# CLASSIFICATION THEOREMS FOR SPACE-LIKE SURFACES IN 4-DIMENSIONAL INDEFINITE SPACE FORMS WITH INDEX 2 

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#### Abstract

Surfaces in 4D Riemannian space forms have been investigated extensively. In contrast, only few results are known for surfaces in 4D neutral indefinite space forms $R_{2}^{4}(c)$. Thus, in this paper we study space-like surfaces in $R_{2}^{4}(c)$ satisfying certain simple geometric properties. In particular, we classify space-like surfaces in $\mathbb{E}_{2}^{4}$ with constant mean and Gauss curvatures and null normal curvature. We also classify Wintgen ideal surfaces in $R_{2}^{4}(c)$ whose Gauss and normal curvatures satisfy $K^{D}=2 K$.


## 1. Introduction

Let $\mathbb{E}_{t}^{m}$ denote the pseudo-Euclidean $m$-space equipped with pseudo-Euclidean metric of index $t$ given by

$$
\begin{equation*}
g_{t}=-\sum_{i=1}^{t} d x_{i}^{2}+\sum_{j=t+1}^{n} d x_{j}^{2} \tag{1.1}
\end{equation*}
$$

where $\left(x_{1}, \ldots, x_{m}\right)$ is a rectangular coordinate system of $\mathbb{E}_{t}^{m}$. We put

$$
\begin{align*}
& S_{s}^{k}(c)=\left\{x \in \mathbb{E}_{s}^{k+1}:\langle x, x\rangle=c^{-1}>0\right\}  \tag{1.2}\\
& H_{s}^{k}(c)=\left\{x \in \mathbb{E}_{s+1}^{k+1}:\langle x, x\rangle=c^{-1}<0\right\} \tag{1.3}
\end{align*}
$$

where $\langle$,$\rangle is the associated inner product. Then S_{s}^{k}(c)$ and $H_{s}^{k}(c)$ are pseudoRiemannian manifolds of constant curvature $c$ and with index $s$, which are known as pseudo-Riemannian $k$-sphere and the pseudo-hyperbolic $k$-space, respectively. The pseudo-Riemannian manifolds $\mathbb{E}_{s}^{k}, S_{s}^{k}(c)$ and $H_{s}^{k}(-c)$ are called indefinite space forms, denoted by $R_{s}^{k}$.

[^0]Surfaces in 4-dimensional Riemannian space forms have been investigated very extensively (see, for instance, [1, 2, 3]). In contrast, only few results are known for surfaces in 4-dimensional neutral indefinite space forms $R_{2}^{4}(c)$ of constant curvature $c$ and index 2. Thus, we study in this paper space-like surfaces in $R_{2}^{4}(c)$ satisfying some simple geometric properties.

In Section 2 of this paper we provide basic definitions and formulas. In Section 3 we completely classify space-like surfaces in $\mathbb{E}_{2}^{4}$ with constant mean and Gauss curvatures and null normal curvature. In Section 4, we present a result of Sasaki and the precise expression of a minimal immersion $\psi_{\mathcal{B}}$ of the hyperbolic plane $H^{2}\left(-\frac{1}{3}\right)$ of curvature $-\frac{1}{3}$ into the unit pseudo-hyperbolic 4 -space $H_{2}^{4}(-1)$ discovered by the first author in [4]. It is known that the immersion $\psi_{\mathcal{B}}$ is a Wintgen ideal surfaces in $H_{2}^{4}(-1)$ whose Gauss and normal curvatures satisfy $K^{D}=2 K$. In the last section, we classify Wintgen ideal surfaces in $R_{2}^{4}(c)$ whose Gauss and normal curvatures satisfy the condition $K^{D}=2 K$. The later result provides us another simple geometric characterization of the minimal immersion $\psi_{\mathcal{B}}: H^{2}\left(-\frac{1}{3}\right) \rightarrow H_{2}^{4}(-1)$.

## 2. Preliminaries

A vector $v$ is called space-like (resp., time-like) if $\langle v, v\rangle>0$ (resp., $\langle v, v\rangle<0$ ). A surface $M$ in a pseudo-Riemannian manifold is called space-like if each nonzero tangent vector is space-like.

Let $R_{2}^{4}(c)$ denote an indefinite space form of constant curvature $c$ and with index 2. The curvature tensor $\tilde{R}$ of $R_{2}^{4}(c)$ is given by

$$
\begin{equation*}
\tilde{R}(X, Y) Z=c\{\langle Y, Z\rangle X-\langle X, Z\rangle Y\} \tag{2.1}
\end{equation*}
$$

for vectors $X, Y, Z$ tangent to $R_{2}^{4}(c)$. Let $\psi: M \rightarrow R_{2}^{4}(c)$ be an isometric immersion of a space-like surface $M$ into $R_{2}^{4}(c)$. Denote by $\nabla$ and $\tilde{\nabla}$ the Levi-Civita connections on $M$ and $R_{2}^{4}(c)$, respectively. For vector fields $X, Y$ tangent to $M$ and $\xi$ normal to $M$, the formulas of Gauss and Weingarten are given respectively by (cf. [1, 2, 10]):

$$
\begin{align*}
& \tilde{\nabla}_{X} Y=\nabla_{X} Y+h(X, Y)  \tag{2.2}\\
& \tilde{\nabla}_{X} \xi=-A_{\xi} X+D_{X} \xi \tag{2.3}
\end{align*}
$$

where $\nabla_{X} Y$ and $A_{\xi} X$ are the tangential components and $h(X, Y)$ and $D_{X} \xi$ the normal components of $\tilde{\nabla}_{X} Y$ and $\tilde{\nabla}_{X} \xi$, respectively.

The shape operator $A$ and the second fundamental form $h$ are related by

$$
\begin{equation*}
\langle h(X, Y), \xi\rangle=\left\langle A_{\xi} X, Y\right\rangle . \tag{2.4}
\end{equation*}
$$

The mean curvature vector $H$ of $M$ in $H_{2}^{4}(-1)$ is defined by $H=\frac{1}{2}$ trace $h$.

The equations of Gauss, Codazzi and Ricci are given respectively by

$$
\begin{align*}
R(X, Y) Z= & c\{\langle Y, Z\rangle X-\langle X, Z\rangle Y\}+A_{h(Y, Z)} X-A_{h(X, Z)} Y  \tag{2.5}\\
& \left(\bar{\nabla}_{X} h\right)(Y, Z)=\left(\bar{\nabla}_{Y} h\right)(X, Z)  \tag{2.6}\\
\langle & \left.R^{D}(X, Y) \xi, \eta\right\rangle=\left\langle\left[A_{\xi}, A_{\eta}\right] X, Y\right\rangle \tag{2.7}
\end{align*}
$$

for vector fields $X, Y, Z$ tangent to $M$, and $\xi, \eta$ normal to $M$, where $\bar{\nabla} h$ is defined by

$$
\left(\bar{\nabla}_{X} h\right)(Y, Z)=D_{X} h(Y, Z)-h\left(\nabla_{X} Y, Z\right)-h\left(Y, \nabla_{X} Z\right),
$$

and $R^{D}$ is the curvature tensor associated with the normal connection $D$, i.e.,

$$
\begin{equation*}
R^{D}(X, Y) \xi=D_{X} D_{Y} \xi-D_{Y} D_{X} \xi-D_{[X, Y]} \xi \tag{2.8}
\end{equation*}
$$

The normal curvature $K^{D}$ is given by

$$
\begin{equation*}
K^{D}=\left\langle R^{D}\left(e_{1}, e_{2}\right) e_{3}, e_{4}\right\rangle \tag{2.9}
\end{equation*}
$$

A surface $M$ in $R_{2}^{4}(c)$ is called parallel (resp., minimal) if $\bar{\nabla} h=0$ (resp., $H=0$ ) holds identically. An immersion $\psi: M \rightarrow R_{2}^{4}(c)$ is called full if the image $\psi(M)$ does not lies in any totally geodesic submanifold of $R_{2}^{4}(c)$. A surface $M$ in $R_{2}^{4}(c)$ is called isotropic if, at each point $p \in M,|h(u, u)|$ is independent of the choice of the unit vector $u \in T_{p} M$.

For an immersion $\psi: M \rightarrow H_{2}^{4}(-1)$, let $L=\iota \circ \psi: M \rightarrow \mathbb{E}_{3}^{5}$ be the composition of $\psi$ with the standard inclusion $\iota: H_{2}^{4}(-1) \rightarrow \mathbb{E}_{3}^{5}$ via (1.2). Since $H_{2}^{4}(-1)$ is totally umbilical with mean curvature one in $\mathbb{E}_{3}^{5}$, we have

$$
\begin{equation*}
\hat{\nabla}_{X} Y=\nabla_{X} Y+h(X, Y)+\langle X, Y\rangle L \tag{2.10}
\end{equation*}
$$

for $X, Y$ tangent to $M$, where $h$ is the second fundamental form of $\psi$ and $\hat{\nabla}$ denotes the Levi-Civita connection of $\mathbb{E}_{3}^{5}$.

## 3. Surfaces with Null Normal Curvature in $\mathbb{E}_{2}^{4}$

Theorem 3.1. Let $M$ be a space-like surface in the pseudo-Euclidean 4-space $\mathbb{E}_{2}^{4}$. If $M$ has constant mean and Gauss curvatures and null normal curvature, then $M$ is congruent to an open part of one of the following six types of surfaces:
(1) A totally geodesic plane in $\mathbb{E}_{2}^{4}$ defined by $(0,0, x, y)$;
(2) a totally umbilical hyperbolic plane $H^{2}\left(-\frac{1}{a^{2}}\right) \subset \mathbb{E}_{1}^{3} \subset \mathbb{E}_{2}^{4}$ given by

$$
(0, a \cosh u, a \sinh u \cos v, a \sinh u \sin v)
$$

where $a$ is a positive number;
(3) A flat surface in $\mathbb{E}_{2}^{4}$ defined by

$$
\frac{1}{\sqrt{2} m}(\cosh (\sqrt{2} m x), \cosh (\sqrt{2} m y), \sinh (\sqrt{2} m x), \sinh (\sqrt{2} m y))
$$

where $m$ is a positive number;
(4) A flat surface in $\mathbb{E}_{2}^{4}$ defined by

$$
\left(0, \frac{1}{a} \cosh (a x), \frac{1}{a} \sinh (a x), y\right),
$$

where $a$ is a positive number;
(5) A flat surface in $\mathbb{E}_{2}^{4}$ defined by

$$
\left(\frac{\cosh (\sqrt{2} x)}{\sqrt{2 m r}}, \frac{\cosh (\sqrt{2} y)}{\sqrt{2 m(2 m-r)}}, \frac{\sinh (\sqrt{2} x)}{\sqrt{2 m r}}, \frac{\sinh (\sqrt{2} y)}{\sqrt{2 m(2 m-r)}}\right),
$$

where $m$ and $r$ are positive numbers satisfying $2 m>r>0$;
(6) A surface of negative curvature $-b^{2}$ in $\mathbb{E}_{2}^{4}$ defined by

$$
\begin{aligned}
& \left(\frac{1}{b} \cosh (b x) \cosh (b y), \int_{0}^{y} \cosh (b y) \sinh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y\right. \\
& \left.\frac{1}{b} \sinh (b x) \cosh (b y), \int_{0}^{y} \cosh (b y) \cosh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y\right),
\end{aligned}
$$

where $b$ and $m$ are real numbers satisfying $0<b<m$.
Proof. Assume that $L: M \rightarrow \mathbb{E}_{2}^{4}$ is an isometric immersion of a space-like surface $M$ into $\mathbb{E}_{2}^{4}$. If $M$ is totally geodesic in $\mathbb{E}_{2}^{4}$, we obtain case (1). Thus, from now on, we assume that $M$ is non-totally geodesic in $\mathbb{E}_{2}^{4}$.

Let us choose an orthonormal tangent frame $\left\{e_{1}, e_{2}\right\}$ of the tangent bundle and an orthonormal normal frame $\left\{e_{3}, e_{4}\right\}$ of the normal bundle of $M$ which satisfy

$$
\begin{align*}
& \left\langle e_{1}, e_{1}\right\rangle=\left\langle e_{2}, e_{2}\right\rangle=1,\left\langle e_{1}, e_{2}\right\rangle=0  \tag{3.1}\\
& \left\langle e_{3}, e_{3}\right\rangle=\left\langle e_{4}, e_{4}\right\rangle=-1,\left\langle e_{3}, e_{4}\right\rangle=0 . \tag{3.2}
\end{align*}
$$

We may also choose $e_{1}, e_{2}$ which diagonalize $A_{e_{3}}$ so that the shape operator satisfies

$$
A_{e_{3}}=\left(\begin{array}{cc}
\alpha & 0  \tag{3.3}\\
0 & \mu
\end{array}\right), A_{e_{4}}=\left(\begin{array}{cc}
\delta & \gamma \\
\gamma & -\delta
\end{array}\right)
$$

for some functions $\alpha, \gamma, \delta, \mu$.
By definition, the normal curvature $K^{D}$ of $M$ is defined by

$$
\begin{equation*}
K^{D}=\left\langle\left[A_{e_{3}}, A_{e_{4}}\right] e_{1}, e_{2}\right\rangle \tag{3.4}
\end{equation*}
$$

For the orthonormal frame $\left\{e_{1}, e_{2}, e_{3}, e_{4}\right\}$, we put

$$
\begin{equation*}
\nabla_{X} e_{1}=\omega_{1}^{2}(X) e_{2}, \quad D_{X} e_{3}=\omega_{3}^{4}(X) e_{4} . \tag{3.5}
\end{equation*}
$$

From (2.3), (3.2) and (3.3) we have

$$
\begin{equation*}
h\left(e_{1}, e_{1}\right)=-\alpha e_{3}-\delta e_{4}, h\left(e_{1}, e_{2}\right)=-\gamma e_{4}, h\left(e_{2}, e_{2}\right)=-\mu e_{3}+\delta e_{4} . \tag{3.6}
\end{equation*}
$$

Thus, the mean curvature vector, the Gauss curvature and the normal curvature are given respectively by

$$
\begin{equation*}
H=-\frac{\alpha+\mu}{2} e_{3}, K=\gamma^{2}+\delta^{2}-\alpha \mu, K^{D}=\gamma(\mu-\alpha) . \tag{3.7}
\end{equation*}
$$

It follows from (3.5), (3.6) and the equation of Codazzi that

$$
\begin{align*}
& e_{1} \gamma-e_{2} \delta=\alpha \omega_{3}^{4}\left(e_{2}\right)-2 \gamma \omega_{1}^{2}\left(e_{2}\right)-2 \delta \omega_{1}^{2}\left(e_{1}\right),  \tag{3.8}\\
& e_{2} \alpha=-\gamma \omega_{3}^{4}\left(e_{1}\right)+\delta \omega_{3}^{4}\left(e_{2}\right)+(\alpha-\mu) \omega_{1}^{2}\left(e_{1}\right),  \tag{3.9}\\
& e_{2} \gamma+e_{1} \delta=\mu \omega_{3}^{4}\left(e_{1}\right)-2 \delta \omega_{1}^{2}\left(e_{2}\right)+2 \gamma \omega_{1}^{2}\left(e_{1}\right),  \tag{3.10}\\
& e_{1} \mu=-\delta \omega_{3}^{4}\left(e_{1}\right)-\gamma \omega_{3}^{4}\left(e_{2}\right)+(\alpha-\mu) \omega_{1}^{2}\left(e_{2}\right) . \tag{3.11}
\end{align*}
$$

Since $M$ has null normal curvature, we may also assume that $\gamma=0$. Thus, by the constancy of mean and Gauss curvatures, we obtain from (3.7) that

$$
\begin{equation*}
\mu=2 m-\alpha, \quad k=\delta^{2}+\alpha^{2}-2 m \alpha \tag{3.12}
\end{equation*}
$$

for some constants $k, m$. Without loss of generality, we may assume $m \geq 0$.
Case (i). $\mu=\alpha$. In this case, $\mu=\alpha=m$ is a constant, which gives $A_{e_{3}}=m I$. Moreover, (3.12) gives

$$
\begin{equation*}
\delta^{2}=m^{2}+k \geq 0 . \tag{3.13}
\end{equation*}
$$

Case (i.1). $m^{2}=-k$. From (3.13), we get $\delta=0$. Hence, $M$ is a totally umbilical surfaces in $\mathbb{E}_{2}^{4}$. Such a surface has parallel second fundamental form. Therefore, after applying Proposition 4.3 of [5], we obtain case (2) of the theorem,

Case (i.2). $m^{2}>-k$. Without loss of generality, we may put $\delta=\sqrt{m^{2}+k}$, which is a nonzero constant. Thus, we find from (3.8)-(3.11) that $\omega_{1}^{2}=\omega_{3}^{4}=0$. Hence, $M$ must be flat. So, we have $k=0$. Because $\omega_{1}^{2}=0$, we may choose coordinates $\{x, y\}$ such that $e_{1}=\partial / \partial x, e_{2}=\partial / \partial y$. The metric tensor is then given by $g=d x^{2}+d y^{2}$. Moreover, we know that the second fundamental form satisfies

$$
\begin{equation*}
h\left(e_{1}, e_{1}\right)=-m e_{3}-m e_{4}, h\left(e_{1}, e_{2}\right)=0, h\left(e_{2}, e_{2}\right)=-m e_{3}+m e_{4} . \tag{3.14}
\end{equation*}
$$

Now, it follows from (2.1), (3.14) that the immersion $L: M \rightarrow \mathbb{E}_{2}^{4}$ satisfies

$$
\begin{align*}
& L_{x x}=-m e_{3}-m e_{4}, L_{x y}=0, L_{y y}=-m e_{3}+m e_{4} \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}=-m L_{x}, \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}=-m L_{y}, \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-m L_{x}, \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4}=m L_{y} \tag{3.15}
\end{align*}
$$

After solving this system and choosing suitable initial conditions, we get case (3).
Case (ii). $\mu \neq \alpha$. It follows from (3.12) that

$$
\begin{equation*}
\mu=2 m-\alpha, \quad \delta=\sqrt{k+2 m \alpha-\alpha^{2}} \tag{3.16}
\end{equation*}
$$

Case (ii.1). $\quad \delta=0$. In this case, we have $k=\alpha^{2}-2 m \alpha$ which is constant. Hence, $\alpha$ is also a constant. Thus, we derive from (3.8)-(3.11) that

$$
\begin{equation*}
\omega_{1}^{2}=0, \quad \alpha \omega_{3}^{4}\left(e_{2}\right)=\mu \omega_{3}^{4}\left(e_{1}\right)=0 \tag{3.17}
\end{equation*}
$$

Therefore, $M$ is flat and $\alpha \mu=0$. Since $M$ is non-totally geodesic, without loss of generality we may assume that $\alpha \neq 0$ and $\mu=0$. Since $\omega_{1}^{2}=0$, we may choose coordinates $\{x, y\}$ such that $e_{1}=\partial / \partial x, e_{2}=\partial / \partial y$. So, we obtain

$$
\begin{equation*}
h\left(e_{1}, e_{1}\right)=-\alpha e_{3}, \quad h\left(e_{1}, e_{2}\right)=h\left(e_{2}, e_{2}\right)=0 \tag{3.18}
\end{equation*}
$$

It follows from (3.17) and (3.18) that immersion $L: M \rightarrow \mathbb{E}_{2}^{4}$ satisfies

$$
\begin{gather*}
L_{x x}=-\alpha e_{3}, \quad L_{x y}=L_{y y}=0 \\
\tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}=-\alpha L_{x}, \quad \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}=0 \tag{3.19}
\end{gather*}
$$

After solving this system and choosing suitable initial conditions, we get case (4).
Case (ii.2). $\delta \neq 0$. We have

$$
\begin{equation*}
\mu=2 m-\alpha, \quad \gamma=0, \quad \delta=\sqrt{k+2 m \alpha-\alpha^{2}} \neq 0 \tag{3.20}
\end{equation*}
$$

Case (ii.2.1). $m \alpha=-k$. In this case, $\alpha$ and $\delta$ are constant. Moreover, we have

$$
\begin{equation*}
\alpha=-\frac{k}{m}, \quad \delta=\frac{\sqrt{-k\left(k+m^{2}\right)}}{m}, \quad \mu=2 m+\frac{k}{m}, \quad \gamma=0 \tag{3.21}
\end{equation*}
$$

Because $\delta$ is a real nonzero number, we must have $-m^{2}<k<0$. Thus, we may put $k=-b^{2}$ with $0<b<m$. Substituting (3.21) into (3.8)-(3.11) yields

$$
\begin{equation*}
\omega_{1}^{2}\left(e_{2}\right)=\omega_{3}^{4}\left(e_{1}\right)=0, \quad \omega_{3}^{4}\left(e_{2}\right)=\frac{2 \sqrt{m^{2}-b^{2}}}{b} \omega_{1}^{2}\left(e_{1}\right) \tag{3.22}
\end{equation*}
$$

Thus, if $f$ is a function satisfying $e_{2}(\ln f)=\omega_{1}^{2}\left(e_{1}\right)$, then we get $\left[f e_{1}, e_{2}\right]=0$, which implies that there exist coordinates $\{x, y\}$ such that

$$
\begin{equation*}
\frac{\partial}{\partial x}=f e_{1}, \quad \frac{\partial}{\partial y}=e_{2} . \tag{3.23}
\end{equation*}
$$

Therefore, the metric tensor is given by

$$
\begin{equation*}
g=f^{2} d x^{2}+d y^{2} . \tag{3.24}
\end{equation*}
$$

Consequently, the Levi-Civita connection satisfies

$$
\begin{equation*}
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial x}=\frac{f_{x}}{f} \frac{\partial}{\partial x}-f f_{y} \frac{\partial}{\partial y}, \quad \nabla_{\frac{\partial}{\partial x} \frac{\partial}{\partial y}}=\frac{f_{y}}{f} \frac{\partial}{\partial x}, \quad \nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y}=0 . \tag{3.25}
\end{equation*}
$$

From (3.22), (3.23) and (3.25), we derive that

$$
\begin{equation*}
\omega_{1}^{2}\left(e_{1}\right)=-\frac{f_{y}}{f}, \omega_{1}^{2}\left(e_{2}\right)=\omega_{3}^{4}\left(e_{1}\right)=0, \omega_{3}^{4}\left(e_{2}\right)=-\frac{2 f_{y} \sqrt{m^{2}-b^{2}}}{b f} . \tag{3.26}
\end{equation*}
$$

Moreover, it follow from (3.24) and $K=-b^{2}$ that $f$ satisfies

$$
\begin{equation*}
f_{y y}=b^{2} f \tag{3.27}
\end{equation*}
$$

By solving (3.27) we obtain $f=u(x) \cosh (b y+v(x))$ for some functions $u(x), v(x)$. After replacing $x$ by an anti-derivative of $u(x)$, we find from (3.24) and (3.25) that

$$
\begin{gather*}
g=\cosh ^{2}(b y+v(x)) d x^{2}+d y^{2},  \tag{3.28}\\
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial x}=v^{\prime} \tanh (b y+v) \frac{\partial}{\partial x}-\frac{b}{2} \sinh (2 b y+2 v) \frac{\partial}{\partial y},  \tag{3.29}\\
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial y}=b \tanh (b y+v) \frac{\partial}{\partial x}, \quad \nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y}=0 .
\end{gather*}
$$

Also, it follows from (3.6), (3.21), and (3.28) that

$$
\begin{align*}
& h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right)=-\frac{\cosh ^{2}(b y+v)}{m}\left\{b^{2} e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}\right\}, \\
& h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right)=0,  \tag{3.30}\\
& h\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right)=\frac{\left(b^{2}-2 m^{2}\right) e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}}{m} .
\end{align*}
$$

Therefore, the immersion $L: M \rightarrow \mathbb{E}_{2}^{4}$ satisfies

$$
\begin{align*}
L_{x x}= & v^{\prime} \tan (b y+v) L_{x}-\frac{b}{2} \sinh (2 b y+2 v) L_{y} \\
& -\frac{\cosh ^{2}(b y+v)}{m}\left\{b^{2} e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}\right\}, \\
L_{x y}= & b \tanh (b y+v) L_{x}, \\
L_{y y}= & \frac{\left(b^{2}-2 m^{2}\right) e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}}{m},  \tag{3.31}\\
\tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}= & -\frac{b^{2}}{m} L_{x}, \quad \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-\frac{b \sqrt{m^{2}-b^{2}}}{m} L_{x}, \\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}= & \frac{b^{2}-2 m^{2}}{m} L_{y}-2 \sqrt{m^{2}-b^{2}} \tanh (b y+v) e_{4}, \\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4}= & \frac{b \sqrt{m^{2}-b^{2}}}{m} L_{y}+2 \sqrt{m^{2}-b^{2}} \tanh (b y+v) e_{3} .
\end{align*}
$$

The compatibility condition of (3.31) is given by $v^{\prime}(x)=0$. Then, after applying a suitable translation in $y$, we may put $v=0$. Therefore, system (3.31) reduces to

$$
\begin{align*}
L_{x x} & =-\frac{b}{2} \sinh (2 b y) L_{y}-\frac{\cosh ^{2}(b y)}{m}\left\{b^{2} e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}\right\}, \\
L_{x y} & =b \tanh (b y) L_{x}, \\
L_{y y} & =\frac{\left(b^{2}-2 m^{2}\right) e_{3}+b \sqrt{m^{2}-b^{2}} e_{4}}{m}, \\
\tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3} & =-\frac{b^{2}}{m} L_{x}, \quad \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-\frac{b \sqrt{m^{2}-b^{2}}}{m} L_{x},  \tag{3.32}\\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3} & =\frac{b^{2}-2 m^{2}}{m} L_{y}-2 \sqrt{m^{2}-b^{2}} \tanh (b y) e_{4}, \\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4} & =\frac{b \sqrt{m^{2}-b^{2}}}{m} L_{y}+2 \sqrt{m^{2}-b^{2}} \tanh (b y) e_{3} .
\end{align*}
$$

Solving the second equation in (3.32) gives

$$
\begin{equation*}
L=A(x) \cosh b y+B(y) \tag{3.33}
\end{equation*}
$$

for some vector-valued functions $A(x), B(y)$. Substituting this into the first, third and fourth equations in (3.32) gives $A^{\prime \prime \prime}(x)=b^{2} A^{\prime}(x)$. Thus, we get

$$
\begin{equation*}
A(x)=c_{5}+c_{1} \cosh (b x)+c_{2} \sinh (b x) \tag{3.34}
\end{equation*}
$$

for some vectors $c_{5}, c_{1}, c_{2}$. Combining this with (3.33) gives

$$
\begin{equation*}
L=\left(c_{5}+c_{1} \cosh (b x)+c_{2} \sinh (b x)\right) \cosh b y+B(y) . \tag{3.35}
\end{equation*}
$$

By substituting (3.35) into the first, third and fifth equations in (3.32), we find

$$
\begin{align*}
& \cosh ^{2}(b y) B^{\prime \prime \prime}-\frac{b}{2} \sinh (2 b y) B^{\prime \prime}+\left(3 b^{2}-4 m^{2}\right) B^{\prime}  \tag{3.36}\\
= & c_{5} b\left(3 b^{2}-4 m^{2}\right) \sinh (b y) .
\end{align*}
$$

A direct computation shows that $B_{p}=-c_{5} \cosh (b y)$ is a particular solution of (3.36). Thus, it follows from (3.35) and (3.36) that

$$
\begin{equation*}
L=\left(c_{1} \cosh (b x)+c_{2} \sinh (b x)\right) \cosh b y+C(y) \tag{3.37}
\end{equation*}
$$

where $C(y)$ satisfies the homogeneous differential equation:

$$
\begin{equation*}
\cosh ^{2}(b y) C^{\prime \prime \prime}(y)-\frac{b}{2} \sinh (2 b y) C^{\prime \prime}(y)+\left(3 b^{2}-4 m^{2}\right) C^{\prime}(y)=0 . \tag{3.38}
\end{equation*}
$$

After solving this differential equation, we have

$$
\begin{align*}
C(y)= & c_{3} \int_{0}^{y} \cosh (b y) \cosh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y  \tag{3.39}\\
& +c_{4} \int_{0}^{y} \cosh (b y) \sinh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y+c_{0}
\end{align*}
$$

for some vectors $c_{3}, c_{4}, c_{5} \in \mathbb{E}_{2}^{4}$. Combining this with (3.37) yields

$$
\begin{aligned}
L= & c_{0}+\left(c_{1} \cosh (b x)+c_{2} \sinh (b x)\right) \cosh b y \\
& +c_{3} \int_{0}^{y} \cosh (b y) \cosh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y \\
& +c_{4} \int_{0}^{y} \cosh (b y) \sinh \left(\frac{4 \sqrt{m^{2}-b^{2}}}{b} \tan ^{-1}\left(\tanh \frac{b y}{2}\right)\right) d y
\end{aligned}
$$

Therefore, after choosing suitable initial conditions, we obtain case (6).
Case (ii.2.2). $m \alpha \neq-k$. By substituting (3.20) into (3.8)-(3.11) we obtain

$$
\begin{align*}
& \omega_{3}^{4}\left(e_{1}\right)=\frac{2\left(k+m^{2}\right) \omega_{1}^{2}\left(e_{2}\right)}{m \sqrt{k+2 m \alpha-\alpha^{2}}}, \omega_{3}^{4}\left(e_{2}\right)=\frac{2\left(k+m^{2}\right) \omega_{1}^{2}\left(e_{1}\right)}{m \sqrt{k+2 m \alpha-\alpha^{2}}}  \tag{3.40}\\
& \omega_{1}^{2}\left(e_{1}\right)=e_{2}(\ln \sqrt{k+m \alpha}), \omega_{2}^{1}\left(e_{2}\right)=e_{1}\left(\ln \sqrt{k+2 m^{2}-m \alpha}\right) \tag{3.41}
\end{align*}
$$

It follows from (3.41) that $\left[e_{1} / \sqrt{k+m \alpha}, e_{2} / \sqrt{k+2 m^{2}-m \alpha}\right]=0$. Thus, there exist coordinates $\{x, y\}$ such that

$$
\begin{equation*}
\frac{\partial}{\partial x}=\frac{e_{1}}{\sqrt{k+m \alpha}}, \frac{\partial}{\partial y}=\frac{e_{2}}{\sqrt{k+2 m^{2}-m \alpha}} . \tag{3.42}
\end{equation*}
$$

Hence, the metric tensor is given by

$$
\begin{equation*}
g=\frac{d x^{2}}{k+m \alpha}+\frac{d y^{2}}{k+2 m^{2}-m \alpha} \tag{3.43}
\end{equation*}
$$

From (3.43) we have

$$
\begin{align*}
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial x} & =\frac{-m \alpha_{x}}{2(k+m \alpha)} \frac{\partial}{\partial x}+\frac{m\left(k+2 m^{2}-m \alpha\right) \alpha_{y}}{2(k+m \alpha)^{2}} \frac{\partial}{\partial y} \\
\nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y} & =\frac{-m \alpha_{y}}{2(k+m \alpha)} \frac{\partial}{\partial x}+\frac{m \alpha_{x}}{2\left(k+2 m^{2}-m \alpha\right)} \frac{\partial}{\partial y}  \tag{3.44}\\
\nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y} & =\frac{-m(k+m \alpha) \alpha_{x}}{2\left(k+2 m^{2}-m \alpha\right)^{2}} \frac{\partial}{\partial x}+\frac{m \alpha_{y}}{2\left(k+2 m^{2}-m \alpha\right)} \frac{\partial}{\partial y}
\end{align*}
$$

It follows from (3.6), (3.20) and (3.23) that the second fundamental form satisfies

$$
\begin{align*}
h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right) & =\frac{-\alpha e_{3}-\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+m \alpha}, h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right)=0  \tag{3.45}\\
h\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right) & =\frac{(\alpha-2 m) e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+2 m^{2}-m \alpha}
\end{align*}
$$

By applying (2.1), (3.24), (3.25) and (3.26) we obtain

$$
\begin{align*}
& L_{x x}= \frac{-m \alpha_{x} L_{x}}{2(k+m \alpha)}+\frac{m\left(k+2 m^{2}-m \alpha\right) \alpha_{y} L_{y}}{2(k+m \alpha)^{2}}-\frac{\alpha e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+m \alpha} \\
& L_{x y}= \frac{-m \alpha_{y} L_{x}}{2(k+m \alpha)}+\frac{m \alpha_{x} L_{y}}{2\left(k+2 m^{2}-m \alpha\right)}, \\
& L_{y y}= \frac{-m(k+m \alpha) \alpha_{x} L_{x}}{2\left(k+2 m^{2}-m \alpha\right)^{2}}+\frac{m \alpha_{y} L_{y}}{2\left(k+2 m^{2}-m \alpha\right)} \\
&+\frac{(\alpha-2 m) e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+2 m^{2}-m \alpha}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}=-\alpha L_{x}+\frac{\left(k+m^{2}\right) \alpha_{x}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{4}  \tag{3.46}\\
& \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}=(\alpha-2 m) L_{y}+\frac{\left(k+m^{2}\right) \alpha_{y}}{(k+m \alpha) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{4}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-\sqrt{k+2 m \alpha-\alpha^{2}} L_{x}-\frac{\left(k+m^{2}\right) \alpha_{x}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{3}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4}=\sqrt{k+2 m \alpha-\alpha^{2}} L_{y}-\frac{\left(k+m^{2}\right) \alpha_{y}}{(k+m \alpha) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{3} .
\end{align*}
$$

After applying (3.46) and a long computation, we find from $\left\langle L_{x x y}, L_{y}\right\rangle=$ $\left\langle L_{x y x}, L_{y}\right\rangle$ and from $\left\langle L_{x y y}, L_{x}\right\rangle=\left\langle L_{y y x}, L_{x}\right\rangle$ that

$$
\begin{equation*}
\alpha_{y}\left\{(k+m \alpha) \alpha_{x}+\left(k+2 m^{2}-m \alpha\right) \alpha_{y}\right\}=0 \tag{3.47}
\end{equation*}
$$

Hence, we have either
(1) $\alpha_{y}=0$ or
(2) $(k+m \alpha) \alpha_{x}+\left(k+2 m^{2}-m \alpha\right) \alpha_{y}=0$.

Case (ii.2.2.a). $\alpha_{y}=0$. In this case, system (3.46) reduces to

$$
\begin{align*}
& L_{x x}=\frac{-m \alpha_{x} L_{x}}{2(k+m \alpha)}-\frac{\alpha e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+m \alpha}, \\
& L_{x y}=\frac{m \alpha_{x} L_{y}}{2\left(k+2 m^{2}-m \alpha\right)}, \\
& L_{y y}=\frac{-m(k+m \alpha) \alpha_{x} L_{x}}{2\left(k+2 m^{2}-m \alpha\right)^{2}}+\frac{(\alpha-2 m) e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+2 m^{2}-m \alpha}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}=-\alpha L_{x}+\frac{\left(k+m^{2}\right) \alpha_{x}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{4},  \tag{3.48}\\
& \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}=(\alpha-2 m) L_{y}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-\sqrt{k+2 m \alpha-\alpha^{2}} L_{x}-\frac{\left(k+m^{2}\right) \alpha_{x}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} e_{3}, \\
& \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4}=\sqrt{k+2 m \alpha-\alpha^{2}} L_{y} .
\end{align*}
$$

Now, after applying (3.48), $\left\langle L_{x x y}, L_{y}\right\rangle=\left\langle L_{x y x}, L_{y}\right\rangle$ and $\left\langle L_{x y y}, L_{y}\right\rangle=\left\langle L_{y y x}, L_{y}\right\rangle$, we obtain that

$$
\begin{align*}
& a_{x x}=-\frac{2 k\left(k+2 m^{2}-m \alpha\right)^{2}-m^{2}\left(2 k+m^{2}+m \alpha\right) \alpha_{x}^{2}}{m\left(k+2 m^{2}-m \alpha\right)(k+m \alpha)}  \tag{3.49}\\
& \alpha_{x}^{2}=\frac{2 k\left(k+2 m^{2}-m \alpha\right)^{2}}{m^{2}\left(k+m^{2}\right)} \tag{3.50}
\end{align*}
$$

Next, by differentiating (3.50) and by applying (3.49), we find $\alpha_{x}=0$. Thus, $\alpha$ is a constant, say $\alpha=r$. Because $\delta \neq 0$, (3.50) gives $k=0$. Therefore, system (3.48) becomes

$$
\begin{aligned}
L_{x x} & =-\frac{r e_{3}+\sqrt{2 m r-r^{2}} e_{4}}{m r} \\
L_{x y} & =0 \\
L_{y y} & =\frac{(r-2 m) e_{3}+\sqrt{2 m r-r^{2}} e_{4}}{2 m^{2}-m r}
\end{aligned}
$$

$$
\begin{aligned}
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3}=-\alpha L_{x}, \quad \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3}=(\alpha-2 m) L_{y} \\
& \tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4}=-\sqrt{2 m r-r^{2}} L_{x}, \quad \tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4}=\sqrt{2 m r-r^{2}} L_{y}
\end{aligned}
$$

After solving this system and choosing suitable initial conditions, we have case (5) of the theorem.

Case (ii.2.2.b). $\quad(k+m \alpha) \alpha_{x}+\left(k+2 m^{2}-m \alpha\right) \alpha_{y}=0$. In this case, we have

$$
\begin{equation*}
\alpha_{y}=\frac{(k+m \alpha) \alpha_{x}}{m \alpha-k-2 m^{2}} \tag{3.51}
\end{equation*}
$$

Thus, system (3.46) becomes

$$
\begin{aligned}
L_{x x} & =-\frac{m \alpha_{x}\left(L_{x}+L_{y}\right)}{2(k+m \alpha)}-\frac{\alpha e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+m \alpha} \\
L_{x y} & =\frac{m \alpha_{x}\left(L_{x}+L_{y}\right)}{2\left(k+2 m^{2}-m \alpha\right)} \\
L_{y y} & =-\frac{m(k+m \alpha) \alpha_{x}\left(L_{x}+L_{y}\right)}{2\left(k+2 m^{2}-m \alpha\right)^{2}}+\frac{(\alpha-2 m) e_{3}+\sqrt{k+2 m \alpha-\alpha^{2}} e_{4}}{k+2 m^{2}-m \alpha}
\end{aligned}
$$

$$
\begin{align*}
\tilde{\nabla}_{\frac{\partial}{\partial x}} e_{3} & =-\alpha L_{x}+\frac{\left(k+m^{2}\right) \alpha_{x} e_{4}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}}  \tag{3.52}\\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{3} & =(\alpha-2 m) L_{y}-\frac{\left(k+m^{2}\right) \alpha_{x} e_{4}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} \\
\tilde{\nabla}_{\frac{\partial}{\partial x}} e_{4} & =-\sqrt{k+2 m \alpha-\alpha^{2}} L_{x}-\frac{\left(k+m^{2}\right) \alpha_{x} e_{3}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}} \\
\tilde{\nabla}_{\frac{\partial}{\partial y}} e_{4} & =\sqrt{k+2 m \alpha-\alpha^{2}} L_{y}+\frac{\left(k+m^{2}\right) \alpha_{x} e_{4}}{\left(k+2 m^{2}-m \alpha\right) \sqrt{k+2 m \alpha-\alpha^{2}}}
\end{align*}
$$

Now, from $L_{x x y}=L_{x y x}$ we find $(k+m \alpha) \alpha_{x}=0$. Also, we find from $L_{x y y}=L_{y y x}$ that $k=0$. Thus, $\alpha$ is a constant and $k=0$. Hence, this case reduces to (ii.2.a).

## 4. Spacelike Minimal Surfaces with Constant Gauss Curvature

From the equation of Gauss, we have
Lemma 4.1. Let $M$ be a space-like minimal surface in $R_{2}^{4}(c)$. Then $K \geq c$. In particular, if $K=c$ holds identically, then $M$ is totally geodesic.

For space-like minimal surfaces in $R_{2}^{4}(c)$, Theorem 1 of [12] implies that $M$ has constant Gauss curvature if and only if it has constant normal curvature.

We recall the following result of Sasaki from [12].

Theorem 4.2. Let $M$ be a space-like minimal surface in $R_{2}^{4}(c)$. If $M$ has constant Gauss curvature, then either
(1) $K=c$ and $M$ is a totally geodesic surface in $R_{2}^{4}(c)$;
(2) $c<0, K=0$ and $M$ is congruent to an open part of the minimal surface defined by $\frac{1}{\sqrt{2}}(\cosh u, \cosh v, 0, \sinh u, \sinh v)$, or
(3) $c<0, K=c / 3$ and $M$ is isotropic.

Let $\mathbf{R}^{2}$ be a plane with coordinates $s, t$. Consider a map $\mathcal{B}: \mathbf{R}^{2} \rightarrow \mathbb{E}_{3}^{5}$ given by

$$
\begin{align*}
& \mathcal{B}(s, t)=\left(\sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{7}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}-\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}\right. \\
& \left.\quad \frac{1}{2}+\frac{t^{2}}{2} e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}+\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}, \sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{1}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}\right) . \tag{4.1}
\end{align*}
$$

The first author proved in [4] that $\mathcal{B}$ defines a full isometric parallel immersion

$$
\begin{equation*}
\psi_{\mathcal{B}}: H^{2}\left(-\frac{1}{3}\right) \rightarrow H_{2}^{4}(-1) \tag{4.2}
\end{equation*}
$$

of the hyperbolic plane $H^{2}\left(-\frac{1}{3}\right)$ of curvature $-\frac{1}{3}$ into $H_{2}^{4}(-1)$.
The following result was also obtained in [4].
Theorem 4.3. Let $\psi: M \rightarrow H_{2}^{4}(-1)$ be a parallel full immersion of a spacelike surface $M$ into $H_{2}^{4}(-1)$. Then $M$ is minimal in $H_{2}^{4}(-1)$ if and only if $M$ is congruent to an open part of the surface defined by

$$
\begin{gathered}
\left(\sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{7}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}-\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}\right. \\
\left.\frac{1}{2}+\frac{t^{2}}{2} e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}+\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}, \sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{1}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}\right) .
\end{gathered}
$$

Combining Theorem 4.2 and Theorem 4.3, we obtain the following.
Theorem 4.4. Let $M$ be a non-totally geodesic space-like minimal surface in $H_{2}^{4}(-1)$. If $M$ has constant Gauss curvature $K$, then either
(1) $K=0$ and $M$ is congruent to an open part of the surface defined by

$$
\frac{1}{\sqrt{2}}(\cosh u, \cosh v, 0, \sinh u, \sinh v),
$$

or
(2) $K=-\frac{1}{3}$ and $M$ is is congruent to an open part of the surface defined by

$$
\begin{gathered}
\left(\sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{7}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}-\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}\right. \\
\left.\frac{1}{2}+\frac{t^{2}}{2} e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}+\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}, \sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{1}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}\right) .
\end{gathered}
$$

## 5. Wintgen Ideal Surfaces Satisfying $K^{D}=-2 K$

In 1979, P. Wintgen [13] proved a basic relationship between Gauss curvature $K$, normal curvature $K^{D}$, and mean curvature vector $H$ of a surface $M$ in a Euclidean 4-space $\mathbb{E}^{4}$; namely,

$$
\begin{equation*}
K+\left|K^{D}\right| \leq\langle H, H\rangle \tag{5.1}
\end{equation*}
$$

with the equality holding if and only if the curvature ellipse is a circle.
The following Wintgen type inequality for space-like surfaces in $R_{2}^{4}(c)$ can be found in [7].

Theorem 5.1. Let $M$ be a space-like surface in a 4-dimensional indefinite space form $R_{2}^{4}(c)$ of constant sectional curvature $c$ and index two. Then we have

$$
\begin{equation*}
K+K^{D} \geq\langle H, H\rangle+c \tag{5.2}
\end{equation*}
$$

at every point. Moreover, the equality sign of (5.2) holds at a point $p \in M$ if and only if, with respect to some suitable orthonormal frame $\left\{e_{1}, e_{2}, e_{3}, e_{4}\right\}$, the shape operator at $p$ satisfies

$$
A_{e_{3}}=\left(\begin{array}{cc}
\mu+2 \gamma & 0  \tag{5.3}\\
0 & \mu
\end{array}\right), A_{e_{4}}=\left(\begin{array}{ll}
0 & \gamma \\
\gamma & 0
\end{array}\right)
$$

Following [6, 9, 11], we call a surface in $R_{2}^{4}(c)$ Wintgen ideal if it satisfies the equality case of (5.2) identically. Wintgen ideal surfaces in $\mathbb{E}_{2}^{4}$ satisfying $|K|=$ $\left|K^{D}\right|$ are classified by the first author in [7] (see [6] for the classification of Wintgen ideal surfaces in $\mathbb{E}^{4}$ satisfying $|K|=\left|K^{D}\right|$.

We need the following existence result.
Theorem 5.2. Let $c$ be a real number and $\gamma$ with $3 \gamma^{2}>-c$ be a positive solution of the second order elliptic differential equation

$$
\begin{align*}
& \frac{\partial}{\partial x}\left(\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}} \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)}\right)  \tag{5.4}\\
& \quad-\frac{\partial}{\partial y}\left(\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right) \gamma_{y}}{2 \gamma\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}\right)=\gamma \sqrt{c+3 \gamma^{2}}
\end{align*}
$$

defined on a simply-connected domain $D \subset \mathbf{R}^{2}$. Then $M_{\gamma}=\left(D, g_{\gamma}\right)$ with the metric

$$
\begin{equation*}
g_{\gamma}=\frac{\sqrt{c+3 \gamma^{2}}}{\gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}\left(d x^{2}+\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{2 \sqrt{3}} d y^{2}\right) \tag{5.5}
\end{equation*}
$$

admits a non-minimal Wintgen ideal immersion $\psi_{\gamma}: M_{\gamma} \rightarrow R_{2}^{4}(c)$ into a complete simply-connected indefinite space form $R_{2}^{4}(c)$ satisfying $K^{D}=2 K$ identically.

Proof. Let $c$ be a real number and $\gamma$ be positive solution of (5.4) with $3 \gamma^{2}>-c$ defined on a simply-connected domain $D$. Consider the surface $M_{\gamma}=\left(D, g_{\gamma}\right)$ with metric $g_{\gamma}$ given by (5.5). Then the Levi-Civita connection of $g_{\gamma}$ satisfies

$$
\begin{align*}
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial x}= & -\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}+\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{y}}{2 \gamma\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{2 \sqrt{3}}} \frac{\partial}{\partial y} \\
\nabla_{\frac{\partial}{\partial x}} \frac{\partial}{\partial y}= & -\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}+\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial y} \\
\nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y}= & \frac{\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{2 \sqrt{3}}\left(c-3 \gamma \sqrt{c+3 \gamma^{2}}\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}  \tag{5.6}\\
& +\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right) \gamma_{y}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial y}
\end{align*}
$$

Let us define a bilinear map: $h: T M \rightarrow N M$ by

$$
\begin{align*}
h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right) & =-\frac{\left(\gamma+\sqrt{c+3 \gamma^{2}}\right) \sqrt{c+3 \gamma^{2}}}{\gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}} e_{3} \\
h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right) & =-\sqrt{c+3 \gamma^{2}} e_{4}  \tag{5.7}\\
h\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right) & =\frac{\left(\gamma-\sqrt{c+3 \gamma^{2}}\right) \sqrt{c+3 \gamma^{2}}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}{\gamma} e_{3}
\end{align*}
$$

where $N M$ is the plane bundle over $M$ spanned by an orthonormal time-like frame $\left\{e_{3}, e_{4}\right\}$. Define a linear metric connection $D$ on $N M$ by

$$
\begin{align*}
D_{\frac{\partial}{\partial x}} e_{3} & =\frac{-3 \gamma \gamma_{y} e_{4}}{\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}} \\
D_{\frac{\partial}{\partial y}} e_{3} & =\frac{3 \gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}} \gamma_{x}}{c+3 \gamma^{2}} e_{4} \\
D_{\frac{\partial}{\partial x}} e_{4} & =\frac{3 \gamma \gamma_{y} e_{3}}{\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}  \tag{5.8}\\
D_{\frac{\partial}{\partial y}} e_{3} & =-\frac{3 \gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}} \gamma_{x}}{c+3 \gamma^{2}} e_{3}
\end{align*}
$$

Then it follows from a very long direct computation that $\left(M_{\gamma}, g_{\gamma}, D, h\right)$ satisfies the equations of Gauss, Codazzi and Ricci. Hence, the fundamental existence and uniqueness theorem of submanifolds implies that, up to rigid motions, there exists a unique isometric immersion from $M_{\gamma}$ into $R_{2}^{4}(c)$ whose second fundamental form and normal connection are given by $h$ and $D$, respectively. By applying (5.5), (5.7) and $c+3 \gamma^{2}>0$ we see that $M$ is a non-minimal Wintgen ideal surface in $R_{2}^{4}(c)$.

Now, we classify Wintgen ideal surfaces in $R_{2}^{4}(c)$ which satisfy $K^{D}=2 K$.

Theorem 5.3. Let $M$ be a Wintgen ideal surface in a complete simply-connected indefinite space form $R_{2}^{4}(c)$ with $c=1,0$ or -1 . If $M$ satisfies $K^{D}=2 K$ identically, then one of following three cases occurs:
(1) $c=0$ and $M$ is a totally geodesic surface in $\mathbb{E}_{2}^{4}$;,
(2) $c=-1$ and $M$ is a minimal surface in $H_{2}^{4}(-1)$ congruent to an open part of $\psi_{\mathcal{B}}: H^{2}\left(-\frac{1}{3}\right) \rightarrow H_{2}^{4}(-1) \subset \mathbb{E}_{3}^{5}$ defined by

$$
\begin{gathered}
\quad\left(\sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{7}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}-\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}\right. \\
\left.\frac{1}{2}+\frac{t^{2}}{2} e^{\frac{2 s}{\sqrt{3}}}, t+\left(\frac{t^{3}}{3}+\frac{t}{4}\right) e^{\frac{2 s}{\sqrt{3}}}, \sinh \left(\frac{2 s}{\sqrt{3}}\right)-\frac{t^{2}}{3}-\left(\frac{1}{8}+\frac{t^{4}}{18}\right) e^{\frac{2 s}{\sqrt{3}}}\right)
\end{gathered}
$$

(3) $M$ is a non-minimal surface in $R_{2}^{4}(c)$ which is congruent to an open part of $\psi_{\gamma}: M_{\gamma} \rightarrow R_{2}^{4}(c)$ associated with a positive solution $\gamma$ of the elliptic differential equation (5.4) as described in Theorem 5.2.

Proof. Let $M$ be a Wintgen surface in $R_{2}^{4}(c)$. Then, according to Theorem 5.1, there exist an orthonormal frame $\left\{e_{1}, e_{2}, e_{3}, e_{4}\right\}$ such that shape operator satisfies (5.3) for some functions $\gamma, \mu$. Thus, the Gauss and normal curvatures are given by

$$
\begin{equation*}
K=c+\gamma^{2}-\mu^{2}-2 \gamma \mu, \quad K^{D}=-2 \gamma^{2} . \tag{5.9}
\end{equation*}
$$

It follows from the condition $K^{D}=2 K$ that $\mu=-\gamma \pm \sqrt{c+3 \gamma^{2}}$. Without loss of generality, we may assume $\gamma \geq 0$.

Case (i). $\mu=-\gamma+\sqrt{c+3 \gamma^{2}}$. We divide this into two subcases.
Case (i.1). $\quad c+3 \gamma^{2}=0$. We have $\mu=-\gamma$ and $c \leq 0$. Thus, $M$ is a minimal surface.

If $c=0$, we get $\gamma=\mu=0$, which implies that $M$ is totally geodesic. So, we get case (1) of the theorem.

If $c=-1$, we have $\gamma=-\mu=\frac{1}{\sqrt{3}}$. Thus, by (5.9) $M$ is a minimal surface with curvature $-\frac{1}{3}$. Hence, we obtain case (2) of the theorem according to Theorem 4.4.

Case (i.2). $\quad c+3 \gamma^{2} \neq 0$. From (5.3) we obtain

$$
\begin{align*}
& h\left(e_{1}, e_{1}\right)=-\left(\gamma+\sqrt{c+3 \gamma^{2}}\right) e_{3} \\
& h\left(e_{1}, e_{2}\right)=-\gamma e_{4}  \tag{5.10}\\
& h\left(e_{2}, e_{2}\right)=\left(\gamma-\sqrt{c+3 \gamma^{2}}\right) e_{3}
\end{align*}
$$

Thus, it follows from Codazzi's equation that

$$
\begin{equation*}
\omega_{1}^{2}\left(e_{1}\right)=\frac{3 \gamma \sqrt{c+3 \gamma^{2}}+c}{2 \gamma\left(c+3 \gamma^{2}\right)} e_{2} \gamma, \quad \omega_{1}^{2}\left(e_{2}\right)=\frac{3 \gamma \sqrt{c+3 \gamma^{2}}-c}{2 \gamma\left(c+3 \gamma^{2}\right)} e_{1} \gamma, \tag{5.11}
\end{equation*}
$$

$$
\begin{equation*}
\omega_{3}^{4}\left(e_{1}\right)=-\frac{3 \gamma e_{2} \gamma}{c+3 \gamma^{2}}, \quad \omega_{3}^{4}\left(e_{2}\right)=\frac{3 \gamma e_{1} \gamma}{c+3 \gamma^{2}} . \tag{5.12}
\end{equation*}
$$

After applying (5.11) we derive that

$$
\left[\frac{\left(c+3 \gamma^{2}\right)^{1 / 4}}{\sqrt{\gamma}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3} / 2}} e_{1}, \frac{\left(c+3 \gamma^{2}\right)^{1 / 4}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3} / 2}}{\sqrt{\gamma}} e_{2}\right]=0 .
$$

Hence there exist coordinates $\{x, y\}$ such that

$$
\begin{align*}
& \frac{\partial}{\partial x}=\frac{\left(c+3 \gamma^{2}\right)^{1 / 4}}{\sqrt{\gamma}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3} / 2}} e_{1}, \\
& \frac{\partial}{\partial y}=\frac{\left(c+3 \gamma^{2}\right)^{1 / 4}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3} / 2}}{\sqrt{\gamma}} e_{2} . \tag{5.13}
\end{align*}
$$

By using (5.13) we know that the metric tensor is given by

$$
\begin{equation*}
g=\frac{\sqrt{c+3 \gamma^{2}}}{\gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}} d x^{2}+\frac{\sqrt{c+3 \gamma^{2}}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}{\gamma} d y^{2} \tag{5.14}
\end{equation*}
$$

which implies that the Levi-Civita connection satisfies

$$
\begin{align*}
\nabla_{\frac{\partial}{}}^{\partial x} \frac{\partial}{\partial x}= & -\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}+\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{y}}{2 \gamma\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{2 \sqrt{3}}} \frac{\partial}{\partial y}, \\
\nabla_{\frac{\partial}{\partial x}}^{\partial y}= & -\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}+c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}+\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial y}, \\
\nabla_{\frac{\partial}{\partial y}} \frac{\partial}{\partial y}= & \frac{\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{2 \sqrt{3}}\left(c-3 \gamma \sqrt{c+3 \gamma^{2}}\right) \gamma_{x}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial x}  \tag{5.15}\\
& +\frac{\left(3 \gamma \sqrt{c+3 \gamma^{2}}-c\right) \gamma_{y}}{2 \gamma\left(c+3 \gamma^{2}\right)} \frac{\partial}{\partial y} .
\end{align*}
$$

From (5.12) and (5.13) we find

$$
\begin{align*}
& \omega_{3}^{4}\left(\frac{\partial}{\partial x}\right)=\frac{-3 \gamma \gamma_{y}}{\left(c+3 \gamma^{2}\right)\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}  \tag{5.16}\\
& \omega_{3}^{4}\left(\frac{\partial}{\partial y}\right)=\frac{3 \gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}{c+3 \gamma^{2}} \gamma_{x} .
\end{align*}
$$

Also, it follows from (5.10) and (5.13) that

$$
\begin{align*}
h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial x}\right) & =-\frac{\left(\gamma+\sqrt{c+3 \gamma^{2}}\right) \sqrt{c+3 \gamma^{2}}}{\gamma\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}} e_{3} \\
h\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}\right) & =-\sqrt{c+3 \gamma^{2}} e_{4}  \tag{5.17}\\
h\left(\frac{\partial}{\partial y}, \frac{\partial}{\partial y}\right) & =\frac{\left(\gamma-\sqrt{c+3 \gamma^{2}}\right) \sqrt{c+3 \gamma^{2}}\left(6 \gamma+2 \sqrt{3 c+9 \gamma^{2}}\right)^{\sqrt{3}}}{\gamma} e_{3}
\end{align*}
$$

Moreover, from (5.10), (5.15) and the equation of Gauss we know that $\gamma$ satisfies the elliptic differential equation (5.4). Consequently, after applying Theorem 5.2 we obtain case (3) of the theorem.

Case (ii). $\quad \mu=-\gamma-\sqrt{c+3 \gamma^{2}}$. After replacing $e_{3}, e_{4}$ by $-e_{3},-e_{4}$, respectively, this reduces to (i).

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[^0]:    Received September 4, 2009, accepted September 7, 2009.
    Communicated by J. C. Yao.
    2000 Mathematics Subject Classification: Primary 53C40; Secondary 53C50.
    Key words and phrases: Gauss curvature, Normal curvature, Wintgen ideal surface, Space-like surface.

