

## A NOTE ON SKEW PRODUCT PRESERVING MAPS ON FACTOR VON NEUMANN ALGEBRAS

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**ABSTRACT.** Let  $\mathcal{A}$  be a factor von Neumann algebra, with unit  $I$ , which contains a nontrivial projection  $P_1$ , and let  $\psi : \mathcal{A} \rightarrow \mathcal{A}$  be a surjective map that satisfies one of the two conditions:  $\psi(A)\psi(P) + \lambda\psi(P)\psi(A) = AP + \lambda PA$  and  $\psi(A)\psi(P) + \lambda\psi(P)\psi(A)^* = AP + \lambda PA^*$  for all  $A \in \mathcal{A}$  and  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ . Then, we determine the concrete form of  $\psi$ .

**1. Introduction.** Let  $\mathcal{R}$  be a  $*$ -ring. The Jordan product, Lie product,  $*$ -Jordan product and  $*$ -Lie product of  $A, B \in \mathcal{R}$  are defined as  $A \circ B = AB + BA$ ,  $[A, B] = AB - BA$ ,  $A \bullet B = AB + BA^*$  and  $[A, B]_* = AB - BA^*$ , respectively. These products play an important role in different fields of research. The additive map

$$\psi : \mathcal{R} \longrightarrow \mathcal{R},$$

defined by  $\psi(A) = AB - BA^*$  for all  $A, B \in \mathcal{R}$ , is a Jordan  $*$ -derivation, that is, it satisfies  $\psi(A^2) = \psi(A)A^* + A\psi(A)$ . The notion of Jordan  $*$ -derivations arose naturally in Šemrl's work [7, 8], where he investigated the problem of representing quadratic functionals with sesquilinear functionals. Let  $\mathcal{H}$  be a complex Hilbert space and  $\mathcal{B}(\mathcal{H})$  all of the bounded linear operators on  $\mathcal{H}$ . Motivated by the theory of rings (and algebras) equipped with a Lie product or a Jordan product, Molnár [5] studied the Lie product and gave a characterization of ideals of  $\mathcal{B}(\mathcal{H})$  in terms of the Lie product. It is shown [5] that, if  $\mathcal{N} \subseteq \mathcal{B}(\mathcal{H})$  is an ideal, then

$$\begin{aligned} \mathcal{N} &= \text{span}\{AB - BA^* : A \in \mathcal{N}, B \in \mathcal{B}(\mathcal{H})\} \\ &= \text{span}\{AB - BA^* : A \in \mathcal{B}(\mathcal{H}), B \in \mathcal{N}\}. \end{aligned}$$

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In particular, every operator in  $\mathcal{B}(\mathcal{H})$  is a finite sum of  $AB - BA^*$  type operators. Later, Beršar and Fsoñer [1] generalized the above results [5] to rings using different methods of involution. Let  $\mathcal{A}$  be a factor von Neumann algebra and

$$\phi : \mathcal{A} \longrightarrow \mathcal{A}$$

the  $*$ -Jordan derivation on  $\mathcal{A}$ . Then, in [11], we showed that  $\phi$  is an additive  $*$ -derivation.

Recall that a map

$$\psi : \mathcal{R} \longrightarrow \mathcal{R}$$

is skew commutativity preserving if, for any  $A, B \in \mathcal{R}$ ,  $[A, B]_* = 0$  implies  $[\psi(A), \psi(B)]_* = 0$ . The problem of characterizing linear (or additive) bijective maps preserving skew commutativity has been studied intensively in various algebras (see [2, 3] and the references therein). More specifically, we say that a map

$$\psi : \mathcal{R} \longrightarrow \mathcal{R}$$

is strong skew commutativity preserving if  $[\psi(A), \psi(B)]_* = [A, B]_*$  for all  $A, B \in \mathcal{R}$ . These maps are also called strong skew Lie product preserving maps in [4]. In [4], Cui and Park proved that, if  $\mathcal{A} \subseteq \mathcal{B}(\mathcal{H})$  is a factor von Neumann algebra, then every strong skew commutativity preserving map  $\psi$  on  $\mathcal{A}$  has the form

$$\psi(A) = \phi(A) + h(A)I \quad \text{for every } A \in \mathcal{A},$$

where  $\phi : \mathcal{A} \rightarrow \mathcal{A}$  is a linear bijective map satisfying  $[\phi(A), \phi(B)]_* = [A, B]_*$  for  $A, B \in \mathcal{A}$  and  $h$  is a real functional on  $\mathcal{A}$  with  $h(0) = 0$ . In particular, if  $\mathcal{A}$  is a type  $I$  factor, then  $\psi(A) = cA + h(A)I$  for every  $A \in \mathcal{A}$ , where  $c \in \{-1, 1\}$ . In addition, Qi and Hou [6] proved that, if  $\mathcal{M}$  is a von Neumann algebra with no central summands of type  $I_1$ , then a surjective map

$$\Phi : \mathcal{M} \longrightarrow \mathcal{M}$$

satisfies

$$\Phi(A)\Phi(B) - \Phi(B)\Phi(A)^* = AB - BA^*$$

for all  $A, B \in \mathcal{M}$  if and only if there exists a self-adjoint element  $Z$  in the center of  $\mathcal{M}$  with  $Z^2 = I$  such that  $\Phi(A) = ZA$  for all  $A \in \mathcal{M}$ .

In [9], we investigated the  $*$ -additivity of

$$\psi : \mathcal{A} \longrightarrow \mathcal{B},$$

where  $\mathcal{A}$  and  $\mathcal{B}$  are two prime  $C^*$ -algebras and  $\mathcal{A}$  contains a nontrivial projection  $P_1$ . We showed that, if  $\psi$  is a unital and bijective map and satisfies

$$\psi(AP + \lambda PA^*) = \psi(A)\psi(P) + \lambda\psi(P)\psi(A)^*$$

for all  $A \in \mathcal{A}$ ,  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ , then  $\psi$  is a  $*$ -additive map, where  $\mathcal{A}$  and  $\mathcal{B}$  are two  $C^*$ -algebras such that  $\mathcal{B}$  is prime. In [10], we investigated the additivity of map

$$\Phi : \mathcal{A} \longrightarrow \mathcal{B},$$

which is bijective, unital and satisfies

$$\Phi(AP + \eta PA^*) = \Phi(A)\Phi(P) + \eta\Phi(P)\Phi(A)^*$$

for all  $A \in \mathcal{A}$  and  $P \in \{P_1, I - P_1\}$ , where  $P_1$  is a nontrivial projection in  $\mathcal{A}$  and  $\eta$  is a non-zero complex number such that  $|\eta| \neq 1$ .

In this paper, we distinguish the concrete form of two types of strong skew-preserving maps on von Neumann algebras. Let  $\mathcal{A}$  be a factor von Neumann algebra (with identity  $I$ ) that contains a nontrivial projection  $P_1$ , and let  $\psi : \mathcal{A} \rightarrow \mathcal{A}$  be a map. First, if  $\psi$  is surjective and satisfies the condition

$$\psi(A)\psi(P) + \lambda\psi(P)\psi(A) = AP + \lambda PA$$

for all  $A \in \mathcal{A}$ ,  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ , then we will show that  $\psi(T) = \alpha T$  for  $\alpha \in \{-1, 1\}$  and for all  $T \in \mathcal{A}$ . Also, if  $\mathcal{A}$  is a von Neumann algebra and  $\psi : \mathcal{A} \rightarrow \mathcal{A}$  is not necessarily a surjective map satisfying the condition

$$\psi(A)\psi(P) + \lambda\psi(P)\psi(A)^* = AP + \lambda PA^*$$

for all  $A \in \mathcal{A}$ ,  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ , then we will show that there exists a  $Z \in \mathcal{A}$  with  $Z^2 = I$  such that  $\psi(A) = AZ$  for all  $A \in \mathcal{A}$ . Note that a subalgebra  $\mathcal{A}$  from  $\mathcal{B}(\mathcal{H})$  is called *von Neumann algebra* when it is closed in the weak topology of operators. A von Neumann algebra  $\mathcal{A}$  is called a *factor* when its center is trivial, i.e.,  $\mathcal{Z}(\mathcal{A}) = \mathbb{C}I$ . It is clear that, if  $\mathcal{A}$  is a factor von Neumann algebra, then  $\mathcal{A}$  is prime, that is, if  $A\mathcal{A}B = \{0\}$ , for  $A, B \in \mathcal{A}$ , then  $A = 0$  or  $B = 0$ .

**2. Statement of the main theorem.** The statement of our main theorems follow.

**Theorem 2.1.** *Let  $\mathcal{A}$  be a factor von Neumann algebra, with identity  $I$ , that contains a nontrivial projection  $P_1$ , and let  $\psi : \mathcal{A} \rightarrow \mathcal{A}$  be a surjective map which satisfies*

$$\psi(A)\psi(P) + \lambda\psi(P)\psi(A) = AP + \lambda PA$$

*for all  $A \in \mathcal{A}$ ,  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ . Then,  $\psi(T) = \alpha T$  for all  $T \in \mathcal{A}$ , where  $\alpha \in \{-1, 1\}$ .*

**Theorem 2.2.** *Let  $\mathcal{A}$  be a von Neumann algebra, with identity  $I$ , that contains a nontrivial projection  $P_1$ , and let*

$$\psi : \mathcal{A} \longrightarrow \mathcal{A}$$

*be a map which satisfies*

$$\psi(A)\psi(P) + \lambda\psi(P)\psi(A)^* = AP + \lambda PA^*$$

*for all  $A \in \mathcal{A}$ ,  $P \in \{P_1, I - P_1\}$  and  $\lambda \in \{-1, 1\}$ . Then, there exists a  $Z \in \mathcal{A}$  with  $Z^2 = I$  such that  $\psi(A) = AZ$  for all  $A \in \mathcal{A}$ .*

For the above-determined  $P_1$ , let  $P_2 = I - P_1$ . By taking  $\mathcal{A}_{ij} = P_i \mathcal{A} P_j$  for  $i, j = 1, 2$ , we can write

$$\mathcal{A} = \sum_{i,j=1,2} \mathcal{A}_{ij}.$$

We also note that each  $\mathcal{A}_{ij}$  is nonempty, and their pairwise intersections are the set of zero.

Note, in addition, that, by the assumptions

$$A \circ B = AB + BA \quad \text{and} \quad [A, B] = AB - BA,$$

for  $A, B \in \mathcal{A}$ , we can show the condition of  $\psi$  in Theorem 2.1 as follows:

$$(2.1) \quad \psi(A) \circ \psi(P) = A \circ P$$

and

$$(2.2) \quad [\psi(A), \psi(P)] = [A, P]$$

for all  $A \in \mathcal{A}$  and  $P \in \{P_1, P_2\}$ . Also, by the assumptions

$$A \bullet B = AB + BA^* \quad \text{and} \quad [A, B]_* = AB - BA^*,$$

for  $A, B \in \mathcal{A}$ , we show the condition of  $\psi$  in Theorem 2.2 as follows:

$$(2.3) \quad \psi(A) \bullet \psi(P) = A \bullet P$$

and

$$(2.4) \quad [\psi(A), \psi(P)]_* = [A, P]_*$$

for all  $A \in \mathcal{A}$  and  $P \in \{P_1, P_2\}$ .

We prove Theorem 2.1 in two steps.

*Step 1.* There exist  $\alpha_i, \beta_i \in \mathbb{C}$  with  $\alpha_i \neq 0$  such that  $\psi(P_i) = \alpha_i P_i + \beta_i I$  for  $i = 1, 2$ .

*Proof.* With simple computation, we can obtain

$$[P_1, [P_1, [A, P_1]]] = [A, P_1]$$

for all  $A \in \mathcal{A}$ . Thus, from equation (2.2), we have

$$[P_1, [P_1, [\psi(A), \psi(P_1)]]] = [\psi(A), \psi(P_1)].$$

Therefore,

$$[P_1, [P_1, [T, \psi(P_1)]]] = [T, \psi(P_1)]$$

for all  $T \in \mathcal{A}$ , as  $\psi$  is surjective.

Let  $K = [T, \psi(P_1)]$ . By simple calculation, from the above equation, we can obtain

$$(2.5) \quad P_1 K - 2P_1 K P_1 + K P_1 = K.$$

Multiplying by  $P_1$  from both sides of equation (2.5), it follows that  $P_1 K P_1 = 0$ . This yields

$$(2.6) \quad P_1 (T \psi(P_1) - \psi(P_1) T) P_1 = 0$$

for all  $T \in \mathcal{A}$ .

Let  $T = X_{11} \in \mathcal{A}_{11}$  in equation (2.6). We can write

$$X_{11} \psi(P_1) P_1 - P_1 \psi(P_1) X_{11} = 0,$$

and thus,

$$X_{11}P_1\psi(P_1)P_1 = P_1\psi(P_1)P_1X_{11}$$

for all  $X_{11} \in \mathcal{A}_{11}$ . Hence, there exists a  $\lambda_1 \in \mathbb{C}$  such that

$$(2.7) \quad P_1\psi(P_1)P_1 = \lambda_1 P_1$$

since  $\mathcal{A}$  is a factor. Replacing  $T$  by  $X_{12} \in \mathcal{A}_{12}$  in equation (2.6), we have

$$P_1X_{12}\psi(P_1)P_1 = 0,$$

and thus,

$$P_1XP_2\psi(P_1)P_1 = 0$$

for all  $X \in \mathcal{A}$ . The primeness of  $\mathcal{A}$  shows that

$$(2.8) \quad P_2\psi(P_1)P_1 = 0.$$

Similarly, by taking  $T = X_{21}$  in equation (2.6), we can obtain

$$(2.9) \quad P_1\psi(P_1)P_2 = 0.$$

Also, from  $P_1K - 2P_1KP_1 + KP_1 = K$ , we can obtain  $P_2KP_2 = 0$ . Therefore,

$$P_2(T\psi(P_1) - \psi(P_1)T)P_2 = 0.$$

Let  $T = X_{22} \in \mathcal{A}_{22}$  in the above equation. Similar to equation (2.7), we can write

$$(2.10) \quad P_2\psi(P_1)P_2 = \lambda_2 P_2$$

for some  $\lambda_2 \in \mathbb{C}$ .

On the other hand, from

$$\psi(P_1) = P_1\psi(P_1)P_1 + P_1\psi(P_1)P_2 + P_2\psi(P_1)P_1 + P_2\psi(P_1)P_2,$$

and from equations (2.7)–(2.10), it follows that

$$\psi(P_1) = \lambda_1 P_1 + \lambda_2 P_2,$$

which yields  $\alpha_1 = \lambda_1 - \lambda_2$  and  $\beta_1 = \lambda_2$ . The result  $\psi(P_1) = \alpha_1 P_1 + \beta_1 I$  is derived.

Now, we show that  $\alpha_1 \neq 0$ . On the contrary, suppose that  $\alpha_1 = 0$ . Then, for all  $B \in \mathcal{A}$ , we have

$$[\psi(B), \psi(P_1)] = [\psi(B), \beta_1 I] = 0.$$

Therefore,

$$[B, P_1] = o \implies BP_1 = P_1B.$$

Multiplying this latter equation on the left and right sides, respectively, by  $P_2$ , we obtain

$$B_{21} = B_{12} = 0$$

for all  $B \in \mathcal{A}$ , which is impossible. Thus,  $\alpha_1 \neq 0$ . Similarly, in this way,  $\psi(P_2) = \alpha_2 P_2 + \beta_2 I$  and  $\alpha_2 \neq 0$  can be obtained.  $\square$

*Step 2.*  $\psi(T) = \alpha T$  for all  $T \in \mathcal{A}$ , where  $\alpha^2 = 1$ .

*Proof.* From Step 1, for all  $T \in \mathcal{A}$ , we have

$$\begin{aligned} TP_1 - P_1T &= \psi(T)\psi(P_1) - \psi(P_1)\psi(T) \\ &= \psi(T)(\alpha_1 P_1 + \beta_1 I) - (\alpha_1 P_1 + \beta_1 I)\psi(T). \end{aligned}$$

Thus,

$$TP_1 - P_1T = \alpha_1 \psi(T)P_1 - \alpha_1 P_1 \psi(T).$$

Multiplying this equation on the left and right sides, respectively, by  $P_2$ , we have

$$\begin{aligned} P_2TP_1 &= \alpha_1 P_2 \psi(T)P_1 \\ P_1TP_2 &= \alpha_1 P_1 \psi(T)P_2. \end{aligned}$$

Therefore,

$$(2.11) \quad \psi(T)_{21} = P_2 \psi(T)P_1 = \alpha T_{21}$$

and

$$(2.12) \quad \psi(T)_{12} = P_1 \psi(T)P_2 = \alpha T_{12},$$

where  $\alpha = 1/\alpha_1$ .

On the other hand,

$$\begin{aligned} TP_1 + P_1T &= \psi(T)\psi(P_1) + \psi(P_1)\psi(T) \\ &= \psi(T)(\alpha_1 P_1 + \beta_1 I) + (\alpha_1 P_1 + \beta_1 I)\psi(T) \\ &= \alpha_1 \psi(T)P_1 + \alpha_1 P_1 \psi(T) + 2\beta_1 \psi(T). \end{aligned}$$

Therefore, from this equation and equations (2.11) and (2.12), we have

$$\begin{aligned} 2T_{11} + T_{21} + T_{12} &= \alpha_1\psi(T)_{11} + \alpha_1\psi(T)_{21} + \alpha_1\psi(T)_{11} \\ &\quad + \alpha_1\psi(T)_{12} + 2\beta_1\psi(T) \\ &= 2\alpha_1\psi(T)_{11} + T_{21} + T_{12} + 2\beta_1\psi(T). \end{aligned}$$

Hence,

$$\begin{aligned} T_{11} &= \alpha_1\psi(T)_{11} + \beta_1\psi(T) \\ &= \alpha_1\psi(T)_{11} + \beta_1(\psi(T)_{11} + \psi(T)_{12} \\ &\quad + \psi(T)_{21} + \psi(T)_{22}). \end{aligned}$$

If  $\beta_1 \neq 0$ , then, from the fact that the set of zero contains the pairwise intersections of  $\mathcal{A}_{ij}$ , we can obtain

$$\psi(T)_{12} = \psi(T)_{21} = \psi(T)_{22} = 0$$

for all  $T \in \mathcal{A}$ . This is a contraction from the surjectivity of  $\psi$ . Thus,  $\beta_1 = 0$ , and we have

$$(2.13) \quad P_1\psi(T)P_1 = \psi(T)_{11} = \alpha T_{11}.$$

Similarly, in this way, we can obtain

$$(2.14) \quad P_2\psi(T)P_2 = \delta T_{22}$$

and also

$$P_1\psi(T)P_2 = \delta T_{12},$$

where  $\delta = 1/\alpha_2$ . Hence, from the above equation and equation (2.12), we have  $\alpha = \delta$  and so  $\alpha_1 = \alpha_2$ . Since

$$\psi(T) = P_1\psi(T)P_1 + P_1\psi(T)P_2 + P_2\psi(T)P_1 + P_2\psi(T)P_2,$$

it follows from equations (2.11)–(2.14) that

$$\psi(T) = \alpha T$$

for all  $T \in \mathcal{A}$ . Thus,  $\psi(P_1) = \alpha P_1$ , and we also have  $\psi(P_1) = \alpha_1 P_1 = P_1/\alpha$ . Finally, this yields  $1/\alpha = \alpha$ , and thus,  $\alpha^2 = 1$ , which completes the proof of Theorem 2.1.  $\square$

Now, we will prove Theorem (2.2) by the following several steps.



*Step 1.* Under the assumptions of Theorem 2.2,  $\psi$  is additive on  $\mathcal{A}$ .

*Proof.* Letting  $A = P = P_1$  in equations (2.3) and (2.4), we have

$$\psi(P_1) \bullet \psi(P_1) = P_1 \bullet P_1$$

and

$$[\psi(P_1), \psi(P_1)]_* = [P_1, P_1]_*.$$

Thus,

$$\begin{aligned}\psi(P_1)^2 + \psi(P_1)\psi(P_1)^* &= 2P_1 \\ \psi(P_1)^2 - \psi(P_1)\psi(P_1)^* &= 0.\end{aligned}$$

Adding these equations, we have

$$(2.15) \quad \psi(P_1)^2 = P_1.$$

On the other hand, for all  $A, B \in \mathcal{A}$ , we have

$$\begin{aligned}(\psi(A+B) - \psi(A) - \psi(B)) \bullet \psi(P_1) \\ = \psi(A+B) \bullet \psi(P_1) - \psi(A) \bullet \psi(P_1) - \psi(B) \bullet \psi(P_1) \\ = (A+B) \bullet P_1 - A \bullet P_1 - B \bullet P_1 \\ = 0\end{aligned}$$

and

$$\begin{aligned}[\psi(A+B) - \psi(A) - \psi(B), \psi(P_1)]_* \\ = [\psi(A+B), \psi(P_1)]_* - [\psi(A), \psi(P_1)]_* - [\psi(B), \psi(P_1)]_* \\ = [A+B, P_1]_* - [A, P_1]_* - [B, P_1]_* \\ = 0.\end{aligned}$$

Therefore,

$$(\psi(A+B) - \psi(A) - \psi(B))\psi(P_1) + \psi(P_1)(\psi(A+B) - \psi(A) - \psi(B))^* = 0$$

and

$$(\psi(A+B) - \psi(A) - \psi(B))\psi(P_1) - \psi(P_1)(\psi(A+B) - \psi(A) - \psi(B))^* = 0.$$

Adding these equations, we have

$$(\psi(A+B) - \psi(A) - \psi(B))\psi(P_1) = 0.$$

Multiplying the above equation by  $\psi(P_1)$  from the right side and using equation (2.15), we have

$$(2.16) \quad (\psi(A+B) - \psi(A) - \psi(B))P_1 = 0.$$

Similarly, we can show that  $\psi(P_2)^2 = P_2$  and

$$(2.17) \quad (\psi(A+B) - \psi(A) - \psi(B))P_2 = 0.$$

Adding equations (2.16) and (2.17), we have

$$\psi(A+B) = \psi(A) + \psi(B). \quad \square$$

*Step 2.*  $\psi(I)^2 = \psi(I)\psi(I)^* = I$  and  $\psi(P_i) = \psi(I)P_i = P_i\psi(I)$  for  $i = 1, 2$ .

*Proof.* First, we show that equations (2.3) and (2.4) hold for  $P = I$ . Letting  $P = P_1$  and  $P = P_2$  in equation (2.3), respectively, we have

$$\psi(A) \bullet \psi(P_1) = A \bullet P_1$$

and

$$\psi(A) \bullet \psi(P_2) = A \bullet P_2$$

for all  $A \in \mathcal{A}$ . Adding these two equations, the equation

$$\psi(A) \bullet (\psi(P_1) + \psi(P_2)) = A \bullet (P_1 + P_2)$$

is inferred, and, from the additivity of  $\psi$ , we have

$$(2.18) \quad \psi(A) \bullet \psi(I) = A \bullet I.$$

In a similar way, we have

$$(2.19) \quad [\psi(A), \psi(I)]_* = [A, I]_*.$$

Let  $A = I$  in equations (2.18) and (2.19). With their aid, we can write

$$\psi(I)^2 + \psi(I)\psi(I)^* = 2I$$

$$\psi(I)^2 - \psi(I)\psi(I)^* = 0.$$

Hence,

$$\psi(I)^2 = \psi(I)\psi(I)^* = I.$$

Letting  $A = I$  and  $P = P_i$  for  $i = 1, 2$  in equations (2.3) and (2.4) we have

$$\psi(I)\psi(P_i) + \psi(P_i)\psi(I)^* = 2P_i$$

and

$$\psi(I)\psi(P_i) - \psi(P_i)\psi(I)^* = 0.$$

These equations yield

$$\psi(I)\psi(P_i) = P_i.$$

Multiplying this equation with  $\psi(I)$  from the left side, and from  $\psi(I)^2 = I$ , we have

$$\psi(P_i) = \psi(I)P_i.$$

Similarly to obtaining  $A = P_i$  for  $i = 1, 2$  in equations (2.18) and (2.19), we can obtain

$$\psi(P_i) = P_i\psi(I). \quad \square$$

*Step 3.* There exists a  $Z \in \mathcal{A}$  with  $Z^2 = I$  such that  $\psi(T) = TZ$  for all  $T \in \mathcal{A}$ .

*Proof.* From equation (2.3) and the fact that  $\psi(P_i) = \psi(I)P_i$  for  $i = 1, 2$ , we have

$$\begin{aligned} TP_1 + P_1T^* &= \psi(T)\psi(P_1) + \psi(P_1)\psi(T)^* \\ &= \psi(T)\psi(I)P_1 + \psi(I)P_1\psi(T)^* \end{aligned}$$

and

$$\begin{aligned} TP_2 + P_2T^* &= \psi(T)\psi(P_2) + \psi(P_2)\psi(T)^* \\ &= \psi(T)\psi(I)P_2 + \psi(I)P_2\psi(T)^* \end{aligned}$$

for every  $T \in \mathcal{A}$ . Adding these two equations, we have

$$T + T^* = \psi(T)\psi(I) + \psi(I)\psi(T)^*.$$

In addition, from equation (2.4), we can similarly obtain

$$T - T^* = \psi(T)\psi(I) - \psi(I)\psi(T)^*.$$

Adding these two latter equations, we can write

$$T = \psi(T)\psi(I).$$

Multiplying this equation with  $\psi(I)$  from the right side and the fact that  $\psi(I)^2 = I$ , we have

$$\psi(T) = T\psi(I).$$

Therefore, by obtaining  $Z = \psi(I)$ , we have  $Z^2 = I$  and  $\psi(T) = TZ$  for all  $T \in \mathcal{A}$ .

This completes the proof of Theorem 2.2. □

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