# Gauss Maps of Ruled Submanifolds and Applications II 

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#### Abstract

The notion of pointwise 1-type Gauss map was derived from the ordinary finite type Gauss map and it gives an interesting geometric properties on surfaces of 3-dimensional Euclidean space. In particular, the helicoid and the right cone of 3-dimensional Euclidean space are characterized by pointwise 1-type Gauss map. Inspired by such a study, in this paper, we completely classify ruled submanifolds of Euclidean space with pointwise 1-type Gauss map.


## 1. Introduction

One of main objects in differential geometry is to study Riemannian manifolds. Due to Nash's embedding theorem, every Riemannian manifolds can be regarded as a submanifold of Euclidean space with sufficiently high codimension. Thus, it is convenient to study Riemannian manifolds by examining submanifolds of Euclidean space with the intrinsic and extrinsic properties. The theory of finite type submanifolds initiated to estimate the total mean curvature of compact submanifolds of Euclidean space in the late 1970s. Inspired by the degree of algebraic varieties, B.-Y. Chen introduced the concept of order and type on submanifolds of Euclidean space:

Let $M$ be a submanifold of $m$-dimensional Euclidean space $\mathbb{E}^{m}$ with an isometric immersion $x$. We can identifying $x$ with the position vector of $\mathbb{E}^{m}$. Let $\Delta$ be the Laplace operator of $M$ in $\mathbb{E}^{m}$. The submanifold $M$ is said to be of finite type if $x$ has a spectral decomposition by $x=x_{0}+x_{1}+\cdots+x_{k}$, where $x_{0}$ is a constant vector and $x_{i}$ are

[^0]the vector fields satisfying $\Delta x_{i}=\lambda_{i} x_{i}$ for some $\lambda_{i} \in \mathbb{R}(i=1,2, \ldots, k)$. In particular, $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{k}$ are different, it is called $k$-type. Many works have been done in studying finite type submanifolds of Euclidean space (see $[2,3,5,10]$ ). This notion of finite type immersion was naturally extended to pseudo-Riemannian manifolds in pseudo-Euclidean space $13,17,19$ and it can be applied to smooth functions, particularly the Gauss map defined on submanifolds of Euclidean space or pseudo-Euclidean space [2,3].

In regarding the Gauss map of finite type, B.-Y. Chen and P. Piccini [6] studied the submanifolds in Euclidean space with finite type Gauss map so that they classified compact surfaces with 1-type Gauss map, that is, $\Delta G=\lambda(G+C)$, where $C$ is a constant vector and $\lambda \in \mathbb{R}$. Since then, quite a few works on ruled surfaces and ruled submanifolds with finite type Gauss map in Euclidean space or pseudo-Euclidean space have been done [12, 13, 15,20 .

However, some surfaces including a helicoid have an interesting property concerning the Gauss map which looks satisfying 1-type Gauss map. As a matter of fact, it is not: The helicoid in $\mathbb{E}^{3}$ parameterized by

$$
x(u, v)=(u \cos v, u \sin v, a v), \quad a \neq 0
$$

has the Gauss map and its Laplacian are respectively given by

$$
G=\frac{1}{\sqrt{a^{2}+u^{2}}}(a \sin v,-a \cos v, u)
$$

and

$$
\Delta G=\frac{2 a^{2}}{\left(a^{2}+u^{2}\right)^{2}} G
$$

The right (or circular) cone $C_{a}$ with parametrization

$$
x(u, v)=(u \cos v, u \sin v, a u), \quad a \geq 0
$$

has the Gauss map

$$
G=\frac{1}{\sqrt{1+a^{2}}}(a \cos v, a \sin v,-1)
$$

which satisfies

$$
\Delta G=\frac{1}{u^{2}}\left(G+\left(0,0, \frac{1}{\sqrt{1+a^{2}}}\right)\right)
$$

(cf. [7/8]). The Gauss maps of examples above are similar to 1-type, but obviously different from the usual sense of 1-type Gauss map. Based on these, we define

Definition 1.1. An oriented $n$-dimensional submanifold $M$ of the Euclidean space $\mathbb{E}^{m}$ is said to have pointwise 1-type Gauss map if it satisfies the condition

$$
\begin{equation*}
\Delta G=f(G+C) \tag{1.1}
\end{equation*}
$$

where $f$ is a non-zero smooth function on $M$ and $C$ some constant vector. In particular, if $C$ is zero, the Gauss map $G$ is said to be of the first kind. Otherwise, it is said to be of the second kind [4,7 9].

In [7, 8], M. Choi et al. proved that a ruled surface in 3-dimensional Euclidean space with pointwise 1-type Gauss map is part of a plane, a circular cylinder, a heilcoid, a cylinder over a plane curve of infinite type or a circular cone. And, in [9, 20], ruled surfaces in pseudo-Euclidean space with pointwise 1-type Gauss map were studied.

Continuing to [14], we now raise a question: Can we completely classify ruled submanifolds in Euclidean space with pointwise 1-type Gauss map of the second kind?

In this paper, we study the problem described in the question above and completely classify ruled submanifolds of Euclidean space with pointwise 1-type Gauss map.

All of geometric objects under consideration are smooth and submanifolds are assumed to be connected unless otherwise stated.

## 2. Preliminaries

Let $x: M \rightarrow \mathbb{E}^{m}$ be an isometric immersion of an $n$-dimensional Riemannian manifold $M$ into $\mathbb{E}^{m}$. Let $\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ be a local coordinate system of $M$. For the components $g_{i j}$ of the Riemannian metric $\langle\cdot, \cdot\rangle$ on $M$ induced from that of $\mathbb{E}^{m}$, we denote by ( $g^{i j}$ ) (respectively, $\mathcal{G}$ ) the inverse matrix (respectively, the determinant) of the component matrix $\left(g_{i j}\right)$. Then the Laplace operator $\Delta$ on $M$ is defined by

$$
\Delta=-\frac{1}{\sqrt{\mathcal{G}}} \sum_{i, j} \frac{\partial}{\partial x_{i}}\left(\sqrt{\mathcal{G}} g^{i j} \frac{\partial}{\partial x_{j}}\right) .
$$

We now choose an adapted local orthonormal frame $\left\{e_{1}, e_{2}, \ldots, e_{m}\right\}$ in $\mathbb{E}^{m}$ such that $e_{1}, e_{2}, \ldots, e_{n}$ are tangent to $M$ and $e_{n+1}, e_{n+2}, \ldots, e_{m}$ normal to $M$. The Gauss map $G: M \rightarrow G(n, m) \subset \mathbb{E}^{N}\left(N={ }_{m} C_{n}\right), G(p)=\left(e_{1} \wedge e_{2} \wedge \cdots \wedge e_{n}\right)(p)$ of $M$ is a smooth map which carries a point $p$ in $M$ to an oriented $n$-plane in $\mathbb{E}^{m}$ by the parallel translation of the tangent space of $M$ at $p$ to an $n$-plane passing through the origin in $\mathbb{E}^{m}$, where $G(n, m)$ is the Grassmannian manifold consisting of all oriented $n$-planes through the origin of $\mathbb{E}^{m}$.

An inner product $\langle\langle\cdot, \cdot\rangle\rangle$ on $G(n, m) \subset \mathbb{E}^{N}$ is defined by

$$
\left\langle\left\langle e_{i_{1}} \wedge \cdots \wedge e_{i_{n}}, e_{j_{1}} \wedge \cdots \wedge e_{j_{n}}\right\rangle\right\rangle=\operatorname{det}\left(\left\langle e_{i_{l}}, e_{j_{k}}\right\rangle\right)
$$

where $l, k$ run over the range $\{1,2, \ldots, n\}$. Then, $\left\{e_{i_{1}} \wedge e_{i_{2}} \wedge \cdots \wedge e_{i_{n}} \mid 1 \leq i_{1}<\cdots<i_{n}\right.$ $\leq m\}$ is an orthonormal basis of $\mathbb{E}^{N}$.

An $(r+1)$-dimensional submanifold $M$ in $\mathbb{E}^{m}$ is called a ruled submanifold if $M$ is foliated by $r$-dimensional totally geodesic submanifolds $E(s, r)$ of $\mathbb{E}^{m}$ along a regular
curve $\alpha=\alpha(s)$ on $M$ defined on an open interval $I$. Thus, a parametrization of a ruled submanifold $M$ in $\mathbb{E}^{m}$ can be given by

$$
x=x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s), \quad s \in I, t_{i} \in I_{i},
$$

where $I_{i}$ 's are some open intervals for $i=1,2, \ldots, r$. For each $s, E(s, r)$ is open in Span $\left\{e_{1}(s), e_{2}(s), \ldots, e_{r}(s)\right\}$, which is the linear span of linearly independent vector fields $e_{1}(s), e_{2}(s), \ldots, e_{r}(s)$ along the curve $\alpha$. We call $E(s, r)$ the rulings and $\alpha$ the base curve of the ruled submanifold $M$. In particular, the ruled submanifold $M$ is said to be cylindrical if $E(s, r)$ is parallel along $\alpha$, or non-cylindrical otherwise.

Definition 2.1. An $(r+1)$-dimensional cylindrical ruled submanifold $M$ is called a generalized circular cylinder $\Sigma_{a} \times \mathbb{E}^{r-1}$ if the base curve $\alpha$ is a circle and the generators of rulings are orthogonal to the plane containing the circle $\alpha$, where $\Sigma_{a}$ is a circular cylinder $S^{1}(a) \times \mathbb{R}$ in $\mathbb{E}^{3}$.

For later use, we need
Lemma 2.2. (1) Given a curve $\alpha$ and orthonormal vector fields $e_{1}, e_{2}, \ldots, e_{n}$ along $\alpha$ in a Riemannian manifold $\bar{M}$ with the Riemannian connection $\bar{D}$, we can always choose orthonormal vector fields $f_{1}, \ldots, f_{n}$ along $\alpha$ such that:
(i) The sets of vectors $\left\{f_{j}(s): 1 \leq j \leq n\right\}$ and $\left\{e_{j}(s): 1 \leq j \leq n\right\}$ generate the same subspace of $T_{\alpha(s)} \bar{M}$.
(ii) The vector fields $(\bar{D} / d s) f_{i}(s)$ are normal to the subspace of $T_{\alpha(s)} \bar{M}$ spanned by $\left\{f_{j}(s): 1 \leq j \leq n\right\}$ for all $i=1,2, \ldots, n$.

## 3. Pointwise 1-type Gauss map of the second kind

In this section, we consider a ruled submanifold $M$ in $\mathbb{E}^{m}$ with pointwise 1-type Gauss map of the second kind. Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{E}^{m}$. Then, the base curve $\alpha$ can be chosen to be orthogonal to the rulings by taking an integral curve of the field of normal directions to the rulings of $M$. Without loss of generality, we may assume that $\alpha$ is a unit speed curve, that is, $\left\langle\alpha^{\prime}(s), \alpha^{\prime}(s)\right\rangle=1$. From now on, the prime / denotes $d / d s$ unless otherwise stated. By Lemma 2.2 , we may choose orthonormal vector fields $e_{1}(s), \ldots, e_{r}(s)$ along $\alpha$ satisfying

$$
\begin{equation*}
\left\langle\alpha^{\prime}(s), e_{i}(s)\right\rangle=0, \quad\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0, \quad \text { for } s \in I \text { and } i, j=1,2, \ldots, r \tag{3.1}
\end{equation*}
$$

A parametrization of $M$ is then obtained as

$$
\begin{equation*}
x=x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s), \quad s \in I . \tag{3.2}
\end{equation*}
$$

In this paper, we always assume that the parametrization (3.2) satisfies the condition (3.1). Then, $M$ has the Gauss map

$$
G=\frac{1}{\left\|x_{s}\right\|} x_{s} \wedge x_{t_{1}} \wedge \cdots \wedge x_{t_{r}}
$$

or, equivalently

$$
\begin{equation*}
G=\frac{1}{q^{1 / 2}}\left(\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}\right) \tag{3.3}
\end{equation*}
$$

where $q$ is the function of $s, t_{1}, t_{2}, \ldots, t_{r}$ defined by $q=\left\langle x_{s}, x_{s}\right\rangle, \Phi$ and $\Psi_{i}(i=1,2, \ldots, r)$ are vector fields along $\alpha$ given by

$$
\Phi=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \quad \text { and } \quad \Psi_{i}=e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

Now, we separate the cases into two typical types of ruled submanifolds which are cylindrical or non-cylindrical.

First of all, we consider the following lemma.
Lemma 3.1. Suppose that a unit speed curve $\alpha(s)$ in an m-dimensional Euclidean space $\mathbb{E}^{m}$ defined on an interval I satisfies

$$
\begin{equation*}
\alpha^{\prime \prime \prime}(s)=f(s)\left(\alpha^{\prime}(s)+C\right) \tag{3.4}
\end{equation*}
$$

where $f$ is a function and $C$ a constant vector in $\mathbb{E}^{m}$. Then, the curve $\alpha$ lies in a 3dimensional Euclidean space. In particular, if the constant vector $C$ is zero, we see that $\alpha$ is a plane curve.

Proof. We fix a point $s_{0} \in I$. Let us denote by $V$ the linear span of $\left\{\alpha^{\prime}\left(s_{0}\right), \alpha^{\prime \prime}\left(s_{0}\right), C\right\}$. Then $V$ is of at most 3 -dimensional space in $\mathbb{E}^{m}$.

For any vector $a$ in the orthogonal complement $V^{\perp}$ of $V$, we consider the function $h_{a}(s)$ defined by $h_{a}(s)=\left\langle a, \alpha^{\prime}(s)\right\rangle$. Then, it follows from (3.4) that

$$
\begin{equation*}
h_{a}^{\prime \prime}(s)=f(s) h_{a}(s) \tag{3.5}
\end{equation*}
$$

Hence, the function $h_{a}(s)$ is a solution of a second order linear differential equation with initial condition $h_{a}\left(s_{0}\right)=h_{a}^{\prime}\left(s_{0}\right)=0$. This shows that the function $h_{a}(s)$ vanishes identically on the interval $I$. Thus, we have $\alpha^{\prime}(s) \in V$ for all $s \in I$, which shows that the curve $\alpha$ lies in a parallel displacement $\alpha\left(s_{0}\right)+V$ of the space $V$. This completes the proof.

Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{E}^{m}$ parameterized by (3.2) satisfying Condition (3.1).

Suppose that the cylindrical ruled submanifold $M$ has pointwise 1-type of the second kind. Equation $\Delta G=f(G+C)$ implies

$$
\begin{equation*}
-\Phi^{\prime \prime}=f(\Phi+C) \tag{3.6}
\end{equation*}
$$

for some non-zero function $f$ and some non-zero constant vector $C$. From equation (3.6), we note that $f$ is a function of $s$. We may assume that $f$ is non-zero on some open interval $I_{1}$. Then, on the interval $I_{1}$, differentiating equation (3.6) with respect to $s$ gives

$$
\begin{equation*}
\frac{f^{\prime}}{f^{2}} \Phi^{\prime \prime}-\frac{1}{f} \Phi^{\prime \prime \prime}-\Phi^{\prime}=\mathbf{0} \tag{3.7}
\end{equation*}
$$

or, equivalently

$$
\frac{f^{\prime}}{f^{2}} \alpha^{\prime \prime \prime}-\frac{1}{f} \alpha^{(4)}-\alpha^{\prime \prime}=\mathbf{0}
$$

which means that $-\frac{1}{f} \alpha^{\prime \prime \prime}-\alpha^{\prime}$ is a constant vector. Namely,

$$
\begin{equation*}
\Delta^{\prime} \alpha^{\prime}=-\alpha^{\prime \prime \prime}=f\left(\alpha^{\prime}+D\right) \tag{3.8}
\end{equation*}
$$

for some constant vector $D$ on $I_{1}$. On $I-I_{1},(3.8)$ holds obviously. Again, using Lemma 3.1, we see that the curve $\alpha$ lies in a 3 -dimensional Euclidean space $\mathbb{E}^{3}$.

Making use of Theorem 3.3 of [9] and Lemma 3.1, we have
Theorem 3.2. An $(r+1)$-dimensional cylindrical ruled submanifold $M$ in $\mathbb{E}^{m}$ has pointwise 1-type Gauss map of the second kind if and only if $M$ is an open part of an ( $r+1$ )-plane or a cylinder over a curve in $\mathbb{E}^{3}$ of infinite type.

We now consider the case of non-cylindrical ruled submanifold with pointwise 1-type Gauss map of the second kind. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold in $\mathbb{E}^{m}$ parameterized by (3.2).

Suppose that $M$ has pointwise 1-type Gauss map of the second kind, that is, the Gauss map $G$ of $M$ satisfies the condition

$$
\Delta G=f(G+C)
$$

for a non-zero function $f$ and some non-zero constant vector $C$. In order to make the matter simple, we introduce the following lemma.

Lemma 3.3. 14] Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold of $\mathbb{E}^{m}$ parametrized by (3.2) satisfying (3.1). Suppose some generators $e_{j_{1}}, e_{j_{2}}, \ldots, e_{j_{k}}$ $(1 \leq k<r)$ of the rulings are constant vectors along $\alpha$. Then, $M$ has pointwise 1-type Gauss map if and only if the ruled submanifold $M_{1}$ has pointwise 1-type Gauss map, where $M_{1}$ is the non-cylindrical ruled submanifold over the base curve $\alpha$ with the rulings generated by $e_{j}$ for $j \neq j_{1}, j_{2}, \ldots, j_{k}$.

According to Lemma 3.3, we may assume that $e_{i}^{\prime} \neq 0$ for all $i=1,2, \ldots, r$. Let us define functions:

$$
u_{i}=\left\langle\alpha^{\prime}, e_{i}^{\prime}\right\rangle, \quad w_{i j}=\left\langle e_{i}^{\prime}, e_{j}^{\prime}\right\rangle, \quad \phi=\left\langle\left\langle\Phi, \Phi^{\prime \prime}\right\rangle\right\rangle \quad \text { and } \quad \varphi_{i}=\left\langle\left\langle\Phi, \Psi^{\prime \prime}\right\rangle\right\rangle
$$

on $M$ for $i, j=1,2, \ldots, r$. Note that $u_{i} \neq 0$ for all $i=1,2, \ldots, r$. Then, $\Delta G=f(G+C)$ can be written as

$$
\begin{align*}
& \left(\frac{\partial q}{\partial s}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime} t_{j}\right)-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right) \\
& +q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime \prime} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} \Psi_{i}  \tag{3.9}\\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)+f\left(q^{3}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)+q^{\frac{7}{2}} C\right) \\
& =\mathbf{0} .
\end{align*}
$$

By putting

$$
\begin{aligned}
P(t)= & \left(\frac{\partial q}{\partial s}\right)^{2}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s} \sum_{j=1}^{r} x_{j} t_{j}-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right) \\
& +q^{2}\left(\phi+\sum_{j=1}^{r} \varphi_{j} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} u_{i} \\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)
\end{aligned}
$$

we may put

$$
\begin{equation*}
f=-\frac{P(t)}{q^{3}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)+q^{\frac{7}{2}} \gamma(s)}, \tag{3.11}
\end{equation*}
$$

where $\gamma(s)=\langle\langle C, \Phi(s)\rangle\rangle$. Substituting equation (3.11) into (3.9), we get

$$
\begin{align*}
& \text { 2) }\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)\left\{-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime} t_{j}\right)+q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime \prime} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} \Psi_{i}\right\}  \tag{3.12}\\
& -\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)\left\{-\frac{3}{2} q \frac{\partial q}{\partial s} \sum_{j=1}^{r} x_{j} t_{j}+q^{2}\left(\phi+\sum_{j=1}^{r} \varphi_{j} t_{j}\right)-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} u_{i}\right\} \\
& =-q^{\frac{1}{2}} \gamma(s)\left\{\left(\frac{\partial q}{\partial s}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime} t_{j}\right)\right.
\end{align*}
$$

$$
\begin{aligned}
& -\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)+q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime \prime} t_{j}\right)+\frac{1}{2} q \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right) \\
& \left.-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} \Psi_{i}-\frac{1}{2} q^{2} \sum_{i=1}^{r} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)\right\}+q^{\frac{1}{2}} C P(t) .
\end{aligned}
$$

Regarding equation (3.12), we have two possible cases whether $q^{\frac{1}{2}}$ is a polynomial or not.
Suppose that $q^{\frac{1}{2}}$ is not a polynomial in $t$. We will show that it is impossible. In this case, the polynomial in the left side must vanish and we can rewrite it as follows:

$$
\begin{align*}
& -\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)+\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) \sum_{k=1}^{r} x_{k} t_{k} \\
& +q^{2}\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)-q^{2}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\phi+\sum_{k=1}^{r} \varphi_{k} t_{k}\right)  \tag{3.13}\\
& -\frac{1}{2} q^{2} \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right) \Psi_{i}\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)+\frac{1}{2} q^{2} \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right) u_{i}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) \\
= & \mathbf{0} .
\end{align*}
$$

We put (3.13) in the following form

$$
\begin{equation*}
-\frac{3}{2}\left(\frac{\partial q}{\partial s}\right) R(t)=q Q(t) \tag{3.14}
\end{equation*}
$$

where

$$
\begin{equation*}
R(t)=\left(\Phi^{\prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime}\right)\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)-\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) \sum_{k=1}^{r} x_{k} t_{k} \tag{3.15}
\end{equation*}
$$

and

$$
\begin{aligned}
Q(t)= & -\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} t_{j} \Psi_{j}^{\prime \prime}\right)\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)+\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right)\left(\phi+\sum_{k=1}^{r} \varphi_{k} t_{k}\right) \\
& +\frac{1}{2} \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right) \Psi_{i}\left(1+\sum_{k=1}^{r} u_{k} t_{k}\right)-\frac{1}{2} \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right) u_{i}\left(\Phi+\sum_{j=1}^{r} t_{j} \Psi_{j}\right) .
\end{aligned}
$$

It yields

$$
R(t)=B(s) q(t)
$$

for some vector field $B(s)$ along $\alpha$ since the degree of the polynomial of $\sqrt{3.15}$ is 2 and $\operatorname{deg} q=2$. We can show that $u_{i}$ and $w_{i j}$ are constant along $\alpha$ for all $i, j=1,2, \ldots, r$. It means $\frac{\partial q}{\partial s}=0$. Thus, we have

$$
\begin{equation*}
\frac{\partial q}{\partial s}=0 \tag{3.16}
\end{equation*}
$$

For the details, see [14]. Then, the right side of equation (3.12) implies

$$
\begin{align*}
& \frac{1}{2} \sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left\{-\gamma(s)\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)+C\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)\right\} \\
= & q\left\{\gamma(s)\left(\Phi^{\prime \prime}+\sum_{j=1}^{r} \Psi_{j}^{\prime \prime} t_{j}-\frac{1}{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} \Psi_{i}-\frac{1}{2} \sum_{i=1}^{r} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(\Phi+\sum_{j=1}^{r} \Psi_{j} t_{j}\right)\right)\right.  \tag{3.17}\\
& \left.-C\left(\phi+\sum_{j=1}^{r} \varphi_{j} t_{j}-\frac{1}{2} \sum_{i=1}^{r} \frac{\partial q}{\partial t_{i}} u_{i}-\frac{1}{2} \sum_{i=1}^{r} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)\right)\right\} .
\end{align*}
$$

By (3.17), we have

$$
\begin{equation*}
\sum_{i=1}^{r}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}=a q(t) \tag{3.18}
\end{equation*}
$$

for a constant $a$ because of (3.16).
We now put

$$
\begin{equation*}
\left(\frac{\partial q}{\partial t_{j}}\right)^{2}=a_{j} q(t)+r_{j}(t) \tag{3.19}
\end{equation*}
$$

for some non-zero constants $a_{j}$ and polynomials $r_{j}(t)=\sum_{k=1}^{r} c_{j k} t_{k}+b_{j}$ in $t$ for $j=$ $1,2, \ldots, r$. According to (3.18), we must have

$$
\begin{equation*}
\sum_{j=1}^{r} r_{j}(t)=0 \tag{3.20}
\end{equation*}
$$

On the other hand, (3.19) with $w_{j j} \neq 0$ gives

$$
\begin{equation*}
4 w_{j j}=a_{j} \tag{3.21}
\end{equation*}
$$

(3.21) implies

$$
r_{j}(t)=\sum_{k=1}^{r} c_{j k} t_{k}+b_{j}=\sum_{k=1}^{r} 8\left(u_{j} w_{j k}-w_{j j} u_{k}\right) t_{k}+4\left(u_{j}^{2}-w_{j j}\right)
$$

for $j=1,2, \ldots, r$. By the definitions of $u_{i}$ and $w_{i i}$, we can see that $b_{j}=4\left(u_{j}^{2}-w_{j j}\right)$ are non-positive constants of the form

$$
b_{j}=-4 \sum_{a=r+1}^{m-1}\left\langle e_{j}^{\prime}, e_{a}\right\rangle^{2}
$$

for $j=1,2, \ldots, r$. By (3.20), we see that $\sum_{j=1}^{r} b_{j}=0$ and hence we have $e_{j}^{\prime}=u_{j} \alpha^{\prime}$, that is, $w_{j k}=u_{j} u_{k}$ for $j=1,2, \ldots, r$ and for $k=1,2, \ldots, r$. Using this, we can see that
$r_{j}(t)=0$, which means that $\left(\frac{\partial q}{\partial t_{j}}\right)^{2}$ are the multiples of $q(t)$ for $j=1,2, \ldots, r$. Thus we see that

$$
\left(\frac{\partial q}{\partial t_{i}}\right)^{2}=4 u_{i}^{2} q(t)
$$

for $i=1,2, \ldots, r$.
Comparing the both sides of the above equation, we have

$$
w_{i j}=u_{i} u_{j}
$$

for $i, j=1,2, \ldots, r$, which yields

$$
q=\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)^{2}
$$

which contradicts that $q^{1 / 2}$ is not a polynomial in $t$. Thus, we have

$$
\begin{equation*}
q=\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)^{2} \tag{3.22}
\end{equation*}
$$

Together with 3.22 and the definitions of $u_{i}$ and $w_{i i}$, we see that

$$
\begin{equation*}
e_{i}^{\prime}=u_{i} \alpha^{\prime} \tag{3.23}
\end{equation*}
$$

for $i=1,2, \ldots, r$. Hence, we have

$$
G=\Phi
$$

from which,

$$
\begin{equation*}
\phi(s)=\left\langle\left\langle\Phi(s), \Phi^{\prime \prime}(s)\right\rangle\right\rangle=-\left\langle\left\langle\Phi^{\prime}(s), \Phi^{\prime}(s)\right\rangle\right\rangle . \tag{3.24}
\end{equation*}
$$

Therefore, $\Delta G=f(G+C)$ yields

$$
\begin{equation*}
\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \Phi^{\prime}-\frac{1}{q} \Phi^{\prime \prime}=f(\Phi+C) \tag{3.25}
\end{equation*}
$$

Taking the inner product to (3.25) with $\Phi$, we obtain

$$
\begin{equation*}
-\frac{1}{q} \phi(s)=f(1+\gamma(s)) . \tag{3.26}
\end{equation*}
$$

On the other hand, the function $\phi(s)$ is reduced to

$$
\begin{equation*}
\phi(s)=-\sum_{a=r+1}^{m-1} u_{a}^{2}(s), \tag{3.27}
\end{equation*}
$$

where $u_{a}(s)=\left\langle\alpha^{\prime}(s), e_{a}^{\prime}(s)\right\rangle$ and $e_{r+1}, e_{r+2}, \ldots, e_{m-1}$ are unit normal vector fields to $M$ along $\alpha$. Note that $\phi(s)=0$ iff $\Phi^{\prime}(s)=\mathbf{0}$.

Suppose $\phi(s) \equiv 0$ on $I$. Then, $\Phi^{\prime}(s) \equiv \mathbf{0}$. This means the Gauss map $G$ is a constant vector field. Hence, $M$ is an open part of an $(r+1)$-plane $\mathbb{E}^{r+1}$.

We now suppose that the open subset $J=\{s \in I \mid \phi(s) \neq 0\}$ is not empty. We may put

$$
\begin{equation*}
f=-\frac{\phi(s)}{q(1+\gamma(s))} . \tag{3.28}
\end{equation*}
$$

Putting $q=\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)^{2}$ and substituting (3.28) into (3.25), we have

$$
\begin{equation*}
\left(\sum_{j=1}^{r} u_{j}^{\prime} t_{j}\right) \Phi^{\prime}-\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right) \Phi^{\prime \prime}=-\frac{\phi}{1+\gamma}\left(1+\sum_{j=1}^{r} u_{j} t_{j}\right)(\Phi+C) \tag{3.29}
\end{equation*}
$$

Considering the constant terms with respect to $t$ and coefficients of terms containing $t_{j}$ in (3.29), we get

$$
\begin{equation*}
\Phi^{\prime \prime}=\frac{\phi}{1+\gamma}(\Phi+C) \quad \text { and } \quad u_{j}^{\prime} \Phi^{\prime}=\mathbf{0} \tag{3.30}
\end{equation*}
$$

for all $j=1,2, \ldots, r$.
Since $\Phi^{\prime}(s) \neq \mathbf{0}, u_{j}^{\prime}(s)=0$ for all $j=1,2, \ldots, r$ and $s \in J$. Thus, $u_{j}$ and $w_{i j}$ are constant on $J$ for all $i, j=1,2, \ldots, r$. Hence, we have

Lemma 3.4. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (3.2) in $\mathbb{E}^{m}$. If $M$ has pointwise 1-type Gauss map of the second kind, then the functions

$$
u_{i}(s)=\left\langle\alpha^{\prime}(s), e_{i}^{\prime}(s)\right\rangle \quad \text { and } \quad w_{i j}(s)=\left\langle e_{i}^{\prime}(s), e_{j}^{\prime}(s)\right\rangle
$$

are constant functions on the open interval $J=\{s \in I \mid \phi(s) \neq 0\}$ for all $i, j=1,2, \ldots, r$.
The following lemma tells us how to choose some suitable normal vector fields along the base curve $\alpha$ for non-cylindrical ruled submanifold of $\mathbb{E}^{m}$ with pointwise 1-type Gauss map of the second kind.

Lemma 3.5. Let $M$ be an (r+1)-dimensional non-cylindrical ruled submanifold parametrized by (3.2) satisfying (3.1) in $\mathbb{E}^{m}$. If $M$ has pointwise 1-type Gauss map of the second kind, we can choose orthonormal frame $\left\{e_{a}\right\}_{a=r+1}^{m-1}$ of the normal space $T_{\alpha(s)} N$ of $M$ along $\alpha(s)$ satisfying

$$
\begin{equation*}
e_{a}^{\prime}(s) \wedge \alpha^{\prime}(s)=\mathbf{0} \tag{3.31}
\end{equation*}
$$

for all $a=r+1, \ldots, m-1$.

Proof. Let $\left\{\bar{e}_{a}\right\}_{a=r+1}^{m-1}$ denote an orthonormal frame of the normal space $T_{\alpha(s)} N$ of $M$ along $\alpha$. If we apply Lemma 2.2 to the normal space $T_{\alpha(s)} N$, then there exists an orthonormal frame $\left\{e_{a}\right\}_{a=r+1}^{m-1}$ of the normal space $T_{\alpha(s)} N$ satisfying

$$
\begin{equation*}
\left\langle e_{a}^{\prime}(s), e_{b}(s)\right\rangle=0 \tag{3.32}
\end{equation*}
$$

for all $a, b=r+1, \ldots, m-1$.
It follows from (3.23) that for all $i=1, \ldots, r$ and $a=r+1, \ldots, m-1$,

$$
\begin{equation*}
\left\langle e_{a}^{\prime}(s), e_{i}(s)\right\rangle=-\left\langle e_{i}^{\prime}(s), e_{a}(s)\right\rangle=0 \tag{3.33}
\end{equation*}
$$

Together with (3.32, (3.33) completes the proof.
We now give the following definition of a generalized right cone.
Definition 3.6. Suppose $\beta=\beta(s)$ is a circle on the unit sphere centered at the origin. Let $\boldsymbol{a}_{2}, \boldsymbol{a}_{3}, \ldots, \boldsymbol{a}_{r}$ be orthonormal constant vectors satisfying $\left\langle\beta^{\prime}(s), \boldsymbol{a}_{i}\right\rangle=\left\langle\beta(s), \boldsymbol{a}_{i}\right\rangle=0$ for all $i=2,3, \ldots, r$ and $s$. A ruled submanifold $M$ parametrized by

$$
\begin{equation*}
x\left(s, t_{1}, t_{2}, \ldots, t_{r}\right)=t_{1} \beta(s)+\sum_{i=2}^{r} t_{i} \boldsymbol{a}_{i}+D \tag{3.34}
\end{equation*}
$$

is called a generalized right cone $C_{a} \times \mathbb{E}^{r-1}$, where $C_{a}$ is a right cone in $\mathbb{E}^{3}, D$ a constant vector and $t_{i} \in I_{i}$ for some open intervals $I_{i}$ and $i=2,3, \ldots, r$.

Remark 3.7. If $\beta=\beta(s)$ is a great circle on the unit sphere in Definition 3.6, a generalized right cone is an $(r+1)$-plane $\mathbb{E}^{r+1}$.

Theorem 3.8. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold in $\mathbb{E}^{m}$. Then, $M$ has pointwise 1-type Gauss map of the second kind if and only if $M$ is an open portion of a generalized right cone.

Proof. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold with pointwise 1-type Gauss map of the second kind.

Suppose $\phi \equiv 0$ on the whole domain $I$ of $\alpha$. In this case, we see that $M$ is part of an $(r+1)$-plane $\mathbb{E}^{r+1}$. It is a special case of a generalized right cone.

Now, we suppose that $M$ is not part of $(r+1)$-plane, that is, $J=\{s \in I \mid \phi(s) \neq 0\}$ is non-empty. Note that

$$
\begin{equation*}
\left\langle\left\langle\Phi^{\prime}, \Phi^{\prime \prime}\right\rangle\right\rangle=-\frac{\phi^{\prime}}{2}, \quad\left\langle\left\langle\Phi^{\prime}, \Phi\right\rangle\right\rangle=0 \quad \text { and } \quad \phi=-\sum_{a=r+1}^{m-1} u_{a}^{2} . \tag{3.35}
\end{equation*}
$$

Then, the first equation in (3.30) implies

$$
-\frac{\phi^{\prime}}{2}=\frac{\phi}{1+\gamma}\left\langle\left\langle\Phi^{\prime}, C\right\rangle\right\rangle=\frac{\phi}{1+\gamma} \gamma^{\prime}
$$

or,

$$
\begin{equation*}
-\frac{1}{2} \frac{\phi^{\prime}}{\phi}=\frac{(1+\gamma)^{\prime}}{(1+\gamma)} \tag{3.36}
\end{equation*}
$$

on $J$. Equation (3.36) yields

$$
\frac{1}{\sqrt{|\phi|}}=\tilde{\lambda}|1+\gamma|
$$

for some positive constant $\tilde{\lambda}$. Therefore, we get

$$
\Phi^{\prime \prime}=\lambda \sqrt{|\phi|^{3}}(\Phi+C)
$$

for some non-zero constant $\lambda$.
On the other hand, according to Lemma 3.5, we have

$$
\begin{equation*}
\alpha^{\prime \prime}=-\sum_{i=1}^{r} u_{i} e_{i}-\sum_{a=r+1}^{m-1} u_{a} e_{a} \quad \text { and } \quad e_{a}^{\prime}=u_{a} \alpha^{\prime} \tag{3.37}
\end{equation*}
$$

for all $a=r+1, \ldots, m-1$. Since $\frac{1}{\lambda \sqrt{|\phi|^{3}}} \Phi^{\prime \prime}-\Phi$ is a constant vector, we have

$$
\begin{gather*}
\mathbf{0}=\sum_{a=r+1}^{m-1}\left\{\left(\frac{\mu}{\lambda \sqrt{|\phi|^{3}}}+1\right) u_{a}-\left(\frac{u_{a}^{\prime}}{\lambda \sqrt{|\phi|^{3}}}\right)^{\prime}\right\} e_{a} \wedge e_{1} \wedge \cdots \wedge e_{r} \\
+\sum_{k=1}^{r} \sum_{a=r+1}^{m-1} u_{k}\left\{\left(\frac{2}{\lambda \sqrt{|\phi|^{3}}}\right) u_{a}^{\prime}+\left(\frac{1}{\lambda \sqrt{|\phi|^{3}}}\right)^{\prime} u_{a}\right\}  \tag{3.38}\\
\quad \times \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{k-1} \wedge e_{a} \wedge e_{k+1} \wedge \cdots \wedge e_{r} .
\end{gather*}
$$

Considering the orthogonality of the vectors in (3.38), we get

$$
\begin{equation*}
\left(\frac{2}{\lambda \sqrt{|\phi|^{3}}}\right) u_{a}^{\prime}+\left(\frac{1}{\lambda \sqrt{|\phi|^{3}}}\right)^{\prime} u_{a}=0 \tag{3.39}
\end{equation*}
$$

for all $a=r+1, \ldots, m-1$.
Since $\phi \neq 0$ on $J$, there exists $u_{b} \neq 0$ for some $b=r+1, \ldots, m-1$. Then, equation (3.39) implies

$$
\frac{3}{4} \frac{|\phi|^{\prime}}{|\phi|}=\frac{u_{b}^{\prime}}{u_{b}}
$$

So we see that

$$
|\phi|^{\frac{3}{4}}=\varepsilon_{b} u_{b} \quad \text { or, } \quad u_{b}^{2}=\frac{1}{\varepsilon_{b}^{2}}|\phi|^{\frac{3}{2}}
$$

for some non-zero real number $\varepsilon_{b}$. Together with (3.35), we can see

$$
\phi=c|\phi|^{\frac{3}{2}}
$$

for some negative constant $c$, which means that the function $\phi$ is constant and hence the function $u_{b}$ is also constant. By continuity, the interval $J$ is the whole domain $I$ of $\alpha$.

It follows from (3.39) that $u_{a}$ are constant for all $a=r+1, \ldots, m-1$. Thus, (3.37) implies

$$
\alpha^{\prime \prime \prime}=-\mu \alpha^{\prime}
$$

for the constant $\mu=\sum u_{i}^{2}+\sum u_{a}^{2}$. According to Lemma 3.1, $\alpha$ is a space curve in $\mathbb{E}^{3}$. By considering the Frenet formula for $\alpha$, we see that the base curve $\alpha$ is a plane curve with non-zero constant curvature and thus it is part of a circle.

Since $e_{i}^{\prime}(s) \neq 0$ for all $i=1,2, \ldots, r$ as indicated in Section 3, we have

$$
\begin{equation*}
\alpha(s)=\frac{1}{u_{1}}\left(e_{1}(s)-\boldsymbol{a}_{1}\right) \quad \text { and } \quad e_{i}(s)=\frac{u_{i}}{u_{1}} e_{1}(s)+\boldsymbol{b}_{i} \tag{3.40}
\end{equation*}
$$

for some constant vectors $\boldsymbol{a}_{1}$ and $\boldsymbol{b}_{i}$ for $i=2,3, \ldots, r$ such that $e_{1}(s), \boldsymbol{b}_{2}, \boldsymbol{b}_{3}, \ldots, \boldsymbol{b}_{r}$ are linearly independent for each $s$. By applying the Gram-Schmidt's method for orthogonalization, we have orthonormal constant vectors $\boldsymbol{a}_{2}, \ldots, \boldsymbol{a}_{r}$ from $\boldsymbol{b}_{2}, \boldsymbol{b}_{3}, \ldots, \boldsymbol{b}_{r}$. In this case, $\left\langle e_{1}(s), \boldsymbol{b}_{i}\right\rangle$ are constant and thus $\left\langle e_{1}(s), \boldsymbol{a}_{i}\right\rangle$ are constant for all $i=2, \ldots, r$.

We put $v_{i}=\left\langle e_{1}, \boldsymbol{a}_{i}\right\rangle$ for all $i=2, \ldots, r$. Define $\beta_{1}(s)=e_{1}(s)-\sum_{i=2}^{r} v_{i} \boldsymbol{a}_{i}$. Then the length $\left\|\beta_{1}(s)\right\|=\sqrt{1-\sum_{i=2}^{r} v_{i}^{2}}$ of the vector field $\beta_{1}$ is a non-zero constant since $e_{1}, \boldsymbol{a}_{2}, \ldots, \boldsymbol{a}_{r}$ are linearly independent. Take $\beta(s)=\frac{\beta_{1}(s)}{\left\|\beta_{1}(s)\right\|}$. After appropriate change of parameters $t_{1}, t_{2}, \ldots, t_{r}$, the parametrization $(3.2)$ for $M$ can be reduced to

$$
\begin{equation*}
x\left(s, \bar{t}_{1}, \bar{t}_{2}, \ldots, \bar{t}_{r}\right)=\bar{t}_{1} \beta(s)+\sum_{i=2}^{r} \bar{t}_{i} \boldsymbol{a}_{i}+D \tag{3.41}
\end{equation*}
$$

for some constant vector $D$.
Since $\alpha$ is a circle, the first equation of (3.40) indicates that the trace of position vectors of $\beta(s)$ is a circle on the unit sphere. Thus, the parametrization given by (3.41) representing the ruled submanifold $M$ is an open part of a generalized right cone.

Conversely, suppose that $M$ is an open part of a generalized right cone $C_{a} \times \mathbb{E}^{r-1}$ parametrized by (3.34). In Introduction, we see that $C_{a}$ has pointwise 1-type Gauss map of the second kind. According to Lemma 3.3 , we see that $M$ has pointwise 1-type Gauss map of the second kind.

From above theorem, we immediately have
Corollary 3.9. 7 Let $M$ be a non-cylindrical ruled surface in $\mathbb{E}^{3}$. Then, $M$ has pointwise 1-type Gauss map of the second kind if and only if $M$ is an open portion of a right cone $C_{a}$.

Combining Theorem 3.7 of [14], Theorem 3.2 and Theorem 3.8, we have

Theorem 3.10 (Classification). An (r+1)-dimensional ruled submanifold $M$ of Euclidean space $\mathbb{E}^{m}$ with pointwise 1-type Gauss map is an open part of one of an $(r+1)$-plane, a generalized circular cylinder $\Sigma_{a} \times \mathbb{E}^{r-1}$, a cylinder over a curve in $\mathbb{E}^{3}$ of infinite type, a generalized helicoid or a generalized right cone $C_{a} \times \mathbb{E}^{r-1}$.

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