# CLASSIFICATION OF RULED SURFACES WITH POINTWISE 1-TYPE GAUSS MAP IN MINKOWSKI 3-SPACE 

Miekyung Choi ${ }^{1}$, Young Ho Kim ${ }^{2}$ and Dae Won Yoon ${ }^{3}$


#### Abstract

We study the ruled surfaces in Minkowski 3 -space with pointwise 1 -type Gauss map. As a result, we introduce some new examples of the ruled surfaces with pointwise 1-type Gauss map in Minkowski 3-space.


## 1. Introduction

The notion of finite type immersions has played an important role in classifying and characterizing the submanifolds in Euclidean space or pseudo-Euclidean space since it was introduced by B.-Y. Chen in the late 1970s ([3, 4]). Also, we can apply it to the smooth maps in Euclidean space or pseudo-Euclidean space naturally ( $[1,2,6]$ ). A smooth map $\phi$ of a submanifold $M$ of Euclidean space or pseudoEuclidean space is said to be of finite type if $\phi$ can be expressed as a finite sum of eigenfunctions of the Laplacian $\Delta$ of $M$, that is, $\phi=\phi_{0}+\sum_{i=1}^{k} \phi_{i}$, where $\phi_{0}$ is a constant function, $\phi_{i}(i=1, \cdots, k)$ non-constant functions satisfying $\Delta \phi_{i}=\lambda_{i} \phi_{i}$, $\lambda_{i} \in \mathbb{R}$.

In this regards, it is worth investigating the classification of the submanifolds in Euclidean space or pseudo-Euclidean space in terms of finite type Gauss map. In general, the Laplacian of 1-type Gauss map of a submanifold in Euclidean or pseudo-Euclidean space satisfies $\Delta G=\lambda(G+\mathbb{C})$ for some constant $\lambda$ and a constant vector $\mathbb{C}$.

On the other hand, the Gauss map of some minimal surfaces such as a helicoid of the first kind, a helicoid of the second kind, a helicoid of the third kind and the conjugate of Enneper's surface of the second kind in Minkowski 3-space satisfies $\Delta G=f G$ for some smooth function $f$. It looks like an eigenvalue problem but the

[^0]function $f$ turns out to be non-constant for such surfaces. For this reason, the notion of pointwise 1-type Gauss map in Euclidean space or pseudo-Euclidean space was initiated: A submanifold $M$ in Euclidean space $\mathbb{E}^{m}$ or pseudo-Euclidean space $\mathbb{E}_{s}^{m}$ of index $s$ is said to have pointwise 1-type Gauss map if
\[

$$
\begin{equation*}
\Delta G=f(G+\mathbb{C}) \tag{1.1}
\end{equation*}
$$

\]

for a nonzero smooth function $f$ and some constant vector $\mathbb{C}$. In particular, if $\mathbb{C}$ is zero, it is said to be of the first kind. Otherwise, it is said to be of the second kind ( $[5,7,8,9,10,11,12,14,15]$ ).

Recently, the present authors have introduced some new examples of the ruled surfaces with pointwise 1-type Gauss map of the second kind in $\mathbb{E}^{3}$ called a cylinder of an infinite type and a rotational ruled surface ([10]). Two of the present authors gave the classification of the ruled surfaces with pointwise 1-type Gauss map of the first kind in Minkowski 3 -space $\mathbb{E}_{1}^{3}$ ([14]).

In the present paper, we mainly focus on a ruled surface in Minkowski 3-space $\mathbb{E}_{1}^{3}$ with pointwise 1-type Gauss map of the second kind. In fact, the class of solution spaces of equation (1.1) could be very big because it could have infinitely many solutions associated with a function $f$ and a constant vector $\mathbb{C}$.

As a consequence, by combining the results in [14], we give a complete classification of ruled surfaces with pointwise 1-type Gauss map in Minkowski 3-space $\mathbb{E}_{1}^{3}$ :

Theorem A (Classification). Let $M$ be a ruled surface in Minkowski 3-space $\mathbb{E}_{1}^{3}$ with pointwise 1-type Gauss map. Then, $M$ is an open part of a Euclidean plane, a Minkowski plane, a hyperbolic cylinder, a Lorentz circular cylinder, a circular cylinder of index 1, a cylinder of an infinite type, a helicoid of the first kind, a helicoid of the second kind, a helicoid of the third kind, the conjugate of Enneper's surface of the second kind, a rotational ruled surface of type I or type II, a transcendental ruled surface, or a $B$-scroll.

## 2. Preliminaries

Let $\mathbb{E}_{1}^{3}$ be Minkowski 3-space with the Lorentz metric $d s^{2}=-d x_{1}^{2}+d x_{2}^{2}+d x_{3}^{2}$, where $\left(x_{1}, x_{2}, x_{3}\right)$ denotes the standard coordinate system in $\mathbb{E}_{1}^{3}$. Let $M$ be a nondegenerate connected surface in $\mathbb{E}_{1}^{3}$. The map $G: M \rightarrow Q^{2}(\epsilon) \subset \mathbb{E}_{1}^{3}$ which maps each point of $M$ to the unit normal vector to $M$ at the point is called the Gauss map of $M$, where $\epsilon(= \pm 1)$ denotes the sign of the vector field $G$ and $Q^{2}(\epsilon)$ is a 2-dimensional space form with constant sectional curvature $\epsilon$.

Now, we define a ruled surface $M$ in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Let $I$ and $J$ be some open intervals in the real line $\mathbb{R}$. Let $\alpha=\alpha(s)$ be a curve in $\mathbb{E}_{1}^{3}$ defined on $I$
and $\beta=\beta(s)$ a transversal vector field with $\alpha^{\prime}(s)$ along $\alpha$. Then, a parametrization of a ruled surface $M$ is given by

$$
x(s, t)=\alpha(s)+t \beta(s), \quad s \in I, \quad t \in J
$$

The curve $\alpha=\alpha(s)$ is called a base curve and $\beta=\beta(s)$ a director vector field. In particular, if $\beta$ is constant, $M$ is said to be cylindrical. Otherwise, it is said to be non-cylindrical.

First, we consider a base curve $\alpha$ is space-like or time-like. Then, the base curve $\alpha$ can be chosen to be orthogonal to the director vector field $\beta$ which can be normalized as $\langle\beta, \beta\rangle= \pm 1$. In this case, we have five different types according to the character of $\alpha$ and $\beta$ as follows: According as the base curve $\alpha$ is space-like or time-like, the ruled surface $M$ is said to be of type $M_{+}$or $M_{-}$, respectively. Also, the ruled surface of type $M_{+}$can be divided into three types. If $\beta$ is space-like, it is said to be of type $M_{+}^{1}$ or $M_{+}^{2}$ if $\beta^{\prime}$ is non-null or null, respectively. When $\beta$ is time-like, $\beta^{\prime}$ is space-like because of the causal vector of $\beta$, which is said to be of type $M_{+}^{3}$. On the other hand, when $\alpha$ is time-like, $\beta$ is always space-like. Accordingly, it is also said to be of type $M_{-}^{1}$ or $M_{-}^{2}$ if $\beta^{\prime}$ is non-null or null, respectively. The ruled surface of type $M_{+}^{1}$ or $M_{+}^{2}$ (resp. $M_{+}^{3}, M_{-}^{1}$ or $M_{-}^{2}$ ) is clearly space-like (resp. time-like).

A curve in $\mathbb{E}_{1}^{3}$ is said to be null or light-like if its tangent vector field is null along it. If the base curve $\alpha$ is null and the director vector field $\beta$ along $\alpha$ is null, then the ruled surface $M$ is called a null scroll. It is evidently a time-like surface.

Other cases such as $\alpha$ is non-null and $\beta$ is null, or $\alpha$ is null and $\beta$ is non-null are reduced to one of the types $M_{ \pm}^{1}, M_{ \pm}^{2}$ and $M_{+}^{3}$, or a null scroll by an appropriate change of the base curve ([13]).

## 3. Cylindrical Ruled Surfaces

In this section, we examine the cylindrical ruled surfaces with pointwise 1-type Gauss map of the second kind in Minkowski 3-space.

Let $M$ be a cylindrical ruled surface in Minkowski 3-space $\mathbb{E}_{1}^{3}$ of type $M_{+}^{1}, M_{-}^{1}$ or $M_{+}^{3}$. For a unit constant vector field $\beta, M$ is parameterized by

$$
x(s, t)=\alpha(s)+t \beta
$$

such that $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=\epsilon_{1}(= \pm 1),\left\langle\alpha^{\prime}, \beta\right\rangle=0$ and $\langle\beta, \beta\rangle=\epsilon_{2}(= \pm 1)$.
We consider two cases separately.
Case 1. Let $M$ be a cylindrical ruled surface of type $M_{+}^{1}$ or $M_{-}^{1}$, i.e., $\epsilon_{2}=1$. Without loss of generality, we may assume that $\alpha(s)=\left(\alpha_{1}(s), \alpha_{2}(s), 0\right)$ is a plane curve parameterized by the arc-length $s$ and the constant vector field $\beta$ is chosen
as $\beta=(0,0,1)$. Then, the Gauss map $G$ and its Laplacian $\Delta G$ of $M$ are given by $G=\alpha^{\prime} \times \beta=\left(-\alpha_{2}{ }^{\prime},-\alpha_{1}{ }^{\prime}, 0\right)$ and $\Delta G=\left(\epsilon_{1} \alpha_{2}^{\prime \prime \prime}, \epsilon_{1} \alpha_{1}^{\prime \prime \prime}, 0\right)$, respectively, where the prime denotes the differentiation with respect to $s$.

Suppose that $M$ has pointwise 1-type Gauss map of the second kind, that is, the Gauss map $G$ satisfies equation (1.1). Then, the third component of the constant vector $\mathbb{C}$ is zero and we have a system of differential equations:

$$
\begin{align*}
& \epsilon_{1} \alpha_{2}^{\prime \prime \prime}=f\left(-\alpha_{2}^{\prime}+c_{1}\right) \\
& \epsilon_{1} \alpha_{1}^{\prime \prime \prime}=f\left(-\alpha_{1}^{\prime}+c_{2}\right) \tag{3.1}
\end{align*}
$$

where $\mathbb{C}=\left(c_{1}, c_{2}, 0\right)$.
First, we consider the case that $M$ is of type $M_{+}^{1}$. Since $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=-\alpha_{1}^{\prime 2}+\alpha_{2}^{\prime 2}=$ 1, we may put

$$
\alpha_{1}^{\prime}(s)=\sinh \theta, \quad \alpha_{2}^{\prime}(s)=\cosh \theta
$$

for a function $\theta=\theta(s)$. Therefore, equation (3.1) can be written as

$$
\begin{aligned}
& \left(\theta^{\prime}\right)^{2} \cosh \theta+\theta^{\prime \prime} \sinh \theta=f\left(-\cosh \theta+c_{1}\right) \\
& \left(\theta^{\prime}\right)^{2} \sinh \theta+\theta^{\prime \prime} \cosh \theta=f\left(-\sinh \theta+c_{2}\right)
\end{aligned}
$$

It follows

$$
\begin{equation*}
\left(\theta^{\prime}\right)^{2}=f\left(-1+c_{1} \cosh \theta-c_{2} \sinh \theta\right) \tag{3.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\theta^{\prime \prime}=f\left(-c_{1} \sinh \theta+c_{2} \cosh \theta\right) \tag{3.3}
\end{equation*}
$$

Suppose $\theta^{\prime} \equiv 0$. Then, obviously $\Delta G=0$. Since $f$ is non-zero, (3.1) implies $\alpha_{1}^{\prime}=c_{2}, \alpha_{2}^{\prime}=c_{1}$ and thus $G=-\mathbb{C}$. Therefore, $M$ is an open part of a Euclidean plane. If the interior $\operatorname{Int}(\mathrm{U})$ of a closed subset $\mathrm{U}=\left\{p \in M \mid \theta^{\prime}(p)=0\right\}$ is nonempty, U must be $M$ by the above argument and connectedness of $M$. Otherwise, if $\theta^{\prime}$ has zeros, the set of zeros of $\theta^{\prime}$ has measure zero.

Now we suppose $\theta^{\prime} \neq 0$. (3.1) shows that $f$ depends only on the parameter $s$, i.e., $f(s, t)=f(s)$. Differentiating (3.2) with respect to $s$ and using (3.2) and (3.3), we obtain

$$
\begin{equation*}
\theta^{\prime}=c \sqrt[3]{f} \tag{3.4}
\end{equation*}
$$

for some non-zero constant $c$. On the other hand, combining (3.2) and (3.3), we get the following differential equation

$$
\begin{equation*}
\left(\frac{\left(\theta^{\prime}\right)^{2}}{f}+1\right)^{2}-\left(\frac{\theta^{\prime \prime}}{f}\right)^{2}=c_{1}^{2}-c_{2}^{2} \tag{3.5}
\end{equation*}
$$

By using (3.4), equation (3.5) gives

$$
\begin{equation*}
\left(c^{2} f^{-\frac{1}{3}}+1\right)^{2}-\left(-\frac{c}{2}\left(f^{-\frac{2}{3}}\right)^{\prime}\right)^{2}=c_{1}^{2}-c_{2}^{2} \tag{3.6}
\end{equation*}
$$

If we put $f^{-\frac{1}{3}}=y$, then equation (3.6) becomes

$$
\left(c^{2} y+1\right)^{2}-\left(c y y^{\prime}\right)^{2}=c_{1}^{2}-c_{2}^{2} .
$$

If $\mathbb{C}$ is null, then the solution of the differential equation $\left(c^{2} y+1\right)^{2}-\left(c y y^{\prime}\right)^{2}=0$ is given by

$$
c^{2} y-\ln \left|c^{2} y+1\right|= \pm c^{3}(s+k)
$$

or, equivalently,

$$
\begin{equation*}
c^{2} f^{-\frac{1}{3}}-\ln \left|c^{2} f^{-\frac{1}{3}}+1\right|= \pm c^{3}(s+k) \tag{3.7}
\end{equation*}
$$

for some constant $k$.
If $\mathbb{C}$ is non-null, the solution of (3.6) is obtained as follows:

$$
\begin{align*}
& \sqrt{\left(c^{2} f^{-\frac{1}{3}}+1\right)^{2}+\left(-c_{1}^{2}+c_{2}^{2}\right)} \\
& \quad-\ln \left(c^{2} f^{-\frac{1}{3}}+1+\sqrt{\left(c^{2} f^{-\frac{1}{3}}+1\right)^{2}+\left(-c_{1}^{2}+c_{2}^{2}\right)}\right)  \tag{3.8}\\
& \quad+\ln \sqrt{\left|-c_{1}^{2}+c_{2}^{2}\right|}= \pm c^{3}(s+k)
\end{align*}
$$

for some constant $k$.
We now consider the case that $M$ is of type $M_{-}^{1}$. Since $-\alpha_{1}^{\prime 2}+\alpha_{2}^{\prime 2}=-1$, we may put

$$
\alpha_{1}^{\prime}(s)=\cosh \theta, \quad \alpha_{2}^{\prime}(s)=\sinh \theta
$$

for a function $\theta=\theta(s)$. As is the previous case, if $\theta^{\prime} \equiv 0, M$ is an open portion of a Minkowski plane. If $\theta^{\prime}$ is non-zero, we get $\theta^{\prime}=c \sqrt[3]{f}$ and

$$
\begin{equation*}
\left(c^{2} f^{-\frac{1}{3}}-1\right)^{2}-\left(-\frac{c}{2}\left(f^{-\frac{2}{3}}\right)^{\prime}\right)^{2}=-c_{1}^{2}+c_{2}^{2} \tag{3.9}
\end{equation*}
$$

for some non-zero constant $c$. In this case, if $\mathbb{C}$ is null or non-null, then its solution is obtained as, respectively,

$$
\begin{equation*}
c^{2} f^{-\frac{1}{3}}+\ln \left|c^{2} f^{-\frac{1}{3}}-1\right|= \pm c^{3}(s+k) \tag{3.10}
\end{equation*}
$$

or

$$
\begin{align*}
& \sqrt{\left(c^{2} f^{-\frac{1}{3}}-1\right)^{2}-\left(-c_{1}^{2}+c_{2}^{2}\right)} \\
& \quad+\ln \left(c^{2} f^{-\frac{1}{3}}-1+\sqrt{\left(c^{2} f^{-\frac{1}{3}}-1\right)^{2}+\left|-c_{1}^{2}+c_{2}^{2}\right|}\right)  \tag{3.11}\\
& \quad-\ln \sqrt{\left|-c_{1}^{2}+c_{2}^{2}\right|}= \pm c^{3}(s+k)
\end{align*}
$$

where $k$ is a constant.
Case 2. Let $M$ be a cylindrical ruled surface of type $M_{+}^{3}$. Then we may assume that $\beta=(1,0,0)$ and $\alpha(s)=\left(0, \alpha_{2}(s), \alpha_{3}(s)\right)$ without loss of generality. Hence, the Gauss map $G$ and its Laplacian $\Delta G$ of $M$ are obtained by $G=\left(0, \alpha_{3}{ }^{\prime},-\alpha_{2}{ }^{\prime}\right)$ and $\Delta G=\left(0,-\alpha_{3}^{\prime \prime \prime}, \alpha_{2}^{\prime \prime \prime}\right)$, respectively.

Suppose that the Gauss map $G$ of $M$ is of pointwise 1-type of the second kind. Then, we have

$$
\begin{aligned}
-\alpha_{3}^{\prime \prime \prime} & =f\left(\alpha_{3}^{\prime}+c_{2}\right), \\
\alpha_{2}^{\prime \prime \prime} & =f\left(-\alpha_{2}^{\prime}+c_{3}\right),
\end{aligned}
$$

where $\mathbb{C}=\left(0, c_{2}, c_{3}\right)$. Since $\alpha(s)$ is parameterized by the arc length, i.e., $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=$ ${\alpha_{2}^{\prime}}^{2}+{\alpha_{3}^{\prime}}^{2}=1$, we may put

$$
\alpha_{2}^{\prime}(s)=\cos \theta, \quad \alpha_{3}^{\prime}(s)=\sin \theta
$$

for a function $\theta=\theta(s)$. Like a similar discussion developed in Case $1, M$ is an open portion of a Minkowski plane when $\theta^{\prime} \equiv 0$. Otherwise, we can have $\theta^{\prime}=c \sqrt[3]{f}$ for some non-zero constant $c$. Moreover, the smooth function $f$ and the constant vector $\mathbb{C}$ satisfy

$$
\begin{equation*}
\sqrt{c_{2}^{2}+c_{3}^{2}-\left(c^{2} f^{-\frac{1}{3}}-1\right)^{2}}-\sin ^{-1}\left(\frac{c^{2} f^{-\frac{1}{3}}-1}{\sqrt{c_{2}^{2}+c_{3}^{2}}}\right)= \pm c^{3}(s+k) \tag{3.12}
\end{equation*}
$$

where $c$ is a non-zero constant and $k$ a constant.
Definition 3.1. ([10]). A cylindrical ruled surface over an infinite type base curve in Minkowski space is called a cylinder of an infinite type.

Thus, we have
Proposition 3.1. Let $M$ be a cylindrical ruled surface over a non-null base curve in Minkowski 3-space $\mathbb{E}_{1}^{3}$ with pointwise 1-type Gauss map of the second kind. If $M$ is not totally geodesic, then the non-zero smooth function $f$ satisfies
one of the equations (3.7), (3.8), (3.10), (3.11) or (3.12) depending upon the types of the base curve.

Combining the results above and [14], we have
Theorem 3.2. Let $M$ be a cylindrical ruled surface over a non-null base curve in Minkowski 3-space $\mathbb{E}_{1}^{3}$ with pointwise 1-type Gauss map of the first kind. Then, the curvature of the base curve is a non-zero constant. In other words, $M$ is an open part of a hyperbolic cylinder, a Lorentz circular cylinder or a circular cylinder of index 1 .

Theorem 3.3. (Classification). Let $M$ be a cylindrical ruled surface over a non-null base curve in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Then, M has pointwise 1-type Gauss map of the second kind if and only if $M$ is an open part of a Euclidean plane, a Minkowski plane or a cylinder of an infinite type satisfying (3.7), (3.8), (3.10), (3.11) or (3.12) up to rigid motion.

## 4. Non-cylindrical Ruled Surfaces

In this section, we classify the non-cylindrical ruled surfaces in Minkowski 3 -space $\mathbb{E}_{1}^{3}$ with pointwise 1 -type Gauss map.

Let $M$ be a non-cylindrical ruled surface of type $M_{+}^{1}, M_{+}^{3}$ or $M_{-}^{1}$ whose Gauss map is of pointwise 1-type of the second kind. Then, $M$ is parameterized by, up to rigid motion,

$$
x(s, t)=\alpha(s)+t \beta(s)
$$

such that $\left\langle\alpha^{\prime}, \beta\right\rangle=0,\langle\beta, \beta\rangle=\epsilon_{2}(= \pm 1)$ and $\left\langle\beta^{\prime}, \beta^{\prime}\right\rangle=\epsilon_{3}(= \pm 1)$. For later use, we define the smooth functions $q, u, Q$ and $R$ as follows:

$$
q=\left\|x_{s}\right\|^{2}=\epsilon_{4}\left\langle x_{s}, x_{s}\right\rangle, u=\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle, Q=\left\langle\alpha^{\prime}, \beta \times \beta^{\prime}\right\rangle, R=\left\langle\beta^{\prime \prime}, \beta \times \beta^{\prime}\right\rangle
$$

where $\epsilon_{4}(= \pm 1)$ is the sign of the coordinate vector field $x_{s}=\frac{\partial x}{\partial s}$. For an orthonormal frame $\left\{\beta, \beta^{\prime}, \beta \times \beta^{\prime}\right\}$ along the base curve $\alpha$, we have

$$
\begin{align*}
& \alpha^{\prime}=\epsilon_{3} u \beta^{\prime}-\epsilon_{2} \epsilon_{3} Q \beta \times \beta^{\prime}, \\
& \beta^{\prime \prime}=-\epsilon_{2} \epsilon_{3} \beta-\epsilon_{2} \epsilon_{3} R \beta \times \beta^{\prime},  \tag{4.1}\\
& \alpha^{\prime} \times \beta=-\epsilon_{3} u \beta \times \beta^{\prime}+\epsilon_{3} Q \beta^{\prime}, \\
& \beta \times \beta^{\prime \prime}=-\epsilon_{3} R \beta^{\prime},
\end{align*}
$$

which imply the smooth function $q$ given by

$$
q=\epsilon_{4}\left(\epsilon_{3} t^{2}+2 u t+\epsilon_{3} u^{2}-\epsilon_{2} \epsilon_{3} Q^{2}\right)
$$

We note that $t$ must be chosen so that $q$ takes positive values.

Furthermore, the Gauss map $G$ and the mean curvature $H$ of $M$ are straightforwardly obtained by, respectively,

$$
\begin{aligned}
& G=q^{-\frac{1}{2}}\left(\epsilon_{3} Q \beta^{\prime}-\left(\epsilon_{3} u+t\right) \beta \times \beta^{\prime}\right), \\
& H=\frac{1}{2} \epsilon_{2} q^{-3 / 2}\left(R t^{2}+\left(2 \epsilon_{3} u R+Q^{\prime}\right) t+u^{2} R+\epsilon_{3} u Q^{\prime}-\epsilon_{3} u^{\prime} Q-\epsilon_{2} Q^{2} R\right)
\end{aligned}
$$

On the other hand, the Laplacian $\Delta G$ of the Gauss map $G$ can be expressed as follows ([14]):

$$
\begin{equation*}
\Delta G=q^{-7 / 2} \epsilon_{4}\left(A_{1} \beta+A_{2} \beta^{\prime}+A_{3} \beta \times \beta^{\prime}\right), \tag{4.2}
\end{equation*}
$$

where we have put

$$
\begin{aligned}
& A_{1}=\epsilon_{2} R t^{5}+\left(2 \epsilon_{2} Q^{\prime}+5 \epsilon_{2} \epsilon_{3} u R\right) t^{4}+\left(-3 \epsilon_{2} \epsilon_{3} \epsilon_{4} u^{\prime} Q-2 Q^{2} R+8 \epsilon_{2} \epsilon_{3} u Q^{\prime}+10 \epsilon_{2} u^{2} R\right) t^{3} \\
& +\left(-4 Q^{2} Q^{\prime}+3 \epsilon_{4} Q^{2} Q^{\prime}-9 \epsilon_{2} \epsilon_{4} u u^{\prime} Q-6 \epsilon_{3} u Q^{2} R+12 \epsilon_{2} u^{2} Q^{\prime}+10 \epsilon_{2} \epsilon_{3} u^{3} R\right) t^{2} \\
& +\left(3 \epsilon_{3} \epsilon_{4} u^{\prime} Q^{3}+\epsilon_{2} Q^{4} R-8 \epsilon_{3} u Q^{2} Q^{\prime}+6 \epsilon_{3} \epsilon_{4} u Q^{2} Q^{\prime}-9 \epsilon_{2} \epsilon_{3} \epsilon_{4} u^{2} u^{\prime} Q-6 u^{2} Q^{2} R+8 \epsilon_{2} \epsilon_{3} u^{3} Q^{\prime}\right. \\
& \left.+5 \epsilon_{2} u^{4} R\right) t+2 \epsilon_{2} Q^{4} Q^{\prime}-3 \epsilon_{2} \epsilon_{4} Q^{4} Q^{\prime}+3 \epsilon_{4} u u^{\prime} Q^{3}+\epsilon_{2} \epsilon_{3} u Q^{4} R-4 u^{2} Q^{2} Q^{\prime}+3 \epsilon_{4} u^{2} Q^{2} Q^{\prime} \\
& -3 \epsilon_{2} \epsilon_{4} u^{3} u^{\prime} Q-2 \epsilon_{3} u^{3} Q^{2} R+2 \epsilon_{2} u^{4} Q^{\prime}+\epsilon_{2} \epsilon_{3} u^{5} R \text {, } \\
& A_{2}=-\epsilon_{3} R^{\prime} t^{5}+\left(u^{\prime} R-\epsilon_{2} \epsilon_{3} Q R^{2}-\epsilon_{3} Q^{\prime \prime}-5 u R^{\prime}\right) t^{4}+\left(u^{\prime \prime} Q+2 \epsilon_{2} \epsilon_{3} Q^{2} R^{\prime}+3 u^{\prime} Q^{\prime}-3 \epsilon_{2} \epsilon_{3} Q Q^{\prime} R\right. \\
& \left.+4 \epsilon_{3} u u^{\prime} R-4 \epsilon_{2} u Q R^{2}-4 u Q^{\prime \prime}-10 \epsilon_{3} u^{2} R^{\prime}\right) t^{3}+\left(-3 \epsilon_{3} u^{\prime 2} Q-2 Q^{3}+\epsilon_{2} u^{\prime} Q^{2} R+2 \epsilon_{3} Q^{3} R^{2}\right. \\
& -4 \epsilon_{2} \epsilon_{3} Q Q^{\prime 2}+\epsilon_{2} \epsilon_{3} Q^{2} Q^{\prime \prime}+3 \epsilon_{3} u u^{\prime \prime} Q+6 \epsilon_{2} u Q^{2} R^{\prime}+9 \epsilon_{3} u u^{\prime} Q^{\prime}-9 \epsilon_{2} u Q Q^{\prime} R+6 u^{2} u^{\prime} R \\
& \left.-6 \epsilon_{2} \epsilon_{3} u^{2} Q R^{2}-6 \epsilon_{3} u^{2} Q^{\prime \prime}-10 u^{3} R^{\prime}\right) t^{2}+\left(-\epsilon_{2} u^{\prime \prime} Q^{3}-\epsilon_{3} Q^{4} R^{\prime}+5 \epsilon_{2} u^{\prime} Q^{2} Q^{\prime}+3 \epsilon_{3} Q^{3} Q^{\prime} R\right. \\
& -6 u u^{\prime 2} Q-4 \epsilon_{3} u Q^{3}+2 \epsilon_{2} \epsilon_{3} u u^{\prime} Q^{2} R+4 u Q^{3} R^{2}-8 \epsilon_{2} u Q Q^{\prime 2}+2 \epsilon_{2} u Q^{2} Q^{\prime \prime}+3 u^{2} u^{\prime \prime} Q \\
& \left.+6 \epsilon_{2} \epsilon_{3} u^{2} Q^{2} R^{\prime}+9 u^{2} u^{\prime} Q^{\prime}-9 \epsilon_{2} \epsilon_{3} u^{2} Q Q^{\prime} R+4 \epsilon_{3} u^{3} u^{\prime} R-4 \epsilon_{2} u^{3} Q R^{2}-4 u^{3} Q^{\prime \prime}-5 \epsilon_{3} u^{4} R^{\prime}\right) t \\
& -\epsilon_{2} \epsilon_{3} u^{\prime 2} Q^{3}+2 Q^{5}-2 u^{\prime} Q^{4} R-\epsilon_{3} Q^{5} R^{2}-\epsilon_{2} \epsilon_{3} u u^{\prime \prime} Q^{3}-u Q^{4} R^{\prime}+5 \epsilon_{2} \epsilon_{3} u u^{\prime} Q^{2} Q^{\prime}+3 u Q^{3} Q^{\prime} R \\
& -3 \epsilon_{3} u^{2} u^{\prime 2} Q-2 u^{2} Q^{3}+\epsilon_{2} u^{2} u^{\prime} Q^{2} R+2 \epsilon_{3} u^{2} Q^{3} R^{2}-4 \epsilon_{2} \epsilon_{3} u^{2} Q Q^{\prime 2}+\epsilon_{2} \epsilon_{3} u^{2} Q^{2} Q^{\prime \prime}+\epsilon_{3} u^{3} u^{\prime \prime} Q \\
& +2 \epsilon_{2} u^{3} Q^{2} R^{\prime}+3 \epsilon_{3} u^{3} u^{\prime} Q^{\prime}-3 \epsilon_{2} u^{3} Q Q^{\prime} R+u^{4} u^{\prime} R-\epsilon_{2} \epsilon_{3} u^{4} Q R^{2}-\epsilon_{3} u^{4} Q^{\prime \prime}-u^{5} R^{\prime} \text {, } \\
& A_{3}=\epsilon_{2} R^{2} t^{5}+\left(\epsilon_{2} Q R^{\prime}+2 \epsilon_{2} Q^{\prime} R+5 \epsilon_{2} \epsilon_{3} u R^{2}\right) t^{4}+\left(2 \epsilon_{3} Q^{2}-3 \epsilon_{2} \epsilon_{3} u^{\prime} Q R-2 Q^{2} R^{2}+\epsilon_{2} Q^{\prime 2}\right. \\
& \left.+\epsilon_{2} Q Q^{\prime \prime}+4 \epsilon_{2} \epsilon_{3} u Q R^{\prime}+8 \epsilon_{2} \epsilon_{3} u Q^{\prime} R+10 \epsilon_{2} u^{2} R^{2}\right) t^{3}+\left(-\epsilon_{2} \epsilon_{3} u^{\prime \prime} Q^{2}-2 Q^{3} R^{\prime}-5 \epsilon_{2} \epsilon_{3} u^{\prime} Q Q^{\prime}\right. \\
& -Q^{2} Q^{\prime} R+6 u Q^{2}-9 \epsilon_{2} u u^{\prime} Q R-6 \epsilon_{3} u Q^{2} R^{2}+3 \epsilon_{2} \epsilon_{3} u Q^{\prime 2}+3 \epsilon_{2} \epsilon_{3} u Q Q^{\prime \prime}+6 \epsilon_{2} u^{2} Q R^{\prime} \\
& \left.+12 \epsilon_{2} u^{2} Q^{\prime} R+10 \epsilon_{2} \epsilon_{3} u^{3} R^{2}\right) t^{2}+\left(4 \epsilon_{2} u^{\prime 2} Q^{2}-2 \epsilon_{2} \epsilon_{3} Q^{4}+3 \epsilon_{3} u^{\prime} Q^{3} R+\epsilon_{2} Q^{4} R^{2}+3 Q^{2} Q^{\prime 2}\right. \\
& -Q^{3} Q^{\prime \prime}-2 \epsilon_{2} u u^{\prime \prime} Q^{2}-4 \epsilon_{3} u Q^{3} R^{\prime}-10 \epsilon_{2} u u^{\prime} Q Q^{\prime}-2 \epsilon_{3} u Q^{2} Q^{\prime} R+6 \epsilon_{3} u^{2} Q^{2}-9 \epsilon_{2} \epsilon_{3} u^{2} u^{\prime} Q R \\
& \left.-6 u^{2} Q^{2} R^{2}+3 \epsilon_{2} u^{2} Q^{\prime 2}+3 \epsilon_{2} u^{2} Q Q^{\prime \prime}+4 \epsilon_{2} \epsilon_{3} u^{3} Q R^{\prime}+8 \epsilon_{2} \epsilon_{3} u^{3} Q^{\prime} R+5 \epsilon_{2} u^{4} R^{2}\right) t+\epsilon_{3} u^{\prime \prime} Q^{4} \\
& +\epsilon_{2} Q^{5} R^{\prime}-3 \epsilon_{2} u^{\prime} Q^{3} Q^{\prime}-\epsilon_{2} Q^{4} Q^{\prime} R+4 \epsilon_{2} \epsilon_{3} u u^{\prime 2} Q^{2}-2 \epsilon_{2} u Q^{4}+3 u u^{\prime} Q^{3} R+\epsilon_{2} \epsilon_{3} u Q^{4} R^{2} \\
& +3 \epsilon_{3} u Q^{2} Q^{\prime 2}-\epsilon_{3} u Q^{3} Q^{\prime \prime}-\epsilon_{2} \epsilon_{3} u^{2} u^{\prime \prime} Q^{2}-2 u^{2} Q^{3} R^{\prime}-5 \epsilon_{2} \epsilon_{3} u^{2} u^{\prime} Q Q^{\prime}-u^{2} Q^{2} Q^{\prime} R+2 u^{3} Q^{2} \\
& -3 \epsilon_{2} u^{3} u^{\prime} Q R-2 \epsilon_{3} u^{3} Q^{2} R^{2}+\epsilon_{2} \epsilon_{3} u^{3} Q^{\prime 2}+\epsilon_{2} \epsilon_{3} u^{3} Q Q^{\prime \prime}+\epsilon_{2} u^{4} Q R^{\prime}+2 \epsilon_{2} u^{4} Q^{\prime} R+\epsilon_{2} \epsilon_{3} u^{5} R^{2} .
\end{aligned}
$$

Now, we suppose that $M$ has pointwise 1-type Gauss map of the second kind. Then, equation (1.1) together with (4.2) gives

$$
\begin{equation*}
q^{-7 / 2} \epsilon_{4}\left(A_{1} \beta+A_{2} \beta^{\prime}+A_{3} \beta \times \beta^{\prime}\right)=f\left\{q^{-1 / 2}\left(\epsilon_{3} Q \beta^{\prime}-\left(\epsilon_{3} u+t\right) \beta \times \beta^{\prime}\right)+\mathbb{C}\right\} . \tag{4.3}
\end{equation*}
$$

If we respectively take the scalar product to equation (4.3) with $\beta, \beta^{\prime}$ and $\beta \times \beta^{\prime}$, then we have the following system of equations:

$$
\begin{gather*}
q^{-7 / 2} \epsilon_{4} A_{1}=f c_{1}  \tag{4.4}\\
q^{-7 / 2} \epsilon_{4} A_{2}=f\left(q^{-1 / 2} \epsilon_{3} Q+c_{2}\right)  \tag{4.5}\\
-q^{-7 / 2} \epsilon_{4} A_{3}=f\left(q^{-1 / 2}\left(\epsilon_{3} u+t\right)-c_{3}\right) \tag{4.6}
\end{gather*}
$$

where $c_{1}=\epsilon_{2}\langle\mathbb{C}, \beta\rangle, c_{2}=\epsilon_{3}\left\langle\mathbb{C}, \beta^{\prime}\right\rangle$ and $c_{3}=-\epsilon_{2} \epsilon_{3}\left\langle\mathbb{C}, \beta \times \beta^{\prime}\right\rangle$. Differentiating the functions $c_{1}, c_{2}$ and $c_{3}$ with respect to $s$, we get

$$
\begin{gather*}
c_{1}^{\prime}=\epsilon_{2} \epsilon_{3} c_{2},  \tag{4.7}\\
c_{1}+c_{2}^{\prime}-\epsilon_{3} c_{3} R=0  \tag{4.8}\\
c_{3}^{\prime}-\epsilon_{2} \epsilon_{3} c_{2} R=0 \tag{4.9}
\end{gather*}
$$

Combining equations (4.4), (4.5) and (4.6), we find

$$
\begin{gather*}
q\left(A_{2} c_{1}-A_{1} c_{2}\right)^{2}-Q^{2} A_{1}^{2}=0  \tag{4.10}\\
q\left(A_{1} c_{3}-A_{3} c_{1}\right)^{2}-A_{1}^{2}\left(\epsilon_{3} u+t\right)^{2}=0  \tag{4.11}\\
q\left(A_{2} c_{3}-A_{3} c_{2}\right)^{2}-\left(A_{2}\left(\epsilon_{3} u+t\right)+\epsilon_{3} Q A_{3}\right)^{2}=0 \tag{4.12}
\end{gather*}
$$

The left hand sides of (4.10), (4.11) and (4.12) are polynomials in $t$ with functions of $s$ as the coefficients. Therefore, all of them as functions of $s$ of polynomials in $t$ must be zero.

First, from the leading coefficient in $t$ of the left hand side of (4.10) with the help of (4.7), we get

$$
\begin{equation*}
c_{1} R=\mathrm{a} \text { constant } \tag{4.13}
\end{equation*}
$$

Also, the leading coefficient in $t$ of the left hand side of (4.11) gives

$$
\begin{equation*}
\epsilon_{3} \epsilon_{4}\left(c_{3}-c_{1} R\right)^{2} R^{2}=R^{2} \tag{4.14}
\end{equation*}
$$

If $\epsilon_{3} \epsilon_{4}=-1$, then $R$ is identically zero. In this case, the leading coefficient of the left hand side of (4.11) implies

$$
\left(c_{3}^{2}+1\right) Q^{\prime 2}=0
$$

from which, $Q$ is a constant. If we consider the coefficient of $t^{8}$ of the left hand side of (4.11), we also get

$$
\left(\left(3 \epsilon_{2} c_{3} u^{\prime}-2 \epsilon_{3} c_{1} Q\right)^{2}+9 u^{\prime 2}\right) Q^{2}=0
$$

If $Q \neq 0$, then we have

$$
\left(3 \epsilon_{2} c_{3} u^{\prime}-2 \epsilon_{3} c_{1} Q\right)^{2}+9 u^{\prime 2}=0
$$

which implies $u^{\prime}=0$. It follows the mean curvature $H$ vanishes on $M$. It contradicts that the Gauss map of $M$ is of pointwise 1-type of the second kind. Thus, $Q=0$. In turn, the mean curvature is zero identically, which is a contradiction, too. Consequently, we have

$$
\epsilon_{3} \epsilon_{4}=1
$$

Case 1. $R$ is not identically zero on $M$.
We now consider an open subset $\mathbf{U}=\{p \in M \mid R(p) \neq 0\}$. Suppose $\mathbf{U}$ is not empty. Then, (4.14) yields

$$
\begin{equation*}
\left(c_{3}-c_{1} R\right)^{2}=1 \tag{4.15}
\end{equation*}
$$

on U. Differentiating equation (4.15), we obtain $c_{3}$ is a constant on a connected component $\mathbf{U}_{0}$ of $\mathbf{U}$ because $c_{1} R$ is a constant. Therefore, (4.9) implies $c_{2}=0$ on $\mathbf{U}_{0}$. In view of (4.7), $c_{1}$ is a constant on $\mathbf{U}_{0}$. So, equation (4.8) yields $R$ is a constant on $\mathbf{U}_{0}$. By continuity of $R$ and connectedness of $M, R$ is a non-zero constant on $M$. Therefore, by (4.7), (4.8) and (4.13), $c_{1}$ and $c_{3}$ are constants and $c_{2}=0$ on $M$. Thus, equations (4.10), (4.11) and (4.12) can be rewritten as follows:

$$
\begin{gather*}
q c_{3}^{2} R^{2} A_{2}^{2}-Q^{2} A_{1}^{2}=0  \tag{4.16}\\
q c_{3}^{2}\left(A_{1}-\epsilon_{3} R A_{3}\right)^{2}-A_{1}^{2}\left(\epsilon_{3} u+t\right)^{2}=0  \tag{4.17}\\
q c_{3}^{2} A_{2}^{2}-\left(A_{2}\left(\epsilon_{3} u+t\right)+\epsilon_{3} Q A_{3}\right)^{2}=0 \tag{4.18}
\end{gather*}
$$

Moreover, combining equations (4.16) and (4.18), we get

$$
\begin{equation*}
Q^{2} A_{1}^{2}-R^{2}\left(A_{2}\left(\epsilon_{3} u+t\right)+\epsilon_{3} Q A_{3}\right)^{2}=0 \tag{4.19}
\end{equation*}
$$

From the leading coefficient of (4.17), we have

$$
\begin{equation*}
c_{3}^{2}\left(1-\epsilon_{3} R^{2}\right)^{2}=1 \tag{4.20}
\end{equation*}
$$

If we examine the coefficients of $t^{10}$ and $t^{9}$ of the left hand side of (4.17) with the help of (4.20), respectively, we get

$$
\begin{equation*}
2 Q^{2}-5 \epsilon_{2} Q^{2}+4 u^{\prime} Q R-6 \epsilon_{3} u^{\prime} Q R+\epsilon_{2} \epsilon_{3} Q^{2} R^{2}-2 \epsilon_{3} Q^{\prime 2}=0 \tag{4.21}
\end{equation*}
$$

and

$$
\begin{equation*}
2 Q Q^{\prime}-u^{\prime \prime} Q R-3 u^{\prime} Q^{\prime} R+\epsilon_{2} \epsilon_{3} Q Q^{\prime} R^{2}=0 \tag{4.22}
\end{equation*}
$$

Furthermore, considering the leading coefficient of the left hand side of (4.19), we obtain

$$
\begin{equation*}
\left(u^{\prime} R-\epsilon_{3} Q^{\prime \prime}\right)^{2}=Q^{2} . \tag{4.23}
\end{equation*}
$$

Without loss of generality, we may assume

$$
u^{\prime} R-\epsilon_{3} Q^{\prime \prime}=Q
$$

From the coefficients of $t^{8}$ and $t^{7}$ of the left hand side of (4.19), respectively, we have

$$
\begin{equation*}
Q^{2}\left(u^{\prime} Q-\epsilon_{3} u^{\prime 2} R-\epsilon_{2} \epsilon_{3} Q^{\prime 2} R\right)=0 \tag{4.24}
\end{equation*}
$$

and

$$
\begin{equation*}
Q^{2}\left(2 u^{\prime} R-\epsilon_{3} Q^{\prime}+2 \epsilon_{3} u R\right)=0 \tag{4.25}
\end{equation*}
$$

Suppose the open subset $\mathbf{O}=\{p \in M \mid Q(p) \neq 0\}$ is not empty. Then, (4.24) and (4.25) imply

$$
\begin{equation*}
u^{\prime} Q-\epsilon_{3} u^{\prime 2} R-\epsilon_{2} \epsilon_{3} Q^{\prime 2} R=0 \tag{4.26}
\end{equation*}
$$

and

$$
\begin{equation*}
2 u^{\prime} R-\epsilon_{3} Q^{\prime}+2 \epsilon_{3} u R=0 \tag{4.27}
\end{equation*}
$$

On the other hand, considering the coefficient of $t^{8}$ of the left hand side of (4.16) with the help of (4.21), (4.22) and (4.26), we obtain

$$
\begin{equation*}
\epsilon_{2} Q^{3}-Q^{3}=0 \tag{4.28}
\end{equation*}
$$

If we think of the non-empty subset $\mathbf{O}, \epsilon_{2}$ must be 1 . Therefore, (4.21) implies

$$
\begin{equation*}
-3 Q^{2}+4 u^{\prime} Q R-6 \epsilon_{3} u^{\prime} Q R+\epsilon_{3} Q^{2} R^{2}-2 \epsilon_{3} Q^{\prime 2}=0 \tag{4.29}
\end{equation*}
$$

Differentiating equation (4.29) with respect to $s$, we obtain

$$
\begin{align*}
& -3 Q Q^{\prime}+2 u^{\prime \prime} Q R+2 u^{\prime} Q^{\prime} R-3 \epsilon_{3} u^{\prime \prime} Q R  \tag{4.30}\\
& -3 \epsilon_{3} u^{\prime} Q^{\prime} R+\epsilon_{3} Q Q^{\prime} R^{2}-2 \epsilon_{3} Q^{\prime} Q^{\prime \prime}=0 .
\end{align*}
$$

Suppose $\epsilon_{3}=1$. Combining equations (4.22) and (4.30), we get $Q Q^{\prime}=0$. Hence, $Q$ is a non-zero constant on a connected component $\mathbf{O}_{0}$ of $\mathbf{O}$. By connectedness of $M$ and continuity of $Q, Q$ is a non-zero constant on $M$. Therefore, equations (4.22) and (4.27) respectively give the following

$$
u^{\prime \prime}=0 \quad \text { and } \quad u^{\prime}+u=0 .
$$

Thus, we have $u^{\prime}=0$. It implies that $Q=0$ because of (4.23), a contradiction. So, $\mathbf{O}$ is empty and thus, $Q \equiv 0$ on $M$.
Let us now assume $\epsilon_{3}=-1$. Then, (4.30) with the help of (4.22) implies

$$
\begin{equation*}
Q^{\prime}\left(3 Q-4 u^{\prime} R-2 Q R^{2}\right)=0 . \tag{4.31}
\end{equation*}
$$

Consider an open subset $\mathbf{O}_{1}=\left\{p \in \mathbf{O} \mid Q^{\prime}(p) \neq 0\right\}$ and suppose $\mathbf{O}_{1}$ is not empty. Then (4.31) gives

$$
\begin{equation*}
3 Q-4 u^{\prime} R-2 Q R^{2}=0 \tag{4.32}
\end{equation*}
$$

Since $u^{\prime} R+Q^{\prime \prime}=Q$, we have the following differential equation from (4.32)

$$
\begin{equation*}
Q^{\prime \prime}-k^{2} Q=0 \tag{4.33}
\end{equation*}
$$

where $k^{2}=\frac{1+2 R^{2}}{4}(k>0)$, a constant. Thus,

$$
\begin{equation*}
u^{\prime} R=\left(1-k^{2}\right) Q, \tag{4.34}
\end{equation*}
$$

or, $Q$ is given by

$$
Q=\tilde{K}_{1} \cosh k s+\tilde{K}_{2} \sinh k s
$$

for some constants $\tilde{K}_{1}$ and $\tilde{K}_{2}$. Together (4.34) with (4.22), we have

$$
2-4\left(1-k^{2}\right)-R^{2}=0,
$$

or, using $k^{2}=\frac{1+2 R^{2}}{4}$, we get

$$
R^{2}=1
$$

on $\mathbf{O}_{1}$. Putting this into (4.26) and using (4.34), we get

$$
Q=0,
$$

which is a contradiction on $\mathbf{O}_{1}$. Therefore, the open subset $\mathbf{O}_{1}$ is empty and $Q$ is a non-zero constant on a connected component $\mathbf{O}_{0}$ of $\mathbf{O}$. Again, connectedness of $M$ and continuity of $Q$ imply $Q$ is a non-zero constant on $M$. Hence, we have $u^{\prime} R=Q$. Together with (4.26), $u^{\prime}=0$ is induced and we get $Q=0$, a
contradiction. Consequently, the open subset $\mathbf{O}$ is empty and $Q \equiv 0$ on $M$. Hence, no matter what cases of $\epsilon_{3}$ may be, we have $Q$ is zero on $M$.

Therefore, $u$ is a non-zero constant by virtue of (4.23) and the first equation of (4.1) with the help of the fact that $\alpha$ is non-null.

On the other hand, it follows from the second equation in (4.1) that

$$
\beta^{\prime \prime \prime}+\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right) \beta^{\prime}=0
$$

Let us give the initial conditions $\beta(0)=\left(a_{1}, 0, a_{2}\right), \beta^{\prime}(0)=\left(b_{1}, b_{2}, 0\right)$ and $\beta^{\prime \prime}(0)=$ $-\epsilon_{2} \epsilon_{3}\left(a_{1}+a_{2} b_{2} R, a_{2} b_{1} R, a_{2}+a_{1} b_{2} R\right)$ of the above differential equation, where $a_{1}, a_{2}, b_{1}$ and $b_{2}$ are some constants satisfying $-a_{1}^{2}+a_{2}^{2}=\epsilon_{2},-b_{1}^{2}+b_{2}^{2}=\epsilon_{3}$ and $a_{1} b_{1}=0$ with $\left(a_{1}, b_{1}\right) \neq(0,0)$.

Considering equation (4.20), we only have the cases: $\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right)>0$ or $\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right)<0$.

First, if $\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right)>0$, then we have $\beta^{\prime \prime \prime}(s)+a^{2} \beta^{\prime}(s)=0$ and its solution $\beta(s)$ is

$$
\begin{align*}
\beta(s)=( & \frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(\epsilon_{3} a_{1} R^{2}+a_{2} b_{2} R\right)+\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(a_{1}+a_{2} b_{2} R\right) \cos a s+\frac{b_{1}}{a} \sin a s \\
& -\frac{\epsilon_{2} \epsilon_{3}}{a^{2}} a_{2} b_{1} R+\frac{\epsilon_{2} \epsilon_{3}}{a^{2}} a_{2} b_{1} R \cos a s+\frac{b_{2}}{a} \sin a s  \tag{4.35}\\
& \left.-\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(\epsilon_{3} a_{2} R^{2}+a_{1} b_{2} R\right)+\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(a_{2}+a_{1} b_{2} R\right) \cos a s\right)
\end{align*}
$$

where $a=\sqrt{\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right)}$.
If $\epsilon_{2} \epsilon_{3}\left(1-\epsilon_{3} R^{2}\right)<0$, then the solution of $\beta^{\prime \prime \prime}(s)-a^{2} \beta^{\prime}(s)=0$ is obtained as follows

$$
\begin{align*}
\beta(s)=( & \frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(\epsilon_{3} a_{1} R^{2}+a_{2} b_{2} R\right)-\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(a_{1}+a_{2} b_{2} R\right) \cosh a s+\frac{b_{1}}{a} \sinh a s \\
& \frac{\epsilon_{2} \epsilon_{3}}{a^{2}} a_{2} b_{1} R-\frac{\epsilon_{2} \epsilon_{3}}{a^{2}} a_{2} b_{1} R \cosh a s+\frac{b_{2}}{a} \sinh a s  \tag{4.36}\\
& \left.\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(\epsilon_{3} a_{2} R^{2}+a_{1} b_{2} R\right)-\frac{\epsilon_{2} \epsilon_{3}}{a^{2}}\left(a_{2}+a_{1} b_{2} R\right) \cosh a s\right)
\end{align*}
$$

where $a=\sqrt{\epsilon_{2} \epsilon_{3}\left(\epsilon_{3} R^{2}-1\right)}$.
By applying the first equation of (4.1), $\alpha^{\prime}=\epsilon_{3} u \beta^{\prime}$ for some non-zero constant $u$. Therefore, we can easily obtain the base curve $\alpha(s)$ by means of $\beta(s)$ of the form (4.35) or (4.36).

Definition 4.1. A non-cylindrical ruled surface $M$ generated by a base curve $\alpha(s)$ and the director vector field $\beta(s)$ satisfying (4.35) or (4.36) is called a rotational ruled surface of type $I$ or a rotational ruled surface of type II, respectively.

Case 2. $R \equiv 0$ on $M$.
From (4.9), we see that $c_{3}$ is a constant. Combining equations (4.7) and (4.8), we have $c_{1}+\epsilon_{2} \epsilon_{3} c_{1}^{\prime \prime}=0$. Thus, depending upon the sign of $\epsilon_{2}$ and $\epsilon_{3}$, we get

$$
\begin{equation*}
c_{1}=K_{1} \cosh s+K_{2} \sinh s \quad \text { or } \quad c_{1}=K_{3} \sin \left(s+s_{0}\right) \tag{4.37}
\end{equation*}
$$

for some constants $K_{i}(i=1,2,3)$ and $s_{0}$.
If we think of the leading coefficient of the left hand side of (4.11) with $c_{1}$ as above, we have

$$
\left(c_{3}^{2}-1\right) Q^{\prime 2}=0
$$

Suppose $c_{3}^{2} \neq 1$. Then $Q^{\prime}=0$, that is, $Q$ is a constant. If $Q=0$, the mean curvature $H$ vanishes on $M$, which is a contradiction. Therefore, $Q$ is a non-zero constant.

If we consider the leading coefficients of the left hand sides of (4.10) and (4.11), respectively, we get

$$
\begin{equation*}
c_{1} u^{\prime \prime}+3 \epsilon_{2} c_{2} u^{\prime}=0 \quad \text { and } \quad\left(3 \epsilon_{2} c_{3} u^{\prime}+2 \epsilon_{3} c_{1} Q\right)^{2}=9 u^{\prime 2} \tag{4.38}
\end{equation*}
$$

Since $c_{1}^{\prime}=\epsilon_{2} \epsilon_{3} c_{2}$, the first equation of (4.38) implies

$$
c_{1} u^{\prime \prime}+3 \epsilon_{3} c_{1}^{\prime} u^{\prime}=0
$$

Suppose $c_{1}$ is non-trivial. Then, the solution of the above differential equation is given by

$$
\begin{equation*}
u^{\prime}=k_{1} c_{1}^{ \pm 3} \tag{4.39}
\end{equation*}
$$

where $k_{1}$ is a constant. Putting (4.37) and (4.39) into the second equation of (4.38), we obtain $k_{1}=0$ and $K_{i}=0(i=1,2,3)$, which is a contradiction. Thus, we have $c_{1}=0$. From (4.4), we get $A_{1}=0$. It implies that $u^{\prime}=0$, that is, $u$ is a constant. Hence we see that the mean curvature $H$ vanishes identically that is again a contradiction. Consequently, we get $c_{3}^{2}=1$.

Now, if we consider the leading coefficient of the left hand side of (4.10) with the help of (4.7), we have

$$
c_{1} Q^{\prime \prime}+2 c_{1}^{\prime} Q^{\prime}=0
$$

Suppose $c_{1}$ is non-zero. Then the solution of the above equation is given by

$$
\begin{equation*}
Q^{\prime}=k_{2} c_{1}^{-2} \tag{4.40}
\end{equation*}
$$

where $k_{2}$ is a constant. Since the coefficient of $t^{9}$ of the left hand side of (4.11) with $c_{3}^{2}=1$ is zero, we have

$$
Q^{\prime}\left(Q Q^{\prime \prime}+Q^{2}+2 \epsilon_{2} \epsilon_{3} Q^{2}\right)=0
$$

Consider an open set $\mathbf{V}=\left\{p \in M \mid Q^{\prime}(p) \neq 0\right\}$. Suppose $\mathbf{V}$ is not empty. Then we get

$$
\begin{equation*}
Q Q^{\prime \prime}+Q^{\prime 2}+2 \epsilon_{2} \epsilon_{3} Q^{2}=0 \tag{4.41}
\end{equation*}
$$

Equations (4.37), (4.40) and (4.41) lead to $k_{2}=0$ and so $Q^{\prime}=0$, a contradiction. Therefore, $\mathbf{V}$ is empty and $Q$ is a non-zero constant on $M$ since $M$ is not minimal. Thus, the fact that the leading coefficient of the left hand side of (4.10) is zero gives $c_{1} u^{\prime \prime}+3 \epsilon_{3} c_{1}^{\prime} u^{\prime}=0$ because of (4.7). If $c_{1}$ is non-trivial, we have a solution of the form (4.39). Since $c_{3}^{2}=1$, from the leading coefficient of (4.11), we also get

$$
\begin{equation*}
c_{1} Q \pm 3 \epsilon_{2} \epsilon_{3} u^{\prime}=0 \tag{4.42}
\end{equation*}
$$

Putting (4.37) and (4.39) into (4.42), we obtain $c_{1}=0$.
Similarly as before, (4.4) yields $u^{\prime}=0$. Therefore, the mean curvature $H$ vanishes, which is a contradiction. As a consequence, the case of $R=0$ can never occur.

Consequently, we have
Theorem 4.1. Let $M$ be a non-cylindrical ruled surface of type $M_{+}^{1}, M_{-}^{1}$ or $M_{+}^{3}$ in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Suppose that $M$ has pointwise 1-type Gauss map of the second kind. Then, $M$ is an open part of a rotational ruled surfaces of type I or type II.

Now we examine a non-cylindrical ruled surface of type $M_{+}^{2}$ or $M_{-}^{2}$ with pointwise 1 -type Gauss map of the second kind.

Let $M$ be a non-cylindrical ruled surface of type $M_{+}^{2}$ or $M_{-}^{2}$. Then, the parametrization for $M$ is given by

$$
x(s, t)=\alpha(s)+t \beta(s)
$$

such that $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=\epsilon_{1}(= \pm 1),\langle\beta, \beta\rangle=1,\left\langle\alpha^{\prime}, \beta\right\rangle=0$ and $\beta^{\prime}$ is null. Let us also put

$$
q=\left\|x_{s}\right\|^{2}=\epsilon_{4}\left\langle x_{s}, x_{s}\right\rangle, \quad u=\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle .
$$

On the other hand, it is easy to see that $\beta \times \beta^{\prime}$ is null. Since the null vector fields $\beta^{\prime}$ and $\beta \times \beta^{\prime}$ are orthogonal, we may take

$$
\begin{equation*}
\beta^{\prime}=\beta \times \beta^{\prime} . \tag{4.43}
\end{equation*}
$$

Moreover, we may assume $\beta(0)=(0,0,1)$. Thus, $\beta(s)$ is given by

$$
\begin{equation*}
\beta(s)=(a s, a s, 1) \tag{4.44}
\end{equation*}
$$

for a non-zero constant $a$. For an orthonormal frame $\left\{\alpha^{\prime}, \beta, \alpha^{\prime} \times \beta\right\}$ along $\alpha$, we have

$$
\begin{equation*}
\beta^{\prime}=\epsilon_{1} u\left(\alpha^{\prime}-\alpha^{\prime} \times \beta\right) \quad \text { and } \quad \alpha^{\prime \prime}=-u \beta+\frac{u^{\prime}}{u} \alpha^{\prime} \times \beta \tag{4.45}
\end{equation*}
$$

Let $\alpha(s)=\left(\alpha_{1}(s), \alpha_{2}(s), \alpha_{3}(s)\right)$. Since $\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle=u$ and $\left\langle\alpha^{\prime}, \beta\right\rangle=0$, we obtain

$$
\begin{equation*}
\alpha_{1}^{\prime}-\alpha_{2}^{\prime}=-\frac{u}{a} \quad \text { and } \quad \alpha_{3}^{\prime}=-u s \tag{4.46}
\end{equation*}
$$

Equation (4.46) together with $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=\epsilon_{1}$ implies

$$
\begin{equation*}
\alpha_{1}^{\prime}+\alpha_{2}^{\prime}=\frac{a \epsilon_{1}}{u}-a u s^{2} \tag{4.47}
\end{equation*}
$$

Combining equations (4.46) and (4.47), we get

$$
\begin{equation*}
\alpha^{\prime}(s)=\left(\frac{1}{2}\left(\frac{a \epsilon_{1}}{u}-a u s^{2}-\frac{u}{a}\right), \frac{1}{2}\left(\frac{a \epsilon_{1}}{u}-a u s^{2}+\frac{u}{a}\right),-u s\right) . \tag{4.48}
\end{equation*}
$$

On the other hand, the Gauss map $G$ of $M$ is given by

$$
\begin{equation*}
G=q^{-\frac{1}{2}}\left(A-t \beta^{\prime}\right) \tag{4.49}
\end{equation*}
$$

where we put $A=\alpha^{\prime} \times \beta$. By a straightforward computation, the Laplacian $\Delta G$ of the Gauss map $G$ can be expressed as ([14])
$\Delta G=q^{-\frac{7}{2}}\left(\left(-2 u^{2} q+u^{\prime \prime} t q-4 \epsilon_{4} u^{2} t^{2}\right)\left(A-t \beta^{\prime}\right)-\epsilon_{4} u \beta^{\prime} q^{2}+3 u^{\prime} t A^{\prime} q-\epsilon_{4} A^{\prime \prime} q^{2}\right)$.
Now we suppose that $M$ has pointwise 1-type Gauss map of the second kind. Then, we obtain

$$
\begin{align*}
& \left(-2 u^{2} q+u^{\prime \prime} t q-4 \epsilon_{4} u^{\prime 2} t^{2}\right)\left(A-t \beta^{\prime}\right)-\epsilon_{4} u \beta^{\prime} q^{2}+3 u^{\prime} t A^{\prime} q-\epsilon_{4} A^{\prime \prime} q^{2} \\
= & f\left(q^{3}\left(A-t \beta^{\prime}\right)+q^{\frac{7}{2}} \mathbb{C}\right) \tag{4.50}
\end{align*}
$$

for some non-zero smooth function $f$ and a constant vector $\mathbb{C}$.
If we take the scalar product to equation (4.50) with $\alpha^{\prime}, \beta$ and $\alpha^{\prime} \times \beta$, respectively, then we obtain the following:

$$
\begin{gather*}
B_{1}=f\left(-q^{3} t u+q^{\frac{7}{2}}\left\langle\mathbb{C}, \alpha^{\prime}\right\rangle\right)  \tag{4.51}\\
B_{2}=f q^{\frac{7}{2}}\langle\mathbb{C}, \beta\rangle  \tag{4.52}\\
B_{3}=f\left(-\epsilon_{1} q^{3}-u q^{3} t+q^{\frac{7}{2}}\left\langle\mathbb{C}, \alpha^{\prime} \times \beta\right\rangle\right) \tag{4.53}
\end{gather*}
$$

where

$$
\begin{aligned}
& B_{1}=2 u^{3} q t-u u^{\prime \prime} t^{2} q+4 \epsilon_{4} u u^{\prime 2} t^{3}+3 \epsilon_{1} \frac{u^{\prime 2}}{u} t q-\epsilon_{1} \epsilon_{4} \frac{u^{\prime \prime}}{u} q^{2}+\epsilon_{1} \epsilon_{4} \frac{u^{\prime 2}}{u^{2}} q^{2} \\
& B_{2}=-3 u u^{\prime} t q+2 \epsilon_{4} u^{\prime} q^{2} \\
& B_{3}=2 \epsilon_{1} u^{2} q-\epsilon_{1} u^{\prime \prime} t q+4 \epsilon_{1} \epsilon_{4} u^{\prime 2} t^{2}+2 u^{3} q t-u u^{\prime \prime} t^{2} q+4 \epsilon_{4} u u^{\prime 2} t^{3}+\epsilon_{1} \epsilon_{4} q^{2} \frac{u^{\prime 2}}{u^{2}}
\end{aligned}
$$

If we put $\mathbb{C}=c_{1} \alpha^{\prime}+c_{2} \beta+c_{3} \alpha^{\prime} \times \beta$, (4.51)-(4.53) imply

$$
\begin{gather*}
\left(c_{2} B_{1}-\epsilon_{1} c_{1} B_{2}\right)^{2} q=u^{2} t^{2} B_{2}^{2},  \tag{4.54}\\
\left(c_{2} B_{3}+\epsilon_{1} c_{3} B_{2}\right)^{2} q=\left(\epsilon_{1}+u t\right)^{2} B_{2}^{2},  \tag{4.55}\\
\left(u t B_{3}-\left(\epsilon_{1}+u t\right) B_{1}\right)^{2}=q\left(c_{1} B_{3}+c_{3} B_{1}\right)^{2}, \tag{4.56}
\end{gather*}
$$

which are polynomials in $t$ with functions of $s$ as the coefficients. Hence, the leading coefficient of the left hand side of (4.54) must be zero, which means $c_{2}^{2} u^{3}\left(u u^{\prime \prime}-\right.$ $\left.2 u^{\prime 2}\right)^{2}=0$. Because $u \neq 0$, we get

$$
\begin{equation*}
c_{2}^{2}\left(u u^{\prime \prime}-2 u^{\prime 2}\right)^{2}=0 . \tag{4.57}
\end{equation*}
$$

Consider an open subset $\mathbf{U}=\left\{p \in M \mid\left(u u^{\prime \prime}-2 u^{\prime 2}\right)(p) \neq 0\right\}$. Suppose $\mathbf{U}$ is not empty. Then, $c_{2}=0$ on $\mathbf{U}$. Therefore, equation (4.54) can be reduced to

$$
\begin{equation*}
B_{2}^{2}\left(c_{1}^{2} q-u^{2} t^{2}\right)=0 . \tag{4.58}
\end{equation*}
$$

Since the leading coefficient of the left hand side of (4.58) must be zero, $u^{6} u^{\prime 2}=0$ on $\mathbf{U}$, from which, $u^{\prime}=0$ on $\mathbf{U}$. It is a contradiction. Thus, $\mathbf{U}$ is empty and we have

$$
\begin{equation*}
u u^{\prime \prime}-2 u^{\prime 2}=0 . \tag{4.59}
\end{equation*}
$$

Suppose there is a point $s_{0} \in$ domain $(\alpha)$ such that $u^{\prime}\left(s_{0}\right)=0$. Then, (4.54) implies $c_{2}=0$. Also, (4.59) gives $u^{\prime \prime}\left(s_{0}\right)=0$. If we evaluate the left hand side of (4.56), it turns out to be zero at $s_{0}$ and thus

$$
\epsilon_{1} c_{1}+\left(c_{1}+c_{3}\right) u\left(s_{0}\right) t=0
$$

It holds for each $t$ and hence $c_{1}=c_{3}=0$, that is, $\mathbb{C}$ is zero vector, which is a contradiction. Therefore, $u^{\prime} \neq 0$ everywhere. From (4.59), we get

$$
\frac{u^{\prime \prime}}{u^{\prime}}-\frac{2 u^{\prime}}{u}=0
$$

from which,

$$
u(s)=\frac{1}{b s+c}
$$

for some constants $b \neq 0$ and $c$. Thus, from (4.48), the base curve $\alpha(s)$ is given by

$$
\begin{align*}
\alpha(s)= & \frac{1}{2}\left(a \epsilon_{1}\left(\frac{b}{2} s^{2}+c s\right)-\frac{a}{2 b} s^{2}+\frac{a c}{b^{2}} s-\left(\frac{a c^{2}}{b^{3}}+\frac{1}{a b}\right) \ln |b s+c|+d_{1},\right. \\
& a \epsilon_{1}\left(\frac{b}{2} s^{2}+c s\right)-\frac{a}{2 b} s^{2}+\frac{a c}{b^{2}} s-\left(\frac{a c^{2}}{b^{3}}-\frac{1}{a b}\right) \ln |b s+c|+d_{2},  \tag{4.60}\\
& \left.-\frac{2}{b} s+\frac{2 c}{b} \ln |b s+c|+d_{3}\right),
\end{align*}
$$

where $d_{i}(i=1,2,3)$ are some integration constants.
Definition 4.2. A ruled surface $M$ generated by the base curve of the form (4.60) and the director vector field given by (4.44) is called a transcendental ruled surface.

Consequently, we have
Theorem 4.2. Let $M$ be a non-cylindrical ruled surface of type $M_{+}^{2}$ or $M_{-}^{2}$ in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Suppose that the Gauss map $G$ of $M$ is of pointwise 1-type of the second kind. Then, $M$ is an open portion of a transcendental ruled surface.

Combining Theorem 4.1, Theorem 4.2 and the results of [14], we have
Theorem 4.3. (Classification). Let $M$ be a non-cylindrical ruled surface over a non-null base curve in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Then, $M$ has pointwise 1-type Gauss map if and only if $M$ is an open part of a helicoid of the first kind, a helicoid of the second kind, a helicoid of the third kind, the conjugate of Enneper's surface of the second kind, a rotational ruled surface of type I or II, or a transcendental ruled surface.

## 5. Null Scrolls

In this section, we study a null scroll with pointwise 1-type Gauss map in Minkowski 3 -space $\mathbb{E}_{1}^{3}$. We mainly focus to prove the following theorem.

Theorem 5.1. Let $M$ be a null scroll with pointwise 1-type Gauss map of the second kind in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Then, $M$ is an open part of a Minkowski plane.

Proof. Let $\alpha=\alpha(s)$ be a null curve in $\mathbb{E}_{1}^{3}$ and $\beta=\beta(s)$ a null vector field satisfying $\left\langle\alpha^{\prime}, \beta\right\rangle=1$ along $\alpha$. For a null scroll $M$ parameterized by

$$
x=x(s, t)=\alpha(s)+t \beta(s)
$$

we have the natural coordinate frame $\left\{x_{s}, x_{t}\right\}$ given by

$$
x_{s}=\alpha^{\prime}+t \beta^{\prime} \quad \text { and } \quad x_{t}=\beta(s)
$$

Furthermore, we may choose an appropriate parameter $s$ in such a way that $u=$ $\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle=0$, which is possible if the base curve $\alpha$ is chosen as a null geodesic of $M$. Again, we define the smooth functions $q$ and $v$ as follows:

$$
q=\left\langle x_{s}, x_{s}\right\rangle \quad \text { and } \quad v=\left\langle\beta^{\prime}, \beta^{\prime}\right\rangle
$$

On the other hand, the Gauss map $G$ of $M$ is determined by

$$
\begin{equation*}
G=x_{s} \times x_{t}=\alpha^{\prime} \times \beta+t \beta^{\prime} \times \beta \tag{5.1}
\end{equation*}
$$

and the Laplacian of the Gauss map $G$ is obtained by ([14])

$$
\begin{equation*}
\Delta G=-2 \beta^{\prime \prime} \times \beta+2 v t \beta^{\prime} \times \beta . \tag{5.2}
\end{equation*}
$$

In terms of the pseudo-orthonormal frame $\left\{\alpha^{\prime}, \beta, \alpha^{\prime} \times \beta\right\}$, we obtain

$$
\begin{equation*}
\beta^{\prime}=-Q \alpha^{\prime} \times \beta, \quad \beta^{\prime} \times \beta=Q \beta \quad \text { and } \quad \beta^{\prime \prime} \times \beta=R \beta-v \alpha^{\prime} \times \beta, \tag{5.3}
\end{equation*}
$$

where $Q=\left\langle\alpha^{\prime}, \beta^{\prime} \times \beta\right\rangle$ and $R=\left\langle\alpha^{\prime}, \beta^{\prime \prime} \times \beta\right\rangle$.
We now suppose that $M$ has pointwise 1-type Gauss map of the second kind. Then, with the help of (5.3), we have

$$
\begin{equation*}
(2 v t Q-2 R) \beta+2 v \alpha^{\prime} \times \beta=f\left(\alpha^{\prime} \times \beta+t Q \beta+\mathbb{C}\right) \tag{5.4}
\end{equation*}
$$

for some non-zero smooth function $f$ and a constant vector $\mathbb{C}$.
If we take the scalar product to equation (5.4) with $\alpha^{\prime}, \beta$ and $\alpha^{\prime} \times \beta$, respectively, then we have the following system of equations:

$$
\begin{equation*}
2 v t Q-2 R=f\left(Q t+c_{2}\right), \tag{5.5}
\end{equation*}
$$

where $c_{1}=\langle\mathbb{C}, \beta\rangle, c_{2}=\left\langle\mathbb{C}, \alpha^{\prime}\right\rangle$ and $c_{3}=\left\langle\mathbb{C}, \alpha^{\prime} \times \beta\right\rangle$. Clearly, (5.6) gives $c_{1}=0$. From (5.7), the function $f$ depends only on the parameter $s$. Therefore, from (5.5), we can obtain

$$
(2 v-f) Q=0 \quad \text { and } \quad 2 R+f c_{2}=0 .
$$

Consider an open subset $\mathbf{U}=\{p \in M \mid Q(p) \neq 0\}$. Suppose that $\mathbf{U}$ is not empty. Then $f=2 v$ on $\mathbf{U}$ which implies $c_{3} f=0$ by (5.7) and thus $c_{3}=0$. Therefore, the constant vector $\mathbb{C}$ can be written as $\mathbb{C}=c_{2} \beta$. Differentiating the constant vector $\mathbb{C}$ with respect to $s$, we have $0=c_{2}^{\prime} \beta(s)+c_{2} \beta^{\prime}(s)$. Since $\beta$ and $\beta^{\prime}$ are linearly independent for each $s, c_{2}$ vanishes, which is a contradiction because $\mathbb{C}$ is not zero vector. Therefore, the open subset $\mathbf{U}$ is empty, that is, $Q=0$. Hence, (5.3) gives $\beta$ is a constant vector. It follows that $R=0$ and $v=0$. Thus, $\Delta G=0$. Since the Gauss map is of pointwise 1 -type of the second kind, we may get $G=-\mathbb{C}$. Thus, the surface $M$ is an open part of a Minkowski plane. Consequently, the proof is completed.

Combining Theorem 5.1 and the results in [14], we have
Theorem 5.2. (Classification). Let $M$ be a null scroll with pointwise 1-type Gauss map in Minkowski 3-space $\mathbb{E}_{1}^{3}$. Then, $M$ is an open part of a Minkowski plane or a B-scroll.

Remark. Summing up all the cases, Theorem 3.2, Theorem 3.3, Theorem 4.3 and Theorem 5.2, we have a complete classification theorem of the ruled surfaces in Minkowski 3-space $\mathbb{E}_{1}^{3}$ with pointwise 1-type Gauss map, which is described in Section 1.


Fig. 1.


Fig. 2.


Fig. 3.

## Acknowledgment

The authors would like to express sincere thanks to the referee for his valuable suggestions and comments toward improvement of the paper.

## References

1. C. Baikoussis and D. E. Blair, On the Gauss map of ruled surfaces, Glasgow Math. J., 34 (1992), 355-359.
2. C. Baikoussis and L. Verstraelen, On the Gauss map of helicoidal surfaces, Rend. Sem. Mat. Messina Ser. II. 16 1993, pp. 31-42.
3. B.-Y. Chen, Total mean curvature and submanifolds of finite type, World Scientific Publ., New Jersey, 1984.
4. B.-Y. Chen, Finite type submanifolds and generalizations, University of Rome, 1985.
5. B.-Y. Chen, M. Choi and Y. H. Kim, Surfaces of revolution with pointwise 1-type Gauss map, J. Korean Math. Soc., 42 (2005), 447-455.
6. B.-Y. Chen and P. Piccinni, Submanifolds with finite type Gauss map, Bull. Austral. Math. Soc., 35 (1987), 161-186.
7. M. Choi and Y. H. Kim, Characterization of the helicoid as ruled surfaces with pointwise 1-type Gauss map, Bull. Korean Math. Soc., 38 (2001), 753-761.
8. M. Choi, D.-S. Kim and Y. H. Kim, Helicoidal surfaces with pointwise 1-type Gauss map, J. Korean Math. Soc., 46 (2009), 215-223.
9. M. Choi, Y. H. Kim, H. Liu and D. W. Yoon, Helicoidal surfaces and their Gauss map in Minkowski 3-space, Bull. Korean Math. Soc., to appear.
10. M. Choi, Y. H. Kim and D. W. Yoon, Classification of ruled surfaces with pointwise 1-type Gauss map, Taiwanese J. Math., to appear.
11. U. Dursun, Hypersurfaces with pointwise 1-type Gauss map, Taiwanese J. Math., 11 (2007), 1407-1416.
12. U-H. Ki, D.-S. Kim, Y. H. Kim and Y.-M. Roh, Surfaces of revolution with pointwise 1-type Gauss map in Minkowski 3-space, Taiwanese J. Math., 13 (2009), 317-338.
13. D.-S. Kim, Y. H. Kim and D. W. Yoon, Finite type ruled surfaces in LorentzMinkowski space, Taiwanese J. Math., 11 (2007), 1-13.
14. Y. H. Kim and D. W. Yoon, Ruled surfaces with pointwise 1-type Gauss map, J. Geom. Phys., 34 (2000), 191-205.
15. Y. H. Kim and D. W. Yoon, On the Gauss map of ruled surfaces in Minkowski space, Rocky Mountain J. Math., 35 (2005), 1555-1581.

Miekyung Choi ${ }^{1}$, Young Ho Kim ${ }^{2}$ and Dae Won Yoon ${ }^{3}$
${ }^{1,3}$ Department of Mathematics education and RINS
Gyeongsang National University
Jinju 660-701
Korea
E-email: mkchoi@knu.ac.kr dwyoon@gnu.ac.kr
${ }^{2}$ Department of Mathematics
Kyungpook National University
Taegu 702-701
Korea
E-email: yhkim@knu.ac.kr


[^0]:    Received April 23, 2009, accepted December 11, 2009.
    Communicated by B. Y. Chen.
    2000 Mathematics Subject Classification: 53A35, 53B30.
    Key words and phrases: Ruled surface, Minimal surface, Null scroll, Minkowski space, Pointwise 1-type Gauss map.
    ${ }^{1}$ This work was supported by the Korea Research Foundation Grant funded by the Korean Government (MOEHRD, Basic Research Promotion Fund)(KRF-2008-355-C00003).
    ${ }^{2}$ Supported by KOSEF R01-2007-000-20014-0 (2007).

