STABLE HYDRODYNAMIC LIMIT FLUCTUATIONS OF A CRITICAL BRANCHING PARTICLE SYSTEM IN A RANDOM MEDIUM

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A particle system in Euclidean space is considered where the particles are subject to spatial motion according to a symmetric stable law and to a critical branching law in the domain of attraction of a stable law. The "branching intensity" may be position-dependent (varying medium) or be given by a realization of a random field (random medium). It is shown that under natural assumptions the hydrodynamic limit fluctuations around the macroscopic flow are the same as those given by the "averaged medium," the limit being a generalized stable Ornstein-Uhlenbeck process. The convergence proof is based on an analysis of a nonlinear integral equation with random coefficients.

1. Introduction. The question of the existence and properties of equilibrium states for critical branching particle systems has been extensively studied; see, for example, the comprehensive presentation by Kerstan, Matthes and Mecke (1982), as well as Dawson, Fleischmann, Foley and Peletier (1986), and the references therein. On the other hand, there is considerable current interest in understanding the behavior of distributed systems in random media (environments); see Kozlov (1985) and Papanicolaou (1983) for recent surveys.

In Dawson and Fleischmann (1983, 1985) a branching model in a random medium was introduced in a discrete space—time setting. Another model with branching was treated by Greven (1985). These branching models have the advantage of being more tractable than most interacting particle models such as those arising in statistical physics and consequently can serve as test cases.

In this article we will study the large scale fluctuations of a critical branching particle system in the spirit of Holley and Stroock (1978), Dawson (1981) and Dittrich (1987). Our model is different in that the branching law may have infinite variance and may also depend on a random medium. The treatment of the random medium involves ideas from the method of averaging [Kozlov (1985) gives a recent review of these ideas].

The model under consideration may be described roughly as follows. The states of the process are counting measures on Euclidean space R^d of dimension

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d, representing a system of particles. These particles independently undergo spherically symmetric stable motions with exponent α , $0 < \alpha \le 2$. In addition each particle branches at rate V > 0 according to a critical law in the domain of attraction of a stable law with exponent $1 + \beta$, $0 < \beta \le 1$. The branching law depends on a parameter h which is location-dependent: If a particle branches at x, then it produces particles with intensity h(x). This yields a branching particle system in a spatially varying medium $h = \{h(x): x \in R^d\}$. Finally, we assume that h is given by a realization of an ergodic random field; we refer to h as the random medium. A more precise description will be given in Section 2.

We denote by N_t the state at time $t\geq 0$, and we consider the counting measure-valued process $N\equiv\{N_t,\,t\geq 0\}$. The distribution of the initial measure N_0 is chosen to depend on a parameter $\varepsilon>0$ in such a way that as $\varepsilon\to 0$ the scaled random measure $\varepsilon^d N_0(\varepsilon^{-1}\cdot)$ converges in distribution to a nonrandom measure Λ_0 . In this sense N_0 and Λ_0 describe the initial spatial distribution of particles at the microscopic and macroscopic levels, respectively.

The long-time behavior of the critical branching particle system N involves a competition between the long-range "mixing" effect of the particle motions and the "clumping" effect of the branching. In order to exhibit this we consider the process viewed in the natural scaling associated with the particle motions, namely

$$X_{\epsilon}^{\epsilon}(A) := \epsilon^{d} N(\epsilon^{-\alpha} t, \epsilon^{-1} A).$$

One objective of this article is to show that in dimensions $d > \alpha/\beta$ we obtain the hydrodynamic limit

$$X_t^{\epsilon}(\cdot) \to S_t \Lambda_0(\cdot)$$
 as $\epsilon \to 0$,

where S_t is the semigroup of the particle motion. In other words in high dimensions the spatial diffusion dominates. In the case of low dimensions (which is not considered in this article) the branching dominates and the system degenerates to zero in this scaling [cf. Dawson and Fleischmann (1983, 1985) and Fleischmann and Gärtner (1988)]. For hydrodynamic limits of another branching model refer also to Dobrushin and Siegmund-Schultze (1982).

The main objective of this article is to investigate the fluctuations around this hydrodynamic limit. In the case $\alpha=2$, $\beta=1$ and the medium is deterministic the limit fluctuations are given by a well-known Gaussian generalized Ornlstein–Uhlenbeck process [see Holley and Stroock (1978), Gorostiza (1983), Bojdecki and Gorostiza (1986) and Dittrich (1987)]. In particular, we consider the rescaled fluctuation process

$$Y_t^{\varepsilon} \coloneqq \varepsilon^{-k} [X_t^{\varepsilon} - S_t \Lambda_0],$$

where $k = (d\beta - \alpha)/(1 + \beta)$. The main result of this article is that Y^{ϵ} converges as $\epsilon \to 0$ to a generalized Ornstein–Uhlenbeck process. This limit process can also be described as the solution of a generalized Langevin equation which is driven by Gaussian noise when $\beta = 1$ and by stable noise when $\beta < 1$.

Our approach can be regarded as a generalization of Dittrich (1987). The key to the proof of these results is an analysis of the scaling properties of a nonlinear

integral equation which enables us to establish the convergence of finite-dimensional distributions. [This key scaling argument is given in the proof of Proposition (5.3.1).] Concerning tightness, due to the lack of finite second moments, we cannot apply the usual methods based on the increasing process of martingales [as in Holley and Stroock (1978) and Gorostiza (1983)]. Instead we develop a method to verify Aldous' tightness criterion based on a moment inequality involving the characteristic functional.

The outline of the article is as follows. In Section 2 we describe in detail the branching system in a random medium. In Section 3 we introduce the state spaces for the process N and the limiting fluctuation process Y. The precise statements of the main results described above are given in Section 4 and the proofs are given in Section 5.

2. Description of the system. The motion of the particles is assumed to be a spherically symmetric stable process of R^d with exponent α , $0 < \alpha \le 2$.

Fix β , $0 < \beta \le 1$. A varying medium, denoted by $h = \{h(x): x \in \mathbb{R}^d\}$, is an element of the space H of real-valued measurable functions on \mathbb{R}^d satisfying $0 \le h \le (1+\beta)^{-1}$ (ther reason for imposing the upper bound is related to the branching law to be introduced). Let \mathscr{H} be the σ -algebra of subsets of H generated by the mappings $h \to \int_{\mathbb{R}^d} f(x)h(x) dx$ with $f \in L^1(\mathbb{R}^d)$.

The random medium is prescribed by a probability measure m on (H, \mathcal{H}) , not concentrated on $h \equiv 0$. We assume that m is stationary and ergodic with respect to the translation group on R^d . The "classical" medium h = const. is included as a special case. Another family of examples is given by $h(x) = (2/\pi(1+\beta))|\arctan g(x)|$, where g is any continuous regular stationary Gaussian random field on R^d .

To each $h \in H$ is associated the family $g^h \equiv \{g^h(x): x \in R^d\}$ of offspring generating functions

(2.1)
$$g^{h}(x,s) = s + h(x)(1-s)^{1+\beta}$$
$$= s + h(x) \sum_{n=0}^{\infty} {1+\beta \choose n} (-s)^{n}, \quad |s| < 1,$$

where s is a complex number. Here and in the following $z^{1+\beta} = \exp[(1+\beta)\log z]$ is always understood in the sense of the principal branch of the logarithm.

 $g^h(x)$ describes the *local branching law*: If a particle located at x splits, then n particles are produced with probability

$$p_n^h(x) = \begin{cases} h(x) \binom{1+\beta}{n} (-1)^n & \text{if } n \neq 1, \\ 1 - h(x) (1+\beta) & \text{if } n = 1. \end{cases}$$

Note that the branching law is critical, that is, it has mean 1, and it belongs to the domain of normal attraction of a stable law with exponent $1 + \beta$ [if h(x) > 0]; thus for $\beta < 1$ it has infinite variance. In the case $\beta = 1$, $p_n^h(x) = 0$ for n > 2 and the variance is 2h(x). The condition $p_1^h(x) \ge 0$ requires $h(x) \le (1 + \beta)^{-1}$. Since the probability that more than one particle is produced at x is

 $\beta h(x)$, we may interpret h(x) as the branching intensity or fertility at x. Note also that $h(x) = p_0^h(x)$ is the probability of extinction at x.

Each particle has an exponentially distributed lifetime with parameter V > 0, and at the end of its life it branches according to g^h given above. The offspring appear at the site where their parent branches.

The motions, lifetimes and branchings of all particles are independent of each other and of the random medium (except for the branching intensity). It is assumed that the random medium is sampled first (and fixed for all time) according to m and then the system evolves in the resulting varying medium.

The description of the system is completed by specifying an initial distribution of particles which may possibly depend on the random medium.

Let $N_i(A)$ denote the number of particles in the system which lie in the Borel set $A \subset \mathbb{R}^d$ at time $t \geq 0$. For fixed $h \in H$ and in the case of a finite number of initial particles the existence of process N follows from Moyal (1962, 1964) and Ikeda, Nagasawa and Watanabe (1968a, b, 1969). In the next section we make precise the formulation of N as a measure-valued process and its extension to an appropriate class of infinite measures.

3. State spaces for the basic processes.

3.1. The counting process N. The first objective of this section is to introduce a space $\mathcal{N}_p(\mathbb{R}^d)$ of infinite counting measures with an appropriate topology which serves as the state space for N.

Given a topological space M we define the following spaces of real-valued functions on M: C(M) is the space of continuous bounded functions on M with the topology induced by the sup-norm $\| \|_{\infty}$; $C_c(M)$ is the subset of functions having compact support, and if M is locally compact $C_0(M)$ is the subset of functions vanishing at infinity. The collection of nonnegative elements of a real function space will be indicated by the index +.

When $M = R^d$ with its usual norm $|\cdot|$, we denote $\phi_p(x) = (1 + |x|^2)^{-p}$, $x \in \mathbb{R}^d$, $p \ge 0$, and we define the spaces

$$C_p(R^d) = \left\{\phi \in C(R^d) \colon \|\phi/\phi_p\|_{\infty} < \infty\right\}, \qquad p > 0,$$
 $C_{p,0}(R^d) = \left\{\phi \in C(R^d) \colon \phi/\phi_p \in C_0(R^d)\right\}, \qquad p > 0.$

Note that $C_{p,\,0}(R^d)\subset C_p(R^d)\subset C_0(R^d)$. For $p\geq 0$, let $\mathscr{M}_p(R^d)$ denote the space of nonnegative Radon measures μ on R^d such that $\int_{R^d}\phi_p\,d\mu<\infty$. We equip $\mathscr{M}_p(R^d)$ with the p-vague topology, that is, the smallest topology making the maps $\mu \to \int \phi \ d\mu$ continuous for all $\phi \in$ $C_c(R^d)_+ \cup \{\phi_p\}$. We denote $\langle \mu, \phi \rangle = \int \phi \, d\mu$ for $\phi \in C_p(R^d)$ and $\mu \in \mathcal{M}_p(R^d)$. The space $\mathcal{M}_p(R^d)$ is Polish but not locally compact. The Lebesgue measure on R^d belongs to $\mathcal{M}_p(R^d)$ for p > d/2. We denote by $\mathcal{N}_p(R^d)$ the space of counting measures in $\mathcal{M}_p(R^d)$, that is, the integer-valued measures finite on bounded Borel sets.

On $C_p(\mathbb{R}^d)$, p > 0, we define the norm

$$\Box \phi \Box_p = \|\phi/\phi_p\|_{\infty}, \qquad \phi \in C_p(\mathbb{R}^d),$$

making $C_p(R^d)$ a Banach space, and we consider the dual $C_p'(R^d)$ with norm

$$\Box \xi \Box_{-p} = \sup \left\{ |\langle \xi, \phi \rangle| \colon \phi \in C_p(R^d), \Box \phi \Box_p \le 1 \right\}, \qquad \xi \in C_p'(R^d).$$

Then $\mathcal{M}_p(R^d)$ can be regarded as a subspace of $C_p'(R^d)$. It is easy to verify that $\Box \mu \Box_{-p} = \langle \mu, \phi_p \rangle$ for all $\mu \in \mathcal{M}_p(R^d)$. The topology induced on $\mathcal{M}_p(R^d)$ by the norm $\Box \cdot \Box_{-p}$ is stronger than the *p*-vague topology.

Occasionally, we will use some of the function space symbols introduced above (and below) for the corresponding spaces of complex-valued functions.

Let $\{S_t^{\alpha}, t \geq 0\}$ denote the semigroup determined by the spherically symmetric stable process on R^d with exponent α , $0 < \alpha \leq 2$ [see, e.g., Mijnheer (1975)]. By definition this is a time-homogeneous strong Markov process whose transition probability has a density $p_t^{\alpha}(x, y) \equiv p_t^{\alpha}(y - x)$ with characteristic function given by

$$\int_{\mathbb{R}^d} e^{ix \cdot y} p_t^{\alpha}(x) \ dx = \exp\{-t|y|^{\alpha}\}, \qquad y \in \mathbb{R}^d$$

(denotes the inner product in \mathbb{R}^d). The case $\alpha=2$ corresponds to the Wiener process with variance parameter 2. We will suppress the α in the notation of the semigroup and the transition density. We have

For $\mu \in \mathcal{M}_p(\mathbb{R}^d)$ with $p \geq 0$, we also define the Borel measure

$$S_t \mu(A) = \int_{R^d} \left(\int_A p_t(y - x) \, dy \right) \mu(dx), \qquad t > 0,$$

$$S_0 \mu = \mu.$$

Note that $\langle S_t \mu, \phi \rangle = \langle \mu, S_t \phi \rangle$ if $\phi \in C_p(\mathbb{R}^d)$ and $\mu \in \mathcal{M}_p(\mathbb{R}^d)$. Throughout the remainder of this article, p is a fixed number satisfying

$$(3.1a) p > d/2 in all cases,$$

(3.1b)
$$p < (d + \alpha)/2$$
 in the case $\alpha < 2$.

This condition is imposed in order to guarantee that Lebesgue measure belongs to $\mathcal{M}_p(R^d)$ and that S_t maps $\mathcal{M}_p(R^d)$ into itself. The appropriate technical results given below are based on results of Iscoe (1986). [Note, however, that Iscoe defines $\phi_p(x) = (1 + |x|^p)^{-1}$ so that his results translate to our setting with 2p instead of p.]

(3.2) LEMMA. If p satisfies (3.1), then S_t is a bounded linear operator from $(C_p(R^d), \Box \cdot \Box_p)$ into itself, and from $(\mathscr{M}_p(R^d), \Box \cdot \Box_{-p})$ into itself.

PROOF. Since $\phi \in C_p(\mathbb{R}^d)$ is uniformly continuous, then $S_t \phi$ is continuous and bounded. We have

$$|S_t\phi(x)| = |S_t(\phi_p\phi/\phi_p)(x)| \le \Box\phi\Box_p S_t\phi_p(x).$$

By Iscoe (1986), Corollary 2.4, $S_s\phi_p \leq C_t\phi_p$ for $s \leq t$ and some constant C_t . Hence $S_t\phi \in C_p(R^d)$ and $\Box S_t\phi\Box_p \leq C_t\Box\phi\Box_p$.

For $\mu \in \mathcal{M}_p(\mathbb{R}^d)$, by duality and the previous result,

$$\begin{split} \Box S_t \mu \Box_{-p} &= \sup \left\{ |\langle \mu, S_t \phi \rangle| \colon \phi \in C_p(\mathbb{R}^d), \Box \phi \Box_p \le 1 \right\} \\ &\leq \sup \left\{ \Box \mu \Box_{-p} \Box S_t \phi \Box_p \colon \Box \phi \Box_p \le 1 \right\} \le C_t \Box \mu \Box_{-p}. \end{split}$$

(3.3) Lemma. If p satisfies (3.1), and $\phi \in C_p(R^d)$ is such that $\lim_{|x| \to \infty} \phi(x)/\phi_p(x)$ exists $[e.g., \phi \in C_c(R^d)_+ \cup \{\phi_p\}]$, then $t \to S_t \phi$ is a continuous curve in $(C_p(R^d), \Box \cdot \Box_p)$. Also, for $\mu \in \mathcal{M}_p(R^d)$, $t \to S_t \mu$ is p-vaguely continuous.

PROOF. Let $0 \le s < t \le T$. By the proof of (3.2) and Iscoe (1986), Corollary 2.4,

$$\Box S_t \phi - S_s \phi \Box_p = \Box S_s (S_{t-s} \phi - \phi) \Box_p \leq C_T \Box S_{t-s} \phi - \phi \Box_p$$

for some constant C_T , and $\Box S_{t-s}\phi - \phi\Box_p \to 0$ by Iscoe (1986), Lemma 2.4. By the first part of the lemma and dominated convergence,

$$\langle S_t \mu - S_s \mu, \phi \rangle = \langle \mu, S_t \phi - S_s \phi \rangle \to 0$$
 as $t - s \to 0$,

for
$$\phi \in C_c(\mathbb{R}^d)_+ \cup \{\phi_p\}$$
. \square

The infinitesimal generator of $\{S_t, t \geq 0\}$ is given by the fractional power of the Laplacian, $\Delta_{\alpha} \equiv -(-\Delta)^{\alpha/2}$ [see Pazy (1983)].

Given a topological space M, we denote by $D_M \equiv D(R_+, M)$ the space of functions from R_+ into M which are right-continuous and have left limits. D_M is equipped with a Skorohod-type topology when M is a metric space [see Ethier and Kurtz (1986)].

We now return to the counting process N described in the previous section. Following the method of Iscoe (1986), this process can be extended to include initial measures $\mu \in \mathcal{N}_p(R^d)$ and with sample paths in $D_{\mathcal{M}_p(R^d)}$ [see the proof of Lemma (5.2.3)]. Then N is a time-homogeneous Markov process with characteristic functional

$$F(t,\phi,\mu) = E\left[\exp\{-i\langle N_t,\phi\rangle\}|N_0=\mu\right], \qquad t \geq 0, \ \phi \in C_p(R^d), \ \mu \in \mathcal{N}_p(R^d),$$
 which has the following properties:

(3.4) (multiplicative property)

$$F(t, \phi, \mu_1 + \mu_2) = F(t, \phi, \mu_1)F(t, \phi, \mu_2), \qquad \mu_1, \mu_2 \in N_p(\mathbb{R}^d),$$

(3.5) (branching property) $v(t) \equiv \{v(t, x) := F(t, \phi, \delta_x), x \in \mathbb{R}^d\}$ satisfies the integral equation

$$v(t) = e^{-Vt} S_t e^{-i\phi} + \int_0^t [S_r g^h(v(t-r))] V e^{-Vr} dr, \qquad t \ge 0$$

Property (3.4) is a consequence of the independence of the systems generated by different initial particles. Property (3.5) is a nonlinear renewal-type equation which arises as follows: If we start one particle at location x, this particle moves according to the stable law and does not branch before time t (first term on the right-hand side), or it splits at time r at rate V according to g^h and all its offspring evolve independently after birth in the same way (second term).

3.2. The limit fluctuation process Y. In the remainder of this section we discuss some technical points concerning the Schwartz space of tempered distributions, $\mathcal{S}'(R^d)$, which serves as the state space for the limiting fluctuation process Y. Further information on this material can be found in Gelfand and Vilenkin (1964), Itô (1984) and Mitoma (1983).

Let $\mathcal{S}(R^d)$ denote the space of infinitely differentiable functions which are rapidly decreasing at infinity together with all their derivatives. $\mathcal{S}(R^d)$ is topologized by either of the two following sequences of increasing norms:

$$\begin{split} \|\phi\|_n &= \left(\sum_{|k|=0}^n \int_{R^d} (1+|x|^2)^n |D^k \phi(x)|^2 dx\right)^{1/2}, \\ \|\|\phi\|\|_n &= \max_{0 \le |k| \le n} \sup_{x \in R^d} (1+|x|^2)^n |D^k \phi(x)|, \qquad \phi \in \mathscr{S}(R^d), \end{split}$$

 $n=0,1,2,\ldots$, where $k=(k_1,\ldots,k_d)$, $|k|=k_1+\cdots+k_d$, $D^k=\partial^{|k|}/\partial x_1^{k_1}\ldots\partial x_d^{k_d}$. The $\|\cdot\|_n$ are Hilbert norms, and for each $n\geq 1$ there exist constants c_n and $l_n>0$ such that $\|\|\phi\|\|_n\leq c_n\|\|\phi\|\|_{n+l_n}$ and $\|\phi\|\|_n\leq c_n\|\|\phi\|\|_{n+l_n}$, $\phi\in \mathscr{S}(R^d)$. We will denote by $S_n(R^d)$ the completion of $\mathscr{S}(R^d)$ with respect to $\|\cdot\|_n$. Then $\mathscr{S}(R^d)=\bigcap_{n=0}^\infty \mathscr{S}_n(R^d)$ is a countably Hilbert nuclear space; this fact makes possible the application of the results of Hida (1980) and Mitoma (1983). Note that $\mathscr{S}(R^d)\subset C_{p,0}(R^d)$ for all p>0.

The topological dual of $\mathscr{S}(R^d)$ will be designated by $\mathscr{S}'(R^d)$; this is the space of tempered distributions. We have $\mathscr{S}'(R^d) = \bigcup_{n=0}^{\infty} \mathscr{S}'_n(R^d)$, where $\mathscr{S}'_n(R^d)$ is the dual of $\mathscr{S}_n(R^d)$, with the dual-norm written $\|\cdot\|_{-n}$. We will denote by $\langle \cdot, \cdot \rangle$ the duality on $(\mathscr{S}'(R^d), \mathscr{S}(R^d))$ and on $(\mathscr{S}'_n(R^d), \mathscr{S}_n(R^d))$, as well as other dualities.

If $\mathscr{S}'(R^d)_+$ designates the space of nonnegative members of $\mathscr{S}'(R^d)$, then $\mathscr{S}'(R^d)_+ = \bigcup_{p \geq 0} \mathscr{M}_p(R^d)$. Moreover, for each $p \geq 0$ there exists $n \geq 1$ such that $\mathscr{M}_p(R^d) \subset \mathscr{S}'_n(R^d)$; the topology induced on $\mathscr{M}_p(R^d)$ by $\mathscr{S}'_n(R^d)$ is weaker than the p-vague topology, and still weaker is the one induced by $\mathscr{S}'(R^d)$ [since $\mathscr{S}'(R^d)$ is not the strict inductive limit of the $\mathscr{S}'_n(R^d)$].

The space $D_{\mathcal{S}'(\mathbb{R}^d)}$ is equipped with a Skorohod-type topology [see Mitoma (1983)].

(3.6) Remarks. 1. As is now common with models of the present type, an appropriate setting for the fluctuation limit theorem is the nuclear triple

$$\mathscr{S}(R^d) \subset L^2(R^d) \subset \mathscr{S}'(R^d).$$

Since the Brownian motion semigroup and the Laplacian map $\mathcal{S}(R^d)$ into itself, the latter is therefore a natural space of test functions for our model when the particles undergo Brownian motion [e.g., Holley and Stroock (1978)]. However, a technical difficulty arises from the fact that, contrary to the Brownian motion case, the semigroup and the infinitesimal generator of the symmetric stable process with exponent $\alpha < 2$ do not map $\mathcal{S}(R^d)$ into itself. To verify this, note that (3.2) implies that $S_t(\mathcal{S}(R^d)) \subset C_{p,0}(R^d)$ for all p which satisfy (3.1). Moreover, $S_t\phi$ is infinitely differentiable for all $\phi \in \mathcal{S}(R^d)$, since the stable density $p_t(x)$ is infinitely differentiable in x. Nevertheless $S_t(\mathcal{S}(R^d))$ is not contained in $\mathcal{S}(R^d)$ for $\alpha < 2$ because in this case fast decay at infinity fails due to the long tails of the stable law. Moreover, the domain of Δ_{α} contains $\mathcal{S}(R^d)$ and $\Delta_{\alpha}(\mathcal{S}(R^d)) \subset C_{p,0}(R^d)$ but $\Delta_{\alpha}(\mathcal{S}(R^d)) \not\subset \mathcal{S}(R^d)$ again due to the long tails of the stable law [cf. Dawson and Gorostiza (1988)]. Since S_t and Δ_{α} do not map $\mathcal{S}(R^d)$ into itself for $\alpha < 2$, extra work is needed in order to use $\mathcal{S}(R^d)$ as a space of test functions. We will come back to this point in Remark (4.7).

- 2. An alternative approach would be to define an ad hoc nuclear space F of test functions which is invariant under S_t and Δ_{α} . Such a space can be constructed but turns out to be smaller than $\mathcal{S}(R^d)$. Consequently, F' is larger than $\mathcal{S}'(R^d)$ thus yielding a weaker result than the one obtained by working with $\mathcal{S}'(R^d)$.
- **4. Statement of the main results.** In order to set up the scaling limit we introduce a parameter ε , $0 < \varepsilon < 1$, and a corresponding family of processes $N^{\varepsilon} \equiv \{N_t, t \geq 0\}$. For fixed $h \in H$ we denote by P_{ε}^h the distribution of N^{ε} on $D_{\mathscr{L}'(R^d)}$. Then the distribution of N^{ε} in the random medium is given by

$$(4.1) P_{\varepsilon}^{m}(\cdot) = \int_{\mathcal{H}} P_{\varepsilon}^{h}(\cdot) m(dh);$$

and the expectation with respect to P_{ε}^{m} is given by

(4.2)
$$E(\cdot) = \int_{H} E_{h}(\cdot) m(dh),$$

where E_h denotes the expectation with respect to P_{ε}^h [see Dellacherie and Meyer (1978)].

We will assume that as $\varepsilon \to 0$ the scaled initial random measure $\varepsilon^d N_0^\varepsilon(\varepsilon^{-1} \cdot)$ converges weakly under P_ε^m to a nonrandom measure $\Lambda_0 \in \mathscr{M}_p(R^d)$. An example of this is $N_0^\varepsilon = \text{Poisson}$ particle system with intensity measure $\varepsilon^{-d} \Lambda_0(\varepsilon \cdot)$, with $\Lambda_0 \in \mathscr{M}_p(R^d)$. The index ε is introduced only in connection with this assumption on N_0^ε ; the ε has no bearing on the transition probability of N^ε ; hence (3.4) and (3.5) hold under each P_ε^h .

Given N^{ε} we define the $C_p'(R^d)$ -valued process $Y^{\varepsilon} \equiv \{Y_t, t \geq 0\}$ by

$$(4.3) \langle Y_t^{\varepsilon}, \phi \rangle = \varepsilon^{-k} [\langle X_t^{\varepsilon}, \phi \rangle - \langle \Lambda_t, \phi \rangle], t \geq 0, \phi \in C_p(\mathbb{R}^d),$$

where

$$\begin{split} \langle X_t^\varepsilon, \phi \rangle &= \varepsilon^d \langle N_{\varepsilon^- \alpha_t}^\varepsilon, \phi(\varepsilon \cdot) \rangle, \\ (4.4) \qquad & \Lambda_t = S_t \Lambda_0 \quad \text{and} \quad k = \frac{(d\beta - \alpha)}{(1 + \beta)}. \end{split}$$

 Y^{ε} is the *large scale fluctuation process* around the macroscopic deterministic flow Λ . The distribution of Y^{ε} in $D_{\mathscr{S}'(R^d)}$ is determined by P^h_{ε} for the varying medium h, and by P^m_{ε} for the random medium. Note that the space-time scaling involved in (4.4) preserves the distribution of the particle motion, and that d > 2k because $\beta \leq 1$ and $\alpha > 0$.

We will denote weak convergence by \Rightarrow .

Let us recall our assumptions: $d > \alpha/\beta$, p satisfies (3.1a) and (3.1b), N_0^{ε} may depend on the random medium and takes values in $\mathcal{M}_p(R^d)$ and $\varepsilon^d N_0^{\varepsilon}(\varepsilon^{-1} \cdot) \Rightarrow \Lambda_0$ under P_{ε}^m , where $\Lambda_0 \in \mathcal{M}_p(R^d)$ is deterministic. In addition, we will assume that $E\langle N_0^{\varepsilon}, \phi_p \rangle < \infty$ and $E\langle X_0^{\varepsilon}, \phi^{\varepsilon} \rangle \to 0$ as $\varepsilon \to 0$ if $\phi^{\varepsilon} \to 0$ in $C_p(R^d)$, where E is given by (4.2) [see, however, statement 3 in Remark (5.3.8)]. Note that Y_0^{ε} is a $C_p'(R^d)$ -valued random variable. We will now formulate our main result.

(4.5) Theorem. (a) If $Y_0^{\varepsilon} \Rightarrow Y_0$ under P_{ε}^m as $\varepsilon \to 0$, where Y_0 is a $C_p'(R^d)$ -valued random variable, and $\sup_{\varepsilon} E \Box Y_0^{\varepsilon} \Box_{-p}^{1+\beta} < \infty$, then $Y^{\varepsilon} \Rightarrow Y$ under P_{ε}^m in $D_{\mathcal{S}'(R^d)}$ as $\varepsilon \to 0$, where $Y \equiv \{Y_t, t \geq 0\}$ is an $\mathcal{S}'(R^d)$ -valued time-homogeneous Markov process whose transition characteristic functional is given by

$$\begin{split} E\left[\exp\{-i\langle Y_t, \phi \rangle\} | Y_s \right] \\ &\equiv E\left[\exp\{-i\langle Y_t, \phi \rangle\} | \langle Y_s, \psi \rangle, \psi \in \mathscr{S}(R^d) \right] \\ &= \exp\left\{-i\langle Y_s, S_{t-s}\phi \rangle + \left\langle \Lambda_s, V \overline{h} \int_s^t S_{t-r} (iS_{r-s}\phi)^{1+\beta} dr \right\rangle \right\}, \\ &0 \leq s \leq t, \phi \in \mathscr{S}(R^d), \end{split}$$

where $\overline{h} = \int_H h(0)m(dh)$ is the mean of the random medium.

- (b) If the assumptions hold with respect to P_{ε}^h for m-almost all h (in particular if N_0^{ε} is independent of the random medium), then the limit result in (a) is valid for each fixed varying medium h in a set of m-measure 1, the limit process Y being the same as in (a).
- (4.7) Remark. The notation $\langle Y_s, S_{t-s}\phi \rangle$ in (4.6) needs justification because $S_{t-s}\phi$ lies in $C_{p,0}(R^d)$ but not necessarily in $\mathscr{S}(R^d)$. It can be shown that $\langle Y_s, \cdot \rangle$ has a continuous linear extension to $C_{p,0}(R^d)$, and therefore $\phi \to \langle Y_s, S_{t-s}\phi \rangle$ defines a linear random functional on $\mathscr{S}(R^d)$ [see Lemma (3.2) and Remark (3.6)]. Then, by the regularization theorem [e.g., Itô (1984)], this linear random functional has a regular version, that is, a version which is an $\mathscr{S}'(R^d)$ -valued

random variable. We denote the regular version also by $\langle Y_s, S_{t-s}\phi \rangle$, $\phi \in \mathcal{S}(R^d)$, and this is what is meant in (4.6). It should be noted that the previous argument does not allow us to assert that Y is a $C'_{p,0}(R^d)$ -valued process because the regularization theorem does not hold on $C_{p,0}(R^d)$. This and related questions involve other techniques, and they are discussed in detail in Dawson and Gorostiza (1988).

The fluctuation limit theorem, more precisely Proposition (5.3.1) below, also yields the hydrodynamic behavior of the system, that is, the law of large numbers, namely:

(4.8) COROLLARY. $\langle X_t^{\varepsilon}, \phi \rangle \Rightarrow \langle \Lambda_t, \phi \rangle$ under P_{ε}^m as $\varepsilon \to 0$ for each $\phi \in C_p(\mathbb{R}^d)$, and a corresponding a.s. result for condition (b) above.

In the next theorem we will give some properties of the limit fluctuation process Y.

(4.9) Theorem. (i) Y satisfies the generalized Langevin equation

$$dY_t = \Delta_{\alpha} Y_t dt + dZ_t, \qquad t \ge 0,$$

where $Z \equiv \{Z_t, t \geq 0\}$ is an $\mathcal{S}'(R^d)$ -valued process with independent increments such that

$$egin{aligned} E \expig\{ -i(\langle Z_t, \phi
angle - \langle Z_s, \phi
angle) ig\} \ &= \expig\{ V ar{h} \int_s^t \langle \Lambda_r, (i\phi)^{1+eta}
angle \, dr ig\}, \qquad 0 \leq s \leq t, \, \phi \in \mathscr{S}(R^d). \end{aligned}$$

(ii) For $t \geq 0$ and $\phi \in \mathcal{S}(R^d)_+$, $\langle Y_t, \phi \rangle - \langle Y_o, S_t \phi \rangle$ has a distribution with characteristic functional

$$\begin{split} E \exp \big\{ iu(\langle Y_t, \phi \rangle - \langle Y_0, S_t \phi \rangle) \big\} \\ &= \exp \bigg\{ -c(\phi) |u|^{1+\beta} \bigg[1 - i \operatorname{sgn} u \tan \frac{\pi(1+\beta)}{2} \bigg] \bigg\}, \qquad u \in R, \end{split}$$

that is, a Gaussian distribution if $\beta = 1$, and if $\beta < 1$ it is a stable distribution with index $1 + \beta$, and anisotropy parameter -1 [see, e.g., Breiman (1968), page 204]. The normalization constant $c(\phi)$ is given by

$$c(\phi) = -\cos\left(\frac{\pi(1+\beta)}{2}\right) \left\langle \Lambda_0, V\bar{h} \int_0^t S_{t-r}(S_r \phi)^{1+\beta} dr \right\rangle.$$

(iii) $Y_t - S_t Y_0$ is spatially homogeneous for each $t \ge 0$ if and only if $\Lambda_0(dx) = \lambda dx$, $\lambda > 0$. In this case, if $Y_0 = 0$ and $\beta = 1$ (the Gaussian case), the spectral measure of Y_t is given by

$$V\overline{h}\lambda|x|^{-\alpha}(1-\exp\{-2t|x|^{\alpha}\}) dx.$$

(iv) There exists $n \ge 1$ such that $Y_t - S_t Y_0$ takes values in $S'_n(\mathbb{R}^d)$ for all $t \ge 0$.

- (4.10) Remarks. 1. Under condition (b) in (4.5) the tightness proof is basically the same as in case (a), but simpler. We will do it only for (a). The condition $\sup_{\epsilon} E \Box Y_0^{\epsilon} \Box_{-p}^{1+\beta} < \infty$ can be replaced by the slightly weaker one: $\sup_{\epsilon} E \Box Y_0^{\epsilon} \Box_{-p}^{1+\theta} < \infty$ for some $0 < \theta < \beta$ but this complicates the calculations.
- 2. Under each of the two cases in (4.5), if the random medium is only invariant (not necessarily ergodic) with respect to the translation group, then the results hold with \bar{h} replaced by $E[h|\mathcal{I}]$, where \mathcal{I} is the σ -field of invariant sets.
- 3. In view of (4.9)(i) and (4.9)(ii) we refer to Y as a generalized stable Ornstein-Uhlenbeck process. More precisely, $Y-SY_0$ is such a process. In the case $\beta=1$, $Y-SY_0$ is a special case of a generalized Gaussian Ornstein-Uhlenbeck process with continuous paths [see Holley and Stroock (1978) and Dawson (1981)]; Bojdecki and Gorostiza (1986) discuss a more general class of such processes; in this case the continuity can be shown by an extension of the Dudley-Fernique theorem [cf. Mitoma (1981)]. On the other hand, for $\beta < 1$, Y has discontinuous paths. Hence the properties of Y being Gaussian or strictly stable, and correspondingly continuous or discontinuous, depend on whether the branching law has finite second moment or not ($\beta=1$ or $\beta<1$, respectively), and these properties are independent of whether the particle motion is Brownian motion or a strictly stable process ($\alpha=2$ or $\alpha<2$). It should also be noticed that (4.6) shows that the conditional distribution of the stable process $Y-SY_0$ is also stable; this is generally not true for symmetric stable non-Gaussian processes [Adler, Cambanis and Samorodnitsky (1987)].
- 4. The asymmetric behavior of $Y SY_0$ in the case $\beta < 1$ occurs because the branching law is the distribution of a nonnegative random variable and belongs to the domain of attraction of an asymmetric stable law.
- 5. If N_0^{ε} is a Poisson system with intensity measure $\varepsilon^{-d}\Lambda_0(\varepsilon \cdot)$, then $\varepsilon^{-d/2}(\varepsilon^dN_0^{\varepsilon}(\varepsilon^{-1}\cdot)-\Lambda_0)$ converges weakly to a Gaussian white noise on R^d determined by Λ_0 , and therefore $Y_0=0$ because d>2k. In this case the normalization which yields a nondegenerate limit at times t>0 is too strong to yield limit fluctuations at time t=0.
- 6. In the Gaussian case ($\beta=1$) and Brownian motion ($\alpha=2$), the Langevin equation in (4.9)(i) can be obtained by applying Theorem 3.6 in Bojdecki and Gorostiza (1986). For $\alpha<2$ and $\beta=1$ the equation is formally the same, just replacing Δ by Δ_{α} . However, for $\alpha<2$ (and $\beta\leq1$) the Langevin equation cannot be interpreted in the usual way, that is,

$$\langle Y_t, \phi \rangle = \langle Y_0, \phi \rangle + \int_0^t \langle Y_s, \Delta_{\alpha} \phi \rangle ds + \langle Z_t, \phi \rangle, \qquad t \geq 0, \phi \in \mathscr{S}(R^d),$$

because Δ_{α} does not map $\mathcal{S}(R^d)$ into itself. A generalized interpretation of the equation is necessary. This turned out to be a fairly technical problem, and we preferred to discuss it separately [Dawson and Gorostiza (1988)].

5. Proofs.

5.1. A nonlinear integral equation with random coefficients. The proof of convergence is based on properties of solutions of an integral equation. In this section we will derive these properties.

Fix T > 0, $\phi \in C_p(\mathbb{R}^d)$, L > 0 and $0 < \lambda \le L$, and assume that for each h in a subset H_0 of H and $0 < \varepsilon < 1$ we have a bounded complex-valued function $u_{\varepsilon}(t) \equiv u_{\varepsilon}(t,x) \equiv u_{\varepsilon}(t,\phi\lambda,x) \equiv u_{\varepsilon}(t,\phi\lambda)$, $0 \le t \le T$, $x \in \mathbb{R}^d$ (we omit the dependence on h in the notation) which satisfies the integral equation

$$\begin{aligned} u_{\varepsilon}(t) &= S_t \big[\varepsilon^{k-d} \big(1 - \exp \big\{ -i \varepsilon^{d-k} \phi \lambda \big\} \big) \big] \\ &- \varepsilon^k V \! \int_0^t \! S_{t-s} \big[h(\varepsilon^{-1} \cdot) u_{\varepsilon}^{1+\beta}(s) \big] \, ds, \qquad 0 \le t \le T. \end{aligned}$$

(5.1.2) Lemma. There exist $0 < \varepsilon_{T,L} < 1$ and $0 < K < \infty$ such that $\sup\{|u_{\varepsilon}(t,x)|: t \leq T, x \in R^d, \varepsilon < \varepsilon_{T,L}, h \in H_0\} \leq \text{const.} \|\phi\|_{\infty} \lambda \equiv K\lambda.$

PROOF. Let $a_{\epsilon}(t) = \sup\{|u_{\epsilon}(t, x)| : x \in \mathbb{R}^d\}$. Then, since S_t is a contraction and $0 \le h \le 1$, (5.1.1) implies

$$0 \le a_{\varepsilon}(t) \le \|\phi\|_{\infty} \lambda + \varepsilon^k V \int_0^t a_{\varepsilon}^{1+\beta}(s) \, ds, \qquad 0 \le t \le T.$$

The equation

$$b_{arepsilon}(t) = \|\phi\|_{\infty} \lambda + arepsilon^k V \! \int_0^t \! b_{arepsilon}^{1+eta}\! \left(s
ight) d\!s, \qquad 0 \leq t \leq T,$$

has unique solution $b_{\varepsilon}(t) = \|\phi\|_{\infty} \lambda (1 - \varepsilon^k V \beta \|\phi\|_{\infty}^{\beta} \lambda^{\beta} t)^{-1/\beta}$ for all $0 < \lambda \le L$ provided that

(5.1.3)
$$\varepsilon < 1/(V\beta T \|\phi\|_{\infty}^{\beta} L^{\beta})^{1/k}.$$

Clearly, there exists $\varepsilon_{T,L} > 0$ such that $b_{\varepsilon}(t) \leq \text{const.} \|\phi\|_{\infty} \lambda$ for $0 \leq t \leq T$ and $\varepsilon < \varepsilon_{T,L}$.

By a comparison theorem [Miller (1971), page 121], $a_{\epsilon}(t) \leq b_{\epsilon}(t)$ for $0 \leq t \leq T$; hence the assertion follows. \Box

(5.1.4) LEMMA.

$$\sup\{|u_{\varepsilon}(t)| \colon \varepsilon < \varepsilon_{T,L}, \, h \in H_0\} \leq \text{const. } S_t |\phi|\lambda, \qquad 0 \leq t \leq T.$$

PROOF. By (5.1.5) and (5.1.2),

$$\sup_{k} |u_{\varepsilon}(t)| \leq S_{t} |\phi| \lambda + \varepsilon^{k} V K^{\beta} \lambda^{\beta} \int_{0}^{t} S_{t-s} \sup_{k} |u_{\varepsilon}(s)| \, ds, \qquad 0 \leq t \leq T.$$

The comparison theorem used in the previous proof does not apply in this case; henceforth we assume $\varepsilon_{T,L}$ is small enough so that $c \equiv \varepsilon_{T,L}^k V K^{\beta} L^{\beta} T < 1$ where $K = \text{const.} \|\phi\|_{\infty}$ [this is consistent with (5.1.3)]; then by iteration of the inequality it can be shown that

$$\sup_{t} |u_{\varepsilon}(t)| \leq (1-c)^{-1} S_{t} |\phi| \lambda, \qquad 0 \leq t \leq T.$$

(5.1.5) Lemma. There exist positive constants K_1 and K_2 such that for all $t \leq T$ and $\varepsilon < \varepsilon_{T,L}$,

$$\sup_{h} |u_{\varepsilon}(t) - iS_{t}\phi\lambda| \leq K_{1}\varepsilon^{k}S_{t}|\phi|\lambda^{1+\beta} \leq K_{2}\varepsilon^{k}\lambda.$$

Proof. Since

$$|\exp\left\{-i\varepsilon^{d-k}\phi\lambda\right\}-1+i\varepsilon^{d-k}\phi\lambda|\leq \varepsilon^{2(d-k)}|\phi|^2\lambda^2\leq \mathrm{const.}\ \varepsilon^{2(d-k)}|\phi|\lambda^{1+\beta},$$

then

(a)
$$|S_t \{ \varepsilon^{k-d} (1 - e^{-i\varepsilon^{d-k}\phi\lambda}) \} - iS_t \phi\lambda | \le \text{const. } \varepsilon^{d-k} S_t |\phi| \lambda^{1+\beta}.$$

Then, by (5.1.1), (5.1.2), (5.1.4) and (a),

$$\begin{split} |u_{\varepsilon}(t)-iS_{t}\phi\lambda| &\leq |S_{t}\big\{\varepsilon^{k-d}(1-e^{-i\varepsilon^{d-k}\phi\lambda})\big\} - iS_{t}\phi\lambda| + \varepsilon^{k}V\int_{0}^{t}S_{t-s}|u_{\varepsilon}(s)|^{1+\beta}\,ds \\ &\leq \mathrm{const.}\,\,\varepsilon^{d-k}S_{t}|\phi|\lambda^{1+\beta} + \mathrm{const.}\,\,\varepsilon^{k}VK^{\beta}\lambda^{\beta}\int_{0}^{t}S_{t-s}S_{s}|\phi|\lambda\,ds \\ &\leq \mathrm{const.}\,\,\varepsilon^{d-k}S_{t}|\phi|\lambda^{1+\beta} + \mathrm{const.}\,\,\varepsilon^{k}S_{t}|\phi|\lambda^{1+\beta} \leq \mathrm{const.}\,\,\varepsilon^{k}S_{t}|\phi|\lambda^{1+\beta}, \end{split}$$

the last inequality because d - k > k. This proves the first asserted inequality; the second one is obvious. \square

(5.1.6) Proposition. There exists $\gamma > 1$ such that for all $t \leq T$ and $\varepsilon < \varepsilon_{T, L}$,

$$\begin{split} \sup_{h} \left| u_{\varepsilon}(t) - i S_{t} \phi \lambda + \varepsilon^{k} V \int_{0}^{t} S_{t-s} \left[h(\varepsilon^{-1} \cdot) (i S_{s} \phi \lambda)^{1+\beta} \right] ds \right| \\ \leq \text{const. } \varepsilon^{\gamma k} S_{t} |\phi| \lambda^{1+\beta}. \end{split}$$

PROOF. Let $A = \{(s, y) \in [0, T] \times R^d: |S_s\phi(y)| \le 2K_2\varepsilon^k\}$, where K_2 is the constant in (5.1.5). Then, by (5.1.5), on A,

$$|u_{\varepsilon}(s, y)| \leq |u_{\varepsilon}(s, y) - iS_s\phi(y)\lambda| + |S_s\phi(y)\lambda| \leq K_2\varepsilon^k\lambda + 2K_2\varepsilon^k\lambda \leq \text{const. } \varepsilon^k\lambda.$$
 Hence, by (5.1.4), on A ,

(a)
$$\begin{aligned} |u_{\varepsilon}(s, y)|^{1+\beta} &= |u_{\varepsilon}(s, y)| \, |u_{\varepsilon}(s, y)|^{\beta} \\ &\leq \mathrm{const.} \, S_{s} |\phi|(y) \lambda \varepsilon^{k\beta} \lambda^{\beta} \leq \mathrm{const.} \, \varepsilon^{k\beta} S_{s} |\phi|(y) \lambda^{1+\beta}. \end{aligned}$$

Also, on A,

$$\begin{split} |S_s \phi(y) \lambda|^{1+\beta} &\leq |S_s \phi(y)| \, |S_s \phi(y)|^{\beta} \lambda^{1+\beta} \\ &\leq \operatorname{const.} \varepsilon^{k\beta} S_s |\phi|(y) \lambda^{1+\beta}. \end{split}$$

Using (a) and (b), we have

$$\begin{split} \varepsilon^k V \! \int_0^t \! ds \int_{R^d} \! dy \, \mathbf{1}_A(s, y) p_{t-s}(\cdot, y) |u_\varepsilon^{1+\beta}(s, y) - \left(i S_s \phi(y) \lambda\right)^{1+\beta} | \\ & \leq \operatorname{const.} \varepsilon^{k(1+\beta)} \lambda^{1+\beta} \int_0^t \! ds \int_{R^d} \! dy p_{t-s}(\cdot, y) S_s |\phi|(y), \end{split}$$

so

$$\begin{aligned} (c) \qquad & \varepsilon^{k} V \! \int_{0}^{t} \! ds \! \int_{R^{d}} \! dy \mathbf{1}_{A}(s, y) p_{t-s}(\cdot, y) |u_{\varepsilon}^{1+\beta}(s, y) - \left(i S_{s} \phi(y) \lambda\right)^{1+\beta} | \\ & \leq \operatorname{const.} \varepsilon^{k(1+\beta)} S_{\epsilon} |\phi| \lambda^{1+\beta}. \end{aligned}$$

On the other hand, on A^c , $|iS_s\phi(y)\lambda| > 2K_2\varepsilon^k\lambda$, but by (5.1.5) $|u_{\varepsilon}(s,y) - iS_s\phi(y)\lambda| \le K_2\varepsilon^k\lambda$. Hence on A^c both $u_{\varepsilon}(s,y)$ and $iS_s\phi(y)\lambda$ lie on the same half-plane $\mathrm{Im}(z) > 0$ or $\mathrm{Im}(z) < 0$. On these half-planes the function $z^{1+\beta}$ is Lipschitz-continuous on compact sets. Now, A^c is bounded since $\sup_{s\le T} S_s |\phi|(y) \to 0$ as $|y| \to \infty$, due to (3.2) and Iscoe (1986), Corollary 2.4. Hence A^c has compact closure. Therefore by (5.1.5), on A^c ,

$$|u_{\varepsilon}^{1+\beta}(s, y) - (iS_{s}\phi(y)\lambda)^{1+\beta}| \leq \text{const.}|u_{\varepsilon}(s, y) - iS_{s}\phi(y)\lambda|$$

$$\leq \text{const.} \ \varepsilon^{k}S_{s}|\phi|(y)\lambda^{1+\beta},$$

consequently

$$\begin{split} \varepsilon^k V \! \int_0^t \!\! ds \! \int_{R^d} \!\! dy 1_{A^c}\!(s,\,y) p_{t-s}\!(\cdot\,,\,y) |u_\varepsilon^{1+\beta}\!(s,\,y) - \big(iS_s\!\phi(\,y)\lambda\,\big)^{1+\beta}| \\ & \leq \mathrm{const.} \, \varepsilon^{2k} \lambda^{1+\beta} \int_0^t \!\! ds \! \int_{R^d} \!\! dy p_{t-s}\!(\cdot\,,\,y) S_s |\phi|(\,y), \end{split}$$

SO

(d)
$$\begin{aligned} \varepsilon^k V \! \int_0^t \! ds \! \int_{R^d} \! dy \mathbf{1}_{A^c}(s, y) p_{t-s}(\cdot, y) |u_{\varepsilon}^{1+\beta}(s, y) - \left(i S_s \phi(y) \lambda\right)^{1+\beta} | \\ & \leq \mathrm{const.} \ \varepsilon^{2k} S_t |\phi| \lambda^{1+\beta}. \end{aligned}$$

Finally, using (5.1.1), (5.1.5)(a), (c) and (d), we have

$$\begin{split} \sup_{h} \left| u_{\epsilon}(t) - iS_{t}\phi\lambda + \varepsilon^{k}V \int_{0}^{t} S_{t-s} \left[h(\varepsilon^{-1} \cdot)(iS_{s}\phi\lambda)^{1+\beta} \right] ds \right| \\ & \leq |S_{t} \left\{ \varepsilon^{k-d} (1 - e^{i\varepsilon^{d-k}\phi\lambda}) \right\} - iS_{t}\phi\lambda | \\ & + \varepsilon^{k}V \int_{0}^{t} ds \int_{R^{d}} dy 1_{A}(s, y) p_{t-s}(\cdot, y) |(iS_{s}\phi(y)\lambda)^{1+\beta} - u_{\epsilon}^{1+\beta}(s, y)| \\ & + \varepsilon^{k}V \int_{0}^{t} ds \int_{R^{d}} dy 1_{A^{c}}(s, y) p_{t-s}(\cdot, y) |(iS_{s}\phi(y)\lambda)^{1+\beta} - u_{\epsilon}^{1+\beta}(s, y)| \\ & \leq \operatorname{const.} (\varepsilon^{d-k} + \varepsilon^{k(1+\beta)} + \varepsilon^{2k}) S_{t} |\phi| \lambda^{1+\beta} \\ & \leq \operatorname{const.} \varepsilon^{\gamma k} S_{\epsilon} |\phi| \lambda^{1+\beta}, \end{split}$$

where $\gamma = \min\{(d-k)/k, 1+\beta, 2\}; \ \gamma > 1 \text{ because } d-k > k \text{ and } \beta > 0. \ \Box$

REMARK. The randomness in (5.1.1) is due to h. The only fact which we used about h is that it is bounded by 1, and therefore the previous results hold for all realizations of h in H_0 .

5.2. Properties of the process N. We will derive here some properties we need of the process $N^{\varepsilon} \equiv N \equiv \{N_t, t \geq 0\}$ under P_{ε}^h (h fixed). We denote by E_h the expectation with respect to P_{ε}^h .

(5.2.1) LEMMA. For
$$0 \le s < t$$
, $\phi \in C_n(\mathbb{R}^d)$ and $\lambda \in \mathbb{R}$,

(i)
$$E_h[\exp\{-i\lambda\langle N_t, \phi\rangle\}|N_s] = \exp\{\langle N_s, \log[1-w(t-s)]\rangle\},$$

where $w(t) \equiv w(t, \phi \lambda, x)$ satisfies the equation

$$(ext{ii}) \qquad w(t) = S_t ig[1 - e^{-i\phi\lambda} ig] - V \!\! \int_0^t \!\! S_{t-r} ig[h w^{1+eta}(r) ig] \, dr, \qquad t \geq 0,$$

and w(t, 0, x) = 0. Moreover,

(iii)
$$E_h \exp\{-i\lambda \langle N_t, \phi \rangle\} = E_h \exp\{\langle N_0, \log[1 - w(t)] \rangle\}.$$

PROOF. By (3.3),

(a)
$$v(t) \equiv v(t, \phi \lambda, x) \coloneqq E_h \big[\exp\{-i \langle N_t, \phi \lambda \rangle | N_0 = \delta_x \} \big],$$
$$t > 0, x \in \mathbb{R}^d.$$

satisfies

$$v(t) = e^{-Vt} S_t e^{-i\phi\lambda} + V \int_0^t e^{-V(t-s)} S_{t-s} g^h(v(s)) ds, \qquad t \ge 0.$$

The generator of the semigroup $\{e^{-Vt}S_t, t \geq 0\}$ is $\Delta_{\alpha} - VI$, where I is the identity operator. Hence v also satisfies the following variation-of-constants equation with respect to the semigroup S_t [see, e.g., Pazy (1983)]:

$$v(t) = S_t e^{-i\phi\lambda} + V \int_0^t S_{t-s} [g^h(v(s)) - v(s)] ds, \qquad t \geq 0.$$

By (2.1), $g^h(v(s)) - v(s) = h(1 - v(s))^{1+\beta}$; then setting

(b)
$$w(t) \equiv w(t, \phi \lambda, x) := 1 - v(t), \quad t \ge 0, x \in \mathbb{R}^d,$$

we see that w satisfies (ii) and w(t, 0, x) = 0.

Using the multiplicative property (3.4), the time homogeneity of N, (a) and (b), we have

$$\begin{split} E_h \big[& \exp\{-i\lambda\langle N_t, \phi \rangle\} | N_s \big] \\ &= \exp\{\langle N_s, \log E_h \big[\exp\{-i\lambda\langle N_t, \phi \rangle\} | N_s = \delta. \big] \rangle\} \\ &= \exp\{\langle N_s, \log[1 - w(t-s)] \rangle\}. \end{split}$$

Then (iii) follows by taking expectation. \Box

We have stated the previous lemma as if we already knew that N_t takes values in $\mathcal{M}_p(R^d)$. In fact this will follow from the next results. For the moment we may assume the validity of the lemma for $\phi \in C_c(R^d)$. Moreover, the assumption $E\langle N_0, \phi_p \rangle < \infty$ implies $E_h\langle N_0, \phi_p \rangle < \infty$ for *m*-almost all *h* [see (4.2)]; we may assume for simplicity that this holds for all *h*.

(5.2.2) LEMMA. For
$$0 \le s < t$$
 and $\phi \in C_p(\mathbb{R}^d)$,

(i)
$$E_h[\langle N_t, \phi \rangle | N_s] = \langle N_s, S_{t-s} \phi \rangle$$

and

(ii)
$$E_h\langle N_t, \phi \rangle = E_h\langle N_0, S_t \phi \rangle < \infty.$$

Proof. By (5.2.1)(i),

$$E_{h}[\langle N_{t}, \phi \rangle | N_{s}] = i \partial / \partial \lambda (\exp\{\langle N_{s}, \log[1 - w(t - s)] \rangle\})|_{\lambda = 0}$$
$$= i \langle N_{s}, -\partial / \partial \lambda (w(t - s)) \rangle|_{\lambda = 0},$$

because w(t,0,x)=0, and from (5.2.1)(ii) we find $\partial/\partial\lambda(w(t-s))|_{\lambda=0}=iS_{t-s}\phi$, because $\beta>0$. This yields (i), and (ii) follows from (i) by taking expectation. The results hold for $\phi\in C_p(R^d)$ due to (3.2) and the assumption $E_h\langle N_0,\phi_p\rangle<\infty$. \square

Since $E_h\langle N_t, \phi_p \rangle < \infty$, then $\langle N_t, \phi_p \rangle < \infty$ a.s. Hence for each $t \geq 0$, $E_h N_t$ and N_t take values in $\mathscr{M}_p(R^d)$ a.s. Since $\mathscr{M}_p(R^d) \subset \mathscr{S}_n'(R^d)$ for some $n \geq 1$, then for each $t \geq 0$, N_t and $N_t - E_h N_t$ take values in $\mathscr{S}_n'(R^d)$ a.s. Moreover, we have the following result.

(5.2.3) LEMMA. The processes N and $N - E_h N$ can be realized in $D_{\mathscr{S}'_n(R^d)}$ for some $n \geq 1$, and they have the strong Markov property (under P_s^h).

PROOF. We assume the filtration generated by N is completed and made right-continuous in the usual way. The result for N follows by an application of Theorem 9.4 in Chapter 1 of Blumenthal and Getoor (1968) to the $\mathcal{M}_p(R^d)$ -valued Markov process N. However, since $\mathcal{M}_p(R^d)$ is not locally compact (in the p-vague topology), in order to apply this theorem we embed $\mathcal{M}_p(R^d)$ in the locally compact space $\mathcal{M}_p(\dot{R}^d)$ introduced by Iscoe (1986).

Let $P_t(\mu, d\nu)$ denote the transition probability of the time-homogeneous Markov process N in the medium h. For $\Phi \in C(\mathcal{M}_p(\dot{R}^d))$ we define

$$T_t\Phi(\mu)=\int_{\mathscr{N}_p(R^d)}\Phi(\nu)P_t(\mu,d
u), \qquad t\geq 0,\, \mu\in\mathscr{N}_p(R^d).$$

As in Iscoe (1986) this definition can be extended to $\mu \in \mathcal{N}_p(\dot{R}^d)$, and it can be shown that $T_t\colon C_0(\mathcal{N}_p(\dot{R}^d))\to C_0(\mathcal{N}_p(\dot{R}^d))$, and $T_t\Phi\to\Phi$ as $t\to 0$ uniformly on $\mathcal{N}_p(\dot{R}^d)$ for each $\Phi\in C_0(\mathcal{N}_p(\dot{R}^d))$. Then the quoted theorem applies, so that N can be realized in $D_{\mathcal{N}_p(\dot{R}^d)}$ and has the strong Markov property. As in Iscoe (1986) it can be shown that N is $\mathcal{N}_p(R^d)$ -valued. Then from the results in Section 2, N can also be realized in $D_{\mathcal{S}_p'(R^d)}$ for some $n\geq 1$ and the strong Markov property is maintained in this space.

The result for $N - E_h N$ now follows from the fact that by (5.2.2)(ii) and (3.3) $t \to E_h N_t = S_t E_h N_0$ is p-vaguely continuous. \square

(5.2.4) Lemma. Let $\phi \in C_p(\mathbb{R}^d)$ and assume ϕ is in the domain of Δ_{α} . Then the process

$$\langle N_t, \phi \rangle - \int_0^t \langle N_r, \Delta_{\alpha} \phi \rangle dr, \qquad t \geq 0,$$

is a P_{ε}^{h} -martingale with respect to the filtration $\mathscr{F}_{s} = \sigma(\langle N_{r}, \psi \rangle; r \leq s, \psi \in \mathscr{S}(\mathbb{R}^{d})), s \geq 0.$

PROOF. Integrability follows from the previous results. The integral in t exists because of (5.2.3). Let s < t. Using the Markov property of N, (5.2.2)(i) and $d/dt(S_t) = S_t \Delta_a$, we have

$$\begin{split} E_h \bigg[\langle N_t, \phi \rangle - \int_0^t \langle N_r, \Delta_\alpha \phi \rangle \, dr | \mathscr{F}_s \bigg] \\ &= E_h \big[\langle N_t, \phi \rangle | N_s \big] - \int_0^s \langle N_r, \Delta_\alpha \phi \rangle \, dr - \int_s^t E_h \big[\langle N_r, \Delta_\alpha \phi \rangle | N_s \big] \, dr \\ &= \langle N_s, S_{t-s} \phi \rangle - \int_0^s \langle N_r, \Delta_\alpha \phi \rangle \, dr - \left\langle N_s, \int_s^t S_{r-s} \Delta_\alpha \phi \, dr \right\rangle \\ &= \langle N_s, \phi \rangle - \int_0^s \langle N_r, \Delta_\alpha \phi \rangle \, dr. \end{split}$$

Replacing N by $N - E_h N$ in (5.2.4) also yields a martingale. Note that the assumptions in (5.2.4) hold for $\phi \in \mathcal{S}(\mathbb{R}^d)$.

- 5.3. Convergence of finite-dimensional distributions.
- (5.3.1) Proposition. If $\langle Y_0^{\varepsilon}, \phi \rangle \Rightarrow Y_0(\phi)$ under P_{ε}^m as $\varepsilon \to 0$ for each $\phi \in C_n(\mathbb{R}^d)$, then

$$\left(\langle Y_{t_{t}}^{\varepsilon}, \phi_{1} \rangle, \dots, \langle Y_{t_{s}}^{\varepsilon}, \phi_{n} \rangle\right) \Rightarrow \left(Y_{t_{t}}(\phi_{1}), \dots, Y_{t_{s}}(\phi_{n})\right)$$

under P_{ε}^m as $\varepsilon \to 0$ for all $n \ge 1$, $0 \le t_1 < \cdots < t_n$ and $\phi_1, \ldots, \phi_n \in C_p(\mathbb{R}^d)$, where $Y_{t_k}(\phi_k)$, $k = 0, 1, \ldots, n$, denote the limit real random variables $(t_0 = 0)$.

PROOF. It suffices to prove that

(a)
$$E_{\epsilon} \exp \left\{ -i \sum_{j=0}^{k} \langle Y_{t_{j}}^{\epsilon}, \phi_{j} \rangle \right\} \rightarrow E \exp \left\{ -i \sum_{j=0}^{k} Y_{t_{j}}(\phi_{j}) \right\}$$
 as $\epsilon \rightarrow 0$

for all $k \geq 0$, $0 = t_0 < t_1 < \cdots < t_k$ and $\phi_0, \phi_1, \ldots, \phi_k \in C_p(\mathbb{R}^d)$, where E_{ε} denotes expectation under P_{ε}^m and E expectation under the limit distribution.

For k = 0 (a) holds by hypothesis. Assuming (a) is true for k = n - 1, we will show it holds for k = n.

Using the Markov property of N under P_{ϵ}^{h} , by (5.2.1) with $\lambda = 1$, we have [with the notation of (5.2.1), $N \equiv N^{\epsilon}$]

$$\begin{split} E_h \exp & \left\{ -i \sum_{j=0}^n \left\langle N_{t_j}, \phi_j \right\rangle \right\} \\ &= E_h \Biggl(\exp \left\{ -i \sum_{j=0}^{n-1} \left\langle N_{t_j}, \phi_j \right\rangle \right\} E_h \Bigl[\exp \left\{ -i \left\langle N_{t_n}, \phi_n \right\rangle \right\} | N_{t_{n-1}} \Bigr] \Biggr) \\ &= E_h \Biggl(\exp \Biggl\{ -i \sum_{j=0}^{n-1} \left\langle N_{t_j}, \phi_j \right\rangle + \left\langle N_{t_{n-1}}, \log \bigl[1 - w(t_n - t_{n-1}, \phi_n) \bigr] \right\rangle \Biggr) \Biggr). \end{split}$$

Hence, by (4.4),

$$\begin{split} E_h \exp & \left\{ -i \sum_{j=0}^n \left\langle \varepsilon^{-k} X_{t_j}^\varepsilon, \phi_j \right\rangle \right\} \\ &= E_h \exp \left\{ -i \sum_{j=0}^n \left\langle N_{\varepsilon^{-a} t_j}, \varepsilon^{d-k} \phi_j(\varepsilon \cdot) \right\rangle \right\} \\ &= E_h \exp & \left\{ -i \sum_{j=0}^{n-1} \left\langle N_{\varepsilon^{-a} t_j}, \varepsilon^{d-k} \phi_j(\varepsilon \cdot) \right\rangle \right. \\ &+ & \left. \left\langle N_{\varepsilon^{-a} t_{n-1}}, \log \left[1 - w \left(\varepsilon^{-a} (t_n - t_{n-1}), \varepsilon^{d-k} \phi_n(\varepsilon \cdot) \right) \right] \right\rangle \right\}, \end{split}$$

and therefore

$$\begin{split} E_h \exp & \left\langle -i \sum_{j=0}^n \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right\rangle \\ &= E_h \exp & \left\langle -i \sum_{j=0}^{n-1} \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right. \\ & \left. + \left\langle X_{t_{n-1}}^{\varepsilon}, \varepsilon^{-d} \log \left[1 - w \left(\varepsilon^{-\alpha} (t_n - t_{n-1}), \varepsilon^{d-k} \phi_n (\varepsilon \cdot), \varepsilon^{-1} \cdot \right) \right] \right\rangle \right\rangle. \end{split}$$

Hence, denoting

(b)
$$z(t)\equiv z(t,x)\equiv z(t,\phi_n,x)\coloneqq w\big(\varepsilon^{-\alpha}t,\,\varepsilon^{d-k}\phi_n(\varepsilon\,\cdot\,),\,x\big),$$
 we have

$$\begin{split} E_h \exp & \left\{ -i \sum_{j=0}^n \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right\} \\ (c) & = E_h \exp & \left\{ -i \sum_{j=0}^{n-1} \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right. \\ & \left. + \left\langle X_{t_{n-1}}^{\varepsilon}, \varepsilon^{-d} \log \left[1 - z \left(t_n - t_{n-1}, \varepsilon^{-1} \cdot \right) \right] \right\rangle \right\}, \end{split}$$

where z satisfies, by (5.2.1)(ii),

$$\begin{split} z(t) &= S_{\epsilon^{-\alpha}t} \big[1 - \exp \big\{ -i \varepsilon^{d-k} \phi_n(\varepsilon \, \cdot \,) \big\} \big] - V \! \int_0^{\epsilon^{-\alpha}t} \! S_{\epsilon^{-\alpha}t-s} \big[\, h z^{1+\beta}(\varepsilon^{\alpha}s \,) \big] \, ds \\ & \leq (\mathrm{d}) \qquad = S_{\epsilon^{-\alpha}t} \big[1 - \exp \big\{ -i \varepsilon^{d-k} \phi_n(\varepsilon \, \cdot \,) \big\} \big] \\ & \qquad \qquad - \varepsilon^{-\alpha} V \! \int_0^t \! S_{\epsilon^{-\alpha}(t-r)} \big[\, h z^{1+\beta}(r) \big] \, dr, \qquad 0 \leq t \leq T. \end{split}$$

Now, using the self-similarity of the stable density, that is,

$$p_t(x) = \varepsilon^{-d} p_{\varepsilon^{-\alpha}t}(\varepsilon^{-1}x), \qquad t > 0, x \in \mathbb{R}^d,$$

which is a simple consequence of the stable characteristic function,

$$S_{\epsilon}f(x) = (S_{\epsilon^{-\alpha}\epsilon}f(\epsilon \cdot))(\epsilon^{-1}x);$$

hence by (d), z satisfies

$$\begin{split} z(t,\,\varepsilon^{-1}\,\cdot\,) &= S_t \big[1 - \exp\big\{-i\varepsilon^{d-k}\phi_n\big\}\big] \\ &- \varepsilon^{-\alpha} V \!\int_0^t \!\! S_{t-s} \big[\,h(\,\varepsilon^{-1}\,\cdot\,)z^{1+\beta}(s,\,\varepsilon^{-1}\,\cdot\,)\big]\,ds, \qquad 0 \leq t \leq T. \end{split}$$

Letting

(f)
$$u_{\varepsilon}(t) \equiv u_{\varepsilon}(t, \phi_{n}, x) := \varepsilon^{k-d} z(t, \phi_{n}, \varepsilon^{-1} x),$$

by (c) we have

$$\begin{split} E_h \exp & \left\langle -i \sum_{j=0}^n \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right\rangle \\ (\mathbf{g}) & = E_h \exp & \left\langle -i \sum_{j=0}^{n-1} \left\langle \varepsilon^{-k} X_{t_j}^{\varepsilon}, \phi_j \right\rangle \right. \\ & \left. + \left\langle X_{t_{n-1}}^{\varepsilon}, \varepsilon^{-d} \log \left[1 - \varepsilon^{d-k} u_{\epsilon}(t_n - t_{n-1}) \right] \right\rangle \right\rangle, \end{split}$$

where, observing that $\varepsilon^{-\alpha+k-d+(d-k)(1+\beta)} = \varepsilon^k$, by (e) u_{ε} satisfies

$$(5.3.2) \qquad u_{\varepsilon}(t) = S_{t} \left[\varepsilon^{k-d} \left(1 - \exp\left\{ -i \varepsilon^{d-k} \phi_{n} \right\} \right) \right] \\ - \varepsilon^{k} V \int_{0}^{t} S_{t-s} \left[h(\varepsilon^{-1} \cdot) u_{\varepsilon}^{1+\beta}(s) \right] ds, \qquad 0 \le t \le T.$$

Note that this coincides with (5.1.1) if $\phi_n = \phi$ and $\lambda = 1$, and that (5.2.1)(a) and (5.2.1)(b), (b) and (f) imply the boundedness required of u_{ε} in Section 5.1. By (4.3), (g) and (4.2) we have

$$\begin{split} C_{\varepsilon} &:= E_{\varepsilon} \exp \left\langle -i \sum_{j=0}^{n} \left\langle Y_{t_{j}}^{\varepsilon}, \phi_{j} \right\rangle \right\rangle \\ &(5.3.3) &= E_{\varepsilon} \exp \left\langle -i \sum_{j=0}^{n-1} \left\langle Y_{t_{j}}^{\varepsilon}, \phi_{j} \right\rangle + \left\langle X_{t_{n-1}}^{\varepsilon}, \varepsilon^{-d} \log \left[1 - \varepsilon^{d-k} u_{\varepsilon}(\tau) \right] \right\rangle \\ &+ i \varepsilon^{-k} \left\langle \Lambda_{t_{n-1}}, S_{\tau} \phi_{n} \right\rangle \right\rangle, \end{split}$$

where we have denoted $\tau = t_n - t_{n-1}$ and u_{ε} solves (5.3.2). This can be written

(h)
$$C_{\epsilon} = E_{\epsilon} \exp\{-iA_{1,\epsilon} - iA_{2,\epsilon} + G_{\epsilon} - B_{1,\epsilon} + B_{2,\epsilon}\}$$

with

$$\begin{split} A_{1,\,\varepsilon} &= \sum_{j=0}^{n-1} \left\langle Y_{t_j}^{\varepsilon}, \phi_j \right\rangle, \qquad A_{2,\,\varepsilon} &= \left\langle Y_{t_{n-1}}^{\varepsilon}, S_{\tau} \phi_n \right\rangle, \\ G_{\varepsilon} &= \left\langle X_{t_{n-1}}^{\varepsilon}, V \int_0^{\tau} S_{\tau-s} \left[h(\varepsilon^{-1} \cdot) (iS_s \phi_n)^{1+\beta} \right] ds \right\rangle, \\ B_{1,\,\varepsilon} &= \varepsilon^{-k} \left\langle X_{t_{n-1}}^{\varepsilon}, u_{\varepsilon}(\tau) - iS_{\tau} \phi_n + \varepsilon^k V \int_0^{\tau} S_{\tau-s} \left[h(\varepsilon^{-1} \cdot) (iS_s \phi_n)^{1+\beta} \right] ds \right\rangle \end{split}$$

and

$$B_{2,\,\epsilon} = \langle X^{\epsilon}_{t_{n-1}}, \, \epsilon^{-d} \big\{ \log \big[1 - \epsilon^{d-k} u_{\epsilon} \! \big(\tau \big) \big] \, + \, \epsilon^{d-k} u_{\epsilon} \! \big(\tau \big) \big\} \rangle.$$

We will show that

$$C_{\varepsilon} \to C := E \exp\{-iA_1 - iA_2 + G\}$$
 as $\varepsilon \to 0$,

with

$$A_1 = \sum_{j=1}^{n-1} Y_{t_j}(\phi_j), \qquad A_2 = Y_{t_{n-1}}(S_{\tau}\phi_n)$$

and

$$G = \left\langle \Lambda_{t_{n-1}}, V \overline{h} \int_0^{\tau} S_{\tau-s} (i S_s \phi_n)^{1+\beta} ds \right\rangle,$$

where E is the expectation with respect to the limit distribution. Note that G is nonrandom.

The convergence statements below refer to P_{ε}^{m} . Suppose for the moment that we have already proved the following lemma.

(5.3.4) Lemma. $B_{1,\epsilon}$ and $B_{2,\epsilon}$ converge to 0 in probability as $\epsilon \to 0$ and G_{ϵ} converges to G in probability as $\epsilon \to 0$.

Now, since $A_{1,\,\varepsilon}$ and $A_{2,\,\varepsilon}$ are real, $|\exp\{G_{\varepsilon}-B_{1,\,\varepsilon}+B_{2,\,\varepsilon}\}|=1$, and hence $|\exp\{G_{\varepsilon}-G-B_{1,\,\varepsilon}+B_{2,\,\varepsilon}\}|\leq e^{|G|}$, which is bounded. Therefore for any $\delta>0$, by (h)

$$\begin{split} |C_{\epsilon} - E_{\epsilon} \exp\{-iA_{1,\,\epsilon} - iA_{2,\,\epsilon} + G\}| \\ & \leq e^{|G|} E_{\epsilon} |\exp\{G_{\epsilon} - G - B_{1,\,\epsilon} + B_{2,\,\epsilon}\} - 1| \\ & \leq \operatorname{const.} \left(P_{\epsilon}^{m} \left[|G_{\epsilon} - G - B_{1,\,\epsilon} + B_{2,\,\epsilon}| > \delta\right] + e^{\delta} - 1\right). \end{split}$$

Taking limits as $\varepsilon \to 0$, using (5.3.4) and then letting δ tend to 0, we have

$$C_{\varepsilon} - E_{\varepsilon} \exp\{-iA_{1,\varepsilon} - iA_{2,\varepsilon} + G\} \to 0 \text{ as } \varepsilon \to 0.$$

Then, to finish the proof of the proposition, it suffices to show that

$$E_{\varepsilon} \exp\{-iA_{1,\varepsilon} - iA_{2,\varepsilon} + G\} \rightarrow C \text{ as } \varepsilon \rightarrow 0,$$

but this follows from the induction hypothesis and Lemma 3.2.

PROOF OF LEMMA (5.3.4). By the induction hypothesis, $\langle Y_{t_{n-1}}^{\epsilon}, \phi \rangle \Rightarrow Y_{t_{n-1}}(\phi)$ as $\epsilon \to 0$, $\phi \in C_p(\mathbb{R}^d)$. Hence

(i)
$$\langle X_{t_{n-1}}^{\varepsilon}, \phi \rangle \Rightarrow \langle \Lambda_{t_{n-1}}, \phi \rangle \text{ as } \varepsilon \to 0, \phi \in C_p(\mathbb{R}^d).$$

By Proposition (5.1.6), for $\varepsilon < \varepsilon_{T,1}$, we have

$$|B_{1,\,\epsilon}| \leq \mathrm{const.}\; \epsilon^{(\gamma-1)k} \langle X_{t_{n-1}}^{\epsilon}, S_{\tau} |\phi_n| \rangle,$$

where $\gamma > 1$. Then, by (2.1) and (i), $B_{1,\epsilon} \Rightarrow 0$ as $\epsilon \rightarrow 0$.

Now we turn to $B_{2, \epsilon}$. According to (5.1.2), $\epsilon^{d-k}u_{\epsilon}(\tau) \to 0$ as $\epsilon \to 0$ uniformly in x and h, since d-k > 0. Hence, for sufficiently small $\epsilon < \epsilon_{T,1}$,

$$|B_{2,\,\varepsilon}| \leq \mathrm{const.} \langle X_{t_{n-1}}^{\varepsilon}, \, \varepsilon^{-d} \big(\varepsilon^{d-k} |u_{\varepsilon}(\tau)| \big)^2 \rangle,$$

and then, by (5.1.4),

$$|B_{2,\,arepsilon}| \leq \mathrm{const.}\; \varepsilon^{d-2k} \langle X_{t_{n-1}}^{arepsilon}, \left(S_{\tau}|\phi_n|\right)^2 \rangle \leq \mathrm{const.}\; \varepsilon^{d-2k} \langle X_{t_{n-1}}^{arepsilon}, S_{\tau}|\phi_n| \rangle.$$

Hence, since d > 2k, by (3.2) and (i), $B_{2, \epsilon} \Rightarrow 0$ as $\epsilon \rightarrow 0$.

In order to show that $G_{\epsilon} \Rightarrow G$, suppose that we have already proved the following lemma.

(5.3.5) LEMMA. For
$$\phi \in C_p(\mathbb{R}^d)$$
,

$$\Box \int_0^{\tau} S_{\tau-s} \left[h(\varepsilon^{-1} \cdot) (iS_s \phi)^{1+\beta} \right] ds - \overline{h} \int_0^{\tau} S_{\tau-s} (iS_s \phi)^{1+\beta} ds \Box_p \to 0 \quad as \ \varepsilon \to 0$$

for m a.e. h.

[We will see in Lemma (5.3.6) that the integrals above belong to $C_{p,0}(\mathbb{R}^d)$.] Let

$$(j) \quad g_{\varepsilon} = \int_{0}^{\tau} S_{\tau-s} \left[h(\varepsilon^{-1} \cdot) (iS_{s}\phi)^{1+\beta} \right] ds, \qquad g = \overline{h} \int_{0}^{\tau} S_{\tau-s} (iS_{s}\phi)^{1+\beta} ds.$$

By (i) we have $\langle X_{t_{n-1}}^{\varepsilon}, g \rangle \Rightarrow \langle \Lambda_{t_{n-1}}, g \rangle$ as $\varepsilon \to 0$. Therefore to show that $G_{\varepsilon} \Rightarrow G$, it suffices to prove that $\langle X_{t_{n-1}}^{\varepsilon}, g_{\varepsilon} - g \rangle \Rightarrow 0$ as $\varepsilon \to 0$. Now (see Section 3.1)

$$|\langle X_{t_{n-1}}^{\epsilon}, g_{\epsilon} - g \rangle| \leq \square X_{t_{n-1}}^{\epsilon} \square_{-p} \square g_{\epsilon} - g \square_{p},$$

and by (i) $\Box X_{t_{n-1}}^{\epsilon} \Box_{-p} \Rightarrow \Box \Lambda_{t_{n-1}} \Box_{-p}$; hence the result follows from (5.3.5). \Box

Before proving Lemma (5.3.5) we need the following lemma.

(5.3.6) Lemma. For each $0<\varepsilon<1,$ g_{ε} and g defined by (j) belong to $C_{p,0}(R^d)$.

PROOF. g is a special case of g_{ε} with $h \equiv \overline{h}$. We will prove first the continuity of g_{ε} . Since $|iS_s\phi|^{\beta} \leq ||S_s\phi||_{\infty}^{\beta} \leq ||\phi||_{\infty}^{\beta}$, then

$$|S_{\tau-s}h(\varepsilon^{-1}\cdot)(iS_s\phi)^{1+\beta}| \leq \text{const.} ||\phi||_{\infty}^{1+\beta}.$$

Let $f_s \equiv (iS_s\phi)^{1+\beta}$. For each $0 < \delta < \tau$ and $x, z \in \mathbb{R}^d$, since $0 \le h \le 1$, by (k), $|g_s(x+z) - g_s(x)|$

(1)
$$\leq \int_0^{\tau - \delta} \int_{R^d} |f_s(y)| |p_{\tau - s}(x + z, y) - p_{\tau - s}(x, y)| dy ds$$
$$+ 2\delta \operatorname{const.} ||\phi||_{\infty}^{1 + \beta}.$$

The stable density $p_s(x)$ is uniformly continuous in $\{s: s \geq \delta\} \times R^d$. Hence for any such δ above there is an $\eta > 0$ such that if $|z| < \eta$ the first term on the right in (1) can be estimated by

$$\delta \int_0^{\tau - \delta} \int_{R^d} |f_s(y)| \, dy \, ds \le \delta \operatorname{const.} \int_0^{\tau} \int_{R^d} S_s |\phi|(y) \, dy \, ds$$
$$= \delta \operatorname{const.} \tau \int_{R^d} |\phi|(y) \, dy = \operatorname{const.} \delta.$$

Hence the continuity of g_{ϵ} is obtained.

For all $0 \le s \le \tau$, by (k) we have

$$\text{(m)} \quad |S_{\tau-s}h(\varepsilon^{-1} \cdot) \big(iS_s \phi\big)^{1+\beta} |/\phi_p \leq \text{const.}(S_{\tau}|\phi|)/\phi_p \leq \text{const.} \square S_{\tau}|\phi|\square_p,$$

which is finite by (3.2).

By (3.2),

$$|S_s \phi|^{1+\beta}/\phi_p \le |S_s \phi|^{\beta} |S_s \phi|/\phi_p \le \square S_s \phi \square_p |S_s \phi|^{\beta},$$

and since $S_s \phi \in C_0(\mathbb{R}^d)$, then $|S_s \phi|^{1+\beta} \in C_{p,0}(\mathbb{R}^d)$; consequently

(n)
$$\lim_{|x| \to \infty} |x|^{2p} |S_s \phi(x)|^{1+\beta} = 0.$$

Now, $|S_{\tau-s}h(\varepsilon^{-1}\cdot)(iS_s\phi)^{1+\beta}| \leq S_{\tau-s}|S_s\phi|^{1+\beta}$; therefore by (n) and Iscoe (1986), Proposition 2.3,

$$\lim_{|x|\to\infty}|x|^{2p}|S_{\tau-s}h(\varepsilon^{-1}\cdot)(iS_s\phi)^{1+\beta}(x)|=0,$$

so

(o)
$$\lim_{|x|\to\infty} |S_{\tau-s}h(\varepsilon^{-1}\cdot)(iS_s\phi)^{1+\beta}(x)|/\phi_p(x)=0.$$

The dominated convergence theorem can be applied due to (m), hence (o) implies $g_{\varepsilon} \in C_{p,0}(\mathbb{R}^d)$. \square

PROOF OF LEMMA (5.3.5). It follows from the proof of (5.3.6) that $(1 + |x|^2)^p |g_{\epsilon}(x) - g(x)|$ is uniformly continuous on R^d uniformly in ϵ . Hence, to prove the lemma, it suffices to show that $(1 + |x|^2)^p |g_{\epsilon}(x) - g(x)| \to 0$, m a.e. for fixed $x \in R^d$, and therefore it is enough to prove that $g_{\epsilon}(x) - g(x) \to 0$, m a.e. for fixed x (the exceptional set may depend on x, but only countably many x are needed).

From (j) we have, with $f_s \equiv (iS_s\phi)^{1+\beta}$,

$$g_{\varepsilon}(x) - g(x) = \int_0^{\tau} \int_{R^d} \left(h(\varepsilon^{-1}y) - \overline{h} \right) f_s(y) p_{\tau - s}(y - x) \, dy \, ds$$
$$= \varepsilon^d \int_0^{\tau} \int_{R^d} \left(h(z) - \overline{h} \right) f_s(\varepsilon z) p_{\tau - s}(\varepsilon z - x) \, dz \, ds.$$

Hence

$$|g_{\varepsilon}(x) - g(x)| = \varepsilon^d \left| \int_{\mathbb{R}^d} (h(z) - \overline{h}) \theta(\varepsilon z) dz \right|,$$

where

$$\theta(z) = \int_0^{\tau} p_s(z-x) f_s(z) ds.$$

Note that $\theta(z) \ge 0$ and $\int_{\mathbb{R}^d} \theta(z) dz = \tau$. Now,

$$\left| \varepsilon^d \right| \int_{R^d} (h(z) - \overline{h}) \theta(\varepsilon z) dz \right| \leq A_{\varepsilon} + B_{\varepsilon},$$

where

$$egin{aligned} A_{arepsilon} &= arepsilon^d \! \int_{R^d} \! |h(z) - \overline{h}| heta(arepsilon z) 1_{\{|arepsilon z| > lpha\}} \, dz, \ &B_{arepsilon} &= arepsilon^d \! \left| \int_{R^d} \! (h(z) - \overline{h}) heta(arepsilon z) 1_{\{|arepsilon z| \leq lpha\}} \, dz
ight|, \end{aligned}$$

and a > 0 is arbitrary.

Given arbitrary $\delta > 0$, taking a sufficiently large, we have

$$A_{\varepsilon} \leq 2 \int_{\mathbb{R}^d} \theta(y) 1_{\{|y| \geq a\}} dy < \delta.$$

Define $F: H \to R$ by $F(h) = (h(0) - \overline{h})$. Hence $F(T_z h) = (h(z) - \overline{h})$, where T_z is the translation by z on H. Thus

$$B_{\varepsilon} = \varepsilon^{d} \int_{R^{d}} F(T_{z}h) \theta(\varepsilon z) 1_{\{|\varepsilon z| \leq a\}} dz.$$

The function $\theta(y)1_{\{|y| \le a\}}$ satisfies the conditions of Tempel'man's individual ergodic theorem [Tempel'man (1972), Corollary 7.2]. By this theorem,

$$B_{\varepsilon} \to \int_{\mathbb{R}^d} \theta(y) 1_{\{|y| \le a\}} dy \hat{F}(h)$$

as $\varepsilon \to 0$ for m a.e. h. Since m is ergodic with respect to T_z , then \hat{F} is constant on H, and we have $\hat{F} \equiv 0$.

Since δ is arbitrary we have the desired result. \square

The proof of Proposition (5.3.1) is finally complete. \square

(5.3.7) COROLLARY. The distributions of the limit random vectors $(Y_{t_0}(\phi_0), Y_{t_1}(\phi_1), \dots, Y_{t_n}(\phi_n))$ in Proposition (5.3.1) satisfy the relations

$$egin{aligned} E \exp iggl(-i \sum_{j=0}^{n} Y_{t_{j}}(\phi_{j}) iggr) \ &= E \exp iggl(-i iggl[\sum_{j=0}^{n-1} Y_{t_{j}}(\phi_{j}) + Y_{t_{n-1}}(S_{t_{n}-t_{n-1}}\phi_{n}) iggr] \ &+ iggl(\Lambda_{t_{n-1}}, V \overline{h} iggr)_{t_{n-1}}^{t_{n}} S_{t_{n}-r} (i S_{r-t_{n-1}}\phi_{n})^{1+eta} \, dr iggr) iggr\} \end{aligned}$$

and

$$\begin{split} E\left[\exp\left\{-iY_{t_{n}}(\phi_{n})\right\}|Y_{t_{0}}(\phi_{0}),\ldots,Y_{t_{n-1}}(\phi_{n-1});\;\phi_{0},\ldots,\phi_{n-1}\in C_{p}(R^{d})\right]\\ (\mathrm{ii}) &=\exp\left\{-iY_{t_{n-1}}\!\!\left(S_{t_{n}-t_{n-1}}\!\!\phi_{n}\right)+\left\langle\Lambda_{t_{n-1}},V\!\bar{h}\int_{t_{n-1}}^{t_{n}}\!\!S_{t_{n}-r}\!\!\left(iS_{r-t_{n-1}}\!\!\phi_{n}\right)^{1+\beta}dr\right\rangle\right\},\\ &\phi_{n}\in C_{p}(R^{d}). \end{split}$$

PROOF. (i) follows directly from the argument before Lemma (5.3.4). (i) implies that $Y_{t_n}(\phi_n) - Y_{t_{n-1}}(S_{t_n-t_{n-1}}\phi_n)$ and $\{Y_{t_0}(\phi_0), \ldots, Y_{t_{n-1}}(\phi_{n-1})\}$ are independent, which yields (ii). \square

- (5.3.8) Remarks. 1. Expression (i) in the proof is the law of large numbers for N^{ε} (Corollary 4.8).
- 2. According to Proposition (5.3.1), the finite-dimensional distributions of Y^{ε} converge if and only if the initial fluctuations converge.
- 3. For Proposition (5.3.1) the assumptions $E\langle N_0^{\varepsilon}, \phi_p \rangle < \infty$ and $E\langle X_0^{\varepsilon}, \phi^{\varepsilon} \rangle \to 0$ if $\phi^{\varepsilon} \to 0$ are not needed.
- 4. A weaker form of Lemma (5.3.5), namely mean square convergence, can be proved using a statistical ergodic theorem [Tempel'man (1972) and Fleischmann (1978)]. This suffices for the proof of assertion (a) of Theorem (4.5).
- 5. Once it is proved that $Y^{\varepsilon} \Rightarrow Y$ in $D_{\mathscr{S}'(R^d)}$ (Section 5.4), we can assert that the finite-dimensional distributions of the process Y satisfy

$$\left(\langle Y_{t_0}, \phi_0 \rangle, \dots, \langle Y_{t_n}, \phi_n \rangle\right) = \left(Y_{t_0}(\phi_0), \dots, Y_{t_n}(\phi_n)\right)$$

in distribution for $0=t_0 < t_1 < \cdots < t_n, \phi_0, \ldots, \phi_n \in \mathcal{S}(\mathbb{R}^d)$, and therefore we may write (5.3.7) (i) and (5.3.7)(ii) with $Y_{t_j}(\phi_j)$ replaced by $\langle Y_{t_j}, \phi_j \rangle$. Note that (5.3.7)(ii) implies the Markov property of Y, and therefore we have (4.6), where the meaning of $\langle Y_s, S_{t-s}\phi \rangle$ has been explained in Remark (4.7).

5.4. Tightness and convergence in $D_{\mathscr{S}'(R^d)}$ [Proof of Theorem (4.5)]. In this section we complete the proof that Y^{ε} converges under P^m_{ε} in $D_{\mathscr{S}'(R^d)}$ as $\varepsilon \to 0$ to an $\mathscr{S}'(R^d)$ -valued process whose finite-dimensional distributions are given by (5.3.7). Since (5.3.1) yields the uniqueness of subsequential limits, it suffices to

establish the tightness of $\{Y^{\epsilon_n}\}_n$ in $D_{\mathscr{S}'(R^d)}$ for $\epsilon_n \to 0$. By Mitoma (1983) it suffices to prove that the sequence $\{\langle Y^{\epsilon_n}, \phi \rangle\}_n$ is tight in D_R for each $\phi \in \mathscr{S}(R^d)$. According to the criterion of Aldous (1978), to prove the latter, it suffices to show that

- (i) for each fixed t, $\{\langle Y_t^{\epsilon_n}, \phi \rangle\}_n$ is tight in R,
- (ii) given stopping times τ_n bounded by T and $\delta_n \to 0$ as $n \to \infty$, then $\langle Y_{\tau_n + \delta_n}^{\epsilon_n}, \phi \rangle \langle Y_{\tau_n}^{\epsilon_n}, \phi \rangle \to 0$ in probability as $n \to \infty$.

Condition (i) follows immediately from the convergence of finite-dimensional distributions [Proposition (5.3.1)]. Condition (ii) is established in Proposition (5.4.8) below.

We will recall first some general facts. Let X be a random variable and $0 < \theta < \infty$. If $\int_K^\infty r^\theta P[|X| \ge r] \, dr < \infty$ for an arbitrary $0 < K < \infty$, then $E|X|^{1+\theta} < \infty$ and

$$E|X|^{1+\theta} \leq K^{1+\theta} + (1+\theta) \int_{K}^{\infty} r^{\theta} P[|X| \geq r] dr.$$

If f_X denotes the characteristic function of X, then

$$(5.4.1) P[|X| \ge r] \le Cr \int_0^{1/r} [1 - \operatorname{Re} f_X(\lambda)] d\lambda,$$

where C is a constant [see, e.g., Breiman (1968)]. Therefore in order to show that $E|X|^{1+\theta} < \infty$, it suffices to prove that

(5.4.2)
$$\int_{K}^{\infty} r^{1+\theta} \int_{0}^{1/r} \left[1 - \operatorname{Re} f_{X}(\lambda)\right] d\lambda \, dr < \infty$$

for some K, and combining (5.4.1) and (5.4.2), we have

$$(5.4.3) \quad E|X|^{1+\theta} \le K^{1+\theta} + C(1+\theta) \int_{K}^{\infty} r^{1+\theta} \int_{0}^{1/r} [1 - \operatorname{Re} f_{X}(\lambda)] \ d\lambda \ dr$$

for arbitrary K. These results hold also for conditional expectation (with conditional characteristic function).

In order to apply Aldous' tightness criterion we need the following lemmas. The constants T, L and $\varepsilon_{T,L}$ are those of Section 5.1.

(5.4.4) LEMMA. For $\phi \in C_p(\mathbb{R}^d)$, $\varepsilon < \varepsilon_{T,L}$, $0 < t \le T$, $K \ge 1/L$, $0 < \theta < \beta$ and h fixed,

(i)
$$E_{h} \Big[|\langle Y_{t}^{\varepsilon}, \phi \rangle|^{1+\theta} | X_{0}^{\varepsilon} \Big] \\ \leq K^{1+\theta} + \text{const. } K^{\theta-\beta} \Big[\langle X_{0}^{\varepsilon}, S_{t} | \phi | \rangle + H_{t}^{\varepsilon}(\phi) \Big],$$

where

$$(\mathrm{ii}) \qquad H^{\varepsilon}_t \coloneqq \sup_{\lambda \geq 0} \lambda^{-(1+\beta)} \big(1 - \cos \langle Y^{\varepsilon}_0, S_t \phi \rangle \lambda \big) = K_{\beta} |\langle Y^{\varepsilon}_0, S_t \phi \rangle|^{1+\beta},$$

 K_{β} is a constant depending on β , and

$$\begin{array}{ll} \text{(iii)} & & E_h \big[|\langle X_t^{\varepsilon}, \phi \rangle|^{1+\theta} | X_0^{\varepsilon} \big] \\ & \leq K^{1+\theta} + \text{const.} \ K^{\theta-\beta} \big[\varepsilon^k \! \langle X_0^{\varepsilon}, S_t \! | \phi | \rangle + J_t^{\varepsilon} \! \big(\varepsilon^k \! \phi \big) \big], \end{array}$$

where

$$\begin{aligned} J_t^{\varepsilon}(\psi) &\coloneqq \sup_{\lambda \geq 0} \lambda^{-(1+\beta)} \big[1 - \cos \big(\langle Y_0^{\varepsilon}, \psi \rangle + \langle \Lambda_t, \varepsilon^{-k} \psi \rangle \big) \lambda \big] \\ &= K_{\beta} |\langle Y_0^{\varepsilon}, \psi \rangle + \langle \Lambda_t, \varepsilon^{-k} \psi \rangle|^{1+\beta}, \qquad \psi \in C_p(\mathbb{R}^d). \end{aligned}$$

PROOF. Let $f(\lambda) = E_h[\exp\{i\lambda\langle Y_t^{\epsilon}, \phi\rangle\}|X_0^{\epsilon}], \ \lambda \geq 0$. Then, using (5.2.1) and proceeding as in Section 5.3 [see (5.3.2) and (5.3.3)], we have

(a)
$$f(\lambda) = \exp\{\langle X_0^{\varepsilon}, \varepsilon^{-d} \log[1 - \varepsilon^{d-k} u_{\varepsilon}(t, -\phi \lambda)] \rangle - i\lambda \langle \varepsilon^{-k} \Lambda_0, S_t \phi \rangle\},$$

where $u_{\varepsilon}(t) \equiv u_{\varepsilon}(t, \psi), \psi \in C_0(\mathbb{R}^d)$, satisfies

$$\begin{split} u_{\epsilon}(t) &= S_t \big[\varepsilon^{k-d} \big(1 - \exp \big\{ -i \varepsilon^{d-k} \psi \big\} \big) \big] \\ &- \varepsilon^k V \! \int_0^t \! S_{t-s} \big[h(\varepsilon^{-1} \cdot) u_{\epsilon}^{1+\beta}(s) \big] \, ds. \end{split}$$

We can write (a) as

(c)
$$f(\lambda) = \exp\{\langle X_0^{\varepsilon}, A_{\varepsilon}(\lambda) + B_{\varepsilon}(\lambda) + C_{\varepsilon}(\lambda) \rangle\} \exp\{iD_{\varepsilon}\lambda\},$$

where

$$egin{aligned} A_{arepsilon}(\lambda) &= arepsilon^{-d} iggl[\log iggl\{ 1 - arepsilon^{d-k} u_{arepsilon}(t, -\phi \lambda) iggr\} + arepsilon^{d-k} u_{arepsilon}(t, -\phi \lambda) iggr], \ B_{arepsilon}(\lambda) &= -arepsilon^{-k} iggl\{ u_{arepsilon}(t, -\phi \lambda) - i S_{t}(-\phi \lambda) + arepsilon^{t} V \int_{0}^{t} S_{t-s} iggl[h(arepsilon^{-1} \cdot) (-i S_{s} \phi)^{1+eta} iggr] \ ds iggr\}, \ C_{arepsilon}(\lambda) &= V \int_{0}^{t} S_{t-s} iggl[h(arepsilon^{-1} \cdot) (-i S_{s} \phi \lambda)^{1+eta} iggr] \ ds, \ D_{arepsilon} &= \langle Y_{0}^{arepsilon}, S_{t} \phi
angle. \end{aligned}$$

Hence
$$f(\lambda) = (1 + Z)(1 + W) = 1 + Z(1 + W) + W$$
, where

$$Z = \exp\{\langle X_0^{\epsilon}, A_{\epsilon}(\lambda) + B_{\epsilon}(\lambda) + C_{\epsilon}(\lambda) \rangle\} - 1, \qquad W = \exp\{iD_{\epsilon}\lambda\} - 1,$$

and therefore

(d)
$$1 - \operatorname{Re} f(\lambda) = -\operatorname{Re} Z(1 + W) - \operatorname{Re} W$$
$$\leq |Z| |1 + W| - \operatorname{Re} W \leq |Z| - \operatorname{Re} W.$$

We will now estimate |Z| and $-\operatorname{Re} W$.

By (a) and (c), $|\exp\{\langle X_0^{\epsilon}, A_{\epsilon}(\lambda) + B_{\epsilon}(\lambda) + C_{\epsilon}(\lambda)\rangle\}| \le 1$. If ξ is a complex number such that $|e^{\xi}| \le 1$, then $|e^{\xi} - 1| \le 2|\xi|$. Hence

(e)
$$|Z| \leq 2\langle X_0^{\varepsilon}, |A_{\varepsilon}(\lambda)| + |B_{\varepsilon}(\lambda)| + |C_{\varepsilon}(\lambda)| \rangle.$$

Now, by (5.1.2) and (5.1.4),

$$\begin{split} \langle X_0^{\varepsilon}, |A_{\varepsilon}(\lambda)| \rangle &\leq \text{const. } \varepsilon^{-d} \varepsilon^{2(d-k)} \langle X_0^{\varepsilon}, |u_{\varepsilon}(t, -\phi \lambda)|^2 \rangle \\ &\leq \text{const. } \varepsilon^{d-2k} \langle X_0^{\varepsilon}, (S_t |\phi| \lambda)^2 \rangle \leq \text{const. } \varepsilon^{d-2k} \langle X_0^{\varepsilon}, S_t |\phi| \rangle \lambda^{1+\beta}, \end{split}$$

with $\lambda \leq L$; recall that d > 2k. By (5.1.6),

$$\langle X_0^{\varepsilon}, |B_{\varepsilon}(\lambda)| \rangle \leq \text{const. } \varepsilon^{(\gamma-1)k} \langle X_0^{\varepsilon}, S_t |\phi| \rangle \lambda^{1+\beta},$$

where $\gamma > 1$. Finally,

$$\langle X_0^{\epsilon}, |C_{\epsilon}(\lambda)| \rangle \leq \left\langle X_0^{\epsilon}, V \int_0^t S_{t-s} |S_s \phi|^{1+\beta} ds \right\rangle \lambda^{1+\beta}$$

$$\leq \text{const.} \|\phi\|_{\infty}^{\beta} \langle X_0^{\epsilon}, S_t |\phi| \rangle \lambda^{1+\beta}.$$

Putting these results into (e), we obtain

$$|Z| \leq \operatorname{const.}\langle X_0^{\varepsilon}, S_t | \phi | \rangle \lambda^{1+\beta} \quad \text{with } \lambda \leq L.$$

For Re W we have

$$-\operatorname{Re} W = -\operatorname{Re} \exp\{iD_{\epsilon}\lambda\} + 1$$

= $-\cos D_{\epsilon}\lambda + 1 = \lambda^{1+\beta}(1 - \cos D_{\epsilon}\lambda)/\lambda^{1+\beta} \le H_{\epsilon}^{\epsilon}(\phi)\lambda^{1+\beta}$,

where $H_t^{\varepsilon}(\phi)$ is given by (ii).

Taking these inequalities into (d), we obtain

$$1 - \operatorname{Re} f(\lambda) \leq \operatorname{const.} \left[\langle X_0^{\varepsilon}, S_t | \phi | \rangle + H_t^{\varepsilon}(\phi) \right] \lambda^{1+\beta}.$$

Therefore, by (5.4.3), for $0 \le \theta < \beta$ and $K \ge 1/L$ (so that $\lambda \le L$) we have $E_h[|\langle Y_t^{\varepsilon}, \phi \rangle|^{1+\theta}|X_0^{\varepsilon}]$

$$\leq K^{1+\theta} + \text{const.} [\langle X_0^{\epsilon}, S_t | \phi | \rangle + H_t^{\epsilon}(\phi)] \int_K^{\infty} r^{1+\theta} \int_0^{1/r} \lambda^{1+\beta} d\lambda dr$$

$$= K^{1+\theta} + \text{const.} K^{\theta-\beta} [\langle X_0^{\epsilon}, S_t | \phi | \rangle + H_t^{\epsilon}(\phi)],$$

which proves (i).

(iii) and (iv) are obtained similarly to (i) and (ii), using

$$E\left[\exp\{i\langle X_t^{\varepsilon},\phi\rangle\}|X_0^{\varepsilon}\right] = E\left[\exp\{i\left(\langle \cdot,\varepsilon,\varepsilon^k\phi\rangle + \lambda\langle\Lambda_t,\phi\rangle\right)\}|X_0^{\varepsilon}\right]. \quad \Box$$

(5.4.5) LEMMA. Let $\phi^{\epsilon} \in C_p(\mathbb{R}^d)$ and assume $\phi^{\epsilon} \to 0$ as $\epsilon \to 0$. Then for $0 < \theta < \beta$ the following limits hold in $L^1(H, m)$:

(i)
$$\sup_{0 \le t \le T} E_h \Big[|\langle Y_t^{\epsilon}, \phi^{\epsilon} \rangle|^{1+\theta} |X_0^{\epsilon}| \Big] \to 0 \quad in \ probability \left(P_{\epsilon}^h \right) \ as \ \epsilon \to 0$$

and

(ii)
$$\sup_{0 \le t \le T} E_h \Big[|\langle X_t^{\epsilon}, \phi^{\epsilon} \rangle|^{1+\theta} | X_0^{\epsilon} \Big] \to 0 \quad \text{in probability } \Big(P_{\epsilon}^h \Big) \text{ as } \epsilon \to 0.$$

PROOF. Let $\delta > 0$. By (5.4.4)(i) and (5.4.4)(ii), (3.2) and (4.1),

$$\begin{split} \int_{H} & P_{\epsilon}^{h} \bigg[\sup_{t \leq T} E_{h} \big[|\langle Y_{t}^{\epsilon}, \phi^{\epsilon} \rangle|^{1+\theta} | X_{0}^{\epsilon} \big] > \delta \bigg] m(dh) \\ & \leq \int_{H} & P_{\epsilon}^{h} \bigg[K^{1+\theta} + \operatorname{const.} K^{\theta-\beta} \bigg[\langle X_{0}^{\epsilon}, \sup_{t \leq T} S_{t} | \phi^{\epsilon} | \rangle + \sup_{t \leq T} H_{t}^{\epsilon} (\phi^{\epsilon}) \bigg] > \delta \bigg\} m(dh) \\ & \leq 1_{\{K^{1+\theta} > \delta_{1}\}} + P_{\epsilon}^{m} \Big[K^{\theta-\beta} \langle X_{0}^{\epsilon}, | \phi^{\epsilon} | \rangle > \delta_{2} \Big] \\ & + P_{\epsilon}^{m} \bigg[K^{\theta-\beta} \sup_{t \leq T} H_{t}^{\epsilon} (\phi^{\epsilon}) > \delta_{3} \bigg], \end{split}$$

where δ_1 , δ_2 and δ_3 are some positive constants. Now,

$$\left\langle X_0^{\epsilon}, |\phi^{\epsilon}| \right\rangle \leq \square X_0^{\epsilon} \square_{-p} \square \phi^{\epsilon} \square_p, \qquad \square X_0^{\epsilon} \square_{-p} \Rightarrow \square \Lambda_0 \square_{-p} \quad \text{and} \quad \square \phi^{\epsilon} \square_p \rightarrow 0$$

by assumption; hence $P_{\varepsilon}^m[K^{\theta-\beta}\langle X_0^{\varepsilon}, |\phi^{\varepsilon}| \rangle > \delta_2] \to 0$. Since

$$\sup_{t \leq T} H_t^{\varepsilon}(\phi^{\varepsilon}) \leq \operatorname{const.} \square Y_0^{\varepsilon} \square_{-p}^{1+\beta} \square \phi^{\varepsilon} \square_p^{1+\beta} \quad \text{and} \quad \square Y_0^{\varepsilon} \square_{-p} \Rightarrow \square Y_0 \square_{-p},$$

then

$$P_{\varepsilon}^{m}\bigg[K^{\theta-\beta}\sup_{t\leq T}H_{t}^{\varepsilon}(\phi^{\varepsilon})>\delta_{3}\bigg]\rightarrow0.$$

Hence assertion (i) will be proved if we can let $K \to 0$.

Recall that the results we have used were proved for fixed $\phi \in C_n(\mathbb{R}^d)$,

$$(a) \quad \varepsilon < 1/\big(V\!\beta T\|\phi\|_{\infty}^{\beta}L^{\beta}\big)^{1/k} \quad \text{and} \quad \varepsilon_{T,\,L} < \text{const.}/\big(VT\|\phi\|_{\infty}^{\beta}L^{\beta}\big)^{1/k}$$

[see (5.1.3) and (5.1.4)], and $K \geq 1/L$. If we now let $\phi = \phi^{\epsilon}$ depend on ϵ , we may also let $L = L_{\epsilon}$ depend on ϵ in such a way that conditions (a) are not altered. For example, we may define $L_{\epsilon} = \|\phi^{\epsilon}\|_{\infty}^{-1}$. Now, since $\phi^{\epsilon} \to 0$ in $C_{p}(R^{d})$ implies $\|\phi^{\epsilon}\|_{\infty} \to 0$, then $L_{\epsilon} \to \infty$ and therefore we may let $K \to 0$ as desired.

Assertion (ii) is proved similarly, using (5.4.4)(iii) and (5.4.4)(iv). □

We will need the following result, which can be proved from (4.2): If \mathscr{G} is a sub- σ -algebra of the basic σ -algebra, then for any $A \in \mathscr{G}$,

(5.4.6)
$$\int_{A} E[\cdot|\mathscr{G}] dP = \int_{H} \int_{A} E_{h}[\cdot|\mathscr{G}] dP_{\varepsilon}^{h} m(dh).$$

(5.4.7) Lemma. Let $\phi^{\varepsilon} \in C_p(R^d)$ and assume that ϕ^{ε} is in the domain of Δ_{α} , $\Delta_{\alpha}\phi^{\varepsilon} \in C_p(R^d)$, $\phi^{\varepsilon} \to 0$ and $\Delta_{\alpha}\phi^{\varepsilon} \to 0$. Then for T > 0 and $0 < \theta < \beta$,

(i)
$$E\left[\sup_{0 \le t \le T} |\langle Y_t^{\varepsilon}, \phi^{\varepsilon} \rangle|^{1+\theta} | X_0^{\varepsilon}\right] \to 0$$
 in probability (P_{ε}^m) as $\varepsilon \to 0$

and

(ii)
$$E\left[\sup_{0 < t \le T} |\langle X_t^{\varepsilon}, \phi^{\varepsilon} \rangle|^{1+\theta} | X_0^{\varepsilon} \right] \to 0$$
 in probability (P_{ε}^m) as $\varepsilon \to 0$.

PROOF. First we will estimate $E[\sup_{0 \le t \le T} |\langle Y_t^{\varepsilon}, \phi \rangle|^{1+\theta} |X_0^{\varepsilon}], \phi \in \mathcal{S}(\mathbb{R}^d)$. We start by showing that

(a)
$$\langle Y_t^{\epsilon}, \phi \rangle - \int_0^t \langle Y_s^{\epsilon}, \Delta_{\alpha} \phi \rangle \, ds, \qquad t \geq 0,$$

is a P_{ε}^{m} -martingale. By (4.3) and (4.4),

$$\begin{split} \langle Y_t^{\varepsilon}, \phi \rangle &= \langle \varepsilon^{d-k} N_{\varepsilon^{-\alpha}t}, \phi(\varepsilon \cdot) \rangle - \langle \varepsilon^{-k} \Lambda_0, S_t \phi \rangle \\ &= \varepsilon^{d-k} \bigg[\langle N_{\varepsilon^{-\alpha}t}, \phi(\varepsilon \cdot) \rangle - \int_0^{\varepsilon^{-\alpha}t} \langle N_s, (\Delta_{\alpha}) \phi(\varepsilon \cdot) \rangle \, ds \bigg] \\ &+ \varepsilon^{d-k} \int_0^{\varepsilon^{-\alpha}t} \langle N_s, (\Delta_{\alpha}) \phi(\varepsilon \cdot) \rangle \, ds - \bigg\langle \varepsilon^{-k} \Lambda_0, \int_0^t S_s \Delta_{\alpha} \phi \, ds \bigg\rangle - \langle \varepsilon^{-k} \Lambda_0, \phi \rangle, \end{split}$$

where we have used $d/ds(S_s)=S_s\Delta_\alpha$. Since the term in brackets is a P^h_ϵ -martingale for fixed h because of (5.2.4), and the last term is independent of t, we have that the difference of the remaining terms constitutes a P^h_ϵ -martingale. Since $\varepsilon^{-\alpha}(\Delta_\alpha)\phi(\varepsilon\,\cdot\,)=(\Delta_\alpha\phi)(\varepsilon\,\cdot\,)$ [this follows from the self-similarity of $p_t(x)$], then

$$\begin{split} \varepsilon^{d-k} \int_0^{\varepsilon^{-a}t} \langle N_s, (\Delta_\alpha) \phi(\varepsilon \cdot) \rangle ds &- \left\langle \varepsilon^{-k} \Lambda_0, \int_0^t S_s \Delta_\alpha \phi \, ds \right\rangle \\ &= \varepsilon^{d-k} \int_0^t \langle N_{\varepsilon^{-a}s}, (\Delta_\alpha \phi)(\varepsilon \cdot) \rangle \, ds - \left\langle \varepsilon^{-k} \Lambda_0, \int_0^t S_s \Delta_\alpha \phi \, ds \right\rangle \\ &= \int_0^t \langle Y_s^{\varepsilon}, \Delta_\alpha \phi \rangle \, ds, \end{split}$$

and therefore (a) is a P_{ε}^h -martingale for each h. Using (5.4.6), one can then show that (a) is a P_{ε}^m -martingale.

Applying the conditional form of Doob's inequality to the martingale (a), we have

$$\begin{split} E\left[\sup_{0\leq t\leq T}|\langle Y^{\epsilon}_{t},\phi\rangle|^{1+\theta}|X^{\epsilon}_{0}\right] \\ &\leq \mathrm{const.} \bigg\langle E\left[\sup_{0\leq t\leq T}\left|\langle Y^{\epsilon}_{t},\phi\rangle - \int_{0}^{t}\langle Y^{\epsilon}_{s},\Delta_{a}\phi\rangle\,ds\,\right|^{1+\theta}|X^{\epsilon}_{0}\right] \\ &+ E\left[\sup_{0\leq t\leq T}\left|\int_{0}^{t}\langle Y^{\epsilon}_{s},\Delta_{a}\phi\rangle\,ds\,\right|^{1+\theta}|X^{\epsilon}_{0}\right] \bigg\rangle \\ &\leq \mathrm{const.} \bigg\langle E\left[\left|\langle Y^{\epsilon}_{T},\phi\rangle - \int_{0}^{T}\langle Y^{\epsilon}_{s},\Delta_{a}\phi\rangle\,ds\,\right|^{1+\theta}|X^{\epsilon}_{0}\right] \\ &+ E\left[\left(\int_{0}^{T}|\langle Y^{\epsilon}_{s},\Delta_{a}\phi\rangle|\,ds\,\right)^{1+\theta}|X^{\epsilon}_{0}\right] \bigg\rangle \\ &\leq \mathrm{const.} \bigg\langle E\left[\left|\langle Y^{\epsilon}_{T},\phi\rangle\right|^{1+\theta}|X^{\epsilon}_{0}\right] \\ &+ E\left[\left(\int_{0}^{T}|\langle Y^{\epsilon}_{s},\Delta_{a}\phi\rangle|\,ds\,\right)^{1+\theta}|X^{\epsilon}_{0}\right] \bigg\rangle, \end{split}$$

and, using Hölder's inequality we obtain

$$\begin{split} E \left[\sup_{0 \leq t \leq T} |\langle Y_t^{\varepsilon}, \phi \rangle|^{1+\theta} |X_0^{\varepsilon} \right] \\ & \leq \operatorname{const.} \left\{ E \left[|\langle Y_T^{\varepsilon}, \phi \rangle|^{1+\theta} |X_0^{\varepsilon} \right] + \int_0^T E \left[|\langle Y_s^{\varepsilon}, \Delta_{\alpha} \phi \rangle|^{1+\theta} |X_0^{\varepsilon} \right] ds \right\} \\ & \leq \operatorname{const.} \left\{ E \left[|\langle Y_T^{\varepsilon}, \phi \rangle|^{1+\theta} |X_0^{\varepsilon} \right] + \sup_{0 \leq t \leq T} E \left[|\langle Y_t^{\varepsilon}, \Delta_{\alpha} \phi \rangle|^{1+\theta} |X_0^{\varepsilon} \right] \right\}. \end{split}$$

Therefore assertion (i) will follow if we show that

$$\sup_{0 \le t \le T} E\left[|\langle Y_t^{\varepsilon}, \psi^{\varepsilon} \rangle|^{1+\theta} | X_0^{\varepsilon} \right] \to 0$$

in probability as $\varepsilon \to 0$, where $\psi^{\varepsilon} \to 0$ in $C_p(\mathbb{R}^d)$.

Let $\delta > 0$. We have

$$P \Big\{ \sup_{t \le T} E \big[|\langle Y^{\epsilon}_t, \psi^{\epsilon} \rangle|^{1+\theta} |X^{\epsilon}_0 \big] > \delta \Big\} = P \Big[\sup_{t \le T} d\mu^{\epsilon}_t / dP > \delta \Big],$$

where μ_t^{ϵ} is the measure

$$\mu_t^\epsilon(\,\cdot\,) = \int_H\!\int_\cdot E_hig[|\langle Y_t^\epsilon,\psi^\epsilon
angle|^{1\,+\, heta}|X_0^\epsilon]\;dP_\epsilon^h m(\,dh)$$

[see (5.4.6)] and $d\mu_t^{\epsilon}/dP$ in the Radon-Nikodym derivative. Now,

$$\mu_t^{\varepsilon}(\,\cdot\,) \leq \nu^{\varepsilon}(\,\cdot\,) \equiv \int_H \int_{\cdot} \sup_{t < T} E_h \Big[|\langle Y_t^{\varepsilon}, \psi^{\varepsilon} \rangle|^{1 + \theta} |X_0^{\varepsilon}] \, dP_{\varepsilon}^h m(\,dh\,),$$

and, using (5.4.4)(i) and (3.2), we obtain

$$\nu^{\epsilon}(\,\cdot\,) \leq \int_{\cdot} \mathrm{const.}\Big(K^{\,1+\theta} + K^{\,\theta-\beta}\Big[\langle\,X_0^{\,\epsilon},|\psi^{\epsilon}|\,\rangle \,+\, \sup_{t\,<\,T} H_t^{\epsilon}(\,\psi^{\epsilon})\,\Big]\Big)\,dP.$$

Hence $d\mu_t^{\epsilon}/dP = (d\mu_t^{\epsilon}/d\nu^{\epsilon})(d\nu^{\epsilon}/dP) \le d\nu^{\epsilon}/dP$. Therefore

$$P\!\!\left\langle \sup_{t\leq T} E\!\left[|\langle Y_t^{\epsilon}, \psi^{\epsilon} \rangle|^{1+\theta} | X_0^{\epsilon} \right] > \delta \right\rangle \leq P(A),$$

where $A \equiv [d\nu^{\epsilon}/dP > \delta]$; so

$$u^{\epsilon}(A) = \int_{A} (d\nu^{\epsilon}/dP) dP \geq \delta P(A),$$

and, using the previous results, we have

$$\begin{split} P \Big\langle \sup_{t \leq T} E \left[|\langle Y_t^{\varepsilon}, \psi^{\varepsilon} \rangle|^{1+\theta} | X_0^{\varepsilon} \right] &> \delta \Big\rangle \\ &\leq \delta^{-1} \nu^{\varepsilon}(A) \\ &\leq \delta^{-1} \operatorname{const.} \Big\langle K^{1+\theta} + K^{\theta-\beta} \Big[E \langle X_0^{\varepsilon}, |\psi^{\varepsilon}| \rangle + E \sup_{t \leq T} H_t^{\varepsilon}(\psi^{\varepsilon}) \Big] \Big\rangle. \end{split}$$

Since $\psi^{\varepsilon} \to 0$, by assumption we have $E\langle X_0^{\varepsilon}, |\psi^{\varepsilon}| \rangle \to 0$. Also, $E \sup_{t \le T} H_t^{\varepsilon}(\psi^{\varepsilon}) \to 0$ since

$$\sup_{t\leq T} H^{\epsilon}_t(\psi^{\epsilon}) \leq \mathrm{const.} \Box Y^{\epsilon}_0\Box^{1+\beta}_{-p}\Box \psi^{\epsilon}\Box^{1+\beta}_{p} \quad \text{and} \quad \sup_{\epsilon} E\Box Y^{\epsilon}_0\Box^{1+\beta}_{-p} < \infty.$$

Finally, we can let $K \to 0$. Assertion (i) is proved.

Assertion (ii) is proved similarly, using (5.4.4)(iii), (5.4.4)(iv) and (5.4.5)(ii). \square

(5.4.8) Lemma. For $0 < \varepsilon < 1$, let τ_{ε} be stopping times bounded by T with respect to the filtration $(\sigma\{X_s^{\varepsilon}, s \leq t; h\})_{t \geq 0}$, and $\delta_{\varepsilon} > 0$ such that $\delta_{\varepsilon} \to 0$ as

 $\varepsilon \to 0$. Then for each $\phi \in \mathcal{S}(\mathbb{R}^d)$,

$$\langle Y_{\tau_{\epsilon}+\delta_{\epsilon}}^{\epsilon}, \phi \rangle - \langle Y_{\tau_{\epsilon}}^{\epsilon}, \phi \rangle \to 0 \quad in \, probability \left(P_{\epsilon}^{m}\right) \, as \, \, \epsilon \to 0.$$

PROOF. Recall we denote $P \equiv P_{\varepsilon}^{m}$ and E the corresponding expectation. Let $\tau \equiv \tau_{\varepsilon}$, $\delta \equiv \delta_{\varepsilon}$, $\eta > 0$. Since

$$\begin{split} P\big[|\langle Y^{\varepsilon}_{\tau+\delta}, \phi \rangle - \langle Y^{\varepsilon}_{\tau}, \phi \rangle| > \eta\big] &= E\big(P\big[|\langle Y^{\varepsilon}_{\tau+\delta}, \phi \rangle - \langle Y^{\varepsilon}_{\tau}, \phi \rangle| > \eta|X^{\varepsilon}_{0}, h\big]\big) \\ &= E\big(P^{h}_{\varepsilon}\big[|\langle Y^{\varepsilon}_{\tau+\delta}, \phi \rangle - \langle Y^{\varepsilon}_{\tau}, \phi \rangle| > \eta|X^{\varepsilon}_{0}\big]\big), \end{split}$$

by dominated convergence it suffices to prove that

(a)
$$P_{\varepsilon}^{h}[|\langle Y_{\tau+\delta}^{\varepsilon}, \phi \rangle - \langle Y_{\tau}^{\varepsilon}, \phi \rangle| > \eta |X_{0}^{\varepsilon}] \to 0$$
 in probability (P) as $\varepsilon \to 0$.

As a consequence of (5.2.3) the process $(Y^{\epsilon}, h) \equiv \{(Y^{\epsilon}_{t}, h), t \geq 0\}$ is strong Markov with respect to the given filtration; hence we have

$$(b) \begin{array}{c} P_{\varepsilon}^{h} \big[|\langle Y_{\tau+\delta}^{\varepsilon}, \phi \rangle - \langle Y_{\tau}^{\varepsilon}, \phi \rangle| > \eta |X_{0}^{\varepsilon} \big] \\ = E_{h} \big[P_{\varepsilon}^{h} \big[|\langle Y_{\tau+\delta}^{\varepsilon}, \phi \rangle - \langle Y_{\tau}^{\varepsilon}, \phi \rangle| > \eta |X_{\tau}^{\varepsilon} \big] |X_{0}^{\varepsilon} \big]. \end{array}$$

Now, by (5.4.1),

$$(\mathrm{c}) \quad P_{\epsilon}^{h}\big[|\langle Y_{\tau+\delta}^{\epsilon}, \phi \rangle - \langle Y_{\tau}^{\epsilon}, \phi \rangle| > \eta |X_{\tau}^{\epsilon}\big] \leq C\eta \int_{0}^{1/\eta} [1 - \mathrm{Re}\,f(\lambda)] \; d\lambda,$$

where

$$f(\lambda) := E_h \left[\exp\{i\lambda \langle Y_{\tau+\delta}^{\epsilon}, \phi \rangle - \langle Y_{\tau}^{\epsilon}, \phi \rangle\} | X_{\tau}^{\epsilon} \right].$$

Proceeding as in Section 5.3 [see (5.3.2) and (5.3.3)], with the strong Markov property, $f(\lambda)$ is given by

$$egin{aligned} f(\lambda) &= \expig\{-i\lambda\langle Y^arepsilon_ au, \phi
angle + \langle X^arepsilon_ au, arepsilon^{-d}\logig[1-arepsilon^{d-k}u_arepsilon(\delta,-\phi\lambda)ig]
angle \ &-i\lambda\langlearepsilon^{-k}\Lambda_ au, S_\delta\phi
angleig\}, \end{aligned}$$

where $u_s(t) \equiv u_s(t, \psi), \psi \in \mathcal{S}(\mathbb{R}^d)$, satisfies

$$u_{\epsilon}(t) = S_t \left[\epsilon^{k-d} \left(1 - \exp\left\{ -i \epsilon^{d-k} \psi \right\} \right) \right] - \epsilon^k V \int_0^t S_{t-s} \left[h(\epsilon^{-1} \cdot) u_{\epsilon}^{1+\beta}(s) \right] ds;$$

in particular, $u_{\epsilon}(0,\psi) = \epsilon^{k-d}(1-\exp\{-i\epsilon^{d-k}\psi\})$, and therefore

$$\log[1-\varepsilon^{d-k}u_{\varepsilon}(0,-\phi\lambda)]=i\varepsilon^{d-k}\phi\lambda.$$

 $f(\lambda)$ can then be written as

$$f(\lambda) = \exp\{i\langle Y_{\tau}^{\epsilon}, S_{\delta}\phi - \phi\rangle\lambda + \langle X_{\tau}^{\epsilon}, a_{1}^{\epsilon}(\lambda) - a_{2}^{\epsilon}(\lambda) - a_{3}^{\epsilon}(\lambda)\rangle\},\$$

where

$$\begin{split} &a_1^{\epsilon}(\lambda) = \varepsilon^{-d} \big\{ \log \big[1 - \varepsilon^{d-k} u_{\epsilon}(\delta, -\phi \lambda) \big] + \varepsilon^{d-k} u_{\epsilon}(\delta, -\phi \lambda) \big\}, \\ &a_2^{\epsilon}(\lambda) = \varepsilon^{-d} \big\{ \log \big[1 - \varepsilon^{d-k} u_{\epsilon}(0, -S_{\delta} \phi \lambda) \big] + \varepsilon^{d-k} u_{\epsilon}(0, -S_{\delta} \phi \lambda) \big\}, \\ &a_3^{\epsilon}(\lambda) = \varepsilon^{-k} \big[u_{\epsilon}(\delta, -\phi \lambda) - u_{\epsilon}(0, -S_{\delta} \phi \lambda) \big]. \end{split}$$

By (a), (b) and (c) we have to show that

$$\int_0^{1/\eta} (1 - E_h[\operatorname{Re} f(\lambda) | X_0^{\epsilon}]) d\lambda \to 0 \quad \text{in probability } (P) \text{ as } \epsilon \to 0,$$

and since $f(\lambda)$ is bounded, it suffices to prove that

$$(\mathrm{d}) \qquad \qquad E\bigg[\sup_{0 \leq t \leq T} |\langle Y^{\varepsilon}_t, S_{\delta} \phi - \phi \rangle|^{1+\theta} |X^{\varepsilon}_0\bigg] \to 0 \quad \text{in probability as } \varepsilon \to 0,$$

(e)
$$E\left[\sup_{0 \le \lambda \le 1/\eta} \sup_{0 \le t \le T} |\langle X_t^{\epsilon}, a_j^{\epsilon}(\lambda) \rangle|^{1+\theta} |X_0^{\epsilon}\right] \to 0 \quad \text{in probability as } \epsilon \to 0,$$

for j=1,2,3. (d) holds by (5.4.7)(i) since $S_{\delta}\phi - \phi \to 0$ and $\Delta_{\alpha}(S_{\delta}\phi - \phi) \to 0$ in $C_p(R^d)$. Indeed, $\phi \in \mathcal{S}(R^d)$ implies $\Delta_{\alpha}\phi \in C_{p,0}(R^d)$ [Dawson and Gorostiza (1988)], and $\Delta_{\alpha}(S_{\delta}\phi - \phi) = S_{\delta}\Delta_{\alpha}\phi - \Delta_{\alpha}\phi \to 0$ as $\delta \to 0$ by (3.3). For (e) with j=1, proceeding as in the proof of (5.3.4), we have

$$\varepsilon^{-d} \Big\{ \log \Big[1 - \varepsilon^{d-k} u_{\varepsilon}(t, -\phi \lambda) \Big] + \varepsilon^{d-k} u_{\varepsilon}(t, -\phi \lambda) \Big\} \le \text{const. } \varepsilon^{d-2k} S_{t} |\phi| \lambda$$

$$\le \text{const. } \varepsilon^{d-2k} S_{t} \psi,$$

where $\psi \in \mathcal{S}(R^d)_+$ is such that $|\phi| \leq \psi$ and $\lambda \leq 1/\eta$. We can now apply (5.4.7)(ii) with $\phi^{\varepsilon} = \varepsilon^{d-2k} S_t \psi$, hence (e) holds for j=1, since d>2k. The proof for j=2 is similar. For j=3 by (5.1.6) we have

$$\begin{split} |a_3^{\epsilon}(\lambda)| &\leq \epsilon^{-k} \bigg| u_{\epsilon}(\delta, -\phi\lambda) + iS_{\delta}\phi\lambda + \epsilon^{k}V \int_0^{\delta} S_{\delta-s} \Big[h(\epsilon^{-1} \cdot) (-iS_{s}\phi\lambda)^{1+\beta} \Big] ds \bigg| \\ &+ V \int_0^{\delta} S_{\delta-s} |S_{s}\phi\lambda|^{1+\beta} ds + \epsilon^{-k}| - iS_{\delta}\phi\lambda - u(0, S_{\delta}\phi\lambda)| \\ &\leq \mathrm{const.} \Big\{ \epsilon^{(\gamma-1)k} S_t |\phi| \lambda^{1+\beta} + \delta S_{\delta} |\phi| \lambda^{1+\beta} + \epsilon^{d-2k} (S_{\delta} |\phi|)^2 \lambda^2 \Big\} \\ &\leq \mathrm{const.} \Big\{ \epsilon^{(\gamma-1)k} S_t \psi + \delta S_{\delta} \psi + \epsilon^{d-2k} S_{\delta} \psi \Big\}, \end{split}$$

where ψ is as above and $\lambda \leq 1/\eta$. Again we can apply (5.4.7)(ii) since $\gamma > 1$, d > 2k and $\delta \to 0$. \square

5.5. Properties of Y [Proof of Theorem (4.9)]. We have already mentioned that (4.9)(i) is proved in Dawson and Gorostiza (1988) [Remark (4.10.6)].

The results (4.9)(ii) and (4.9)(iii) are straightforward. We will give only an outline of the proofs.

The form (4.9)(ii) of the stable characteristic function is standard [see Mijnheer (1975)], and it follows easily from (4.6).

If $\Lambda_0(dx) = \lambda dx$, $\lambda > 0$, it is easy to verify that $\langle \Lambda_0, S_{t-r}(iS_r\phi)^{1+\beta} \rangle = \langle \Lambda_0, (iS_r\phi)^{1+\beta} \rangle$. The translation invariance of the stable distribution implies that $\langle \Lambda_0, (iS_r\phi)^{1+\beta} \rangle$ is invariant under translations of ϕ . By (4.6) this yields the spatial homogeneity of $Y_t - S_t Y_0$.

Now suppose that $Y_t - S_t Y_0$ is spatially homogeneous. Then, denoting $\phi_h = \phi(h + \cdot)$, $h \in \mathbb{R}^d$, by (4.6) we have

$$\int_0^t \langle \Lambda_0, S_{t-r}(iS_r \phi)^{1+\beta} \rangle dr = \int_0^t \langle \Lambda_0, S_{t-r}(iS_r \phi_h)^{1+\beta} \rangle dr,$$

and by the translation invariance of the stable distribution,

$$\int_0^t \langle \Lambda_0, S_{t-r}(iS_r\phi)^{1+eta} \rangle dr = \int_0^t \langle \Lambda_{0, h}, S_{t-r}(iS_r\phi)^{1+eta} \rangle dr,$$

where $\Lambda_{0,h}(dx) := \Lambda_0(dx - h)$. Differentiating the latter expression at t = 0 gives $\langle \Lambda_0, (i\phi)^{1+\beta} \rangle = \langle \Lambda_{0,h}, (i\phi)^{1+\beta} \rangle$. This implies that Λ_0 is invariant under translations, and hence is of the form $\Lambda_0(dx) = \lambda dx$, $\lambda > 0$.

For $\Lambda_0(dx) = \lambda dx$, $Y_0 = 0$ and $\beta = 1$, from (4.6) we have

$$E\langle Y_t, \phi \rangle^2 = 2\lambda V \overline{h} \int_0^t \int_{\mathbb{R}^d} (S_r \phi(x))^2 dx dr = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \phi(x) \phi(y) k_t(x, y) dx dy,$$

where

$$k_t(x, y) = 2\lambda V \overline{h} \int_0^t p_{2r}(y - x) dr,$$

and $p_t(y)$ denotes the stable density. The spectral density is given by

$$\sigma_t(x) = \int_{R^d} e^{-ix \cdot y} k_t(0, y) dy = \lambda V \overline{h} |x|^{-\alpha} (1 - \exp\{-2t|x|^{\alpha}\}).$$

To prove (4.9)(iv), since $\mathcal{S}(R^d)$ is a countably Hilbert nuclear space, by Hida (1980), Chapter 3, Theorem 3.1, it suffices to show that for all $t \geq 0$, the characteristic functional of $Y_t - S_t Y_0$ is continuous in some norm $\|\cdot\|_n$. By (4.6) it suffices to prove that $\int_0^t \langle \Lambda_0, S_{t-r}(S_r \phi)^{1+\beta} \rangle dr$ is continuous in ϕ in some norm $\|\cdot\|_n$. By (3.2) and $\langle \Lambda_0, \phi_p \rangle = \Box \Lambda_0 \Box_{-p}$ (Section 3.1),

$$\begin{split} \langle \Lambda_0, |S_{t-r}(S_r \phi)|^{1+\beta} \rangle &\leq \|\phi\|_{\infty}^{\beta} \langle \Lambda_0, S_t |\phi| \rangle \\ &\leq \mathrm{const.} \|\phi\|_{\infty}^{\beta} \Box \phi \Box_p \langle \Lambda_0, \phi_p \rangle \\ &= \mathrm{const.} \Box \lambda_0 \Box_{-p} \|\phi\|_{\infty}^{\beta} \Box \phi \Box_p \\ &\leq \mathrm{const.} \|\|\phi\|_m^{1+\beta}, \end{split}$$

for some $m \ge 1$, because $\|\phi\|_{\infty} = \|\|\phi\|\|_{0}$ and $\Box \phi \Box_{p} \le \|\|\phi\|\|_{m}$ for $m \ge p$. This yields the assertion for some n > m (see Section 3.2). \Box

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