PAPERS COMMUNICATED

99. Theory of Connections in the Generalized Finsler Manifold, III.

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I propose, in this paper, to extend the results of my investigation,¹⁾ so that a theory on the Finsler space by E. Cartan²⁾ can be deduced from mine by specialization.

1. In a generalized Finsler manifold, whose every element (x^{ν}, p^{ν}) is a point x^{ν} (in an *n*-dimensional manifold) associated to a system of the line elements $p^{\nu}, p^{\nu}, \ldots, p^{\nu}$, we define a connection from an element (x^{ν}, p^{ν}) to its consecutive element $(x^{\nu} + dx^{\nu}, p^{\nu} + dp^{\nu})$ for an ordinary vector v^{ν} as follows:

(1. 1)
$$\delta v^{\nu} = dv^{\nu} + v^{\lambda} (\Gamma^{\nu}_{\lambda\mu} dx^{\mu} + \sum_{i}^{j} \bigwedge_{\lambda\mu}^{i} dp^{i})$$

and for a pseudo-vector of the class k: r^{ν}

(1. 2)
$$\delta \mathbf{r}^{\nu} = d\mathbf{r}^{\nu} + \mathbf{r}^{\lambda} (\Gamma_{\lambda \mu}^{\nu} dx^{\mu} + \sum_{i} \bigwedge_{\lambda \mu}^{i} dx^{\mu}) - k \mathbf{r}^{\nu} \Gamma_{\mu} dx^{\mu} ,$$

where $\Gamma_{\lambda\mu}^{\nu}$, Γ_{μ} and $\bigwedge_{\lambda\mu}^{i}$ depend upon x^{ν} , p^{ν} ,, p^{ν} . The variations dp^{ν} are independent not only of one another but also of dx^{ν} . For a transformation of coordinate systems $\overline{x}^{\nu} = \overline{x}^{\nu}(x^{\lambda})$ the new parameters $\Gamma_{\lambda\mu}^{\nu}$ and Γ_{μ} of the connection have the forms

respectively and the other parameters $\dot{\wedge}_{\lambda\mu}^{\nu}$

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¹⁾ A. Kawaguchi, Theory of connections in the generalized Finsler manifold, Proc. 7 (1931), 211-214; and Die Differentialgeometrie in der verallgemeinerten Mannigfaltigkeit, Rendiconti del Circolo Matematico di Palermo, 56 (1932), 245-276.

²⁾ E. Cartan, Sur les espaces de Finsler, Comptes Rendus, Paris, 196 (1933), 582-586.

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where $P^{\nu}_{a} = \frac{\partial \overline{x}^{\nu}}{\partial x^{a}}, \qquad Q^{a}_{\nu} = \frac{\partial x^{a}}{\partial \overline{x}^{\nu}},$

and α is a function of position which appears in the formula of transformation of pseudo-vectors.¹⁾ (1. 3) shows that $\Gamma^{\nu}_{\lambda\mu}$ are not transformed as the parameters of an affine connection and (1. 5) that $\bigwedge^{i}_{\lambda\mu}$ are affinors of the class -i.

In accordance with the rules of transformation we can derive r affinors $\overset{i}{P}_{\lambda\mu}^{\nu}$ from the parameters $\Gamma_{\lambda\mu}^{\nu}$, Γ_{μ} , $\overset{i}{\wedge}_{\lambda\mu}^{\nu}$ and that of the base-connections:²⁾

(1. 6)
$$\dot{P}_{\lambda\mu}^{\nu} = \Gamma_{\lambda\mu}^{\nu} - i\delta_{\lambda}^{\nu}\Gamma_{\mu} - \dot{\Gamma}_{\lambda\mu}^{i} + i\delta_{\lambda}^{\nu}\dot{\Gamma}_{\mu}^{i} - \sum_{j}\dot{\bigwedge}_{\lambda\tau}^{j}(\dot{\Gamma}_{\pi\mu}^{\tau} - j\delta_{\pi}^{\tau}\dot{\Gamma}_{\mu}^{j})p^{\pi}$$

2. There are in this connection not only the covariant derivatives $\nabla_{\mu}v^{\nu}$, which are the coefficients of dx^{ν} in δv^{ν} , but also many others $\overset{i}{\nabla}_{\mu}v^{\nu}$, the coefficients of the covariant differential $\overset{i}{\delta}p^{\nu}$ with respect to the base-connections:

(2. 1)
$$\nabla_{\mu}v^{\nu} = \frac{\partial v^{\nu}}{\partial x^{\mu}} + I^{\nu}_{\lambda\mu}v^{\lambda} - \sum_{i} \overset{i}{\nabla}_{\lambda}v^{\nu} (I^{i}_{\pi\mu} - iI^{i}_{\mu}\delta^{\lambda}_{\pi})p^{i}_{\pi},$$

(2. 2)
$$\nabla^{i}_{\mu}v^{\nu} = \frac{\partial v^{\nu}}{\partial p^{\mu}} + \bigwedge^{i}_{\lambda\mu}v^{\lambda} .$$

We have namely

(2. 3)
$$\delta v^{\nu} = dx^{\mu} \nabla_{\mu} v^{\nu} + \sum_{i} \overset{i}{\delta} \overset{i}{p}^{\mu} \overset{i}{\nabla}_{\mu} v^{\nu}.$$

These covariant derivatives $\bigtriangledown_{\mu}v^{\nu}$ as well as $\stackrel{i}{\bigtriangledown}_{\mu}v^{\nu}$ are all affinors, being contravariant for the upper suffix and covariant for the lower. The affinor $\stackrel{i}{\bigtriangledown}_{\mu}v^{\nu}$ is of the class -i, when v^{ν} is an ordinary vector.

3. There exist many curvature tensors of our connection:

(3. 1)
$$R^{\nu}_{\lambda\mu\omega} = K^{\nu}_{\lambda\mu\omega} - \sum_{i} \overset{i}{R}^{\nu}_{\lambda\mu\rho} (\overset{i}{\Gamma}^{\rho}_{\tau\omega} - i\delta^{\rho}_{\tau} \overset{i}{\Gamma}^{j}_{\omega}) \overset{i}{p}^{\tau} + \sum_{i} \overset{i}{R}^{\nu}_{\lambda\nu\rho} (\overset{i}{\Gamma}^{\rho}_{\tau\mu} - i\delta^{\rho}_{\tau} \overset{i}{\Gamma}^{j}_{\mu}) \overset{i}{p}^{\tau} - \sum_{i,j} \overset{i}{R}^{i,j}_{\lambda\rho\tau} (\overset{i}{\Gamma}^{\rho}_{\sigma\mu} - i\delta^{\rho}_{\sigma} \overset{i}{\Gamma}^{j}_{\mu}) (\overset{j}{\Gamma}^{\tau}_{\beta\omega} - j\delta^{\tau}_{\beta} \overset{j}{\Gamma}^{j}_{\omega}) \overset{i}{p}^{\sigma} \overset{i}{p}^{\beta} ,$$

1) For a pseudo-vector $\mathbf{r}^{\mathbf{v}}$ of the class k

$$r^{\nu} = \alpha^k P^{\nu}_{\beta} r^{\beta}$$
.

2) A. Kawaguchi, loc. cit.

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(3. 2)
$$\hat{R}^{i}_{\lambda\mu\omega} = \hat{K}^{i}_{\lambda\mu\omega} - \sum_{j} \hat{R}^{j,i}_{\lambda\rho\omega} (\hat{I}^{j}_{\tau\mu} - j\delta^{\rho}_{\tau}\hat{I}^{j}_{\mu}) p^{\tau},$$

(3. 3)
$$\overset{i,j}{R}_{\lambda\mu\omega}^{\nu} = \frac{\partial \bigwedge_{\lambda\mu}^{\nu}}{\partial p^{\omega}} - \frac{\partial \bigwedge_{\lambda\omega}^{\nu}}{\partial p^{\mu}} + \bigwedge_{\mu\omega}^{j} \bigwedge_{\lambda\mu}^{i} - \bigwedge_{\mu\mu}^{i} \bigwedge_{\lambda\omega}^{j},$$

which are all affinors, where

(3. 4)
$$K^{\nu}_{\lambda\mu\omega} = \frac{\partial \Gamma^{\nu}_{\lambda\mu}}{\partial x^{\omega}} - \frac{\partial \Gamma^{\nu}_{\lambda\omega}}{\partial x^{i\ell}} + \Gamma^{\nu}_{\rho\omega}\Gamma^{\rho}_{\lambda\mu} - \Gamma^{\nu}_{\rho\mu}\Gamma^{\rho}_{\lambda\omega},$$

(3. 5)
$$\vec{K}_{\lambda\mu\omega}^{i} = \frac{\partial \Gamma_{\lambda\mu}^{\nu}}{\partial p^{\omega}} - \frac{\partial \bigwedge_{\lambda\omega}^{i}}{\partial x^{\mu}} + \bigwedge_{\rho\omega}^{i} \Gamma_{\lambda\mu}^{\rho} - \Gamma_{\rho\mu}^{\nu} \bigwedge_{\lambda\omega}^{i} ,$$

which are not affinors. $\overset{i}{R}_{\lambda\mu\omega}^{\nu}$ is of the class -i and $\overset{i,j}{R}_{\lambda\mu\omega}^{\nu}$ of the class -(i+j). It follows after some calculation

$$(3. 6) \qquad (\delta_{2}\delta_{1} - \delta_{1}\delta_{2})v^{\nu} = \{K^{\nu}_{,\lambda\mu\omega}d_{1}x^{\mu}d_{2}x^{\omega} + \sum_{i} \overset{i}{K}^{\nu}_{,\lambda\mu\omega}(d_{1}x^{\mu}d_{2}\overset{i}{p}^{\omega} - d_{1}\overset{i}{p}^{\omega}d_{2}x^{\mu}) \\ + \sum_{i,j} \overset{i}{R}^{\nu}_{,\lambda\mu\omega}d_{1}\overset{i}{p}^{\mu}d_{2}\overset{j}{p}^{\omega}\}v^{\lambda} \\ = \{R^{\nu}_{,\lambda\mu\omega}d_{1}x^{\mu}d_{2}x^{\omega} + \sum_{i} \overset{i}{R}^{\nu}_{,\lambda\mu\omega}(d_{1}x^{\mu}\overset{i}{\delta}^{j}_{2}\overset{i}{p}^{\omega} - \overset{i}{\delta}^{i}_{1}\overset{i}{p}^{\omega}d_{2}x^{\mu}) \\ + \sum_{i,j} \overset{i}{R}^{i,j}_{,\lambda\mu\omega}\overset{i}{\delta}^{i}_{1}\overset{j}{p}^{\mu}\overset{j}{\delta}^{j}_{2}\overset{j}{p}^{\omega}\}v^{\lambda}$$

for the surrounding of the vector v^{ν} about an infinitesimal circuit.

4. On the other hand, we have

$$(4. 1) \qquad (\stackrel{j}{\nabla}_{\omega}\stackrel{i}{\nabla}_{\mu}-\stackrel{i}{\nabla}_{\mu}\stackrel{j}{\nabla}_{\omega})v^{\nu}=\stackrel{i,j}{R^{\nu}}_{\lambda\mu\omega}v^{\lambda}-\stackrel{j}{\wedge}_{\mu\omega}\stackrel{i}{\nabla}_{\pi}v^{\nu}+\stackrel{i}{\wedge}_{\omega\mu}\stackrel{j}{\nabla}_{\pi}v^{\nu}\\ =\stackrel{(i,j)}{(R^{\nu}}_{\lambda\mu\omega}-\stackrel{j}{\wedge}_{\mu\omega}\stackrel{i}{\wedge}_{\lambda\pi}^{\nu}+\stackrel{i}{\wedge}_{\omega\mu}\stackrel{j}{\wedge}_{\lambda\pi}^{\nu})v^{\lambda}-\left(\stackrel{j}{\wedge}_{\mu\omega}\stackrel{i}{\partial}\stackrel{i}{v}-\stackrel{i}{\wedge}_{\mu\omega}\stackrel{i}{\partial}\stackrel{j}{v}\stackrel{j}{\rho}\right),\\ (4. 2) \qquad (\stackrel{i}{\nabla}_{\rho}\nabla_{\mu}-\nabla_{\mu}\stackrel{i}{\nabla}_{\rho})v^{\nu}=\stackrel{R^{\nu}}{N^{\nu}}_{\lambda\mu\rho}v^{\lambda}+\stackrel{i}{P^{\lambda}}_{\rho\mu}\stackrel{i}{\nabla}_{\lambda}v^{\nu}-\stackrel{i}{\wedge}\stackrel{i}{\lambda}_{\mu\rho}^{\mu}\nabla_{\lambda}v^{\nu}$$

(4. 2)
$$(\bigtriangledown_{p}\lor_{\mu}-\bigtriangledown_{\mu}\lor_{p})v = K_{\lambda\mu\rho}v^{\mu}+P_{\rho\mu}\lor_{\lambda}v^{\nu}-\bigwedge_{\mu\rho}\lor_{\lambda}v^{\nu}$$

 $-\sum_{j}\bigvee_{\lambda}v^{\nu}\left(\frac{\partial I_{\pi\mu}}{\partial p^{\rho}}-j\delta_{\pi}^{\lambda}\frac{\partial I_{\mu}}{\partial p^{\rho}}\right)^{j}p^{\pi},$

(4. 3)
$$(\nabla_{\omega}\nabla_{\mu} - \nabla_{\mu}\nabla_{\omega})v^{\nu} = R^{\nu}_{\lambda\mu\omega}v^{\lambda} + 2T^{\lambda}_{\omega\mu}\nabla_{\lambda}v^{\nu} + \sum_{i} \dot{Q}^{\lambda}_{\tau\omega\mu}\dot{p}^{\tau}\dot{\nabla}_{\lambda}v^{\nu},$$

where

(4. 4)
$$T_{*\omega\mu}^{\lambda} = \Gamma_{[\omega\mu]}^{\lambda} - \sum_{j} (I_{\pi[\mu}^{j} - j\delta_{\pi}^{\tau} I_{[\mu]}^{j}) p^{\pi} \bigwedge_{\omega]\tau}^{j} = P_{[\omega\mu]}^{i} - I_{[\omega\mu]}^{i} + i\delta_{[\omega}^{\lambda} (\Gamma_{\mu]} - I_{\mu]}^{i}),$$

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$$(4. 5) \qquad \frac{1}{2} \dot{Q}_{\tau\tau\mu\omega}^{\lambda} = \frac{\partial \dot{I}_{\tau[\mu}^{\lambda}}{\partial x^{\omega]}} - i \delta_{\tau}^{\lambda} \frac{\partial \dot{I}_{[\mu}}{\partial x^{\omega]}} - \dot{I}_{\pi[\mu}^{\lambda} \dot{I}_{[\tau]\omega]}^{i} + \sum_{j} (\dot{I}_{\rho[\mu}^{j} - j \delta_{\rho}^{\pi} \dot{I}_{[\mu]}^{j}) p^{\omega} \\ \times \frac{\partial}{\partial p^{\pi}} (\dot{I}_{|\tau|\mu]}^{\lambda} - i \delta_{|\tau|}^{\lambda} \dot{I}_{\omega]}^{i}).$$

 $T^{\lambda}_{\cdot \omega \mu}$ and $\tilde{Q}^{\lambda}_{\cdot \tau \mu \omega}$ are also affinors, of which the former is the *torsion* tensor of the connection and the latter are the *curvature tensors* of the base-connections.

5. In Cartan's case we may assume $I_{\mu}^{1}=0$, as there exists an integral invariant $s=\int \sqrt{F(x,x')}$, which will be taken as the parameter for the line element $p^{\nu}=\frac{dx^{\nu}}{ds}$. Let $\bigwedge_{\lambda\mu\nu}=g_{\lambda\sigma}\bigwedge_{\mu\nu}^{i}$ be symmetric and $\bigwedge_{\lambda\mu\nu}^{1}p^{\nu}=0$, being $g_{\lambda\mu}=F_{x'}\lambda_{x'}\mu$. For the parameters of the base-connection we take

(5. 1)
$$\Gamma^{\nu}_{\lambda\mu} = \Gamma^{\nu}_{\lambda\mu} - \bigwedge^{1}_{\lambda\tau} \Gamma^{1}_{\pi\mu} p^{\pi} = \Gamma^{\nu}_{\lambda\mu} - \bigwedge^{1}_{\lambda\tau} \Gamma^{\tau}_{\pi\mu} p^{\pi}$$

and assume $\delta g_{\lambda\mu} = 0$. Then the curvature tensors (3.1)-(3.3) are nothing but those introduced by Cartan. In this case we have

(5. 2)
$$\dot{P}^{\nu}_{\lambda\mu} = 0$$
,

$$(5. 3) T^{\nu}_{\lambda\mu} = 0$$

(5. 4)
$$\overset{1}{Q}^{\lambda}_{\tau \tau \mu \omega} p^{\tau} = R^{\nu}_{\tau \tau \mu \omega} p^{\tau}$$

From (5.3) we can conclude that Cartan's connection is symmetric or has no torsion.

1) Cartan, loc. cit. Cartan uses the notation $C_{\lambda\mu}^{\nu}$ for our $\Lambda_{\lambda\mu}^{\nu}$.

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