e^{i t}\right) d t r d r=\pi f(0)
$$

If the nontangential limit of $f(z)$ at $x \in \partial D_{*}$ exists, we denote it by NT- $\lim _{z \rightarrow x} f(z)$. If $f \in L^{1}\left(\partial D_{*}\right)$ then the harmonic extension of $f$ to $D_{*}$, given by the Poisson integral, has nontangential limits equal to $f$ a.e. We will follow the usual convention of using the same letter $f$ to denote the harmonic extension. If $f$ is harmonic in $D_{*}$, let $\widetilde{f}$ denote the harmonic conjugate of $f$ that vanishes at 0 .

Define

$$
\begin{aligned}
\mathcal{T}= & \left\{\theta \in L^{\infty}\left(\partial D_{*}\right):\|\theta\|_{\infty} \leq \pi / 2, \theta \not \equiv \pi / 2, \text { and } \theta \not \equiv-\pi / 2\right\}, \\
\mathcal{B}= & \left\{\theta: \theta \text { is harmonic in } D_{*} \text { and }|\theta(z)|<\pi / 2 \text { for all } z \in D_{*}\right\}, \\
\mathcal{H}= & \left\{\left(h, \mu_{0}\right): h \text { is harmonic in } D_{*}, h(z)>0\right. \\
& \text { for all } \left.z \in D_{*},\|h\|_{1}=\pi h(0)=1 \text { and } \mu_{0} \in \mathbb{R}\right\}
\end{aligned}
$$

and
$\mathcal{R}=\left\{\mu: \mu\right.$ is harmonic in $D_{*}$ and its harmonic conjugate $\widetilde{\mu}(z)>-1$ for all $\left.z \in D_{*}\right\}$.
The following theorem relates these spaces. See (2.23), (2.24) and Corollary 2.5 for additional formulae.

THEOREM 2.1. There are one-to-one correspondences

$$
\begin{aligned}
\mathcal{T} \leftrightarrow \mathcal{B}, & & \theta(x) \leftrightarrow \theta(z) ; \\
\mathcal{H} \leftrightarrow \mathcal{R}, & & \left(h(z), \mu_{0}\right) \leftrightarrow \mu(z) ; \\
\mathcal{B} \leftrightarrow \mathcal{H}, & & \theta(z) \leftrightarrow\left(h(z), \mu_{0}\right) ;
\end{aligned}
$$

given by

$$
\begin{align*}
\theta(z) & =\operatorname{Re} \int_{\partial D_{*}} \frac{x+z}{x-z} \theta(x) \frac{|d x|}{2 \pi}  \tag{2.9}\\
\theta(x) & =\operatorname{NT}_{z \rightarrow x} \theta(z) \quad \text { a.e. }|d x|  \tag{2.10}\\
\mu(z) & =\mu_{0}-\pi \widetilde{h}(z)  \tag{2.11}\\
h(z) & =(\widetilde{\mu}(z)+1) / \pi \quad \text { and } \quad \mu_{0}=\mu(0)  \tag{2.12}\\
h(z) & =\frac{e^{\widetilde{\theta}(z)} \cos \theta(z)}{\pi \cos \theta(0)} \quad \text { and } \quad \mu_{0}=\tan \theta(0) \quad \text { and }  \tag{2.13}\\
\theta(z) & =-\arg \left(h(z)+i \widetilde{h}(z)-i \mu_{0} / \pi\right) \tag{2.14}
\end{align*}
$$

## Moreover,

$$
\begin{align*}
\mu(z) & =\pi h(z) \tan \theta(z) \\
& =\frac{1}{2} \lim _{r \uparrow 1} \int_{|x|=r} \operatorname{Re}\left(\frac{x+z}{x-z}\right) h(x) \tan \theta(x)|d x| \tag{2.15}
\end{align*}
$$

and

$$
\begin{equation*}
\theta(z)=-\arg (h(z)-i \mu(z) / \pi) \tag{2.16}
\end{equation*}
$$

Proof. The subject of analytic and harmonic functions on the disk and their boundary values has a long history. An eminently readable reference for background material on this subject is given in the first three introductory chapters of Hoffman (1962).

Nontangential limits give the correspondence between $\mathcal{T}$ and $\mathcal{B}$. If $\theta \in \mathcal{B}$, then $\theta$ has a nontangential limit at almost every $x \in \partial D_{*}$, which we will call $\theta(x)$. The limit function $\theta(x) \in L^{\infty}\left(\partial D_{*}\right)$, and $\|\theta\|_{\infty} \leq \pi / 2$. Moreover, since $\frac{1}{2 \pi} \operatorname{Re} \frac{x+z}{x-z}$ is the Poisson kernel on $\partial D_{*}$ for $z \in D_{*}$, we have that

$$
\begin{equation*}
\theta(z)+i \tilde{\theta}(z)=\int_{\partial D_{*}} \frac{x+z}{x-z} \theta(x) \frac{|d x|}{2 \pi} . \tag{2.17}
\end{equation*}
$$

In fact, if $\theta$ is any function in $L^{\infty}$ bounded by $\pi / 2$ then the right-hand side (2.17) defines an analytic function on $D_{*}$ whose real part is harmonic on $D_{*}$, bounded by $\pi / 2$ and has nontangential limit function $\theta(x)$, a.e. Since $\int_{\partial D_{*}} \theta(x) \frac{|d x|}{2 \pi}=\theta(0)$, we have $\theta(x) \not \equiv \pi / 2$ and $\theta(x) \not \equiv-\pi / 2$ a.e. if and only if $|\theta(0)|<\pi / 2$ and by the maximum principle, this occurs if and only if $|\theta(z)|<\pi / 2$ for all $z \in D_{*}$.

If $\left(h, \mu_{0}\right) \in \mathcal{H}$, then $\mu$ defined by (2.11) is harmonic on $D_{*}$, with $\mu(0)=\mu_{0}$, and $h(z)=(\widetilde{\mu}(z)+1) / \pi$, since $\pi h(0)=1$ and $\widetilde{h}=h(0)-h$. Since $h>0$, we conclude that $\tilde{\mu}>-1$ and $\mu \in \mathcal{R}$. If $\mu \in \mathcal{R}$, and if $h$ is given by (2.12) then it is easy to verify that $\left(h, \mu_{0}\right) \in \mathcal{H}$. This proves the one-to-one correspondence between functions in $\mathcal{H}$ and $\mathcal{R}$.

The proof for the correspondence between $\mathcal{B}$ and $\mathcal{H}$, (2.15)-(2.16), as well as useful formulae for the corresponding harmonic conjugates are presented in the next two lemmas.

LEMMA 2.2. $\quad$ There is a one-to-one correspondence between $\mathcal{B}$ and $\mathcal{H}, \theta \leftrightarrow$ $\left(h, \mu_{0}\right)$, given by

$$
\begin{align*}
& \theta+i \widetilde{\theta}=i \log \left(h+i \widetilde{h}-i \mu_{0} / \pi\right)-i \log \left(\left(\sqrt{1+\mu_{0}^{2}}\right) / \pi\right) \quad \text { and }  \tag{2.18}\\
& h+i \widetilde{h}=\frac{e^{-i(\theta+i \widetilde{\theta})}}{\pi \cos \theta(0)}+i \frac{\tan \theta(0)}{\pi} \quad \text { and } \quad \mu_{0}=\tan \theta(0) . \tag{2.19}
\end{align*}
$$

Proof. If $\left(h, \mu_{0}\right) \in \mathcal{H}$ then the right-hand side of (2.18) defines an analytic function $S\left(h, \mu_{0}\right)(z)$ on $D_{*}$ with

$$
\operatorname{Re} S\left(h, \mu_{0}\right)(z)=-\arg \left(h+i \widetilde{h}-i \mu_{0} / \pi\right) \in(-\pi / 2, \pi / 2)
$$

and $S\left(h, \mu_{0}\right)(0)=-\arg \left(1-i \mu_{0}\right)$, which is purely real. Thus, $S\left(h, \mu_{0}\right)=\theta+i \tilde{\theta}$ for some $\theta \in \mathcal{B}$. Likewise, if $\theta \in \mathcal{B}$ then the right-hand side of the first equation in (2.19) defines an analytic function, $T(\theta)(z)$, on $D_{*}$ with $\operatorname{Re} T(\theta)(z)=$ $e^{\widetilde{\theta}(z)} \cos \theta(z) /(\pi \cos \theta(0))>0$ and $\operatorname{Re} T(\theta)(0)=1 / \pi$. Setting $\mu_{0}=\tan \theta(0)$ we conclude that if $h \equiv \operatorname{Re} T(\theta)$ then $\left(h, \mu_{0}\right) \in \mathcal{H}$. Moreover, it is straightforward to verify that, given $\left(h, \mu_{0}\right) \in \mathcal{H}$, if $\theta$ is defined by (2.18) then

$$
h=\operatorname{Re} T(\theta) \quad \text { and } \quad \mu_{0}=\tan \theta(0)
$$

Alternatively, given $\theta \in \mathcal{B}$, if $\left(h, \mu_{0}\right)$ is defined by (2.19) then

$$
\theta=\operatorname{Re} S\left(h, \mu_{0}\right)
$$

This proves the one-to-one correspondence in Lemma 2.2.
The equality in (2.16) of Theorem 2.1 follows immediately from (2.14) and (2.11). The first equality in (2.15) of Theorem 2.1 follows by taking real and imaginary parts in (2.19), then applying (2.11). The second equality in (2.15) follows from the Poisson integral formula on the circle of radius $r<1$ because $\mu$ is harmonic by (2.11).

This completes the proof of Theorem 2.1.
The next lemma relates $\mu \in \mathcal{R}$ to both $h$ and $\theta$ via a Mobius transformation. It will be used in the proof of Theorem 3.15.

Lemma 2.3. Suppose $\left(h, \mu_{0}\right) \in \mathcal{H}, \theta \in \mathcal{B}$, and $\mu \in \mathcal{R}$ with $\left(h, \mu_{0}\right) \leftrightarrow \theta \leftrightarrow \mu$. If $\phi$ is a one-to-one analytic map of $D_{*}$ onto $D_{*}$ then

$$
\begin{equation*}
\theta \circ \phi \in \mathcal{B} \leftrightarrow\left(\frac{h \circ \phi}{\|h \circ \phi\|_{1}}, \frac{\mu(\phi(0))}{\|h \circ \phi\|_{1}}\right) \in \mathcal{H} . \tag{2.20}
\end{equation*}
$$

Proof. First observe that if $f$ is harmonic then $(f+i \tilde{f}) \circ \phi-i \tilde{f}(\phi(0))$ is analytic with imaginary part vanishing at 0 , so that

$$
\begin{equation*}
\widetilde{f \circ \phi}=\tilde{f} \circ \phi-\tilde{f}(\phi(0)) \tag{2.21}
\end{equation*}
$$

Evaluating the real part of (2.19) at $z=\phi(0)$ we obtain

$$
\begin{equation*}
\|h \circ \phi\|_{1}=\pi h(\phi(0))=\frac{e^{\widetilde{\theta}(\phi(0))} \cos \theta(\phi(0))}{\cos \theta(0)} \tag{2.22}
\end{equation*}
$$

Set $h_{1}=h \circ \phi /\|h \circ \phi\|_{1}=h \circ \phi / \pi h(\phi(0))$. Then composing (2.19) with $\phi$ and using (2.21) and (2.11),

$$
\begin{aligned}
h_{1}+i \widetilde{h}_{1} & =\frac{h \circ \phi+i \widetilde{h} \circ \phi-i \widetilde{h}(\phi(0))}{\|h \circ \phi\|_{1}} \\
& =\frac{\exp (-i(\theta+i \widetilde{\theta}) \circ \phi)}{\|h \circ \phi\|_{1} \pi \cos \theta(0)}+\frac{i}{\pi}\left(\frac{\tan \theta(0)-\pi \widetilde{h}(\phi(0))}{\|h \circ \phi\|_{1}}\right) \\
& =\frac{\exp (-i(\theta \circ \phi+i \widetilde{\theta \circ \phi}))}{\pi \cos \theta(\phi(0))}+\frac{i \mu(\phi(0))}{\pi\|h \circ \phi\|_{1}} .
\end{aligned}
$$

By (2.19) and (2.9), the correspondence between $\left(h, \mu_{0}\right) \in \mathcal{H}, \mu \in \mathcal{R}$, and $\theta \in \mathcal{T}$ can also be written as

$$
\begin{align*}
h(z) & =\operatorname{Re}\left(\frac{\exp \left(-i \int_{\partial D_{*}} \frac{x+z}{x-z} \theta(x) \frac{|d x|}{2 \pi}\right)}{\pi \cos \left(\int_{\partial D_{*}} \theta(x) \frac{|d x|}{2 \pi}\right)}\right) \text { and }  \tag{2.23}\\
\mu_{0} & =\tan \left(\int_{\partial D_{*}} \theta(x) \frac{|d x|}{2 \pi}\right) \\
\mu(z) & =-\pi \operatorname{Im}\left(\frac{\exp \left(-i \int_{\partial D_{*}} \frac{x+z}{x-z} \theta(x) \frac{|d x|}{2 \pi}\right)}{\pi \cos \left(\int_{\partial D_{*}} \theta(x) \frac{|d x|}{2 \pi}\right)}\right) . \tag{2.24}
\end{align*}
$$

We would like to have a similar formula for $\mu$ and $\theta$ in terms of $h$, but the situation is a little more complicated for boundary values of positive harmonic functions. A function $h$ is positive and harmonic on $D_{*}$ if and only if

$$
\begin{equation*}
h(z)=\int_{\partial D_{*}} \operatorname{Re}\left(\frac{x+z}{x-z}\right) \sigma(d x) \tag{2.25}
\end{equation*}
$$

for some positive finite (regular Borel) measure $\sigma$ on $\partial D_{*}$. The measures $h(r x)|d x|$ converge weakly to $\sigma(d x)$ as $r \uparrow 1$. The function $h$ has a nontangential limit at almost every $x \in \partial D_{*}$, which we will call $h(x)$, but $h(z)$ is not necessarily the Poisson integral of $h(x)$. In fact, $h \rightarrow+\infty$ radially $\sigma_{s}$-a.e., where $\sigma_{s}$ is the singular component of the Radon-Nikodym decomposition of $\sigma$ with respect to the length measure $|d x|$ on $\partial D_{*}$. It is true, however, that a harmonic function $f$ has nontangential limits $f(x)$ a.e. and satisfies

$$
\begin{equation*}
f(z)+i \tilde{f}(z)=\int_{\partial D_{*}} \frac{x+z}{x-z} f(x) \frac{|d x|}{2 \pi} \tag{2.26}
\end{equation*}
$$

if and only if

$$
\begin{equation*}
\lim _{r \uparrow 1} \int_{\partial D_{*}}|f(r x)-f(x)||d x|=0 \tag{2.27}
\end{equation*}
$$

Given a function $f$ defined on $\partial D_{*}$ which is integrable $|d x|$, if we define $f(z)$ for $z \in D_{*}$ via (2.26) then $f$ satisfies (2.27). See Hoffman (1962), pages 32-33.

If for some $p>1$,

$$
\begin{equation*}
\sup _{r<1} \int_{\partial D_{*}}|f(r x)|^{p}|d x|<\infty \tag{2.28}
\end{equation*}
$$

or if

$$
\sup _{r<1} \int_{\partial D_{*}}|(f+i \tilde{f})(r x)||d x|<\infty
$$

then (2.27) holds. See Hoffman (1962), pages 33 and 51.
Example 2.4. A good example to keep in mind is

$$
\begin{equation*}
h(z)=\frac{1}{\pi} \operatorname{Re}\left(\frac{1+z}{1-z}\right) . \tag{2.29}
\end{equation*}
$$

Then $h(x)=0$ for $x \in \partial D_{*} \backslash\{1\}$. So $h$ cannot be the Poisson integral of its boundary values. Nevertheless, if $\theta \leftrightarrow(h, 0)$ then since $\theta$ is bounded, it satisfies (2.28), and hence satisfies (2.27). In fact, $\theta(x)=-\pi / 2$ for $x \in \partial D_{*}$ with $\operatorname{Im} x>0$ and $\theta(x)=\pi / 2$ for $x \in \partial D_{*}$ with $\operatorname{Im} x<0$, so that

$$
\theta(z)+i \widetilde{\theta}(z)=i \log \frac{1+z}{1-z}=\int_{\partial D_{*}} \frac{x+z}{x-z} \theta(x)|d x| /(2 \pi)
$$

If $h$ satisfies (2.27), where $\left(h, \mu_{0}\right) \in \mathcal{H} \leftrightarrow \theta \in \mathcal{B}$, then we can recover $\theta$ directly from the boundary values of $h$ and $\mu_{0}$. A similar result holds for $\mu$. The following corollary will be used later to interpret $\mu(z)$ as a "rotation rate" about the point $z \in D_{*}$.

Corollary 2.5. $\quad$ Suppose $\left(h, \mu_{0}\right) \in \mathcal{H} \leftrightarrow \theta(z) \in \mathcal{B} \leftrightarrow \theta(x) \in \mathcal{T} \leftrightarrow \mu \in \mathcal{R}$.
(i) If h satisfies (2.27) then for $z \in D_{*}$

$$
\begin{equation*}
\theta(z)=-\arg \left(\int_{\partial D_{*}} \frac{x+z}{x-z} h(x) \frac{|d x|}{2 \pi}-i \mu_{0} / \pi\right) \tag{2.30}
\end{equation*}
$$

(ii) If $h(z) \tan \theta(z)$ or $\widetilde{h}(z)$ satisfy (2.27), then

$$
\begin{align*}
\mu_{0} & =\mu(0)=\frac{1}{2} \int_{\partial D_{*}} h(x) \tan \theta(x)|d x| \quad \text { and }  \tag{2.31}\\
\mu(z) & =\frac{1}{2} \int_{\partial D_{*}} \operatorname{Re}\left(\frac{x+z}{x-z}\right) h(x) \tan \theta(x)|d x|  \tag{2.32}\\
& =\frac{1}{2} \int_{\partial D_{*}} h\left(\frac{x+z}{1+\bar{z} x}\right) \tan \theta\left(\frac{x+z}{1+\bar{z} x}\right)|d x| . \tag{2.33}
\end{align*}
$$

PROOF. (i) follows from the discussion above and (2.18).
(ii) Note that since $\mu=\mu_{0}-\pi \widetilde{h}(z)=\pi h(z) \tan \theta(z)$, for $z \in D_{*}$, it follows that $h(z) \tan \theta(z)$ satisfies (2.27) if and only if $\widetilde{h}(z)$ satisfies (2.27). Equations (2.31) and (2.32) follow from (2.11), (2.15), and (2.26). Finally, equation (2.33) follows from (2.32) and a change of variables.

REMARK 2.6. (i) The maps $\left(h, \mu_{0}\right) \rightarrow \theta$ and $\theta \rightarrow\left(h, \mu_{0}\right)$ are continuous under the topologies of uniform convergence on compact subsets of $D_{*}$ and ( $D_{*}, \mathbb{R}$ ).
(ii) For functions in $\mathcal{B}$, uniform convergence on compact subsets of $D_{*}$ is equivalent to pointwise bounded convergence in $D_{*}$ and is also equivalent to weak-* convergence (of the corresponding boundary value functions) in $L^{\infty}\left(\partial D_{*}\right)$, as elements of the dual space of $L^{1}\left(\partial D_{*}\right)$. But this convergence is not equivalent to pointwise bounded a.e. convergence on $\partial D_{*}$. For example, if $\theta_{k}(z)=-\arg \left(1+z^{k} / 2\right)$, then $\theta_{k} \leftrightarrow\left(h_{k}, 0\right)$, with $h_{k}=\operatorname{Re}\left(1+z^{k} / 2\right)$. The functions $\theta_{k}$ converge to 0 , uniformly on compact subsets of $D_{*}$, pointwise boundedly on $D_{*}$, and weak-* on $\partial D_{*}$. However, $\theta_{k}$ does not contain a subsequence converging pointwise on any subarc in $\partial D_{*}$.
(iii) The function $\theta$ is a constant function if and only if $h \equiv 1 / \pi$ and $\mu_{0}=\tan \theta$. It is tempting to extend the definition of $\mathcal{T}$ to include $\theta \equiv \pi / 2$ by saying $\theta \equiv \pi / 2$ corresponds to $h \equiv 1 / \pi$ and $\mu_{0}=+\infty$. However, we would lose the continuity of the correspondence. Indeed if $\left(h, \mu_{n}\right),\left(g, \mu_{n}\right) \in \mathcal{H}$ with $\mu_{n} \rightarrow+\infty$ and $g \neq h$, let $\theta_{2 n} \leftrightarrow\left(h, \mu_{2 n}\right)$ and $\theta_{2 n+1} \leftrightarrow\left(g, \mu_{2 n+1}\right)$. Then $\theta_{n}$ converges to $\pi / 2$ uniformly on compact subsets of $D_{*}$, but the corresponding elements of $\mathcal{H}$ do not converge.
(iv) If the pair $\left(h, \mu_{0}\right)$ corresponds to $\theta$, then $\left(h(\bar{z}),-\mu_{0}\right)$ corresponds to $-\theta(\bar{z})$. This follows from Lemma 2.2 since $f$ is analytic if and only if $\overline{f(\bar{z})}$ is analytic. But $\left(h,-\mu_{0}\right)$ does not correspond to $-\theta$, unless $h \equiv 1 / \pi$. Indeed, if ( $h,-\mu_{0}$ ) does correspond to $-\theta$ then

$$
-(\theta+i \widetilde{\theta})=i \log \left(h+i \tilde{h}-i\left(-\mu_{0}\right) / \pi\right)-i \log \sqrt{1+\mu_{0}^{2}} / \pi
$$

Adding this equation to (2.18), we obtain

$$
0=i \log \left((h+i \widetilde{h})^{2}+\mu_{0}^{2} / \pi^{2}\right)-2 i \log \sqrt{1+\mu_{0}^{2}} / \pi
$$

and thus $h+i \tilde{h}$ is constant. Since $\left(h, \mu_{0}\right) \in \mathcal{H}$, we have $h \equiv h(0)=1 / \pi$.
(v) Equation (2.30) fails for the example $\theta \leftrightarrow(h, 0) \in \mathcal{H}$ where $h$ is given by (2.29).

EXAMPLE 2.7. Let $F=\phi+i \widetilde{\phi}=\sqrt{\log \left(1-z^{2}\right)}$. We claim we can choose the branch of the square root so that $F$ is analytic on $D_{*}$, with $\phi$ continuous on $\bar{D}_{*}$ and $\widetilde{\phi}$ not bounded above or below. By Theorem 2.1 and the definition of $\mathcal{R}$, there is no $\left(h, \mu_{0}\right) \in \mathcal{H}$ so that $\phi=\mu$, where $\mu \leftrightarrow\left(h, \mu_{0}\right)$. In fact, there do not exist any $a, b \in \mathbb{R}, b \neq 0$ and $\left(h, \mu_{0}\right) \in \mathcal{H}$ such that $a+b \phi=\mu$. To see the
claim, we set $g(z)=(\log (1-z)) / z$. Then $g$ is analytic on a simply connected neighborhood of $\bar{D}_{*} \backslash\{1\}$ and nonvanishing, and hence has an analytic square root $k$. Then $F(z) \equiv z k\left(z^{2}\right)$ is analytic on a neighborhood of $\bar{D}_{*} \backslash\{ \pm 1\}$ and satisfies $F(z)^{2}=\log \left(1-z^{2}\right)$. Thus, $\phi$ and $\widetilde{\phi}$ are continuous and smooth on $\bar{D}_{*} \backslash\{ \pm 1\}$. Since $\phi^{2}-\widetilde{\phi}^{2}=\log \left|1-z^{2}\right| \rightarrow-\infty$ as $z \rightarrow \pm 1$, we conclude $\widetilde{\phi}^{2} \rightarrow \infty$ as $z \rightarrow \pm 1$. But $2 \phi \tilde{\phi}=\arg \left(1-z^{2}\right)$ is bounded, so we must have $\phi \rightarrow 0$ as $z \rightarrow \pm 1$. Thus, $\phi$ is continuous on $\bar{D}_{*}$, and $\widetilde{\phi}$ is unbounded. Since $F$ is odd, $\widetilde{\phi}$ is neither bounded above nor below.

EXAmple 2.8. Consider the harmonic function $\phi(z)=\operatorname{Re} z$ in $D_{*}$ with boundary values $\phi\left(e^{i t}\right)=\cos t, 0 \leq t<2 \pi$. If $a, b \in \mathbb{R}$, with $b \neq 0$, set $\mu=$ $a+b \phi=a+b \operatorname{Re} z$. Then $\tilde{\mu}=b \operatorname{Im} z>-1$ for all $z \in D_{*}$ if and only if $|b| \leq 1$. By the equivalence of $\mathcal{R}$ and $\mathcal{H}$ given in Theorem 2.1, $\mu=a+b \phi$ corresponds to some $\left(h, \mu_{0}\right) \in \mathcal{H}$ if and only if $|b| \leq 1$.

If $\phi$ is harmonic on $D_{*}$ and if $\tilde{\phi}$ is bounded, then for $a, b \in \mathbb{R}$ with $b \neq 0$, the function $\mu=a+b \phi$ has harmonic conjugate $b \widetilde{\phi}$. So for sufficiently small $b$, we have $\tilde{\mu}>-1$ which implies $\mu \in \mathcal{R}$ and $a+b \phi \leftrightarrow\left(h, \mu_{0}\right) \in \mathcal{H}$ for some $\left(h, \mu_{0}\right)$. Since $\widetilde{\mu}(0)=0$, we have that $\inf \widetilde{\phi}<0<\sup \widetilde{\phi}$ so that for $|b|$ sufficiently large $\mu=a+b \phi$ fails to be in $\mathcal{R}$. So in some sense, membership in $\mathcal{R}$ depends on the "oscillation" of the harmonic function on $D_{*}$, but not its mean. The next proposition gives a more precise version. Its proof is elementary, but it will be useful for understanding our (later) description of rotation rates and stationary distributions for ORBMs.

Proposition 2.9. Suppose $\phi$ is (real-valued and) harmonic in $D_{*}$. Set

$$
K_{-}=\inf _{z \in D_{*}} \widetilde{\phi}(z) \quad \text { and } \quad K_{+}=\sup _{z \in D_{*}} \widetilde{\phi}(z)
$$

If $a, b \in \mathbb{R}$ with $-1 /\left|K_{+}\right| \leq b \leq 1 /\left|K_{-}\right|$, then there is a unique $\left(h, \mu_{0}\right) \in \mathcal{H}$ such that

$$
\begin{equation*}
a+b \phi(z)=\mu(z) \tag{2.34}
\end{equation*}
$$

where $\mu$ and $\left(h, \mu_{0}\right)$ are related as in Theorem 2.1. Conversely, if $b<-1 /\left|K_{+}\right|$ or $b>1 /\left|K_{-}\right|$then there do not exist any $a \in \mathbb{R}$ and $\left(h, \mu_{0}\right) \in \mathcal{H}$ such that (2.34) holds.

In the statement of Proposition 2.9, we allow the possibility that $K_{+}$is infinite, in which case we interpret $1 /\left|K_{+}\right|$as equal to zero. A similar statement holds for $\left|K_{-}\right|$.

Proof of Proposition 2.9. Note that $K_{-} \leq 0 \leq K_{+}$since $\widetilde{\phi}(0)=0$. If $b \in \mathbb{R}$ and if $-1 /\left|K_{+}\right| \leq b \leq 1 /\left|K_{-}\right|$, set $\mu=a+b \phi$. Then $\widetilde{\mu}(z)=b \widetilde{\phi}(z) \geq-1$.

Since $\widetilde{\mu}(0)=0$, the maximum principle implies that $\widetilde{\mu}(z)>-1$ for all $z \in D_{*}$, so that $\mu \in \mathcal{R}$. The corresponding $\left(h, \mu_{0}\right) \in \mathcal{H}$ is given by (2.12) of Theorem 2.1.

Conversely if $\left(h, \mu_{0}\right) \in \mathcal{H}$ corresponds to $\mu=a+b \phi \in \mathcal{R}$ as in Theorem 2.1, then $\widetilde{\mu}(z)=b \widetilde{\phi}(z)>-1$. But this implies $b \geq-1 / \sup \widetilde{\phi}(z)$ and $b \leq 1 /|\inf \widetilde{\phi}(z)|$.

If a real-valued function is slightly better than continuous, then its harmonic conjugate is continuous and hence bounded. For a function $f: \partial D_{*} \rightarrow \mathbb{R}$, we define the modulus of continuity of $f$ by $\omega_{f}(a)=\sup _{|s-t|<a}\left|f\left(e^{i s}\right)-f\left(e^{i t}\right)\right|$. We say that $f$ is Dini continuous if $\int_{0}^{b}\left(\omega_{f}(a) / a\right) d a<\infty$ for some $b>0$. If $f$ is Dini continuous then $\tilde{f}$ is continuous and, therefore, bounded; see Garnett (2007), Theorem III.1.3.

Theorem 2.10. Suppose that $\theta \in \mathcal{T},\left(h, \mu_{0}\right) \in \mathcal{H}$, and $\mu \in \mathcal{R}$ correspond to each other as in Theorem 2.1. See also (2.23) and (2.24).
(i) If $\theta$ is Dini continuous on $\partial D_{*}$, then $h$ and $\mu$ extend to be continuous on $\bar{D}_{*}$. If $\mu$ is Dini continuous on $\partial D_{*}$, then $h$ is continuous on $\bar{D}_{*}$ and $\theta$ is continuous on $\bar{D}_{*} \backslash Z$, where $Z=\left\{x \in \partial D_{*}: h(x)=\mu(x)=0\right\}$. Similarly, if $h$ is Dini continuous on $\partial D_{*}$, then $\mu$ is continuous on $\bar{D}_{*}$, and $\theta$ is continuous on $\bar{D}_{*} \backslash Z$. In each of these cases, $h$ and $\widetilde{h}$ satisfy (2.27), so that the conclusions of Corollary 2.5 hold.
(ii) Suppose that $\omega$ is an increasing continuous concave function on $[0, \pi / 2]$ such that $\omega(0)=0, \omega(\pi / 2)=\pi / 4$, and $\int_{0}^{\pi / 2} \frac{\omega(a)}{a} d a=\infty$. Then there exists $\theta \in \mathcal{T}$ such that its modulus of continuity $\omega_{\theta}(a)=\omega(a)$ for $a \in[0, \pi / 2]$ and both $h$ and $\mu$ are unbounded.

Proof. (i) By Garnett (2007), Theorem III.1.3, if $\theta$ is Dini continuous then the harmonic conjugate $\widetilde{\theta}$ is continuous on $\bar{D}_{*}$. Hence, $F(z)=\exp (\widetilde{\theta}(z)-i \theta(z))$ is continuous and so is $h+i \widetilde{h}$ by (2.19). Hence, $h$ and $\mu=\mu_{0}-\pi \widetilde{h}$ are continuous. The remaining statements in (i) follow from (2.11), (2.12) and (2.18) and Garnett (2007), Corollary III.1.4. In each of the cases in (i), $h$ and $\widetilde{h}$ are continuous on $\bar{D}_{*}$, and hence satisfy (2.27).
(ii) We give here an example based on Garnett (2007), page 101. Suppose that $\omega$ is increasing and concave on $[0, \pi / 2]$ with $\omega(0)=0, \omega(\pi / 2)=\pi / 4$, and

$$
\begin{equation*}
\int_{0}^{\pi / 2}(\omega(t) / t) d t=\infty \tag{2.35}
\end{equation*}
$$

Set

$$
\alpha(t)= \begin{cases}\omega(t), & \text { if } 0 \leq t \leq \pi / 2 \\ \omega(\pi-t), & \text { if } \pi / 2 \leq t \leq \pi \\ 0, & \text { if }-\pi<t<0\end{cases}
$$

For $0 \leq x<y \leq \pi$, write $x=t y, 0<t<1$, and so $y-x=(1-t) y$. Since $\omega$ is concave and $\alpha(0)=\omega(0)=0$,

$$
t \alpha(y) \leq \alpha(x) \quad \text { and } \quad(1-t) \alpha(y) \leq \alpha(y-x)
$$

Adding these inequalities, we obtain $\alpha(y)-\alpha(x) \leq \alpha(y-x)$. Since $\alpha(\pi)=0$, replacing $\alpha(t)$ by $\alpha(\pi-t)$ in the above argument, we also have that $\alpha(x)-\alpha(y) \leq$ $\alpha(y-x)$. If $x<0<y<\pi$ with $|x-y|<\pi / 2$, then

$$
\alpha(y)-\alpha(x)=\alpha(y) \leq \alpha(y+|x|)=\alpha(y-x)
$$

Set $\theta\left(e^{i t}\right)=-\alpha(t)$. Then $\theta \in \mathcal{T}$, because $|\alpha| \leq \pi / 4$, and $\omega_{\theta}(a)=\omega_{\alpha}(a)=\omega(a)$ for $0 \leq a \leq \pi / 2$.

Let $b(r)=\cos ^{-1}\left(\frac{1+r}{2}\right)$. Then for $r \in(0,1)$,

$$
\begin{aligned}
\tilde{\theta}(r) & =-\frac{1}{2 \pi} \int_{0}^{\pi} \operatorname{Im}\left(\frac{e^{i t}+r}{e^{i t}-r}\right) \alpha(t) d t \\
& \geq \frac{1}{2 \pi} \int_{b(r)}^{\pi} \frac{2 r \sin t}{\left|e^{i t}-r\right|^{2}} \alpha(t) d t
\end{aligned}
$$

Since $\left|e^{i t}-1\right| \geq\left|e^{i t}-r\right|$ when $\cos t \leq(1+r) / 2$, we have that

$$
\tilde{\theta}(r) \geq-\frac{r}{2 \pi} \int_{b(r)}^{\pi} \operatorname{Im}\left(\frac{e^{i t}+1}{e^{i t}-1}\right) \alpha(t) d t=\frac{r}{2 \pi} \int_{b(r)}^{\pi} \frac{\alpha(t)}{\tan t / 2} d t
$$

which increases to $+\infty$ as $r \rightarrow 1$. So $\widetilde{\theta}(r)$ is not bounded above. Because $\theta$ is continuous on $\partial D_{*}$ with $\theta(1)=0, \theta(z)$ extends to be continuous on $\bar{D}_{*}$ and $\cos \theta(r) \rightarrow 1$ as $r \rightarrow 1$, so by (2.13) $h$ is also unbounded.

Theorem 2.10(ii) implies that if $\theta \in \mathcal{T}$ is not Dini continuous on $\partial D_{*}$, then $h$ and $\mu$ may not be extended continuously to $\bar{D}_{*}$. The next proposition examines the situation when $\theta$ is as large as possible on an interval of $\partial D_{*}$.

Proposition 2.11. Suppose $I$ is an open arc in $\partial D_{*}$, and suppose $\theta \in \mathcal{T} \leftrightarrow$ $\left(h, \mu_{0}\right) \in \mathcal{H}$.
(i) If $\theta(x)=\pi / 2$ a.e. on $I$, then $f=h+i \widetilde{h}-i \mu_{0} / \pi$ extends to be analytic in a neighborhood of $D_{*} \cup I$ with $h=0$ on $I$. The same conclusion holds if $\theta(x)=$ $-\pi / 2$ a.e. on $I$.
(ii) If $h$ extends to be continuous on $D_{*} \cup I$ with $h=0$ on $I$, then $f=h+i \widetilde{h}-$ $i \mu_{0} / \pi$ extends to be analytic in a neighborhood of $D_{*} \cup I$ with at most one zero $e^{i t_{0}} \in I$. If $f \neq 0$ on $I$, then $\theta \equiv \pi / 2$ or $\theta \equiv-\pi / 2$ on I. If $f\left(e^{i t_{0}}\right)=0$ for some $e^{i t_{0}} \in I$, then $\theta\left(e^{i t}\right)=-\pi / 2$ for $e^{i t} \in I$ with $t<t_{0}$ and $\theta\left(e^{i t}\right)=\pi / 2$ for $e^{i t} \in I$ with $t>t_{0}$.

Proof. (i) Suppose $\theta(x)=\pi / 2$ a.e. on $I$. For $z \in D_{*}$ set $F(z)=\theta(z)-$ $\pi / 2+i \widetilde{\theta}(z)$. Then by (2.17)

$$
\begin{align*}
F(z) & =\int_{\partial D_{*}} \frac{x+z}{x-z}(\theta(x)-\pi / 2) \frac{|d x|}{2 \pi} \\
& =\int_{\partial D_{*} \backslash I} \frac{x+z}{x-z}(\theta(x)-\pi / 2) \frac{|d x|}{2 \pi} . \tag{2.36}
\end{align*}
$$

The right-hand side of (2.36) defines an analytic function on $\mathbb{C} \backslash\left(\partial D_{*} \backslash I\right)$. By (2.19), $f \equiv h+i \widetilde{h}-i \mu_{0} / \pi$ extends to be analytic in a neighborhood of $D_{*} \cup I$. Also by (2.36)

$$
\operatorname{Re} F(z)=\theta-\pi / 2=\int_{\partial D_{*} \backslash I} \frac{1-|z|^{2}}{|x-z|^{2}}(\theta(x)-\pi / 2) \frac{|d x|}{2 \pi} .
$$

If $y \in I$, then $\frac{1-|z|^{2}}{|x-z|^{2}} \rightarrow 0$ uniformly in $x \in \partial D_{*} \backslash I$ as $z \rightarrow y$. Thus, $\operatorname{Re} F(z)=$ $\theta(z)-\pi / 2 \rightarrow 0$ as $z \rightarrow y \in I$. Taking real part of (2.19),

$$
h(z)=\frac{e^{\widetilde{\theta}(z)} \cos \theta(z)}{\pi \cos \theta(0)},
$$

so by the continuity of $\theta$ and $\tilde{\theta}$ on $D_{*} \cup I$, we have $h \rightarrow 0$ as $z \rightarrow y \in I$.
To prove (ii), suppose that $h$ extends to be continuous on $D_{*} \cup I$ with $h=0$ on $I$. By the Schwarz reflection principle $f=h+i \widetilde{h}-i \mu_{0} / \pi$ extends analytically across $I$. By the Cauchy-Riemann equations,

$$
\frac{\partial}{\partial t} \operatorname{Im} f\left(e^{i t}\right)=\left.\frac{\partial}{\partial r} \operatorname{Re} f\left(r e^{i t}\right)\right|_{r=1}=\frac{\partial h}{\partial r} \leq 0
$$

on $I$ since $h=0$ on $I$ and $h>0$ on $D_{*}$. Since $\operatorname{Re} f=0$ on $I, \operatorname{Im} f$ cannot be constant on any subarc of $I$, and thus $f$ is a one-to-one map of the arc $I$ onto a subarc of the imaginary axis, and (ii) follows from (2.30).
3. Main results. This section contains only statements of the main results of this paper. The proofs will be given in Section 4. First, in Section 3.1 we establish results when the domain $D$ is smooth and the angle of reflection $\theta$ is $C^{2}$ and nontangential everywhere, that is, $\theta$ lies in a closed subinterval of $(-\pi / 2, \pi / 2)$. Theorem 3.1 summarizes results on existence and uniqueness of ORBMs, and Theorem 3.2 considers ORBMs on the disk $D_{*}$ and establishes the probabilistic interpretation of the quantity $\left(h(z), \mu_{0}\right)$ corresponding to $\theta \in \mathcal{T}$, as specified in Theorem 2.1. ORBMs in $D_{*}$ with general reflection angles $\theta \in \mathcal{T}$ are constructed in Section 3.2. The focus of Section 3.3 (in particular, see Theorem 3.12) is the case when the reflection vector field is tangential at every point, which leads to a process referred to as excursion reflected Brownian motion (ERBM). Lastly, in Section 3.4 (specifically, Theorems 3.15-3.18 therein) we construct ORBMs in simply connected domains using conformal mappings and then show, in the case of simply connected bounded Jordan domains, that they can also be obtained as suitable limits of ORBMs in $C^{2}$ domains.
3.1. Smooth $D$ and $C^{2}$-smooth nontangential $\theta$. We start with a theorem on existence and uniqueness of ORBM in the simplest case, when the domain is smooth and the angle of reflection is smooth and takes values in a closed subinterval of $(-\pi / 2, \pi / 2)$. The result is essentially known.

THEOREM 3.1. Assume that $D \subset \mathbb{C}$ is a bounded open set with $C^{2}$ boundary, and a function $\theta: \partial D \rightarrow(-\pi / 2, \pi / 2)$ is $C^{2}$.
(i) [Harrison, Landau and Shepp (1985), Theorem 2.6] The submartingale problem (2.3) has a unique solution which defines a strong Markov process.
(ii) The strong Markov process defined by the Skorokhod equation (2.1) is continuous and has the same distribution as the process defined by the submartingale problem (2.3).
(iii) [Kim, Kim and Yun (1998)] The ORBM obtained in (i) and (ii) can also be constructed by using the nonsymmetric Dirichlet form approach.

It follows from the results in Harrison, Landau and Shepp (1985) that if $\theta$ is $C^{1}$ then the ORBM $X$ in the unit disc $D_{*}$ has a unique stationary distribution with the density $h$ given by (2.23). The stationary distribution was characterized in Harrison, Landau and Shepp (1985) in terms of a partial differential equation in $D_{*}$ with appropriate boundary conditions. In Theorem 3.2(ii), we will show a partial converse, namely, that the stationary distribution characterizes an ORBM up to a real number that represents the "rotation rate" of $X$ about 0 .

Under the assumptions of Theorem 3.1, the ORBM $X$ is continuous, a.s. Consider a fixed $z \in D_{*}$. Since $X_{t} \neq z$ for all $t>0$, a.s. (even if $X_{0}=z$ ), we can uniquely define the function $t \rightarrow \arg \left(X_{t}-z\right)$ by choosing its continuous version and making an arbitrary convention that $\arg \left(X_{1}-z\right) \in[0,2 \pi)$.

Since $h$ is the density of the stationary measure of $X$ and $\theta$ is the reflection angle, (2.31) suggests that $\mu_{0}$ represents one half of the speed of rotation of $X$ about 0 . Hence, one might hope that $\lim _{t \rightarrow \infty} \arg X_{t} / t$ is equal to a constant multiple of $\mu_{0}$, a.s. Unfortunately, this simple interpretation of $\mu_{0}$ is false because $\arg X_{t}$ behaves like a Cauchy process [see Bertoin and Werner (1994), Spitzer (1958)] and, therefore, the law of large numbers does not hold for $\arg X_{t}$. We will identify $\mu_{0}$ with the speed of rotation using two other representations in Theorem 3.2(ii)-(iii). We need the following definitions to state the representations. First of all, recall that a random variable has the Cauchy distribution if its density is $1 /\left(\pi\left(1+x^{2}\right)\right)$ for $x \in \mathbb{R}$. Next we will define a new measure of winding speed which does not include large windings if they occur during a single excursion from the boundary. Recall definitions related to excursions from Section 2.2. We will say that $\mathrm{e}_{s}$ belongs to the family $\varepsilon_{t}^{L}$ of excursions with "large winding number" if $s+\zeta\left(\mathrm{e}_{s}\right) \leq t$ and $\left|\arg X_{s}-\arg X_{s+\zeta\left(\mathrm{e}_{s}\right)-}\right|>2 \pi$, where $X_{u-}$ denotes
the left-hand limit. For $z \in D_{*}$, let

$$
\begin{align*}
\arg ^{*} X_{t}= & \arg X_{t}-\sum_{s: \mathrm{e}_{s} \in \mathcal{E}_{t}^{L}}\left(\arg X_{s+\zeta\left(\mathrm{e}_{s}\right)-}-\arg X_{s}\right),  \tag{3.1}\\
\arg ^{*}\left(X_{t}-z\right)= & \arg \left(X_{t}-z\right) \\
& -\sum_{s: \mathrm{e}_{s} \in \mathcal{E}_{t}^{L}}\left(\arg \left(X_{s+\zeta\left(\mathrm{e}_{s}\right)-}-z\right)-\arg \left(X_{s}-z\right)\right) .
\end{align*}
$$

THEOREM 3.2. In parts (i)-(iii), we assume that a $C^{2}$ function $\theta: \partial D_{*} \rightarrow$ $(-\pi / 2, \pi / 2)$ is given.
(i) [Harrison, Landau and Shepp (1985), Theorem 2.18] The density of the stationary measure for $X$ defined in (2.1) is a positive harmonic function $h$ in $D_{*}$ given by (2.23) [see also (2.19)].
(ii) With probability $1, X$ is continuous and, therefore, $\arg X_{t}$ is well defined for $t>0$. Let $\mu_{0} \in \mathbb{R}$ be given by (2.23). For every $z \in \bar{D}_{*}$, the distributions of $\frac{1}{t} \arg X_{t}-\mu_{0}$ under $\mathbb{P}_{z}$ converge to the Cauchy distribution when $t \rightarrow \infty$.
(iii) For every $y \in \bar{D}_{*}$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{t}=\mu_{0}, \quad \mathbb{P}_{y} \text {-a.s. } \tag{3.3}
\end{equation*}
$$

The formula holds more generally. For any $y, z \in D_{*}$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*}\left(X_{t}-z\right)=\mu(z), \quad \mathbb{P}_{y} \text {-a.s. } \tag{3.4}
\end{equation*}
$$

where $\mu(z)$ is given by (2.24).
(iv) Conversely, suppose we are given any $\mu_{0} \in \mathbb{R}$ and a harmonic function $h$ in $D_{*}$, that is, $C^{2}$ in $\bar{D}_{*}$, positive on $\bar{D}_{*}$, and satisfies $h(0)=1 / \pi$. Let $\theta \leftrightarrow\left(h, \mu_{0}\right)$. Then for every $x_{0} \in \bar{D}_{*}$, there exists a unique in distribution process $X$ satisfying (2.1) with this $\theta$. Its stationary distribution has density $h$ and (3.3) holds.

REMARK 3.3. (i) We could have defined the family $\mathcal{E}_{t}^{L}$ of excursions $\mathrm{e}_{s}$ with a "large winding number" as those satisfying $s+\zeta\left(\mathrm{e}_{s}\right) \leq t$ and $\mid \arg X_{s}-$ $\arg X_{s+\zeta\left(\mathrm{e}_{s}\right)-} \mid>a$, where $a>0$ is not necessarily $2 \pi$. It turns out that (3.3) holds for any $a>0$. The limit in (3.3) holds for any value of $a$ because the only thing that matters in (3.1) is that the large jumps of the Cauchy-like process $\arg X$ are removed. The "remaining part" of this process satisfies the law of large numbers and has mean $\mu_{0} t$, no matter how large the threshold for the "large jumps" is. We have chosen $a=2 \pi$ because this value has a natural geometric interpretation and is invariant, in a sense, under conformal mappings.
(ii) We will prove (3.4) using (3.23) and a purely analytic argument. Formula (3.4) has the same heuristic meaning as (2.31) as a rotation rate, except that it represents the sum (integral) of infinitesimally small increments of the angle around $z$, not 0 .
(iii) In view of Theorem 2.1, if the rotation rate $\mu(z)$ is known for all $z \in D_{*}$, it completely determines $\theta$ and $h$. Moreover, due to the harmonic character of $\mu(z)$, if this function is known in an arbitrarily small nonempty open subset of $D_{*}$, this also determines $\theta$ and $h$.
(iv) Theorem 2.1 and the definition of the function space $\mathcal{R}$ show which harmonic functions $\mu(z)$ represent rotation rates for an ORBM. See also Proposition 2.9. Roughly speaking, $\mu(z)$ represents rotation rates for an ORBM if its oscillation over $\bar{D}_{*}$ is not too large. There is no restriction, however, on the average value of $\mu(z)$. If $\mu(z)$ and $\mu_{1}(z)$ represent the rotation rates for two ORBM's, and $\mu(z)=c+\mu_{1}(z)$ for some constant $c$ and all $z$ then $\tilde{\mu}=\tilde{\mu}_{1}$. By (2.12) of Theorem 2.1, the corresponding stationary densities are the same for both ORBMs.
(v) Parts (ii) and (iii) of Theorem 3.2 are similar in spirit to Le Gall and Yor (1986), Theorem 7.1, although that paper is concerned with Brownian motion with drift, not reflection.
3.2. ORBMs on $D_{*}$ with general reflection angles $\theta$. Suppose $\theta \in \mathcal{T}$. Then $\theta \not \equiv \pi / 2$ and $\theta \not \equiv-\pi / 2$, although $\theta$ could be tangential on a strict subset of the boundary $\partial D_{*}$. In Theorem 3.5, we show that ORBMs on the disk $D_{*}$ associated with $\theta$ can be obtained as limits of ORBMs on $D_{*}$ with $C^{2}$ angles of reflection, which are well defined by Theorem 3.1. Then in Theorem 3.8 we establish a conformal invariance property for such ORBMs. If there do exist points on the boundary at which $\theta$ is tangential, the associated ORBM will not in general be continuous, and thus one has to carefully define the topology in which the above limit procedure can be carried out.

We start by introducing some relevant notation to define this topology. Let

$$
\begin{equation*}
N_{\theta}^{+}=\left\{x \in \partial D_{*}: \theta(x)=\pi / 2\right\}, \quad N_{\theta}^{-}=\left\{x \in \partial D_{*}: \theta(x)=-\pi / 2\right\} . \tag{3.5}
\end{equation*}
$$

Since we identify functions in $\mathcal{T}$ that are equal to each other a.e.,

$$
\begin{equation*}
\left|N_{\theta}^{+}\right|<2 \pi \quad \text { and } \quad\left|N_{\theta}^{-}\right|<2 \pi \tag{3.6}
\end{equation*}
$$

We will say that $x \in \operatorname{Int} N_{\theta}^{+}$if $\theta \equiv \pi / 2$ a.e. in some neighborhood of $x$. The definition of $\operatorname{Int} N_{\theta}^{-}$is analogous. For $x=e^{i \alpha} \in \operatorname{Int} N_{\theta}^{+}$, let $\alpha^{+}$be the largest real number such that $\left\{e^{i t}: t \in\left[\alpha, \alpha^{+}\right)\right\} \subset \operatorname{Int} N_{\theta}^{+}$, and let $\beta^{+}(x)=e^{i \alpha^{+}}$. Similarly, for $x=e^{i \alpha} \in \operatorname{Int} N_{\theta}^{-}$, let $\alpha^{-}$be the smallest real number such that $\left\{e^{i t}: t \in\left(\alpha^{-}, \alpha\right]\right\} \subset$ Int $N_{\theta}^{-}$, and let $\beta^{-}(x)=e^{i \alpha^{-}}$.

We recall below the definition of the $M_{1}$ topology introduced by Skorokhod in Skorohod (1956). We will use the $M_{1}$ topology rather than the more popular $J_{1}$ topology because we will be concerned with convergence of continuous processes to (possibly) discontinuous processes. In the $J_{1}$ topology, a sequence of continuous processes cannot converge to a discontinuous process. We will also define an $M_{1}^{\mathcal{T}}$ topology, appropriate for our setting.

DEFINITION 3.4. (i) Suppose that $0<T<\infty$ and $x:[0, T] \rightarrow \mathbb{R}^{n}$ is a càdlàg function. The graph $\Gamma_{x}$ is the set consisting of all pairs ( $a, t$ ) such that $0 \leq t \leq T$ and $a \in[x(t-), x(t)]$ (here $[x(t-), x(t)]$ is the line segment between the left-hand limit $x(t-)$ and $x(t)$ in $\left.\mathbb{R}^{n}\right)$. A pair of functions $\{(y(s), t(s)), s \in[0,1]\}$ is a parametric representation of $\Gamma_{x}$ if $y$ is continuous, $t$ is continuous and nondecreasing, and $(v, u) \in \Gamma_{x}$ if and only if $(v, u)=(y(s), t(s))$ for some $s \in[0,1]$. We say that $x_{n}$ converge to $x$ in $M_{1}$ topology if there exist parametric representations $\{(y(s), t(s)), s \in[0,1]\}$ of $\Gamma_{x}$ and $\left\{\left(y_{n}(s), t_{n}(s)\right), s \in[0,1]\right\}$ of $\Gamma_{x_{n}}$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sup _{s \in[0,1]}\left|\left(y_{n}(s), t_{n}(s)\right)-(y(s), t(s))\right|=0 \tag{3.7}
\end{equation*}
$$

(ii) If $x:[0, \infty) \rightarrow \mathbb{R}^{n}$ then we say that $x_{n}(t)$ converge to $x(t)$ in $M_{1}$ topol$o g y$ if they converge to $x$ on $[0, T]$ in $M_{1}$ topology for every $0<T<\infty$. This is equivalent to the following statement. There exist parametric representations $\{(y(s), t(s)), s \in[0, \infty)\}$ of $\Gamma_{x}$ and $\left\{\left(y_{n}(s), t_{n}(s)\right), s \in[0, \infty)\right\}$ of $\Gamma_{x_{n}}$ such that for every $T \in(0, \infty)$,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sup _{s \in[0, T]}\left|\left(y_{n}(s), t_{n}(s)\right)-(y(s), t(s))\right|=0 \tag{3.8}
\end{equation*}
$$

(iii) Consider $\theta \in \mathcal{T}$. We will say that $x:[0, \infty) \rightarrow \bar{D}_{*}$ belongs to $\mathcal{A}_{\theta}$ if it is càdlàg and satisfies the following conditions. For all $t \geq 0, x_{t-} \neq x_{t}$ if and only if $x_{t-} \in \operatorname{Int} N_{\theta}^{+} \cup \operatorname{Int} N_{\theta}^{-}$. Moreover, if $x_{t-} \in \operatorname{Int} N_{\theta}^{+}$then $x_{t}=\beta^{+}\left(x_{t-}\right)$. If $x_{t-} \in$ Int $N_{\theta}^{-}$then $x_{t}=\beta^{-}\left(x_{t-}\right)$. Let $\mathcal{A}_{\mathcal{T}}=\bigcup_{\theta \in \mathcal{T}} \mathcal{A}_{\theta}$.
(iv) Assume that $\theta \in \mathcal{T}$ and $x \in \mathcal{A}_{\theta}$. If $x_{t-}=e^{i \alpha} \in \operatorname{Int} N_{\theta}^{+}$and $x_{t}=\beta^{+}\left(x_{t-}\right)=$ $e^{i \alpha^{+}}$, then we let $\left[x_{t-}, x_{t}\right]_{\theta}=\left\{e^{i t}: t \in\left[\alpha, \alpha^{+}\right]\right\}$be the arc on $\partial D_{*}$ between $x_{t-}$ and $x_{t}$. Thus, $\theta\left(e^{i s}\right)=\pi / 2$ for a.e. $e^{i s} \in\left[x_{t-}, x_{t}\right]_{\theta}$. Similarly, if $x_{t-}=e^{i \alpha} \in \operatorname{Int} N_{\theta}^{-}$ and $x_{t}=\beta^{-}\left(x_{t-}\right)=e^{i \alpha^{-}}$, then we let $\left[x_{t-}, x_{t}\right]_{\theta}=\left\{e^{i t}: t \in\left[\alpha^{-}, \alpha\right]\right\}$.

We define the graph $\Gamma_{x}^{\theta}$ as the set of all pairs ( $a, t$ ) such that $a=x_{t}$ if $x$ is continuous at $t$ and $a \in\left[x_{t-}, x_{t}\right]_{\theta}$ if $x_{t-} \neq x_{t}$. A pair of functions $\{(y(s), t(s)), s \in$ $[0, \infty)\}$ is a parametric representation of $\Gamma_{x}^{\theta}$ if $y$ is continuous, $t$ is continuous and nondecreasing, and $(v, u) \in \Gamma_{x}^{\theta}$ if and only if $(v, u)=(y(s), t(s))$ for some $s \in[0, \infty)$. Suppose that $x_{n} \in \mathcal{A}_{\theta_{n}}$ for some $\theta_{n} \in \mathcal{T}, n \geq 1$, and $x \in \mathcal{A}_{\theta}$ for some $\theta \in \mathcal{T}$. We say that $x_{n}$ converge to $x$ in $M_{1}^{\mathcal{T}}$ topology if there exist parametric representations $\{(y(s), t(s)), s \in[0, \infty)\}$ of $\Gamma_{x}^{\theta}$ and $\left\{\left(y_{n}(s), t_{n}(s)\right), s \in[0, \infty)\right\}$ of $\Gamma_{x_{n}}^{\theta_{n}}$ such that for every $T \in(0, \infty)$,

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sup _{s \in[0, T]}\left|\left(y_{n}(s), t_{n}(s)\right)-(y(s), t(s))\right|=0 \tag{3.9}
\end{equation*}
$$

Some càdlàg functions $x$ (e.g., continuous functions) belong to more than one family $\mathcal{A}_{\theta}$. We leave it to the reader to check that the definitions in (iv) are not affected by the choice of $\mathcal{A}_{\theta}$.

We will extend the definition of $t \rightarrow \arg X_{t}$ to (some) processes that are not continuous. Although it is impossible to define a continuous version of $t \rightarrow \arg X_{t}$
for a process $X$ that is discontinuous, we will define a functional $\left\{X_{t}, t \geq 0\right\} \rightarrow$ $\left\{\boldsymbol{\operatorname { a r g }} X_{t}, t \geq 0\right\}$ in a way that reflects the structure of jumps in a natural way, leading to heuristically appealing results. The functional arg will be defined relative to $\theta$ but the dependence will be suppressed in the notation. Consider a function $x \in \mathcal{A}_{\theta}$ such that $x_{t} \neq 0$ for all $t \geq 0$. Consider any parametric representation $\{(y(s), t(s)), s \in[0, \infty)\}$ of $\Gamma_{x}^{\theta}$ and let $s \rightarrow \arg y(s)$ be the continuous version of $\arg y$ with $\arg y(0) \in[0,2 \pi)$. We let $\arg x_{u}=\arg y(s)$ where $s=\sup \{r: t(r)=u\}$. It is elementary to check that this definition of $\arg x_{u}$ does not depend on the choice of parametric representation $\{(y(s), t(s)), s \in[0, \infty)\}$ of $\Gamma_{x}^{\theta}$.

Recall the definitions (3.1)-(3.2) and notation introduced in the paragraph preceding them. We define $\mathbf{a r g}^{*}$ in an analogous way. For $z \in D_{*}$, let

$$
\begin{aligned}
\arg ^{*} X_{t} & =\arg X_{t}-\sum_{s: \mathrm{e}_{s} \in \mathcal{E}_{t}^{L}}\left(\arg X_{s+\zeta\left(\mathrm{e}_{s}\right)-}-\arg X_{s}\right), \\
\arg ^{*}\left(X_{t}-z\right) & =\arg \left(X_{t}-z\right)-\sum_{s: \mathrm{e}_{s} \in \mathcal{E}_{t}^{L}}\left(\arg \left(X_{s+\zeta\left(\mathrm{e}_{s}\right)-}-z\right)-\arg \left(X_{s}-z\right)\right) .
\end{aligned}
$$

THEOREM 3.5. Consider $\theta \in \mathcal{T}$. There exists a sequence of $C^{2}$ functions $\theta_{k}$ : $\partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ which converges to $\theta$ in weak-* topology as elements of the dual space of $L^{1}\left(\partial D_{*}\right)$, that is,

$$
\lim _{k \rightarrow \infty} \int_{\partial D_{*}} f(x) \theta_{k}(x)|d x|=\int_{\partial D_{*}} f(x) \theta(x)|d x| \quad \text { for every } f \in L^{1}\left(\partial D_{*}\right)
$$

Fix such a sequence $\left\{\theta_{k}\right\}$ and let $X^{k}$ be defined by the following SDE analogous to (2.1),

$$
\begin{equation*}
X_{t}^{k}=z_{k}+B_{t}+\int_{0}^{t} \mathbf{v}_{\theta_{k}}\left(X_{s}^{k}\right) d L_{s}^{k} \quad \text { for } t \geq 0 \tag{3.10}
\end{equation*}
$$

Assume that $z_{k} \rightarrow z_{0} \in D_{*}$ as $k \rightarrow \infty, z_{0} \neq 0$, and recall (3.6).
(i) [Burdzy and Marshall (1993), Theorem 1.1] $X^{k}$,s converge weakly in $M_{1}^{\mathcal{T}}$ topology to a conservative Markov process $X$ on $\bar{D}_{*}$ such that $X_{0}=z_{0}$, a.s. Moreover, there is a càdlàg version of $X$ and for this version, $X \in \mathcal{A}_{\theta}$, a.s. The process $\left\{X_{t} ; t \in\left[0, \sigma_{\partial D_{*}}\right)\right\}$, where $\sigma_{\partial D_{*}}:=\inf \left\{t>0: X_{t} \in \partial D_{*}\right\}$, is Brownian motion killed upon leaving $D_{*}$.
(ii) $X^{k}$ 's converge to $X$ in the sense of finite dimensional distributions.
(iii) The Markov process $X$ has a stationary measure whose density $h$ is given by (2.23).
(iv) The functional $\left\{x_{s}, s \in[0, \infty)\right\} \rightarrow\left\{\arg x_{s}, s \in[0, \infty)\right\}$ is a continuous mapping from the set $\mathcal{A}_{\mathcal{T}}$ equipped with $M_{1}^{\mathcal{T}}$ topology to the set of càdlàg functions equipped with the $M_{1}$ topology. For every $t \geq 0$, the distributions of $\arg X_{t}^{k}$ converge to the distribution of $\arg X_{t}$.
(v) Let $\mu_{0}$ be as in (2.23). Then for every $z \in \bar{D}_{*}$, the distributions of $\frac{1}{t} \arg X_{t}-\mu_{0}$ under $\mathbb{P}_{z}$ converge to the Cauchy distribution when $t \rightarrow \infty$.
(vi) For every $y \in \bar{D}_{*}, \mathbb{P}_{y}$-a.s.,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \mathbf{a r g}^{*} X_{t}=\mu_{0} \tag{3.11}
\end{equation*}
$$

Moreover, for any $y, z \in D_{*}$,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*}\left(X_{t}-z\right)=\mu(z), \quad \mathbb{P}_{y} \text {-a.s. } \tag{3.12}
\end{equation*}
$$

where $\mu(z)$ is the harmonic function defined by (2.24).
(vii) Assume that $\theta \in \mathcal{T} \leftrightarrow\left(h, \mu_{0}\right) \in \mathcal{H}$. Then for every $x \in \partial D_{*}, x \in \Gamma_{X}^{\theta}$ with probability 1 if and only if

$$
\begin{align*}
\int_{0}^{1} & e^{-\widetilde{\theta}(r x)} \cos \theta(r x) \frac{d r}{1-r} \\
& =\int_{0}^{1} \frac{h(r x) /(\pi \cos \theta(0))}{h(r x)^{2}+\left(\widetilde{h}(r x)-\mu_{0} / \pi\right)^{2}} \frac{d r}{1-r}<\infty \tag{3.13}
\end{align*}
$$

(viii) Suppose that $\theta, \bar{\theta}_{k} \in \mathcal{T}$ and $\bar{\theta}_{k}$ converge to $\theta$ in weak-* topology. Let $\bar{X}^{k}$,s have their distributions determined by $\bar{\theta}_{k}$ 's in the same way as $X$ 's distribution is determined by $\theta$. Assume that $\bar{X}_{0}^{k}=z_{k}, X_{0}=z_{0}$ and $z_{k} \rightarrow z_{0}$ as $k \rightarrow \infty$. Then $\bar{X}^{k}$ converge weakly to $X$ in $M_{1}^{\mathcal{T}}$ topology.

We will call the process $X$ obtained in Theorem 3.5 ORBM with reflection angle $\theta$.

REMARK 3.6. (i) Note that the distribution of $X$ in Theorem 3.5(i) does not depend on the approximating sequence $\theta_{k}$ because if we have two sequences $\left\{\theta_{k}\right\}$ and $\left\{\bar{\theta}_{k}\right\}$ converging to $\theta$ then we can apply the theorem to the sequence $\theta_{1}, \bar{\theta}_{1}, \theta_{2}, \bar{\theta}_{2}, \ldots$.
(ii) Suppose that $z_{0} \in D_{*}, \mu_{0} \in \mathbb{R}$, and $h$ is positive and harmonic in $D_{*}$ with $h(0)=1 / \pi$. By Theorem 2.1, we can find $\theta \in \mathcal{T} \leftrightarrow\left(h, \mu_{0}\right) \in \mathcal{H}$. Let $X$ be the process corresponding to $z_{0}$ and $\theta$ as in Theorem 3.5. Then $X$ has a stationary distribution with the density $h$ and $\mu_{0}$ is the rate of rotation of $X$ in the sense of Theorem 3.5(v)-(vi).
(iii) Theorem 3.5 establishes existence of ORBM for all angles $\theta$ of oblique reflection. ORBMs can be uniquely parametrized either by $\theta \in \mathcal{T}$ or by pairs $\left(h, \mu_{0}\right) \in \mathcal{H}$. We will write $X \leftrightarrow \theta$ or $X \leftrightarrow\left(h, \mu_{0}\right)$.
(iv) If $\theta=\pi / 2$ a.e. on an open arc $I \subset \partial D_{*}$ then as in the proof of Proposition 2.11, $\theta+i \widetilde{\theta}$ extends to be analytic across $I$, and hence so does $G=e^{i(\theta+i \widetilde{\theta})}$. In this case, for $x \in I$,

$$
\begin{equation*}
\lim _{r \rightarrow 1} \frac{e^{-\tilde{\theta}(r x)} \cos \theta(r x)}{r-1}=\operatorname{Re} \lim _{r \rightarrow 1} \frac{G(r x)-G(x)}{r x-x} x=\operatorname{Re} G^{\prime}(x) x \tag{3.14}
\end{equation*}
$$

Thus, the integral in (3.13) is finite for each $x \in I$. A similar statement holds if $\theta=-\pi / 2$ a.e. on $I$.

Note that the process $X$ itself will not hit a fixed point $x \in I$. The reason is that $X$ has only a countable number of excursions from the boundary of $\partial D_{*}$ and the distribution of the location of the endpoint of an excursion has a density. Hence, with probability 1 , no excursion will end at $x$. If an excursion ends at a point in $I$, the process $X$ will jump at that time to an end of the interval where $\theta=\pi / 2$ a.e. Thus, $X$ itself will avoid $x$ forever but the same argument shows that $x \in \Gamma_{X}^{\theta}$ with probability 1 because $\Gamma_{X}^{\theta}$ contains the arcs between the endpoints of excursions hitting points inside $I$ and the points to which $X$ jumps at those times.
(v) Let $v$ be the positive measure on $\partial D_{*}$ defined by $h(z)=\int_{\partial D_{*}} K_{x}(z) \nu(d x)$, where $K_{x}(z)$ is the Poisson kernel for $z \in D_{*}$. Fix $x \in \partial D_{*}$ and write

$$
h(r x)=c \frac{1+r}{1-r}+\int_{\partial D_{*}} \frac{1-r^{2}}{|y-r x|^{2}} d \sigma(y)
$$

where $\sigma$ is a positive measure with $\sigma(\{x\})=0$ and $c=v(\{x\})$. Then

$$
\begin{equation*}
\lim _{r \rightarrow 1}(1-r) h(r x)=2 c \tag{3.15}
\end{equation*}
$$

as can be seen by splitting the integral into $\int_{I}+\int_{\partial D_{*} \backslash I}$ where $x \in I$ and $\sigma(I)<\varepsilon$. If $c=v(\{x\})>0$, then

$$
\int_{0}^{1} \frac{h(r x)}{h(r x)^{2}+\left(\widetilde{h}(r x)-\mu_{0} / \pi\right)^{2}} \frac{d r}{1-r} \leq \int_{0}^{1} \frac{1}{(1-r) h(r x)} d r<\infty
$$

and so $x \in \Gamma_{X}^{\theta}$ with probability 1 by Theorem 3.5 (vii), where $X \leftrightarrow\left(h, \mu_{0}\right)$.
(vi) The condition $v(\{x\})>0$ is stronger than the integrability condition (3.13). For example, if $h(z)=\frac{1}{\pi} \operatorname{Re}(1-z)^{-p}$, with $0<p<1$, then $\int_{0}^{1} \frac{1}{(1-r) h(r)} d r<$ $\infty$ so that (3.13) holds at $x=1$. However, by (3.15), the corresponding positive measure $v$ satisfies $v(\{1\})=0$.
(vii) Suppose $\mu_{0}=0$. If $h(x)=h(\bar{x})$ for all $x \in \partial D_{*}$, where $\bar{x}$ denotes the complex conjugate of $x$, then $\widetilde{h}(r)=0$ for $-1<r<1$. In this case, the integral in (3.13) is finite for $x=1$ if and only if

$$
\begin{equation*}
\int_{0}^{1} \frac{1}{(1-r) h(r)} d r<\infty \tag{3.16}
\end{equation*}
$$

Condition (3.13) can be restated. Set $f=\operatorname{Re}\left(1 /\left(h+i \tilde{h}-i \mu_{0} / \pi\right)\right)$. Then $f$ is harmonic and positive, so there is a positive measure $d \sigma$ such that

$$
f(z)=\int_{\partial D_{*}} \frac{1-|z|^{2}}{|w-z|^{2}} d \sigma(w)
$$

Proposition 3.7. Condition (3.13) holds for $x \in \partial D_{*}$ if and only if

$$
\begin{equation*}
\int_{\partial D_{*}} \frac{1}{|w-x|} d \sigma(w)<\infty . \tag{3.17}
\end{equation*}
$$

For example, suppose $E$ is a closed subset of $\partial D_{*}$ of positive length. Let $f(y)=$ $\operatorname{dist}(y, E)^{p}$ for $y \in \partial D_{*}$, where $p \in(0,1)$ is fixed. Then it is not hard to verify that $f \in C^{p}\left(\partial D_{*}\right)$, that is, $f$ is Hölder-continuous with exponent $p$ on $\partial D_{*}$. Let the harmonic extension of $f$ to $D_{*}$ be also denoted by $f$. Thus, the function $f+i \widetilde{f}$ is analytic on $D_{*}$, extends to be continuous on the closed disk $\bar{D}_{*}$, and hence the zero set $Z=\left\{y \in \partial D_{*}: f(y)=\tilde{f}(y)=0\right\} \subset E$ has zero length [see Hoffman (1962), page 51]. Set $h+i \widetilde{h}=1 /(f+i \widetilde{f})$. Then $h$ is positive and harmonic on $D_{*}$. Since $f \in C^{p}\left(\partial D_{*}\right), \tilde{f} \in C^{p}\left(\partial D_{*}\right)$ by Theorem II.3.2 in Garnett and Marshall (2005). Thus, $h=f /\left(f^{2}+\widetilde{f}^{2}\right)$ is continuous up to $\partial D_{*} \backslash Z$, and so $h$ tends to 0 as $z \rightarrow E \backslash Z$. The function $h$ tends to a positive number at each point of $\partial D_{*} \backslash E$. The positive measure $\sigma(d y)=f(y)|d y|$ on $\partial D_{*}$ satisfies (3.17) for each $x \in E$, since $f(y) \leq|x-y|^{p}$ for every $x \in E$. Let $\theta \in \mathcal{T} \leftrightarrow(h, 0) \in \mathcal{H}$ and $X \leftrightarrow(h, 0)$ be the corresponding ORBM. By Theorem 3.5(vii) and Proposition 3.7, for every $x \in E, x \in \Gamma_{X}^{\theta}$ with probability 1 . Note also that the integral in (3.17) is infinite for each point $x \in \partial D_{*} \backslash E$, since $f$ is positive and continuous there. So for every $x \in \partial D_{*} \backslash E$, this ORBM does not hit $x$ with probability 1 . The function $\theta$ is continuous on $\partial D_{*} \backslash Z$, and $|\theta|<\pi / 2$ off $E$. We can take $E$ to have no interior in $\partial D_{*}$, so $|\theta|<\pi / 2$ on a dense open set.

Recall that if $f: D_{*} \rightarrow D_{*}$ is a conformal map of $D_{*}$ onto itself, then there exist $\theta_{0} \in[0,2 \pi)$ and $w_{0} \in D_{*}$ such that $f(z)=e^{i \theta_{0}} \frac{z-w_{0}}{1-\overline{w_{0}} z}$. So in particular $f$ extends continuously to $\bar{D}_{*}$ as a smooth homeomorphism. The following result establishes conformal invariance of ORBM on the unit disk.

Theorem 3.8. Suppose $\theta \in \mathcal{T}$ and $X$ is an ORBM on $D_{*}$ with reflection angle $\theta$. Suppose $f: D_{*} \rightarrow D_{*}$ is a conformal map of $D_{*}$ onto $D_{*}$. Define for $t \in[0, \infty)$,

$$
\begin{equation*}
c(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s \quad \text { and } \quad Y_{t}=f\left(X_{c^{-1}(t)}\right) \tag{3.18}
\end{equation*}
$$

Then $Y$ is an ORBM on $D_{*}$ with reflection angle $\theta \circ f^{-1} \in \mathcal{T}$. Equivalently, if $\left(h, \mu_{0}\right) \in \mathcal{H} \leftrightarrow \theta$, then $Y$ is the ORBM on $D_{*}$ parametrized by $\left(\bar{h}, \bar{\mu}_{0}\right) \in \mathcal{H}$, where $\bar{h}(z)=h\left(f^{-1}(z)\right) /\left(\pi h\left(f^{-1}(0)\right)\right)$ and $\bar{\mu}_{0}=\mu\left(f^{-1}(z)\right) / h\left(f^{-1}(0)\right)$. Here, $\mu(w)$ is the harmonic function defined by (2.24).
3.3. Excursion reflected Brownian motions. We now address the question that was left unanswered in Section 3.2, namely whether there exists a process on $D_{*}$ associated with a purely tangential angle of reflection, for example, $\theta \equiv \pi / 2$. In Theorem 3.12, we will show that such a process does indeed exist and can be obtained as a suitable limit of ORBMs in $D_{*}$ corresponding to angles of reflection $\theta \in \mathcal{T}$. We refer to this process as excursion reflected Brownian motion (ERBM).

We will first define ERBM more generally, in a bounded simply connected domain $D$ with variable excursion intensity $\nu(d x)$, where $v$ is a measure on $\partial D$. Our
construction resembles a process introduced in Chen and Fukushima (2008), Chen, Fukushima and Ying (2007), Fukushima and Tanaka (2005) and called "Brownian motion extended by darning" (BMD), and defined simultaneously in Lawler (2006) under the name of ERBM. We will use some concepts from excursion theory reviewed in Section 2.2.

Definition 3.9. Suppose that $v(d x)$ is a finite positive measure on $\partial D$. Let $H^{x}$ be the standard Brownian excursion law in $D$ for excursions starting at $x \in \partial D$. If $D=D_{*}$ then we normalize the $\sigma$-finite measures $H^{x}, x \in \partial D_{*}$, so that all of them can be obtained from $H^{1}$ by rotation around 0 . Let $\Delta$ be a cemetery state and $\mathcal{C}=\mathcal{C}_{D}$ denote the family of all functions $\omega:[0, \infty) \rightarrow \bar{D} \cup\{\boldsymbol{\Delta}\}$ such that $\omega(0) \in$ $\partial D, \omega$ is continuous up to its lifetime $\zeta<\infty$, and $\omega(t)=\boldsymbol{\Delta}$ for $t \geq \zeta$. Let $\lambda$ denote the Lebesgue measure on $\mathbb{R}_{+}=[0, \infty)$ and let $\mathcal{P}$ be the Poisson point process on $\mathbb{R}_{+} \times \mathcal{C}$ with characteristic measure $\lambda \times \int_{\partial D} H^{x} v(d x)$. With probability 1 , there are no two points with the same first coordinate so the elements of $\mathcal{P}$ may be unambiguously denoted by $\left(t, \mathrm{e}_{t}\right)$. Let

$$
\zeta_{t}=\inf \left\{s>0: \mathrm{e}_{t}(s)=\Delta\right\}
$$

Let $\sigma_{v}=\sum_{s \leq v} \zeta_{s}$ and $\sigma_{v-}=\sum_{u<v} \zeta_{u}$ for $v \geq 0$.
Let $D^{\partial}:=D \cup\{\partial\}$ be a one-point compactification of $D$ obtained by identifying the usual boundary $\partial D$ with a single point $\partial$.

If $D=D_{*}$, then the lifetimes of excursions of the process $\mathcal{P}$ have the same structure as those of the symmetric reflected Brownian motion (with the normal reflection), so $\sigma_{v}<\infty$ for all $v<\infty$ and $\lim _{v \rightarrow \infty} \sigma_{v}=\infty$, a.s. For all domains $D$ for which the last two statements are true, with probability 1 , for every $t \geq 0$, the formula $r=\inf \left\{v \geq 0: \sigma_{v} \geq t\right\}$ defines a unique $r \geq 0$. For $t \geq 0$, let

$$
X_{t}= \begin{cases}\mathrm{e}_{r}\left(t-\sigma_{r-}\right), & \text { if } \sigma_{r-}<\sigma_{r} \text { and } t \in\left[\sigma_{r-}, \sigma_{r}\right), \\ \partial, & \text { otherwise } .\end{cases}
$$

With probability one, $X$ is a conservative process taking values in $D^{\partial}$. We will call the process $X$ (or its distribution) excursion reflected Brownian motion (ERBM) in $D$ with excursion intensity $\nu$. In general, $X$ is not a Hunt process on $\bar{D}$ as it does not have the quasi-left continuity property at the first hitting time of $\partial D$, which is a predictable stopping time. However, $X$ is a conservative continuous Hunt process on $D^{\partial}$.

REMARK 3.10. (i) If $H^{x}$ is a standard Brownian excursion law in $D$ and $c>0$ is a constant, then $c H^{x}$ is also a standard Brownian excursion law in $D$. We talked about "the" standard excursion laws above because all standard excursion laws in a simply connected domain corresponding to a given boundary point are constant multiples of each other.
(ii) For any strictly positive function $a(x)$ on the boundary of $D$, ERBM corresponding to $\left(a(x) v(d x),(1 / a(x)) H^{x}\right)_{x \in \partial D}$ has the same distribution as ERBM determined by $\left(v(d x), H^{x}\right)_{x \in \partial D}$. Hence, one has to specify both $v$ and the normalization of the excursion laws $H^{x}$ to identify ERBM uniquely.
(iii) It may be surprising at the first sight but it is easy to see that for any constant $c>0,\left(v(d x), H^{x}\right)_{x \in \partial D}$ and $\left(c v(d x), H^{x}\right)_{x \in \partial D}$ define the same ERBM. So we may assume that $v$ is a probability measure.
(iv) Combining the last two remarks, it is easy to check that if ERBM $X$ can be represented by $\left(v(d x), H^{x}\right)_{x \in \partial D}$ and also by $\left(v_{1}(d x), H_{1}^{x}\right)_{x \in \partial D}$, then

$$
\left(v_{1}(d x), H_{1}^{x}\right)_{x \in \partial D} \equiv\left(c a(x) \nu(d x),(1 / a(x)) H^{x}\right)_{x \in \partial D}
$$

for some positive function $a(x)$ and some positive constant $c$.
(v) When $D$ is the unit ball $D_{*}$, the ERBM in $D_{*}$ with excursion intensity $v$ being the uniform measure on $\partial D_{*}$ has the same distribution as the BMD studied in Chen and Fukushima (2008), Chen, Fukushima and Ying (2007), Fukushima and Tanaka (2005); see Chen and Fukushima (2012), Remark 7.6.4, where this identification is proved when $D$ is the exterior of the unit ball. When $D$ is the exterior of the unit ball, the process also has the same distribution as the ERBM introduced in Lawler (2006); see Chen and Fukushima (2015), Example 6.3.
(vi) To make things simple, we will assume in theorems on ERBM that $\partial D$ is a Jordan curve (in other words, $D$ is a simply connected Jordan domain). This is equivalent to saying that if $f: D_{*} \rightarrow D$ is a one-to-one and onto analytic mapping then $f$ can be extended to be continuous and one-to-one on $\bar{D}_{*}$. We believe that all our results hold for arbitrary bounded simply connected domains because "exotic" points on the boundary are negligible from the point of view of excursion theory.
(vii) The reader who wishes to learn more about potential theoretic properties of domains and their relationship to geometric properties may consult Ohtsuka (1970) for a discussion of "prime ends." The Martin boundary is presented in Doob (1984); in particular, the identification of the Martin boundary and prime ends is mentioned in Doob (1984), 1 XII 3. The Martin topology and boundary in simply connected planar domains are conformally invariant; see Pommerenke (1975), Theorem 9.6.
(viii) If $D$ is a Jordan domain and $x \in \partial D$, then the Martin kernel $K_{x}(\cdot)$ is the unique, up to a multiplicative constant, positive harmonic function in $D$ that vanishes everywhere on the boundary except at $x$. The density of the expected occupation measure for $H^{x}$ is a constant multiple of the Martin kernel $K_{x}(\cdot)$ by Burdzy (1987), Proposition 3.4.

Proposition 3.11. Suppose $D \subset \mathbb{C}$ is a bounded simply connected Jordan domain:
(i) Let $X$ be an ERBM constructed from $\left(\nu, H^{x}\right)_{x \in D}$, where $v$ is a probability measure on $\partial D$. Then $X$ has a unique stationary distribution whose density is proportional to $h(y)=\int_{\partial D} K_{x}(y) v(d x)$.
(ii) For every positive harmonic function $h$ in $D$ with $\|h\|_{L^{1}(D)}=1$ there exists an ERBM $X$ with the stationary density $h$.

We say that a real-valued function $f$ defined on a subset $S$ of $\mathbb{R}^{n}$ is Lipschitz with constant $\lambda<\infty$ if $|f(x)-f(y)| \leq \lambda|x-y|$ for all $x, y \in S$. It follows from the definitions that a Lipschitz function is Dini continuous.

THEOREM 3.12. (i) Consider a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow$ $(-\pi / 2, \pi / 2)$ and let $X^{k}$ be defined by

$$
\begin{equation*}
X_{t}^{k}=x_{k}+B_{t}+\int_{0}^{t} \mathbf{v}_{\theta_{k}}\left(X_{s}^{k}\right) d L_{s}^{k} \quad \text { for } t \geq 0 \tag{3.19}
\end{equation*}
$$

Let $\left(h_{k}, \mu_{0, k}\right) \leftrightarrow \theta_{k}$ as in Lemma 2.2. We make the following assumptions:
(a) $\theta_{k}$ converge to $\pi / 2$ almost everywhere.
(b) For some $c_{1}>-\pi / 2$ and all $x$ and $k, \theta_{k}(x) \geq c_{1}$.
(c) There exist $\lambda<\infty$ and $c_{2}>0$ such that $h_{k}$ restricted to $\partial D_{*}$ is Lipschitz with constant $\lambda$ for every $k$, and $h_{k}(x)>c_{2}$ for every $x$ and $k$.
(d) There is a finite measure $\nu(d x)$ on $\partial D_{*}$ such that $h_{k}(x) d x \rightarrow \nu(d x)$ weakly as measures on $\partial D_{*}$, when $k \rightarrow \infty$.
(e) $\lim _{k \rightarrow \infty} \operatorname{dist}\left(x_{k}, \partial D_{*}\right)=0$.

Then the processes $X^{k}$ converge in the sense of finite dimensional distributions to ERBM $X$ corresponding to $\left(v(d x), H^{x}\right)_{x \in \partial D_{*}}$, where all $H^{x}$ are obtained from $H^{1}$ by rotation around 0 .
(ii) Conversely, suppose that $h$ is harmonic in $D_{*}$, Lipschitz on $\bar{D}_{*}$ and positive on $\bar{D}_{*}$. Then there exists a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ satisfying conditions (a)-(e) with $v(d x)=h(x) d x$ on $\partial D_{*}$. ORBMs $X^{k}$ corresponding to $\theta_{k}$ 's converge in the sense of finite dimensional distributions to an ERBM X with the stationary density $h$.

REMARK 3.13. (i) The roles of $\pi / 2$ and $-\pi / 2$ in Theorem 3.12 can be reversed by replacing $\theta_{k}(x)$ with $-\theta_{k}(\bar{x})$. See Remark 2.6(iv).
(ii) It is easy to see from Theorems 3.8 and 3.12 that if $f: D_{*} \rightarrow D_{*}$ is a conformal map and $X$ is an ERBM on $D_{*}$ corresponding to $\left(\nu(d x), H^{x}\right)_{x \in \partial D_{*}}$, then $f(X)$ is a time-change of ERBM on $D_{*}$ corresponding to $\left(v(d x) \circ f^{-1}, H^{x}\right)_{x \in \partial D_{*}}$.
(iii) Suppose that there exists $\lambda<\infty$ such that $h_{k}$ restricted to $\partial D_{*}$ is Lipschitz with constant $\lambda$ for every $k$. Then it is elementary to show that there exists $c_{2}>0$ such that $h_{k}(x)>c_{2}$ for every $x$ and $k$ if and only if there exists $\lambda_{1}<\infty$ such that $1 / h_{k}$ restricted to $\partial D_{*}$ is Lipschitz with constant $\lambda_{1}$ for every $k$.

EXAMPLE 3.14. Theorem 3.12 has many assumptions so it deserves a simple example to illustrate it. Suppose $h(x)$ and $1 / h(x)$ are positive Lipschitz continuous functions on $\partial D_{*}$ with $\|h\|_{L^{1}\left(D_{*}\right)}=1$. Let $h(z)$ be the harmonic extension of $h$
to $D_{*}$. Suppose also that $\mu_{0, k} \rightarrow \infty$ as $k \rightarrow \infty$. Then $\left(h, \mu_{0, k}\right) \leftrightarrow \theta_{k} \in \mathcal{T}$ as in Theorem 2.1. If $h_{k}=h$ for all $k$, then ( $h_{k}, \mu_{0, k}$ ) and $\theta_{k}$ satisfy assumptions (a)-(e) of Theorem 3.12.
3.4. ORBMs in simply connected domains. We will use conformal mappings to construct ORBMs in arbitrary simply connected domains. In the following, we will usually use $X$ to denote ORBM in the disk $D_{*}$ and $Y$ to denote ORBM in other domains.

THEOREM 3.15. Suppose that $f$ is a one-to-one analytic function mapping $D_{*}$ onto a simply connected domain $D \subset \mathbb{C}$. Suppose that $\theta \in \mathcal{T}, \theta \leftrightarrow(h, \mu)$, let $\bar{h}=h \circ f^{-1}$ and assume that $\bar{h}$ is in $L^{1}(D)$. Let $X \leftrightarrow \theta$ be ORBM in $D_{*}$ and define

$$
\begin{align*}
c(t) & =\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s \quad \text { for } t \geq 0  \tag{3.20}\\
\zeta & =\inf \{t \geq 0: c(t)=\infty\}  \tag{3.21}\\
Y_{t} & =f\left(X_{c^{-1}(t)}\right) \quad \text { for } t \in[0, \zeta) . \tag{3.22}
\end{align*}
$$

We will call $Y$ an ORBM in D. The following hold:
(i) With probability $1, \zeta=\infty$.
(ii) The process $Y$ is an extension of killed Brownian motion in $D$ in the sense that for every $t \geq 0$ and $\tau_{t}=\inf \left\{s \geq t: Y_{s} \in \partial D\right\}$, the process $\left\{Y_{s}, s \in\left[t, \tau_{t}\right)\right\}$ is Brownian motion killed upon exiting $D$.
(iii) The process $Y$ has a stationary distribution with density $\widehat{h}=\bar{h} /\|\bar{h}\|_{L^{1}(D)}$.
(iv) Recall that $\mu$ is the function given by (2.24). For $z \in D$, let arg $^{*}\left(Y_{t}-z\right)=$ $\boldsymbol{\operatorname { a r g }}^{*}\left(X_{c^{-1}(t)}-f^{-1}(z)\right)$ for all $t$. Then, for every $z \in D$, a.s.

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(Y_{t}-z\right)}{t}=\frac{\mu\left(f^{-1}(z)\right)}{\|\bar{h}\|_{L^{1}(D)}} \tag{3.23}
\end{equation*}
$$

(v) Suppose that $\mu_{0} \in \mathbb{R}$ and $\widehat{h}$ is a positive harmonic function in $D$ with $\|\widehat{h}\|_{L^{1}(D)}=1$. Then there exists an ORBM Y in $D$ with the following properties:
(a) The stationary distribution of $Y$ is $\widehat{h}(x) d x$.
(b) Set $g=f^{-1}$ and define

$$
\begin{align*}
b(t) & :=\int_{0}^{t} \mid\left(\left.g^{\prime}\left(Y_{s}\right)\right|^{2} d s, \quad t \geq 0,\right.  \tag{3.24}\\
X_{t} & :=g\left(Y_{b^{-1}(t)}\right), \quad t \geq 0,  \tag{3.25}\\
\arg ^{*} Y_{t} & :=\arg ^{*} X_{b(t)}, \quad t \geq 0 . \tag{3.26}
\end{align*}
$$

Since $\widehat{h} \circ f$ is a positive harmonic function on $D_{*},\|\widehat{h} \circ f\|_{1}=\pi \widehat{h} \circ$ $f(0)<\infty$. Set $h_{1}=\widehat{h} \circ f /\|\widehat{h} \circ f\|_{1}$ and let $\mu \in \mathcal{R} \leftrightarrow\left(h_{1}, \mu_{0}\right) \in \mathcal{H}$. Then
$\underline{X}$ is the ORBM in $D_{*}$ parametrized by $\left(h_{1}, \mu_{0}\right)$ and (3.23) holds with $\bar{h}=h_{1} \circ f^{-1}=\widehat{h} /\|\widehat{h} \circ f\|_{1}$.
(vi) (Consistence) If $D$ has a smooth boundary and $\theta$ is $C^{2}$ then the distribution of $Y$ is the same as that of the process identified in Theorem 3.1(ii) relative to $\theta \circ f^{-1}$.

REMARK 3.16. (i) The quantity $\boldsymbol{\operatorname { a r g }}\left(Y_{t}-z\right)$ has a natural interpretation when $Y$ is continuous, namely, $\boldsymbol{\operatorname { a r g }}\left(Y_{t}-z\right)-\arg \left(Y_{0}-z\right)$ is the number of windings of $Y$ around $z$ over the time interval $[0, t]$. The quantity $\arg ^{*}\left(Y_{t}-z\right)$ is obtained from $\boldsymbol{\operatorname { a r g }}\left(Y_{t}-z\right)$ by discarding (the windings of) all excursions of $Y$ which make more than a full loop around $z$ (from endpoint to endpoint of the excursion, not within the excursion). Our definition of $\mathcal{E}_{s}^{L}$ was chosen to make this simple geometric interpretation of arg $^{*}\left(Y_{t}-z\right)$ possible.

Unfortunately, when $Y$ is not continuous, $\mathbf{a r g}^{*}\left(Y_{t}-z\right)$ does not have a simple intuitive interpretation because the definition of $\arg$ in $D_{*}$ depends on $\theta$.
(ii) The process $Y$ constructed in Theorem 3.15 will be called ORBM in $D$. The family of ORBMs in $D$ can be parametrized either in terms of pairs $(\theta, f)$ or triplets $\left(\widehat{h}, \mu_{0}, f\right)$, so we will write $Y \leftrightarrow(\theta, f)$ or $Y \leftrightarrow\left(\widehat{h}, \mu_{0}, f\right)$. The function $f$ provides a way to parametrize $\partial D$, in a sense.
(iii) If $\mu \in \mathcal{R} \leftrightarrow\left(h, \mu_{0}\right) \in \mathcal{H}$ then we say that $\mu \circ f^{-1}(z)$ is the rotation rate about $z \in D$ for the process $Y$ given by (3.22). If $\mu_{1}$ is a harmonic function defined on $D$, let $\tilde{\mu}_{1}$ be the harmonic conjugate of $\mu_{1}$ vanishing at $f(0)$. Then $\widetilde{\mu_{1} \circ f}$ is a harmonic function on $D_{*}$ vanishing at 0 and $\widetilde{\mu_{1} \circ f}=\widetilde{\mu_{1}} \circ f$. By Theorems 2.1 and $3.15, \mu_{1}$ is a rotation rate for an ORBM if and only if $\tilde{\mu}_{1}(z)>-1$ for all $z \in D$.
(iv) Suppose that $f$ is a conformal mapping from a bounded simply connected planar domain $D_{1}$ to another bounded simply connected planar domain $D_{2}$. Let $\mathcal{K}\left(D_{1}, D_{2}, f\right)$ be the family of positive integrable harmonic functions $h$ in $D_{1}$ such that $h \circ f^{-1} \in L^{1}\left(D_{2}\right)$. By Theorems 3.8 and 3.15, $f$ establishes a correspondence between a subfamily of ORBMs on $D_{1}$ that have the density of stationary distribution in $\mathcal{K}\left(D_{1}, D_{2}, f\right)$ and a subfamily of ORBMs on $D_{2}$ that have the density of stationary distribution in $\mathcal{K}\left(D_{2}, D_{1}, f^{-1}\right)$. The subfamilies are nonempty because they always contain normally reflected Brownian motions. Theorem 3.20 below gives some sufficient conditions on the integrability of positive harmonic functions in domains. The correspondence between ORBMs on different planar domains need not extend to all ORBMs on either side because the assumption $\bar{h} \in L^{1}(D)$ of Theorem 3.15 does not hold for some $h$ and $f$; see Example 4.1 below.
(v) There exist processes in $D$ that are extensions of Brownian motion in $D$, which have a stationary density and a "limiting rate of rotation" $\mu_{0}$ and which are not ORBMs. An example of such a process is the conformal image of reflected Brownian motion in $D_{*}$ with diffusion on the boundary [see a Ph.D. thesis Card (2009) devoted to this class of processes].

THEOREM 3.17. Suppose that $D \subset \mathbb{C}$ is a simply connected bounded Jordan domain and $f$ is a conformal mapping from $D_{*}$ onto $D$, which, by Carathéodory's theorem, necessarily extends to a homeomorphism from $\bar{D}_{*}$ onto $\bar{D}$. Consider a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ and processes $X^{k}$ which satisfy (3.19) and assumptions (a)-(e) of Theorem 3.12. Let $\left(h_{k}, \mu_{k}\right) \leftrightarrow \theta_{k}$ and let $c_{k}(t), \zeta_{k}$ and $Y^{k}$ be defined relative to $\theta_{k}, f$ and $X^{k}$ as in Theorem 3.15.

Let $v, h$ and $X$ be defined as in Theorem 3.12. Let $c(t), \zeta$ and $Y$ be defined relative to $\theta, f$ and $X$ as in Theorem 3.15. In (i)-(iv) below, $\vec{h}:=h \circ f^{-1}$ is assumed to be in $L^{1}(D)$ :
(i) Almost surely, $\zeta_{k}=\infty$ for every $k \geq 1$ and $\zeta=\infty$.
(ii) The process $Y$ is an ERBM in $D$ corresponding to $\left(\bar{v}(d x), \bar{H}^{x}\right)_{x \in \partial D}$ with excursion intensity $\bar{v}$ defined by $\bar{v}(A)=v\left(f^{-1}(A)\right)$ for $A \subset \partial D$, and excursion laws $\bar{H}^{x}$ normalized so that the density of the expected occupation time for $\bar{H}^{x}$ is the Martin kernel $K_{x}(\cdot)$ in D normalized by $K_{x}(f(0))=1$.
(iii) Processes $Y^{k}$ converge to $Y$ in the sense of convergence of finite dimensional distributions.
(iv) The process $Y$ has a stationary distribution with the density $\widehat{h}=\bar{h}$ / $\|\bar{h}\|_{L^{1}(D)}$.
(v) For every positive harmonic function $\widehat{h}$ in $D$ with $\|\widehat{h}\|_{L^{1}(D)}=1$ such that $\widehat{h} \circ f$ is Lipschitz on $\bar{D}_{*}$ and strictly positive on $\partial D_{*}$, there is a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ such that $Y^{k}$ and $Y$ can be constructed as in the initial part of the theorem and the stationary measure for ERBM Y has density $\widehat{h}$.

The next two theorems show that ORBM in an arbitrary domain (possibly with a fractal boundary) can be approximated by ORBMs in smooth domains where the oblique angle of reflection has a natural interpretation. This provides a justification of the name "obliquely reflected Brownian motion" for processes in domains with rough boundaries.

THEOREM 3.18. Suppose that $D \subset \mathbb{C}$ is a simply connected Jordan domain, $y_{0} \in D$ and $f$ is a conformal mapping from $D_{*}$ onto $D$ which, necessarily, has a one-to-one continuous extension to $\bar{D}_{*}$. Let $D_{k}$ be simply connected domains with smooth boundaries such that $y_{0} \in D_{k} \subset D_{k+1} \subset D$ for all $k$ and $\bigcup_{k} D_{k}=D$. Let $f_{k}: D_{*} \rightarrow D_{k}$ be conformal mappings such that $f_{k}^{-1}\left(y_{0}\right)=f^{-1}\left(y_{0}\right)$ and $f_{k} \rightarrow f$ as $k \rightarrow \infty$.

Suppose that $\mu_{0} \in \mathbb{R}, \bar{h} \in L^{1}(D)$ is positive and harmonic with $\|\bar{h}\|_{L^{1}(D)}=$ 1, and $\bar{h} \circ f$ is strictly positive on $\partial D_{*}$. Let $Y$ be the process constructed as in Theorem 3.15(v), relative to $D, f, \mu_{0}$ and $\bar{h}$, with $Y_{0}=y_{0}$. Let $\bar{h}_{k}=\bar{h} /\|\bar{h}\|_{L^{1}\left(D_{k}\right)}$. Let $Y^{k}$ be defined in the same way that $Y$ was defined, relative to $D_{k}, f_{k}, \mu_{0}$ and $\bar{h}_{k}$, with $Y_{0}^{k}=y_{0}$. Then $Y^{k}$ converge weakly to $Y$ in $M_{1}^{\mathcal{T}}$ topology.

The following concrete example shows how one can approximate a general ORBM in $D$ by ORBMs in an increasing sequence of smooth domains with smooth reflection angles. Suppose that $Y \leftrightarrow(\theta, f) \leftrightarrow\left(\bar{h}, \mu_{0}, f\right)$. Take any strictly increasing sequence of positive numbers $r_{k}$ that increases to 1 . Let $D_{k}=$ $f\left(B\left(0, r_{k}\right)\right)$ and $f_{k}(z)=f\left(z / r_{k}\right)$. It is easy to see that $D_{k}$ is a smooth subdomain of $D$ and $f_{k}$ is a conformal mapping from $B\left(0, r_{k}\right)$ to $D$. Clearly, $h_{k}(z):=$ $\bar{h}\left(f\left(r_{k} z\right)\right)$ is a positive harmonic function on $D_{*}$ that is smooth on $\bar{D}_{*}$. By Theorem 2.1, $\theta_{k} \leftrightarrow\left(h_{k} / h_{k}(0), \mu_{0}\right)$ is smooth on $\partial D_{*}$. Thus, $\bar{\theta}_{k}(w)=\theta_{k}\left(f^{-1}(w) / r_{k}\right) \in$ $(-\pi / 2, \pi / 2)$ defines a smooth function on $\partial D_{k}$. Let $Y^{k}$ be the ORBM on $D_{k}$ with reflection angle $\bar{\theta}_{k}$ constructed in Theorem 3.1(ii). Theorem 3.18 asserts that $Y^{k}$ converge weakly to ORBM $Y$ on $D$ in $M_{1}^{\mathcal{T}}$ topology.

THEOREM 3.19. Suppose that $D \subset \mathbb{C}$ is a simply connected Jordan domain, $y_{0} \in D$ and $f: D_{*} \rightarrow D$ is a conformal mapping which, necessarily, has a one-to-one continuous extension to $\bar{D}_{*}$. Let $D_{k}$ be simply connected domains with smooth boundaries such that $y_{0} \in D_{k} \subset D_{k+1} \subset D$ for all $k$ and $\bigcup_{k} D_{k}=D$. Let $f_{k}: D_{*} \rightarrow D_{k}$ be one-to-one analytic functions such that $f_{k}^{-1}\left(y_{0}\right)=f^{-1}\left(y_{0}\right)$ and $f_{k} \rightarrow f$ as $k \rightarrow \infty$.

Suppose that $\theta: \partial D \rightarrow(-\pi / 2, \pi / 2)$ is a continuous function. Let $\theta_{*}: \partial D_{*} \rightarrow$ $(-\pi / 2, \pi / 2)$ be defined by $\theta_{*}=\theta \circ f$. Let $Y$ be ORBM in $D$, such that $Y \leftrightarrow\left(\theta_{*}, f\right)$ and $Y_{0}=y_{0}$.

For every $k$, let $g_{k}: \partial D_{k} \rightarrow \partial D$ be a measurable function such that for every $x \in \partial D_{k}, g_{k}(x)=y \in \partial D$ and $|x-y|=\operatorname{dist}(x, \partial D)$. Let $\theta_{k}(x)=\theta\left(g_{k}(x)\right)$ for $x \in \partial D_{k}$. Let $Y^{k}$ be the ORBM in $D_{k}$ such that $Y^{k} \leftrightarrow\left(\theta_{k}, f_{k}\right)$ and $Y_{0}^{k}=y_{0}$. Then $Y^{k}$ 's converge weakly in $M_{1}$ topology to $Y$.

The assumption that $\bar{h} \in L^{1}(D)$ applied in Theorem 3.15 is sufficient but not necessary. We will sketch an argument illustrating this claim in Example 4.2 below. In other words, the construction given in Theorem 3.15 generates a process $Y_{t}$ for all $t \geq 0$ for some domains $D$ and functions $\bar{h}$ such that $\|\bar{h}\|_{L^{1}(D)}=\infty$. Of course, in such a case no constant multiple of $\bar{h}(x) d x$ can be the stationary (probability) distribution for $Y$, although it can be an invariant measure.

In view of the assumption of integrability of $\bar{h}$ made in Theorems 3.15 and 3.18, it would be useful to have an effective tool to check whether a given harmonic function is in $L^{1}(D)$. We do not have such a test and we doubt that a universal test of this kind exists. We do have some sufficient conditions for integrability of positive harmonic functions. First, recall Theorem 2.10. It contains a criterion for a harmonic function $h$ in $D_{*}$ corresponding to an angle of oblique reflection $\theta$ to be bounded. A "push" $h \circ f^{-1}$ of such function to a bounded simply connected domain is also bounded, and hence integrable. Second, Theorem 3.20 below presents some examples of domains where all positive harmonic functions are integrable.

Recall that a function $\psi: \mathbb{R} \rightarrow \mathbb{R}$ is Lipschitz, with constant $\lambda<\infty$, if $\mid \psi(x)-$ $\psi(y)|\leq \lambda| x-y \mid$ for all $x, y \in \mathbb{R}$. A domain $D \subset \mathbb{R}^{2}$ is said to be Lipschitz,
with constant $\lambda$, if there exists $\delta>0$ such that, for every $x \in \partial D$, there exists an orthonormal basis $\left(e_{1}, e_{2}\right)$ and a Lipschitz function $\psi: \mathbb{R} \rightarrow \mathbb{R}$, with constant $\lambda$, such that

$$
\{y \in \mathcal{B}(x, \delta) \cap D\}=\left\{y \in \mathcal{B}(x, \delta): \psi\left(\left\langle y, e_{1}\right\rangle\right)<\left\langle y, e_{2}\right\rangle\right\} .
$$

We recall the definition of a John domain following Aikawa (2000). Let $\delta_{D}(x)=\operatorname{dist}(x, \partial D)$ and $x_{0} \in D$. We say that $D$ is a John domain with John constant $c_{J}>0$ if each $x \in D$ can be joined to $x_{0}$ by a rectifiable curve $\gamma$ such that $\delta_{D}(y) \geq c_{J} \ell(\gamma(x, y))$ for all $y \in \gamma$, where $\gamma(x, y)$ is the subarc of $\gamma$ from $x$ to $y$ and $\ell(\gamma(x, y))$ is the length of $\gamma(x, y)$. The first two parts of the following theorem follow from Theorems 1 and 2 of Aikawa (2000).

THEOREM 3.20. (i) [Aikawa (2000), Theorem 1] If $D \subset \mathbb{R}^{2}$ is a bounded John domain with John constant $c_{J} \geq 7 / 8$ then all positive harmonic functions in $D$ are in $L^{1}(D)$.
(ii) [Aikawa (2000), Theorem 2] If $D \subset \mathbb{R}^{2}$ is a bounded Lipschitz domain with constant $\lambda<1$ then all positive harmonic functions in $D$ are in $L^{1}(D)$.
(iii) There exists a bounded Lipschitz domain $D$ with constant $\lambda=1$ and $a$ positive harmonic function $h$ in $D$ which is not in $L^{1}(D)$.

## 4. Proofs.

Proof of Theorem 3.1. (i) This part is a special case of Harrison, Landau and Shepp (1985), Theorem 2.6.
(ii) Let $X$ be the unique pathwise solution of (2.1). Then by Itô's formula, $f\left(X_{t}\right)-\frac{1}{2} \int_{0}^{t} \Delta f\left(X_{s}\right) d s$ is a submartingale under $\mathbb{P}_{z}$ for every $z \in \bar{D}$ and $f \in \mathcal{C}$. Thus, in view of (i), $\left(X, \mathbb{P}_{z}\right)$ is the unique solution to the submartingale problem (2.3).
(iii) This part is known; see, for example, Kim, Kim and Yun (1998). For the reader's convenience, we give a sketch of the Dirichlet form approach to the construction of ORBM. The argument given below works in higher dimensions as well. In $C^{2}$-smooth domains with $C^{2}$-smooth reflection angle, it is enough to construct ORBM locally nearly the boundary and then patch the pieces together. Thus, by locally flattening the boundary, we may and do assume that $D=\mathbb{H}$, the upper half space. Let $\mathbf{v}(x)=\left(v_{1}(x), 1\right)$ for $x \in \partial \mathbb{H}$ with $v_{1}(x):=\tan \theta(x)$. Consider a nonsymmetric bilinear form $(\mathcal{E}, \mathcal{F})$ on $L^{2}(\mathbb{H}, d z)$, where

$$
\begin{aligned}
\mathcal{F} & =\left\{f \in L^{2}(\mathbb{H}, d z): \nabla f \in L^{2}(\mathbb{H}, d z)\right\}, \\
\mathcal{E}(f, g) & =\int_{\mathbb{H}} \nabla f(z) \cdot \nabla g(z) d z-\int_{\partial \mathbb{H}} v_{1}(x) \frac{\partial f(x, 0)}{\partial x} g(x, 0) d x \quad \text { for } f, g \in \mathcal{F} .
\end{aligned}
$$

Let $\varepsilon^{0}(f, g)=\int_{\mathbb{H}} \nabla f(z) \cdot \nabla g(z) d z$, and for $\alpha>0$,

$$
\begin{aligned}
& \mathcal{E}_{\alpha}^{0}(f, g):=\mathcal{E}^{0}(f, g)+\alpha(f, g)_{L^{2}(\mathbb{H} ; d z)} \quad \text { and } \\
& \mathcal{E}_{\alpha}(f, g):=\mathcal{E}(f, g)+\alpha(f, g)_{L^{2}(\mathbb{H} ; d z)} .
\end{aligned}
$$

Observe that for $f \in C_{c}^{2}(\overline{\mathbb{H}})$, by the integration by parts formula,

$$
\begin{aligned}
\left|\int_{\partial \mathbb{H}} v_{1}(x) \frac{\partial f(x, 0)}{\partial x} f(x, 0) d x\right| & =\frac{1}{2}\left|\int_{\partial \mathbb{H}} v_{1}^{\prime}(x) f(0, x)^{2} d x\right| \\
& \leq \frac{1}{2}\left\|v_{1}^{\prime}\right\|_{\infty}\|f(x, u)\|_{L^{2}(\partial \mathbb{H}, d x)}^{2}
\end{aligned}
$$

By the boundary trace theorem, for every $\varepsilon>0$ there is $C_{\varepsilon}>0$ such that

$$
\|f(x, u)\|_{L^{2}(\partial \mathbb{H}, d x)}^{2} \leq \varepsilon \mathcal{E}^{0}(f, f)+C_{\varepsilon}\|f\|_{L^{2}(\mathbb{H} ; d z)}^{2} \quad \text { for } f \in \mathcal{F} .
$$

It follows from the above two displays that there are constants $\alpha>0$ and $C_{0} \geq 1$ such that

$$
C_{0}^{-1} \mathcal{E}_{1}^{0}(f, f) \leq \mathcal{E}_{\alpha}(f, f) \leq C_{0} \varepsilon_{1}^{0}(f, f)
$$

for every $f \in C_{c}^{2}(\overline{\mathbb{H}})$ and hence for every $f \in \mathcal{F}$. On the other hand, for $f, g \in$ $C_{c}^{2}(\overline{\mathbb{H}})$,

$$
\begin{aligned}
-\int_{\partial \mathbb{H}} & v_{1}(x) \frac{\partial f(x, 0)}{\partial x} g(x, 0) d x \\
= & -\int_{\partial \mathbb{H}} v_{1}(x) \int_{0}^{\infty} \frac{\partial}{\partial y}\left(\frac{\partial f(x, y)}{\partial x} g(x, y)\right) d y d x \\
(4.1)= & -\int_{\mathbb{H}} v_{1}(x) \frac{\partial f(x, y)}{\partial x} \frac{\partial g(x, y)}{\partial y} d y d x-\int_{\mathbb{H}} v_{1}(x) g(x, y) \frac{\partial^{2} f(x, y)}{\partial x \partial y} d y d x \\
= & \int_{\mathbb{H}} v_{1}(x)\left(\frac{\partial f(x, y)}{\partial x} \frac{\partial g(x, y)}{\partial y}-\frac{\partial f(x, y)}{\partial y} \frac{\partial g(x, y)}{\partial x}\right) d y d x \\
& -\int_{\mathbb{H}} v_{1}^{\prime}(x) g(x, y) \frac{\partial g(x, y)}{\partial y} d y d x .
\end{aligned}
$$

Thus, with $C_{1}=2\|v\|_{\infty},+\left\|v^{\prime}\right\|_{\infty}$,
$\left|\int_{\partial \mathbb{H}} v_{1}(x) \frac{\partial f(x, 0)}{\partial x} g(x, 0) d x\right| \leq C_{1} \varepsilon_{1}^{0}(f, f)^{1 / 2} \varepsilon_{1}^{0}(g, g)^{1 / 2} \quad$ for $f, g \in C_{c}^{2}(\overline{\mathbb{H}})$.
Hence, the bilinear form $(\mathcal{E}, \mathcal{F})$ satisfies the sector condition: there is a constant $C_{2} \geq 1$ such that

$$
|\mathcal{E}(f, g)| \leq C_{2} \mathcal{E}_{\alpha}(f, f)^{1 / 2} \mathcal{E}_{\alpha}(g, g)^{1 / 2} \quad \text { for } f, g \in \mathcal{F}
$$

Moreover, by increasing the value of $\alpha$ if needed, we have from (4.1) that for every $f \in C_{c}^{2}(\overline{\mathbb{H}})$,

$$
\mathcal{E}(f, f-(0 \vee f) \wedge 1) \geq 0 \quad \text { and } \quad \mathcal{E}_{\alpha}(f-(0 \vee f) \wedge 1, f) \geq 0
$$

Thus, $(\mathcal{E}, \mathcal{F})$ is a regular nonsymmetric Dirichlet form on $L^{2}(\overline{\mathbb{H}} ; d z)$. Let $X$ be the Hunt process on $\overline{\mathbb{H}}$ associated with $(\mathcal{E}, \mathcal{F})$. Then one can use stochastic analysis
for nonsymmetric Dirichlet forms to show that $X$ satisfies the SDE (2.1) for quasievery starting point $x \in \overline{\mathbb{H}}$ [see Kim, Kim and Yun (1998)]. Since $X$ behaves like Brownian motion inside $\mathbb{H}$, we can refine the result to allow $X$ to start from every point $x \in \mathbb{H}$ and conclude that (2.1) holds for such $X$.

Proof of Theorem 3.2. (i) This part of our theorem is a special case of Harrison, Landau and Shepp (1985), Theorem 2.18.
(ii) Almost sure continuity of $X$ follows from (2.1).

Recall that we are assuming that $\theta: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ and $\theta \in C^{2}$. It follows from Garnett and Marshall (2005), Corollary II.3.3, that $h$ is $C^{2-\varepsilon}$ on $\bar{D}_{*}$ for every $\varepsilon>0$.

Let $Q$ denote the probability measure on $D_{*}$ with density $h(z)$. We will show that

$$
\begin{equation*}
\mathbb{E}_{Q}\left[\int_{0}^{1} g\left(X_{s}\right) d L_{s}\right]=\int_{\partial D_{*}} g(x)(h(x) / 2) d x \tag{4.2}
\end{equation*}
$$

for every continuous function $g$ on $\partial D_{*}$. Fix any continuous function $g$ on $\partial D_{*}$. Its harmonic extension to $\bar{D}_{*}$ (also denoted $g$ ) is continuous on $\bar{D}_{*}$. Then for $\varepsilon \in(0,1)$,

$$
\begin{equation*}
\mathbb{E}_{Q}\left[\int_{0}^{1} \frac{1}{\varepsilon} \mathbf{1}_{\left\{1-\varepsilon<\left|X_{s}\right|<1\right\}} g\left(X_{s}\right) d s\right]=\int_{D_{*}} \frac{1}{\varepsilon} \mathbf{1}_{\{1-\varepsilon<|z|<1\}} g(z) h(z) d z . \tag{4.3}
\end{equation*}
$$

By continuity and boundedness of $g$ and $h$, the limit of the right-hand side, as $\varepsilon \rightarrow 0$, is equal to $\int_{\partial D_{*}} g(x) h(x) d x$. It is standard to show that $\int_{0}^{1} \frac{1}{\varepsilon} \mathbf{1}_{\left\{1-\varepsilon<\left|X_{s}\right|<1\right\}} \times$ $g\left(X_{s}\right) d s$ converges to $2 \int_{0}^{1} g\left(X_{s}\right) d L_{s}$ in distribution as $\varepsilon \rightarrow 0$. We claim that the family

$$
\begin{equation*}
\left\{\int_{0}^{1} \frac{1}{\varepsilon} \mathbf{1}_{\left\{1-\varepsilon<\left|X_{s}\right|<1\right\}} g\left(X_{s}\right) d s, \varepsilon \in(0,1 / 2)\right\} \tag{4.4}
\end{equation*}
$$

is uniformly integrable. Since $g$ is bounded, it suffices to prove uniform integrability of the family $\left\{\int_{0}^{1} \frac{1}{\varepsilon} \mathbf{1}_{\left\{1-\varepsilon<\left|X_{s}\right|<1\right\}} d s, \varepsilon \in(0,1 / 2)\right\}$. If we denote by $\mathcal{L}_{t}^{a}$ the local time of the two-dimensional Bessel process on $[0,1]$ reflected at 1 , then the distribution of $\int_{0}^{1} \frac{1}{\varepsilon} \mathbf{1}_{\left\{1-\varepsilon<\left|X_{s}\right|<1\right\}} d s$ is the same as $\frac{1}{\varepsilon} \int_{1-\varepsilon}^{1} \mathcal{L}_{1}^{a} d a$. The last random variable is stochastically majorized by $\sup \left\{\mathcal{L}_{1}^{a}: a \in[1 / 2,1]\right\}$ for every $\varepsilon \in(0,1 / 2)$. A version of the Trotter and Ray-Knight theorems shows that $\mathcal{L}_{1}^{a}$ is a diffusion in $a$, $\operatorname{so} \sup \left\{\mathcal{L}_{1}^{a}: a \in[1 / 2,1]\right\}$ is an almost surely finite random variable. Therefore, the family in (4.4) is uniformly integrable. Taking $\varepsilon \rightarrow 0$ in (4.3) yields (4.2). It follows that the Revuz measure of $L$ is $\frac{1}{2} h(x) d x$ on $\partial D_{*}$, relative to the invariant measure $h(z) d z$ on $D_{*}$.

We will now provide a representation of $X$ using a map which is locally conformal. Let $D_{-}=\{z \in \mathbb{C}: \operatorname{Re} z<0\}$ be the left half-plane and $f(z)=\exp (z)$ the exponential function that maps $D_{-}$onto $D_{*} \backslash\{0\}$. For $x \in \partial D_{-}$such that
$f(x)=z \in \partial D_{*}$, define $\widehat{\mathbf{v}}(x)=i \tan \theta(z)-1$. Note that $\widehat{\mathbf{v}}(x)$ is a periodic $C^{2}-$ smooth function on $\partial D_{-}$with period $2 \pi i$. Suppose that $\widehat{x}_{0} \in \bar{D}_{-}$and $\widehat{B}$ is a twodimensional Brownian motion. It is known [see Lions and Sznitman (1984), Theorem 4.3] that there is a pathwise unique solution $(\widehat{X}, \widehat{L})$ to the following Skorokhod SDE:

$$
\begin{equation*}
\widehat{X}_{t}=\widehat{x}_{0}+\widehat{B}_{t}+\int_{0}^{t} \widehat{\mathbf{v}}\left(\widehat{X}_{s}\right) d \widehat{L}_{s} \tag{4.5}
\end{equation*}
$$

where $\widehat{X}$ is a continuous process that takes values in $\bar{D}_{-}$and $\widehat{L}$ is a continuous nondecreasing real-valued process with $\widehat{L}_{0}=0$ that increases only when $\widehat{X}_{t} \in$ $\partial D_{-}$. The process $\widehat{X}$ is an ORBM in $D_{-}$with the oblique angle of reflection $\theta \circ f$. The Itô formula yields

$$
\begin{align*}
f\left(\widehat{X}_{t}\right) & =f\left(\widehat{X}_{0}\right)+\int_{0}^{t} f^{\prime}\left(\widehat{X}_{s}\right) d \widehat{B}_{s}+\int_{0}^{t} f^{\prime}\left(\widehat{X}_{s}\right) \widehat{\mathbf{v}}\left(\widehat{X}_{s}\right) d \widehat{L}_{s} \\
& =f\left(\widehat{X}_{s}\right)+\int_{0}^{t} f^{\prime}\left(\widehat{X}_{s}\right) d \widehat{B}_{s}+\int_{0}^{t} \mathbf{v}_{\theta}\left(f\left(\widehat{X}_{s}\right)\right) d \widehat{L}_{s} \tag{4.6}
\end{align*}
$$

where $f^{\prime}$ is interpreted as the Jacobian of $f$. Let

$$
\begin{equation*}
c(t)=\int_{0}^{t}\left|f^{\prime}\left(\widehat{X}_{s}\right)\right|^{2} d s \tag{4.7}
\end{equation*}
$$

It is not hard to show that $c(t)<\infty$ for every $t>0$, a.s. It follows that

$$
c^{-1}(t):=\inf \{s>0: c(s)>t\}
$$

is well defined for every $t>0$ and the process $X_{t}:=f\left(\widehat{X}_{c^{-1}(t)}\right)$ satisfies (2.1) with Brownian motion $B_{t}:=\int_{0}^{c^{-1}(t)} f^{\prime}\left(\widehat{X}_{s}\right) d \widehat{B}_{s}$ and $L:=\widehat{L}$. So $X$ is an ORBM in $D_{*}$ with the oblique angle of reflection $\theta$. The exponential function $f(z)=\exp (z)$ : $D_{-} \rightarrow D_{*}$ is neither one-to-one nor onto $D_{*}$, but it is locally conformal and maps $\partial D_{-}$onto $\partial D_{*}$ so we will refer to the fact that $X_{t}$ is an ORBM as conformal invariance of ORBM.

Let $\sigma_{t}=\inf \left\{s \geq 0: L_{s}>t\right\}=\widehat{\sigma}_{t}=\inf \left\{s \geq 0: \widehat{L}_{s}>t\right\}, A_{t}=\arg X_{\sigma_{t}}$ and $\widehat{A}_{t}=$ $\operatorname{Im} \widehat{X}_{\widehat{\sigma}_{t}}$ for $t \geq 0$. Then $\widehat{A}$ and $A$ are indistinguishable processes.

It follows from the uniqueness of the deterministic Skorohod problem that the process $\bar{X}_{t}:=\widehat{X}_{t}-i \int_{0}^{t} \tan \theta\left(\widehat{X}_{s}\right) d \widehat{L}_{s}$ is a normally reflected Brownian motion in the left half-plane $D_{-}$. Hence, if we let $C_{t}=\operatorname{Im} \bar{X}\left(\widehat{\sigma}_{t}\right)$ for $t \geq 0$, then $C_{t}$ is a Cauchy process with the initial value $C_{0}=\operatorname{Im} \bar{X}_{\bar{S}}=\arg X_{S}$, where $\bar{S}:=\inf \{t>0$ : $\left.\bar{X}_{t} \in \partial D_{-}\right\}$and $S:=\inf \left\{t>0: X_{t} \in \partial D_{*}\right\}$. Clearly, $C_{0}$ depends only on the initial starting point of $X$ and is independent of the reflection angle $\theta$. We have

$$
\begin{equation*}
A_{t}=\widehat{A}_{t}=C_{t}+\int_{0}^{\widehat{\sigma}_{t}} \tan \left(\theta\left(\widehat{X}_{s}\right)\right) d \widehat{L}_{s}=C_{t}+\int_{0}^{\sigma_{t}} \tan \left(\theta\left(X_{s}\right)\right) d L_{s} \tag{4.8}
\end{equation*}
$$

For $u \geq 0$, define

$$
\begin{equation*}
T_{u}=\inf \left\{t>u: X_{t} \in \partial D_{*}\right\} \tag{4.9}
\end{equation*}
$$

with the convention $\inf \varnothing:=\infty$. We obtain from (4.8)

$$
\begin{aligned}
\arg X_{t} & =A_{L_{t}}+\arg X_{t}-\arg X_{T_{t}} \\
& =C_{L_{t}}+\int_{0}^{t} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}+\arg X_{t}-\arg X_{T_{t}} \\
& =C_{t}+\left(C_{L_{t}}-C_{t}\right)+\int_{0}^{t} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}+\left(\arg X_{t}-\arg X_{T_{t}}\right) .
\end{aligned}
$$

Hence,

$$
\begin{align*}
\frac{1}{t} \arg X_{t}-\mu_{0}= & \frac{1}{t} C_{t}+\frac{1}{t}\left(C_{L_{t}}-C_{t}\right)+\left(\frac{1}{t} \int_{0}^{t} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}-\mu_{0}\right)  \tag{4.10}\\
& +\frac{1}{t}\left(\arg X_{t}-\arg X_{T_{t}}\right)
\end{align*}
$$

By (4.2), $\mathbb{E}_{Q}\left[L_{1}\right]=\int_{\partial D_{*}}(h(x) / 2) d x=1$. It follows from these remarks, (2.31), (4.2) with $g(x)=\tan \theta(x)$ and the limit-quotient theorem for additive functionals [see, e.g., Revuz and Yor (1999), Theorem X 3.12] that for every $z \in \bar{D}_{*}, \mathbb{P}_{z}$-a.s.,

$$
\begin{align*}
\lim _{t \rightarrow \infty} \frac{1}{t} \int_{0}^{t} \tan \left(\theta\left(X_{S}\right)\right) d L_{s} & =\mathbb{E}_{Q}\left[\int_{0}^{1} \tan \theta\left(X_{s}\right) d L_{s}\right]  \tag{4.11}\\
& =\int_{\partial D_{*}}(1 / 2) \tan \theta(x) h(x) d x=\mu_{0} \\
\lim _{t \rightarrow \infty} \frac{1}{t} L_{t} & =1 \tag{4.12}
\end{align*}
$$

Fix an arbitrarily small $\varepsilon>0$ and any $z \in \bar{D}_{*}$ and let

$$
\begin{equation*}
p_{1}(t)=\mathbb{P}_{z}\left(\left|\arg X_{t}-\arg X_{T_{t}}\right|>\varepsilon t\right) \tag{4.13}
\end{equation*}
$$

We will argue that $p_{1}(t)$ is small for large $t$. Let $T_{u}^{\prime}=\sup \left\{t \in[0, u]: X_{t} \in \partial D_{*}\right\}$ with the convention $\sup \varnothing=0$. By the Markov property applied at time $t$ and the symmetry of Brownian motion,

$$
\mathbb{P}_{z}\left(\arg X_{T_{t}^{\prime}}-\arg X_{t}>0\right)=\mathbb{P}_{z}\left(\arg X_{T_{t}^{\prime}}-\arg X_{t}<0\right)=1 / 2
$$

This and the Markov property applied at time $t$ imply that

$$
\begin{equation*}
\mathbb{P}_{z}\left(\left|\arg X_{T_{t}^{\prime}}-\arg X_{T_{t}}\right|>\varepsilon t\right) \geq p_{1}(t) / 2 \tag{4.14}
\end{equation*}
$$

For a fixed $u>0$, the Cauchy process $C$ is continuous at time $u$, a.s. Let $\delta>0$ be so small that

$$
\mathbb{P}\left(\sup _{1-\delta \leq u, v \leq 1+\delta}\left|C_{u}-C_{v}\right| \geq \varepsilon / 2\right)<\varepsilon .
$$

Then, by scaling, for any $t>0$,

$$
\begin{equation*}
\mathbb{P}\left(\sup _{(1-\delta) t \leq u, v \leq(1+\delta) t}\left|C_{u}-C_{v}\right| \geq \varepsilon t / 2\right)<\varepsilon . \tag{4.15}
\end{equation*}
$$

By (4.12), we can find $t_{1}$ so large that for $t \geq t_{1}$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(L_{t} \in((1-\delta) t,(1+\delta) t)\right) \geq 1-\varepsilon \tag{4.16}
\end{equation*}
$$

The jumps of $A$ have the same size as those of $C$ and occur at the same time because the last integral in (4.8) is a continuous function of $t$. If the events in (4.14) and (4.16) occur then $C$ has a jump of size greater than $\varepsilon t$ at a time $s=$ $L_{t} \in((1-\delta) t,(1+\delta) t)$. The probability of this event is greater than $p_{1}(t) / 2-\varepsilon$, by (4.14) and (4.16). However, by (4.15), this probability is less than $\varepsilon$. Hence, $p_{1}(t) / 2<2 \varepsilon$ and, therefore, $p_{1}(t)<4 \varepsilon$ for sufficiently large $t$. This and (4.13) imply that, for sufficiently large $t$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\frac{1}{t}\left|\arg X_{t}-\arg X_{T_{t}}\right|>\varepsilon\right)<4 \varepsilon . \tag{4.17}
\end{equation*}
$$

Another consequence of (4.15) and (4.16) is that $\left|C_{L_{t}}-C_{t}\right| \leq \varepsilon t$ with probability greater than $1-2 \varepsilon$ for large $t$. Thus, for sufficiently large $t$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\frac{1}{t}\left|C_{L_{t}}-C_{t}\right|>\varepsilon\right)<2 \varepsilon \tag{4.18}
\end{equation*}
$$

It follows from (4.11) that for sufficiently large $t$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\left|\frac{1}{t} \int_{0}^{t} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}-\mu_{0}\right|>\varepsilon\right)<\varepsilon . \tag{4.19}
\end{equation*}
$$

Note that $\left(C_{t}-C_{0}\right) / t$ has the Cauchy distribution. Since $\varepsilon>0$ is arbitrarily small, the last observation, (4.10), (4.17), (4.18) and (4.19) imply that the distributions of $\frac{1}{t} \arg X_{t}-\mu_{0}$ converge to the Cauchy distribution as $t \rightarrow \infty$.
(iii) Consider a modification of the process $C$ which is left continuous with right limits. For $t \geq 0$, let
$\Lambda_{t}=\sum_{s \leq t}\left(C_{t+}-C_{t}\right) \mathbf{1}_{\left\{\left|C_{t+}-C_{t}\right|>2 \pi\right\}}, \quad C_{t}^{*}=C_{t}-C_{0}-\Lambda_{t}=C_{t}-\arg X_{S}-\Lambda_{t}$.
The process $C^{*}$ is a Cauchy process with jumps larger than $2 \pi$ removed and starts from $C_{0}^{*}=0$. It is elementary to see that $C_{t}^{*}$ is a zero mean martingale and a Lévy process. Hence, the law of large numbers holds for $C^{*}$, that is, a.s.,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} C_{t}^{*} / t=0 \tag{4.20}
\end{equation*}
$$

Note that the jumps removed from $C$ correspond to increments of $\arg X$ in the sum on the right-hand side of (3.1). Thus,

$$
\begin{equation*}
\arg ^{*} X_{\sigma(t)}=C_{t}^{*}+\int_{0}^{\sigma(t)} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}+\arg X_{S} \tag{4.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{t} \arg ^{*} X_{\sigma(t)}=\frac{1}{t} C_{t}^{*}+\frac{1}{t} \int_{0}^{\sigma(t)} \tan \left(\theta\left(X_{s}\right)\right) d L_{s}+\frac{1}{t} \arg X_{S} \tag{4.22}
\end{equation*}
$$

It follows from (4.12) that, a.s.,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sigma(t) / t=1 \tag{4.23}
\end{equation*}
$$

This, (4.11), (4.20) and (4.22) imply that for every $z \in \bar{D}_{*}, \mathbb{P}_{z}$-a.s.,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{\sigma(t)}=\mu_{0} \tag{4.24}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} T_{t} / t=1 \quad \text { a.s. } \tag{4.25}
\end{equation*}
$$

First note that since $\int_{0}^{\infty} 1_{\left\{X_{s} \in D_{*}\right\}} d L_{s}=0$, we have by (4.12) that $\lim _{t \rightarrow \infty} T_{t}=\infty$. For every $\varepsilon>0, L_{t}-\varepsilon<L_{T_{t}} \leq L_{t}$ so

$$
\frac{L_{t}-\varepsilon}{t}<\frac{L_{T_{t}}}{T_{t}} \frac{T_{t}}{t} \leq \frac{L_{t}}{t} .
$$

This together with (4.12) establishes the claim (4.25). Combining (4.23), (4.24) and (4.25) yields

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{T_{t}}=\mu_{0} \tag{4.26}
\end{equation*}
$$

Next we will argue that (4.26) implies that $\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{t}=\mu_{0}$ by using excursion theory. Recall that $H^{x}$ denotes the excursion law for Brownian motion in $D_{*}$. We will estimate the $H^{x}$-measure of the family $F_{a}$ of excursions with the property that $|\arg \mathrm{e}(0)-\arg \mathrm{e}(\zeta-)| \leq 2 \pi$ and $\sup _{t \in[0, \zeta(\mathrm{e}))}|\arg \mathrm{e}(0)-\arg \mathrm{e}(t)| \geq a$, for $a \geq 4 \pi$. Note that this quantity does not depend on $x$. Let $\widehat{H}^{x}$ be the excursion law for Brownian motion in $D_{-}=\{z \in \mathbb{C}: \operatorname{Re} z<0\}$ starting from $x \in \partial D_{-}$. Excursion laws are conformally invariant, up to a multiplicative constant [see Burdzy (1987), Proposition 10.1]. The exponential function $f(z)=\exp (z)$ maps $D_{-}$onto $D_{*} \backslash\{0\}$ and is locally conformal, up to the boundary. Hence, for some constant $c_{4}$, $H^{x}\left(F_{a}\right)=c_{4} \widehat{H}^{y}\left(\widehat{F}_{a}\right)$, where $\widehat{F}_{a}$ is the family of excursions with the property that $|\operatorname{Ime}(0)-\operatorname{Ime}(\zeta-)| \leq 2 \pi$ and $\sup _{t \in[0, \zeta(\mathrm{e}))}|\operatorname{Ime}(0)-\operatorname{Ime}(t)| \geq a$. If we normalize all excursion laws as in (2.5), then it is easy to check that $c_{4}=1$ (although our argument does not depend on the value of this constant). Thus, the equality $H^{x}\left(F_{a}\right)=\widehat{H}^{y}\left(\widehat{F}_{a}\right)$ holds for all $x \in \partial D_{*}$ and $y \in \partial D_{-}$. By Burdzy [(1987), Theorem 5.1(v)], for some $c_{5}<\infty$,

$$
\begin{equation*}
\widehat{H}^{x}\left(\sup _{t \in[0, \zeta(\mathrm{e}))}|\operatorname{Im} \mathrm{e}(0)-\operatorname{Im} \mathrm{e}(t)| \geq a\right) \leq c_{5} / a . \tag{4.27}
\end{equation*}
$$

It is easy to see that if Brownian motion starts in $D_{-}$from a point $z$ with $|\operatorname{Im} z|>a$ with $a \geq 4 \pi$ then the chance that it will exit $D_{-}$through the line segment on the imaginary axis between $-2 \pi i$ and $2 \pi i$ is bounded above by $c_{6} / a$. This, (4.27) and
the strong Markov property of $\widehat{H}^{x}$ applied at the time $\inf \{t \in[0, \zeta(\mathrm{e})): \mid \operatorname{Ime}(0)-$ $\operatorname{Im} \mathrm{e}(t) \mid \geq a\}$ imply that

$$
\begin{equation*}
H^{x}\left(F_{a}\right)=\widehat{H}^{x}\left(\widehat{F}_{a}\right) \leq c_{5} c_{6} / a^{2}=c_{7} / a^{2} \tag{4.28}
\end{equation*}
$$

Fix some $\alpha>0$. By the exit system formula (2.4), the probability that there exists an excursion $\mathrm{e}_{t}$ of $X$ such that $L_{t}>s$ and $\mathrm{e}_{t}$ belongs to $F_{\alpha L_{t}}$ is equal to

$$
\int_{s}^{\infty} H^{X_{\sigma(u)}}\left(F_{\alpha u}\right) d u \leq \int_{s}^{\infty} c_{7} /(\alpha u)^{2} d u=c_{8} /\left(\alpha^{2} s\right)
$$

This quantity goes to 0 as $s \rightarrow \infty$, so for every fixed $\alpha>0$, with probability 1 , there is $s_{\alpha}=s_{\alpha}(\omega)<\infty$ such that there are no excursions $\mathrm{e}_{t} \in F_{\alpha L_{t}}$ with $L_{t}>s_{\alpha}$.

Fix an arbitrarily small $\alpha>0$ and suppose that $t_{1}$ is so large that $\frac{1}{t} \arg ^{*} X_{T_{t}} \leq$ $\mu_{0}+\alpha$ and $L_{t} / t \leq 2$ for all $t>t_{1}$. If $\frac{1}{u}$ arg* $^{*} X_{u} \geq \mu_{0}+5 \alpha$ for some $u>t_{1}$, then $\left|\arg ^{*} X_{u}-\arg ^{*} X_{T_{u}}\right| \geq 4 \alpha u \geq 2 \alpha L_{u}$. This means that an excursion starting at $T_{u}$ belongs to $F_{2 \alpha L_{u}}=F_{2 \alpha L_{T_{u}}}$. Since there are no such excursions beyond some $s_{2 \alpha}$, it follows that limsup $\operatorname{sim}_{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{T_{t}} \leq \mu_{0}+5 \alpha$, a.s. This holds for all rational $\alpha>$ 0 simultaneously, a.s., so $\lim \sup _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{T_{t}} \leq \mu_{0}$, a.s. The matching lower bound for liminf can be proved analogously. We conclude that for every $z \in D_{*}$, $\mathbb{P}_{z}$-a.s., $\lim _{t \rightarrow \infty} \frac{1}{t} \arg ^{*} X_{T_{t}}=\mu_{0}$.

The proof of (3.4) will be combined with the proof of Theorem 3.15(iv) given below.
(iv) Since $h$ is $C^{2}$ on $\bar{D}_{*}$, it follows from (2.18) that $\theta(z)$ is $C^{2}$ on $\bar{D}_{*}$, and hence $\theta(x)$ is $C^{2}$ on $\partial D_{*}$. Moreover $H=h+i \widetilde{h}$ is $C^{2-\varepsilon}$, by Corollary II.3.3 in Garnett and Marshall (2005). By assumption, $h$ is positive and continuous on $\partial D_{*}$. Thus, $H\left(\bar{D}_{*}\right)$ is a compact subset of $\{\operatorname{Re} z>0\}$ and so by (2.18), $\sup _{x}|\theta(x)|<\pi / 2$. We can now apply parts (i) and (iii) of the theorem to see that part (iv) holds.

Proof of Theorem 3.5. Fix a Borel measurable function $\theta: \partial D_{*} \rightarrow$ $[-\pi / 2, \pi / 2]$. First we need to prove that there exists a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ which converges to $\theta$ in weak-* topology. For this, we extend $\theta$ harmonically to $\bar{D}_{*}$ and then we let $\theta_{k}\left(e^{i t}\right)=\theta\left(e^{i t}(1-1 / k)\right)$. Then $\theta_{k}$ 's converge to $\theta$ in weak-* topology. See Hoffman (1962), page 33.
(i) This was essentially proved in Burdzy and Marshall (1993), Theorem 1.1. That theorem was concerned with ORBM in a half-plane while the present result is set in a disc. Theorem 3.5(i) can be proved just like Burdzy and Marshall (1993), Theorem 1.1, by repeating the arguments given in Burdzy and Marshall (1993) with some minor adjustments. We omit the proof to save space. The Markov property of $X$ follows from that of $X^{k}$ and the convergence of finite dimensional distributions. Since for each $k$, the subprocess of $X^{k}$ before hitting $\partial D_{*}$ is Brownian motion in $D_{*}$ before hitting $\partial D_{*}$, the same claim applies to the subprocess of $X$ before hitting $\partial D_{*}$.

The transition probabilities are the same for each process $\left|X^{k}\right|$ so the process $|X|$ has the same transition probabilities. It follows that $X$ is conservative.
(ii) This claim was shown in the proof of Burdzy and Marshall (1993), Theorem 1.1, although it was not a part of the statement of that theorem. See Step 4 on page 214 of Burdzy and Marshall (1993).
(iii) Suppose that $\left(h_{k}, \mu_{k}\right) \leftrightarrow \theta_{k}$ and $X^{k}$ solves the SDE (3.10) except that the initial distribution for $X^{k}$ is the stationary distribution $h_{k}(z) d z$. According to Remark 2.6, the measures $h_{k}(z) d z$ converge to $h(z) d z$. It is easy to see that part (i) of this theorem implies that $X^{k}$ 's converge weakly to a process $X$ satisfying the SDE (2.1), with the initial distribution $h(z) d z$. For every $t \geq 0$ and $k \geq 1$, the distribution of $X_{t}^{k}$ is $h_{k}(z) d z$. Hence, for every $t \geq 0$, the distribution of $X_{t}$ is $h(z) d z$. This shows that $h$ is a stationary distribution for $X$ satisfying (2.1).

We next show uniqueness of the stationary distribution. As observed in (2.2), for every reflection angle field $\theta$, the radial part $|X|$ of $X$ is a two-dimensional Bessel process confined to $[0,1]$ by reflection at 1 . This easily implies that for any initial distribution of $X$, the distribution of $X_{1}$ has a strictly positive density inside $\mathcal{B}(0,1 / 2)$. If there were more than one invariant measure, at least two of them (say, $Q_{1}$ and $Q_{2}$ ) would be mutually singular by Birkhoff's ergodic theorem Sinaĭ (1994). We have just shown that the Lebesgue measure restricted to $\mathcal{B}(0,1 / 2)$ (let us call it $Q_{3}$ ) is absolutely continuous with respect to the distribution of $X_{1}$, so that in particular, $Q_{3} \ll Q_{1}$ and $Q_{3} \ll Q_{2}$. Since $Q_{1} \perp Q_{2}$ by assumption, there exists a set $A \subset \mathcal{B}(0,1 / 2)$ such that $Q_{1}(A)=0$ and $Q_{2}(\mathcal{B}(0,1 / 2) \backslash A)=0$. Therefore, one must have $Q_{3}(A)=Q_{3}(\mathcal{B}(0,1 / 2) \backslash A)=0$ which contradicts the fact that $Q_{3}(\mathcal{B}(0,1 / 2)) \neq 0$.
(iv) The first claim follows easily from the definitions. The second claim follows from the first claim and part (i) of the theorem.
(v) Since $\theta_{k}$ are smooth, (4.10) holds for $X^{k}$ 's, that is,

$$
\begin{align*}
& \frac{1}{t} \arg X_{t}^{k}-\mu_{k} \\
& \quad=\frac{1}{t} C_{t}^{k}+\frac{1}{t}\left(C_{L_{t}^{k}}^{k}-C_{t}^{k}\right)+\left(\frac{1}{t} \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}-\mu_{k}\right)  \tag{4.29}\\
& \quad+\frac{1}{t}\left(\arg X_{t}^{k}-\arg X_{T_{t}^{k}}^{k}\right)
\end{align*}
$$

where the symbols with the superscript or subscript $k$ denote objects analogous to those in (4.10). Since $X^{k}$,s converge to $X$ weakly, we can assume that all these processes are constructed on a single probability space and $X_{t}^{k} \rightarrow X_{t}$, a.s., for every fixed $t$, as $k \rightarrow \infty$. In view of (4.29), we can write

$$
\begin{align*}
\frac{1}{t} \arg X_{t}-\mu_{0}= & \left(\frac{1}{t} \arg X_{t}-\frac{1}{t} \arg X_{t}^{k}\right) \\
& -\left(\mu_{0}-\mu_{k}\right)+\frac{1}{t}\left(C_{t}^{k}-C_{0}^{k}\right)+\frac{1}{t}\left(C_{L_{t}^{k}}^{k}-C_{t}^{k}\right) \tag{4.30}
\end{align*}
$$

$$
\begin{aligned}
& +\left(\frac{1}{t} \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}-\mu_{k}\right) \\
& +\frac{1}{t}\left(\arg X_{t}^{k}-\arg X_{T_{t}^{k}}^{k}\right)+\frac{1}{t} \arg X_{S^{k}}^{k}
\end{aligned}
$$

where $S^{k}=\inf \left\{t>0: X_{t}^{k} \in \partial D_{*}\right\}$. The distribution of $\frac{1}{t}\left(C_{t}^{k}-C_{0}^{k}\right)$ is Cauchy for every $k$ and $t$ so it suffices to show that all other terms on the right-hand side of (4.30) are small for large $t$ and $k$.

Fix an arbitrarily small $\varepsilon>0$. Note that (4.17) and (4.18) do not depend on $\theta$ so we can apply them for all $\theta_{k}$. Hence, we can find $t_{1}$ so large that for $t \geq t_{1}$,

$$
\mathbb{P}\left(\left|\frac{1}{t}\left(C_{L_{t}^{k}}^{k}-C_{t}^{k}\right)+\frac{1}{t}\left(\arg X_{t}^{k}-\arg X_{T_{t}^{k}}^{k}\right)\right| \geq \varepsilon\right)<\varepsilon .
$$

We will assume without loss of generality that $X_{0}^{k}=z \neq 0$, a.s., for all $k$. (The case $z=0$ can be dealt with by applying the Markov property at time $t=1$.) Then $\arg X_{S^{k}}^{k}$ has the same distribution for each $k \geq 1$ and so by taking $t_{1}$ larger if needed,

$$
\mathbb{P}\left(\left|\frac{1}{t} \arg X_{S^{k}}^{k}\right| \geq \varepsilon\right)<\varepsilon \quad \text { for all } k \geq 1 \text { and } t \geq t_{1}
$$

Recall that $X_{t}^{k} \rightarrow X_{t}$, a.s. By Remark $2.6(\mathrm{vi}), \mu_{k} \rightarrow \mu_{0}$. Thus, for a fixed $t$, we can make $k$ so large that

$$
\mathbb{P}\left(\left|\frac{1}{t} \arg X_{t}-\frac{1}{t} \arg X_{t}^{k}\right|+\left|\mu_{0}-\mu_{k}\right| \geq \varepsilon\right)<\varepsilon
$$

Hence, it will suffice to prove that for a fixed $\varepsilon>0$, some $t_{1}$ and $k_{1}$, all $t \geq t_{1}$, $k \geq k_{1}$ and $z_{k} \in \bar{D}_{*}$,

$$
\begin{equation*}
\mathbb{P}_{z_{k}}\left(\left|\frac{1}{t} \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}-\mu_{k}\right|>\varepsilon\right)<\varepsilon \tag{4.31}
\end{equation*}
$$

If we let $Q_{k}(d x)=h_{k}(x) d x$ then by (4.11),

$$
\begin{equation*}
\mathbb{E}_{Q_{k}}\left[\frac{1}{t} \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right]=\mu_{k} \tag{4.32}
\end{equation*}
$$

Hence, to complete the proof of part (iv) of the theorem, it will suffice to show that

$$
\begin{equation*}
\operatorname{Var}\left(\frac{1}{t} \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right) \leq c_{1} / t \tag{4.33}
\end{equation*}
$$

We will split the rest of the proof of this part of the theorem into steps.
Step 1. We will recall some results from Burdzy and Marshall (1993), Lemmas $2.2-2.3$, but we will change the notation.

We will say that $D \subset \mathbb{C}$ is a monotone domain if $D$ is open, connected and for every $z \in D$ and $b>0$ we have $z+i b \in D$.

Let $\mathbb{H}=\{z \in \mathbb{C}: \operatorname{Im} z>0\}$ be the upper half-plane. Suppose that $\theta: \partial \mathbb{H} \rightarrow$ $[-\pi / 2, \pi / 2]$ is a Borel measurable function and suppose $\theta$ is not equal almost everywhere either to $\pi / 2$ or to $-\pi / 2$. Then there exists a univalent analytic mapping $g$ of $\mathbb{H}$ onto a monotone domain $D=g(\mathbb{H})$ such that for almost all $x \in \partial \mathbb{H}, g(x)$ and $g^{\prime}(x)$ exist, $g^{\prime}(x) \neq 0$ and $\arg g^{\prime}(x)=\theta(x)$. We choose $g$ so that $\lim _{|z| \rightarrow \infty}|g(z)|=\infty$. We construct $g$ as follows. Let $\theta: \mathbb{H} \rightarrow \mathbb{R}$ be the bounded harmonic extension of our original function $\theta: \partial \mathbb{H} \rightarrow[-\pi / 2, \pi / 2]$ and let $\widetilde{\theta}$ be the harmonic conjugate of $\theta$ such that $\tilde{\theta}(i)=0$. Define $g: \mathbb{H} \rightarrow \mathbb{C}$ by setting $g(i)=i$ and

$$
g^{\prime}(z)=\exp (i(\theta(z)+i \widetilde{\theta}(z)))
$$

Then $g$ is one-to-one on $\mathbb{H}$ because $\operatorname{Re} g^{\prime}(z)>0$. [See Burdzy and Marshall (1993).] By abuse of notation, we will use the same symbol $\theta$ to denote real functions on both $\partial D_{*}$ and $\partial \mathbb{H}$. Specifically, for $z \in \partial \mathbb{H}$, we let $\theta(z)=\theta(\exp (i z))$, where $\theta(\exp (i z))$ refers to the function $\theta \in \mathcal{T}$ introduced in the assumptions of Theorem 3.5. Hence, in this proof, $\theta: \partial \mathbb{H} \rightarrow \mathbb{R}$ is a periodic function with period $2 \pi$. It follows that $g$ is also periodic with period $2 \pi$, up to an additive constant. That is, $g(z+2 \pi)=g(z)+d$ for all $z \in \mathbb{H}$, where $d=g(i+2 \pi)-g(i)$. The constant $d$ is nonzero since $\operatorname{Re} g^{\prime}>0$.

Suppose that $\theta_{k}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ are $C^{2}$-functions which converge weak-*to $\theta$ as $k \rightarrow \infty$. Let $g_{k}$ and $D_{k}:=g_{k}(\mathbb{H})$ correspond to $\theta_{k}$ in the same way as $g$ and $D=g(\mathbb{H})$ correspond to $\theta$. Note that $g_{k}(z+2 \pi)=g_{k}(z)+d_{k}$ for some constant $d_{k}$. Moreover, if $\varepsilon>0$, then $g_{k}(z+i \varepsilon)$ converges to $g(z+i \varepsilon)$ uniformly in $z \in \mathbb{R}$ and $d_{k} \rightarrow d$. Indeed, by weak-* convergence of $\theta_{k} \in \mathcal{T}$, we conclude uniform convergence of $\theta_{k}(z)+i \widetilde{\theta}(z)$ on the compact set $\left\{z:|z|=e^{-\varepsilon}\right\}$, and hence $g_{k}^{\prime}$ converges uniformly on $I=\{z: 0 \leq \operatorname{Re} z \leq 2 \pi, \operatorname{Im} z=\varepsilon\}$. Integration then shows that $d_{k} \rightarrow d$, and hence $g_{k}$ converges uniformly to $g$ on $\mathbb{R}+i \varepsilon$. Let $f(z)=\exp \left(i g^{-1}(z)\right)$, for $z \in D$ and $f_{k}(z)=\exp \left(i g_{k}^{-1}(z)\right)$ for $z \in D_{k}$. Then $f$ and $f_{k}$ are locally conformal maps of $D$ and $D_{k}$ onto $D_{*} \backslash\{0\}$ which are periodic with periods $d$ and $d_{k}$, respectively.

The monotone domains $D_{k}$ converge to $D$ in the following sense:
(a) If $B$ is open and such that $B \cap \partial D \neq \varnothing$, there is a $k_{0}=k_{0}(B)$ such that $B \cap \partial D_{k} \neq \varnothing$ for all $k \geq k_{0}$.
(b) If $B$ is connected and open, with $B \cap D \neq \varnothing$ and $B \subset D_{k}$ for infinitely many $k$, then $B \subset D$.
(c) If $K$ is compact and $K \subset D$, then $K \subset D_{k}$ for all $k \geq k_{0}=k_{0}(K)$.

We invoke conformal invariance of ORBM as in (4.5)-(4.7). For $x \in \partial D_{k}$ such that $f_{k}(x)=z \in \partial D_{*}$, let $\widehat{\mathbf{v}}_{k}(x)=i \sec \theta_{k}(z)$. In other words, $\widehat{\mathbf{v}}_{k}$ is the conformal (inverse) image of the vector of reflection $\mathbf{v}_{\theta_{k}}$. Suppose that $\widehat{B}$ is a two-dimensional Brownian motion and consider the Skorokhod SDE

$$
\begin{equation*}
\widehat{X}_{t}^{k}=\widehat{x}_{k}+\widehat{B}_{t}+\int_{0}^{t} \widehat{\mathbf{v}}_{k}\left(\widehat{X}_{s}^{k}\right) d \widehat{L}_{s}^{k} \tag{4.34}
\end{equation*}
$$

where $\widehat{L}^{k}$ is the local time of $\widehat{X}^{k}$ on $\partial D_{k}$. The process $\widehat{X}^{k}$ is reflected Brownian motion in $D_{k}$ with the oblique angle of reflection $\theta_{k}$. If $c_{k}(t)=\int_{0}^{t}\left|f_{k}^{\prime}\left(\widehat{X}_{s}^{k}\right)\right|^{2} d s$ then the process $X_{t}^{k}=f_{k}\left(\widehat{X}_{c_{k}(t)}^{k}\right)$ is reflected Brownian motion in $D_{*}$ with the oblique angle of reflection $\theta_{k}$.

Let $K_{k}=f_{k}^{-1}(\partial \mathcal{B}(0,1 / 2))$. Note that $K_{k}$ is the image under the map $g_{k}$ of the horizontal line $\{z: \operatorname{Im} z=\ln 2\}$, and so $K_{k}$ is an analytic curve. Let $a_{k}=\operatorname{Re} d_{k}=$ $\operatorname{Re}\left(g_{k}(2 \pi)-g_{k}(0)\right)$ and for $z \in \partial D_{k}$, let $R_{k}(z)=\left\{x \in \partial D_{k}:|\operatorname{Re} x-\operatorname{Re} z| \geq a_{k}\right\}$. Let $\widehat{T}^{k}(A)=\inf \left\{t \geq 0: \widehat{X}_{t}^{k} \in A\right\}$. We will show that for every $\theta$ there exists $p_{1}>0$ such that for every approximating sequence $\left\{\theta_{k}\right\}$ there exists $k_{1}$ such that for any $k \geq k_{1}$ and $z_{k} \in \partial D_{k}$,

$$
\begin{equation*}
\mathbb{P}_{z_{k}}\left(\widehat{T}^{k}\left(K_{k}\right)<\widehat{T}^{k}\left(R_{k}\right)\right) \geq p_{1} \tag{4.35}
\end{equation*}
$$

Let $[x, z]$ denote the line segment between $x, z \in \mathbb{C}$. For every $\theta$ there exist $a, b \in(0, \infty)$ such that for every approximating sequence $\left\{\theta_{k}\right\}$ there exists $k_{1}$ such that for any $k \geq k_{1}$ and $z \in \partial D_{k}$ we have $a_{k} \geq a$ and $K_{k} \cap[z, z+i b] \neq \varnothing$.

With probability greater than $p_{2}>0$, Brownian motion starting from 0 will hit the line $\{z: \operatorname{Im} z=2 b\}$ before hitting the lines $\{z:|\operatorname{Re} z|=a / 2\}$, and then it will cross the imaginary axis before hitting any of the lines $\{z:|\operatorname{Re} z|=a\}$ or $\{z: \operatorname{Im} z=b\}$. Since $\int_{0}^{t} \widehat{\mathbf{v}}_{k}\left(\widehat{X}_{s}^{k}\right) d \widehat{L}_{s}^{k}$ is a purely imaginary number with nonnegative imaginary part, this implies that with probability greater than $p_{2}$, the process $\widehat{X}^{k}$ starting from $z_{k} \in \partial D_{k}$ will hit the line $\left\{z: \operatorname{Im} z-\operatorname{Im} z_{k}=2 b\right\}$ before hitting the lines $\left\{z:\left|\operatorname{Re} z-\operatorname{Re} z_{k}\right|=a / 2\right\}$, and then it will cross the line $\left\{z: \operatorname{Re} z=\operatorname{Re} z_{k}\right\}$ before hitting any of the lines $\left\{z:\left|\operatorname{Re} z-\operatorname{Re} z_{k}\right|=a\right\}$ or $\left\{z: \operatorname{Im} z-\operatorname{Im} z_{k}=b\right\}$. If the trajectory of $\widehat{X}^{k}$ follows a path described above then, in view of the definitions of $a$ and $b$, it will cross $K_{k}$ before hitting $R_{k}$. We conclude that (4.35) holds with $p_{1}=p_{2}>0$.

Let

$$
\begin{aligned}
T^{k}(A) & =\inf \left\{t \geq 0: X_{t}^{k} \in A\right\} \\
T_{b}^{k} & =T^{k}(\mathcal{B}(0,1 / 2)), \\
T_{*}^{k} & =\inf \left\{t \geq 0: X_{t}^{k} \in \partial D_{*},\left|\arg X_{t}^{k}-\arg X_{0}^{k}\right| \geq 2 \pi\right\} .
\end{aligned}
$$

By the conformal invariance of ORBM, (4.35) implies that

$$
\begin{equation*}
\mathbb{P}_{z_{k}}\left(T_{b}^{k}<T_{*}^{k}\right) \geq p_{1} \quad \text { for all } k \text { and } z_{k} \in \partial D_{*} . \tag{4.36}
\end{equation*}
$$

Step 2. We will estimate the variance of $\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}$.
Let $S_{1}^{k}=T^{k}\left(\partial D_{*} \cup \partial \mathcal{B}(0,1 / 2)\right)$. The probability that Brownian motion will make a loop in the annulus $D_{*} \backslash \mathcal{B}(0,1 / 2)$ (that is, $\arg X^{k}$ will increase or decrease by $2 \pi$ ) before hitting the boundary of the annulus is less than $p_{3}<1$. This implies that, for any $z \in \bar{D}_{*}$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\left|\arg X_{S_{1}^{k}}^{k}-\arg X_{0}^{k}\right| \leq 2 \pi\right) \geq 1-p_{3} . \tag{4.37}
\end{equation*}
$$

This and an easy inductive argument based on the strong Markov property applied at the times when consecutive loops are completed shows that there exists $n$ so large that for any $z \in \bar{D}_{*}$,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\left|\arg X_{S_{1}^{k}}^{k}-\arg X_{0}^{k}\right| \geq n 2 \pi\right) \leq p_{1} / 4 \tag{4.38}
\end{equation*}
$$

where $p_{1}$ is as in (4.36). Fix such an $n$ and let

$$
\begin{aligned}
S_{2}^{k} & =\inf \left\{t \geq 0:\left|\arg X_{t}^{k}-\arg X_{0}^{k}\right| \geq(n+1) 2 \pi\right\} \\
S_{3}^{k} & =\inf \left\{t \geq S_{2}^{k}:\left|\arg X_{t}^{k}-\arg X_{S_{k}^{2}}^{k}\right| \geq n 2 \pi\right\}, \\
S_{4}^{k} & =\inf \left\{t \geq 0:\left|\arg X_{t}^{k}-\arg X_{0}^{k}\right| \geq(2 n+1) 2 \pi\right\}, \\
S_{5, j}^{k} & =\inf \left\{t \geq 0:\left|\arg X_{t}^{k}-\arg X_{0}^{k}\right| \geq j(2 n+2) 2 \pi\right\} .
\end{aligned}
$$

By (4.36), we have for $z \in \partial D_{*}$,

$$
\mathbb{P}_{z}\left(T_{b}^{k} \leq T_{*}^{k} \wedge S_{2}^{k}\right)+\mathbb{P}_{z}\left(S_{2}^{k} \leq T_{b}^{k} \leq T_{*}^{k}\right) \geq p_{1}
$$

It follows that either

$$
\begin{equation*}
\mathbb{P}_{z}\left(T_{b}^{k} \leq T_{*}^{k} \wedge S_{2}^{k}\right) \geq p_{1} / 2 \tag{4.39}
\end{equation*}
$$

or

$$
\begin{equation*}
\mathbb{P}_{z}\left(S_{2}^{k} \leq T_{b}^{k} \leq T_{*}^{k}\right) \geq p_{1} / 2 \tag{4.40}
\end{equation*}
$$

Suppose that the last estimate holds. By (4.38) and the strong Markov property applied at $S_{2}^{k}$,

$$
\mathbb{P}_{z}\left(S_{2}^{k} \leq S_{3}^{k} \leq T_{b}^{k} \leq T_{*}^{k}\right) \leq p_{1} / 4
$$

so, in view of (4.40),

$$
\mathbb{P}_{z}\left(S_{2}^{k} \leq T_{b}^{k} \leq S_{3}^{k} \wedge T_{*}^{k}\right) \geq p_{1} / 4
$$

It follows from this and (4.39) that

$$
\mathbb{P}_{z}\left(T_{b}^{k} \leq S_{3}^{k}\right) \geq p_{1} / 4,
$$

and, therefore, for $z \in \partial D_{*}$,

$$
\mathbb{P}_{z}\left(T_{b}^{k} \leq S_{4}^{k}\right) \geq p_{1} / 4
$$

We combine this with (4.37) using the strong Markov property at $S_{1}^{k}$ to see that for $z \in \bar{D}_{*}$,

$$
\mathbb{P}_{z}\left(T_{b}^{k} \leq S_{5,1}^{k}\right) \geq\left(1-p_{3}\right) p_{1} / 4=: p_{4}>0 .
$$

Applying the strong Markov property repeatedly at $S_{5, j}^{k}$ 's, we see that for $z \in \bar{D}_{*}$ and $j \geq 1$,

$$
\mathbb{P}_{z}\left(T_{b}^{k} \geq S_{5, j}^{k}\right) \leq\left(1-p_{4}\right)^{j}
$$

In other words,

$$
\begin{equation*}
\mathbb{P}_{z}\left(\left|\arg X_{T_{b}^{k}}^{k}-\arg X_{0}^{k}\right| \geq j(2 n+2) 2 \pi\right) \leq\left(1-p_{4}\right)^{j} \tag{4.41}
\end{equation*}
$$

Let $X^{0}$ be the ORBM corresponding to $\theta \equiv 0$. It is easy to see that

$$
\arg X_{t}^{k}-\arg X_{0}^{k}-\int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}
$$

has the same distribution as $\arg X_{t}^{0}-\arg X_{0}^{0}$. The estimate (4.41) applies to $X^{0}$; to prove that, one can apply the same argument as the one for $X^{k}$ 's or a direct elementary proof. Since

$$
\begin{aligned}
& \int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k} \\
& \quad=\left(\arg X_{t}^{k}-\arg X_{0}^{k}\right)-\left(\arg X_{t}^{k}-\arg X_{0}^{k}-\int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right)
\end{aligned}
$$

and (4.41) applies to both quantities within parentheses, we obtain for $z \in \bar{D}_{*}$ and $j \geq 1$,

$$
\begin{aligned}
& \mathbb{P}_{z}\left(\left|\int_{0}^{T_{b}^{k}} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right| \geq 2 j(2 n+2) 2 \pi\right) \\
& \quad \leq \mathbb{P}_{z}\left(\left|\arg X_{T_{b}^{k}}^{k}-\arg X_{0}^{k}\right| \geq j(2 n+2) 2 \pi\right) \\
& \quad+\mathbb{P}_{z}\left(\left|\arg X_{T_{b}^{k}}^{k}-\arg X_{0}^{k}-\int_{0}^{T_{b}^{k}} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right| \geq j(2 n+2) 2 \pi\right) \\
& \quad \leq \\
& \quad 2\left(1-p_{4}\right)^{j} .
\end{aligned}
$$

This implies that for some $c_{2}<\infty$ and all $z \in \bar{D}_{*}$ and all $k$,

$$
\begin{equation*}
\mathbb{E}_{z}\left[\left|\int_{0}^{T_{b}^{k}} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3}\right] \leq c_{2} \tag{4.42}
\end{equation*}
$$

Let $V_{0}=U_{1}=0$, and for $m \geq 1$,

$$
\begin{aligned}
V_{m} & =\inf \left\{t \geq U_{m}: X_{t}^{k} \in \mathcal{B}(0,1 / 2)\right\}, \\
U_{m+1} & =\inf \left\{t \geq V_{m}: X_{t}^{k} \notin \mathcal{B}(0,3 / 4)\right\} .
\end{aligned}
$$

Since $\mathbb{P}\left(U_{m+1}-V_{m}>1 \mid \mathcal{F}_{V_{m}}\right)>p_{5}>0$, we have

$$
\begin{equation*}
\mathbb{P}\left(U_{m} \leq 1\right) \leq c_{3}\left(1-p_{5}\right)^{m} \tag{4.43}
\end{equation*}
$$

Note that the local time $L^{k}$ does not increase on intervals [ $V_{m}, U_{m+1}$ ]. Hence,

$$
\begin{equation*}
\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}=\sum_{m=1}^{\infty} \int_{U_{m} \wedge 1}^{V_{m} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k} \tag{4.44}
\end{equation*}
$$

and, therefore,

$$
\begin{aligned}
\left|\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3}= & \left|\sum_{m=1}^{\infty} \int_{U_{m} \wedge 1}^{V_{m} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3} \\
\leq & 3 \sum_{m=1}^{\infty} \sum_{i \leq m} \sum_{j \leq m}\left|\mathbf{1}_{\left\{U_{m}<1\right\}} \int_{U_{m} \wedge 1}^{V_{m} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{2} \\
& \times\left|\mathbf{1}_{\left\{U_{i}<1\right\}} \int_{U_{i} \wedge 1}^{V_{i} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right| \\
& \times\left|\mathbf{1}_{\left\{U_{j}<1\right\}} \int_{U_{j} \wedge 1}^{V_{j} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{2} \\
\leq & 3 \sum_{m=1}^{\infty} \sum_{i \leq m} \sum_{j \leq m}\left[\mathbf{1}_{\left\{U_{m}<1\right\}}\left|\int_{U_{m} \wedge 1}^{V_{m} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3}\right. \\
& +\mathbf{1}_{\left\{U_{i}<1\right\}}\left|\int_{U_{i} \wedge 1}^{V_{i} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3} \\
& \left.+\mathbf{1}_{\left\{U_{j}<1\right\}}\left|\int_{U_{j} \wedge 1}^{V_{j} \wedge 1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3}\right]
\end{aligned}
$$

This, (4.42) and (4.43) imply that for some $c_{4}<\infty$, all $z \in \bar{D}_{*}$ and all $k$,

$$
\begin{equation*}
\mathbb{E}_{z}\left[\left|\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right|^{3}\right] \leq 3 \sum_{m=1}^{\infty} \sum_{i \leq m} \sum_{j \leq m} 3 c_{3}\left(1-p_{5}\right)^{m} c_{2}<c_{4} \tag{4.45}
\end{equation*}
$$

Step 3. For a fixed $z \in \bar{D}_{*}$ and all $k$, the processes $\left\{\left|X_{t}^{k}\right|, t \geq 0\right\}$ have the same distribution, that of 2-dimensional Bessel process on [0, 1], reflected at 1. Hence, $\mathbb{P}_{z}\left(\left|X_{1 / 2}^{k}\right|<1 / 4\right)>p_{6}$, where $p_{6}$ does not depend on $z \in \bar{D}_{*}$ and $k$. This and the Markov property at time $1 / 2$ can be used to show that the density of the distribution of $X_{1}^{k}$ under $\mathbb{P}_{z}$ is greater than $c_{5}>0$ on $\mathcal{B}(0,1 / 2)$, where $c_{5}$ does not depend on $z \in \bar{D}_{*}$ and $k$.

Let $\mathbb{P}_{x}^{k}$ denote the distribution of the process $X^{k}$ starting from $x$. Consider $z \in$ $\bar{D}_{*}$. We will construct a process $X^{k}$ with distribution $\mathbb{P}_{z}^{k}$ in a special way. First we will construct i.i.d. random vectors $A_{1}, A_{2}, A_{3}, \ldots$ The distribution of each $A_{j}$ is partly continuous, with density $c_{5}$ in $\mathcal{B}(0,1 / 2)$. With probability $1-c_{5} \pi / 4$, $A_{j}$ takes value $\boldsymbol{\Delta}$ (the cemetery state). Let $q_{1}^{k}$ be the density of $X_{1}^{k}$ under the distributions $\mathbb{P}_{z}^{k}$. Let $B_{1}$ be a random vector with density $q_{1}^{k}(x)-c_{5} \mathbf{1}_{\mathcal{B}(0,1 / 2)}(x)$ on $D_{*}$. With probability $c_{5} \pi / 4, B_{1}$ takes value $\Delta$. We construct $B_{1}$ so that it is equal to $\boldsymbol{\Delta}$ if and only if $A_{1} \neq \boldsymbol{\Delta}$. Moreover, we make the conditional distribution of $B_{1}$ given $\left\{B_{1} \neq \boldsymbol{\Delta}\right\}$ independent of $A_{j}$ 's.

In the following construction, the expression "Markov bridge" will refer to the Markov bridge corresponding to $\mathbb{P}^{k}$. If $A_{1} \in \mathcal{B}(0,1 / 2)$ then we let $\left\{X_{t}^{k}, 0 \leq t \leq 1\right\}$
be the Markov bridge between the points in time-space $(0, z)$ and $\left(1, A_{1}\right)$. If $A_{1}=$ $\boldsymbol{\Delta}$ then we let $\left\{X_{t}^{k}, 0 \leq t \leq 1\right\}$ be the Markov bridge between the points $(0, z)$ and $\left(1, B_{1}\right)$, otherwise independent of $A_{j}$ 's and $B_{1}$.

We continue by induction. Suppose that $\left\{X_{t}^{k}, 0 \leq t \leq n\right\}$ has been defined. Let $q_{n+1}^{k}\left(X_{n}^{k}, x\right)$ be the density of $X_{1}^{k}$ under the distribution $\mathbb{P}_{X_{n}^{k}}^{k}$. Let $B_{n+1}$ be a random vector with density $q_{n+1}^{k}\left(X_{n}^{k}, x\right)-c_{5} \mathbf{1}_{\mathcal{B}(0,1 / 2)}(x)$ on $D_{*}$. With probability $c_{5} \pi / 4$, this random vector takes value $\boldsymbol{\Delta}$. We construct $B_{n+1}$ so that it is equal to $\boldsymbol{\Delta}$ if and only if $A_{n+1} \neq \boldsymbol{\Delta}$. Moreover, we make the conditional distribution of $B_{n+1}$ given $\left\{B_{n+1} \neq \boldsymbol{\Delta}\right\}$ independent of $A_{j}$ 's and $\left\{X_{t}^{k}, 0 \leq t \leq n\right\}$, except that it has the density $q_{n+1}^{k}\left(X_{n}^{k}, x\right)-c_{5} \mathbf{1}_{\mathcal{B}(0,1 / 2)}(x)$ on $D_{*}$.

If $A_{n+1} \in \mathcal{B}(0,1 / 2)$ then we let $\left\{X_{t}^{k}, n \leq t \leq n+1\right\}$ be the Markov bridge between the points in time-space ( $n, X_{n}^{k}$ ) and $\left(n+1, A_{n+1}\right)$, otherwise independent of $A_{j}$ 's and $\left\{X_{t}^{k}, 0 \leq t \leq n\right\}$. If $A_{n+1}=\boldsymbol{\Delta}$ then we let $\left\{X_{t}^{k}, n \leq t \leq n+1\right\}$ be the Markov bridge between $\left(n, X_{n}^{k}\right)$ and $\left(n+1, B_{n+1}\right)$, otherwise independent of $A_{j}$ 's and $\left\{X_{t}^{k}, 0 \leq t \leq n\right\}$. It is easy to check that this inductive construction yields a process $\left\{X_{t}^{k}, t \geq 0\right\}$ with distribution $\mathbb{P}_{z}^{k}$.

Let $\Gamma_{n}^{k}=\int_{n}^{n+1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}$. Let $\mathbf{A}_{n}=\bigcup_{j=1}^{n}\left\{A_{n} \neq \boldsymbol{\Delta}\right\}$ and note that $\mathbb{P}\left(\mathbf{A}_{n}^{c}\right)=\left(1-c_{5} \pi / 4\right)^{n}=: c_{6}^{n}$, where $c_{6}<1$. If $\mathbf{A}_{n}$ holds then the trajectory of $\left\{X_{t}^{k}, n \leq t \leq n+1\right\}$ does not depend on $X_{1}^{k} . \operatorname{Hence}, \operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}}\right)=0$. We have

$$
\begin{aligned}
\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k}\right) & =\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}}+\Gamma_{n}^{k} \mathbf{1}_{n}^{c}\right)=\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}}\right)+\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}^{c}}\right) \\
& =\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}^{c}}\right)=\mathbb{E}\left(\Gamma_{1}^{k} \Gamma_{n}^{k} \mathbf{1}_{n}^{c}\right)-\mathbb{E} \Gamma_{1}^{k} \mathbb{E}\left(\Gamma_{n}^{k} \mathbf{1}_{\mathbf{A}_{n}^{c}}\right),
\end{aligned}
$$

so, in view of (4.45), for some $c_{10}<1$,

$$
\begin{aligned}
\left|\operatorname{Cov}\left(\Gamma_{1}^{k}, \Gamma_{n}^{k}\right)\right| & \leq\left(\mathbb{E}\left|\Gamma_{1}^{k}\right|^{3}\right)^{1 / 3}\left(\mathbb{E}\left|\Gamma_{n}^{k}\right|^{3}\right)^{1 / 3}\left(\mathbb{E} \mathbf{1}_{\mathbf{A}_{n}^{c}}^{3}\right)^{1 / 3}+\mathbb{E} \Gamma_{1}^{k}\left(\mathbb{E}\left(\Gamma_{n}^{k}\right)^{2}\right)^{1 / 2}\left(\mathbb{E} \mathbf{1}_{\mathbf{A}_{n}^{c}}^{2}\right)^{1 / 2} \\
& \leq c_{7} c_{4}^{n / 3}+c_{8} c_{4}^{n / 2} \leq c_{9} c_{10}^{n}
\end{aligned}
$$

It is easy to see that the estimate applies also to $n=1$ (possibly with new values of the constants). This implies that

$$
\begin{aligned}
& \operatorname{Var}\left(\int_{0}^{n} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right) \\
& \quad=\sum_{i=1}^{n} \sum_{j=1}^{n} \operatorname{Cov}\left(\int_{i-1}^{i} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}, \int_{j-1}^{j} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right) \\
& \quad \leq \sum_{i=1}^{n} \sum_{j=1}^{n} c 9 c_{10}^{|i-j|} \leq c_{11} n .
\end{aligned}
$$

It is elementary to check that the estimate also applies with noninteger upper limit, that is, for any $t>1$,

$$
\operatorname{Var}\left(\int_{0}^{t} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}\right) \leq c_{11} t
$$

This completes the proof of (4.33), and hence the proof of part (v) of the theorem.
(vi) The claim follows from the ergodic theorem if we show that under the stationary distribution $h(x) d x$,

$$
\begin{equation*}
\mathbb{E}_{h}\left[\arg ^{*} X_{1}\right]=\mu_{0} \tag{4.46}
\end{equation*}
$$

Recall that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \mu_{k}=\mu_{0} \tag{4.47}
\end{equation*}
$$

Theorem 3.2(iii) implies that

$$
\begin{equation*}
\mathbb{E}_{h_{k}}\left[\arg ^{*} X_{1}^{k}\right]=\mu_{k} \tag{4.48}
\end{equation*}
$$

It follows easily from definitions of $\arg { }^{*}$ and $\arg ^{*}$, and Theorem 3.5(iv) that $\arg ^{*} X_{1}^{k} \rightarrow \mathbf{a r g}^{*} X_{1}$ in distribution. Hence, in view of (4.47)-(4.48), the proof of (4.46) will be complete if we prove that the family $\left\{\arg ^{*} X_{1}^{k}\right\}_{k \geq 1}$ is uniformly integrable.

The following formula can be derived in the same way as (4.10) has been derived:

$$
\begin{equation*}
\arg ^{*} X_{1}^{k}=C_{L_{1}^{k}}^{*}+\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}+\left(\arg X_{1}^{k}-\arg X_{T_{1}^{k}}^{k}\right) \tag{4.49}
\end{equation*}
$$

Here $C^{*}$ is a Cauchy process with jumps larger than $2 \pi$ removed.
Recall that $S^{k}=\inf \left\{t>0: X_{t}^{k} \in \partial D_{*}\right\}$ and $T_{1}^{k}=\inf \left\{t>1: X_{t}^{k} \in \partial D_{*}\right\}$. So by the Markov property of $X^{k}$, under the stationary measure $h_{k}(x) d x,\left\{X_{s}^{k}-X_{1}^{k} ; 1 \leq\right.$ $\left.s \leq T_{1}^{k}\right\}$ has the same distribution as that of $\left\{X_{s}^{k}-X_{0}^{k} ; 0 \leq s \leq S^{k}\right\}$. It follows from the paragraph following (2.2) that under the stationary measure $h_{k}(x) d x$, $Y^{k}=\left|X^{k}\right|$ is a stationary 2-dimensional Bessel process in $(0,1]$ reflected at 1. Let $\sigma_{k}(a, b]=\int_{\{a<|x| \leq b\}} h_{k}(x) d x$. Then $\sigma_{k}(d r)$ is the stationary probability distribution of $Y^{k}$ so it is independent of $k$. This and the rotational invariance of Brownian motion imply that the distribution of $\arg X_{0}^{k}-\arg X_{S^{k}}^{k}$ does not depend on $k$. By an earlier remark, the distribution of $\arg X_{1}^{k}-\arg X_{T_{1}^{k}}^{k}$ is the same so it does not depend on $k$ either. Hence, the family $\left\{\arg X_{1}^{k}-\arg X_{T_{1}^{k}}^{k}\right\}, k \geq 1$, is uniformly integrable. The distribution of $L_{1}^{k}$ does not depend on $k$ so the same applies to $C_{L_{1}^{k}}^{*}$. Random variables $\int_{0}^{1} \tan \left(\theta_{k}\left(X_{s}^{k}\right)\right) d L_{s}^{k}$ are uniformly integrable by (4.45). All these remarks taken together with (4.49) show that the family $\left\{\arg ^{*} X_{1}^{k}\right\}_{k \geq 1}$ is uniformly integrable. This completes the proof of part (vi) of the theorem.
(vii) An explicit integral test was given in Burdzy and Marshall (1992): The ORBM in $D_{+}=\{z: \operatorname{Im} z>0\}$ with angle of reflection $\theta$ hits 0 with positive probability if and only if

$$
\begin{equation*}
\int_{0}^{1} \frac{1}{y} \operatorname{Re} \exp (i(\theta(i y)+i \widetilde{\theta}(i y))) d y<\infty \tag{4.50}
\end{equation*}
$$

where $\theta(z)$ is the bounded harmonic extension of $\theta$ to $D_{+}$and $\widetilde{\theta}$ is the harmonic conjugate of $\theta$ vanishing at $z=i$. In Burdzy and Marshall (1992) there was the added assumption that $\theta \in C^{1+\varepsilon}$, for some $\varepsilon>0$, except possibly at 0 . As noted in Burdzy and Marshall (1993), the same result holds if we only assume $\theta$ is measurable and $|\theta| \leq \pi / 2$. One way to transfer this result to $\theta \in \mathcal{T}$ is to set $\theta_{1}(t)=\theta\left(e^{i t}\right)$, for $t \in \mathbb{R}$, and $\theta_{1}(z)=\theta\left(e^{i z}\right)$ for $z \in D_{+}$as before. Then

$$
\int_{0}^{1} \frac{1}{y} \operatorname{Re} \exp \left(i\left(\theta_{1}(i y)+i \widetilde{\theta}_{1}(i y)\right)\right) d y=\int_{0}^{1} \frac{1}{y} \operatorname{Re} \exp \left(i(\theta+i \tilde{\theta})\left(e^{-y}\right)\right) d y
$$

Setting $r=e^{-y}$, we have $y=\ln 1 / r \sim 1-r$ on $\left[e^{-1}, 1\right]$ and so ORBM hits 1 with positive probability in $D_{*}$ if and only if the left-hand side of (3.13) is finite for $x=1$. By (2.19),

$$
1 /\left(h+i \widetilde{h}-i \mu_{0} / \pi\right)=\pi \cos \theta(0) e^{i(\theta+i \tilde{\theta})}
$$

and by taking real parts, the two integrals in (3.13) are equal.
Suppose that for some $z_{0} \in D_{*}$ and $x \in \partial D_{*}, \mathbb{P}_{z_{0}}\left(x \in \Gamma_{X}^{\theta}\right)>0$. A simple coupling argument shows that for some $r>0$ and $p>0, \mathbb{P}_{z}\left(x \in \Gamma_{X[0,1]}^{\theta}\right) \geq p$ for all $z \in \mathcal{B}\left(z_{0}, r\right) \subset D_{*}$. Since for every $k \geq 1, X_{t}$ returns to $\mathcal{B}\left(z_{0}, r\right)$ for some $t \geq k$ with probability one, we have by this "renewal property" that $\mathbb{P}_{z}\left(x \in \Gamma_{X}^{\theta}\right)=1$ for all $z \in \mathcal{B}\left(z_{0}, r\right)$.
(viii) Let $\rho$ denote the Prokhorov distance between probability measures [Billingsley (1999), Appendix III]. For any stochastic processes $V$ and $Z$, we will write $\rho(V, Z)$ to denote the distance between their distributions relative to $M_{1}$ distance between trajectories. For every $k$, one can find a sequence $\left(\theta_{k}^{n}\right)_{n \geq 1}$ of $C^{2}$ functions with values in $(-\pi / 2, \pi / 2)$ which converges to $\bar{\theta}_{k}$ as $n \rightarrow \infty$ in weak-* topology. Recall that $\bar{X}_{k}$ are defined relative to $\bar{\theta}_{k}$ in the same way that $X$ is defined relative to $\theta$. Processes $X^{k}$ are defined by (3.10) relative to $\theta_{k}$. By part (i) of the theorem, one can find a sequence $\theta_{k}^{n_{k}}: \partial D_{*} \rightarrow(-\pi / 2, \pi / 2)$ with the following properties. Let $X^{k, n_{k}}$ be the solution to (3.10) relative to $\theta_{k}^{n_{k}}$. Then $\rho\left(X^{k, n_{k}}, \bar{X}^{k}\right)<1 / k$. Moreover, we can choose $n_{k}$ 's so large that the sequence $\left(\theta_{k}^{n_{k}}\right)_{k \geq 1}$ converges to $\theta$ in weak-* topology. Since the sequence $\left(\theta_{1}, \theta_{1}^{n_{1}}, \theta_{2}, \theta_{2}^{n_{2}}, \theta_{3}, \theta_{3}^{n_{3}}, \ldots\right)$ converges to $\theta$, the sequence of processes $X^{1}, X^{1, n_{1}}, X^{2}, X^{2, n_{2}}, X^{3}, X^{3, n_{3}}, \ldots$ converges in distribution to a process $X^{\prime}$, by part (i) of the theorem. We must have $X=X^{\prime}$ in distribution, because $\left(\theta_{k}\right)_{k \geq 1}$ is a subsequence of $\left(\theta_{1}, \theta_{1}^{n_{1}}, \theta_{2}, \theta_{2}^{n_{2}}, \theta_{3}, \theta_{3}^{n_{3}}, \ldots\right)$. We see that $\rho\left(X^{k, n_{k}}, X\right) \rightarrow 0$. Since $\rho\left(X^{k, n_{k}}, \bar{X}^{k}\right)<1 / k$, we obtain $\rho\left(\bar{X}^{k}, X\right) \rightarrow 0$ as $k \rightarrow \infty$.

Proof of Proposition 3.7. The integral in (3.13) is equal to

$$
\int_{0}^{1} \frac{1}{1-r} \int_{\partial D_{*}} \frac{1-r^{2}}{|z-r x|^{2}} d \sigma(z) d r=\int_{\partial D_{*}} \int_{0}^{1} \frac{1+r}{|1-r x \bar{z}|^{2}} d r d \sigma(z)
$$

Let $w=x \bar{z}$. Then $|w|=1$ and

$$
\frac{1}{|1-r w|^{2}}=\frac{1}{(1-r w)(1-r \bar{w})}=\frac{1}{\bar{w}-w}\left(\frac{-w}{1-r w}-\frac{\bar{w}}{1-r \bar{w}}\right) .
$$

So

$$
\int_{0}^{1} \frac{1}{|1-r w|^{2}} d r=\frac{1}{\bar{w}-w} \ln \frac{1-w}{1-\bar{w}}=\frac{\arg (1-w)}{|1-w| \sin \arg (1-w)} \sim \frac{1}{|1-w|}
$$

Thus, the integral in (3.13) is finite if and only if (3.17) holds.
Proof of Proposition 3.11. An application of the Riemann mapping theorem shows that it suffices to prove the proposition for $D=D_{*}$.
(i) The expected occupation measure for an excursion law $H^{x}$ is a constant multiple of $K_{x}(\cdot)$ by (2.8). According to the definition, the ERBM is a "mixture" of excursion laws. This easily implies that the stationary distribution for $X$ has the density that is proportional to $\int_{\partial D_{*}} K_{x}(y) \nu(d x)$.
(iii) The function $h$ has a representation $h(y)=\int_{\partial D_{*}} K_{x}(y) \nu(d x)$. If one constructs an ERBM corresponding to $v$ then the stationary measure of this process is $h$ by part (i) of the proposition.

Proof of Theorem 3.12. (i) Since $\lim _{k \rightarrow \infty} \operatorname{dist}\left(x_{k}, \partial D_{*}\right)=0$, every subsequence of $x_{k}$ contains a further subsequence that converges to some point in $\partial D_{*}$. We will assume that the whole sequence $x_{k}$ converges to a point $x_{\infty} \in \partial D_{*}$. We will show that the limit distribution of $X^{k}$ does not depend on $x_{\infty}$. Hence, the result holds for every sequence satisfying $\lim _{k \rightarrow \infty} \operatorname{dist}\left(x_{k}, \partial D_{*}\right)=0$.

As was noted in the paragraph following (2.2), for any $r_{0} \in[0,1], t \geq 0$ and $\theta_{1}, \theta_{2} \in \mathcal{T}$, if $X^{k}$ is an ORBM in $D_{*}$ with the angle of reflection $\theta_{k}$ and $\left|X_{0}^{k}\right|=r_{0}$ for $k=1,2$, then the distributions of $\left|X_{t}^{1}\right|$ and $\left|X_{t}^{2}\right|$ are identical. Suppose that $X$ is an ORBM. Then $\mathbb{P}\left(\left|X_{t}\right| \in[1-\varepsilon, 1]\right) \leq c \varepsilon$ for some $c$ and all $\varepsilon \geq 0$. Fix an arbitrary $\varepsilon \in(0,1)$. Let $\mathcal{E}_{\varepsilon}^{*}=\left\{\mathrm{e}^{1}, \mathrm{e}^{2}, \ldots\right\}$ be the set of all excursions of $X$ from $\partial D_{*}$ which enter the ball $\mathcal{B}(0,1-\varepsilon)$, ordered according to their starting times. Let $S_{n}=S_{n}(\varepsilon)=\inf \left\{t \geq 0: \mathrm{e}_{t}^{n} \in \mathcal{B}(0,1-\varepsilon)\right\}$. It follows from the rotation invariance of Brownian motion that the distribution of $\left\{\exp \left(-i \arg \mathrm{e}_{S_{n}}^{n}\right) \mathrm{e}_{t}^{n}, t \geq S_{n}\right\}$ (the excursion rotated about 0 so that $\mathrm{e}_{S_{n}}^{n}$ is mapped to $1-\varepsilon \in \mathbb{R}$ ) does not depend on $n, \theta$ or the value taken by $S_{n}$.

Since the process $\left\{\mathrm{e}_{t}^{n}, t \geq S_{n}\right\}$ is Brownian motion killed upon hitting $\partial D_{*}$, its trajectory has modulus of continuity $c(\omega) \sqrt{2 r|\log r|}$, where $c(\omega)$ is finite for almost all $\omega$ [see Karatzas and Shreve (1991), Theorem 2.9.25]. If we time-reverse
$\mathrm{e}^{n}$ and rotate it so that it starts from 0 , then it will have the distribution $H^{0}$ conditioned by $\left\{\exists t>0: \mathrm{e}_{t} \in \mathcal{B}(0,1-\varepsilon)\right\}$. Hence, the claim about the modulus continuity can be extended as follows. The modulus of continuity of $\left\{\mathrm{e}_{t}^{n}, t \in(0, \zeta)\right\}$ is $c_{1}(\omega) \sqrt{2 r|\log r|}$, where $c_{1}(\omega)$ is finite for almost all $\omega$. This easily implies that for any sequence of random variables $V_{k}$ which converges to 0 in distribution, processes $\left\{\exp \left(-i V_{k}\right) \mathrm{e}_{t}^{n}, t \geq 0\right\}$ converge to $\left\{\mathrm{e}_{t}^{n}, t \geq 0\right\}$ in distribution in the Skorokhod topology as $k \rightarrow \infty$. Note that no assumptions on the joint distribution of $V_{k}$ and $\left\{\mathrm{e}_{t}^{n}, t \geq 0\right\}$ are needed.

Recall that $h_{k}(0)=1 / \pi$ for any $\left(h_{k}, \mu_{0, k}\right) \in \mathcal{H}$. Hence, $\int_{\partial D_{*}} h_{k}(x) d x=2$ and, therefore, $\nu\left(\partial D_{*}\right)=2$. It follows that $\nu / 2$ is a probability distribution on $\partial D_{*}$.

Let $\varepsilon_{\varepsilon}^{k}$ be defined relative to $X^{k}$ in the same way as $\mathcal{E}_{\varepsilon}^{*}$ has been defined relative to a generic $X$. We will suppress both $\varepsilon$ and $k$ in the notation for excursions, that is, we will write $\varepsilon_{\varepsilon}^{k}=\left\{\mathrm{e}^{1}, \mathrm{e}^{2}, \ldots\right\}$. In view of the opening remarks of this proof, it is routine to show that in order to prove part (i) of the theorem, it is sufficient to show that for any fixed $\varepsilon \in(0,1)$ and $n$, the joint distribution of $\left(\mathrm{e}_{0}^{1}, \mathrm{e}_{0}^{2}, \ldots, \mathrm{e}_{0}^{n}\right)$ converges to that of a sequence of $n$ i.i.d. random variables with distribution $v / 2$, as $k \rightarrow \infty$.

Let $\sigma_{t}^{k}=\inf \left\{s \geq 0: L_{s}^{k}>t\right\}$ and $A_{t}^{k}=\arg X_{\sigma_{t}^{k}}^{k}$, with the convention that $\arg X_{\sigma_{t}^{k}}^{k} \in[0,2 \pi)$. By abuse of notation, we define $\theta_{k}$ for real $x$ by $\theta_{k}(x)=\theta_{k}\left(e^{i x}\right)$. Let $B$ be Brownian motion in $\mathbb{C}$ starting at the origin and $S^{k}=\inf \left\{t>0: x_{k}+B_{t} \in\right.$ $\left.\partial D_{*}\right\}$. Let $\widehat{A}_{0}^{k}=\arg \left(x_{k}+B_{S^{k}}\right)$. Since $x_{k} \rightarrow x_{\infty} \in \partial D_{*}, a_{0}:=\lim _{k \rightarrow \infty} \widehat{A}_{0}^{k}=$ $\arg x_{\infty}$ a.s. Let $C_{t}$ be a Cauchy process with $C_{0}=0$ that is independent of $B$, and let $\widehat{A}_{t}^{k}$ be the solution to the SDE

$$
\begin{equation*}
\widehat{A}_{t}^{k}=\widehat{A}_{0}^{k}+C_{t}+\int_{0}^{t} \tan \theta_{k}\left(\widehat{A}_{s}^{k}\right) d s \tag{4.51}
\end{equation*}
$$

Clearly, $\widehat{A}_{0}^{k}$ has the same distribution as $A_{0}^{k}$. Let $\bar{A}_{t}^{k} \in[0,2 \pi)$ be the unique number such that $\bar{A}_{t}^{k}=\widehat{A}_{t}^{k}+j 2 \pi$ for some integer $j$. Then, by the conformal invariance of ORBMs presented in (4.5)-(4.7), the distribution of $\left\{\bar{A}_{t}^{k}, t \geq 0\right\}$ is the same as that of $\left\{A_{t}^{k}, t \geq 0\right\}$.

To incorporate our assumptions on $h_{k}$ and $1 / h_{k}$, we first note that by (2.15) and (2.11)

$$
\begin{equation*}
\tan \theta_{k}(z)=\frac{\mu_{k}(z)}{\pi h_{k}(z)}=\frac{\mu_{0, k}}{\pi h_{k}(z)}-\frac{\tilde{h}_{k}(z)}{h_{k}(z)}, \tag{4.52}
\end{equation*}
$$

for $z \in D_{*}$. If $f$ is Lipschitz with constant $\lambda$, then its modulus of continuity satisfies $\omega_{f}(\delta) \leq \lambda \delta$. By Garnett (2007), Theorem III.1.3, the modulus of continuity of $\tilde{f}$ satisfies

$$
\omega_{\tilde{f}}(\delta) \leq C \lambda \delta(1+\log \pi / \delta)
$$

where $C$ is a constant not depending on $f$ or $\delta$. So by assumption (c), $\tilde{h}_{k}$ are Dini continuous on $\bar{D}_{*}$, with constants depending only on $\lambda$, not $k$. We also conclude
that each $\theta_{k}$ and $\widetilde{\theta}_{k}$ are Dini continuous on $D_{*}$ and, therefore, on $\bar{D}_{*}$, by (2.18). In particular, (4.52) holds for $x \in \partial D_{*}$.

By a change of variables,

$$
\widehat{A}_{t / \mu_{0, k}}^{k}=\widehat{A}_{0}^{k}+C_{t / \mu_{0, k}}+\int_{0}^{t} \tan \theta_{k}\left(\widehat{A}_{r / \mu_{0, k}}^{k}\right) \frac{d r}{\mu_{0, k}}
$$

$$
\begin{equation*}
=\widehat{A}_{0}^{k}+C_{t / \mu_{0, k}}-\frac{1}{\mu_{0, k}} \int_{0}^{t} \frac{\widetilde{h}_{k}\left(A_{\left.r / \mu_{0, k}\right)}^{k}\right)}{h_{k}\left(A_{r / \mu_{0, k}}^{k}\right)} d r+\frac{1}{\pi} \int_{0}^{t} \frac{1}{h_{k}\left(A_{r / \mu_{0, k}}^{k}\right)} d r . \tag{4.53}
\end{equation*}
$$

By assumption (d), $h_{k}(z)=\int K_{z}(x) h_{k}(x)|d x|$ converges to $\int K_{z} v(d x): \equiv h(z)$, where $K_{z}$ is the Poisson kernel for $z \in D_{*}$. Since each $h_{k}$ is Lipschitz with constant $\lambda$ on $\partial D_{*}$ and, therefore, on $\bar{D}_{*}$, we have that $|h(z)-h(w)| \leq \lambda|z-w|$ for $z, w \in D_{*}$. Thus, $h$ extends to be Lipschitz with constant $\lambda$ on $\bar{D}_{*}$ and so $\nu(d x)=h(x)|d x|$.

Recall from Remark 3.13(iii) that the assumption (c) implies that all functions $1 / h_{k}$ are Lipschitz with the same constant. Without loss of generality, we will assume that the Lipshitz constant for $1 / h_{k}$ is $\lambda$. It follows that $1 / h$ is Lipshitz with constant $\lambda$.

Recall that $a_{0}:=\lim _{k \rightarrow \infty} \widehat{A}_{0}^{k}=\arg x_{\infty}$. By abuse of notation, let $h(x)=h\left(e^{i x}\right)$ for real $x$ and let $a_{t}$ be the solution to

$$
\begin{equation*}
a_{t}=a_{0}+\int_{0}^{t} \frac{1}{\pi h\left(a_{s}\right)} d s \tag{4.54}
\end{equation*}
$$

Let $t_{1}$ be such that $a_{t_{1}}=a_{0}+2 \pi$. Since

$$
\frac{\partial}{\partial t} v\left(\left[a_{0}, a_{t}\right]\right) / 2=(1 / 2) \frac{\partial}{\partial t} \int_{a_{0}}^{a_{t}} h(b) d b=(1 / 2) \frac{h\left(a_{t}\right)}{\pi h\left(a_{t}\right)}=\frac{1}{2 \pi}
$$

and $v\left(\left[a_{0}, a_{t_{1}}\right]\right) / 2=v\left(\left[a_{0}, a_{0}+2 \pi\right]\right) / 2=1$, we must have $t_{1}=2 \pi$. Hence, for $0 \leq s \leq t \leq 2 \pi$,

$$
\begin{equation*}
\nu\left(\left[a_{s}, a_{t}\right]\right) / 2=\frac{t-s}{2 \pi} . \tag{4.55}
\end{equation*}
$$

It follows from (4.53)-(4.54) that

$$
\widehat{A}_{t / \mu_{0, k}}^{k}-a_{t}=F_{t}^{k}+\frac{1}{\pi} \int_{0}^{t}\left(\frac{1}{h\left(A_{r / \mu_{0, k}}^{k}\right)}-\frac{1}{h\left(a_{r}\right)}\right) d r
$$

where

$$
\begin{align*}
F_{t}^{k}= & \widehat{A}_{0}^{k}-a_{0}+C_{t / \mu_{0, k}}-\frac{1}{\mu_{0, k}} \int_{0}^{t} \frac{\widetilde{h}_{k}\left(A_{\left.r / \mu_{0, k}\right)}^{k}\right.}{h_{k}\left(A_{r / \mu_{0, k}}^{k}\right)} d r  \tag{4.56}\\
& +\frac{1}{\pi} \int_{0}^{t}\left(\frac{1}{h_{k}\left(A_{r / \mu_{0, k}}^{k}\right)}-\frac{1}{h\left(A_{r / \mu_{0, k}}^{k}\right)}\right) d r .
\end{align*}
$$

Since $1 / h$ is Lipschitz with constant $\lambda$,

$$
\left|\widehat{A}_{t / \mu_{0, k}}^{k}-a_{t}\right| \leq \sup _{0 \leq s \leq 2 \pi}\left|F_{s}^{k}\right|+\frac{\lambda}{\pi} \int_{0}^{t}\left|\widehat{A}_{r / \mu_{0, k}}^{k}-a_{r}\right| d r
$$

for $0 \leq t \leq 2 \pi$. By Grönwall's inequality [see Bellman (1943)],

$$
\begin{equation*}
\left|\widehat{A}_{t / \mu_{0, k}}-a_{t}\right| \leq\left(\sup _{0 \leq s \leq 2 \pi}\left|F_{s}^{k}\right|\right) e^{\lambda t / \pi} \tag{4.57}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \sup _{0 \leq s \leq 2 \pi}\left|F_{s}^{k}\right|=0 \tag{4.58}
\end{equation*}
$$

in probability. By the definition of $a_{0}, \lim _{k} \widehat{A}_{0}^{k}-a_{0}=0$. By assumption (a), $\theta_{k}(0)=\int \theta_{k}(x)|d x| / 2 \pi$ converges to $\pi / 2$. But then $\mu_{0, k}=\tan \theta_{k}(0)$ converges to $+\infty$. Thus, $\sup _{0 \leq t \leq 2 \pi} C_{t / \mu_{0, k}}=0$, a.s. Since $\widetilde{h}_{k}$ and $1 / h_{k}$ are Dini continuous on $D_{*}$ with constant depending only on $\lambda$, and $h_{k}(0)=1 / \pi$ and $\widetilde{h}_{k}(0)=0$, we have that $\widetilde{h}_{k} / h_{k}$ is bounded on $\partial D_{*}$ by a constant independent of $k$. Thus, the first integral in (4.56) also tends to 0 .

If $\beta_{n}(f)$ denotes the $n$th Cesaro mean of $f$ on $\partial D_{*}$, then for continuous $f$, $\beta_{n}(f)$ converges uniformly on $\partial D_{*}$ to $f$, with the difference $\left\|\beta_{n}(f)-f\right\|_{\infty}$ depending only on the modulus of continuity of $f$ and $n$; see Hoffman (1962), page 18 . Since $1 / h_{k}$ and $1 / h$ are Lipschitz with constant $\lambda$, given $\varepsilon>0$ we can choose $n$ so that

$$
\begin{equation*}
\left\|1 / h_{k}-\beta_{n}\left(1 / h_{k}\right)\right\|_{\infty}<\varepsilon \quad \text { and } \quad\left\|1 / h-\beta_{n}(1 / h)\right\|_{\infty}<\varepsilon \tag{4.59}
\end{equation*}
$$

By assumption (d), $h_{k}$ converges to $h$, uniformly on compact subsets of $D_{*}$, and since $1 / h_{k}$ are uniformly bounded, $1 / h_{k}$ converges to $1 / h$ uniformly on compact subsets of $D_{*}$. Since $1 / h_{k}$ are uniformly bounded, this also implies $1 / h_{k}$ converges to $1 / h$ weak-* and, therefore, for $k$ sufficiently large, and $n$ fixed,

$$
\begin{equation*}
\left\|\beta_{n}\left(1 / h_{k}\right)-\beta_{n}(1 / h)\right\|_{\infty}<\varepsilon \tag{4.60}
\end{equation*}
$$

By (4.59), (4.60) and the triangle inequality, $1 / h_{k}$ converges uniformly to $1 / h$. We conclude that the second integral in (4.56) tends to 0 as well, proving the claim.

We will need a generalization of the above results (4.57) and (4.58). Let $D_{u}=\{z \in \mathbb{C}: \operatorname{Im} z>0\}$ be the upper half-plane. Let $H^{x}$ be the excursion law for Brownian motion in $D_{*}$, for excursions starting from $x \in \partial D_{*}$ and let $\widehat{H}^{x}$ be the excursion law for Brownian motion in $D_{u}$, for excursions starting from $x \in \partial D_{u}$. The measure $\widehat{H}^{0}(\mathrm{e}(\zeta-) \in d x)$ is the distribution of the end point of the excursion under $\widehat{H}^{0}$. It is also the Lévy measure for the Cauchy process. Let

$$
\mu_{\varepsilon}(d x)=\widehat{H}^{0}\left(\sup _{t \in[0, \zeta)} \operatorname{Ime}_{t}<|\log (1-\varepsilon)|, \mathrm{e}(\zeta-) \in d x\right)
$$

The measure $\mu_{\varepsilon}$ is the Lévy measure for a pure jump process, say $C_{t}^{\varepsilon}$, similar to the Cauchy process, except that it has fewer big jumps. We can choose a right continuous version of $C^{\varepsilon}$, and so $\sup _{0 \leq s \leq t}\left|C_{s}^{\varepsilon}\right| \rightarrow 0$, a.s., as $t \rightarrow 0$. We let $\widehat{A}_{t}^{k, \varepsilon}$ be the solution to the equation analogous to (4.51),

$$
\begin{equation*}
\widehat{A}_{t}^{k, \varepsilon}=\widehat{A}_{0}^{k, \varepsilon}+C_{t}^{\varepsilon}+\int_{0}^{t} \tan \theta_{k}\left(\widehat{A}_{s}^{k, \varepsilon}\right) d s \tag{4.61}
\end{equation*}
$$

An argument analogous to that showing (4.57) and (4.58) proves that for every fixed $\varepsilon>0$,

$$
\begin{equation*}
\sup _{0 \leq s \leq 2 \pi}\left|\widehat{A}_{s / \mu_{0, k}}^{k, \varepsilon}-a_{s}\right| \rightarrow 0 \tag{4.62}
\end{equation*}
$$

in probability, as $k \rightarrow \infty$.
Recall the definition of $\mathcal{E}_{\varepsilon}^{k}=\left\{\mathrm{e}^{1}, \mathrm{e}^{2}, \ldots\right\}$ from the beginning of the proof. We claim that for any fixed $\varepsilon \in(0,1)$ and $n$, the joint distribution of $\left(\mathrm{e}_{0}^{1}, \mathrm{e}_{0}^{2}, \ldots, \mathrm{e}_{0}^{n}\right)$ converges to that of a sequence of $n$ i.i.d. random variables with distribution $v / 2$, as $k \rightarrow \infty$.

We will present a special construction of $\left(\mathrm{e}_{0}^{1}, \mathrm{e}_{0}^{2}, \ldots, \mathrm{e}_{0}^{n}\right)$. The heuristic meaning of the construction is the following. Excursions that reach $\mathcal{B}(0,1-\varepsilon)$ occur as a Poisson process with constant intensity on the local time scale. If we have already observed $\mathrm{e}^{1}, \mathrm{e}^{2}, \ldots, \mathrm{e}^{m}$, the next excursion will occur after an exponential waiting time on the local time scale, where the local time has the same distribution as the process $\widehat{A}_{t}^{k, \varepsilon}$. This process, suitably rescaled, behaves like the function $a_{t}$ according to (4.62). By (4.55), a point on the boundary chosen in a uniform manner on the $a_{t}$ scale has the distribution $v / 2$. We will also need a fact that, on small time intervals, exponential density is almost constant. The process $\widehat{A}_{t}^{k, \varepsilon}$ represents rapid rotation along the unit circle and the exponential clock will chose a point on the circle according to the distribution very close to $\nu / 2$, because the almost constant exponential density (on small intervals) is transformed into the density of $\nu / 2$ by the function $a_{t}$.

Suppose that excursions $\mathrm{e}^{1}, \mathrm{e}^{2}, \ldots, \mathrm{e}^{m}$ have been already generated, for some $m \geq 0$. If $m \geq 1$, let $T_{m}$ be the time when $\mathrm{e}^{m}$ ended. If $m=0$, then we take $T_{0}$ to be the first hitting time of $\partial D_{*}$ by $X^{k}$. Unless stated otherwise, every new random object introduced below will be assumed to be independent from all random objects constructed so far.

By conformal invariance of excursion laws,

$$
H^{x}\left(\exists t \in[0, \zeta): \mathrm{e}_{t} \in \mathcal{B}(0,1-\varepsilon)\right)=\widehat{H}^{0}\left(\exists t \in[0, \zeta): \operatorname{Im~}_{t} \geq|\log (1-\varepsilon)|\right)
$$

and the last quantity is equal to $1 /|\log (1-\varepsilon)|$ [see Burdzy (1987) for the justification of both claims].

Consider an exponential random variable $\alpha$ with density $f_{\alpha}(t)$ and expected value $|\log (1-\varepsilon)|$, independent of objects constructed so far. For every $\delta>0$, there exists $c_{3}>0$ so small that for any interval $\left[t, t+c_{3}\right]$ and any $s_{1}, s_{2} \in\left[t, t+c_{3}\right]$,
we have $f_{\alpha}\left(s_{1}\right) / f_{\alpha}\left(s_{2}\right) \in(1-\delta, 1+\delta)$. We generate an integer-valued random variable $N$, such that $\mathbb{P}(N=j)=\mathbb{P}\left(\alpha \in\left[j 2 \pi / \mu_{0, k},(j+1) 2 \pi / \mu_{0, k}\right]\right)$ for $j \geq 0$. We consider a solution to (4.61) with $\widehat{A}_{0}^{k, \varepsilon}=\arg X_{T_{m}}^{k}+N 2 \pi / \mu_{0, k}$. We generate a random variable $\alpha^{\prime}$ with the same distribution as $\alpha$ conditioned to be in $[N, N+1$ ). Note that we can take $\delta>0$ so small and then let $k$ be so large that, in view of (4.55) and (4.62), the distribution of $\exp \left(i \widehat{A}_{\alpha^{\prime}-N}^{k, \varepsilon}\right)$ is arbitrarily close to $\nu / 2$.

We generate an excursion $\overline{\mathrm{e}}^{m+1}$ with the (probability) distribution $H^{0}(\cdot \mid \exists t \in$ $\left.[0, \zeta): \mathrm{e}_{t} \in \mathcal{B}(0,1-\varepsilon)\right)$. We let $\widehat{\mathrm{e}}_{t}^{m+1}=\exp \left(i \widehat{A}_{\alpha^{\prime}-N}^{k, \varepsilon}\right) \overline{\mathrm{e}}_{t}^{m+1}$.

In view of the preceding remarks, the distribution of $\widehat{\mathrm{e}}_{0}^{m+1}$ is arbitrarily close to $v / 2$, conditional on the trajectories of $\mathrm{e}^{1}, \ldots, \mathrm{e}^{m}$, if $k$ is arbitrarily large. According to our construction, the joint distribution of ( $\mathrm{e}^{1}, \ldots, \mathrm{e}^{m}, \widehat{\mathrm{e}}^{m+1}$ ) is the same as that of $\left(\mathrm{e}^{1}, \ldots, \mathrm{e}^{m}, \mathrm{e}^{m+1}\right)$. We conclude that for any fixed $\varepsilon \in(0,1)$ and $n$, the joint distribution of $\left(\mathrm{e}_{0}^{1}, \mathrm{e}_{0}^{2}, \ldots, \mathrm{e}_{0}^{n}\right)$ converges to that of a sequence of $n$ i.i.d. random variables with distribution $\nu / 2$, as $k \rightarrow \infty$. This completes the proof of part (i) of the theorem.
(ii) We will generalize Example 3.14. Suppose that $h$ is positive on $\bar{D}_{*}$, harmonic in $D_{*}$ and Lipschitz on $\bar{D}_{*}$. Then $1 / h$ is Lipschitz on $\bar{D}_{*}$. Set $h_{k}(z)=$ $h((1-1 / k) z)$ and suppose $\mu_{0, k} \rightarrow \infty$. Then $\left(h_{k}, \mu_{0, k}\right) \leftrightarrow \theta_{k} \in \mathcal{T}$ as in Theorem 2.1, satisfy the assumptions of part (i) and the conclusions of that part of the theorem with the given $h$.

Proof of Theorem 3.15. (i) Suppose that $X_{0}$ has the stationary distribution with density $h$. Then for every $t>0$,

$$
\begin{aligned}
\mathbb{E}[c(t)] & =\mathbb{E}\left[\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s\right]=\int_{0}^{t} \mathbb{E}\left[\left|f^{\prime}\left(X_{s}\right)\right|^{2}\right] d s=\int_{0}^{t} \int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h(x) d x d s \\
& =t \int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h(x) d x=t \int_{D} \bar{h}(x) d x<\infty
\end{aligned}
$$

It follows that under the stationary distribution, $\zeta=\infty$, a.s. This implies that $\zeta=$ $\infty, \mathbb{P}_{x}$-a.s., for almost all $x \in D_{*}$.

Consider an $x \in D_{*}$ and $r>0$ so small that $\overline{\mathcal{B}(x, r)} \subset D_{*}$. The exit distributions from $\mathcal{B}(x, r)$ are mutually absolutely continuous for any two points $y, z \in \mathcal{B}(x, r)$. Let $T$ be the exit time from $\mathcal{B}(x, r)$. It is easy to see that $c(T)<\infty, \mathbb{P}_{y}$-a.s., for every $y \in \mathcal{B}(x, r)$. Since $\zeta=\infty, \mathbb{P}_{y}$-a.s., for at least one $y \in \mathcal{B}(x, r)$, it follows that this claim holds for all $y \in \mathcal{B}(x, r)$. The claim holds for all balls such that $\overline{\mathcal{B}(x, r)} \subset D_{*}$ so $\zeta=\infty, \mathbb{P}_{y}$-a.s., for all $y \in D_{*}$.
(ii) This part follows easily from conformal invariance of Brownian motion killed upon leaving a domain.
(iii) This claim follows from the interpretation of the stationary distribution as the long time occupation measure, the definition of $\widehat{h}$ and the "clock" $c(t)$. We sketch the easy argument. For an arbitrarily small $\varepsilon>0$ and $x, y \in D_{*}$ we can find
$r>0$ so small that

$$
\begin{aligned}
& \lim _{t \rightarrow \infty} \frac{\int_{0}^{t} \mathbf{1}_{\left\{Y_{t} \in \mathcal{B}(f(x), r)\right\}} d s}{\int_{0}^{t} \mathbf{1}_{\left\{Y_{t} \in \mathcal{B}(f(y), r)\right\}} d s} \\
& \quad \leq \lim _{t \rightarrow \infty} \frac{\sup _{z \in f^{-1}(\mathcal{B}(f(x), r))}\left|f^{\prime}(z)\right|^{2} \int_{0}^{t} \mathbf{1}_{\left\{X_{t} \in f^{-1}(\mathcal{B}(f(x), r))\right\}} d s}{\inf _{z \in f^{-1}(\mathcal{B}(f(y), r))}\left|f^{\prime}(z)\right|^{2} \int_{0}^{t} \mathbf{1}_{\left\{X_{t} \in f^{-1}(\mathcal{B}(f(y), r))\right\}} d s} \\
& \quad \leq \lim _{t \rightarrow \infty} \frac{\sup _{z \in f^{-1}(\mathcal{B}(f(x), r))}\left|f^{\prime}(z)\right|^{2}(1+\varepsilon)\left|f^{\prime}(x)\right|^{-2} \int_{0}^{t} \mathbf{1}_{\left\{X_{t} \in \mathcal{B}(x, r)\right\}} d s}{\inf _{z \in f^{-1}(\mathcal{B}(f(y), r))}\left|f^{\prime}(z)\right|^{2}(1-\varepsilon)\left|f^{\prime}(y)\right|^{-2} \int_{0}^{t} \mathbf{1}_{\left\{X_{t} \in \mathcal{B}(y, r)\right\}} d s} \\
& \quad \leq \lim _{t \rightarrow \infty} \frac{\sup _{z \in f^{-1}(\mathcal{B}(f(x), r))}\left|f^{\prime}(z)\right|^{2}(1+\varepsilon)\left|f^{\prime}(x)\right|^{-2} \sup _{z \in \mathcal{B}(x, r)} h(z)}{\inf _{z \in f^{-1}(\mathcal{B}(f(y), r))}\left|f^{\prime}(z)\right|^{2}(1-\varepsilon)\left|f^{\prime}(y)\right|^{-2} \inf _{z \in \mathcal{B}(y, r)} h(z)} .
\end{aligned}
$$

If we let $\varepsilon, r \rightarrow 0$, then the right-hand side converges to $h(x) / h(y)$. Hence, the limsup of the left-hand side is at most $h(x) / h(y)$. A similar argument shows that the liminf of the left-hand side is at least $h(x) / h(y)$. This implies that the stationary density for $Y$ is proportional to $h \circ f^{-1}$. Hence, it must be equal to $\widehat{h}$.
(iv) It follows from the definition of the "clock" $c(t)$ and the ergodic theorem that, a.s.,

$$
\lim _{t \rightarrow \infty} \frac{c(t)}{t}=\int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h(x) d x=\|\bar{h}\|_{L^{1}(D)}
$$

We have already proved (3.3). That claim and the above formula imply for $z=$ $f(0)$,

$$
\begin{align*}
\lim _{t \rightarrow \infty} \frac{\boldsymbol{\operatorname { r r g }}^{*}\left(Y_{t}-z\right)}{t} & =\lim _{t \rightarrow \infty} \frac{\boldsymbol{\operatorname { r r g }}^{*} X_{c^{-1}(t)}}{t} \\
& =\lim _{t \rightarrow \infty} \frac{\boldsymbol{\operatorname { r g }}^{*} X_{t}}{c(t)}=\lim _{t \rightarrow \infty} \frac{\boldsymbol{\operatorname { a r g }}^{*} X_{t}}{t} \cdot \frac{t}{c(t)}  \tag{4.63}\\
& =\lim _{t \rightarrow \infty} \frac{\boldsymbol{\operatorname { r g }}^{*} X_{t}}{t} \lim _{t \rightarrow \infty} \frac{t}{c(t)}=\frac{\mu_{0}}{\|\bar{h}\|_{L^{1}(D)}}=\frac{\mu(0)}{\|\bar{h}\|_{L^{1}(D)}}
\end{align*}
$$

Next we prove (3.4). Suppose that $f=\tau$ is a one-to-one analytic map of $D_{*}$ onto $D_{*}$ such that $\tau(0)=z$, as in Lemma 2.3. Then $\tau$ is a Möbius transformation. Let $\widehat{h}=h \circ \tau /\|h \circ \tau\|_{1}, \widehat{\mu}_{0}=\mu(z) /\|h \circ \tau\|_{1}$, and $\widehat{\theta}=\theta \circ \tau$. Then by Lemma 2.3, $\widehat{\theta} \leftrightarrow\left(\widehat{h}, \widehat{\mu}_{0}\right)$. If $\bar{h}=\widehat{h} \circ \tau^{-1}=h /\|h \circ \tau\|_{1}$ then

$$
\|\bar{h}\|_{1}=1 /\|h \circ \tau\|_{1} .
$$

By (4.63),

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(X_{t}-z\right)}{t}=\frac{\widehat{\mu}_{0}}{\|\bar{h}\|_{1}}=\widehat{\mu}_{0}\|h \circ \tau\|_{1}=\mu(z) \tag{4.64}
\end{equation*}
$$

Finally, we prove (3.23) in full generality along the same lines as in (4.63). For any $z \in D$, by (3.4),

$$
\begin{aligned}
\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(Y_{t}-z\right)}{t} & =\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(X_{c^{-1}(t)}-f^{-1}(z)\right)}{t}=\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(X_{t}-f^{-1}(z)\right)}{c(t)} \\
& =\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(X_{t}-f^{-1}(z)\right)}{t} \cdot \frac{t}{c(t)} \\
& =\lim _{t \rightarrow \infty} \frac{\arg ^{*}\left(X_{t}-f^{-1}(z)\right)}{t} \lim _{t \rightarrow \infty} \frac{t}{c(t)}=\frac{\mu\left(f^{-1}(z)\right)}{\|\bar{h}\|_{1}}
\end{aligned}
$$

(v) Let $\theta$ correspond to $\left(h, \mu_{0}\right)$. Let $Y$ be constructed as in (3.20)-(3.22). Then it is easy to see that $Y$ satisfies conditions (a) and (b) of part (v).
(vi) This follows directly from the Itô formula and Theorem 3.1.

We now present an example showing that a conformal mapping may not always map an ORBM in one planar domain to another ORBM, in the sense of Theorem 3.15.

EXAMPLE 4.1. Let $S$ be a two-dimensional infinite wedge with corner at the origin 0 and angle $0<\alpha<2 \pi$. Consider $\theta_{1}, \theta_{2} \in(-\pi / 2, \pi / 2)$ and suppose that each $\theta_{k}$ represents the angle of reflection on one of the two sides of the wedge, measured from the inward normal toward the origin 0. In Varadhan and Williams (1985), it was shown that there exists a strong Markov process that behaves like Brownian motion in the interior of the wedge and reflects instantaneously at the boundary with the oblique angle of reflection given by $\theta_{k}$. This process, called obliquely reflected Brownian motion in Varadhan and Williams (1985), is characterized as the unique solution to the corresponding submartingale problem away from the vertex. It was shown Varadhan and Williams (1985) that the process enters 0 in a finite time and then stays there forever (i.e., it cannot be continued as a Markov process beyond that time) if and only if $\beta:=\left(\theta_{1}+\theta_{2}\right) / \alpha \geq 2$. Let $D$ be an acute triangle obtained by truncation of the infinite wedge $S$. Assume that $\theta_{1}$ and $\theta_{2}$ are such that $\beta \geq 2$, set $\theta_{3}=0$ on the edge opposite to 0 , and assume that the analogues of $\beta$ at the other two vertices are strictly less than 2 . Let $f$ be a conformal mapping from the unit disk $D_{*}$ onto the Jordan domain $D$ and note that it extends to a homeomorphism from $\bar{D}_{*}$ onto $\bar{D}$. Let $\theta(x)$ be the pre-image of the $\theta$-function on $\partial D$ by $f$. Then $\theta$ is a piecewise constant function on $\partial D_{*}$ taking values in $(-\pi / 2, \pi / 2)$. Thus by Theorem 3.5 , the ORBM $X$ in $D_{*}$ with reflection angle $\theta$ is a continuous, conservative Markov process having stationary distribution $h(x) d x$. Consequently, $Z_{t}=f\left(X_{t}\right)$ is a continuous, conservative Markov process on $\bar{D}$. The process $Z$ is an extension of killed Brownian motion in $D$ modulo a time change in the sense that for every $t \geq 0$ and $\tau_{t}=\inf \left\{s \geq t: Z_{s} \in \partial D\right\}$, the process $\left\{Z_{s}, s \in\left[t, \tau_{t}\right)\right\}$ is a time change of Brownian motion killed upon exiting $D$. Let $\widehat{\tau}_{t}=\inf \left\{s \geq t: Z_{s}=0\right\}$ for $t \geq 0$. Then the process $\left\{Z_{s}, s \in\left[t, \widehat{\tau}_{t}\right)\right\}$ is a
time change of the obliquely reflected Brownian motion in $D$ killed upon hitting 0 . More precisely, let $x_{0}=f^{-1}(0), \sigma_{x_{0}}=\inf \left\{t \geq 0: X_{t}=x_{0}\right\}, c(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s$ and $c^{-1}(t)=\inf \{s: c(s)>t\}$. Then $Y_{t}=f\left(X_{c^{-1}(t)}\right), t \in\left[0, \sigma_{x_{0}}\right)$, is obliquely reflected Brownian motion in $D$ killed upon hitting 0 . The result in Varadhan and Williams (1985) and Theorem 3.15 imply that $c\left(\sigma_{x_{0}}\right)=\int_{0}^{\sigma_{x_{0}}}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s<\infty$ but $\int_{0}^{\sigma_{x_{0}}+\varepsilon}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s=\infty$ a.s. for every $\varepsilon>0$, and that $h \circ f^{-1} \notin L^{1}(D)$.

Proof of Theorem 3.17. (i) The argument given in the proof of Theorem 3.15(i) which shows that $\zeta=\infty$, a.s., applies verbatim in the present case because we have assumed that $\|\bar{h}\|_{L^{1}(D)}<\infty$.

Every harmonic function $h_{k}$ is bounded because $\theta_{k}$ is continuous and takes values in $(-\pi / 2, \pi / 2)$. Hence, the function $\bar{h}_{k}:=h_{k} \circ f^{-1}$ is also bounded. Since $D$ is bounded, it follows that $\left\|\bar{h}_{k}\right\|_{L^{1}(D)}<\infty$. Once again, the argument given in the proof of Theorem 3.15(i) applies and shows that $\zeta_{k}=\infty$, a.s., for all $k$.
(ii) Recall the representation of $X$ as the Poisson point process on the space $\mathbb{R}_{+} \times \mathcal{C}_{D_{*}}$ (see Definition 3.9). Excursion laws are conformally invariant in the sense of the transformation in (3.20)-(3.22) by Burdzy (1987), Proposition 10.1, so $Y$ can be represented as a Poisson point process on $\mathbb{R}_{+} \times \mathcal{C}_{D}$. In other words, $Y$ is an ERBM and it only remains to identify the corresponding $\left(\bar{v}(d x), \bar{H}^{x}\right)_{x \in \partial D}$. We can arbitrarily set the excursion intensity $\bar{v}$ to be $\bar{\nu}(A)=v\left(f^{-1}(A)\right)$ for $A \subset \partial D$, in view of Remark 3.10(ii).

We will find the matching normalization for $\bar{H}^{x}$. Fix some $z \in D$ and suppose that $r>0$ is very small. The Green function $G_{x}(\cdot)$ in $D$ has the property that

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{\inf _{y \in \partial \mathcal{B}(z, r)} G_{y}(z)}{\sup _{y \in \partial \mathcal{B}(z, r)} G_{y}(z)}=\lim _{r \rightarrow 0} \frac{\inf _{y \in \partial \mathcal{B}(z, r)} G_{y}(z)}{|\log r|}=1 . \tag{4.65}
\end{equation*}
$$

Let $T_{A}$ denote the hitting time of $A$. Recall that $G_{x}(\cdot)$ is the density of the expected occupation time for Brownian motion in $D$ killed upon exiting from $D$. Also, by Remark $3.10(\mathrm{v})$, the density of the expected occupation time for $\bar{H}^{x}$ is $\bar{c}_{x} K_{x}(\cdot)$. Hence, for $x \in \partial D$, by the strong Markov property of $\bar{H}^{x}$,

$$
\bar{c}_{x} K_{x}(z)=\int_{\partial \mathcal{B}(z, r)} G_{y}(z) \bar{H}^{x}\left(X\left(T_{\partial \mathcal{B}(z, r)}\right) \in d y\right)
$$

This and (4.65) imply that, as $r \rightarrow 0$,

$$
\begin{equation*}
|\log r| \bar{H}^{x}\left(T_{\partial \mathcal{B}(z, r)}<\infty\right)=\bar{c}_{x} K_{x}(z)+o(1) . \tag{4.66}
\end{equation*}
$$

An analogous formula holds for excursion laws $H^{x}$ in $D_{*}$, with the corresponding constants $c_{x}$ equal to each other, by rotation invariance. Let $N(d x, z, r, D, t)$ be the number of excursions of the ERBM in $D$ (here $D$ can be also $D_{*}$ ), which started from $d x \subset \partial D$ before time $t$ and hit $\partial \mathcal{B}(z, r)$ before their lifetime. It is easy to see that

$$
\begin{equation*}
\lim _{r \downarrow 0, \varepsilon \downarrow 0} \lim _{t \rightarrow \infty} \frac{N(d x, z, r, D, t)}{N(d x, z, r(1+\varepsilon), D, t)}=1 . \tag{4.67}
\end{equation*}
$$

By the ergodic theorem,

$$
\lim _{r \rightarrow 0} \lim _{t \rightarrow \infty} \frac{N\left(d x, 0, r, D_{*}, t\right)}{N\left(d y, 0, r, D_{*}, t\right)}
$$

exists and is equal to $v(d x) / \nu(d y)$. The fact that small balls are mapped by $f$ onto regions very close to balls, (4.67), and the definition of $Y$ as a transform of $X$ imply that for $Y$ we have

$$
\lim _{r \rightarrow 0} \lim _{t \rightarrow \infty} \frac{N(d x, f(0), r, D, t)}{N(d y, f(0), r, D, t)}=\frac{v\left(f^{-1}(d x)\right)}{v\left(f^{-1}(d y)\right)}=\frac{\bar{v}(d x)}{\bar{v}(d y)}
$$

This in turn implies that all $\bar{c}_{x}$ in (4.66) must be equal to each other so, in view of Remark 3.10(iii), we may take all of them to be equal to 1 .
(iii) The processes $X^{k}$ converge to $X$ in the sense of finite dimensional distributions according to Theorem 3.12. A stronger assertion follows from the proof of that theorem. Fix some $\varepsilon>0$ and let $\mathrm{e}^{k, n}$ be the $n$th excursion of the process $X^{k}$ which hits the ball $\mathcal{B}(0,1-\varepsilon)$, and let $T_{\varepsilon}^{k, n}$ be the hitting time of the ball. Then the joint distributions of $\left\{\mathrm{e}_{t}^{k, n}, t \in\left[T_{\varepsilon}^{k, n}, \zeta\right)\right\}, n \geq 1, \varepsilon>0, \varepsilon \in \mathbb{Q}$, converge as $k \rightarrow \infty$, in the Skorokhod topology. By the Skorokhod lemma, we can assume that $\left\{\mathrm{e}_{t}^{k, n}, t \in\left[T_{\varepsilon}^{k, n}, \zeta\right)\right\}, n \geq 1, \varepsilon>0, \varepsilon \in \mathbb{Q}$, converge a.s., as $k \rightarrow \infty$, in the Skorokhod topology. Hence, $X_{t}^{k} \rightarrow X_{t}$ for almost all $t \geq 0$ simultaneously, a.s.

The function $f$ is Lipschitz continuous inside every disc $\mathcal{B}(0,1-\rho), \rho \in(0,1)$. This implies that for every $\varepsilon>0$ and $n$, the images of the excursions $f\left(\mathrm{e}_{t}^{k, n}\right)$ converge as $k \rightarrow \infty$, a.s., in the Skorokhod topology over their lifetimes to the corresponding excursion of $Y$. It will suffice to show that for every fixed $t>0$, the clocks $c_{k}(t)$ converge to $c(t)$ in probability (note that the clocks are monotone functions).

Let

$$
\begin{align*}
c(t) & =\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s \quad \text { for } t \geq 0 \\
Y(t) & =f\left(X_{c^{-1}(t)}\right) \quad \text { for } t \in[0, \infty)  \tag{4.68}\\
c_{k}(t) & =\int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} d s \quad \text { for } t \geq 0, \\
Y^{k}(t) & =f\left(X_{c_{k}^{-1}(t)}^{k}\right) \quad \text { for } t \in[0, \infty) . \tag{4.69}
\end{align*}
$$

Then $Y$ and $Y^{k}$,s have distributions as specified in the statement of the theorem.
We will assume for a moment that $X_{0}^{k}$ 's and $X_{0}$ have stationary distributions. Let $D_{\varepsilon}=D_{*} \backslash \mathcal{B}(0,1-\varepsilon)$. By assumption (i),

$$
\begin{equation*}
\int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h(x) d x=\int_{D} h \circ f^{-1} d x<\infty \tag{4.70}
\end{equation*}
$$

By assumption $D$ is bounded, so that $\int_{D_{*}}\left|f^{\prime}\right|^{2} d x=\operatorname{Area}(D)<\infty$ and by the proof of Theorem 3.12, $h_{k}$ converges uniformly to $h$. Thus,

$$
\begin{equation*}
\sup _{k} \int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h_{k}(x) d x<\infty, \tag{4.71}
\end{equation*}
$$

and, moreover,

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h(x) d x=0 \tag{4.72}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \sup _{k} \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h_{k}(x) d x=0 . \tag{4.73}
\end{equation*}
$$

For $\varepsilon>0$ (suppressed in the notation), let

$$
\begin{array}{cl}
\bar{c}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}} d s, & \widehat{c}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s \\
\bar{c}_{k}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}} d s, & \widehat{c}_{k}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s
\end{array}
$$

Fix some $t \geq 0$ and arbitrarily small $p_{1}, \delta>0$. It follows from (4.70)-(4.73) that there exists $\varepsilon_{1}>0$ such that for $\varepsilon \in\left(0, \varepsilon_{1}\right)$ and all $k$,

$$
\begin{aligned}
\mathbb{E}[\bar{c}(t)] & =\mathbb{E}\left[\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}} d s\right]=\int_{0}^{t} \mathbb{E}\left[\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}}\right] d s \\
& =\int_{0}^{t} \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h(x) d x d s=t \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h(x) d x d s<p_{1} \delta
\end{aligned}
$$

and

$$
\begin{aligned}
\mathbb{E} \bar{c}_{k}(t) & =\mathbb{E} \int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}} d s=\int_{0}^{t} \mathbb{E}\left(\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}}\right) d s \\
& =\int_{0}^{t} \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h_{k}(x) d x d s=t \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h_{k}(x) d x d s<p_{1} \delta .
\end{aligned}
$$

It follows that for $\varepsilon \in\left(0, \varepsilon_{1}\right)$ and all $k$,

$$
\begin{equation*}
\mathbb{P}(\bar{c}(t) \geq \delta) \leq p_{1} \quad \text { and } \quad \mathbb{P}\left(\bar{c}_{k}(t) \geq \delta\right) \leq p_{1} \tag{4.74}
\end{equation*}
$$

For almost all $s>0, X_{s}^{k} \rightarrow X_{s}$, a.s., and $\mathbb{P}\left(X_{s} \in \partial \mathcal{B}(0,1-\varepsilon)\right)=0$. Hence, for almost all $s>0$, a.s.,

$$
\lim _{k \rightarrow \infty}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}}=\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}},
$$

and, therefore, by the bounded convergence theorem, a.s.,

$$
\begin{aligned}
\lim _{k \rightarrow \infty} \widehat{c}_{k}(t) & =\lim _{k \rightarrow \infty} \int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s \\
& =\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s=\widehat{c}(t)
\end{aligned}
$$

This and (4.74) imply that for every fixed $t>0$, a.s.,

$$
\lim _{k \rightarrow \infty} c_{k}(t)=c(t)
$$

because $\delta$ and $p_{1}$ can be chosen arbitrarily close to 0 .
We can remove the assumption that the processes are in the stationary distribution as in the proof of Theorem 3.15(i).
(iv) This can be proved just as part (iii) of Theorem 3.15.
(v) Let $h^{*}=\widehat{h} \circ f$. Then $h^{*}$ is a positive harmonic function in $D_{*}$ and so $\left\|h^{*}\right\|_{1}=\pi h^{*}(0)<\infty$. Let $h=h^{*} /\left\|h^{*}\right\|_{1}$. By assumption, $h$ is Lipschitz continuous on $\bar{D}_{*}$ and strictly positive on $\partial D_{*}$. Let $h_{k}(z)=\left(1-2^{-k}\right)^{1 / 2} h\left(\left(1-2^{-k}\right) z\right)$. Then $h_{k}$ is a sequence of positive harmonic functions in $D_{*}$ with $L^{1}$ norm equal to 1 and $C^{2}$ on $\bar{D}_{*}$, such that $h_{k} \rightarrow h$ uniformly on compact subsets of $D_{*}$, and both $h_{k}$ and $1 / h_{k}$ are $\lambda$-Lipschitz on $\partial D_{*}$ for some $\lambda>0$ when $k$ is sufficiently large. Let $\mu_{0, k}=k$, and let $\theta_{k}$ correspond to $\left(h_{k}, \mu_{0, k}\right)$. Let $Y^{k}$,s and $Y$ be constructed as in the statement of Theorem 3.17. Then it is easy to see that the stationary distribution for ERBM $Y$ has density $\widehat{h}$.

Proof of Theorem 3.18. Let $D_{*}^{k}=f^{-1}\left(D_{k}\right)$. It is easy to see that $D_{*}^{k}$ converge to $D_{*}$ in the sense that for every $r<1$ there exists $k_{0}$ such that $\mathcal{B}(0, r) \subset$ $D_{*}^{k}$ for $k \geq k_{0}$. Set $x_{0}=f^{-1}\left(y_{0}\right)=f_{k}^{-1}\left(y_{0}\right), a_{0}=f(0)$ and $a_{k}=f_{k}(0)$. Then $a_{k} \rightarrow a_{0}$. Let $h_{k}=\bar{h} \circ f_{k} /\left\|\bar{h} \circ f_{k}\right\|_{1}=\bar{h} \circ f_{k} /\left(\pi \bar{h}\left(a_{k}\right)\right)$, and let $\theta_{k} \leftrightarrow\left(h_{k}, \mu_{0}\right)$. Note that $h_{k}$ are smooth and bounded on $\bar{D}_{*}$ and therefore $\theta_{k}$ are smooth on $\partial D_{*}$ and take values in $(-\pi / 2, \pi / 2)$. Let $h=\bar{h} \circ f /\|\bar{h} \circ f\|_{1}=\bar{h} \circ f /\left(\pi \bar{h}\left(a_{0}\right)\right)$, and let $\theta \leftrightarrow$ $\left(h, \mu_{0}\right)$. Then $h_{k}$ converges to $h$ uniformly on compact subsets of $D_{*}$ and by (2.18), $\theta_{k}(z)$ converges to $\theta(z)$ uniformly on compact subsets of $D_{*}$. Since the closed unit ball in $L^{\infty}\left(\partial D_{*} ;|d x|\right)=L^{1}\left(\partial D_{*} ;|d x|\right)^{*}$ is compact in the weak-* topology, it follows that $\theta_{k}$ converges to $\theta$ in the in the weak-* topology in $L^{\infty}\left(\partial D_{*} ;|d x|\right)$. Let $X^{k}$ be the solution to (2.1) corresponding to $\theta_{k}$ and starting from $x_{0}=f^{-1}\left(y_{0}\right)$ and let $X$ be constructed as in Theorem 3.5, relative to $\theta$ and also starting from $x_{0}=f^{-1}\left(y_{0}\right)$. Let

$$
\begin{equation*}
c(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s \quad \text { and } \quad Y(t)=f\left(X_{c^{-1}(t)}\right) \tag{4.75}
\end{equation*}
$$

$$
\text { for } t \in[0, \infty)
$$

$$
\begin{aligned}
& c_{k}(t)=\int_{0}^{t}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} d s \quad \text { and } \quad Y^{k}(t)=f_{k}\left(X_{c_{k}^{-1}(t)}^{k}\right) \\
& \text { for } t \in[0, \infty)
\end{aligned}
$$

Then $Y$ and $Y^{k}$,s have distributions as specified in the statement of the theorem.
We will assume for a moment that $X_{0}^{k}$ 's and $X_{0}$ have stationary distributions. According to Theorem 3.5(i), the processes $\left\{X_{s}^{k}, 0 \leq s \leq t\right\}$ converge weakly to $\left\{X_{s}, 0 \leq s \leq t\right\}$ in $M_{1}^{\mathcal{T}}$ topology. By the Skorokhod theorem, we can assume that
all these processes are defined on the same probability space and $\left\{X_{s}^{k}, 0 \leq s \leq t\right\}$ converge almost surely to $\left\{X_{s}, 0 \leq s \leq t\right\}$ in $M_{1}^{\mathcal{T}}$ topology.

Let $D_{\varepsilon}=D_{*} \backslash \mathcal{B}(0,1-\varepsilon)$. We have

$$
\begin{align*}
\int_{D_{*}}\left|f^{\prime}(x)\right|^{2} h(x) d x & =\frac{1}{\pi \bar{h}\left(a_{0}\right)} \int_{D} \bar{h} d x<\infty,  \tag{4.77}\\
\sup _{k} \int_{D_{*}}\left|f_{k}^{\prime}(x)\right|^{2} h_{k}(x) d x & =\frac{1}{\pi \bar{h}\left(a_{k}\right)} \sup _{k} \int_{D_{k}} \bar{h} d x<\infty \tag{4.78}
\end{align*}
$$

and, moreover, as in (4.72) and (4.73)

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \sup _{k} \int_{D_{\varepsilon}}\left|f_{k}^{\prime}(x)\right|^{2} h_{k}(x) d x=0 \tag{4.79}
\end{equation*}
$$

For $\varepsilon>0$ (suppressed in the notation), let

$$
\begin{array}{cc}
\bar{c}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}} d s, & \widehat{c}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s, \\
\bar{c}_{k}(t)=\int_{0}^{t}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}} d s, & \widehat{c}_{k}(t)=\int_{0}^{t}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s .
\end{array}
$$

Fix some $t \geq 0$ and arbitrarily small $p_{1}, \delta>0$. It follows from (4.79)-(4.80) that there exists $\varepsilon_{1}>0$ such that for $\varepsilon \in\left(0, \varepsilon_{1}\right)$ and all $k$,

$$
\begin{aligned}
\mathbb{E}[\bar{c}(t)] & =\mathbb{E}\left[\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}} d s\right]=\int_{0}^{t} \mathbb{E}\left[\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in D_{\varepsilon}\right\}}\right] d s \\
& =\int_{0}^{t} \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h(x) d x d s=t \int_{D_{\varepsilon}}\left|f^{\prime}(x)\right|^{2} h(x) d x d s<p_{1} \delta
\end{aligned}
$$

and

$$
\begin{aligned}
\mathbb{E}\left[\bar{c}_{k}(t)\right] & =\mathbb{E}\left[\int_{0}^{t}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}} d s\right]=\int_{0}^{t} \mathbb{E}\left[\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in D_{\varepsilon}\right\}}\right] d s \\
& =\int_{0}^{t} \int_{D_{\varepsilon}}\left|f_{k}^{\prime}(x)\right|^{2} h_{k}(x) d x d s=t \int_{D_{\varepsilon}}\left|f_{k}^{\prime}(x)\right|^{2} h_{k}(x) d x d s<p_{1} \delta .
\end{aligned}
$$

It follows that for $\varepsilon \in\left(0, \varepsilon_{1}\right)$ and all $k$,

$$
\begin{equation*}
\mathbb{P}(\bar{c}(t) \geq \delta) \leq p_{1} \quad \text { and } \quad \mathbb{P}\left(\bar{c}_{k}(t) \geq \delta\right) \leq p_{1} \tag{4.81}
\end{equation*}
$$

For any fixed $\varepsilon>0$, there is $k_{0} \geq 1$ such that

$$
\begin{equation*}
\sup _{x \in \mathcal{B}(0,1-\varepsilon)}\left(\left|f^{\prime}(x)\right|^{2} h(x) \vee \sup _{k \geq k_{0}}\left|f_{k}^{\prime}(x)\right|^{2} h_{k}(x)\right)<\infty . \tag{4.82}
\end{equation*}
$$

For every fixed $s>0, X_{s}^{k} \rightarrow X_{s}$, a.s., and $\mathbb{P}\left(X_{s} \in \partial \mathcal{B}(0,1-\varepsilon)\right)=0$. Hence, for every fixed $s>0$, a.s.,

$$
\lim _{k \rightarrow \infty}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}}=\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}}
$$

and, therefore, by the bounded convergence theorem, a.s.,

$$
\begin{aligned}
\lim _{k \rightarrow \infty} \widehat{c}_{k}(t) & =\lim _{k \rightarrow \infty} \int_{0}^{t}\left|f_{k}^{\prime}\left(X_{s}^{k}\right)\right|^{2} \mathbf{1}_{\left\{X_{s}^{k} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s \\
& =\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} \mathbf{1}_{\left\{X_{s} \in \mathcal{B}(0,1-\varepsilon)\right\}} d s=\widehat{c}(t)
\end{aligned}
$$

This and (4.81) imply that for every fixed $t>0$, a.s.,

$$
\begin{equation*}
\lim _{k \rightarrow \infty} c_{k}(t)=c(t) \tag{4.83}
\end{equation*}
$$

because $\delta$ and $p_{1}$ can be chosen arbitrarily close to 0 .
It follows easily from the definition (3.9) of convergence in $M_{1}^{\mathcal{T}}$ topology and continuity of $f$ on $\bar{D}_{*}$ that convergence of $X^{k}$ to $X$ in $M_{1}^{\mathcal{T}}$ topology implies convergence of $f\left(X^{k}\right)$ to $f(X)$ in $M_{1}^{\mathcal{T}}$ topology. This is because the transformation $f$ affects only the first components of the pairs $\left(y_{n}(s), t_{n}(s)\right)$ and $(y(s), t(s))$ in (3.9). When the clocks are changed, the second components are affected as well. Then we use (4.83) to conclude that $Y^{k}$ converge to $Y$ in $M_{1}^{\mathcal{T}}$ topology.

We can remove the assumption that the processes are in the stationary distribution as in the proof of Theorem 3.15(i).

Proof of Theorem 3.8. Take a sequence of $C^{2}$ functions $\theta_{k}: \partial D_{*} \rightarrow$ $(-\pi / 2, \pi / 2)$ that converges to $\theta \in \mathcal{T}$ in weak-* topology as elements of the dual space of $L^{1}\left(\partial D_{*}\right)$. Let $X^{k}$ be ORBM on $D_{*}$ that satisfies (3.10). By Theorem $3.5(\mathrm{i}), X^{k}$ converges weakly in $M_{1}^{\mathcal{T}}$-topology to $X$, so does $f\left(X^{k}\right)$ to $f(X)$. Define

$$
c_{k}(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}^{k}\right)\right|^{2} d s \quad \text { and } \quad c(t)=\int_{0}^{t}\left|f^{\prime}\left(X_{s}\right)\right|^{2} d s
$$

By an argument similar to that proving (4.83), we can show that $\lim _{k \rightarrow \infty} c_{k}(t)=$ $c(t)$ a.s. for every fixed $t>0$. Consequently by the argument as in the second to the last paragraph in the proof of Theorem 3.18, $f\left(X_{c_{k}^{-1}(t)}^{k}\right)$ converges weakly in $M_{1}^{\mathcal{T}}$ topology to $f\left(X_{c^{-1}(t)}\right)$. It is easy to see that $f\left(X_{c^{-1}(t)}\right)$ has stationary distribution with density $\bar{h}$. Since $f$ is smooth on $\bar{D}_{k}$ and $\theta_{k} \circ f^{-1}$ converges to $\theta \circ f^{-1} \in \mathcal{T}$ in weak-* topology as elements of the dual space of $L^{1}\left(\partial D_{*}\right)$, it follows from Theorem 3.5 that $f\left(X_{c^{-1}(t)}\right)$ is the ORBM on $D_{*}$ with reflection angle $\theta \circ f^{-1}$.

Proof of Theorem 3.19. This theorem can be proved just like Theorem 3.18. All we have to check is whether the following claims hold: (4.77), (4.78), (4.79), (4.80) and (4.82). They are all easily seen to hold in the present context.

Example 4.2. We will sketch an example of a bounded domain $D$, an oblique angle of reflection $\theta$ and the corresponding ORBM with a stationary measure whose density $h$ is not in $L^{1}(D)$. The construction is a typical fractal-type argument; a construction similar in spirit can be found in Section 4 of Bass and Burdzy (1992). We will not supply a formal proof because it would require a lot of space and the claim is rather specialized.

Let $D_{0}=(0,1)^{2}$, and for $k \geq 1$ and small $r_{k} \in\left(0,2^{-k-2}\right)$ (to be specified later), let

$$
\begin{aligned}
D_{k} & =\mathcal{B}\left(2^{-k}-i 2^{-k}, 2^{-k-2}\right) \\
D_{k}^{\prime} & =\left(2^{-k}-r_{k}, 2^{-k}+r_{k}\right) \times\left(-2^{-k}, 2^{-k}\right) \\
D & =D_{0} \cup \bigcup_{k \geq 1}\left(D_{k} \cup D_{k}^{\prime}\right)
\end{aligned}
$$

The boundary $\partial D$ is smooth except for a countable number of points. We will specify the reflection angle relative to the inward normal vector $\mathbf{n}$ at each boundary point where $\mathbf{n}$ is well defined. For all points $x \in \partial D \cap\left(\partial D_{0} \cup \partial D_{k}\right), k \geq 0$, we let $\theta(x)=0$. In other words, the reflection is in the normal direction at the points on the boundary of the square $D_{0}$ and on the (arcs of the) circles $\partial D_{k}$.

It remains to define the angle of reflection for the part of $\partial D$ which lies on the sides of very thin channels $D_{k}^{\prime}$. To make the example simple, we let the angle of reflection be $\pi / 2$ or $-\pi / 2$, at $x \in \partial D \cap \partial D_{k}^{\prime}, k \geq 1$, so that the reflected process is pushed down towards $D_{k}$. It would be more accurate to say that the process is teleported to $D_{k}$ if it hits the side of a channel $\partial D \cap \partial D_{k}^{\prime}$ because it has a jump that takes it to $\partial D_{k}$.

Heuristically speaking, the ratio of the average amounts of time spent by ORBM in $D_{k}$ and $D_{0}$ can be made arbitrarily large by making $r_{k}$ sufficiently small. The reason is that ORBM will jump to $D_{k}$ when it hits the boundary of $D_{k}^{\prime}$. Going the other way is much harder - the process has to go though the very thin channel connecting $D_{k}$ and $D_{0}$ without hitting the sides of the channel. Let $a_{k}$ be the ratio of the average amounts of time spent by ORBM in $D_{k}$ and $D_{0}$. If we make all $a_{k} \geq 1$, then $\sum_{k \geq 1} a_{k}=\infty$ and it follows that there is no stationary probability distribution for ORBM. Every stationary measure has to have infinite mass.

It is clear that the ORBM described above is well defined as long as it does not hit $(0,0)$. An elementary argument can be used to show that the ORBM will not hit $(0,0)$ at a finite time, a.s., if we make the channels sufficiently thin (i.e., $r_{k}$ 's sufficiently small).

Proof of Theorem 3.20. Parts (i) and (ii) are special cases of Theorems 1 and 2 of Aikawa (2000).

For part (iii), let $D$ be the image of the unit disk by the map $F(z)=\sqrt{1-z}$ and let $h(w)=\operatorname{Re}((1+z) /(1-z))$ where $z=F^{-1}(w)$. Then for the region $C$ in the
disk given by $1-|z|^{2}>|1-z|$ (an approximate cone),

$$
\int_{D} h(w) d w=\int_{C} \operatorname{Re}((1+z) /(1-z))\left|F^{\prime}(z)\right|^{2} d z \geq \int_{C}|1-z|^{-2} d z / 4
$$

since $\operatorname{Re}((1+z) /(1-z))=\left(1-|z|^{2}\right) /|1-z|^{2}$. This latter integral is infinite by integrating in polar coordinates centered at $z=1$.

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