Research Article

Application of Reproducing Kernel Method for Solving Nonlinear Fredholm-Volterra Integrodifferential Equations

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This paper investigates the numerical solution of nonlinear Fredholm-Volterra integro-differential equations using reproducing kernel Hilbert space method. The solution $u(x)$ is represented in the form of series in the reproducing kernel space. In the mean time, the n-term approximate solution $u_n(x)$ is obtained and it is proved to converge to the exact solution $u(x)$. Furthermore, the proposed method has an advantage that it is possible to pick any point in the interval of integration and as well the approximate solution and its derivative will be applicable. Numerical examples are included to demonstrate the accuracy and applicability of the presented technique. The results reveal that the method is very effective and simple.

1. Introduction

In recent years, there has been a growing interest in the integrodifferential equations (IDEs) which are a combination of differential and Fredholm-Volterra integral equations. IDEs are often involved in the mathematical formulation of physical phenomena. IDEs can be encountered in various fields of science such as physics, biology, and engineering. These kinds of equations can also be found in numerous applications, such as biomechanics, electromagnetic, elasticity, electrodynamics, fluid dynamics, heat and mass transfer, and oscillation theory [1–4].

The purpose of this paper is to extend the application of the reproducing kernel Hilbert space (RKHS) method to solve the nonlinear Fredholm-Volterra IDE which is as follows:

$$u'(x) + f(x)N(x, u(x)) + \int_a^b k_1(x,t)G_1(u(t))dt + \int_a^x k_2(x,t)G_2(u(t))dt = g(x), \quad a \le x, t \le b,$$

(1.1)

subject to the initial condition

$$u(a) = \alpha,$$

(1.2)

where a, b, α are real finite constants, $u \in W_2^2[a, b]$ is an unknown function to be determined, $f, g \in W_2^1[a, b]$, $k_1(x,t)$, $k_2(x,t)$ are continuous functions on $[a, b] \times [a, b]$, $G_1(w)$, $G_2(y)$, $N(x, z)$ are continuous terms in $W_2^1[a, b]$ as $w = w(x)$, $y = y(x)$, $z = z(x) \in W_2^2[a, b]$, $a \le x \le b$, $-\infty < w, y, z < \infty$ and are depending on the problem discussed, and $W_2^1[a, b], W_2^2[a, b]$ are reproducing kernel spaces.

In general, nonlinear Fredholm-Volterra IDEs do not always have solutions which we can obtain using analytical methods. In fact, many of real physical phenomena encountered are almost impossible to solve by this technique. Due to this, some authors have proposed numerical methods to approximate the solutions of nonlinear Fredholm-Volterra IDEs. To mention a few, in [5] the authors have discussed the Taylor polynomial method for solving IDEs (1.1) and (1.2) when $N(x, u(x)) = 1$, $G_1(u(t)) = u(t)$, and $G_2(u(t)) = [u(t)]^q$, where $q \in \mathbb{N}$. The triangular functions method has been applied to solve the same equations when $N(x, u(x)) = u(x)$, $G_1(u(t)) = [u(t)]^p$, and $G_2(u(t)) = [u(t)]^q$, where $p, q \in \mathbb{N}$ as described in [6]. Furthermore, the operational matrix with block-pulse functions method is carried out in [7] for the aforementioned IDEs in the case $N(x, u(x)) = 1$, $G_1(u(t)) = [u(t)]^p$, and $G_2(u(t)) = [u(t)]^q$, where $p, q \in \mathbb{N}$. Recently, the Hybrid Legendre polynomials and Block-Pulse functions approach for solving IDEs (1.1) and (1.2) when $N(x, u(x)) = 1$, $G_1(u(t)) = [u(t)]^p$, and $G_2(u(t)) = [u(t)]^q$, where $p, q \in \mathbb{N}$ are proposed in [8]. The numerical solvability of Fredholm and Volterra IDEs and other related equations can be found in [9–11] and references therein. However, none of previous studies propose a methodical way to solve these equations. Moreover, previous studies require more effort to achieve the results, they are not accurate, and usually they are developed for special types of IDEs (1.1) and (1.2).

Reproducing kernel theory has important application in numerical analysis, differential equations, integral equations, probability and statistics, and so on [12–14]. Recently, using the RKHS method, the authors in [15–29] have discussed singular linear two-point boundary value problems, singular nonlinear two-point periodic boundary value problems, nonlinear system of boundary value problems, initial value problems, singular integral equations, nonlinear partial differential equations, operator equations, and fourth-order IDEs.

The outline of the paper is as follows: several reproducing kernel spaces are described in Section 2. In Section 3, a linear operator, a complete normal orthogonal system, and some essential results are introduced. Also, a method for the existence of solutions for (1.1) and (1.2) based on reproducing kernel space is described. In Section 4, we give an iterative method to solve (1.1) and (1.2) numerically in RKHS. Various numerical examples are presented in Section 5. This paper ends in Section 6 with some concluding remarks.

2. Several Reproducing Kernel Spaces

In this section, several reproducing kernels needed are constructed in order to solve (1.1) and (1.2) using RKHS method. Before the construction, we utilize the reproducing kernel concept. Throughout this paper $\mathbb{C}$ is the set of complex numbers, $L^2[a,b] = \{u \mid \int_a^b u^2(x)dx < \infty\}$, $l^2 = \{A \mid \sum_{i=1}^{\infty} (A_i)^2 < \infty\}$, and the superscript (n) in $u^{(n)}(t)$ denotes the n-th derivative of $u(t)$.

Definition 2.1 (see [18]). Let E be a nonempty abstract set. A function $K : E \times E \to \mathbb{C}$ is a reproducing kernel of the Hilbert space H if

(1) for each $t \in E$, $K(\cdot, t) \in H$,

(2) for each $t \in E$ and $\varphi \in H$, $\langle \varphi(\cdot), K(\cdot, t) \rangle = \varphi(t)$.

The last condition is called "the reproducing property": the value of the function φ at the point t is reproducing by the inner product of φ with $K(\cdot, t)$. A Hilbert space which possesses a reproducing kernel is called a RKHS [18].

Next, we first construct the space $W_2^2[a,b]$ in which every function satisfies the initial condition (1.2) and then utilize the space $W_2^1[a,b]$.

Definition 2.2 (see [30]). $W_2^2[a,b] = \{u : u, u' \text{ are absolutely continuous on } [a,b], u, u', u'' \in L^2[a,b], \text{ and } u(a) = 0\}$. The inner product and the norm in $W_2^2[a,b]$ are defined, respectively, by

$$\langle u, v \rangle_{W_2^2} = u(a)v(a) + u'(a)v'(a) + \int_a^b u''(y)v''(y)dy \tag{2.1}$$

and $\|u\|_{W_2^2} = \sqrt{\langle u, u \rangle_{W_2^2}}$, where $u, v \in W_2^2[a,b]$.

Definition 2.3 (see [23]). $W_2^1[a,b] = \{u : u \text{ is absolutely continuous on } [a,b] \text{ and } u, u' \in L^2[a,b]\}$. The inner product and the norm in $W_2^1[a,b]$ are defined, respectively, by $\langle u, v \rangle_{W_2^1} = \int_a^b u(t)v(t) + u'(t)v'(t)dt$ and $\|u\|_{W_2^1} = \sqrt{\langle u, u \rangle_{W_2^1}}$, where $u, v \in W_2^1[a,b]$.

In [23], the authors have proved that the space $W_2^1[a,b]$ is a complete reproducing kernel space and its reproducing kernel function is given by

$$T_x(y) = \frac{1}{2\sinh(b-a)} \left[\cosh(x+y-b-a) + \cosh(|x-y|-b+a) \right]. \tag{2.2}$$

From the definition of the reproducing kernel spaces $W_2^1[a,b]$ and $W_2^2[a,b]$, we get $W_2^1[a,b] \supset W_2^2[a,b]$.

The Hilbert space $W_2^2[a,b]$ is called a reproducing kernel if for each fixed $x \in [a,b]$ and any $u(y) \in W_2^2[a,b]$, there exist $K(x,y) \in W_2^2[a,b]$ (simply $K_x(y)$) and $y \in [a,b]$ such that $\langle u(y), K_x(y) \rangle_{W_2^2} = u(x)$. The next theorem formulates the reproducing kernel function $K_x(y)$.

Theorem 2.4. *The Hilbert space $W_2^2[a,b]$ is a reproducing kernel and its reproducing kernel function $K_x(y)$ can be written as*

$$K_x(y) = \begin{cases} \sum_{i=1}^{4} p_i(x)y^{i-1}, & y \le x, \\ \sum_{i=1}^{4} q_i(x)y^{i-1}, & y > x, \end{cases} \tag{2.3}$$

where $p_i(x)$ and $q_i(x)$ are unknown coefficients of $K_x(y)$.

Proof. Through several integrations by parts for (2.1), we obtain $\langle u(y), K_x(y)\rangle_{W_2^2} = \sum_{i=0}^{1} u^{(i)}(a)(K_x^{(i)}(a) + (-1)^i K_x^{(3-i)}(a)) + \sum_{i=0}^{1}(-1)^{1-i}u^{(i)}(b)K_x^{(3-i)}(b) + \int_a^b u(y)K_x^{(4)}(y)dy$. Since $K_x(y) \in W_2^2[a,b]$, it follows that $K_x(a) = 0$. Also, since $u \in W_2^2[a,b]$, one obtains $u(a) = 0$. Thus, if $K_x^{(i)}(b) = 0$, $i = 2,3$, and $K_x'(a) - K_x''(a) = 0$, then $\langle u(y), K_x(y)\rangle_{W_2^2} = \int_a^b u(y)K_x^{(4)}(y)dy$. Now, for each $x \in [a,b]$, if $K_x(y)$ also satisfies $K_x^{(4)}(y) = \delta(x-y)$, where δ is the dirac-delta function, then $\langle u(y), K_x(y)\rangle_{W_2^2} = u(x)$. Obviously, $K_x(y)$ is the reproducing kernel function of the space $W_2^2[a,b]$.

Next, we give the expression of the reproducing kernel function $K_x(y)$. The characteristic equation of $K_x^{(4)}(y) = \delta(x-y)$ is $\lambda^4 = 0$, and their characteristic values are $\lambda = 0$ with 4 multiple roots. So, let the kernel $K_x(y)$ be as defined in (2.3).

On the other hand, let $K_x(y)$ satisfy $K_x^{(m)}(x+0) = K_x^{(m)}(x-0)$, $m = 0,1,2$. Integrating $K_x^{(4)}(y) = \delta(x-y)$ from $x-\varepsilon$ to $x+\varepsilon$ with respect to y and letting $\varepsilon \to 0$, we have the jump degree of $K_x^{(3)}(y)$ at $y = x$ given by $K_x^{(3)}(x-0) - K_x^{(3)}(x+0) = -1$. Through the last descriptions the unknown coefficients of (2.3) can be obtained. This completes the proof. ☐

By using Mathematica 7.0 software package, the representation of the reproducing kernel function $K_x(y)$ is provided by

$$K_x(y) = \begin{cases} \dfrac{1}{6}(y-a)(2a^2 - y^2 + 3x(2+y) - a(6+3x+y)), & y \le x, \\ \dfrac{1}{6}(x-a)(2a^2 - x^2 + 3y(2+x) - a(6+x+3y)), & y > x. \end{cases} \tag{2.4}$$

The following corollary summarizes some important properties of the reproducing kernel function $K_x(y)$.

Corollary 2.5. *The reproducing kernel function $K_x(y)$ is symmetric, unique, and $K_x(x) \ge 0$ for any fixed $x \in [a,b]$.*

Proof. By the reproducing property, we have $K_x(y) = \langle K_x(\xi), K_y(\xi)\rangle = \langle K_y(\xi), K_x(\xi)\rangle = K_y(x)$ for each x and y. Now, let $K_x^1(y)$ and $K_x^2(y)$ be all the reproducing kernels of the space $W_2^2[a,b]$; then $K_x^1(y) = \langle K_x^1(\xi), K_y^2(\xi)\rangle = \langle K_y^2(\xi), K_x^1(\xi)\rangle = K_y^2(x) = K_x^2(y)$. Finally, we note that $K_x(x) = \langle K_x(\xi), K_x(\xi)\rangle = \|K_x(\xi)\|^2 \ge 0$. ☐

3. Introduction to a Linear Operator and a Normal Orthogonal System in $W_2^2[a, b]$

In this section, we construct an orthogonal function system of $W_2^2[a, b]$. Also, representation of the solution of (1.1) and (1.2) is given in the reproducing kernel space $W_2^2[a, b]$.

To do this, we define a differential operator $L : W_2^2[a, b] \rightarrow W_2^1[a, b]$ such that $Lu(x) = u'(x)$. After homogenization of the initial condition (1.2), IDEs (1.1) and (1.2) can be converted into the equivalent form as follows:

$$Lu(x) = F(x, u(x), Tu(x), Su(x)), \quad a \le x \le b,$$

$$u(a) = 0,$$

(3.1)

such that $F(x, u(x), Tu(x), Su(x)) = g(x) - f(x)N(x, u(x)) - Tu(x) - Su(x)$, $Tu(x) = \int_a^b k_1(x, t)G_1(u(t))dt$, and $Su(x) = \int_a^x k_2(x, t)G_2(u(t))dt$, where $u(x) \in W_2^2[a, b]$ and $F(x, w, y, z) \in W_2^1[a, b]$ for $x \in [a, b]$ and $w = w(x)$, $y = y(x)$, $z = z(x) \in W_2^2[a, b]$, $-\infty < w, \ y, z < \infty$. It is easy to show that L is a bounded linear operator from $W_2^2[a, b]$ to $W_2^1[a, b]$.

Now, we construct an orthogonal function system of $W_2^2[a, b]$. Let $\varphi_i(x) = T_{x_i}(x)$ and $\psi_i(x) = L^*\varphi_i(x)$, where $\{x_i\}_{i=1}^{\infty}$ is dense on $[a, b]$ and L^* is the adjoint operator of L. From the properties of the reproducing kernel $T_x(y)$, we have $\langle u(x), \varphi_i(x) \rangle_{W_2^1} = \langle u(x), T_{x_i}(x) \rangle_{W_2^1} = u(x_i)$ for every $u(x) \in W_2^1[a, b]$. In terms of the properties of $K_x(y)$, one obtains $\langle u(x), \psi_i(x) \rangle_{W_2^2} = \langle u(x), L^*\varphi_i(x) \rangle_{W_2^2} = \langle Lu(x), \varphi_i(x) \rangle_{W_2^1} = Lu(x_i), i = 1, 2, \cdots$.

It is easy to see that $\psi_i(x) = L^*\varphi_i(x) = \langle L^*\varphi_i(x), K_x(y) \rangle_{W_2^2} = \langle \varphi_i(x), L_y K_x(y) \rangle_{W_2^1} = L_y K_x(y)|_{y=x_i}$. Thus, $\psi_i(x)$ can be expressed in the form $\psi_i(x) = L_y K_x(y)|_{y=x_i}$, where L_y indicates that the operator L applies to the function of y.

Theorem 3.1. *For (3.1), if $\{x_i\}_{i=1}^{\infty}$ is dense on $[a, b]$, then $\{\psi_i(x)\}_{i=1}^{\infty}$ is the complete function system of the space $W_2^2[a, b]$.*

Proof. Clearly, $\psi_i(x) \in W_2^2[a, b]$. For each fixed $u(x) \in W_2^2[a, b]$, let $\langle u(x), \psi_i(x) \rangle_{W_2^2} = 0, i = 1, 2, \ldots$, which means that $\langle u(x), \psi_i(x) \rangle_{W_2^2} = \langle u(x), L^*\varphi_i(x) \rangle_{W_2^2} = \langle Lu(x), \varphi_i(x) \rangle_{W_2^1} = Lu(x_i) = 0$. Note that $\{x_i\}_{i=1}^{\infty}$ is dense on $[a, b]$; therefore, $Lu(x) = 0$. It follows that $u(x) = 0$ from the existence of L^{-1}. So, the proof of the theorem is complete. $\square$

The orthonormal function system $\{\overline{\psi}_i(x)\}_{i=1}^{\infty}$ of the space $W_2^2[a, b]$ can be derived from Gram-Schmidt orthogonalization process of $\{\psi_i(x)\}_{i=1}^{\infty}$ as follows:

$$\overline{\psi}_i(x) = \sum_{k=1}^{i} \beta_{ik} \psi_k(x),$$

(3.2)

where β_{ik} are orthogonalization coefficients given as $\beta_{11} = 1/\|\psi_1\|$, $\beta_{ii} = 1/d_{ik}$, and $\beta_{ij} = -(1/d_{ik}) \sum_{k=j}^{i-1} c_{ik}\beta_{kj}$ for $j < i$ in which $d_{ik} = \sqrt{\|\psi_i\|^2 - \sum_{k=1}^{i-1} c_{ik}^2}$, $c_{ik} = \langle \psi_i, \overline{\psi}_k \rangle_{W_2^2}$, and $\{\psi_i(x)\}_{i=1}^{\infty}$ is the orthonormal system in the space $W_2^2[a, b]$.

Theorem 3.2. *For each $u(x) \in W_2^2[a,b]$, the series $\sum_{i=1}^{\infty} \langle u(x), \overline{\psi}_i(x) \rangle \overline{\psi}_i(x)$ is convergent in the norm of $W_2^2[a,b]$. On the other hand, if $\{x_i\}_{i=1}^{\infty}$ is dense on $[a,b]$ and $u(x)$ is the exact solution of (3.1), then*

$$u(x) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} F(x_k, u(x_k), Tu(x_k), Su(x_k)) \overline{\psi}_i(x). \tag{3.3}$$

Proof. Since $u(x) \in W_2^2[a,b]$, $\sum_{i=1}^{\infty} \langle u(x), \overline{\psi}_i(x) \rangle \overline{\psi}_i(x)$ is the Fourier series expansion about normal orthogonal system $\{\overline{\psi}_i(x)\}_{i=1}^{\infty}$, and $W_2^2[a,b]$ is the Hilbert space, then the series $\sum_{i=1}^{\infty} \langle u(x), \overline{\psi}_i(x) \rangle \overline{\psi}_i(x)$ is convergent in the sense of $\| \cdot \|_{W_2^2}$. On the other hand, using (3.2), we have

$$
\begin{aligned}
u(x) &= \sum_{i=1}^{\infty} \langle u(x), \overline{\psi}_i(x) \rangle_{W_2^2} \overline{\psi}_i(x) = \sum_{i=1}^{\infty} \left\langle u(x), \sum_{k=1}^{i} \beta_{ik} \psi_k(x) \right\rangle_{W_2^2} \overline{\psi}_i(x) \\
&= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \langle u(x), \psi_k(x) \rangle_{W_2^2} \overline{\psi}_i(x) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \langle u(x), L^* \varphi_k(x) \rangle_{W_2^2} \overline{\psi}_i(x) \\
&= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \langle Lu(x), \varphi_k(x) \rangle_{W_2^1} \overline{\psi}_i(x) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} Lu(x_k) \overline{\psi}_i(x).
\end{aligned}
\tag{3.4}
$$

But since $u(x)$ is the exact solution of (3.1), then $Lu(x) = F(x, u(x), Tu(x), Su(x))$ and

$$u(x) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} F(x_k, u(x_k), Tu(x_k), Su(x_k)) \overline{\psi}_i(x). \tag{3.5}$$

So, the proof of the theorem is complete. □

Note that we denote to the approximate solution of $u(x)$ by

$$u_n(x) = \sum_{i=1}^{n} \sum_{k=1}^{i} \beta_{ik} F(x_k, u(x_k), Tu(x_k), Su(x_k)) \overline{\psi}_i(x). \tag{3.6}$$

Theorem 3.3. *If $u(x) \in W_2^2[a,b]$, then there exists $M > 0$ such that $\|u^{(i)}(x)\|_C \leq M \|u(x)\|_{W_2^2}$, $i = 0, 1$, where $\|u(x)\|_C = \max_{a \leq x \leq b} |u(x)|$.*

Proof. For any $x, y \in [a,b]$, we have $u^{(i)}(x) = \langle u(y), K_x^{(i)}(y) \rangle_{W_2^2}$, $i = 0, 1$. By the expression of $K_x(y)$, it follows that $\|K_x^{(i)}(y)\|_{W_2^2} \leq M_i$, $i = 0, 1$. Thus, $|u^{(i)}(x)| = |\langle u(x), K_x^{(i)}(x) \rangle_{W_2^2}| \leq \|u(x)\|_{W_2^2} \|K_x^{(i)}(x)\|_{W_2^2} \leq M_i \|u(x)\|_{W_2^2}$, $i = 0, 1$. Hence, $\|u^{(i)}(x)\|_C \leq M \|u(x)\|_{W_2^2}$, $i = 0, 1$, where $M = \max\{M_0, M_1\}$. The proof is complete. □

Corollary 3.4. *The approximate solution $u_n(x)$ and its derivative $u_n'(x)$ are uniformly convergent.*

Proof. By Theorems 3.2 and 3.3, for any $x \in [a, b]$, we get

$$|u_n(x) - u(x)|$$
$$= \left| \langle u_n(x) - u(x), K_x(x) \rangle_{W_2^2} \right| \leq \|K_x(x)\|_{W_2^2} \|u_n(x) - u(x)\|_{W_2^2} \leq M_0 \|u_n(x) - u(x)\|_{W_2^2}.$$
$$(3.7)$$

On the other hand,

$$|u_n'(x) - u'(x)|$$
$$= \left| \langle u_n(x) - u(x), K_x'(x) \rangle_{W_2^2} \right| \leq \|K_x'(x)\|_{W_2^2} \|u_n(x) - u(x)\|_{W_2^2} \leq M_1 \|u_n(x) - u(x)\|_{W_2^2}.$$
$$(3.8)$$

Hence, $|u_n^{(i)}(x) - u^{(i)}(x)| \leq \|u_n^{(i)}(x) - u^{(i)}(x)\|_C \leq M_i \|u_n(x) - u(x)\|_{W_2^2}$, where M_0 and M_1 are positive constants. Hence, if $\|u_n(x) - u(x)\|_{W_2^2} \to 0$ as $n \to \infty$, the approximate solutions $u_n(x)$ and $u_n'(x)$ converge uniformly to the exact solution $u(x)$ and its derivative, respectively. □

4. Iterative Method and Convergence Theorem

In this section, an iterative method of obtaining the solution of (3.1) is presented in the reproducing kernel space $W_2^2[a, b]$.

First of all, we will mention the following remark in order to solve (1.1) and (1.2) numerically. If (1.1) is linear, then the exact and approximate solutions can be obtained directly from (3.5) and (3.6), respectively. On the other hand, if (1.1) is nonlinear, then the exact and approximate solutions can be obtained using the following iterative method.

According to (3.5), the representation of the solution of (1.1) can be denoted by

$$u(x) = \sum_{i=1}^{\infty} A_i \overline{\psi}_i(x), \qquad (4.1)$$

where $A_i = \sum_{k=1}^{i} \beta_{ik} F(x_k, u(x_k), Tu(x_k), Su(x_k))$. In fact, A_i, $i = 1, 2, \ldots$, in (4.1) are unknown, and we will approximate A_i using known B_i. For a numerical computation, we define initial function $u_0(x_1) = 0$ and the n-term approximation to $u(x)$ by

$$u_n(x) = \sum_{i=1}^{n} B_i \overline{\psi}_i(x), \qquad (4.2)$$

where the coefficients B_i are given as

$$B_1 = \beta_{11}F(x_1, u_0(x_1), Tu_0(x_1), Su_0(x_1)),$$

$$u_1(x) = B_1\overline{\psi}_1(x),$$

$$B_2 = \sum_{k=1}^{2}\beta_{2k}F(x_k, u_{k-1}(x_k), Tu_{k-1}(x_k), Su_{k-1}(x_k)),$$

$$u_2(x) = \sum_{i=1}^{2}B_i\overline{\psi}_i(x),$$

$$\vdots$$

$$u_{n-1}(x) = \sum_{i=1}^{n-1}B_i\overline{\psi}_i(x),$$

$$B_n = \sum_{k=1}^{n}\beta_{nk}F(x_k, u_{k-1}(x_k), Tu_{k-1}(x_k), Su_{k-1}(x_k)).$$

(4.3)

We mention here the following remark: in the iterative process of (4.2), we can guarantee that the approximation $u_n(x)$ satisfies the initial condition (1.2).

Now, the approximate solution $u_n^N(x)$ can be obtained by taking finitely many terms in the series representation of $u_n(x)$ and

$$u_n^N(x) = \sum_{i=1}^{N}\sum_{k=1}^{i}\beta_{ik}F(x_k, u_{n-1}(x_k), Tu_{n-1}(x_k), Su_{n-1}(x_k))\overline{\psi}_i(x).$$

(4.4)

Next, we will prove that $u_n(x)$ in the iterative formula (4.2) is convergent to the exact solution $u(x)$ of (1.1).

Lemma 4.1. *If* $\|u_n(x) - u(x)\|_{W_2^2} \to 0$, $x_n \to y$ *as* $n \to \infty$ *and* $F(x, v, w, z)$ *is continuous in* $[a, b]$ *with respect to* x, v, w, z, *for* $x \in [a, b]$ *and* $v, w, z \in (-\infty, \infty)$, *then as* $n \to \infty$, *one has* $F(x_n, u_{n-1}(x_n), Tu_{n-1}(x_n), Su_{n-1}(x_n)) \to F(y, u(y), Tu(y), Su(y))$.

Proof. Firstly, we will prove that $u_{n-1}(x_n) \to u(y)$ in the sense of $\|\cdot\|_{W_2^2}$. Since

$$\begin{aligned}\left|u_{n-1}(x_n) - u(y)\right| &= \left|u_{n-1}(x_n) - u_{n-1}(y) + u_{n-1}(y) - u(y)\right| \\ &\leq \left|u_{n-1}(x_n) - u_{n-1}(y)\right| + \left|u_{n-1}(y) - u(y)\right|,\end{aligned}$$

(4.5)

by reproducing kernel property of $K_x(y)$, we have $u_{n-1}(x_n) = \langle u_{n-1}(x), K_{x_n}(x)\rangle$ and $u_{n-1}(y) = \langle u_{n-1}(x), K_y(x)\rangle$. Thus, $|u_{n-1}(x_n) - u_{n-1}(y)| = |\langle u_{n-1}(x), K_{x_n}(x) - K_y(x)\rangle_{W_2^2}| \leq \|u_{n-1}(x)\|_{W_2^2}$ $\|K_{x_n}(x) - K_y(x)\|_{W_2^2}$. From the symmetry of $K_x(y)$, it follows that $\|K_{x_n}(x) - K_y(x)\|_{W_2^2} \to 0$ as $n \to \infty$. Hence, $|u_{n-1}(x_n) - u_{n-1}(y)| \to 0$ as soon as $x_n \to y$.

On the other hand, by Corollary 3.4, for any $y \in [a, b]$, it holds that $|u_{n-1}(y) - u(y)| \to 0$. Therefore, $u_{n-1}(x_n) \to u(y)$ in the sense of $\|\cdot\|_{W_2^2}$ as $x_n \to y$ and $n \to \infty$.

Thus, by means of the continuation of G_1, G_2, and N, it is obtained that $G_1(u_{n-1}(x_n)) \rightarrow G_1(u(y))$, $G_2(u_{n-1}(x_n)) \rightarrow G_2(u(y))$, and $N(x_n, u_{n-1}(x_n)) \rightarrow N(y, u(y))$ as $n \rightarrow \infty$. This shows that $Tu_{n-1}(x_n) \rightarrow Tu(y)$ and $Su_{n-1}(x_n) \rightarrow Su(y)$ as $n \rightarrow \infty$. Hence, the continuity of F gives the result. □

Lemma 4.2. $\{u_n\}_{n=1}^{\infty}$ *in* (4.2) *is monotonically increasing in the sense of the norm of* $W_2^2[a, b]$.

Proof. By Theorem 3.1, $\{\overline{\psi}_i\}_{i=1}^{\infty}$ is the complete orthonormal system in the space $W_2^2[a, b]$. Hence, we have $\|u_n\|_{W_2^2}^2 = \langle u_n(x), u_n(x)\rangle_{W_2^2} = \langle \sum_{i=1}^{n} B_i\overline{\psi}_i(x), \sum_{i=1}^{n} B_i\overline{\psi}_i(x)\rangle_{W_2^2} = \sum_{i=1}^{n} (B_i)^2$. Therefore, $\|u_n\|_{W_2^2}$ is monotonically increasing. □

Lemma 4.3. *One has* $Lu_n(x_j) = F(x_j, u_{j-1}(x_j), Tu_{j-1}(x_j), Su_{j-1}(x_j)), j \leq n$.

Proof. The proof will be obtained by mathematical induction as follows: if $j \leq n$, then $Lu_n(x_j) = \sum_{i=1}^{n} B_i L\overline{\psi}_i(x_j) = \sum_{i=1}^{n} B_i \langle L\overline{\psi}_i(x), \varphi_j(x)\rangle_{W_2^1} = \sum_{i=1}^{n} B_i \langle \overline{\psi}_i(x), L_j^*\varphi(x)\rangle_{W_2^2} = \sum_{i=1}^{n} B_i \langle \overline{\psi}_i(x), \psi_j(x)\rangle_{W_2^2}$. Thus,

$$Lu_n(x_j) = \sum_{i=1}^{n} B_i \langle \overline{\psi}_i(x), \psi_j(x)\rangle_{W_2^2}. \tag{4.6}$$

Multiplying both sides of (4.6) by β_{jl}, summing for l from 1 to j, and using the orthogonality of $\{\overline{\psi}_i(x)\}_{i=1}^{\infty}$ yield that

$$\sum_{l=1}^{j} \beta_{jl} Lu_n(x_l) = \sum_{i=1}^{n} B_i \left\langle \overline{\psi}_i(x), \sum_{l=1}^{j} \beta_{jl}\psi_l(x)\right\rangle_{W_2^2} = \sum_{i=1}^{n} B_i \left\langle \overline{\psi}_i(x), \overline{\psi}_j(x)\right\rangle_{W_2^2} = B_j$$

$$= \sum_{l=1}^{j} \beta_{jl} F(x_l, u_{l-1}(x_l), Tu_{l-1}(x_l), Su_{l-1}(x_l)). \tag{4.7}$$

Now, if $j = 1$, then $Lu_n(x_1) = F(x_1, u_0(x_1), Tu_0(x_1), Su_0(x_1))$. On the other hand, if $j = 2$, then $\beta_{21}Lu_n(x_1) + \beta_{22}Lu_n(x_2) = \beta_{21}F(x_1, u_0(x_1), Tu_0(x_1), Su_0(x_1)) + \beta_{22}F(x_2, u_1(x_2), Tu_1(x_2), Su_1(x_2))$. Thus, $Lu_n(x_2) = F(x_2, u_1(x_2), Tu_1(x_2), Su_1(x_2))$. It is easy to see that $Lu_n(x_j) = F(x_j, u_{j-1}(x_j), Tu_{j-1}(x_j), Su_{j-1}(x_j))$ by using mathematical induction. □

Lemma 4.4. *One has* $Lu_n(x_j) = Lu(x_j), j \leq n$.

Proof. It is clear that on taking limits in (4.2) $u(x) = \sum_{i=1}^{\infty} B_i\overline{\psi}_i(x)$. Therefore, $u_n(x) = P_n u(x)$, where P_n is an orthogonal projector from $W_2^2[a, b]$ to Span$\{\psi_1, \psi_2, \ldots, \psi_n\}$. Thus, $Lu_n(x_j) = \langle Lu_n(x), \varphi_j(x)\rangle_{W_2^1} = \langle u_n(x), L_j^*\varphi(x)\rangle_{W_2^2} = \langle P_n u(x), \psi_j(x)\rangle_{W_2^2} = \langle u(x), P_n\psi_j(x)\rangle_{W_2^2} = \langle u(x), \psi_j(x)\rangle_{W_2^2} = \langle Lu(x), \varphi_j(x)\rangle_{W_2^1} = Lu(x_j)$. □

Theorem 4.5. *Suppose that* $\|u_n\|_{W_2^2}$ *is bounded in* (4.2). *If* $\{x_i\}_{i=1}^{\infty}$ *is dense on* $[a, b]$, *then the n-term approximate solution* $u_n(x)$ *in the iterative formula* (4.2) *converges to the exact solution* $u(x)$ *of* (3.1) *in the space* $W_2^2[a, b]$ *and* $u(x) = \sum_{i=1}^{\infty} B_i\overline{\psi}_i(x)$, *where* B_i *is given by* (4.3).

Proof. First of all, we will prove the convergence of $u_n(x)$. From (4.2), we infer that $u_{n+1}(x) = u_n(x) + B_{n+1}\overline{\psi}_{n+1}(x)$. The orthogonality of $\{\overline{\psi}_i(x)\}_{i=1}^{\infty}$ yields that

$$\|u_{n+1}\|_{W_2^2}^2 = \|u_n\|_{W_2^2}^2 + (B_{n+1})^2 = \|u_{n-1}\|_{W_2^2}^2 + (B_n)^2 + (B_{n+1})^2 = \cdots = \|u_0\|_{W_2^2}^2 + \sum_{i=1}^{n+1}(B_i)^2. \quad (4.8)$$

From Lemma 4.2, the sequence $\|u_n\|_{W_2^2}$ is monotonically increasing. Due to the condition that $\|u_n\|_{W_2^2}$ is bounded, $\|u_n\|_{W_2^2}$ is convergent as $n \to \infty$. Then, there exists a constant c such that $\sum_{i=1}^{\infty}(B_i)^2 = c$. This implies that $B_i = \sum_{k=1}^{i}\beta_{ik}F(x_k, u_{k-1}(x_k), Tu_{k-1}(x_k), Su_{k-1}(x_k)) \in l^2$, $i = 1, 2, \ldots$.

If $m > n$, using $(u_m - u_{m-1}) \perp (u_{m-1} - u_{m-2}) \perp \cdots \perp (u_{n+1} - u_n)$, then one gets

$$\|u_m(x) - u_n(x)\|_{W_2^2}^2 = \|u_m(x) - u_{m-1}(x) + u_{m-1}(x) - \cdots + u_{n+1}(x) - u_n(x)\|_{W_2^2}^2$$
$$= \|u_m(x) - u_{m-1}(x)\|_{W_2^2}^2 + \cdots + \|u_{n+1}(x) - u_n(x)\|_{W_2^2}^2. \quad (4.9)$$

Furthermore, $\|u_m(x) - u_{m-1}(x)\|_{W_2^2}^2 = (B_m)^2$. Consequently, $\|u_m(x) - u_n(x)\|_{W_2^2}^2 = \sum_{i=n+1}^{m}(B_i)^2 \to 0$ as $n, m \to \infty$. Considering the completeness of the space $W_2^2[a, b]$, there exists a $u(x) \in W_2^2[a, b]$ such that $u_n(x) \to u(x)$ as $n \to \infty$ in the sense of $\|\cdot\|_{W_2^2}$.

Secondly, we will prove that $u(x)$ is the solution of (3.1). Since $\{x_i\}_{i=1}^{\infty}$ is dense on $[a, b]$, for any $x \in [a, b]$, there exists subsequence $\{x_{n_j}\}_{j=1}^{\infty}$ such that $x_{n_j} \to x$ as $j \to \infty$. From Lemmas 4.3 and 4.4, it is easy to see that $Lu(x_{n_j}) = F(x_{n_j}, u_{n_j-1}(x_k), Tu_{n_j-1}(x_k), Su_{n_j-1}(x_k))$. Hence, letting $j \to \infty$, by Lemma 4.1 and the continuity of F, we have $Lu(x) = F(x, u(x), Tu(x), Su(x))$. That is, $u(x)$ is the solution of (3.1).

Since $\overline{\psi}_i(x) \in W_2^2[a, b]$, clearly, $u(x)$ satisfies the initial condition (1.2). In other words, $u(x)$ is the solution of (1.1) and (1.2), where $u(x) = \sum_{i=1}^{\infty}B_i\overline{\psi}_i(x)$ and B_i is given by (4.3). The proof is complete. $\qquad\square$

Theorem 4.6. *Assume that $u(x)$ is the solution of (3.1) and $r_n(x)$ is the difference between the approximate solution $u_n(x)$ and the exact solution $u(x)$. Then, $r_n(x)$ is monotonically decreasing in the sense of the norm of $W_2^2[a, b]$.*

Proof. It obvious that $\|r_n(x)\|_{W_2^2}^2 = \|u(x) - u_n(x)\|_{W_2^2}^2 = \|\sum_{i=n+1}^{\infty}B_i\overline{\psi}_i(x)\|_{W_2^2}^2 = \sum_{i=n+1}^{\infty}(B_i)^2$ and $\|r_{n-1}(x)\|_{W_2^2}^2 = \sum_{i=n}^{\infty}(B_i)^2$. Thus, $\|r_n(x)\|_{W_2^2} \le \|r_{n-1}(x)\|_{W_2^2}$; consequently, the difference $r_n(x)$ is monotonically decreasing in the sense of $\|\cdot\|_{W_2^2}$. So, the proof of the theorem is complete. $\qquad\square$

5. Numerical Examples

In this section, some numerical examples are studied to demonstrate the accuracy and applicability of the present method. Results obtained are compared with the exact solution of each example and are found to be in good agreement with each other. In the process of computation, all the symbolic and numerical computations were performed by using Mathematica 7.0 software package.

Table 1: Numerical results for Example 5.1.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.16	0.0256	0.025600985200558	9.85201×10^{-7}	3.84844×10^{-5}
0.32	0.1024	0.102401995676314	1.99568×10^{-6}	1.94890×10^{-5}
0.48	0.2304	0.230403057350893	3.05735×10^{-6}	1.32698×10^{-5}
0.64	0.4096	0.409604197461425	4.19746×10^{-6}	1.02477×10^{-5}
0.80	0.6400	0.640005445257301	5.44526×10^{-6}	8.50821×10^{-6}
0.96	0.9216	0.921606832750570	6.83275×10^{-6}	7.41401×10^{-6}

Table 2: Numerical results of $u'(x)$ for Example 5.1.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.16	0.32	0.319999014799438	9.85201×10^{-7}	3.07875×10^{-6}
0.32	0.64	0.639998004323680	1.99568×10^{-6}	3.11824×10^{-6}
0.48	0.96	0.959996942649100	3.05735×10^{-6}	3.18474×10^{-6}
0.64	1.28	1.279995802538568	4.19746×10^{-6}	3.27927×10^{-6}
0.80	1.60	1.599994554742679	5.44526×10^{-6}	3.40329×10^{-6}
0.96	1.92	1.919993167249420	6.83275×10^{-6}	3.55872×10^{-6}

Example 5.1. Consider the nonlinear Fredholm-Volterra IDE:

$$u'(x) + u(x) - \frac{1}{4}\int_0^1 t u^3(t)dt + \frac{1}{2}\int_0^x x u^2(t)dt = g(x), \quad 0 \le x, t \le 1,$$

$$u(0) = 0, \tag{5.1}$$

where $g(x) = (1/10)x^6 + x^2 + 2x - (1/32)$. The exact solution is $u(x) = x^2$.

Using RKHS method, taking $x_i = (i-1)/(N-1)$, $i = 1, 2, \ldots, N$, with the reproducing kernel function $K_x(y)$ on $[0,1]$, the approximate solution $u_n^N(x)$ is calculated by (4.4). The numerical results at some selected grid points for $N = 26$ and $n = 5$ are given in Table 1.

As we mention, we used the grid nodes mentioned earlier in order to obtain approximate solutions. Moreover, it is possible to pick any point in $[a, b]$ and as well the approximate solution and its derivative will be applicable. Next, the numerical results for Example 5.1 at some selected gird nodes in $[0,1]$ of $u'(x)$ are given in Table 2.

Table 3 shows, a comparison between the absolute errors of our method together with triangular functions method [6], operational matrix with block-pulse functions method [7], and Hybrid Legendre polynomials and block-pulse functions method [8]. As it is evident from the comparison results, it was found that our method in comparison with the mentioned methods is better with a view to accuracy and utilization.

Example 5.2. Consider the nonlinear Fredholm-Volterra IDE:

$$u'(x) + 2x u(x) - \int_0^1 (x - t)u(t)dt - \int_0^x (x + t)u^3(t)dt = g(x), \quad 0 \le x, t \le 1,$$

$$u(0) = 1, \tag{5.2}$$

Table 3: Numerical comparison of absolute error for Example 5.1.

x	Method in [6]	Method in [7]	Method in [8]	RKHS method
0.0	0	0	0	0
0.1	1.6600×10^{-4}	2.1840×10^{-3}	3.1000×10^{-5}	6.1415×10^{-7}
0.2	2.5400×10^{-4}	1.4630×10^{-3}	7.5000×10^{-5}	1.2345×10^{-6}
0.3	2.6200×10^{-4}	1.6780×10^{-3}	1.7100×10^{-4}	1.8671×10^{-6}
0.4	1.9100×10^{-4}	7.2390×10^{-3}	9.4000×10^{-5}	2.5185×10^{-6}
0.5	4.1000×10^{-5}	1.6024×10^{-2}	2.2800×10^{-4}	3.1950×10^{-6}
0.6	2.0200×10^{-4}	1.1480×10^{-2}	5.0200×10^{-4}	3.9035×10^{-6}
0.7	2.8300×10^{-4}	4.5150×10^{-3}	5.8300×10^{-4}	4.6511×10^{-6}
0.8	2.8300×10^{-4}	4.8720×10^{-3}	3.7400×10^{-4}	5.4453×10^{-6}
0.9	2.0200×10^{-4}	1.6682×10^{-2}	4.7000×10^{-5}	6.2939×10^{-6}
1.0	3.7000×10^{-5}	3.0914×10^{-2}	1.4000×10^{-5}	7.2055×10^{-6}

Table 4: Numerical results for Example 5.2.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.0	1	1	0	0
0.1	1.105170918075648	1.105171575587132	6.57511×10^{-7}	5.94941×10^{-7}
0.2	1.221402758160170	1.221404273302208	1.51514×10^{-6}	1.24049×10^{-6}
0.3	1.349858807576003	1.349861385895122	2.57832×10^{-6}	1.91007×10^{-6}
0.4	1.491824697641270	1.491828549248914	3.85161×10^{-6}	2.58181×10^{-6}
0.5	1.648721270700128	1.648726613238682	5.34254×10^{-6}	3.24041×10^{-6}
0.6	1.822118800390509	1.822125866868782	7.06648×10^{-6}	3.87817×10^{-6}
0.7	2.013752707470476	2.013761760093792	9.05262×10^{-6}	4.49540×10^{-6}
0.8	2.225540928492467	2.225552279886237	1.13514×10^{-5}	5.10051×10^{-6}
0.9	2.459603111156949	2.459617154977751	1.40438×10^{-5}	5.70979×10^{-6}
1.0	2.718281828459045	2.718299082481196	1.72540×10^{-5}	6.34740×10^{-6}

where $g(x) = (-(2/3)x + (1/9))e^{3x} + (2x + 1)e^x + ((4/3) - e)x + (8/9)$. The exact solution is $u(x) = e^x$.

Using RKHS method, taking $x_i = (i - 1)/(N - 1)$, $i = 1, 2, \ldots, N$, with the reproducing kernel function $K_x(y)$ on $[0, 1]$, the approximate solution $u_n^N(x)$ is calculated by (4.4). The numerical results at some selected grid points for $N = 26$ and $n = 1$ are given in Table 4.

The comparison among the RKHS solution besides the solutions of triangular functions [6], operational matrix with block-pulse functions solution [7], and exact solutions are shown in Table 5.

Example 5.3. Consider the nonlinear Fredholm-Volterra IDE:

$$u'(x) + u^2(x) - \int_0^1 (t + 1) \sinh(u(t) - 1)dt - \int_0^x x e^{u(t)} dt = g(x), \quad 0 \le x, t \le 1,$$

$$u(0) = 1,$$

(5.3)

Table 5: Numerical comparison of approximate solution for Example 5.2.

x	Exact solution	RKHS method	Method in [6]	Method in [7]
0.0	1.000000	1.000000	1.000000	1.000000
0.1	1.105171	1.105172	1.105223	1.115708
0.2	1.221403	1.221404	1.221494	1.225369
0.3	1.349859	1.349861	1.349971	1.345821
0.4	1.491825	1.491829	1.491933	1.478128
0.5	1.648721	1.648727	1.648795	1.675019
0.6	1.822119	1.822126	1.822484	1.839752
0.7	2.013753	2.013762	2.014465	2.020740
0.8	2.225541	2.225552	2.226719	2.219630
0.9	2.459603	2.459617	2.461507	2.438278
1.0	2.718280	2.718300	2.721505	2.678832

Table 6: Numerical results for Example 5.3.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.16	1.148420005118273	1.148420410550930	4.05433×10^{-7}	3.53035×10^{-7}
0.32	1.277631736598280	1.277632408078983	6.71481×10^{-7}	5.25567×10^{-7}
0.48	1.392042087776024	1.392042943683538	8.55908×10^{-7}	6.14857×10^{-7}
0.64	1.494696241836107	1.494697231133553	9.89297×10^{-7}	6.61872×10^{-7}
0.80	1.587786664902119	1.587787753995242	1.08909×10^{-6}	6.85919×10^{-7}
0.96	1.672944473242426	1.672945638649806	1.16541×10^{-6}	6.96620×10^{-7}

where $g(x) = (\ln(x + 1) + 1)^2 + (x + 1)^{-1} - (e/2)(x + 2)x^2 - (2/3)$. The exact solution is $u(x) = 1 + \ln(x + 1)$.

Using RKHS method, taking $x_i = (i - 1)/(N - 1)$, $i = 1, 2, \ldots, N$, with the reproducing kernel function $K_x(y)$ on $[0, 1]$, the approximate solution $u_n^N(x)$ is calculated by (4.4). The numerical results at some selected grid points for $N = 51$ and $n = 1$ are given in Table 6.

Example 5.4. Consider the nonlinear Fredholm-Volterra IDE:

$$u'(x) + \cos(u(x)) - \int_0^1 (t + x)u(t)dt - \int_0^x te^{u(t)}dt = g(x), \quad 0 \le x, t \le 1,$$

$$u(0) = 1,$$

(5.4)

where $g(x) = \cos(1 - x) + e^{1-x}(1 + x) - (1/2)x - (7/6) - e$. The exact solution is $u(x) = 1 - x$.

Using RKHS method, taking $x_i = (i - 1)/(N - 1)$, $i = 1, 2, \ldots, N$, with the reproducing kernel function $K_x(y)$ on $[0, 1]$, the approximate solution $u_n^N(x)$ is calculated by (4.4). The numerical results at some selected grid points for $N = 51$ and $n = 5$ are given in Table 7.

Table 7: Numerical results for Example 5.4.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.16	0.84	0.840000573907808	5.73908×10^{-7}	6.83224×10^{-7}
0.32	0.68	0.680001166402260	1.16640×10^{-6}	1.71530×10^{-6}
0.48	0.52	0.520001763060855	1.76306×10^{-6}	3.39050×10^{-6}
0.64	0.36	0.360002349461090	2.34946×10^{-6}	6.52628×10^{-6}
0.80	0.20	0.200002911180463	2.91118×10^{-6}	1.45559×10^{-5}
0.96	0.04	0.040003433796472	3.43380×10^{-6}	8.58449×10^{-5}

Table 8: Numerical results for Example 5.5.

x	Exact solution	Approximate solution	Absolute error	Relative error
0.16	0.16	0.159999994448885	5.55112×10^{-9}	3.46945×10^{-8}
0.32	0.32	0.319999983346655	1.66533×10^{-8}	5.20417×10^{-8}
0.48	0.48	0.479999983346655	1.66533×10^{-8}	3.46945×10^{-8}
0.64	0.64	0.639999777955395	2.22045×10^{-7}	3.46945×10^{-7}
0.80	0.80	0.799999555910790	4.44089×10^{-7}	5.55112×10^{-7}
0.96	0.96	0.959999555910790	4.44089×10^{-7}	4.62593×10^{-7}

Example 5.5. Consider the nonlinear Fredholm-Volterra IDE:

$$u'(x) + xu^2(x) - \int_0^1 (t+x)\left(1 + u^2(t)\right)dt - \int_0^x x\cos(u(t))dt = g(x), \quad 0 \le x, t \le 1,$$

$$u(0) = 0, \tag{5.5}$$

where $g(x) = -x\sin(x) + x^3 - (4/3)x + (1/4)$. The exact solution is $u(x) = x$.

Using RKHS method, taking $x_i = (i-1)/(N-1)$, $i = 1, 2, \ldots, N$ with the reproducing kernel function $K_x(y)$ on $[0,1]$, the approximate solution $u_n^N(x)$ is calculated by (4.4). The numerical results at some selected grid points for $N = 26$ and $n = 5$ are given in Table 8.

6. Conclusion

In this paper, the RKHS method was employed to solve the nonlinear Fredholm-Volterra IDEs (1.1) and (1.2). The solution $u(x)$ and the approximate solution $u_n(x)$ are represented in the form of series in the space $W_2^2[a, b]$. Moreover, the approximate solution and its derivative converge uniformly to the exact solution and its derivative, respectively. Meanwhile, the error of the approximate solution is monotonically decreasing in the sense of the norm of $W_2^2[a, b]$.

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