# AN EXTREMAL LENGTH CRITERION FOR THE PARABOLICITY OF RIEMANNIAN SPACES 

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#### Abstract

It is the purpose of this paper to show that a given Riemannian space satisfying a regularity condition is parabolic if and only if the extremal distance of a fixed ball in the space from the ideal boundary of the space is infinite.

We will also show that the harmonic modulus of a space bounded by two sets of boundary components coincides with the extremal distance between the two sets.


## Statements of Main Results

1. Regularity condition. Throughout this paper we denote by $R$ a noncompact $C^{\infty}$ Riemannian space with the ideal boundary $\beta$. We always assume that $R$ is orientable and connected. Let $A$ be the complement of a regular subregion of $R$ with the relative boundary $\alpha$. We also assume that $A-\alpha$ is connected. We consider the following regularity condition for $R$ (more precisely, for $A$ ):

For any nonconstant harmonic function $u$ defined on a region $\Omega \subset A$, the set $\{x \in \Omega||\nabla u(x)|=0\}$ has zero capacity.

This condition is always satisfied if the dimension of $R$ is two. This is also true, for example, when the metric tensor $g_{i j}$ is real analytic on $A-\alpha$. A typical case is furnished by a locally flat $A-\alpha$.

In this paper we only consider those spaces $R$ for which the above regularity condition is met.
2. Extremal length. Let $\rho$ be a density, i.e. a nonnegative Borel function on $A$, and let $\Gamma$ be a family of curves $\gamma$ which issue from a point in $\alpha$ and lie in $A-\alpha$. We define the harmonic extremal length, or simply the extremal length of $\Gamma$, by

$$
\begin{equation*}
\lambda(\Gamma)=\sup _{\rho \neq 0} \frac{L(\Gamma, \rho)^{2}}{V(A, \rho)}, \tag{1}
\end{equation*}
$$

where $V(A, \rho)=\int_{A} \rho^{2} d V$ and $L(\Gamma, \rho)=\inf _{\Gamma} \int_{\gamma} \rho d s$. Here $d V$ and $d s$ are the volume and the line element

We are particularly interested in the family $\Gamma_{\beta} \subset \Gamma$ of all curves $\gamma \in \Gamma$ terminating at $\beta$.
3. Parabolicity. We call $R$ parabolic, $R \in O_{G}$, if $R$ carries no nonconstant positive superharmonic function. The main object of this
paper is to prove:
Theorem 1. The space $R$ is parabolic if and only if $\lambda\left(\Gamma_{\beta}\right)=\infty$.
4. Moduli. Let $\Omega$ be a regular subregion of $R$ with relative boundary $\beta_{\Omega} \subset A-\alpha$, and let $u_{\Omega}$ be the continuous function on $\bar{\Omega} \cap A$ which is harmonic in the interior of $\bar{\Omega} \cap A$ with $u_{\Omega} \mid=0$ and $u_{\Omega} \mid \beta_{\Omega}=1$. The constant $\mu_{\rho}$ given by

$$
\begin{equation*}
\log \mu_{\Omega}=1 / \int_{\bar{\Omega}_{\cap A}} d u_{\Omega} \wedge * d u_{\Omega} \tag{2}
\end{equation*}
$$

is called the harmonic modulus, or simply the modulus of $\bar{\Omega} \cap A$ with respect to $\alpha$. It is easy to see that

$$
\begin{equation*}
\mu_{\Omega} \leqq \mu_{\Omega^{\prime}} \tag{3}
\end{equation*}
$$

for $\Omega \subset \Omega^{\prime}$. Therefore, we can define $\mu_{R}$, the harmonic modulus of $A$ with respect to $\alpha$, as the directed limit

$$
\begin{equation*}
\mu_{R}=\lim _{\Omega \rightarrow R} \mu_{\Omega} \tag{4}
\end{equation*}
$$

It is again easy to see that $u_{R}=\lim _{\Omega \rightarrow R} u_{\Omega}$ exists and is continuous on $A$, harmonic on $A-\alpha$ with $u_{R} \mid \alpha=0$. Moreover,

$$
\begin{equation*}
\log \mu_{R}=1 / \int_{A} d u_{R} \wedge * d u_{R} \tag{5}
\end{equation*}
$$

It can be seen that $R \in 0_{G}$ if and only if $\mu_{R}=\infty$ (Glasner [3]). Thus Theorem 1 may be considered as a special case of

Theorem 2. The following identity is valid:

$$
\begin{equation*}
\lambda\left(\Gamma_{\beta}\right)=\log \mu_{R} \tag{6}
\end{equation*}
$$

The proof will be given in 5-9.

## Parabolic Case

5. A general inequality. We start with proving

$$
\begin{equation*}
\lambda\left(\Gamma_{\beta}\right) \geqq \log \mu_{R} \tag{7}
\end{equation*}
$$

Let $\Gamma_{B_{g}}$ be the family of curves $\gamma \in \Gamma$ which lie in $\bar{\Omega} \cap A$ and terminate at a point of $\beta_{\Omega}$. Define $\rho$ as $\left(\log \mu_{\Omega}\right)\left|\nabla u_{\Omega}\right|$ in the interior of $\bar{\Omega} \cap A$ and as zero elsewhere in $R$. For $\gamma \in \Gamma_{\beta_{\Omega}}$,

$$
\int_{\gamma} \rho d s=\int_{\gamma}\left(\log \mu_{\Omega}\right)\left|\nabla u_{\Omega}\right| d s \geqq\left(\log \mu_{\Omega}\right) \int_{\gamma} \frac{d h}{d s} d s=\log \mu_{\Omega} .
$$

Therefore

$$
L\left(\Gamma_{\beta_{\Omega}}, \rho\right)=\inf _{\Gamma} \int_{\gamma} \rho d s \geqq \log \mu_{\Omega}
$$

By (2) we also obtain

$$
V(A, \rho)=\int_{\bar{\Omega} \cap A}\left(\log \mu_{\Omega}\right)^{2}\left|\nabla u_{\Omega}\right|^{2} d V=\left(\log \mu_{\Omega}\right)^{2} \int_{\bar{\Omega}} d_{\cap A} u_{\Omega} \wedge * d u_{\Omega}=\log \mu_{\Omega},
$$

and infer by (1) that

$$
\begin{equation*}
\lambda\left(\Gamma_{\beta_{\Omega}}\right) \geqq \log \mu_{\Omega} \tag{8}
\end{equation*}
$$

Since every $\gamma \in \Gamma_{\beta}$ contains a $\gamma^{\prime} \in \Gamma_{\beta_{\Omega}}$, we can easily see that $\lambda\left(\Gamma_{\beta}\right) \geqq \gamma\left(\Gamma_{\beta_{\Omega}}\right)$ (cf. Ahlfors-Sario [1, p. 222]). Thus (8) implies that $\lambda\left(\Gamma_{\beta}\right) \geqq \log \mu_{\Omega}$ for every $\Omega$. On letting $\Omega \rightarrow R$ we obtain (7).
6. Now suppose that $R \in 0_{G}$. Then since $\mu_{R}=\infty$, (7) implies that

$$
\begin{equation*}
\lambda\left(\Gamma_{\beta}\right)=\log \mu_{R}=\infty \tag{9}
\end{equation*}
$$

In order to complete the proofs of Theorems 1 and 2 , we have only to show the validity of (6) under the assumption $R \notin 0_{G}$. Note that in our discussion thus far we have not made any use of the regularity condition.

## Hyperbolic case

7. $u$-lines. Hereafter we assume that $R \notin 0_{G}$. Then $u_{R}$, to be denoted simply by $u$, is not constant on $A$. Since $u \mid \alpha=0$ and $u \mid A-\alpha>0$, we infer that $|\nabla u|$ can be extended continuously to all of $A$ and that $\mid \nabla u \| \alpha \neq 0$.

For each $x \in \alpha$ we consider the unique curve $l_{x}$ issuing from $x$ and such that $l_{x}-x \subset A-\alpha, * d u=0$ on $l_{x},|\nabla u| \neq 0$ on $l_{x}$. Moreover we require that $l_{x}$ either terminates at $\beta$ or at a point of $A$ at which $|\nabla u|=0$. Such an $l_{x}$ will be called a $u$-line. As $y$ traces $l_{x}, u(y)$ increases. Thus we can classify points of $\alpha$ as follows:

$$
\begin{aligned}
& \alpha_{0}=\left\{x \in \alpha \mid \lim _{y \rightarrow \beta, y \in l_{x}} u(y)<1\right\}, \\
& \alpha_{1}=\left\{x \in \alpha \mid \lim _{y \rightarrow \beta, y \in l_{x}} u(y)=1\right\},
\end{aligned}
$$

with

$$
\begin{equation*}
\alpha=\alpha_{0} \cup \alpha_{1} \tag{10}
\end{equation*}
$$

8. Vanishing surface area. We denote by $d S$ the surface element of $\alpha$. We wish to show that

$$
\begin{equation*}
S\left(\alpha_{0}\right)=\int_{\alpha_{0}} d S=0 \tag{11}
\end{equation*}
$$

Let $F_{-1}$ be the set of points $x \in \alpha$ such that $l_{x}$ terminates at some point of $R$. Clearly $F_{-1} \subset \alpha_{0}$, and we set $F_{0}=\alpha_{0}-F_{-1}$. By the regularity condition in $\S 1$, we see that $S\left(F_{-1}\right)=0$ (cf. Brelot-Choquet [2]). Therefore we only have to show that $S\left(F_{0}\right)=0$. Let

$$
F_{n}=\left\{x \in F_{0} \left\lvert\, \lim _{y \rightarrow \beta, y \in l_{x}}(1-u(y)) \geqq \frac{1}{n}\right.\right\} \quad(n=1,2, \cdots) .
$$

Since $F_{0}=\bigcup_{1}^{\infty} F_{n}$, it is sufficient to show that $S\left(F_{n}\right)=0$.
We can find a positive harmonic function $\omega$ in the interior of $A$ with the following properties (cf. Nakai [4]): (a) $\omega$ has the boundary values 0 on $\alpha$, (b) $\lim _{y \rightarrow \beta, y \in l_{x}} \omega(y)=\infty$ for $x \in F_{0}$, (c) $\int_{A}\left|\nabla \omega_{c}\right|^{2} d V \leqq c$, with $\omega_{c}=\min (\omega, c)$ for every positive number $c$.

Fix a $c>0$ arbitrarily and a point $y_{x} \in l_{x}$ with $w_{c}\left(y_{x}\right)=c$ for each $x \in F_{n}$.

Set $v=1-u$ on A. In a neighborhood of a point in $\alpha$ with respect to $A$ we may incorporate $v$ into a coordinate system, say $v=x^{1}$, while $x^{2}, \cdots, x^{m}$ are $m-1$ linearly independent parameters for $\alpha$. Then

$$
|\Delta v|^{2}=g^{11}\left(\frac{\partial v}{\partial x^{1}}\right)^{2}=g^{11}
$$

Since $* d v=|\nabla v| d S=\sqrt{g^{11}} d S$ on $\alpha, S\left(F_{n}\right)=0 \quad$ is equivalent to $\int_{F_{n}} * d v=0$. Observe that

$$
\begin{aligned}
c \int_{F_{n}} * d v & \leqq \int_{F_{n}}\left(\int_{x}^{y_{x}} \frac{\partial \omega_{c}}{\partial v} d v\right) * d v \\
& =\int_{F_{n}} \int_{x}^{y_{x}}\left|\frac{\partial \omega_{c}}{\partial v}\right| d v \wedge * d v \\
& =\int_{F_{n}} \int_{x}^{y_{x}}\left|\frac{\partial \omega_{c}}{\partial x^{1}}\right| g^{11} d V
\end{aligned}
$$

By the Schwarz inequality we have

$$
\int_{F_{n}} \int_{x}^{y_{x}}\left|\frac{\partial \omega_{c}}{\partial x^{1}}\right| g^{11} d V \leqq\left(\int_{F_{n}} \int_{x}^{y_{x}}\left|\frac{\partial \omega_{c}}{\partial x^{1}}\right|^{2} g^{11} d V\right)^{1 / 2}\left(\int_{F_{n}} \int_{x}^{y_{x}} g^{11} d V\right)^{1 / 2}
$$

$$
\begin{aligned}
& \leqq\left(\int_{A}\left|\nabla \omega_{c}\right|^{2} d V\right)^{1 / 2}\left(\int_{A}|\nabla v|^{2} d V\right)^{1 / 2} \\
& \leqq \sqrt{c}\left(\int_{A} d u \wedge * d u\right)^{1 / 2}
\end{aligned}
$$

From this we infer that

$$
\left|\int_{F_{n}} * d v\right| \leqq 1 / \sqrt{\mu_{R} c}
$$

Since the number $c$ can be arbitrarily large, we have $\int_{F_{n}} * d v=0$, and (11) follows.
9. Let $\rho$ be a density with $\rho \not \equiv 0$ on $A$. Since

$$
d u \wedge * d u=|\nabla u| d V
$$

we can compute

$$
\begin{aligned}
V(A, \rho) & =\int_{A} \rho^{2} d V=\int_{A} \frac{\rho^{2}}{|\nabla u|^{2}} d u \wedge * d u \\
& \geqq \int_{\alpha_{1}}\left(\int_{l_{x}} \frac{\rho^{2}}{|\nabla u|^{2}} d u\right) * d u \\
& =\int_{\alpha_{1}}\left(\int_{l_{x}} \frac{\rho^{2}}{|\nabla u|^{2}} d u \cdot \int_{l_{x}} 1^{2} d u\right) * d u \\
& \geqq \int_{\alpha_{1}}\left(\int_{l_{x}} \frac{\rho}{|\nabla u|} d u\right)^{2} * d u
\end{aligned}
$$

On $l_{x}\left(x \in \alpha_{1}\right)$ we have $d u=|\nabla u| d s$, and thus

$$
V(A, \rho) \geqq \int_{\alpha_{1}}\left(\int_{l_{x}} \rho d s\right)^{2} * d u
$$

From $l_{x} \in \Gamma_{\beta}$ for $x \in \alpha_{1}$ we obtain $\int_{l_{x}} \rho d s \geqq L\left(\Gamma_{\beta}, \rho\right)$, and therefore

$$
\begin{equation*}
V(A, \rho) \geqq L\left(\Gamma_{\beta}, \rho\right)^{2} \int_{\alpha_{1}} * d u \tag{12}
\end{equation*}
$$

On the other hand, by (11), we have $\int_{\alpha_{1}} * d u=\int_{\alpha} * d u$. Take an arbitrary regular region $\Omega$ with $\beta_{\Omega} \subset A-\alpha$. Then

$$
\int_{\alpha} * d u=\lim _{\Omega \rightarrow R} \int_{\alpha} * d u_{\Omega} .
$$

Here we see that

$$
\int_{\alpha} * d u_{\Omega}=\int_{\beta_{\Omega}} * d u_{\Omega}=\int_{\beta_{\Omega-\alpha}} u_{\Omega} * d u_{\Omega}=\int_{\bar{\Omega}_{\cap A}} d u_{\Omega} \wedge * d u_{\Omega},
$$

and infer that

$$
\int_{\Omega} * d u=\lim _{\Omega \rightarrow R} \int_{\bar{\Omega}_{\cap A}} d u_{\Omega} \wedge * d u_{\Omega}=\int_{A} d u \wedge * d u
$$

This together with (5) and (12) implies the inequality

$$
\log \mu_{R} \geqq \frac{L\left(\Gamma_{\beta}, \rho\right)^{2}}{V(A, \rho)}
$$

Since $\rho$ was arbitrary, we now conclude that

$$
\log \mu_{R} \geqq \lambda\left(\Gamma_{\beta}\right)
$$

We combine this with (7) and obtain (6).
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