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# ON THE GAUSS MAP OF TRANSLATION SURFACES IN MINKOWSKI 3-SPACE 

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#### Abstract

In this article, we study translation surfaces in the 3-dimensional Minkowski space whose Gauss map $G$ satisfies the condition $\Delta G=A G, A \in$ $\operatorname{Mat}(3, \mathbb{R})$, where $\Delta$ denotes the Laplacian of the surface with respect to the induced metric and $\operatorname{Mat}(3, \mathbb{R})$ the set of $3 \times 3$ real matrices, and also obtain the complete classification theorem for those.


## 1. Introduction

As is well-known, the theory of Gauss map is always one of interesting topics in Euclidean space and pseudo-Euclidean space and it has been investigated from the various viewpoints by many differential geometers [1, 2, 4, 7, 8, 10, 11].
F. Dillen, J. Pas and L. Verstraelen [10] studied surfaces of revolution in Euclidean 3 -space $\mathbb{E}^{3}$ such that its Gauss map $G$ satisfies the condition

$$
\begin{equation*}
\Delta G=A G, \quad A=\left(a_{i j}\right) \in \operatorname{Mat}(3, \mathbb{R}) \tag{1.1}
\end{equation*}
$$

where $\Delta$ denotes the Laplacian of the surface with respect to the induced metric and $\operatorname{Mat}(3, \mathbb{R})$ the set of $3 \times 3$ real matrices. On the other hand, C. Baikoussis and D. E. Blair [3] investigated the ruled surfaces in $\mathbb{E}^{3}$ satisfying the condition (1.1). C. Baikoussis and L. Verstraelen [4, 5, 6] studied the helicoidal surfaces, the translation surfaces and the spiral surfaces in $\mathbb{E}^{3}$ satisfying the condition (1.1). Also, for the Lorentz version, S. M. Choi [8, 9] completely classified the surfaces of revolution and the ruled surfaces with non-null base curve satisfying the condition (1.1) in Minkowski 3 -space $\mathbb{E}_{1}^{3}$. Furthermore, L. J. Alías, A. Ferrández, P. Lucas and M. A. Meroño [2] studied the ruled surfaces with null ruling satisfying the condition (1.1) in Minkowski 3 -space $\mathbb{E}_{1}^{3}$. On the other hand, condition (1.1) is a

[^0]special case of a finite type Gauss map introduced by B. Y. Chen [7]. Recently, Y. H. Kim and the author [13] studied the ruled surfaces with pointwise 1-type Gauss map in $\mathbb{E}_{1}^{3}$ and obtained a new characterization of minimal ruled surfaces. In [11], D.-S. Kim, Y. H. Kim and the author obtained the complete classification theorem of ruled surfaces with 1-type Gauss map in Minkowski $m$-space $\mathbb{E}_{1}^{m}$ and also characterized the extended $B$-scroll with Gauss map.

In this article, we investigate the Lorentz version of the translation surfaces satisfying condition (1.1) and prove the following theorem:

Theorem. The only translation surfaces in Minkowski 3 -space $\mathbb{E}_{1}^{3}$ whose Gauss map satisfies (1.1) are the Euclidean plane $\mathbb{R}^{2}$, the Minkowski plane $\mathbb{R}_{1}^{2}$, the Lorentz circular cylinder $\mathbb{S}_{1}^{1} \times \mathbb{R}$, the hyperbolic cylinder $\mathbb{H}^{1} \times \mathbb{R}$ and the circular cylinder of index $1, \mathbb{R}_{1}^{1} \times \mathbb{S}^{1}$.

To prove this theorem, we use the reasoning first developed by C. Baikoussis and L. Verstraelen in [5], in which they classified translation surfaces satisfying the condition (1.1) in $\mathbb{E}^{3}$.

For the study of the translation surfaces in Minkowski 3 -space $\mathbb{E}_{1}^{3}$, I. V. de Woestijne [15] studied minimal translation surfaces, and H. Liu [14] investigated the translation surfaces with constant mean curvature or constant Gauss curvature.

Throughout this paper, we assume that all objects are smooth and all surfaces are pseudo-Riemannian, unless otherwise specified.

## 2. Preliminaries

An $m$-dimensional vector space $L=L_{1}^{m}$ with scalar product $\langle$,$\rangle of index 1$ is called a Lorentz vector space. In particular, if $L=\mathbb{E}_{1}^{m}, m \geq 2$, it is called a Minkowski $m$-space. A vector $X$ of $L_{1}^{m}$ is said to be space-like if $\langle X, X\rangle>0$ or $X=0$, time-like if $\langle X, X\rangle<0$ and light-like or null if $\langle X, X\rangle=0$ and $X \neq 0$. A curve in $L_{1}^{m}$ is called space-like (time-like or null, respectively) if its tangent vector is space-like (time-like or null, respectively).

Let $X=\left(X_{i}\right)$ and $Y=\left(Y_{i}\right)$ be the vectors in a 3-dimensional Lorentz vector space $L_{1}^{3}$. Then the scalar product of $X$ and $Y$ is defined by

$$
\begin{equation*}
\langle X, Y\rangle=-X_{1} Y_{1}+X_{2} Y_{2}+X_{3} Y_{3} \tag{2.1}
\end{equation*}
$$

which is called a Lorentz product. Furthermore, a Lorentz cross product $X \times Y$ is given by

$$
\begin{equation*}
X \times Y=\left(-X_{2} Y_{3}+X_{3} Y_{2}, X_{3} Y_{1}-X_{1} Y_{3}, X_{1} Y_{2}-X_{2} Y_{1}\right) \tag{2.2}
\end{equation*}
$$

Let $M^{2}$ be a pseudo-Riemannian surface in Minkowski 3 -space $\mathbb{E}_{1}^{3}$. The map $G: M^{2} \longrightarrow Q^{2}(\varepsilon) \subset \mathbb{E}_{1}^{3}$ which sends each point of $M^{2}$ to the unit normal vector
to $M^{2}$ at the point is called the Gauss map of surface $M^{2}$, where $\varepsilon(= \pm 1)$ denotes the sign of the vector field $G$ and $Q^{2}(\varepsilon)$ is a 2-dimensional space form as follows:

$$
Q^{2}(\epsilon)= \begin{cases}\mathbb{S}_{1}^{2}(1)=\left\{X \in \mathbb{E}_{1}^{3} \mid\langle X, X\rangle=1\right\} \quad \text { if } \quad \epsilon=1 \\ \mathbb{H}^{2}(-1)=\left\{X \in \mathbb{E}_{1}^{3} \mid\langle X, X\rangle=-1\right\} \quad \text { if } \quad \epsilon=-1\end{cases}
$$

$\mathbb{S}_{1}^{2}(1)$ is called the de Sitter space, $\mathbb{H}^{2}(-1)$ the hyperbolic space in $\mathbb{E}_{1}^{3}$. It is wellknown that in terms of local coordinates $\left\{x_{i}\right\}$ of $M^{2}$, the Laplacian can be written as:

$$
\begin{equation*}
\Delta=-\frac{1}{\sqrt{|\mathcal{G}|}} \sum_{i, j} \frac{\partial}{\partial x^{i}}\left(\sqrt{|\mathcal{G}|} g^{i j} \frac{\partial}{\partial x^{j}}\right) \tag{2.3}
\end{equation*}
$$

where $\mathcal{G}=\operatorname{det}\left(g_{i j}\right),\left(g^{i j}\right)=\left(g_{i j}\right)^{-1}$ and $\left(g_{i j}\right)$ are the components of the metric of $M^{2}$ with respect to $\left\{x_{i}\right\}$.

## 3. Translation Surfaces in Minkowski 3-Spaces

Let $x: M^{2} \longrightarrow \mathbb{E}_{1}^{3}$ be a translation surface in $\mathbb{E}_{1}^{3}$. Then $M^{2}$ is parametrized by

$$
\begin{equation*}
x(u, v)=(u, v, \tilde{f}(u)+\tilde{g}(v)) \tag{3.1}
\end{equation*}
$$

$\tilde{f}$ and $\tilde{g}$ being smooth functions of the variables $u$ and $v$, respectively, and we have the natural frame $\left\{x_{u}, x_{v}\right\}$ given by

$$
x_{u}=\frac{\partial x}{\partial u}=(1,0, f), \quad x_{v}=\frac{\partial x}{\partial v}=(0,1, g),
$$

where $f=d \tilde{f} / d u, g=d \tilde{g} / d v$. Accordingly, the induced pseudo-Riemannian metric on $M^{2}$ is obtained by $g_{11}=\left\langle x_{u}, x_{u}\right\rangle=f^{2}-1, g_{12}=\left\langle x_{u}, x_{v}\right\rangle=f g$ and $g_{22}=$ $\left\langle x_{v}, x_{v}\right\rangle=1+g^{2}$. Since the surface is non-degenerate, $\operatorname{det}\left(g_{i j}\right)=f^{2}-g^{2}-1 \neq 0$. For later use, we define smooth function $\omega$ as:

$$
\begin{equation*}
\omega=\left\|x_{u} \times x_{v}\right\|^{2}=\varepsilon\left\langle x_{u} \times x_{v}, x_{u} \times x_{v}\right\rangle=\varepsilon\left(-f^{2}+g^{2}+1\right) \tag{3.2}
\end{equation*}
$$

where $\varepsilon$ denotes the sign of the vector $x_{u} \times x_{v}$ in $\mathbb{E}_{1}^{3}$. Then the Gauss map $G$ of the surface $M^{2}$ is given by

$$
\begin{equation*}
G=\left(G_{1}, G_{2}, G_{3}\right)=\frac{1}{\left\|x_{u} \times x_{v}\right\|} x_{u} \times x_{v}=\frac{1}{\omega^{\frac{1}{2}}}(f,-g, 1) . \tag{3.3}
\end{equation*}
$$

If we make use of (2.3) together with such function $\omega$, the Laplacian $\Delta$ of $M^{2}$ can be expressed as follows:

$$
\begin{align*}
\Delta= & \frac{1}{\omega}\left\{\varepsilon\left(1+g^{2}\right) \frac{\partial^{2}}{\partial u^{2}}+\varepsilon\left(f^{2}-1\right) \frac{\partial^{2}}{\partial v^{2}}-2 \varepsilon f g \frac{\partial^{2}}{\partial u \partial v}\right.  \tag{3.4}\\
& \left.+\frac{\left(f^{2}-1\right) g^{\prime}+\left(1+g^{2}\right) f^{\prime}}{\omega}\left(f \frac{\partial}{\partial u}-g \frac{\partial}{\partial v}\right)\right\}
\end{align*}
$$

By a straightforward computation, the Laplacian $\Delta G$ of the Gauss map $G$ with the help of (3.3) turns out to be

$$
\begin{align*}
\Delta G_{1}= & \frac{1}{\omega^{\frac{7}{2}}}\left\{\left(f^{2}+\varepsilon \omega\right)^{2}\left(4 \varepsilon f f^{\prime 2}+\omega f^{\prime \prime}\right)\right. \\
& +\varepsilon f\left(g^{2}-\varepsilon \omega\right)\left[4 g^{2} g^{\prime 2}-\varepsilon \omega\left(g^{\prime 2}+g g^{\prime \prime}\right)\right] \\
& \left.+4 \varepsilon f f^{\prime} g^{2} g^{\prime}\left(\varepsilon \omega+2 f^{2}\right)-f f^{\prime} g^{\prime} \omega\left(f^{2}+\varepsilon \omega\right)\right\}, \\
\Delta G_{2}= & \frac{1}{\omega^{\frac{7}{2}}}\left\{\varepsilon g\left(f^{2}+\varepsilon \omega\right)\left[-4 f^{2} f^{\prime 2}-\varepsilon \omega\left(f^{\prime 2}+f f^{\prime \prime}\right)\right]\right.  \tag{3.5}\\
& +\left(g^{2}-\varepsilon \omega\right)^{2}\left(-4 \varepsilon g g^{\prime 2}+\omega g^{\prime \prime}\right) \\
& \left.+4 \varepsilon f^{2} f^{\prime} g g^{\prime}\left(\varepsilon \omega-2 g^{2}\right)-f^{\prime} g g^{\prime} \omega\left(g^{2}-\varepsilon \omega\right)\right\}, \\
\Delta G_{3}= & \frac{1}{\omega^{\frac{7}{2}}}\left\{\varepsilon\left(f^{2}+\varepsilon \omega\right)\left[3 f^{2}+f^{\prime} g^{2} g^{\prime}+\varepsilon \omega f f^{\prime \prime}+\left(\varepsilon \omega+f^{2}\right) f^{\prime 2}\right]\right. \\
& \left.+\varepsilon\left(g^{2}-\varepsilon \omega\right)\left[3 g^{2}+f^{2} f^{\prime} g^{\prime}-\varepsilon \omega g g^{\prime \prime}+\left(g^{2}-\varepsilon \omega\right) g^{\prime 2}\right]\right\} .
\end{align*}
$$

Before going into the study of translation surfaces with condition (1.1), let us examine some examples of surfaces in $\mathbb{E}_{1}^{3}$ satisfying that condition. They will be parts of our classifications of translation surfaces.

Example 3.1. Euclidean plane $\mathbb{R}^{2}$, or Minkowski plane $\mathbb{R}_{1}^{2}$.
In these cases the Gauss map is a constant normal time-like or space-like vector $G$, so $\Delta G=0$. Thus, the Euclidean plane $\mathbb{R}^{2}$ or the Minkowski plane $\mathbb{R}_{1}^{2}$ satisfies (1.1) with $A=0$.

Example 3.2. Lorentz circular cylinder $\mathbb{S}_{1}^{1} \times \mathbb{R}$.
Let $-x_{1}^{2}+x_{3}^{2}=r^{2}, r>0$, be the Lorentz circular cylinder. We consider this surface parametrized by $x(u, v)=\left(x_{1}=u, x_{2}=v, x_{3}= \pm \sqrt{r^{2}+u^{2}}\right)$. The Gauss map $G$ is given by $G=\left( \pm u / r, 0, \sqrt{r^{2}+u^{2}} / r\right)$ and the Laplacian is $\Delta G=$ $\left(1 / r^{2}\right) G$. Thus, the Lorentz circular cylinder $\mathbb{S}_{1}^{1} \times \mathbb{R}$ satisfies (1.1) with

$$
A=\left[\begin{array}{ccc}
\frac{1}{r^{2}} & a_{12} & 0 \\
0 & a_{22} & 0 \\
0 & a_{32} & \frac{1}{r^{2}}
\end{array}\right]
$$

Example 3.3. Hyperbolic cylinder $\mathbb{H}^{1} \times \mathbb{R}$.
Let $-x_{1}^{2}+x_{3}^{2}=-r^{2}, r>0$, be the hyperbolic cylinder and consider this surface parametrized by $x(u, v)=\left(x_{1}=u, x_{2}=v, x_{3}= \pm \sqrt{u^{2}-r^{2}}\right)$. The Gauss map $G$ is $G=\left( \pm u / r, 0, \sqrt{u^{2}-r^{2}} / r\right)$, and the Laplacian is $\Delta G=-\left(1 / r^{2}\right) G$. Thus, the hyperbolic cylinder $\mathbb{H}^{1} \times \mathbb{R}$ satisfies (1.1) with

$$
A=\left[\begin{array}{ccc}
-\frac{1}{r^{2}} & a_{12} & 0 \\
0 & a_{22} & 0 \\
0 & a_{32} & -\frac{1}{r^{2}}
\end{array}\right] .
$$

Example 3.4. Circular cylinder of index $1, \mathbb{R}_{1}^{1} \times \mathbb{S}^{1}$.
Let $x_{2}^{2}+x_{3}^{2}=r^{2}, r>0$, be the circular cylinder of index 1 and consider this surface parametrized by $x(u, v)=\left(x_{1}=u, x_{2}=v, x_{3}= \pm \sqrt{r^{2}-v^{2}}\right)$. The Gauss map $G$ is $G=\left(0, \pm v / r, \sqrt{r^{2}-v^{2}} / r\right)$, and the Laplacian $\Delta G$ of the Gauss map $G$ can be expressed as $\Delta G=\left(1 / r^{2}\right) G$. Thus, the circular cylinder of index $1, \mathbb{R}_{1}^{1} \times \mathbb{S}$, satisfies (1.1) with

$$
A=\left[\begin{array}{ccc}
a_{11} & 0 & 0 \\
a_{21} & \frac{1}{r^{2}} & 0 \\
a_{31} & 0 & \frac{1}{r^{2}}
\end{array}\right]
$$

## 4. Proof of the Theorem

We now assume that the surface $M^{2}$ satisfies condition (1.1). Then, combining (3.3) and (3.5), we have

$$
\begin{equation*}
+4 \varepsilon f^{2} f^{\prime} g g^{\prime}\left(\varepsilon \omega-2 g^{2}\right)-\omega\left(g^{2}-\varepsilon \omega\right) f^{\prime} g g^{\prime}=\omega^{3}\left(a_{21} f-a_{22} g+a_{23}\right) \tag{4.2}
\end{equation*}
$$

$$
\begin{align*}
& \left(f^{2}+\varepsilon \omega\right)^{2}\left(4 \varepsilon f f^{\prime 2}+\omega f^{\prime \prime}\right)+\varepsilon f\left(g^{2}-\varepsilon \omega\right)\left\{4 g^{2} g^{\prime 2}-\varepsilon \omega\left(g^{\prime 2}+g g^{\prime \prime}\right)\right\} \\
& +4 \varepsilon f f^{\prime} g^{2} g^{\prime}\left(\varepsilon \omega+2 f^{2}\right)-\omega\left(f^{2}+\varepsilon \omega\right) f f^{\prime} g^{\prime}=\omega^{3}\left(a_{11} f-a_{12} g+a_{13}\right), \tag{4.1}
\end{align*}
$$

$$
\left(g^{2}-\varepsilon \omega\right)^{2}\left(-4 \varepsilon g g^{\prime 2}+\omega g^{\prime \prime}\right)-\varepsilon g\left(f^{2}+\varepsilon \omega\right)\left\{4 f^{2} f^{\prime 2}+\varepsilon \omega\left(f^{\prime 2}+f f^{\prime \prime}\right)\right\}
$$

$$
\begin{align*}
& \varepsilon\left(f^{2}+\varepsilon \omega\right)\left\{3 f^{2}+f^{\prime} g^{2} g^{\prime}+\varepsilon \omega f f^{\prime \prime}+\left(\varepsilon \omega+f^{2}\right) f^{\prime 2}\right\}  \tag{4.3}\\
& +\varepsilon\left(g^{2}-\varepsilon \omega\right)\left\{3 g^{2}+f^{2} f^{\prime} g^{\prime}-\varepsilon \omega g g^{\prime \prime}+\left(g^{2}-\varepsilon \omega\right) g^{2}\right\} \\
& =\omega^{3}\left(a_{31} f-a_{32} g+a_{33}\right)
\end{align*}
$$

Furthermore, (4.3) can be rewritten in the form

$$
\begin{align*}
& \omega\left(f^{2}+\varepsilon \omega\right) f f^{\prime \prime}=-\varepsilon\left(f^{2}+\varepsilon \omega\right)\left\{3 f^{2}+f^{\prime} g^{2} g^{\prime}+\left(\varepsilon \omega+f^{2}\right) f^{\prime 2}\right\} \\
& -\varepsilon\left(g^{2}-\varepsilon \omega\right)\left\{3 g^{2}+f^{2} f^{\prime} g^{\prime}-\varepsilon \omega g g^{\prime \prime}+\left(g^{2}-\varepsilon \omega\right) g^{\prime 2}\right\}  \tag{4.4}\\
& +\omega^{3}\left(a_{31} f-a_{32} g+a_{33}\right),
\end{align*}
$$

which implies from (4.1) and (4.4)

$$
\begin{equation*}
A_{1} f^{\prime 2}+B_{1} f^{\prime}=\Gamma_{1}, \tag{4.5}
\end{equation*}
$$

where we put

$$
\begin{align*}
A_{1}= & \left(f^{2}+\varepsilon \omega\right)^{2}\left(3 \varepsilon f^{2}-\omega\right),  \tag{4.6}\\
B_{1}= & g^{2} g^{\prime}\left\{4 \varepsilon f^{4}-\varepsilon\left(g^{2}+1\right)\left(-3 f^{2}+g^{2}+1\right)\right\}, \\
\Gamma_{1}= & \omega^{3} f\left(a_{11} f-a_{12} g+a_{13}\right)-\omega^{3}\left(f^{2}+\varepsilon \omega\right)\left(a_{31} f-a_{32} g+a_{33}\right) \\
& +3 \varepsilon f^{2}\left(f^{2}+\varepsilon \omega\right)^{2}+\varepsilon\left(g^{2}-\varepsilon \omega\right)\left\{\left(f^{2}+\varepsilon \omega\right)\left[3 g^{2}-\varepsilon \omega g g^{\prime \prime}+\left(g^{2}-\varepsilon \omega\right) g^{\prime 2}\right]\right. \\
& \left.+\varepsilon \omega f^{2}\left(g^{\prime 2}+g g^{\prime \prime}\right)-4 f^{2} g^{2} g^{\prime 2}\right\} .
\end{align*}
$$

Also, it follows using (4.2) and (4.4) that

$$
\begin{equation*}
A_{2} f^{\prime 2}+B_{2} f^{\prime}=\Gamma_{2}, \tag{4.7}
\end{equation*}
$$

where we set

$$
\begin{align*}
A_{2}= & -3 \varepsilon f^{2} g\left(f^{2}+\varepsilon \omega\right), \\
B_{2}= & g g^{\prime}\left(\omega\left(3 f^{2}+\omega\right)-6 \varepsilon f^{2} g^{2}\right), \\
\Gamma_{2}= & \omega^{3}\left(a_{21} f-a_{22} g+a_{23}\right)+\omega^{3} g\left(a_{31} f-a_{32} g+a_{33}\right)  \tag{4.8}\\
& +\left(g^{2}-\varepsilon \omega\right)^{2}\left(3 \varepsilon g g^{\prime 2}-\omega g^{\prime \prime}\right)-3 \varepsilon f^{2} g\left(f^{2}+\varepsilon \omega\right) \\
& -\varepsilon g^{2}\left(g^{2}-\varepsilon \omega\right)\left(3 g-\varepsilon \omega g^{\prime \prime}\right) .
\end{align*}
$$

In case

$$
\begin{equation*}
A_{1} B_{2}-A_{2} B_{1}=0, \tag{4.9}
\end{equation*}
$$

from (4.6) and (4.8), we see that

$$
\begin{align*}
& \left(f^{2}+\varepsilon \omega\right)^{2}\left(3 \varepsilon f^{2}-\omega\right)\left(3 f^{2} \omega+\omega^{2}-6 \varepsilon f^{2} g^{2}\right) g g^{\prime} \\
& +3 \varepsilon f^{2} g^{3} g^{\prime}\left(f^{2}+\varepsilon \omega\right)\left\{4 \varepsilon f^{4}-\varepsilon\left(g^{2}+1\right)\left(-3 f^{2}+g^{2}+1\right)\right\}=0 . \tag{4.10}
\end{align*}
$$

Thus, the function $f(u)$ satisfies a nontrivial polynomial whose coefficients depend exclusively on the function $g$ and its derivative $g^{\prime}$. Consequently, $f$ must be constant. We will consider this situation further in the last step of the proof.
In case

$$
\begin{equation*}
A_{1} B_{2}-A_{2} B_{1} \neq 0, \tag{4.11}
\end{equation*}
$$

from (4.5) and (4.9) we have

$$
\begin{equation*}
\left(A_{1} \Gamma_{2}-A_{2} \Gamma_{1}\right)^{2}=\left(A_{1} B_{2}-A_{2} B_{1}\right)\left(B_{2} \Gamma_{1}-B_{1} \Gamma_{2}\right) \tag{4.12}
\end{equation*}
$$

Substituting (4.6) and (4.8) in (4.12), again we obtain a nontrivial polynomial in $f$ whose coefficients now depend exclusively on the functions $g, g^{\prime}$ and $g^{\prime \prime}$. Hence, $f$ must be constant. Now, we consider the situation that $f$ is constant. If $f$ is identically zero, then $\tilde{f}$ is constant, say, $c$. Thus, $M^{2}$ is a ruled surface in $\mathbb{E}_{1}^{3}$ and the position vector $x$ can be written in the following form:

$$
\begin{equation*}
x(u, v)=(u, v, c+\tilde{g}(v))=\alpha(v)+u \beta \tag{4.13}
\end{equation*}
$$

where $\alpha(v)=(0, v, c+\tilde{g}(v))$ is a space-like curve and $\beta=(1,0,0)$ is a time-like unit constant vector along $\alpha$ orthogonal to it. Consequently, the surface $M^{2}$ is locally the Minkowski plane $\mathbb{R}_{1}^{2}$ (Example 3.1) or the circular cylinder of index 1, $\mathbb{R}_{1}^{1} \times \mathbb{S}^{1}$ (Example 3.4) according to Proposition 3.2 of [9]. Lastly, we assume that $f$ is a nonzero constant. From (4.1) and (4.2) we obtain the following equations:

$$
\begin{align*}
\varepsilon f\left(g^{2}-\varepsilon \omega\right)\left\{4 g^{2} g^{\prime 2}-\varepsilon \omega\left(g^{\prime 2}+g g^{\prime \prime}\right)\right\} & =\omega^{3} f\left(a_{11} f-a_{12} g+a_{13}\right),  \tag{4.14}\\
\left(g^{2}-\varepsilon \omega\right)^{2}\left(\omega g^{\prime \prime}-4 \varepsilon g g^{\prime 2}\right) & =\omega^{3}\left(a_{21} f-a_{22} g+a_{23}\right) .
\end{align*}
$$

Considering (4.14) as a system of equations in $g^{2}$ and $g^{\prime \prime}$, we observe that since $f \neq 0$, its unique solution is

$$
\begin{align*}
& g^{\prime 2}=-\frac{\omega^{2}}{f\left(g^{2}-\varepsilon \omega\right)^{2}}\left\{\left(g^{2}-\varepsilon \omega\right)\left(a_{11} f-a_{12} g+a_{13}\right)\right. \\
&\left.+f g\left(a_{21} f-a_{22} g+a_{23}\right)\right\},  \tag{4.15}\\
& g^{\prime \prime}=-\frac{\varepsilon \omega}{f\left(g^{2}-\varepsilon \omega\right)^{2}}\left\{f\left(4 g^{2}-\varepsilon \omega\right)\left(a_{21} f-a_{22} g+a_{23}\right)\right. \\
&\left.+4 g\left(g^{2}-\varepsilon \omega\right)\left(a_{11} f-a_{12} g+a_{13}\right)\right\}
\end{align*}
$$

Substituting (4.15) in (4.3) yields a nontrivial polynomial in $g$ with constant coefficients. Hence, $g$ must be constant, which gives $\Delta G=0$. Consequently, $M^{2}$ is a nondegenerate plane, i.e., a Euclidean plane $\mathbb{R}^{2}$ or a Minkowski plane $\mathbb{R}_{1}^{2}$ (Example 3.1).

Now, we come back to relations (4.1), (4.2) and (4.3) and work as above to find from these the function $g$. Thus, we can rewrite (4.3) in the form

$$
\begin{align*}
& \omega\left(g^{2}-\varepsilon \omega\right) g g^{\prime \prime}=\varepsilon\left(f^{2}+\varepsilon \omega\right)\left\{3 f^{2}+f^{\prime} g^{2} g^{\prime}+\varepsilon \omega f f^{\prime \prime}+\left(\varepsilon \omega+f^{2}\right) f^{\prime 2}\right\}  \tag{4.16}\\
& +\varepsilon\left(g^{2}-\varepsilon \omega\right)\left\{3 g^{2}+f^{2} f^{\prime} g^{\prime}+\left(g^{2}-\varepsilon \omega\right) g^{\prime 2}\right\}-\omega^{3}\left(a_{31} f-a_{32} g+a_{33}\right)
\end{align*}
$$

Then, we combine (4.16) and (4.1) to obtain

$$
\begin{equation*}
A_{3} g^{2}+B_{3} g^{\prime}=\Gamma_{3} \tag{4.17}
\end{equation*}
$$

where

$$
\begin{align*}
& A_{3}=3 \varepsilon f g^{2}\left(g^{2}-\varepsilon \omega\right) \\
& B_{3}=f f^{\prime}\left\{6 \varepsilon f^{2} g^{2}+\omega\left(3 g^{2}-\varepsilon \omega\right)\right\} \tag{4.18}
\end{align*}
$$

and

$$
\begin{align*}
\Gamma_{3}= & \omega^{3}\left(a_{11} f-a_{12} g+a_{13}\right)-\omega^{3} f\left(a_{31} f-a_{32} g+a_{33}\right) \\
& -\left(f^{2}+\varepsilon \omega\right)^{2}\left(3 \varepsilon f f^{\prime 2}+\omega f^{\prime \prime}\right)+3 \varepsilon f g^{2}\left(g^{2}-\varepsilon \omega\right)  \tag{4.19}\\
& +\varepsilon f^{2}\left(f^{2}+\varepsilon \omega\right)\left(3 f+\varepsilon \omega f^{\prime \prime}\right) .
\end{align*}
$$

Also, it follows from (4.16) and (4.2) that

$$
\begin{equation*}
A_{4} g^{2}+B_{4} g^{\prime}=\Gamma_{4} \tag{4.20}
\end{equation*}
$$

where we put

$$
\begin{align*}
& A_{4}=\left(g^{2}-\varepsilon \omega\right)^{2}\left(-3 \varepsilon g^{2}-\omega\right) \\
& B_{4}=f^{2} f^{\prime}\left(\varepsilon \omega-2 g^{2}\right)\left(3 \varepsilon g^{2}+\omega\right), \tag{4.21}
\end{align*}
$$

and

$$
\begin{align*}
\Gamma_{4}= & \omega^{3}\left(a_{21} f-a_{22} g+a_{23}\right)+\omega^{3}\left(g^{2}-\varepsilon \omega\right)\left(a_{31} f-a_{32} g+a_{33}\right)  \tag{4.22}\\
& -3 \varepsilon g^{2}\left(g^{2}-\varepsilon \omega\right)^{2}+\varepsilon\left(f^{2}+\varepsilon \omega\right)\left\{\left(g^{2}-\varepsilon \omega\right)\left[-3 f^{2}-\varepsilon \omega f f^{\prime \prime}-\left(f^{2}+\varepsilon \omega\right) f^{\prime 2}\right]\right. \\
& \left.+\varepsilon \omega g^{2}\left(f^{\prime 2}+f f^{\prime \prime}\right)+4 f^{2} g^{2} f^{\prime 2}\right\} .
\end{align*}
$$

Now, by using (4.17) and (4.20), when

$$
A_{3} B_{4}-A_{4} B_{3}=0,
$$

from (4.18) and (4.21) we have that

$$
\begin{aligned}
& \left(g^{2}-\varepsilon \omega\right)^{2}\left(3 \varepsilon g^{2}+\varepsilon \omega\right)\left(3 g^{2} \omega-\varepsilon \omega^{2}+6 \varepsilon f^{2} g^{2}\right) \\
& +3 \varepsilon f^{3} g^{2} f^{\prime}\left(g^{2}-\varepsilon \omega\right)\left[-4 \varepsilon g^{4}+\varepsilon\left(f^{2}-1\right)\left(-3 g^{2}+f^{2}-1\right)\right]=0
\end{aligned}
$$

Thus, the function $g$ satisfies a nontrivial polynomial whose coefficients depend exclusively on the function $f$ and its derivative $f^{\prime}$. Consequently, $g$ must be constant. When

$$
A_{3} B_{4}-A_{4} B_{3} \neq 0,
$$

from (4.17) and (4.20) we have that

$$
\begin{equation*}
\left(A_{3} \Gamma_{4}-A_{4} \Gamma_{3}\right)^{2}=\left(A_{3} B_{4}-A_{4} B_{3}\right)\left(B_{4} \Gamma_{3}-B_{3} \Gamma_{4}\right) . \tag{4.23}
\end{equation*}
$$

Inserting (4.18), (4.19), (4.21) and (4.22) in (4.23), again we obtain a nontrivial polynomial in $g$ whose coefficients now depend exclusively on the functions $f, f^{\prime}$ and $f^{\prime \prime}$. Hence, $g$ must be constant.

If $g$ is identically zero, then $\tilde{g}$ is constant, say, $c$. Thus, in this case $M^{2}$ is a ruled surface in $\mathbb{E}_{1}^{3}$ and the position vector field $x$ takes the form

$$
x(u, v)=(u, v, \tilde{f}(u)+c)=\alpha(u)+v \beta,
$$

where $\alpha(u)=(u, 0, \tilde{f}(u)+c)$ is a space-like or time-like curve and $\beta=(0,1,0)$ is a space-like unit constant vector along $\alpha$ orthogonal to it. Consequently, the surface $M^{2}$ is locally the Euclidean plane $\mathbb{R}^{2}$, the Minkowski plane $\mathbb{R}_{1}^{2}$ (Example 3.1), the hyperbolic cylinder $\mathbb{H}^{1} \times \mathbb{R}$ (Example 3.3) or the Lorentz circular cylinder $\mathbb{S}_{1}^{1} \times \mathbb{R}$ (Example 3.2) according to Proposition 3.1 of [9].

Finally, if $g$ is a nonzero constant, we obtain again, as above, that $f$ is constant, and thus $M^{2}$ is a nondegenerate plane, i.e., a Euclidean plane $\mathbb{R}^{2}$ or a Minkowski plane $\mathbb{R}_{1}^{2}$ (Example 3.1). This completes the proof.

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