

11. The Robin Problems on Domains with Many Tiny Holes

By Satoshi KAIZU

Department of Information Mathematics, University
of Electro-Communications

(Communicated by Kôzaku YOSIDA, M. J. A., Feb. 12, 1985)

1. Introduction. Let Ω be a bounded domain of \mathbf{R}^N , $N \geq 2$, with smooth boundary Γ and with the outer unit normal vector ν_Γ on Γ . Let \mathbf{R}^N be divided into an infinitely many number of cubes C_i^ε , $i \in N$, with volume of $(2\varepsilon)^N$ and let $B_i(r^\varepsilon)$ be a closed ball of radius r^ε ($< \varepsilon$) centered in C_i^ε . From Ω we remove all such balls and obtain a $D^\varepsilon(\subset \Omega)$ with n^ε holes. Under the case $n^\varepsilon \rightarrow \infty$, $r^\varepsilon \rightarrow 0$, a parameter, which determines the behavior of the Laplacian on D^ε with the Dirichlet condition, is known by M. Kac [2], J. Rauch and M. Taylor [3], D. Cioranescu and F. Murat [1]. That parameter is given by $\lim n^\varepsilon (r^\varepsilon)^{N-2}$ for $N \geq 3$ and $\lim n^\varepsilon / |\log r^\varepsilon|$ for $N=2$, ε means values of a fixed sequence decreasing to zero. Now, we show a different parameter is important for the Robin problems.

From Ω we remove all balls $B_i(r^\varepsilon)$ such that $\text{dist}(B_i(r^\varepsilon), \Gamma) \geq \varepsilon$ and obtain R^ε with n^ε holes. Let α be a positive constant and ν^ε the outer unit normal vector on ∂R^ε . We consider the Robin problem: for $f \in L^2(\Omega)$ find $u^\varepsilon \in H^1(R^\varepsilon)$ such that

$$(1) \quad \begin{aligned} -\Delta u^\varepsilon &= f && \text{a.e. in } R^\varepsilon, \\ \frac{\partial u^\varepsilon}{\partial \nu^\varepsilon} + \alpha u^\varepsilon &= 0 && \text{a.e. on } \partial R^\varepsilon. \end{aligned}$$

Theorem 1. Let u^ε be the solution of (1) and $\tilde{u}^\varepsilon \in H^1(\Omega)$ an extension of u^ε to be harmonic in F^ε , $F^\varepsilon = \Omega \setminus R^\varepsilon$. Assume that $r^\varepsilon \rightarrow 0$ and $n^\varepsilon \rightarrow \infty$ with the conditions $\eta = \lim n^\varepsilon (r^\varepsilon)^{N-1}$, $0 < \eta < \infty$. Then \tilde{u}^ε converges weakly in $H^1(\Omega)$ to the solution of the problem:

$$(2) \quad \begin{aligned} -\Delta u + \frac{\alpha S_N \eta u}{|\Omega|} &= f && \text{a.e. in } \Omega, \\ \frac{\partial u}{\partial \nu_\Gamma} + \alpha u &= 0 && \text{a.e. on } \Gamma. \end{aligned}$$

Here $|\Omega|$ means the volume of Ω and S_N means the surface area of the unit sphere of \mathbf{R}^N .

2. Abstract scheme. Let Ω be the same domain as in Section 1. We introduce a certain limit of the minus Laplacian, which corresponds to one of versions for theorem 1.2 of Cioranescu-Murat [1] in the case of Robin condition.

For a subdomain G of Ω we regard all functions of $L^2(G)$ as functions of $L^2(\Omega)$ vanishing outside G . In this section R^ε means a subdomain of Ω satisfying (a.1) below. Let $a^\varepsilon: H^1(R^\varepsilon) \times H^1(R^\varepsilon) \rightarrow \mathbf{R}$ be a bilinear form defined by

$$\alpha^\varepsilon(v, w) = \int_{R^\varepsilon} \nabla v \cdot \nabla w dx + \alpha \int_{\partial R^\varepsilon} v w d\sigma.$$

Let us consider (1). For (1) we have

$$(3) \quad \alpha^\varepsilon(u^\varepsilon, v) = \int_{R^\varepsilon} f v dx \quad \text{for all } v \in H^1(R^\varepsilon).$$

We consider the behavior of the solution of (3). We use the norm of $H^1(R^\varepsilon)$ defined by $\|v\|_{H^1(R^\varepsilon)}^2 = \|\nabla v\|_{L^2(R^\varepsilon)^N}^2 + \|v\|_{L^2(R^\varepsilon)}^2$.

We set $F^\varepsilon = \Omega \setminus R^\varepsilon$. We assume the following conditions.

(a.1) F^ε does not meet Γ and its interior kernel is a nonempty set with locally Lipschitz boundary.

(a.2) There exists a family of extension maps $T^\varepsilon : H^1(R^\varepsilon) \rightarrow H^1(\Omega)$ such that

- (i) $\limsup \|T^\varepsilon\|_{L(H^1(R^\varepsilon), H^1(\Omega))} < \infty$,
- (ii) if $\limsup \alpha^\varepsilon(v^\varepsilon, v^\varepsilon) < \infty$ then $T^\varepsilon v^\varepsilon - v^\varepsilon \rightarrow 0$ in $L^2(\Omega)$.

(a.3) There exists a constant c_1 such that

$$\alpha^\varepsilon(v, v) \geq c_1 \|v\|_{H^1(R^\varepsilon)}^2 \text{ for all } v \in H^1(R^\varepsilon) \text{ and } \varepsilon > 0.$$

(a.4) There exists a family $\{\theta^\varepsilon \in W^{1,\infty}(R^\varepsilon)\}_\varepsilon$ satisfying the conditions below.

- (i) $\theta^\varepsilon = 1$ a.e. on Γ .
- (ii) Set $\tilde{\theta}^\varepsilon = T^\varepsilon \theta^\varepsilon$. Then $\tilde{\theta}^\varepsilon \xrightarrow{w} 1$ in $H^1(\Omega)$ as $\varepsilon \rightarrow 0$.
- (iii) $\limsup \alpha^\varepsilon(\theta^\varepsilon, \theta^\varepsilon) < \infty$.
- (iv) There exists $\bar{\eta} \in W^{1,1}(\Omega)^*$ such that

$$-\nabla \cdot (\chi^\varepsilon \nabla \theta^\varepsilon) + \alpha \theta^\varepsilon \delta(\partial F^\varepsilon) \xrightarrow{s} \bar{\eta} \quad \text{in } W^{1,1}(\Omega)^*,$$

$W^{1,1}(\Omega)^*$ means the dual space of $W^{1,1}(\Omega)$, χ^ε means the characteristic function of R^ε and $\delta(\partial F^\varepsilon)$ means the measure defined by

$$\langle \delta(\partial F^\varepsilon), v \rangle = \int_{\partial F^\varepsilon} v d\sigma \quad \text{for all } v \in W^{1,1}(\Omega).$$

Theorem 2. *Let u^ε be the solution of (3) and set $\tilde{u}^\varepsilon = T^\varepsilon u^\varepsilon$. Under all the conditions from (a.1) to (a.4), \tilde{u}^ε converges weakly in $H^1(\Omega)$ to u , where u is the solution of*

$$(4) \quad \begin{aligned} -\Delta u + \bar{\eta} u &= f && \text{a.e. in } \Omega, \\ \frac{\partial u}{\partial \nu_r} + \alpha u &= 0 && \text{a.e. on } \Gamma. \end{aligned}$$

Proof. Substituting $v = u^\varepsilon$ into (3) and using (a.2) and (a.3) we can see

$$\limsup \| \tilde{u}^\varepsilon \|_{H^1(\Omega)} < \infty.$$

We can choose a weakly convergent subsequence $\{u_n\}_n$ such that $\tilde{u}_n \xrightarrow{w} u$ in $H^1(\Omega)$, where $u_n = u^{\varepsilon_n}$. It suffices to show

$$(5) \quad \int_{\Omega} \nabla u \cdot \nabla \zeta dx + \langle \bar{\eta}, u \zeta \rangle + \alpha \int_{\Gamma} u \zeta d\sigma = \int_{\Omega} f \zeta dx$$

for all $\zeta \in C^\infty(\bar{\Omega})$, where $\langle \cdot, \cdot \rangle$ denotes the dual pair between $W^{1,1}(\Omega)^*$ and $W^{1,1}(\Omega)$. We can substitute $v = \theta_n \zeta$ into (3), where $\theta_n = \theta^{\varepsilon_n}$. We set $R_n = R^{\varepsilon_n}$ and $F_n = F^{\varepsilon_n}$. Thus

$$(6) \quad \begin{aligned} & \int_{R_n} \theta_n \nabla u_n \cdot \nabla \zeta dx + \int_{R_n} \zeta \nabla u_n \cdot \nabla \theta_n dx + \alpha \int_{\partial F_n} \zeta u_n \theta_n d\sigma + \alpha \int_{\Gamma} u_n \zeta d\sigma \\ &= \int_{R_n} f \theta_n \zeta dx. \end{aligned}$$

Here we have used the equality $u_n = u_n \theta_n$ a.e. on Γ . By the definition of u_n and the conditions (a.2)-(ii), (a.4)-(ii) and (a.4)-(iii) we have $\nabla \tilde{u}_n \xrightarrow{w} \nabla u$, $\tilde{\theta}_n \nabla \zeta \xrightarrow{s} \nabla \zeta$ in $L^2(\Omega)^N$, $\tilde{\theta}_n - \theta_n \xrightarrow{s} 0$ in $L^2(\Omega)$ and $\tilde{\theta}_n \xrightarrow{s} 1$ in $L^2(\Omega)$, $\tilde{u}_n \xrightarrow{s} u$ in $L^2(\Gamma)$. Thus,

$$\begin{aligned} & \int_{R_n} \theta_n (\nabla u_n \cdot \nabla \zeta - f \zeta) dx + \alpha \int_{\Gamma} u_n \zeta d\sigma \\ &= \int_{\Omega} \tilde{\theta}_n (\nabla \tilde{u}_n \cdot \nabla \zeta - f \zeta) dx + \alpha \int_{\Gamma} \tilde{u}_n \zeta d\sigma + \int_{\Omega} (\theta_n - \tilde{\theta}_n) [\nabla \tilde{u}_n \cdot \nabla \zeta - f \zeta] dx \\ & \longrightarrow \int_{\Omega} (\nabla u \cdot \nabla \zeta - f \zeta) dx + \alpha \int_{\Gamma} u \zeta d\sigma. \end{aligned}$$

Notice $\chi_n \zeta \nabla u_n = \chi_n [\nabla (\zeta \tilde{u}_n) - \tilde{u}_n \nabla \zeta]$, where $\chi_n = \chi^{\varepsilon n}$. Therefore,

$$\begin{aligned} & \int_{R_n} \zeta \nabla u_n \cdot \nabla \theta_n dx + \alpha \int_{\partial F_n} u_n \theta_n \zeta d\sigma \\ &= \langle -\nabla \cdot (\chi_n \nabla \tilde{\theta}_n) + \alpha \theta_n \delta(\partial F_n), \tilde{u}_n \zeta \rangle - \int_{R_n} u_n \nabla \theta_n \cdot \nabla \zeta dx. \end{aligned}$$

Using (a.2)-(ii) and (a.4)-(ii) we have $u_n \nabla \zeta \xrightarrow{s} u \nabla \zeta$, $(u_n - \tilde{u}_n) \nabla \zeta \xrightarrow{s} 0$ in $L^2(\Omega)^N$, $\nabla \tilde{\theta}_n \xrightarrow{w} 0$ in $L^2(\Omega)^N$. Clearly, $\zeta \tilde{u}_n \xrightarrow{w} \zeta u$ in $W^{1,1}(\Omega)$. Also, using (a.4)-(iv) we get

$$\begin{aligned} & \int_{R_n} u_n \nabla \theta_n \cdot \nabla \zeta dx = \int_{\Omega} \tilde{u}_n \nabla \tilde{\theta}_n \cdot \nabla \zeta dx + \int_{\Omega} (u_n - \tilde{u}_n) \nabla \tilde{\theta}_n \cdot \nabla \zeta dx \longrightarrow 0, \\ & \int_{R_n} \zeta \nabla u_n \cdot \nabla \theta_n dx + \alpha \int_{\partial F_n} u_n \theta_n \zeta d\sigma \longrightarrow \langle \bar{\eta}, u \zeta \rangle. \end{aligned}$$

Therefore we get (5). The proof is completed. q.e.d.

3. Proof of Theorem 1. We introduce special functions θ^ε such that $\theta^\varepsilon = 1$ on $\Omega \setminus \cup \{B_i(\varepsilon) : 1 \leq i \leq n^\varepsilon\}$, $\Delta \theta^\varepsilon = 0$ on $\cup \{B_i(\varepsilon) \setminus B_i(r^\varepsilon) : 1 \leq i \leq n^\varepsilon\}$ and $\partial \theta^\varepsilon / \partial \nu^\varepsilon + \alpha \theta^\varepsilon = 0$ on $\partial F^\varepsilon (\subset \partial R^\varepsilon)$. By a similar way to that in [1] we can see functions θ^ε satisfies (a.4) with $\bar{\eta} = \alpha S_N \eta / |\Omega|$. We show (a.3). We set $E_i = C_i^\varepsilon \cap R^\varepsilon$, $1 \leq i \leq n^\varepsilon$. Choose ε so small that

$$2 \int_{r^\varepsilon}^{N^{1/2\varepsilon}} r^{N-1} dr \max \left\{ \int_{r^\varepsilon}^{N^{1/2\varepsilon}} \rho^{1-N} d\rho, (r^\varepsilon)^{1-N} \right\} \leq \frac{3}{\eta} \left(\frac{N}{2} \right)^N |\Omega|.$$

Here we use $n^\varepsilon \sim |\Omega| / (2\varepsilon)^N$, $\varepsilon \rightarrow 0$. For the equality,

$$|v(r, \omega)|^2 = \left[\int_{r^\varepsilon}^r \frac{\partial v}{\partial \rho} d\rho + v(r^\varepsilon, \omega) \right]^2, \quad v \in H^1(R^\varepsilon), \quad \omega \in \partial B_i(1),$$

using the Schwarz inequalities twice on the left hand side, multiplying both sides by r^{N-1} , integrating them over $(r^\varepsilon, \rho_i(\omega))$, $\rho_i(\omega) = \sup \{r \geq 0 : r\omega \in E_i\}$, next, over the unit sphere $\partial B_i(1)$, we get the inequality

$$\int_{E_i} |v|^2 dx \leq 3 \left(\frac{N}{2} \right)^N \frac{|\Omega|}{\eta} \left[\int_{E_i} |\nabla v|^2 dx + \int_{\partial B_i(r^\varepsilon)} |v|^2 d\sigma \right].$$

We denote by G_i a non-empty set $C_i^\varepsilon \cap R^\varepsilon$, $i > n^\varepsilon$. For sufficiently small ε we have a Lipschitz function h_i such that $G_i = \{(x', x_N) : |x_i| \leq \varepsilon, 1 \leq i \leq N-1, 0 \leq x_N \leq h_i(x')\}$ and $C_i^\varepsilon \cap \Gamma = \{(x', h_i(x') : |x_i| \leq \varepsilon, 1 \leq i \leq N-1\}$, $x' = (x_1, \dots, x_{N-1})$. Similarly, for the same v as the above we have a constant C_1 independent of ε such that

$$\int_{G_i} |v|^2 dx \leq C_i \varepsilon \left(\int_{G_i} |\nabla v|^2 dx + \int_{C_i^\varepsilon \cap \Gamma} |v|^2 d\sigma \right).$$

By these inequalities we can see (a.3). The family of extensions: $H^1(R^\varepsilon) \rightarrow H^1(\Omega)$ in Theorem 1 is uniformly bounded with respect to their operator norm (cf. [3]). So (a.2)-(i) holds. The condition (a.1) also holds. Lastly, (a.2)-(ii) follows from the fact such that the linear operator: $L^2(\partial B_i(r^\varepsilon)) \ni v \rightarrow v_r \in L^2(\partial B_i(r^\varepsilon))$ is bounded with norm one for $0 < r < 1$. Here $v_r(r^\varepsilon, \omega) = \tilde{v}(x_i^\varepsilon + r r^\varepsilon \omega)$, where x_i^ε denotes the center of $B_i(r^\varepsilon)$. The conditions of Theorem 2 are all verified. q.e.d.

References

- [1] D. Cioranescu and F. Murat: Un terme étrange venu d'ailleurs I, II. Nonlinear partial differential equations and their applications. Vol. II-III, Pitman, Boston, London, Melbourne (1982).
- [2] M. Kac: Probabilistic methods in some problems of scattering theory. Rocky Mountain J. of Math., 4, 511-538 (1974).
- [3] J. Rauch and M. Taylor: Potential and scattering theory on wildly perturbed domains. J. Funct. Anal., 18, 27-59 (1975).