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# CLIFFORD ALGEBRAS AND THEIR APPLICATIONS TO LIE GROUPS AND SPINORS 

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#### Abstract

We discuss some well-known facts about Clifford algebras: matrix representations, Cartan's periodicity of 8 , double coverings of orthogonal groups by spin groups, Dirac equation in different formalisms, spinors in $n$ dimensions, etc. We also present our point of view on some problems. Namely, we discuss the generalization of the Pauli theorem, the basic ideas of the method of averaging in Clifford algebras, the notion of quaternion type of Clifford algebra elements, the classification of Lie subalgebras of specific type in Clifford algebra, etc.


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## Introduction

Clifford algebra was invented by W. Clifford [11]. In his research, he combined Hamilton's quaternions [23] and Grassmann's exterior algebra [27]. Further development of the theory of Clifford algebras is associated with a number of famous mathematicians and physicists - R. Lipschitz, T. Vahlen, E. Cartan [9], E. Witt, C. Chevalley, M. Riesz [42] and others. Dirac equation [15], [16] had a great influence on the development of Clifford algebra. Also note the article [5].
Nowadays Clifford algebra is used in different branches of modern mathematics and physics. There are different applications of Clifford algebra in physics, analysis, geometry, computer science, mechanics, robotics, signal and image processing, etc.
In this text, we discuss some well-known facts about Clifford algebras: matrix representations, Cartan's periodicity of 8 , double coverings of orthogonal groups by spin groups, Dirac equation in different formalisms, spinors in $n$ dimensions, etc. We also present our point of view on some problems. Namely, we discuss the generalization of the Pauli theorem, the basic ideas of the method of averaging in Clifford algebras, the notion of quaternion type of Clifford algebra elements, the classification of Lie subalgebras of specific type in Clifford algebra, etc.
We recommend a number of classical books on Clifford algebras and applications [31], [25], [26], [10], [40], [7], [29], [8], [20], [13], [20], [24], [67], [68], [18], etc.

We can recommend also the book [35] and a course of lectures [55] for Russian readers.

## 1. Definition of Clifford Algebra

### 1.1. Clifford Algebra as a Quotient Algebra

In [31], one can find five different (equivalent) definitions of Clifford algebra. We will discuss two definitions of Clifford algebra in this work. Let us start with the definition of Clifford algebra as a quotient algebra [10].

Definition 1. Let we have a vector space $V$ of arbitrary finite dimension $n$ over the field $\mathbb{R}$ and a quadratic form $Q: V \rightarrow \mathbb{R}$. Consider the tensor algebra

$$
T(V)=\bigoplus_{k=0}^{\infty} \bigotimes^{k} V
$$

and the two-sided ideal $I(V, Q)$ generated by all elements of the form $x \otimes x-Q(x) e$ for $x \in V$, where $e$ is the identity element. Then the following quotient algebra

$$
C \ell(V, Q)=T(V) / I(V, Q)
$$

is called a real Clifford algebra.

### 1.2. Clifford Algebra with Fixed Basis

Now let us discuss definition of the real Clifford algebra with fixed basis which is more useful for calculations and different applications.

Definition 2. Let $n$ be a natural number and $E$ be a linear space of dimension $2^{n}$ over the field of real numbers $\mathbb{R}$ with the basis enumerated by the ordered multiindices with a length between 0 and $n$

$$
e, e_{a_{1}}, e_{a_{1} a_{2}}, \ldots, e_{1 \ldots n}
$$

where $1 \leq a_{1}<a_{2}<\cdots<a_{n} \leq n$. Let us introduce the operation of multiplication on $E$

- with the properties of distributivity and associativity

$$
\begin{aligned}
& U(\alpha V+\beta W)=\alpha U V+\beta U W, \quad U, V, W \in E, \quad \alpha, \beta \in \mathbb{R} \\
& (\alpha U+\beta V) W=\alpha U W+\beta V W, \quad U, V, W \in E, \quad \alpha, \beta \in \mathbb{R} \\
& U(V W)=(U V) W, \quad U, V, W \in E
\end{aligned}
$$

- $e$ is the identity element

$$
U e=e U=U, \quad U \in E
$$

- $e_{a}, a=1, \ldots, n$ are generators

$$
e_{a_{1}} e_{a_{2}} \cdots e_{a_{n}}=e_{a_{1} \ldots a_{n}}, \quad 1 \leq a_{1}<a_{2}<\cdots<a_{n} \leq n
$$

- generators satisfy

$$
e_{a} e_{b}+e_{b} e_{a}=2 \eta_{a b} e
$$

where

$$
\begin{equation*}
\eta=\left\|\eta_{a b}\right\|=\operatorname{diag}(\underbrace{1, \ldots, 1}_{p}, \underbrace{-1, \ldots,-1}_{q}, \underbrace{0, \ldots, 0}_{r}), \quad p+q+r=n \tag{1}
\end{equation*}
$$

is a diagonal matrix with $p$ times $1, q$ times -1 , and $r$ times 0 on the diagonal.
The linear space $E$ with such operation of multiplication is called real Clifford algebra $C \ell_{p, q, r}$.

Example 1. In the case $r=0$, we obtain nondegenerate Clifford algebra $C \ell_{p, q}:=$ $C \ell_{p, q, 0}$. The quadratic form $Q$ in Definition 1 is nondegenerate in this case.

Example 2. In the case $r=0, q=0$, we obtain Clifford algebra $C \ell_{n}:=C \ell_{n, 0,0}$ of Euclidian space. The quadratic form $Q$ in Definition 1 is positive definite in this case.

Example 3. In the case $p=q=0, r=n$, we obtain Grassmann algebra $\Lambda_{n}:=$ $C \ell_{0,0, n}$. In this case Clifford multiplication is called exterior multiplication and it is denoted by $\wedge$. The generators of Grassmann algebra satisfy conditions $e_{a} \wedge e_{b}=$ $-e_{b} \wedge e_{a}, a, b=1, \ldots, n$.

Any element of the real Clifford algebra $C \ell_{p, q, r}$ has the form

$$
\begin{equation*}
U=u e+\sum_{a=1}^{n} u_{a} e_{a}+\sum_{a<b} u_{a b} e_{a b}+\cdots+u_{1 \ldots n} e_{1 \ldots n} \tag{2}
\end{equation*}
$$

where $u, u_{a}, u_{a b}, \ldots, u_{1 \ldots n} \in \mathbb{R}$ are real numbers.
Also we consider complexified Clifford algebras $\mathbb{C} \otimes C \ell_{p, q, r}$. Any element of the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{p, q, r}$ has the form (2), where $u, u_{a}, u_{a b}, \ldots$, $u_{1 \ldots n} \in \mathbb{C}$ are complex numbers.
Also we consider complex Clifford algebras. In Definition 1, we must take vector space $V$ over the field of complex numbers $\mathbb{C}$ in this case. In Definition 2, we must take vector space $E$ over the field of complex numbers $\mathbb{C}$ and it is sufficient to consider matrix $\eta=\operatorname{diag}(1, \ldots, 1,0, \ldots, 0), p+r=n$, with $p$ times 1 and $r$ times 0 on the diagonal instead of the matrix (1) in this case. The most popular case is $C \ell\left(\mathbb{C}^{n}\right)$, when the quadratic form $Q$ is nondegenerate and $\eta$ is the identity matrix.

### 1.3. Examples in Small Dimensions

Example 4. In the case of $C \ell_{0}$, arbitrary Clifford algebra element has the form $U=u e$, where $e^{2}=e$. We obtain the isomorphism $C \ell_{0} \cong \mathbb{R}$.

Example 5. In the case of $C \ell_{1}$, arbitrary Clifford algebra element has the form $U=u e+u_{1} e_{1}$, where $e$ is the identity element and $e_{1}^{2}=e$. We obtain the isomorphism with double numbers: $C \ell_{1} \cong \mathbb{R} \oplus \mathbb{R}$.

Example 6. In the case of $C \ell_{0,1}$, arbitrary Clifford algebra element has the form $U=u e+u_{1} e_{1}$, where $e$ is the identity element and $e_{1}^{2}=-e$. We obtain the isomorphism with complex numbers: $C \ell_{0,1} \cong \mathbb{C}$.

Example 7. In the case of $\mathrm{Cl}_{0,2}$, arbitrary Clifford algebra element has the form $U=u e+u_{1} e_{1}+u_{2} e_{2}+u_{12} e_{12}$. We can easily verify the following relations

$$
\begin{aligned}
\left(e_{1}\right)^{2} & =\left(e_{2}\right)^{2}=-e, & \left(e_{12}\right)^{2} & =e_{1} e_{2} e_{1} e_{2}=-e_{1} e_{1} e_{2} e_{2}=-e \\
e_{1} e_{2} & =-e_{2} e_{1}=e_{12}, & e_{2} e_{12} & =-e_{12} e_{2}=e_{1}, \quad e_{12} e_{1}=-e_{1} e_{12}=e_{2}
\end{aligned}
$$

Using the following substitution

$$
e_{1} \rightarrow \mathrm{i}, \quad e_{2} \rightarrow \mathrm{j}, \quad e_{12} \rightarrow \mathrm{k}
$$

where $\mathrm{i}, \mathrm{j}$, and k are imaginary units of quaternions, we obtain the isomorphism $C l_{0,2} \simeq \mathbb{H}$.

Recall that $\mathbb{H}$ is an associative division algebra. An arbitrary quaternion has the form

$$
q=a 1+b \mathrm{i}+c \mathrm{j}+d \mathrm{k} \in \mathbb{H}, \quad a, b, c, d \in \mathbb{R}
$$

where 1 is the identity element, $\mathrm{i}^{2}=\mathrm{j}^{2}=\mathrm{k}^{2}=-1, \mathrm{ij}=-\mathrm{ji}=\mathrm{k}, \mathrm{jk}=-\mathrm{kj}=\mathrm{i}$, $\mathrm{ki}=-\mathrm{ik}=\mathrm{j}$.
If $q \neq 0$, then $q^{-1}=\frac{1}{\|q\|^{2}} \bar{q}$, where

$$
\bar{q}:=a-b \mathrm{i}-c \mathrm{j}-d \mathrm{k}, \quad\|q\|:=\sqrt{q \bar{q}}=\sqrt{a^{2}+b^{2}+c^{2}+d^{2}} .
$$

Note that $C \ell_{2,0} \cong C \ell_{1,1} \nsubseteq C \ell_{0,2}$ (see Section 3.1).
Example 8. Let us consider the Pauli matrices

$$
\sigma_{0}=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad \sigma_{1}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad \sigma_{2}=\left(\begin{array}{rr}
0 & -\mathrm{i} \\
\mathrm{i} & 0
\end{array}\right), \quad \sigma_{3}=\left(\begin{array}{rr}
1 & 0 \\
0 & -1
\end{array}\right) .
$$

W. Pauli [39] have introduced these matrices to describe the spin of the electron. It can be easily verified that

$$
\begin{aligned}
& \sigma_{1} \sigma_{2}=\mathrm{i} \sigma_{3}, \quad \sigma_{2} \sigma_{3}=\mathrm{i} \sigma_{1}, \quad \sigma_{3} \sigma_{1}=\mathrm{i} \sigma_{2} \\
& \left(\sigma_{a}\right)^{\dagger}=\sigma_{a}, \quad \operatorname{tr}\left(\sigma_{a}\right)=0, \quad\left(\sigma_{a}\right)^{2}=\sigma_{0}, \\
& a, b=1,2,3 \text {. }
\end{aligned}
$$

Using the substitution
$e \rightarrow \sigma_{0}, \quad e_{a} \rightarrow \sigma_{a}, a=1,2,3, \quad e_{a b} \rightarrow \sigma_{a} \sigma_{b}, a<b, \quad e_{123} \rightarrow \sigma_{1} \sigma_{2} \sigma_{3}$
we obtain the isomorphism

$$
C \ell_{3} \cong \operatorname{Mat}(2, \mathbb{C}) .
$$

The matrices

$$
\left\{\sigma_{0}, \quad \sigma_{1}, \quad \sigma_{2}, \quad \sigma_{3}, \quad \mathrm{i} \sigma_{1}, \quad \mathrm{i} \sigma_{2}, \quad \mathrm{i} \sigma_{3}, \quad \mathrm{i} \sigma_{0}\right\}
$$

constitute a basis of $\operatorname{Mat}(2, \mathbb{C})$.
Example 9. Let us consider the Dirac gamma-matrices [15], [16]

$$
\begin{array}{ll}
\gamma_{0}=\left(\begin{array}{rrrr}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & -1 & 0 \\
0 & 0 & 0 & -1
\end{array}\right), & \gamma_{1}=\left(\begin{array}{rrrr}
0 & 0 & 0 & 1 \\
0 & 0 & 1 & 0 \\
0 & -1 & 0 & 0 \\
-1 & 0 & 0 & 0
\end{array}\right) \\
\gamma_{2}=\left(\begin{array}{rrrr}
0 & 0 & 0 & -\mathrm{i} \\
0 & 0 & \mathrm{i} & 0 \\
0 & \mathrm{i} & 0 & 0 \\
-\mathrm{i} & 0 & 0 & 0
\end{array}\right), & \gamma_{3}=\left(\begin{array}{rrrr}
0 & 0 & 1 & 0 \\
0 & 0 & 0 & -1 \\
-1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0
\end{array}\right) .
\end{array}
$$

These matrices satisfy conditions

$$
\begin{array}{rlrl}
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a} & =2 \eta_{a b} \mathbf{1}, & a, b & =0,1,2,3, \\
& \eta & =\left\|\eta_{a b}\right\|=\operatorname{diag}(1,-1,-1,-1) \\
\operatorname{tr} \gamma_{a} & =0, & \gamma_{a}^{\dagger} & =\gamma_{0} \gamma_{a} \gamma_{0},
\end{array} r a l l, 2,3 .
$$

Using the substitution $e_{a+1} \rightarrow \gamma_{a}, \quad a=0,1,2,3$, we obtain the isomorphism

$$
\mathbb{C} \otimes C \ell_{1,3} \cong \operatorname{Mat}(4, \mathbb{C})
$$

## 2. Gradings and Conjugations

### 2.1. Gradings

Any Clifford algebra element $U \in C \ell_{p, q, r}$ has the form

$$
U=u e+\sum_{a} u_{a} e_{a}+\sum_{a<b} u_{a b} e_{a b}+\cdots+u_{1 \ldots n} e_{1 \ldots n}=\sum_{A} u_{A} e_{A}, \quad u_{A} \in \mathbb{R}
$$

where we denote arbitrary ordered multi-index by $A=a_{1} \ldots a_{k}$. Denote its length by $|A|=k$.

Definition 3. The subspace

$$
C \ell_{p, q, r}^{k}=\left\{\sum_{|A|=k} u_{A} e_{A}\right\}
$$

is called a subspace of grade $k$.
We have

$$
C \ell_{p, q, r}=\bigoplus_{k=0}^{n} C \ell_{p, q, r}^{k}, \quad \operatorname{dim} C \ell_{p, q, r}^{k}=C_{n}^{k}=\frac{n!}{k!(n-k)!}
$$

Let us consider projection operations onto subspaces of grade $k$

$$
U \in C \ell_{p, q, r} \rightarrow\langle U\rangle_{k} \in C \ell_{p, q, r}^{k}
$$

Note that for arbitrary element $U \in C \ell_{p, q, r}$ we have

$$
U=\sum_{k=0}^{n}\langle U\rangle_{k}
$$

The Clifford algebra $C \ell_{p, q, r}$ is a $Z_{2}$-graded algebra. It can be represented in the form of the direct sum of even and odd subspaces

$$
C \ell_{p, q, r}=C \ell_{p, q, r}^{(0)} \oplus C \ell_{p, q, r}^{(1)}
$$

where

$$
C \ell_{p, q, r}^{(0)}=\bigoplus_{k=0 \bmod 2} C \ell_{p, q, r}^{k}, \quad C \ell_{p, q, r}^{(1)}=\bigoplus_{k=1 \bmod 2} C \ell_{p, q, r}^{k}
$$

We have

$$
C \ell_{p, q, r}^{(i)} C \ell_{p, q, r}^{(j)} \subset C \ell_{p, q, r}^{(i+j) \bmod 2}, \quad i=0,1
$$

and

$$
\operatorname{dim} C \ell_{p, q, r}^{(0)}=\operatorname{dim} C \ell_{p, q, r}^{(1)}=2^{n-1}
$$

Note that $C \ell_{p, q, r}^{(0)}$ is a subalgebra of $C \ell_{p, q, r}$.

### 2.2. Center of Clifford Algebra

We have the following well-known theorem about the center of Clifford algebra $\operatorname{Cen}\left(C \ell_{p, q}\right):=\left\{U \in C \ell_{p, q} ; U V=V U\right.$ for all $\left.V \in C \ell_{p, q}\right\}$.

Theorem 1. The center of the Clifford algebra $C \ell_{p, q}$ is
$\operatorname{Cen}\left(C \ell_{p, q}\right)= \begin{cases}C \ell_{p, q}^{0}=\{u e ; u \in \mathbb{R}\}, & \text { if } n \text { is even } \\ C \ell_{p, q}^{0} \oplus C \ell_{p, q}^{n}=\left\{u e+u_{1 \ldots n} e_{1 \ldots n} ; u, u_{1 \ldots n} \in \mathbb{R}\right\}, & \text { if } n \text { is odd } .\end{cases}$

Proof: Let us represent element $U$ in the form

$$
U=U^{(0)}+U^{(1)}, \quad U^{(i)} \in C \ell_{p, q}^{(i)}, \quad i=0,1
$$

We have conditions $U V=V U$ for all $V \in C \ell_{p, q}$. We obtain

$$
U^{(i)} e_{k}=e_{k} U^{(i)}, \quad k=1, \ldots, n, \quad i=0,1
$$

We represent $U^{(0)}$ in the form $U^{(0)}=A^{(0)}+e_{1} B^{(1)}$, where $A^{(0)} \in C \ell_{p, q}^{(0)}$ and $B^{(1)} \in C \ell_{p, q}^{(1)}$ do not contain $e_{1}$. For $k=1$ we obtain

$$
\left(A^{(0)}+e_{1} B^{(1)}\right) e_{1}=e_{1}\left(A^{(0)}+e_{1} B^{(1)}\right)
$$

Using $A^{(0)} e_{1}=e_{1} A^{(0)}$ and $e_{1} B^{(1)} e_{1}=-e_{1} e_{1} B^{(1)}$, we obtain $B^{(1)}=0$. Acting similarly for $e_{2}, \ldots, e_{n}$, we obtain $U^{(0)}=u e$.
We represent $U^{(1)}$ in the form $U^{(1)}=A^{(1)}+e_{1} B^{(0)}$, where $A^{(1)} \in C \ell_{p, q}^{(1)}$ and $B^{(0)} \in C \ell_{p, q}^{(0)}$ do not contain $e_{1}$. For $k=1$ we obtain

$$
\left(A^{(1)}+e_{1} B^{(0)}\right) e_{1}=e_{1}\left(A^{(1)}+e_{1} B^{(0)}\right)
$$

Using $A^{(1)} e_{1}=-e_{1} A^{(1)}$ and $e_{1} B^{(0)} e_{1}=e_{1} e_{1} B^{(0)}$, we obtain $A^{(1)}=0$. Acting similarly for $e_{2}, \ldots, e_{k}$, we obtain $U^{(1)}=u_{1 \ldots n} e_{1 \ldots n}$ in the case of odd $n$ and $U^{(1)}=0$ in the case of even $n$.

### 2.3. Operations of Conjugations

Definition 4. The following operation (involution) in the Clifford algebra $C \ell_{p, q, r}$

$$
\widehat{U}:=\left.U\right|_{e_{a} \rightarrow-e_{a}}, \quad U \in C \ell_{p, q, r}
$$

is called grade involution or main involution.
It can be verified that

$$
\widehat{U}=\sum_{k=0}^{n}\langle\widehat{U}\rangle_{k}=\sum_{k=0}^{n}(-1)^{k}\langle U\rangle_{k}
$$

We have

$$
\widehat{\widehat{U}}=U, \quad \widehat{U V}=\widehat{U} \widehat{V}, \quad \lambda \widehat{U+\mu} V=\lambda \widehat{U}+\mu \widehat{V}, \quad U, V \in C \ell_{p, q}, \lambda, \mu \in \mathbb{R}
$$

Definition 5. The following operation (anti-involution) in the Clifford algebra $C \ell_{p, q, r}$

$$
\widetilde{U}:=\left.U\right|_{e_{a_{1} \ldots a_{k}} \rightarrow e_{a_{k}} \ldots e_{a_{1}}}, \quad U \in C \ell_{p, q, r}
$$

is called reversion.

We have

$$
\widetilde{U}=\sum_{k=0}^{n}\langle\widetilde{U}\rangle_{k}=\sum_{k=0}^{n}(-1)^{\frac{k(k-1)}{2}}\langle U\rangle_{k}
$$

and

$$
\widetilde{\widetilde{U}}=U, \quad \widetilde{U V}=\widetilde{V} \widetilde{U}, \quad \lambda \widetilde{U+\mu V}=\lambda \widetilde{U}+\mu \widetilde{V}, \quad U, V \in C l_{p, q}, \lambda, \mu \in \mathbb{R}
$$

Definition 6. A superposition of reversion and grade involution is called Clifford conjugation.

We do not use individual notation for Clifford conjugation and use notation $\widehat{\widetilde{U}}$. The operation of Clifford conjugation corresponds to the operation of complex conjugation of complex numbers in the case $C \ell_{0,1} \cong \mathbb{C}$ and quaternion conjugation in the case $C \ell_{0,2} \cong \mathbb{H}$.
We have

$$
\widehat{\widetilde{U}}=\sum_{k=0}^{n}\langle\widehat{\widetilde{U}}\rangle_{k}=\sum_{k=0}^{n}(-1)^{\frac{k(k+1)}{2}}\langle U\rangle_{k} .
$$

Definition 7. The following operation in the complexified Clifford algebra $\mathbb{C} \otimes$ $C \ell_{p, q}$

$$
\bar{U}:=\left.U\right|_{u_{a_{1} \ldots a_{k}} \rightarrow \bar{u}_{a_{1} \ldots a_{k}}}, \quad U \in \mathbb{C} \otimes C \ell_{p, q, r}
$$

where we take complex conjugation of complex numbers $u_{a_{1} \ldots a_{k}}$, is called complex conjugation in Clifford algebra.

We have
$\overline{\bar{U}}=U, \quad \overline{U V}=\overline{U V}, \quad \overline{\lambda U+\mu V}=\bar{\lambda} \bar{U}+\bar{\mu} \bar{V}, \quad U, V \in \mathbb{C} \otimes C \ell_{p, q}, \lambda, \mu \in \mathbb{C}$.
An important operation of Hermitian conjugation in $\mathbb{C} \otimes C \ell_{p, q}$ will be considered in Section 3.4.

### 2.4. Quaternion Types of Clifford Algebra Elements

The operation of grade involution uniquely determines two (even and odd) subspaces of the Clifford algebra:

$$
C \ell_{p, q, r}^{(j)}:=\bigoplus_{k=j \bmod 2} C \ell_{p, q, r}^{k}=\left\{U \in C \ell_{p, q, r} ; \widehat{U}=(-1)^{j} U\right\}, \quad j=0,1
$$

In a similar way, operations of grade involution and reversion uniquely determine the following four subspaces. This is symbolically shown in Table 1 and Table 2. Instead of $(*)$ and depending on the case, the "plus" or "minus" signs should be understood.

Definition 8. The following four subspaces of $C \ell_{p, q, r}$

$$
\overline{\boldsymbol{j}}:=\bigoplus_{k=j \bmod 4} C \ell_{p, q, r}^{k}=\left\{U \in C \ell_{p, q, r} ; \widehat{U}=(-1)^{j} U, \widetilde{U}=(-1)^{\frac{j(j-1)}{2}} U\right\}
$$

are called subspaces of quaternion types $j=0,1,2,3$.

Table 1. Subspaces of quaternion types in $C \ell_{p, q, r}$.

| $C \ell_{p, q, r}$ | $\overline{\mathbf{0}}$ | $\overline{\mathbf{1}}$ | $\overline{\mathbf{2}}$ | $\overline{\mathbf{3}}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\widehat{U}=(*) U$ | + | - | + | - |
| $\widetilde{U}=(*) U$ | + | + | - | - |

We have

$$
C \ell_{p, q, r}=\overline{\mathbf{0}} \oplus \overline{\mathbf{1}} \oplus \overline{\mathbf{2}} \oplus \overline{\mathbf{3}}, \quad C \ell_{p, q, r}^{(0)}=\overline{\mathbf{0}} \oplus \overline{\mathbf{2}}, \quad C \ell_{p, q, r}^{(1)}=\overline{\mathbf{1}} \oplus \overline{\mathbf{3}}
$$

Grade involution, reversion, and complex conjugation uniquely determine eight subspaces of the complexified Clifford algebra.
We have

$$
\mathbb{C} \otimes C \ell_{p, q, r}=\overline{\mathbf{0}} \oplus \overline{\mathbf{1}} \oplus \overline{\mathbf{2}} \oplus \overline{\mathbf{3}} \oplus \mathrm{i} \overline{\mathbf{0}} \oplus \mathrm{i} \overline{\mathbf{1}} \oplus \mathrm{i} \overline{\mathbf{2}} \oplus \mathrm{i} \overline{\mathbf{3}}
$$

Table 2. Subspaces of quaternion types in $\mathbb{C} \otimes C \ell_{p, q, r}$.

| $\mathbb{C} \otimes C \ell_{p, q, r}$ | $\overline{\mathbf{0}}$ | $\overline{\mathbf{1}}$ | $\overline{\mathbf{2}}$ | $\overline{\mathbf{3}}$ | $\mathrm{i} \overline{\mathbf{0}}$ | $\mathrm{i} \overline{\mathbf{1}}$ | $\mathrm{i} \overline{\mathbf{2}}$ | $\mathrm{i} \overline{\mathbf{3}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\widehat{U}=(*) U$ | + | - | + | - | + | - | + | - |
| $\widetilde{U}=(*) U$ | + | + | - | - | + | + | - | - |
| $\bar{U}=(*) U$ | + | + | + | + | - | - | - | - |

The subspaces of quaternion types have the following dimensions

$$
\begin{align*}
& \operatorname{dim} \overline{\mathbf{0}}=\sum_{k} C_{n}^{4 k}=2^{n-2}+2^{\frac{n-2}{2}} \cos \frac{\pi n}{4} \\
& \operatorname{dim} \overline{\mathbf{1}}=\sum_{k} C_{n}^{4 k+1}=2^{n-2}+2^{\frac{n-2}{2}} \sin \frac{\pi n}{4} \\
& \operatorname{dim} \overline{\mathbf{2}}=\sum_{k} C_{n}^{4 k+2}=2^{n-2}-2^{\frac{n-2}{2}} \cos \frac{\pi n}{4}  \tag{3}\\
& \operatorname{dim} \overline{\mathbf{3}}=\sum_{k} C_{n}^{4 k+3}=2^{n-2}-2^{\frac{n-2}{2}} \sin \frac{\pi n}{4} .
\end{align*}
$$

We denote the commutator of two Clifford algebra elements $U, V$ by $[U, V]:=$ $U V-V U$ and the anticommutator by $\{U, V\}:=U V+V U$.

Theorem 2 ([48], [56]). We have the following properties

$$
\left.\begin{array}{rlrl}
{[\overline{\mathbf{j}, \mathbf{j}}]} & \subset \overline{\mathbf{2}}, & {[\overline{\mathbf{j}}, \overline{\mathbf{2}}]} & \subset \overline{\mathbf{j}}, \\
{[\overline{\mathbf{0}}, \overline{\mathbf{1}}]} & \subset \overline{\mathbf{3}}, & {[\overline{\mathbf{0}}, \overline{\mathbf{3}}]} & \subset \overline{\mathbf{1}}, \\
\{\overline{\mathbf{j}}, \overline{\mathbf{j}}\} & \subset \overline{\mathbf{0}}, & \{\overline{\mathbf{j}}, \overline{\mathbf{j}}] & \subset \overline{\mathbf{j}}, \\
\{\overline{\mathbf{1}}, \overline{\mathbf{2}}\} & \subset \overline{\mathbf{3}}, & \{\overline{\mathbf{2}}, \overline{\mathbf{3}}\} & \subset \overline{\mathbf{1}},
\end{array}\right)\{\overline{\mathbf{3}}, \overline{\mathbf{1}}\} \subset \overline{\mathbf{2}} .
$$

By Theorem 2, the Clifford algebra $C \ell_{p, q, r}$ is a $\mathbb{Z}_{2} \times \mathbb{Z}_{2}$-graded algebra w.r.t. the operation of commutator and w.r.t. the operation of anticommutator.
The notion of quaternion type was introduced by the author in the brief report [48] and the paper [56]. Further development of this concept is given in [52], [57], [61], [66], see also books [35], [55].
Subspaces of quaternion types are useful in different calculations (see [52], [57], [61]). Here and below we omit the sign of the direct sum to simplify notation: $\overline{\mathbf{0}} \oplus \overline{\mathbf{1}}=\overline{\mathbf{0 1}}, \overline{\mathbf{0}} \oplus \overline{\mathbf{1}} \oplus \overline{\mathbf{2}} \oplus \overline{\mathbf{3}}=\overline{\mathbf{0 1 2 3}}=C \ell_{p, q}$, etc.
For example, if $U \in \overline{\mathbf{k}}$ for some $k=0,1,2,3$, then (see [57])

$$
U^{m} \in\left\{\begin{array}{l}
\overline{\mathbf{k}}, \text { if } m \text { is odd } \\
\overline{\mathbf{0}}, \text { if } m \text { is even, }
\end{array} \quad \sin U \in \overline{\mathbf{k}}, \quad \cos U \in \overline{\mathbf{0}}\right.
$$

For arbitrary element $U \in C \ell_{p, q}$ we have (see [52])

$$
U \widetilde{U}, \widetilde{U} U \in \overline{\mathbf{0 1}}, \quad U \widehat{\tilde{U}}, \widehat{\widetilde{U}} U \in \overline{\mathbf{0 3}} .
$$

Using the classification of Clifford algebra elements based on the notion of quaternion type, we present a number of Lie algebras in $\mathbb{C} \otimes C \ell_{p, q}$ (see Section 5.4 and [66]).

## 3. Matrix Representations of Clifford Algebras

### 3.1. Cartan's Periodicity of 8. Central and Simple Algebras

Lemma 3 ([31]). We have the following isomorphisms of associative algebras:

1) $C \ell_{p+1, q+1} \cong \operatorname{Mat}\left(2, C \ell_{p, q}\right)$,
2) $C \ell_{p+1, q+1} \cong C \ell_{p, q} \otimes C \ell_{1,1}$
3) $C \ell_{p, q} \cong C \ell_{q+1, p-1}, \quad p \geq 1$,
4) $C \ell_{p, q} \cong C \ell_{p-4, q+4}, \quad p \geq 4$.

Proof: Let $e_{1}, \ldots, e_{n}$ be the generators of $C \ell_{p, q}$ and $\left(e_{+}\right)^{2}=e,\left(e_{-}\right)^{2}=-e$ such that all generators $e_{1}, \ldots, e_{n}, e_{+}, e_{-}$anticommute with each other.

1. We obtain generators of $\operatorname{Mat}\left(2, C \ell_{p, q}\right)$ in the following way

$$
e_{i} \rightarrow\left(\begin{array}{cc}
e_{i} & 0 \\
0 & -e_{i}
\end{array}\right), \quad i=1, \ldots, n, \quad e_{+} \rightarrow\left(\begin{array}{cc}
0 & e \\
e & 0
\end{array}\right), \quad e_{-} \rightarrow\left(\begin{array}{rr}
0 & -e \\
e & 0
\end{array}\right) .
$$

2. $e_{i} e_{+} e_{-}, i=1, \ldots, n$ are generators of $C \ell_{p, q}$ and $e_{+}, e_{-}$are generators of $C \ell_{1,1}$. Each generator of $C \ell_{p, q}$ commutes with each generator of $C \ell_{1,1}$.
3. $e_{1}, e_{i} e_{1}, i=2, \ldots, n$ are generators of $C \ell_{q+1, p-1}$.
4. $e_{i} e_{1} e_{2} e_{3} e_{4}, i=1,2,3,4$ and $e_{j}, j=5, \ldots, n$ are generators of $C \ell_{p-4, q+4}$.

We have the following well-known theorems about isomorphisms between Clifford algebras and matrix algebras.

Theorem 4 (Cartan 1908). We have the following isomorphism of algebras
$C \ell_{p, q} \cong \begin{cases}\operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{R}\right), & \text { if } p-q \equiv 0 ; 2 \bmod 8 \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{R}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{R}\right), & \text { if } p-q \equiv 1 \bmod 8 \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), & \text { if } p-q \equiv 3 ; 7 \bmod 8 \\ \operatorname{Mat}\left(2^{\frac{n-2}{2}}, \mathbb{H}\right), & \text { if } p-q \equiv 4 ; 6 \bmod 8 \\ \operatorname{Mat}\left(2^{\frac{n-3}{2}}, \mathbb{H}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-3}{2}}, \mathbb{H}\right), & \text { if } p-q \equiv 5 \bmod 8 .\end{cases}$
Proof: Using Lemma 3, we obtain isomorphisms for all $C \ell_{p, q}$ (see Table 3). We use notations ${ }^{2} \mathbb{R}:=\mathbb{R} \oplus \mathbb{R}, \quad \mathbb{R}(2):=\operatorname{Mat}(2, \mathbb{R}), \ldots$
We have the isomorphisms (see Section 1.3)

$$
C \ell_{0,0} \cong \mathbb{R}, \quad C \ell_{0,1} \cong \mathbb{C}, \quad C \ell_{1,0} \cong \mathbb{R} \oplus \mathbb{R}, \quad C \ell_{0,2} \cong \mathbb{H}
$$

Using the substitution

$$
e \rightarrow(1,1), \quad e_{1} \rightarrow(\mathrm{i},-\mathrm{i}), \quad e_{2} \rightarrow(\mathrm{j},-\mathrm{j}), \quad e_{3} \rightarrow(\mathrm{k},-\mathrm{k})
$$

we obtain also the isomorphism

$$
C l_{0,3} \cong \mathbb{H} \oplus \mathbb{H} .
$$

Using $C \ell_{p+1, q+1} \cong \operatorname{Mat}\left(2, C \ell_{p, q}\right)$, we get $C \ell_{1,1} \cong \operatorname{Mat}(2, \mathbb{R})$. Using $C \ell_{p+1, q+1} \cong$ $C \ell_{p, q} \otimes C \ell_{1,1}$, we conclude that if we make a step down Table $3(n \rightarrow n+2)$, then the size of corresponding matrix algebra is doubled $(\operatorname{Mat}(k, \ldots) \rightarrow \operatorname{Mat}(2 k, \ldots))$. Using $C \ell_{p, q} \cong C \ell_{q+1, p-1}$, we conclude that Table 3 is symmetric w.r.t. the column " $p-q=1$ ". Using $C \ell_{p, q} \cong C \ell_{p-4, q+4}$, we obtain the symmetry $p-q \leftrightarrow p-q-8$ for each $n$.

Similarly, we can obtain the following isomorphisms for complex Clifford algebras and for even subalgebras of the Clifford algebra.

Theorem 5 ([31]). We have the following isomorphism of algebras $C \ell\left(\mathbb{C}^{n}\right) \cong \mathbb{C} \otimes C \ell_{p, q} \cong \begin{cases}\operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{C}\right), & \text { if } n \text { is even } \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), & \text { if } n \text { is odd } .\end{cases}$

Table 3. Isomorphisms between $C \ell_{p, q}$ and matrix algebras.

| $n \backslash p-q$ | -5 | -4 | -3 | -2 | -1 | 0 | 1 | 2 | 3 | 4 | 5 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | - | - | - | - | - | $\mathbb{R}$ | - | - | - | - | - |
| 1 | - | - | - | - | $\mathbb{C}$ | - | ${ }^{2} \mathbb{R}$ | - | - | - | - |
| 2 | - | - | - | $\mathbb{H}$ | - | $\mathbb{R}(2)$ | - | $\mathbb{R}(2)$ | - | - | - |
| 3 | - | - | ${ }^{2} \mathbb{H}$ | - | $\mathbb{C}(2)-$ | ${ }^{2} \mathbb{R}(2)-$ | $\mathbb{C}(2)$ | - | - |  |  |
| 4 | - | $\mathbb{H}(2)$ | - | $\mathbb{H}(2)$ | - | $\mathbb{R}(4)$ | - | $\mathbb{R}(4)-$ | $\mathbb{H}(2)$ | - |  |
| 5 | $\mathbb{C}(4)-$ | ${ }^{2} \mathbb{H}(2)-$ | $\mathbb{C}(4)-$ | ${ }^{2} \mathbb{R}(4)-$ | $\mathbb{C}(4)$ | - | ${ }^{2} \mathbb{H}(2)$ |  |  |  |  |

Theorem 6 ([31]). We have the following isomorphism of algebras

1) $C \ell_{p, q}^{(0)} \cong C \ell_{p, q-1}, \quad q \geq 1$;
2) $C \ell_{p, q}^{(0)} \cong C \ell_{q, p-1}, \quad p \geq 1 ;$
3) $C \ell_{p, q}^{(0)} \cong C \ell_{q, p}^{(0)}$.

Proof: Let $e_{1}, \ldots, e_{n}$ be the generators of $C \ell_{p, q}$.

1. Then $e_{i} e_{n}, i=1, \ldots, n-1$ are generators of $C e_{p, q}^{(0)}$.
2. Then $\left\{\begin{array}{l}e_{p+i} e_{p}, i=1, \ldots, q \\ e_{j-q} e_{p}, j=q+1, \ldots, n-1\end{array}\right.$ are generators of $C \ell_{p, q}^{(0)}$.
3. Using 1) and 2 ), we get 3 ).

Definition 9. An algebra is simple if it contains no non-trivial two-sided ideals and the multiplication operation is not zero.

Definition 10. A central simple algebra over a field $\mathbb{F}$ is a finite-dimensional associative algebra, which is simple, and for which the center is exactly $\mathbb{F}$.

The following classification of Clifford algebras can be found in [10].

- If $n$ is even, then $C \ell(V, Q)$ is a central simple algebra.
- If $n$ is odd and $\mathbb{F}=\mathbb{C}$, then $C \ell(V, Q)$ is the direct sum of two isomorphic complex central simple algebras.
- If $n$ is odd, $\mathbb{F}=\mathbb{R}$, and $\left(e_{1 \ldots n}\right)^{2}=e$, then $C \ell(V, Q)$ is the direct sum of two isomorphic simple algebras.
- If $n$ is odd, $\mathbb{F}=\mathbb{R}$, and $\left(e_{1 \ldots n}\right)^{2}=-e$, then $C \ell(V, Q)$ is simple with center $\cong \mathbb{C}$.

Note that

$$
\left(e_{1 \ldots n}\right)^{2}=(-1)^{q+\frac{n(n-1)}{2}} e=\left\{\begin{array}{ll}
e, & \text { if } p-q=0,1 \quad \bmod 4 \\
-e, & \text { if } p-q=2,3
\end{array} \bmod 4\right.
$$

and these results agree with Theorem 4.

### 3.2. Clifford Trigonometry Circle and Exterior Signature of Clifford Algebra

In the literature, the Cartan's periodicity of 8 is depicted in the form of a eighthour clock (see Table 4). The clockwise movement by one step corresponds to an increase of $p-q$ by 1 . The clock shows that two Clifford algebra with the same $p-q \bmod 8$ are Morita equivalent (see Theorem 4).

Table 4. Clifford clock or Clifford trigonometry circle.

|  |  |  | 0 |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 7 |  | $\mathbb{R}$ |  | 1 |  |
|  |  | ${ }^{2} \mathbb{R}$ |  | $\mathbb{C}$ |  |  |
| 6 | $\mathbb{R}$ |  |  |  | $\mathbb{H}$ | 2 |
|  |  | $\mathbb{C}$ |  |  |  |  |
|  |  |  |  |  |  |  |
|  | 5 |  | $\mathbb{H}$ |  |  |  |
|  |  |  | 4 |  | 3 |  |

In our opinion, it is more correct to call it not Clifford clock, but Clifford trigonometry circle. The algebras on the clock are in one-to-one correspondence with the values of the function $\sin \frac{\pi(p-q+1)}{4}$. To show this, we do the following calculations (see also [35]).
Let $P$ be the number of basis elements $e_{A}$ of the Clifford algebra $C \ell_{p, q}$ such that $\left(e_{A}\right)^{2}=e$ and $Q$ be the number of basis elements $e_{A}$ of $C \ell_{p, q}$ such that $\left(e_{A}\right)^{2}=$ $-e$. We call $(P, Q)$ exterior signature of Clifford algebra $C \ell_{p, q}, P+Q=2^{n}$.

Theorem 7. We have

$$
P=2^{\frac{n-1}{2}}\left(2^{\frac{n-1}{2}}+\sin \frac{\pi(p-q+1)}{4}\right), \quad Q=2^{\frac{n-1}{2}}\left(2^{\frac{n-1}{2}}-\sin \frac{\pi(p-q+1)}{4}\right)
$$

Note that

$$
\sum_{A}\left(e_{A}\right)^{2}=P-Q=2^{\frac{n+1}{2}} \sin \frac{\pi(p-q+1)}{4}
$$

Finally, we obtain the following theorem.
Theorem 8. Two Clifford algebras $C \ell_{p_{1}, q_{1}}$ and $C \ell_{p_{2}, q_{2}}$ are isomorphic if and only if their exterior signatures coincide $\left(P_{1}, Q_{1}\right)=\left(P_{2}, Q_{2}\right)$ (or, equivalently, $P_{1}-$ $Q_{1}=P_{2}-Q_{2}$.

### 3.3. Trace, Determinant and Inverse of Clifford Algebra Elements

Definition 11. The following projection operation onto subspace $\mathbb{C} \otimes C \ell_{p, q}^{0}$

$$
\operatorname{Tr}(U):=\left.\langle U\rangle\right|_{e \rightarrow 1}
$$

is called a trace of Clifford algebra element.
We have

$$
\operatorname{Tr}(U)=u, \quad U=u e+\sum_{a} u_{a} e_{a}+\cdots+u_{1 \ldots n} e_{1 \ldots n}
$$

Theorem 9 ([54], [35]). The operation of trace of Clifford algebra elements has the following properties

$$
\begin{aligned}
& \operatorname{Tr}(U+V)=\operatorname{Tr}(U)+\operatorname{Tr}(V), \quad \operatorname{Tr}(\lambda U)=\lambda \operatorname{Tr}(U), \quad \operatorname{Tr}(U V)=\operatorname{Tr}(V U) \\
& \operatorname{Tr}(U V W)=\operatorname{Tr}(V W U)=\operatorname{Tr}(W U V), \quad U, V, W \in \mathbb{C} \otimes C \ell_{p, q}, \lambda \in \mathbb{C} \\
& \operatorname{Tr}\left(U^{-1} V U\right)=\operatorname{Tr}(V), \quad \operatorname{Tr}(U)=\operatorname{Tr}(\hat{U})=\operatorname{Tr}(\tilde{U})=\overline{\operatorname{Tr} \bar{U}}
\end{aligned}
$$

We have the following relation with the trace of matrices.
Theorem 10 ([54], [35]). We have

$$
\operatorname{Tr}(U)=\frac{1}{2^{\left[\frac{n+1}{2}\right]}} \operatorname{tr}(\gamma(U))
$$

where

$$
\gamma: \mathbb{C} \otimes C \ell_{p, q} \rightarrow \begin{cases}\operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{C}\right), & \text { if } n \text { is even } \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), & \text { if } n \text { is odd }\end{cases}
$$

is faithful matrix representation of $\mathbb{C} \otimes C \ell_{p, q}$ (of minimal dimension).
Definition 12. The determinant of any faithful matrix representation (of minimal dimension) of element $U$ is called the determinant of Clifford algebra element $U \in$ $\mathbb{C} \otimes C \ell_{p, q}$.

Theorem 11 ([54], [35]). Definition 12 is correct: The determinant does not depend on the choice of matrix representation.

Theorem 12 ([54], [35]). The determinant of Clifford algebra elements has the following properties

- $\operatorname{Det}(U V)=\operatorname{Det}(U) \operatorname{Det}(V), \quad \operatorname{Det}(\lambda U)=\lambda^{2^{\left[\frac{n+1}{2}\right]}} \operatorname{Det}(U)$ $\operatorname{Det}(U)=\operatorname{Det}(\widehat{U})=\operatorname{Det}(\widetilde{U})=\overline{\operatorname{Det}(\bar{U})}, \quad U, V \in \mathbb{C} \otimes C \ell_{p, q}, \lambda \in \mathbb{C}$.
- $U^{-1} \in \mathbb{C} \otimes C \ell_{p, q}$ exists if and only if $\operatorname{Det} U \neq 0$.
- If $U^{-1}$ exists, then
$\operatorname{Det}\left(U^{-1}\right)=(\operatorname{Det} U)^{-1}, \quad \operatorname{Det}\left(U^{-1} V U\right)=\operatorname{Det}(V), \quad V \in \mathbb{C} \otimes C \ell_{p, q}$.

Theorem 13 ([54], [35]). We have the following explicit formulas for the determinant and the inverse of Clifford algebra element $U \in \mathbb{C} \otimes C \ell_{p, q}$ in the cases $n=1, \ldots, 5$

$$
\begin{aligned}
& \operatorname{Det} U= \begin{cases}\left.U\right|_{e \rightarrow 1}, & n=0 \\
\left.U \widehat{U}\right|_{e \rightarrow 1}, & n=1 \\
\left.U \widetilde{\widetilde{U}}\right|_{e \rightarrow 1}, & n=2 \\
\left.U \widetilde{U} \widehat{U} \widehat{\tilde{U}}\right|_{e \rightarrow 1}=\left.U \widehat{\tilde{U}} \widehat{U} \widetilde{U}\right|_{e \rightarrow 1}, & n=3 \\
\left.U \widetilde{U}(\widehat{U} \widehat{\widetilde{U}}) \nabla\right|_{e \rightarrow 1}=\left.U \widehat{\tilde{U}}(\widehat{U} \widetilde{U}) \nabla\right|_{e \rightarrow 1}, & n=4 \\
\left.U \widetilde{U}(\widehat{U} \widehat{\tilde{U}}) \nabla(U \widetilde{U}(\widehat{U} \widehat{\tilde{U}}) \nabla)^{\triangle}\right|_{e \rightarrow 1}, & n=5\end{cases}
\end{aligned}
$$

where $U \nabla=\left.U\right|_{\langle U\rangle_{4} \rightarrow-\langle U\rangle_{4},\langle U\rangle_{5} \rightarrow-\langle U\rangle_{5}}$ and $U^{\Delta}=\left.U\right|_{\langle U\rangle_{5} \rightarrow-\langle U\rangle_{5}}$.
Note that we can introduce the notions of trace and determinant of elements of the real Clifford algebra $C \ell_{p, q}$. These operations have similar properties (see [54], [35]).

### 3.4. Unitary Space on Clifford Algebra

Theorem 14 ([33], [35]). The operation $U, V \in \mathbb{C} \otimes C \ell_{n} \rightarrow(U, V):=\operatorname{Tr}(\overline{\tilde{U}} V)$ is a Hermitian (or Euclidian) scalar product on $\mathbb{C} \otimes C \ell_{n}$ (or $C \ell_{n}$ respectively).

Proof: We must verify

$$
\begin{align*}
& (U, V)=\overline{(V, U)}, \quad(U, \lambda V)=\lambda(U, V), \quad(U, V+W)=(U, V)+(U, W) \\
& (U, U) \geq 0, \quad(U, U)=0 \Leftrightarrow U=0 \tag{4}
\end{align*}
$$

for all $U, V, W \in \mathbb{C} \otimes C \ell_{p, q}, \lambda \in \mathbb{C}$. To prove (4) it is sufficient to prove that the basis of $\mathbb{C} \otimes C \ell_{p, q}$ is orthonormal

$$
\left(e_{i_{1} \ldots i_{k}}, e_{j_{1} \ldots j_{l}}\right)=\operatorname{Tr}\left(e_{i_{k}} \cdots e_{i_{1}} e_{j_{1}} \cdots e_{j_{l}}\right)=\left\{\begin{array}{l}
1, \text { if }\left(i_{1}, \ldots, i_{k}\right)=\left(j_{1}, \ldots, j_{l}\right) \\
0, \text { if }\left(i_{1}, \ldots, i_{k}\right) \neq\left(j_{1}, \ldots, j_{l}\right)
\end{array}\right.
$$

We have

$$
(U, U)=\sum_{A} u_{A} \overline{u_{A}}=\sum_{A}\left|u_{A}\right|^{2} \geq 0
$$

The theorem is proved.
Definition 13. Let us consider the following operation of Hermitian conjugation in Clifford algebra

$$
U^{\dagger}:=\left.U\right|_{e_{a_{1} \ldots a_{k}} \rightarrow e_{a_{1} \ldots a_{k}}^{-1}, u_{a_{1} \ldots a_{k}} \rightarrow \bar{u}_{a_{1} \ldots a_{k}}}, \quad U \in \mathbb{C} \otimes C \ell_{p, q}
$$

This operation has the following properties

$$
\begin{aligned}
& U^{\dagger \dagger}=U, \quad(U V)^{\dagger}=V^{\dagger} U^{\dagger}, \quad(\lambda U+\mu V)^{\dagger}=\bar{\lambda} U^{\dagger}+\bar{\mu} V^{\dagger} \\
& U, V \in \mathbb{C} \otimes C \ell_{p, q}, \quad \lambda, \mu \in \mathbb{C} .
\end{aligned}
$$

Theorem 15 ([33], [35]). The operation $U, V \in \mathbb{C} \otimes C \ell_{n} \rightarrow(U, V):=\operatorname{Tr}\left(U^{\dagger} V\right)$ is a Hermitian (or Euclidian) scalar product on $\mathbb{C} \otimes C \ell_{p, q}$ (or $C \ell_{p, q}$ respectively).

Proof: The proof is similar to the proof of the previous theorem. Now we have $\left(e_{i_{1}} \cdots e_{i_{k}}, e_{i_{1}} \cdots e_{i_{k}}\right)=\operatorname{Tr}\left(e_{i_{k}}^{-1} \ldots e_{i_{1}}^{-1} e_{i_{1}} \cdots e_{i_{k}}\right)=\operatorname{Tr}(e)=1$.

Note that the Hermitian conjugation in the case of the real Clifford algebra $C \ell_{p, q}$ is called transposition anti-involution. It is considered in [2], [3], [4] in more details. We have the following relation between the Hermitian conjugation of Clifford algebra elements and the Hermitian conjugation of matrices.
Theorem 16 ([33], [35]). We have $\gamma\left(U^{\dagger}\right)=(\gamma(U))^{\dagger}$, where

$$
\gamma: \mathbb{C} \otimes C \ell_{p, q} \rightarrow \begin{cases}\operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{C}\right), & \text { if } n \text { is even } \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), & \text { if } n \text { is odd }\end{cases}
$$

is faithful matrix representation of $\mathbb{C} \otimes C \ell_{p, q}$ such that $\left(\gamma\left(e_{a}\right)\right)^{-1}=\left(\gamma\left(e_{a}\right)\right)^{\dagger}$.
Let us consider the following Lie group in Clifford algebra
$\mathrm{UC} \ell_{p, q}:=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; U^{\dagger} U=e\right\} \cong \begin{cases}\mathrm{U}\left(2^{\frac{n}{2}}\right), & \text { if } n \text { is even } \\ \mathrm{U}\left(2^{\frac{n-1}{2}}\right) \oplus \mathrm{U}\left(2^{\frac{n-1}{2}}\right), & \text { if } n \text { is odd. }\end{cases}$
We call it unitary group in Clifford algebra [33], [35]. All basis elements of Clifford algebra lie in this group by definition $e_{a_{1} \ldots a_{k}} \in \mathrm{U} C \ell_{p, q}$. The corresponding Lie algebra is

$$
\mathfrak{u C \ell _ { p , q }}:=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; U^{\dagger}=-U\right\}
$$

Theorem 17 ([33], [35]). We have the following formulas which can be considered as another (equivalent to Definition 13) definition of Hermitian conjugation
$U^{\dagger}=\left\{\begin{array}{l}\left(e_{1 \ldots p}\right)^{-1} \overline{\widetilde{U}} e_{1 \ldots p}, \text { if } p \text { is odd } \\ \left(e_{1 \ldots p}\right)^{-1} \overline{\widetilde{U}} e_{1 \ldots p}, \text { if } p \text { is even }\end{array} \quad U^{\dagger}=\left\{\begin{array}{l}\left(e_{p+1 \ldots n}\right)^{-1} \overline{\widetilde{U}} e_{p+1 \ldots n}, \text { if } q \text { is even } \\ \left(e_{p+1 \ldots n}\right)^{-1} \overline{\widetilde{U}} e_{p+1 \ldots n}, \text { if } q \text { is odd } .\end{array}\right.\right.$

As an example, we obtain well-known relations $\gamma_{a}^{\dagger}=\gamma_{0} \gamma_{a} \gamma_{0}$ for the Dirac gammamatrices.

Proof: Because of the linearity of the operation ${ }^{\dagger}$ it is sufficient to prove the following formulas

$$
e_{i_{1} \ldots i_{k}}^{\dagger}=(-1)^{(p+1) k} e_{1 \ldots p}^{-1} \widetilde{e_{i_{1} \ldots i_{k}}} e_{1 \ldots p}, \quad e_{i_{1} \ldots i_{k}}^{\dagger}=(-1)^{q k} e_{p+1 \ldots n}^{-1} \widetilde{e_{i_{1} \ldots i_{k}}} e_{p+1 \ldots n}
$$

Let $s$ be the number of common indices in $\left\{i_{1}, \ldots, i_{k}\right\}$ and $\{1, \ldots, p\}$. Then

$$
\begin{aligned}
& (-1)^{(p+1) k} e_{1 \ldots p}^{-1} \widetilde{e_{i_{1} \ldots i_{k}}} e_{1 \ldots p}=(-1)^{(p+1) k} e_{p} \cdots e_{1} e_{i_{k}} \cdots e_{i_{1}} e_{1} \cdots e_{p} \\
& \quad=(-1)^{(p+1) k}(-1)^{k p-s} e_{i_{k}} \cdots e_{i_{1}}=(-1)^{k-s} e_{i_{k}} \cdots e_{i_{1}}=e_{i_{1} \ldots i_{k}}^{-1} \\
& (-1)^{q k} e_{p+1 \ldots n}^{-1} \widetilde{e_{i_{1} \ldots i_{k}}} e_{p+1 \ldots n}=(-1)^{q k}(-1)^{q} e_{n} \cdots e_{p+1} e_{i_{k}} \cdots e_{i_{1}} e_{p+1} \cdots e_{n} \\
& \quad=(-1)^{q k+q}(-1)^{q k-(k-s)}(-1)^{q} e_{i_{k}} \cdots e_{i_{1}}=(-1)^{k-s} e_{i_{k}} \cdots e_{i_{1}}=e_{i_{1} \ldots i_{k}}^{-1} .
\end{aligned}
$$

The theorem is proved.

### 3.5. Primitive Idempotents and Minimal Left Ideals

Definition 14. The element $t \in \mathbb{C} \otimes C \ell_{p, q}, t^{2}=t, t^{\dagger}=t$ is called a Hermitian idempotent. The subset $\mathrm{I}(t)=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; U=U t\right\}$ is called the left ideal generated by $t$.

Definition 15. A left ideal that does not contain other left ideals except itself and the trivial ideal (generated by $t=0$ ) is called a minimal left ideal. The corresponding idempotent is called primitive.

Note that if $V \in \mathrm{I}(t)$ and $U \in \mathbb{C} \otimes C \ell_{p, q}$, then $U V \in \mathrm{I}(t)$.
The left ideal $\mathrm{I}(t)$ is a complex vector space with the orthonormal basis $\tau_{1}, \ldots, \tau_{d}$, $d:=\operatorname{dim} \mathrm{I}(t)$. We have the Hermitian scalar product $(U, V)=\operatorname{Tr}\left(U^{\dagger} V\right)$ on $\mathrm{I}(t), \tau_{k}=\tau^{k},\left(\tau_{k}, \tau^{l}\right)=\delta_{k}^{l}, k, l=1, \ldots, n$. We may define the linear map $\gamma: \mathbb{C} \otimes C \ell_{p, q} \rightarrow \operatorname{Mat}(d, \mathbb{C})$

$$
\begin{equation*}
U \tau_{k}=\gamma(U)_{k}^{l} \tau_{l}, \quad \gamma(U)=\left\|\gamma(U)_{k}^{l}\right\| \in \operatorname{Mat}(d, \mathbb{C}) \tag{5}
\end{equation*}
$$

We have $\gamma(U)_{l}^{k}=\left(\tau^{k}, U \tau_{l}\right)$.
Lemma 18. The linear map $\gamma$ is a representation of Clifford algebra of the dimension $d: \gamma(U V)=\gamma(U) \gamma(V)$.

Proof: $\gamma(U V)_{k}^{m} \tau_{m}=(U V) \tau_{k}=U\left(V \tau_{k}\right)=U \tau_{l} \gamma(V)_{k}^{l}=\gamma(U)_{l}^{m} \gamma(V)_{k}^{l} \tau_{m}$.
Lemma 19. We have $\gamma\left(U^{\dagger}\right)=(\gamma(U))^{\dagger}$.

Proof: Using $(A, U B)=\left(A U^{\dagger}, B\right)$ and $(A, B)=\overline{(B, A)}$ for $(A, B)=\operatorname{Tr}\left(A^{\dagger} B\right)$, we obtain $\gamma(U)_{l}^{k}=\left(U^{\dagger} \tau^{k}, \tau_{l}\right), \overline{\gamma(U)}_{l}^{k}=\left(\tau_{l}, U^{\dagger} \tau^{k}\right)$. Transposing, we get $\left(\gamma(U)_{l}^{k}\right)^{\dagger}$ $=\left(\tau^{k}, U^{\dagger} \tau_{l}\right)$, which coincides with $\gamma\left(U^{\dagger}\right)_{l}^{k}=\left(\tau^{k}, U^{\dagger} \tau_{l}\right)$.

Theorem 20 ([33], [35]). The following elements are primitive idempotents in $\mathbb{C} \otimes C \ell_{p, q}$

$$
\begin{aligned}
& t=\frac{1}{2}\left(e+i^{a} e_{1}\right) \prod_{k=1}^{[n / 2]-1} \frac{1}{2}\left(e+i^{b_{k}} e_{2 k} e_{2 k+1}\right) \in \mathbb{C} \otimes C \ell_{p, q}, \quad t^{2}=t^{\dagger}=t \\
& a=\left\{\begin{array}{l}
0, \text { if } p \neq 0 \\
1, \text { if } p=0
\end{array} \quad b_{k}=\left\{\begin{array}{l}
0,2 k=p \\
1,2 k \neq p
\end{array}\right.\right.
\end{aligned}
$$

Further, we choose a basis of the corresponding minimal left ideal $\mathrm{I}(t)$ (for mote details, see [33], [35]) and obtain the representation of the complexified Clifford algebra (5).
For the real Clifford algebras $C \ell_{p, q}$ the construction is similar, see [1]. Using the idempotent and the basis of the left ideal, we can construct representations of the real Clifford algebra.

## 4. Method of Averaging in Clifford Algebras

### 4.1. Averaging in Clifford Algebras

The method of averaging in Clifford algebras is related to the method of averaging in the representation theory of finite groups [47], [17], [6]. We present a number of theorems which one can find in [65], [61], [35], [33], [37].
Let us consider the Reynolds operator [12] of the Salingaros group [44], [45], [46] $\mathrm{G}_{p, q}:=\left\{ \pm e_{A}\right\}$

$$
\mathrm{F}(U)=\frac{1}{\left|\mathrm{G}_{p, q}\right|} \sum_{g \in \mathrm{G}_{p, q}} g^{-1} U g=\frac{1}{2^{n}} \sum_{A}\left(e_{A}\right)^{-1} U e_{A}, \quad U \in C \ell_{p, q}
$$

Theorem 21 ([65]). The operator $\mathrm{F}(U)$ is the projection onto the center of Clifford algebra $\mathrm{Cl}_{p, q}$

$$
\mathrm{F}(U)=\frac{1}{2^{n}} \sum_{A} e_{A}^{-1} U e_{A}=\left\{\begin{array}{ll}
\langle U\rangle_{0}, & \text { if } n \text { is even } \\
\langle U\rangle_{0}+\langle U\rangle_{n}, & \text { if } n \text { is odd }
\end{array} \quad \mathrm{F}^{2}=\mathrm{F}\right.
$$

Let us consider other operators acting on $U \in C \ell_{p, q}$ that are also related to projection operators.

Theorem 22 ([61]). We have

$$
\begin{aligned}
& \mathrm{F}_{\text {Even }}(U):=\sum_{A:|A|=0 \bmod 2} e_{A}^{-1} U e_{A}=\langle U\rangle_{0}+\langle U\rangle_{n}, \quad \mathrm{~F}_{\text {Even }}^{2}=\mathrm{F}_{\text {Even }} \\
& \mathrm{F}_{\text {Odd }}(U):=\sum_{A:|A|=1 \bmod 2} e_{A}^{-1} U e_{A}=\langle U\rangle_{0}+(-1)^{n+1}\langle U\rangle_{n}, \quad \mathrm{~F}_{\text {Odd }}^{2}=\mathrm{F}_{\text {Odd }} .
\end{aligned}
$$

We have $\mathrm{F}=\frac{1}{2}\left(\mathrm{~F}_{\mathrm{Even}}+\mathrm{F}_{\mathrm{Odd}}\right)$ in the case of even $n$ and $\mathrm{F}=\mathrm{F}_{\mathrm{Even}}=\mathrm{F}_{\mathrm{Odd}}$ in the case of odd $n$.

Theorem 23 ([61]). For $m=0,1, \ldots, n$ we have

$$
\mathrm{F}_{m}(U):=\sum_{A:|A|=m} e_{A}^{-1} U e_{A}=\sum_{k=0}^{n}(-1)^{k m}\left(\sum_{i=0}^{m}(-1)^{i} C_{k}^{i} C_{n-k}^{m-i}\right)\langle U\rangle_{k}
$$

## Particulary

$$
\mathrm{F}_{1}(U):=\sum e_{a}^{-1} U e_{a}=\sum_{k=0}^{n}(-1)^{k}(n-2 k)\langle U\rangle_{k}
$$

Theorem 24 ([37]). Let us consider the operator $\mathrm{F}_{1}$ from the previous theorem which acts several times: $\mathrm{F}_{1}^{l}(U)=\underbrace{\mathrm{F}_{1}\left(\mathrm{~F}_{1}\left(\cdots \mathrm{~F}_{1}(U)\right) \cdots\right) \text {, where } \mathrm{F}_{1}^{0}(U)=U \text {. } \text {. }{ }_{l}(U)}_{l}$ Then

If $\quad n=p+q \quad$ is even, then $\langle U\rangle_{k}=\sum_{l=0}^{n} b_{k l} F_{1}^{l}(U), \quad$ where
$B_{n+1}=\left\|b_{k l}\right\|=A_{n+1}^{-1}, \quad A_{n+1}=\left\|a_{k l}\right\|, \quad a_{k l}=\lambda_{l-1}^{k-1}, \quad \lambda_{k}=(-1)^{k}(n-2 k)$.
If $n=p+q$ is $\quad$ odd, then $\quad\langle U\rangle_{k}+\langle U\rangle_{n-k}=\sum_{l=0}^{\frac{n-1}{2}} g_{k l} \mathrm{~F}_{1}^{l}(U)$, where
$G_{\frac{n+1}{2}}=\left\|g_{k l}\right\|=D_{\frac{n+1}{2}}^{-1}, \quad D_{\frac{n+1}{2}}=\left\|d_{k l}\right\|, \quad d_{k l}=\lambda_{l-1}^{k-1}$.

Theorem 25 ([61]). For $k=1, \ldots, n-1$ we have

$$
\begin{aligned}
\sum_{A:|A|=0 \bmod 4} e_{A}^{-1}\langle U\rangle_{k} e_{A} & =2^{\frac{n-2}{2}} \cos \left(\frac{\pi k}{2}-\frac{\pi n}{4}\right)\langle U\rangle_{k} \\
\sum_{A:|A|=1 \bmod 4} e_{A}^{-1}\langle U\rangle_{k} e_{A} & =(-1)^{k+1} 2^{\frac{n-2}{2}} \sin \left(\frac{\pi k}{2}-\frac{\pi n}{4}\right)\langle U\rangle_{k} \\
\sum_{A:|A|=2 \bmod 4} e_{A}^{-1}\langle U\rangle_{k} e_{A} & =-2^{\frac{n-2}{2}} \cos \left(\frac{\pi k}{2}-\frac{\pi n}{4}\right)\langle U\rangle_{k} \\
\sum_{A:|A|=3 \bmod 4} e_{A}^{-1}\langle U\rangle_{k} e_{A} & =(-1)^{k} 2^{\frac{n-2}{2}} \sin \left(\frac{\pi k}{2}-\frac{\pi n}{4}\right)\langle U\rangle_{k}
\end{aligned}
$$

For $m=0,1,2,3$ we have

$$
\sum_{A:|A|=m \bmod 4} e_{A}^{-1} e_{A}=d_{m}(n) e, \sum_{A:|A|=m \bmod 4} e_{A}^{-1} e_{1 \ldots n} e_{A}=(-1)^{m(n+1)} d_{m}(n) e_{1 \ldots n}
$$

where $d_{m}(n)=\operatorname{dim} \overline{\boldsymbol{m}}$ (see (3)).
Theorem 26 ([65]). Let $\mathrm{M}_{n}$ be the matrix of the size $2^{n}$ with the elements $m_{A B}=$ $\left.e_{A} e_{B} e_{A}^{-1} e_{B}^{-1}\right|_{e \rightarrow 1}$ (it is the commutator of $e_{A}$ and $e_{B}$ in the Salingaros group). Then we have

$$
\mathrm{F}_{e_{A}}(U):=e_{A}^{-1} U e_{A}=\sum_{B} m_{A B}\langle U\rangle_{e_{B}}, \quad U \in C \ell_{p, q}
$$

where $\langle U\rangle_{e_{B}}$ is the projection of the element $U$ onto the subspace spanned over $e_{B}$.
Using previous theorems, we can solve several classes of commutator equations (see [65], [61])

$$
e_{A} X+\epsilon X e_{A}=Q_{A}, \quad \epsilon \in \mathbb{R} \backslash\{0\}, \quad A \in \mathrm{G}
$$

for some known elements $Q_{A} \in C \ell_{p, q}$ and unknown element $X \in C \ell_{p, q}$, where G is some subset of the set of all ordered multi-indices with a length between 0 and $n$.
One can find other properties of considered operators in [65], [61], [35].

### 4.2. Pauli's Fundamental Theorem, Faithful and Irreducible Representations

Let the set of Clifford algebra elements satisfies the conditions

$$
\begin{equation*}
\left\{\beta_{a} ; a=1, \ldots, n\right\} \in C \ell_{p, q}, \quad \beta_{a} \beta_{b}+\beta_{b} \beta_{a}=2 \eta_{a b} e \tag{6}
\end{equation*}
$$

Then the set

$$
\begin{equation*}
\gamma_{a}=T^{-1} \beta_{a} T \tag{7}
\end{equation*}
$$

for any invertible $T \in C \ell_{p, q}$ satisfies the conditions

$$
\begin{equation*}
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a}=2 \eta_{a b} e \tag{8}
\end{equation*}
$$

Really

$$
\begin{aligned}
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a} & =T^{-1} \beta_{a} T T^{-1} \beta_{b} T+T^{-1} \beta_{b} T T^{-1} \beta_{a} T \\
& =T^{-1}\left(\beta_{a} \beta_{b}+\beta_{b} \beta_{a}\right) T=T^{-1} 2 \eta_{a b} e T=2 \eta_{a b} e
\end{aligned}
$$

But we are interested in another question. Does the element $T$ in (7) exist for every two sets (8) and (6)? Pauli proved the following theorem in 1936.

Theorem 27 (Pauli [38]). Consider two sets of square complex matrices

$$
\gamma_{a}, \quad \beta_{a}, \quad a=1,2,3,4
$$

of size 4. Let these 2 sets satisfy the following conditions

$$
\begin{aligned}
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a} & =2 \eta_{a b} \mathbf{1}, \quad \eta=\operatorname{diag}(1,-1,-1,-1) \\
\beta_{a} \beta_{b}+\beta_{b} \beta_{a} & =2 \eta_{a b} \mathbf{1} .
\end{aligned}
$$

Then there exists a unique (up to multiplication by a complex constant) complex matrix $T$ such that

$$
\gamma_{a}=T^{-1} \beta_{a} T, \quad a=1,2,3,4
$$

This theorem states that the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{1,3}$ has unique (up to equivalence) faithful and irreducible representation of dimension 4 .
Using the modern representation theory, we can obtain the following facts:

- In the case of even $n=p+q, \mathbb{C} \otimes C \ell_{p, q}$ has one faithful and irreducible representation of dimension $2^{\frac{n}{2}}\left(\mathbb{C} \otimes C \ell_{p, q} \cong \operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{C}\right), n\right.$ is even $)$.
- In the case of odd $n=p+q, \mathbb{C} \otimes C \ell_{p, q}$ has two irreducible representations of dimension $2^{\frac{n-1}{2}}$.
- In the case of odd $n=p+q, \mathbb{C} \otimes C \ell_{p, q}$ has two faithful reducible representation of dimension $2^{\frac{n-1}{2}}+2^{\frac{n-1}{2}}=2^{\frac{n+1}{2}}\left(\mathbb{C} \otimes C \ell_{p, q} \cong \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right) \oplus\right.$ $\operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), n$ is odd $)$.
Similarly we can formulate statements for the real Clifford algebra $C \ell_{p, q}$. The results depend on $n \bmod 2$ and $p-q \bmod 8$.
We also want to obtain an algorithm to compute the element $T$ that connects two sets of Clifford algebra elements. We can do this using the method of averaging in Clifford algebra and the operators $\sum_{A \in \mathrm{G}} \beta_{A} U \gamma_{A}^{-1}$, where G is some subset of the set of all ordered multi-indices with a length between 0 and $n$. One can find different properties of these operators in [64].
We have the following theorems.

Theorem 28 (The case of even $n$ [50]). Consider the real $C \ell_{p, q}$ (or the complexified $\left.\mathbb{C} \otimes C \ell_{p, q}\right)$ Clifford algebra with even $n=p+q$. Let two sets of Clifford algebra elements $\gamma_{a}, \beta_{a}, a=1,2, \ldots, n$ satisfy conditions

$$
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a}=2 \eta_{a b} e, \quad \beta_{a} \beta_{b}+\beta_{b} \beta_{a}=2 \eta_{a b} e
$$

Then both sets generate bases of Clifford algebra and there exists an unique (up to multiplication by a real (respectively complex) number) Clifford algebra element $T$ such that

$$
\gamma_{a}=T^{-1} \beta_{a} T, \quad a=1, \ldots, n .
$$

Additionally, we can obtain this element $T$ in the following way

$$
T=H(F):=\frac{1}{2^{n}} \sum_{A} \beta_{A} F\left(\gamma_{A}\right)^{-1}
$$

where $F$ is an element in the set

$$
\begin{array}{ll}
\text { 1) }\left\{\gamma_{A} ;|A|=0 \bmod 2\right\} & \text { if } \beta_{1 \ldots n} \neq-\gamma_{1 \ldots n} \\
\text { 2) }\left\{\gamma_{A} ;|A|=1 \bmod 2\right\} & \text { if } \beta_{1 \ldots n} \neq \gamma_{1 \ldots n}
\end{array}
$$

such that $H(F) \neq 0$.
Let us consider the case of odd $n$. We start with two examples.
Example 10. Let us consider the Clifford algebra $C \ell_{2,1} \simeq \operatorname{Mat}(2, \mathbb{R}) \oplus \operatorname{Mat}(2, \mathbb{R})$ with the generators $e_{1}, e_{2}, e_{3}$. We can take

$$
\gamma_{1}=e_{1}, \quad \gamma_{2}=e_{2}, \quad \gamma_{3}=e_{1} e_{2}
$$

Then $\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a}=2 \eta_{a b} 1$. The elements $\gamma_{1}, \gamma_{2}, \gamma_{3}$ generate not $C \ell_{2,1}$, but $C \ell_{2,0} \simeq$ $\operatorname{Mat}(2, \mathbb{R})$.

Example 11. Let us consider the Clifford algebra $C \ell_{3,0} \simeq \operatorname{Mat}(2, \mathbb{C})$ with the generators $e_{1}, e_{2}, e_{3}$. We can take

$$
\begin{gathered}
\beta_{1}=\sigma_{1}=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right), \quad \beta_{2}=\sigma_{2}=\left(\begin{array}{cc}
0 & -\mathrm{i} \\
\mathrm{i} & 0
\end{array}\right), \quad \beta_{3}=\sigma_{3}=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right) \\
\gamma_{a}=-\sigma_{a}, \quad a=1,2,3 .
\end{gathered}
$$

Then $\gamma_{123}=-\beta_{123}$. Suppose that we have $T \in \operatorname{GL}(2, \mathbb{C})$ such that $\gamma_{a}=$ $T^{-1} \beta_{a} T$. Then

$$
\gamma_{123}=T^{-1} \beta_{1} T T^{-1} \beta_{2} T T^{-1} \beta_{3} T=T^{-1} \beta_{1} \beta_{2} \beta_{3} T=\beta_{123}
$$

and we obtain a contradiction (we use that $\beta_{123}=\sigma_{123}=\mathrm{i}\left(\begin{array}{ll}1 & 0 \\ 0 & 1\end{array}\right)=\mathrm{i} 1$ ).
But we have $T=\mathbf{1}$ such that $\gamma_{a}=-T^{-1} \beta_{a} T$.

Theorem 29 (The case of odd $n$ [50]). Consider the real $C \ell_{p, q}$ (or the complexified $\left.\mathbb{C} \otimes C \ell_{p, q}\right)$ Clifford algebra with odd $n=p+q$. Suppose that two sets of Clifford algebra elements $\gamma_{a}, \beta_{a}, a=1,2, \ldots, n$ satisfy conditions

$$
\gamma_{a} \gamma_{b}+\gamma_{b} \gamma_{a}=2 \eta_{a b} e, \quad \beta_{a} \beta_{b}+\beta_{b} \beta_{a}=2 \eta_{a b} e
$$

Then, in the case of the Clifford algebra of signature $p-q \equiv 1 \bmod 4$, elements $\gamma_{1 \ldots n}$ and $\beta_{1 \ldots n}$ either take the values $\pm e_{1 \ldots n}$ and the corresponding sets generate bases of Clifford algebra (and we have cases 1-2 below) or take the values $\pm e$ and then the sets do not generate bases (and we have cases 3-4 below).
In the case of the Clifford algebra of signature $p-q \equiv 3 \bmod 4$, elements $\gamma_{1 \ldots n}$ and $\beta_{1 \ldots . n}$ either take the values $\pm e_{1 \ldots n}$ and the corresponding sets generate bases of Clifford algebra (and we have cases 1-2 below) or take the values $\pm \mathrm{i} e$ (this is possible only in the case of the complexified Clifford algebra) and then the sets do not generate bases (and we have cases 5-6 below).
There exists an unique (up to multiplication by an invertible element of the center of the Clifford algebra) element $T$ such that

1) $\gamma_{a}=T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=\gamma_{1 \ldots n}$
2) $\gamma_{a}=-T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=-\gamma_{1 \ldots n}$
3) $\gamma_{a}=e_{1 \ldots n} T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=e_{1 \ldots n} \gamma_{1 \ldots n}$
4) $\gamma_{a}=-e_{1 \ldots n} T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=-e_{1 \ldots n} \gamma_{1 \ldots n}$
5) $\gamma_{a}=i e_{1 \ldots n} T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=i e_{1 \ldots n} \gamma_{1 \ldots n}$
6) $\gamma_{a}=-i e_{1 \ldots n} T^{-1} \beta_{a} T, \quad a=1, \ldots, n \quad \Leftrightarrow \quad \beta_{1 \ldots n}=-i e_{1 \ldots n} \gamma_{1 \ldots n}$.

Note that all six cases have the unified notation $\gamma_{a}=\beta_{1 \ldots n}\left(\gamma_{1 \ldots n}\right)^{-1} T^{-1} \beta_{a} T$.
Additionally, in the case of the real Clifford algebra $C \ell_{p, q}$ of signature $p-q \equiv 1$ $\bmod 4$ and the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{p, q}$ of arbitrary signature, the element $T$, whose existence is stated in cases 1-6 of the theorem, equals

$$
T=H_{\text {Even }}(F):=\frac{1}{2^{n-1}} \sum_{A:|A|=0 \bmod 2} \beta_{A} F \gamma_{A}^{-1}
$$

where $F$ is an element of the set $\left\{\gamma_{A}+\gamma_{B} ;|A|=0 \bmod 2,|B|=0 \bmod 2\right\}$.
In the case of the real Clifford algebra $C \ell_{p, q}$ of signature $p-q \equiv 3 \bmod 4$, the element $T$, whose existence is stated in cases 1 and 2 of the theorem, equals $T=H_{\text {Even }}(F)$, where $F$ is an element of the set $\left\{\gamma_{A} ;|A|=0 \bmod 2\right\}$ such that $H_{\text {Even }}(F) \neq 0$.

Using the algorithm to compute the element $T$ in Theorems 28 and 29, we present an algorithm to compute elements of spin groups in [62].

In [34], we present a local variant of Pauli theorem, when two sets of Clifford algebra elements smoothly depend on the point of pseudo-Euclidian space.

## 5. Lie Groups and Lie Algebras in Clifford Algebras

### 5.1. Orthogonal Groups

Let us consider pseudo-orthogonal group $\mathrm{O}(p, q), p+q=n$

$$
\mathrm{O}(p, q):=\left\{A \in \operatorname{Mat}(n, \mathbb{R}) ; A^{\mathrm{T}} \eta A=\eta\right\}, \eta=\operatorname{diag}(\underbrace{1, \ldots, 1}_{p}, \underbrace{-1, \ldots,-1}_{q}) .
$$

It can be proved that (for more details, see [30] and [35])

$$
A \in \mathrm{O}(p, q) \Rightarrow \operatorname{det} A= \pm 1,\left|A_{1 \ldots p}^{1 \ldots p}\right| \geq 1,\left|A_{p+1 \ldots n}^{p+1 \ldots, n}\right| \geq 1, \quad A_{1 \ldots p}^{1 \ldots p}=\frac{A_{p+1 \ldots n}^{p+1 \ldots n}}{\operatorname{det} A}
$$

where $A_{1 \ldots p}^{1 \ldots p}$ and $A_{p+1 \ldots, n}^{p+1 \ldots, n}$ are the minors of the matrix $A$. The group $\mathrm{O}(p, q)$ has the following subgroups

$$
\begin{aligned}
\mathrm{SO}(p, q) & :=\{A \in \mathrm{O}(p, q) ; \operatorname{det} A=1\} \\
\mathrm{SO}_{+}(p, q) & :=\left\{A \in \mathrm{SO}(p, q) ; A_{1 \ldots p}^{1 \ldots p} \geq 1\right\}=\left\{A \in \mathrm{SO}(p, q) ; A_{p+1 \ldots n}^{p+1 \ldots n} \geq 1\right\} \\
& =\left\{A \in \mathrm{O}(p, q) ; A_{1 \ldots p}^{1 \ldots p} \geq 1, A_{p+1 \ldots n}^{p+\ldots n} \geq 1\right\} \\
\mathrm{O}_{+}(p, q) & :=\left\{A \in \mathrm{O}(p, q) ; A_{1 \ldots p}^{1 \ldots p} \geq 1\right\} \\
\mathrm{O}_{-}(p, q) & :=\left\{A \in \mathrm{O}(p, q) ; A_{p+\ldots n}^{p+1 \ldots n} \geq 1\right\} .
\end{aligned}
$$

The group $\mathrm{O}(p, q)$ has four components in the case $p \neq 0, \neq 0$

$$
\begin{gathered}
\mathrm{O}(p, q)=\mathrm{SO}_{+}(p, q) \sqcup \mathrm{O}_{+}(p, q)^{\prime} \sqcup \mathrm{O}_{-}(p, q)^{\prime} \sqcup \mathrm{SO}(p, q)^{\prime} \\
\mathrm{O}_{+}(p, q)=\mathrm{SO}_{+}(p, q) \sqcup \mathrm{O}_{+}(p, q)^{\prime}, \quad \mathrm{O}_{-}(p, q)=\mathrm{SO}_{+}(p, q) \sqcup \mathrm{O}_{-}(p, q)^{\prime} \\
\mathrm{SO}(p, q)=\mathrm{SO}_{+}(p, q) \sqcup \mathrm{SO}(p, q)^{\prime} .
\end{gathered}
$$

Example 12. In the cases $p=0$ or $q=0$, we have orthogonal groups $\mathrm{O}(n):=$ $\mathrm{O}(n, 0) \cong \mathrm{O}(0, n)$, special orthogonal groups $\mathrm{SO}(n):=\mathrm{SO}(n, 0) \cong \mathrm{SO}(0, n)$. The group $\mathrm{O}(n)$ has two connected components $\mathrm{O}(n)=\mathrm{SO}(n) \sqcup \mathrm{SO}(n)^{\prime}$.
Example 13. In the case $(p, q)=(1,3)$, we have Lorentz group $\mathrm{O}(1,3)$, special (or proper) Lorentz group $\mathrm{SO}(1,3)$, orthochronous Lorentz group $\mathrm{O}_{+}(1,3)$, orthochorous (or parity preserving) Lorentz group $\mathrm{O}_{-}(1,3)$, proper orthochronous Lorentz group $\mathrm{SO}_{+}(1,3)$.

Definition 16. A subgroup $H \subset G$ of a group $G$ is called a normal subgroup ( $H \triangleleft G$ ) if $g H^{-1} \subseteq H$ for all $g \in G$.
Definition 17. A quotient group (or factor group) $\frac{G}{H}:=\{g H ; g \in G\}$ is the set of all left cosets ( $\equiv$ right cosets, because $H$ is normal).

All considered subgroups are normal (for example, $\mathrm{SO}_{+}(p, q) \triangleleft \mathrm{O}(p, q)$ ) and

$$
\begin{align*}
\frac{\mathrm{O}(p, q)}{\mathrm{SO}_{+}(p, q)} & =\mathbb{Z}_{2} \times \mathbb{Z}_{2}, \quad \frac{\mathrm{O}(n)}{\mathrm{SO}(n)}=\mathbb{Z}_{2} \\
\frac{\mathrm{O}(p, q)}{\mathrm{SO}(p, q)} & =\frac{\mathrm{O}(p, q)}{\mathrm{O}_{-}(p, q)}=\frac{\mathrm{O}(p, q)}{\mathrm{O}_{+}(p, q)}=\frac{\mathrm{SO}(p, q)}{\mathrm{SO}_{+}(p, q)}=\frac{\mathrm{O}_{-}(p, q)}{\mathrm{SO}_{+}(p, q)}=\frac{\mathrm{O}_{+}(p, q)}{\mathrm{SO}_{+}(p, q)}=\mathbb{Z}_{2} \tag{9}
\end{align*}
$$

Example 14. The group $\mathrm{O}(1,1)$ has four connected components $\mathrm{O}_{+}^{\prime}(1,1), \mathrm{O}_{-}^{\prime}(1,1)$, $\mathrm{SO}^{\prime}(1,1), \mathrm{SO}_{+}(1,1)$ of matrices of the following type respectively (note that $\cosh ^{2} \psi=1+\sinh ^{2} \psi$ and $\left.\cosh \psi \geq 1\right)$

$$
\begin{array}{ll}
\left(\begin{array}{rr}
\cosh \psi & \sinh \psi \\
-\sinh \psi & -\cosh \psi
\end{array}\right), & \left(\begin{array}{rr}
-\cosh \psi & -\sinh \psi \\
\sinh \psi & \cosh \psi
\end{array}\right) \\
\left(\begin{array}{ll}
-\cosh \psi & -\sinh \psi \\
-\sinh \psi & -\cosh \psi
\end{array}\right), & \left(\begin{array}{ll}
\cosh \psi & \sinh \psi \\
\sinh \psi & \cosh \psi
\end{array}\right), \quad \psi \in \mathbb{R} .
\end{array}
$$

### 5.2. Lipschitz and Clifford Groups

Let us consider the group of all invertible elements

$$
C \ell_{p, q}^{\times}:=\left\{U \in C \ell_{p, q} ; V \in C \ell_{p, q} \text { exists }: U V=V U=e\right\}
$$

of dimension $\operatorname{dim} C \ell_{p, q}^{\times}=2^{n}$. The corresponding Lie algebra is $C \ell_{p, q}$ with the Lie bracket $[U, V]=U V-V U$.
Let us consider the adjoint representation

$$
\operatorname{Ad}: C \ell_{p, q}^{\times} \rightarrow \operatorname{Aut} C \ell_{p, q}, \quad T \rightarrow \operatorname{Ad}_{T}, \operatorname{Ad}_{T} U=T U T^{-1}, \quad U \in C \ell_{p, q}
$$

The kernel of Ad is (see Theorem 1)

$$
\begin{aligned}
\operatorname{ker}(\mathrm{Ad}) & =\left\{T \in C \ell_{p, q}^{\times} ; \operatorname{Ad}_{T}(U)=U \quad \text { for all } U \in C \ell_{p, q}\right\} \\
& = \begin{cases}C \ell_{p, q}^{0 \times}, & \text { if } n \text { is even } \\
\left(C \ell_{p, q}^{0} \oplus C \ell_{p, q}^{n}\right)^{\times}, & \text {if } n \text { is odd. }\end{cases}
\end{aligned}
$$

Let us consider the twisted adjoint representation

$$
\widetilde{\operatorname{Ad}}: C \ell_{p, q}^{\times} \rightarrow \operatorname{End} C \ell_{p, q}, \quad T \rightarrow \widetilde{\operatorname{Ad}_{T}}, \quad \widetilde{\operatorname{Ad}_{T}} U=\widehat{T} U T^{-1}, U \in C \ell_{p, q}
$$

The kernel of $\widetilde{A d}$ is

$$
\operatorname{ker}(\widetilde{\mathrm{Ad}})=\left\{T \in C \ell_{p, q}^{\times} ; \widetilde{\operatorname{Ad}}_{T}(U)=U \quad \text { for all } U \in C \ell_{p, q}\right\}=C \ell_{p, q}^{0 \times}
$$

In the Clifford algebra $C \ell_{p, q}$, we have a vector subspace $V=C \ell_{p, q}^{1}$ with a quadratic form $Q(x)$ or a symmetric bilinear form $g(x, x)$

$$
g(x, y)=\frac{1}{2}(Q(x+y)-Q(x)-Q(y))=\left.\frac{1}{2}(x y+y x)\right|_{e \rightarrow 1}, \quad x, y \in C \ell_{p, q}^{1}
$$

Lemma 30. $\widetilde{\mathrm{Ad}}: C \ell_{p, q}^{1 \times} \rightarrow \mathrm{O}(p, q)$ on $V$.
Proof: For $v \in C \ell_{p, q}^{1 \times}$ and $x \in C \ell_{p, q}^{1}$ we have

$$
Q\left(\widetilde{\operatorname{Ad}}_{v} x\right)=\left(\hat{v} x v^{-1}\right)^{2}=\hat{v} x v^{-1} \hat{v} x v^{-1}=x^{2}=Q(x)
$$

because $x^{2} \in C \ell_{p, q}^{0}$.
$\widetilde{\operatorname{Ad}}_{v}$ acts on $V$ as a reflection along $v$ (in the hyperplane orthogonal to $v$ )
$\widetilde{\operatorname{Ad}}_{v} x=\hat{v} x v^{-1}=x-(x v+v x) v^{-1}=x-2 \frac{g(x, v)}{g(v, v)} v, \quad v \in C \ell_{p, q}^{1 \times}, \quad x \in C \ell_{p, q}^{1}$.
Theorem 31 (Cartan-Diedonné). Every orthogonal transformation on a nongenerate space $(V, g)$ is a product of reflections (the number $\leq \operatorname{dim} V$ ) in hyperplanes.

Let us consider the group $\Gamma_{p, q}^{2}:=\left\{v_{1} v_{2} \cdots v_{k} ; v_{1}, \ldots, v_{k} \in C \ell_{p, q}^{1 \times}\right\}$.
Lemma 32. $\widetilde{\operatorname{Ad}}\left(\Gamma_{p, q}^{2}\right)=\mathrm{O}(p, q)$ (surjectivity).
Proof: If $f \in \mathrm{O}(p, q)$, then

$$
\begin{aligned}
f(x) & =\widetilde{\operatorname{Ad}}_{v_{1}} \circ \cdots \circ \widetilde{\operatorname{Ad}}_{v_{k}}(x)=\hat{v_{1}} \cdots \hat{v}_{k} x v_{k}^{-1} \cdots v_{1}^{-1} \\
& ={\widehat{v} \cdot v_{k}}^{v_{1}}\left(v_{1} \cdots v_{k}\right)^{-1}=\widetilde{\operatorname{Ad}}_{v_{1} \cdots v_{k}}(x)
\end{aligned}
$$

for $v_{1}, \ldots, v_{k} \in V^{\times}$and $x \in V$.
Let us consider the group $\Gamma_{p, q}^{1}:=\left\{T \in C \ell_{p, q}^{\times} ; \hat{T} x T^{-1} \in C \ell_{p, q}^{1}\right.$ for all $\left.x \in C \ell_{p, q}^{1}\right\}$ and the norm mapping (norm function) $N: C \ell_{p, q} \rightarrow C \ell_{p, q}, N(U):=\widehat{\widetilde{U}} U$.
Lemma 33. $N: \Gamma_{p, q}^{1} \rightarrow C \ell_{p, q}^{0 \times} \cong \mathbb{R}^{\times}$.
Proof: If $T \in \Gamma_{p, q}^{1}$ and $x \in C \ell_{p, q}^{1}$, then

$$
\widehat{T} x T^{-1}=\widetilde{\widehat{T} x T^{-1}}=\widetilde{T^{-1}} x \widetilde{\widehat{T}}=(\widetilde{T})^{-1} x \widetilde{\widehat{T}}
$$

Since $\widehat{\widehat{\widetilde{T}} T x}=x \widehat{\widetilde{T}} T$, it follows that $\widehat{\widetilde{T}} T \in \operatorname{ker} \widetilde{A d}=C \ell_{p, q}^{0 \times}$.
Lemma 34. $N: \Gamma_{p, q}^{1} \rightarrow \mathbb{R}^{\times}$is a group homomorphism

$$
N(U V)=N(U) N(V), \quad N\left(U^{-1}\right)=(N(U))^{-1}, \quad U, V \in \Gamma_{p, q}^{1}
$$

Proof: We have

$$
N(U V)=\widehat{\widehat{U V}} U V=\widehat{\tilde{V}} \widehat{\widetilde{U}} U V=\widehat{\widetilde{V}} N(U) V=N(U) N(V)
$$

and

$$
e=N(e)=N\left(U U^{-1}\right)=N(U) N\left(U^{-1}\right)
$$

Lemma 35. $\widetilde{\mathrm{Ad}}: \Gamma_{p, q}^{1} \rightarrow \mathrm{O}(p, q)$.
Proof: We have

$$
N(\widehat{T})=\widehat{\widehat{\widehat{T}}} \widehat{T}=\widehat{\widehat{\widetilde{T}} T}=\widehat{N(T)}=N(T)
$$

and

$$
\begin{aligned}
N\left(\widetilde{\operatorname{Ad}}_{T}(x)\right) & =N\left(\widehat{T} x T^{-1}\right)=N(\widehat{T}) N(x) N\left(T^{-1}\right) \\
& =N(T) N(x)(N(T))^{-1}=N(x)
\end{aligned}
$$

Since $N(x)=\widehat{\widetilde{x}} x=-x^{2}=-Q(x)$, it follows that $Q\left(\widetilde{\operatorname{Ad}}_{T}(x)\right)=Q(x)$.
Lemma 36. $\Gamma_{p, q}^{1}=\Gamma_{p, q}^{2}$.
Proof: We know that $\Gamma_{p, q}^{2} \subseteq \Gamma_{p, q}^{1}$. Let us prove that $\Gamma_{p, q}^{1} \subseteq \Gamma_{p, q}^{2}$. If $T \in \Gamma_{p, q}^{1}$, then $\widetilde{\operatorname{Ad}}_{T} \in \mathrm{O}(\underset{\sim}{p, q})$ by Lemma 35. Using Lemma 32, we conclude that $S \in \Gamma_{p, q}^{2}$ exists: $\widetilde{\operatorname{Ad}}_{S}=\widetilde{\operatorname{Ad}}_{T}$. We obtain $\widetilde{\mathrm{Ad}}_{T S^{-1}}=\mathrm{id}$ and $T S^{-1}=\lambda e, \lambda \in \mathbb{R}$. Finally, $T=\lambda S \in \Gamma_{p, q}^{2}$.

Definition 18. The following group is called Lipschitz group

$$
\begin{aligned}
\Gamma_{p, q}^{ \pm}:=\Gamma_{p, q}^{1}=\Gamma_{p, q}^{2} & =\left\{T \in C \ell_{p, q}^{(0) \times} \cup C \ell_{p, q}^{(1) \times} ; T x T^{-1} \in C \ell_{p, q}^{1} \text { for all } x \in C \ell_{p, q}^{1}\right\} \\
& =\left\{v_{1} v_{2} \cdots v_{k} ; v_{1}, \ldots, v_{k} \in C \ell_{p, q}^{1 \times}\right\}
\end{aligned}
$$

Definition 19. The following group is called Clifford group

$$
\Gamma_{p, q}:=\left\{T \in C \ell_{p, q}^{\times} ; T x T^{-1} \in C \ell_{p, q}^{1} \text { for all } x \in C \ell_{p, q}^{1}\right\} \supseteq \Gamma_{p, q}^{ \pm}
$$

So, we have $\widetilde{\operatorname{Ad}}\left(\Gamma_{p, q}^{ \pm}\right)=\mathrm{O}(p, q)$, i.e.,

$$
\begin{equation*}
\text { for any } P=\left\|p_{b}^{a}\right\| \in \mathrm{O}(p, q) \text { there exists } T \in \Gamma_{p, q}^{ \pm}: \widehat{T} e_{a} T^{-1}=p_{a}^{b} e_{b} \tag{10}
\end{equation*}
$$

Let us consider the following subgroup of Clifford group

$$
\Gamma_{p, q}^{+}:=\left\{T \in C \ell_{p, q}^{(0) \times} ; T x T^{-1} \in C \ell_{p, q}^{1} \text { for all } x \in C \ell_{p, q}^{1}\right\} \subset \Gamma_{p, q}^{ \pm}
$$

We have $\widetilde{\operatorname{Ad}}\left(\Gamma_{p, q}^{+}\right)=\operatorname{Ad}\left(\Gamma_{p, q}^{+}\right)=\operatorname{SO}(p, q)$, i.e., for all

$$
\begin{equation*}
P=\left\|p_{b}^{a}\right\| \in \mathrm{SO}(p, q) \text { there exists } T \in \Gamma_{p, q}^{+}: \widehat{T} e_{a} T^{-1}=T e_{a} T^{-1}=p_{a}^{b} e_{b} \tag{11}
\end{equation*}
$$

We can prove the statements in (10) and (11) without the Cartan-Diedonné theorem (see Theorem 31 and Lemmas 32-36) but with the use of the Pauli theorem (see Theorems 28 and 29). One can find this approach in [58] and [35].

### 5.3. Spin Groups

Let us define spin groups as normalized Lipschitz subgroups.
Definition 20. The following groups are called spin groups

$$
\begin{align*}
& \operatorname{Pin}(p, q):=\left\{T \in \Gamma_{p, q}^{ \pm} ; \widetilde{T} T= \pm e\right\}=\left\{T \in \Gamma_{p, q}^{ \pm} ; \widehat{\widetilde{T}} T= \pm e\right\} \\
& \operatorname{Pin}_{+}(p, q):=\left\{T \in \Gamma_{p, q}^{ \pm} ; \widehat{\widetilde{T}} T=+e\right\} \\
& \operatorname{Pin}_{-}(p, q):=\left\{T \in \Gamma_{p, q}^{ \pm} ; \widetilde{T} T=+e\right\}  \tag{12}\\
& \operatorname{Spin}(p, q):=\left\{T \in \Gamma_{p, q}^{+} ; \widetilde{T} T= \pm e\right\}=\left\{T \in \Gamma_{p, q}^{+} ; \widehat{\widetilde{T}} T= \pm e\right\} \\
& \operatorname{Spin}_{+}(p, q):=\left\{T \in \Gamma_{p, q}^{+} ; \widetilde{T} T=+e\right\}=\left\{T \in \Gamma_{p, q}^{+} ; \widehat{\widetilde{T}} T=+e\right\} .
\end{align*}
$$

In the case $p \neq 0$ and $q \neq 0$, we have

$$
\begin{gathered}
\operatorname{Pin}(p, q)=\operatorname{Spin}_{+}(p, q) \sqcup \operatorname{Pin}_{+}(p, q)^{\prime} \sqcup \operatorname{Pin}_{-}(p, q)^{\prime} \sqcup \operatorname{Spin}(p, q)^{\prime} \\
\operatorname{Pin}_{+}(p, q)=\operatorname{Spin}_{+}(p, q) \sqcup \operatorname{Pin}_{+}(p, q)^{\prime}, \operatorname{Pin}_{-}(p, q)=\operatorname{Spin}_{+}(p, q) \sqcup \operatorname{Pin}_{-}(p, q)^{\prime} \\
\operatorname{Spin}(p, q)=\operatorname{Spin}_{+}(p, q) \sqcup \operatorname{Spin}(p, q)^{\prime} .
\end{gathered}
$$

In Euclidian cases, we have two groups:

$$
\begin{gathered}
\operatorname{Pin}(n):=\operatorname{Pin}(n, 0)=\operatorname{Pin}_{-}(0, n), \operatorname{Spin}(n, 0)=\operatorname{Pin}_{+}(n, 0)=\operatorname{Spin}_{+}(n, 0) \\
\operatorname{Pin}(0, n):=\operatorname{Pin}(0, n)=\operatorname{Pin}_{+}(0, n), \operatorname{Spin}(0, n)=\operatorname{Pin}_{-}(0, n)=\operatorname{Spin}_{+}(0, n)
\end{gathered}
$$

All considered subgroups are normal (for example, $\operatorname{Spin}_{+}(p, q) \triangleleft \operatorname{Spin}(p, q)$ ).
All quotient groups are the same as for the group $\mathrm{O}(p, q)$ and its subgroups respectively (see (9)).

Theorem 37. The following homomorphisms are surjective with the kernel $\{ \pm 1\}$

$$
\begin{aligned}
& \widetilde{\mathrm{Ad}}: \operatorname{Pin}(p, q) \rightarrow \mathrm{O}(p, q) \\
& \widetilde{\mathrm{Ad}}: \operatorname{Spin}(p, q) \rightarrow \mathrm{SO}(p, q) \\
& \widetilde{\mathrm{Ad}}: \operatorname{Spin}_{+}(p, q) \rightarrow \mathrm{SO}_{+}(p, q) \\
& \widetilde{\mathrm{Ad}}: \operatorname{Pin}_{+}(p, q) \rightarrow \mathrm{O}_{+}(p, q) \\
& \widetilde{\mathrm{Ad}}: \operatorname{Pin}_{-}(p, q) \rightarrow \mathrm{O}_{-}(p, q)
\end{aligned}
$$

It means that for all

$$
\begin{equation*}
P=\left\|p_{b}^{a}\right\| \in \mathrm{O}(p, q) \text { there exists } \pm T \in \operatorname{Pin}(p, q): \widehat{T} e_{a} T^{-1}=p_{a}^{b} e_{b} \tag{13}
\end{equation*}
$$

and for the other groups similarly.

Proof: Statement for the group $\operatorname{Pin}(p, q)$ follows from the statements of the previous section (see Lemmas 32 and 33). For the other groups statement follows from the theorem on the norm of elements of spin groups which we give below (see [51] and [53]).

Theorem 38 ([51], [53]). The square of the norm of the element $T \in \operatorname{Pin}(p, q)$ in (13) equals

$$
\|T\|^{2}=\operatorname{Tr}\left(T^{\dagger} T\right)=\left\{\begin{aligned}
P_{1 \ldots p}^{1 \ldots p}=P_{p+1 \ldots n}^{p+1 \ldots n}, & T \in \operatorname{Spin}_{+}(p, q) \\
P_{1 \ldots p}^{1 \ldots p}=-P_{p+1 \ldots n}^{p+1 \ldots,} & T \in \operatorname{Pin}_{+}(p, q)^{\prime} \\
-P_{1 \ldots p}^{1 \ldots p}=P_{p+1 \ldots n}^{p+1 \ldots n}, & T \in \operatorname{Pin}_{-}(p, q)^{\prime} \\
-P_{1 \ldots p}^{1 \ldots p}=-P_{p+1 \ldots n}^{p+1 \ldots,} & T \in \operatorname{Spin}(p, q)^{\prime}
\end{aligned}\right.
$$

where $P_{1 \ldots p}^{1 \ldots p}$ and $P_{p+1 \ldots n}^{p+1 \ldots n}$ are the minors of the matrix $P \in \mathrm{O}(p, q)$ that corresponds to the element $T$ by (13).

Theorem 39. We have the isomorphism $\operatorname{Spin}(p, q) \cong \operatorname{Spin}(q, p)$.
Proof: This follows from the isomorphism $C \ell_{p, q}^{(0)} \cong C \ell_{q, p}^{(0)}$ (see Theorem 6).
Example 15. We have $\operatorname{Spin}(1,0)=\operatorname{Spin}(0,1)=\{ \pm e\}=\mathbb{Z}_{2}$.
Example 16. Note that $\operatorname{Pin}(p, q) \nsupseteq \operatorname{Pin}(q, p)$ in general case. For example, $\operatorname{Pin}(1,0)=\left\{ \pm e, \pm e_{1}\right\} \cong \mathbb{Z}_{2} \times \mathbb{Z}_{2}$ and $\operatorname{Pin}(0,1) \cong \mathbb{Z}_{4}$.

Theorem 40. The condition $T x T^{-1} \in C \ell_{p, q}^{1}$ for all $x \in C \ell_{p, q}^{1}$ holds automatically in the cases $n \leq 5$ for all spin groups (12), i.e.,

$$
\operatorname{Pin}(p, q)=\left\{T \in C \ell_{p, q}^{(0)} \cup C \ell_{p, q}^{(1)} ; \widetilde{T} T= \pm e\right\}, \quad n=p+q \leq 5
$$

Proof: If $T \in C \ell_{p, q}^{(0)} \cup C \ell_{p, q}^{(1)}$, then $T x T^{-1} \in C \ell_{p, q}^{1} \oplus C \ell_{p, q}^{3} \oplus C \ell_{p, q}^{5}$. Using $\widetilde{T} T= \pm e$, we get $\widetilde{T x T^{-1}}= \pm T x \widetilde{T}= \pm T x \widetilde{T}$ and $T x T^{-1} \in C \ell_{p, q}^{1} \oplus C \ell_{p, q}^{5}$. The statement is proved for $n \leq 4$.
Suppose that $n=5$ and $T x T^{-1}=v+\lambda e_{1 \ldots 5}, v \in C \ell_{p, q}^{1}, \lambda \in \mathbb{R}^{\times}$. Then

$$
\lambda=\left.\left(T x T^{-1} e_{1 \ldots 5}^{-1}-v e_{1 \ldots 5}^{-1}\right)\right|_{e \rightarrow 1}=\operatorname{Tr}\left(T x T^{-1} e_{1 \ldots 5}^{-1}\right)=\operatorname{Tr}\left(x e_{1 \ldots 5}^{-1}\right)=0
$$

and we obtain a contradiction.
Example 17. If the case $n=6$ the previous theorem is not valid. The element $T=\frac{1}{\sqrt{2}}\left(e_{12}+e_{3456}\right) \in C e_{6,0}^{(0)}$ satisfies $\widetilde{T} T=e$, but $T e_{1} T^{-1}=-e_{23456} \notin C \ell_{6,0}^{1}$.
Theorem 41. $\operatorname{Spin}_{+}(p, q)$ is isomorphic to the following groups in Table 5 in the cases $n=p+q \leq 6$.

Table 5. Isomorphisms between $\operatorname{Spin}_{+}(p, q)$ and matrix Lie groups.

| $p \backslash q$ | 0 | 1 | 2 | 3 | 4 | 5 | 6 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | $\mathrm{O}(1)$ | $\mathrm{O}(1)$ | $\mathrm{U}(1)$ | $\mathrm{SU}(2)$ | ${ }^{2} \mathrm{SU}(2)$ | $\mathrm{Sp}(2)$ | $\mathrm{SU}(4)$ |
| 1 | $\mathrm{O}(1)$ | $\mathrm{GL}(1, \mathbb{R})$ | $\mathrm{SU}(1,1)$ | $\mathrm{Sp}(1, \mathbb{C})$ | $\mathrm{Sp}(1,1)$ | $\mathrm{SL}(2, \mathbb{H})$ |  |
| 2 | $\mathrm{U}(1)$ | $\mathrm{SU}(1,1)$ | ${ }^{2} \mathrm{SU}(1,1)$ | $\mathrm{Sp}(2, \mathbb{R})$ | $\mathrm{SU}(2,2)$ |  |  |
| 3 | $\mathrm{SU}(2)$ | $\mathrm{Sp}(1, \mathbb{C})$ | $\mathrm{Sp}(2, \mathbb{R})$ | $\mathrm{SL}(4, \mathbb{R})$ |  |  |  |
| 4 | ${ }^{2} \mathrm{SU}(2)$ | $\mathrm{Sp}(1,1)$ | $\mathrm{SU}(2,2)$ |  |  |  |  |
| 5 | $\mathrm{Sp}(2)$ | $\mathrm{SL}(2, \mathbb{H})$ |  |  |  |  |  |
| 6 | $\mathrm{SU}(4)$ |  |  |  |  |  |  |

Note that

$$
\begin{aligned}
& \mathrm{U}(1) \simeq \mathrm{SO}(2), \quad \mathrm{SU}(2) \simeq \operatorname{Sp}(1), \quad \mathrm{SL}(2, \mathbb{C}) \simeq \operatorname{Sp}(1, \mathbb{C}) \\
& \mathrm{SU}(1,1) \simeq \operatorname{SL}(2, \mathbb{R}) \simeq \operatorname{Sp}(1, \mathbb{R}) .
\end{aligned}
$$

The Lie groups $\Gamma_{p, q}^{ \pm}, \Gamma_{p, q}^{+}$has the Lie algebra $C \ell_{p, q}^{0} \oplus C \ell_{p, q}^{2}$. All spin groups $\operatorname{Pin}(p, q), \operatorname{Spin}(p, q), \operatorname{Pin}_{+}(p, q), \operatorname{Pin}_{-}(p, q), \operatorname{Spin}_{+}(p, q)$ has the Lie algebra $C \ell_{p, q}^{2}$.
Since Theorem 37 and some facts from differential geometry, it follows that the spin groups are two-sheeted coverings of the orthogonal groups.
The groups $\operatorname{Spin}_{+}(p, q)$ are pathwise connected for $p \geq 2$ or $q \geq 2$. They are nontrivial covering groups of the corresponding orthogonal groups.

Example 18. The group $\operatorname{Spin}_{+}(1,1)=\left\{u e+v e_{12} ; u^{2}-v^{2}=1\right\}$ is not pathwise connected (it is two branches of the hyperbole).

The groups $\operatorname{Spin}_{+}(n), n \geq 3$ and $\operatorname{Spin}_{+}(1, n-1) \cong \operatorname{Spin}_{+}(n-1,1), n \geq 4$ are simply connected. They are the universal covering groups of the corresponding orthogonal groups.

### 5.4. Other Lie Groups and Lie Algebras in Clifford Algebra

Let us consider the following Lie groups and the corresponding Lie algebras (see Table 6).
Isomorphisms for the group $\mathrm{G}_{p, q}^{23 i 01}$ are proved in [67] (see also [66])

$$
\mathrm{G}_{p, q}^{23 i 01} \cong\left\{\begin{array}{l}
\mathrm{U}\left(2^{\frac{n}{2}}\right), \\
\mathrm{U}\left(2^{\frac{n-1}{2}}\right) \oplus \mathrm{U}\left(2^{\frac{n-1}{2}}\right) \\
\mathrm{U}\left(2^{\frac{n-2}{2}}, 2^{\frac{n-2}{2}}\right) \\
\mathrm{U}\left(2^{\frac{n-3}{2}}, 2^{\frac{n-3}{2}}\right) \oplus \mathrm{U}\left(2^{\frac{n-3}{2}}, 2^{\frac{n-3}{2}}\right) \\
\mathrm{GL}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right)
\end{array}\right.
$$

if $p$ is even and $q=0$
if $p$ is odd and $q=0$
if $n$ is even and $q \neq 0$
if $p$ is odd and $q \neq 0$ is even
if $p$ is even and $q$ is odd.

Table 6. Lie groups and Lie algebras in $\mathbb{C} \otimes C \ell_{p, q}$.

|  | Lie group | Lie algebra | dimension |
| :---: | :---: | :---: | :---: |
| 1 | $\left(\mathbb{C} \otimes C \ell_{p, q}\right)^{\times}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; U^{-1}\right.$ exists $\}$ | $\overline{0123} \oplus \mathrm{i} 0123$ | $2^{n+1}$ |
| 2 | $C C_{p, q}^{\times}=\left\{U \in C \ell_{p, q} ; U^{-1}\right.$ exists $\}$ | 0123 | $2^{n}$ |
| 3 | $C C_{p, q}^{(0) \times}=\left\{U \in C \ell_{p, q}^{(0)} ; U^{-1}\right.$ exists $\}$ | $\overline{02}$ | $2^{\text {n-1 }}$ |
| 4 | $\left(\mathbb{C} \otimes C \ell_{p, q}^{(0)}\right)^{\times}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; U^{-1}\right.$ exists $\}$ | $\overline{\mathbf{0 2}} \oplus \mathrm{i} \overline{\mathbf{0 2}}$ | $2^{n}$ |
| 5 | $\left(C \ell_{p, q}^{(0)} \oplus \mathrm{i} C \ell_{p, q}^{(1)} \times \times\left\{U \in C \ell_{p, q}^{(0)} \oplus \mathrm{i} C \ell_{p, q}^{(1)} ; U^{-1}\right.\right.$ exists $\}$ | $\overline{\mathbf{0 2}} \oplus \mathrm{i} \overline{13}$ | $2^{n}$ |
| 6 | $\mathrm{G}_{p, q}^{23 i 01}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; \bar{U} U=e\right\}$ | $\overline{23} \oplus \mathrm{i} \overline{01}$ | $2^{n}$ |
| 7 | $\mathrm{G}_{p, q}^{12 i 03}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; \underline{\hat{U}} U=e\right\}$ | $\overline{12} \oplus \mathrm{i} \overline{03}$ | $2^{n}$ |
| 8 | $\mathrm{G}_{p, q}^{2 \mathrm{ioj}}=\left\{U \in C C_{p, q}^{(0)} ; \tilde{U} U=e\right\}$ | $2 \oplus \mathrm{i} 0$ | $2^{\text {n-1 }}$ |
| 9 | $\mathrm{G}_{p, q}^{23 i 23}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; \tilde{U} U=e\right\}$ | $\overline{23} \oplus \mathrm{i} \overline{23}$ | $2^{n}-2^{\frac{n+1}{2}} \sin \frac{\pi(n+1)}{4}$ |
| 10 | $\mathrm{G}_{p, q}^{12,212}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q} ; \hat{U} U=e\right\}$ | $\overline{12} \oplus \mathrm{i} \overline{12}$ | $2^{n}-2^{\frac{n+1}{2}} \cos \frac{\pi(n+1)}{4}$ |
| 11 | $\mathrm{G}_{p, q}^{212}=\left\{U \in \mathbb{C} \otimes C \ell_{p, q}^{(0)} ; \tilde{U} U=e\right\}$ | $\overline{2} \oplus \mathrm{i} \overline{2}$ | $2^{n-1}-2^{\frac{n}{2}} \cos \frac{\pi n}{4}$ |
| 12 | $\mathrm{G}_{p, q}^{2 i 2}=\left\{U \in C \ell_{p, q}^{(0)} \oplus \mathrm{i} C \ell_{p, q}^{(1)} ; \underline{\tilde{U}} U=e\right\}$ | $\overline{2} \oplus \mathrm{i} \overline{1}$ | $2^{n-1}-2^{\frac{n-1}{2}} \cos \frac{\pi(n+1)}{4}$ |
| 13 | $\mathrm{G}_{p, q}^{223}=\left\{U \in C \chi_{p, q}^{(0)} \oplus \mathrm{i} C_{p, q}^{(1)} ; \overline{\hat{U}} U=e\right\}$ | $\overline{2} \oplus \mathrm{i} \overline{3}$ | $2^{n-1}-2^{\frac{n-1}{2}} \sin \frac{\pi(n+1)}{4}$ |
| 14 | $\mathrm{G}_{p, q}^{23}=\left\{U \in C \ell_{p, q} ; \tilde{U} U=e\right\}$ | 23 | $2^{n-1}-2^{\frac{n-1}{2}} \sin \frac{\pi(n+1)}{4}$ |
| 15 | $\mathrm{G}_{p, q}^{12}=\left\{U \in C \ell_{p, q} ; \hat{U} U=e\right\}$ | 12 | $2^{n-1}-2^{\frac{n-1}{2}} \cos \frac{\pi(n+1)}{4}$ |
| 16 | $\mathrm{G}_{p, q}^{2}=\left\{U \in C e_{p, q}^{(0)} ; \tilde{U} U=e\right\}$ | $\overline{2}$ | $2^{n-2}-2^{\frac{n-2}{2}} \cos \frac{\pi n}{4}$ |

We call $\mathrm{G}_{p, q}^{23 i 01}$ the pseudo-unitary group in Clifford algebra and use it in some problems of the field theory [35], [49], [36].
Some of these Lie groups are considered in [40] and [31]. Some of them are related to automorphism groups of the scalar products on the spinor spaces ([40], [31], [7], [4]). Note that spin group $\operatorname{Spin}_{+}(p, q)$ is a subgroup of all groups in Table 6. The group $\mathrm{G}_{p, q}^{2}$ coincides with $\operatorname{Spin}_{+}(p, q)$ in the cases $n \leq 5$. The Lie algebra of the spin group $C \ell_{p, q}^{2} \in \overline{\mathbf{2}}$ is a Lie subalgebra of all Lie algebras in Table 6. We have $C \ell_{p, q}^{2}=\overline{\mathbf{2}}$ in the cases $n \leq 5$.
The isomorphisms for the group $\mathrm{G}_{p, q}^{2}$ are represented in Tables 7 and 8. There is $n$ $\bmod 8$ in the lines and $p-q \bmod 8$ in the columns.
One can find isomorphisms for all remaining Lie groups and corresponding Lie algebras from Table 6 in a series of papers [60], [63], and [66].

## 6. Dirac Equation and Spinors in n Dimensions

### 6.1. Dirac Equation in Matrix Formalism

In Section 6, we use the notation with upper indices for the Dirac gamma-matrices and the generators of the Clifford algebra because of the useful covariant form of the Dirac equation.

Table 7. Isomorphisms for the group $\mathrm{G}_{p, q}^{2}$ in the cases of even $n$.

| $n \backslash p-q$ | 0 | 2, 6 | 4 |
| :---: | :---: | :---: | :---: |
| 0 | $\begin{aligned} & { }^{2} \mathrm{O}\left(2^{\frac{n-4}{2}}, 2^{\frac{n-4}{2}}\right) \\ & \text { if } p, q \neq 0 \\ & { }^{2} \mathrm{O}\left(2^{\frac{n-2}{2}}\right) \\ & \text { if } p=0 \text { or } q=0 \end{aligned}$ | $\mathrm{O}\left(2^{\frac{n-2}{2}}, \mathbb{C}\right)$ | ${ }^{2} \mathrm{O}\left(2^{\frac{n-4}{2}}, \mathbb{H}\right)$ |
| 2, 6 | $\mathrm{GL}\left(2^{\frac{n-2}{2}}, \mathbb{R}\right)$ | $\begin{aligned} & \mathrm{U}\left(2^{\frac{n-4}{2}}, 2^{\frac{n-4}{2}}\right) \\ & \text { if } p, q \neq 0 \\ & \mathrm{U}\left(\frac{n-2}{2}\right) \\ & \text { if } p=0 \text { or } q=0 \end{aligned}$ | $\mathrm{GL}\left(2^{\frac{n-4}{2}}, \mathbb{H}\right)$ |
| 4 | ${ }^{2} \mathrm{Sp}\left(2^{\frac{n-4}{2}}, \mathbb{R}\right)$ | $\mathrm{Sp}\left(2^{\frac{n-4}{2}}, \mathbb{C}\right)$ | $\begin{aligned} & { }^{2} \mathrm{Sp}\left(2^{\frac{n-6}{2}}, 2^{\frac{n-6}{2}}\right) \\ & \text { if } p, q \neq 0 \\ & { }^{2} \mathrm{Sp}\left(2^{\frac{n-4}{2}}\right) \\ & \text { if } p=0 \text { or } q=0 \end{aligned}$ |

Table 8. Isomorphisms for the group $\mathrm{G}_{p, q}^{2}$ in the cases of odd $n$.

| $n \backslash p-q$ | 1,7 | 3,5 |
| :---: | :--- | :--- |
|  | $\mathrm{O}\left(2^{\frac{n-3}{2}}, 2^{\frac{n-3}{2}}\right)$ |  |
| 1,7 | if $p, q \neq 0$ | $\mathrm{O}\left(2^{\frac{n-3}{2}}, \mathbb{H}\right)$ |
|  | $\mathrm{O}\left(2^{\frac{n-1}{2}}\right)$ |  |
|  | if $p=0$ or $q=0$ |  |
|  |  | $\mathrm{Sp}\left(2^{\frac{n-5}{2}}, 2^{\frac{n-5}{2}}\right)$ |
| 3,5 | $\mathrm{Sp}\left(2^{\frac{n-3}{2}}, \mathbb{R}\right)$ | if $p, q \neq 0$ <br> $\mathrm{Sp}\left(2^{\frac{n-3}{2}}\right)$ |
|  | if $p=0$ or $q=0$. |  |

Let $\mathbb{R}^{1,3}$ be Minkowski space with Cartesian coordinates $x^{\mu}, \mu=0,1,2,3$. The metric tensor of Minkowski space is given by a diagonal matrix

$$
\eta=\operatorname{diag}(1,-1,-1,-1)
$$

We denote partial derivatives by $\partial_{\mu}:=\frac{\partial}{\partial x^{\mu}}$.
The Dirac equation for the electron [15], [16] can be written in the following way

$$
\mathrm{i} \gamma^{\mu}\left(\partial_{\mu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi=0
$$

where $a_{\mu}: \mathbb{R}^{1,3} \rightarrow \mathbb{R}$ is the electromagnetic 4-vector potential, $m \geq 0 \in \mathbb{R}$ is the electron mass, $\psi: \mathbb{R}^{1,3} \rightarrow \mathbb{C}^{4}$ is the wave function (the Dirac spinor) and $\gamma^{\mu}$ are
the Dirac gamma-matrices which satisfy conditions

$$
\gamma^{\mu} \gamma^{\nu}+\gamma^{\nu} \gamma^{\mu}=2 \eta^{\mu \nu} \mathbf{1}, \quad \gamma^{\mu} \in \operatorname{Mat}(4, \mathbb{C})
$$

The Dirac equation is gauge invariant. If we take the expressions

$$
a_{\mu} \rightarrow \dot{a}_{\mu}=a_{\mu}+\lambda(x), \quad \psi \rightarrow \dot{\psi}=\psi \mathrm{e}^{\mathrm{i} \lambda(x)}, \quad \lambda(x) \in \mathbb{R}
$$

then they satisfy the same equation

$$
\begin{gathered}
\mathrm{i} \gamma^{\mu}\left(\partial_{\mu} \dot{\psi}-\mathrm{i} a_{\mu} \psi\right)-m \dot{\psi}=\mathrm{i} \gamma^{\mu}\left(\partial_{\mu}\left(\mathrm{e}^{\mathrm{i} \lambda} \psi\right)-\mathrm{i}\left(a_{\mu}+\partial_{\mu} \lambda\right)\left(\mathrm{e}^{\mathrm{i} \lambda} \psi\right)\right)-m\left(\mathrm{e}^{\mathrm{i} \lambda} \psi\right) \\
=\mathrm{i} \gamma^{\mu}\left(\mathrm{i}\left(\partial_{\mu} \lambda\right) \mathrm{e}^{\mathrm{i} \lambda} \psi+\mathrm{e}^{\mathrm{i} \lambda}\left(\partial_{\mu} \psi\right)-\mathrm{i} a_{\mu} \mathrm{e}^{\mathrm{i} \lambda} \psi-\mathrm{i}\left(\partial_{\mu} \lambda\right) \mathrm{e}^{\mathrm{i} \lambda} \psi\right)-m \mathrm{e}^{\mathrm{i} \lambda} \psi \\
=\mathrm{e}^{\mathrm{i} \lambda}\left(\mathrm{i} \gamma^{\mu}\left(\partial_{\mu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi\right)=0
\end{gathered}
$$

One says that the Dirac equation is gauge invariant with respect to the gauge group

$$
\mathrm{U}(1)=\left\{\mathrm{e}^{\mathrm{i} \lambda} ; \lambda \in \mathbb{R}\right\}
$$

The corresponding Lie algebra is

$$
\mathfrak{u}(1)=\{\mathrm{i} \lambda ; \lambda \in \mathbb{R}\} .
$$

The Dirac equation is relativistic invariant. Let us consider orthogonal transformation of coordinates

$$
x^{\mu} \rightarrow \dot{x}^{\mu}=p_{\nu}^{\mu} x^{\nu}, \quad P=\left\|p_{\nu}^{\mu}\right\| \in \mathrm{O}(1,3)
$$

Then

$$
\partial_{\mu} \rightarrow \dot{\partial}_{\mu}=q_{\mu}^{\nu} \partial_{\nu}, \quad a_{\mu} \rightarrow \dot{a}_{\mu}=q_{\mu}^{\nu} a_{\nu}, \quad Q=\left\|q_{\nu}^{\mu}\right\|=P^{-1}
$$

There are two points of view on transformations of the Dirac gamma-matrices and the wave function (see [69]).
In the first (tensor) approach, we have

$$
\gamma^{\mu} \rightarrow \dot{\gamma}^{\mu}=p_{\nu}^{\mu} \gamma^{\nu}, \quad \psi \rightarrow \dot{\psi}=\psi
$$

In this approach, all expressions are tensors and the Dirac equation is relativistic invariant. The tensor approach is considered in details in [32].
In the second (spinor) approach, we have

$$
\begin{gathered}
\gamma^{\mu} \rightarrow \dot{\gamma}^{\mu}=\gamma^{\mu}, \quad \psi \rightarrow \dot{\psi}=S \psi, \quad S^{-1} \gamma^{\mu} S=p_{\nu}^{\mu} \gamma^{\nu} \\
\left.\mathrm{i} \dot{\gamma}^{\mu}\left(\partial_{\mu} \dot{\psi}-\mathrm{i} \dot{a}_{\mu} \dot{\psi}\right)-m \dot{\psi}=\mathrm{i} \gamma^{\mu}\left(q_{\mu}^{\nu} \partial_{\nu}(S \psi)-\mathrm{i} q_{\mu}^{\nu} a_{\nu} S \psi\right)-m S \psi\right) \\
=S\left(\mathrm{i} S^{-1} q_{\mu}^{\nu} \gamma^{\mu} S\left(\partial_{\nu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi\right)=S\left(\mathrm{i} \gamma^{\nu}\left(\partial_{\nu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi\right)=0
\end{gathered}
$$

In this approach, the Dirac gamma-matrices do not change and the wave function $\psi$ changes as spinor with the aid of the element $S$ of the spin group. The formula $S^{-1} \gamma^{\mu} S=p_{\nu}^{\mu} \gamma^{\nu}$ describes the double cover of the orthogonal group by the spin group. This approach is generally accepted.

### 6.2. Dirac Equation in Formalism of Clifford Algebra

Let us consider the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{1,3}$ with the generators $e^{0}, e^{1}, e^{2}, e^{3}$. In Section 6, we use notation with upper indices for the generators of the Clifford algebra.
We have a primitive idempotent

$$
t=\frac{1}{2}\left(e+e^{0}\right) \frac{1}{2}\left(e+\mathrm{i} e^{12}\right) \longleftrightarrow\left(\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad t^{2}=t=t^{\dagger}
$$

The Dirac spinor is

$$
\psi \longleftrightarrow\left(\begin{array}{llll}
\psi_{1} & 0 & 0 & 0 \\
\psi_{2} & 0 & 0 & 0 \\
\psi_{3} & 0 & 0 & 0 \\
\psi_{4} & 0 & 0 & 0
\end{array}\right) \in \mathrm{I}(t)=\left(\mathbb{C} \otimes C l_{1,3}\right) t
$$

The corresponding left ideal $\mathrm{I}(t)$ is called spinor space.
The Dirac equation can be written in the following form

$$
\mathrm{i} e^{\mu}\left(\partial_{\mu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi=0
$$

where $\psi$ is an element of the left ideal of the Clifford algebra.
All properties of the Dirac equation from the previous section are valid.

### 6.3. Dirac-Hestenes Equation

Let us consider the Minkowski space $\mathbb{R}^{1,3}$ and the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{1,3}$ with the generators $e^{0}, e^{1}, e^{2}, e^{3}$. We have a primitive idemptonent $t=$ $\frac{1}{4}(e+E)(e-\mathrm{i} I)$ and the corresponding left ideal $\mathrm{I}(t)$, where $E:=e^{0}, I:=-e^{12}$, $t^{2}=t=t^{\dagger}$, $\mathrm{i} t=I t, t=E t$.

Lemma 42. For arbitrary $U \in \mathrm{I}(t)$ the equation $X t=U$ has a unique solution $X \in C \ell_{1,3}^{(0)}$ (and a unique solution $X \in C \ell_{1,3}^{(1)}$ ).

Proof: We can choose the orthonormal basis of the left ideal $\mathrm{I}(t)$ of the following form
$\tau_{k}=F_{k} t, \quad k=1,2,3,4, \quad F_{1}=2 e, F_{2}=2 e^{13}, F_{3}=2 e^{03}, F_{4}=2 e^{01} \in C \ell_{1,3}^{0)}$.
We have $U=\left(\alpha^{k}+\mathrm{i} \beta^{k}\right) \tau_{k}$ for some $\alpha^{k}, \beta^{k} \in \mathbb{R}$.
Using it $=I t$, we conclude that $X=F_{k}\left(\alpha^{k}+I \beta^{k}\right) \in C \ell_{1,3}^{(0)}$ is a solution of $X t=U$.

Now let us prove the following statement. If the element $Y \in C \ell_{1,3}^{(0)}$ is a solution of equation $Y t=0$, then $Y=0$. For element $Y t \in \mathrm{I}(t)$ we have

$$
Y t=\frac{1}{2}\left(\left(y-\mathrm{i} y_{12}\right) \tau_{1}+\left(-y_{13}-\mathrm{i} y_{23}\right) \tau_{2}+\left(y_{03}-\mathrm{i} y_{0123}\right) \tau_{3}+\left(y_{01}+\mathrm{i} y_{02}\right) \tau_{4}\right)=0
$$

Using $t=E t$, we conclude that $X=F_{k} E\left(\alpha^{k}+I \beta^{k}\right) \in C \ell_{1,3}^{(1)}$ is also a solution of equation $X t=U$. The proof of uniqueness in this case is similar.

One can find this lemma and similar statements, for example, in [32].
Let us rewrite the Dirac equation $\mathrm{i} e^{\mu}\left(\partial_{\mu} \psi-\mathrm{i} a_{\mu} \psi\right)-m \psi=0$ in the following form

$$
\begin{equation*}
e^{\mu}\left(\partial_{\mu} \psi-\mathrm{i} a_{\mu} \psi\right)+\mathrm{i} m \psi=0, \quad \psi \in \mathrm{I}(t) \tag{14}
\end{equation*}
$$

The Dirac-Hestenes equation [25] is

$$
\begin{equation*}
e^{\mu}\left(\partial_{\mu} \Psi-a_{\mu} \Psi I\right) E+m \Psi I=0, \quad \Psi \in C \ell_{1,3}^{(0)} \tag{15}
\end{equation*}
$$

Theorem 43. The Dirac equation and the Dirac-Hestenes equation are equivalent.
Proof: Let us multiply both sides of the Dirac-Hestenes equation (15) by $t$ on the right. Using $E t=t, I t=\mathrm{i} t$, and $\Psi t=\psi$, we obtain the Dirac equation (14).
We have $\psi \in \mathrm{I}(t)$. Using Lemma 42, we conclude that there exists $\Psi \in C \ell_{1,3}^{(0)}$ such that $\Psi t=\psi$. Using $E t=t$ and $I t=\mathrm{i} t$, we obtain

$$
\underbrace{\left(e^{\mu}\left(\partial_{\mu} \Psi-a_{\mu} \Psi I\right) E+m \Psi I\right)}_{\in C \ell_{1,3}^{(0)}} t=0 .
$$

Using Lemma 42 for the second time, we obtain the Dirac-Hestenes equation (15).

Note that the dimensions of the spinor spaces are the same in two approaches

$$
\operatorname{dim} \mathrm{I}(t)=\operatorname{dim} \mathbb{C}^{4}=8, \quad \operatorname{dim} C \ell_{1,3}^{(0)}=8
$$

The Dirac-Hestenes equation is widely used in applications (see, for example, [43], [18]).

### 6.4. Weyl, Majorana and Majorana-Weyl Spinors

Detailed information on $n$-dimensional spinors (using the methods of Clifford algebra) can be found in [59]. See also [7].
We study the connection between matrix operations (transpose, matrix complex conjugation) and operations in Clifford algebra (reverse, complex conjugation, grade involution), we introduce the notion of additional signature of the Clifford
algebra (for more details, see [59], also [60], [63], [66], where we develop and use these results).
Let us consider chirality operator (pseudoscalar) in $\mathbb{C} \otimes C \ell_{p, q}$

$$
\omega= \begin{cases}e^{1 \ldots n}, p-q=0,1 & \bmod 4 \\ \mathrm{i} e^{1 \ldots n}, p-q=2,3 & \bmod 4\end{cases}
$$

We have

$$
\omega=\omega^{-1}=\omega^{\dagger} .
$$

Let us consider orthogonal idempotents

$$
\begin{gathered}
P_{L}:=\frac{1}{2}(e-\omega), \quad P_{R}:=\frac{1}{2}(e+\omega) \\
P_{L}^{2}=P_{L}, \quad P_{R}^{2}=P_{R}, \quad P_{L} P_{R}=P_{R} P_{L}=0 .
\end{gathered}
$$

In the case of odd $n$, the complexified Clifford algebra $\mathbb{C} \otimes C \ell_{p, q}$ is the direct sum of two ideals
$\mathbb{C} \otimes C \ell_{p, q}=P_{L}\left(\mathbb{C} \otimes C \ell_{p, q}\right) \oplus P_{R}\left(\mathbb{C} \otimes C \ell_{p, q}\right), \quad \mathbb{C} \otimes C \ell_{p, q} \cong{ }^{2} \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right)$.
Let us consider the case of even $n$. For the set of Dirac spinors $E_{D}=\{\psi \in \mathrm{I}(t)\}$ we have

$$
E_{D}=E_{L W} \oplus E_{R W}
$$

where

$$
E_{L W}:=\left\{\psi \in E_{D} ; P_{L} \psi=\psi\right\}=\left\{\psi \in E_{D} ; \omega \psi=-\psi\right\}
$$

is the set of left Weyl spinors and

$$
E_{R W}:=\left\{\psi \in E_{D} ; P_{R} \psi=\psi\right\}=\left\{\psi \in E_{D} ; \omega \psi=\psi\right\}
$$

is the set of right Weyl spinors.
Using the Pauli theorem (Theorems 28 and 29), we obtain existence of the elements $A_{ \pm}$such that

$$
\begin{equation*}
\left(e^{a}\right)^{\dagger}= \pm A_{ \pm}^{-1} e^{a} A_{ \pm} \tag{16}
\end{equation*}
$$

If $n$ is even, then both elements $A_{ \pm}$exist. If $p$ is odd and $q$ is even, then only $A_{+}$ exists. If $p$ is even and $q$ is odd, then only $A_{-}$exists.
We can rewrite (16) in the following way

$$
U^{\dagger}=A_{+}^{-1} \overline{\tilde{U}} A_{+}, \quad U^{\dagger}=A_{-}^{-1} \widehat{\hat{\tilde{U}}} A_{-}, \quad U \in \mathbb{C} \otimes C \ell_{p, q}
$$

The explicit formulas for $A_{ \pm}$are given in Theorem 17.
Let us consider two Dirac conjugations

$$
\psi^{D_{ \pm}}:=\psi^{\dagger}\left(A_{ \pm}\right)^{-1}
$$

Example 19. In the case $(p, q)=(1,3)$ with the gamma-matrices $\gamma^{0}, \gamma^{1}, \gamma^{2}$, $\gamma^{3}$, we obtain the standard Dirac conjugation $\psi^{D_{+}}=\psi^{\dagger} \gamma^{0}$ and one else $\psi^{D-}=$ $\psi^{\dagger} \gamma^{123}$.

The Dirac conjugation is used to define bilinear covariants

$$
j_{ \pm}^{A}=\psi^{D_{ \pm}} e^{A} \psi
$$

The Dirac current $\psi^{D_{+}} e^{\mu} \psi$ is a particular case of the bilinear covariants. Using the Dirac equation, it is not difficult to obtain the law of conservation of the Dirac current

$$
\partial_{\mu}\left(\psi^{D_{+}} e^{\mu} \psi\right)=0
$$

We denote the matrix complex conjugation by $\leftarrow$. It should not be confused with the operation of complex conjugation in the complexified Clifford algebra $\mathbb{C} \otimes$ $C \ell_{p, q}$.
Let us consider the following two operations in $\mathbb{C} \otimes C \ell_{p, q}$

$$
U^{\mathrm{T}}:=\beta^{-1}\left(\beta^{\mathrm{T}}(U)\right), \quad \overleftarrow{U}:=\beta^{-1}(\overleftarrow{\beta(U)}), \quad U \in \mathbb{C} \otimes C \ell_{p, q}
$$

where
$\beta: \mathbb{C} \otimes C \ell_{p, q} \rightarrow \begin{cases}\operatorname{Mat}\left(2^{\frac{n}{2}}, \mathbb{C}\right), & \text { if } n \text { is even } \\ \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right) \oplus \operatorname{Mat}\left(2^{\frac{n-1}{2}}, \mathbb{C}\right), & \text { if } n \text { is odd }\end{cases}$
is the faithful representation of $\mathbb{C} \otimes C \ell_{p, q}$ of the minimal dimension. These two operations depend on the representation $\beta$.
Using the Pauli theorem (Theorems 28 and 29), we obtain existence of the elements $C_{ \pm}$such that

$$
\begin{equation*}
\left(e^{a}\right)^{\mathrm{T}}= \pm C_{ \pm}^{-1} e^{a} C_{ \pm} \tag{17}
\end{equation*}
$$

If $n$ is even, then both elements $C_{ \pm}$exist. If $n=1 \bmod 4$, then only $C_{+}$exists. If $n=3 \bmod 4$, then only $C_{-}$exists.
We can rewrite (17) in the following way

$$
U^{\mathrm{T}}=C_{+}^{-1} \tilde{U} C_{+}, \quad U^{\mathrm{T}}=C_{-}^{-1} \hat{\tilde{U}} C_{-}, \quad U \in \mathbb{C} \otimes C \ell_{p, q}
$$

The explicit formulas for $C_{ \pm}$are given in [59] using the notion of additional signature of the Clifford algebra. Also these elements have the following properties

$$
\begin{gathered}
\left(C_{ \pm}\right)^{\mathrm{T}}=\lambda_{ \pm} C_{ \pm}, \quad \overleftarrow{C_{ \pm}} C_{ \pm}=\lambda_{ \pm} e \\
\lambda_{+}=\left\{\begin{array}{c}
+1, n \equiv 0,1,2 \bmod 8 \\
-1, n \equiv 4,5,6
\end{array} \quad \lambda_{-}=\left\{\begin{array}{l}
+1, n \equiv 0,6,7 \bmod 8 \\
-1, n \equiv 2,3,4 \bmod 8
\end{array}\right.\right.
\end{gathered}
$$

Using the Pauli theorem (Theorems 28 and 29), we obtain existence of the elements $B_{ \pm}$such that

$$
\begin{equation*}
\overleftarrow{e^{a}}= \pm B_{ \pm}^{-1} e^{a} B_{ \pm} \tag{18}
\end{equation*}
$$

If $n$ is even, then both elements $B_{ \pm}$exist. If $p-q=1 \bmod 4$, then only $B_{+}$ exists. If $p-q=3 \bmod 4$, then only $B_{-}$exists.
We can rewrite (18) in the following way

$$
\overleftarrow{U}=B_{+}^{-1} \bar{U} B_{+}, \quad \overleftarrow{U}=B_{-}^{-1} \widehat{\bar{U}} B_{-}, \quad U \in \mathbb{C} \otimes C \ell_{p, q}
$$

The explicit formulas for $B_{ \pm}$are given in [59] using the notion of additional signature of the Clifford algebra. Also these elements have the following properties

$$
\begin{array}{cc}
B_{ \pm}^{T}=\epsilon_{ \pm} B_{ \pm}, & \overleftarrow{B_{ \pm}} B_{ \pm}=\epsilon_{ \pm} e \\
\epsilon_{+}= \begin{cases}+1, p-q \equiv 0,1,2 \bmod 8 \\
-1, p-q \equiv 4,5,6 \bmod 8\end{cases} & \epsilon_{-}=\left\{\begin{array}{l}
+1, p-q \equiv 0,6,7 \bmod 8 \\
-1, p-q \equiv 2,3,4 \bmod 8
\end{array}\right.
\end{array}
$$

We introduce the Majorana conjugation in the following way

$$
\psi^{M_{ \pm}}:=\psi^{\mathrm{T}}\left(C_{ \pm}\right)^{-1}
$$

Example 20. In the case $(p, q)=(1,3)$, we have $\psi^{M_{+}}=\psi^{\dagger}\left(\gamma^{13}\right)^{-1}$ and $\psi^{M_{-}}$ $=\psi^{\dagger}\left(\gamma^{02}\right)^{-1}$.

We introduce the charge conjugation in the following way

$$
\psi^{c h_{ \pm}}:=B_{ \pm} \overleftarrow{\psi}
$$

Example 21. In the case $(p, q)=(1,3)$, we have $\psi^{c h_{+}}=\gamma^{013} \overleftarrow{\psi}$ and $\psi^{c h_{-}}$ $=\gamma^{2} \overleftarrow{\psi}$

We have the following relation between $A_{ \pm}, B_{ \pm}$, and $C_{ \pm}$(when they exist)

$$
\begin{array}{cl}
B_{+}=\widetilde{A_{+}^{-1}} C_{+}, \quad \widehat{\widehat{A_{-}^{-1}} C_{-}}, \quad B_{-}=\widetilde{A_{-}^{-1}} C_{+}, \quad B_{-}=\widehat{\widehat{A_{+}^{-1}} C_{-}} \\
\psi^{c h_{+}}=C_{+}\left(\psi^{D_{+}}\right)^{\mathrm{T}}=C_{-}\left(\psi^{D_{-}}\right)^{\mathrm{T}}, & \psi^{c h_{-}}=C_{-}\left(\psi^{D_{+}}\right)^{\mathrm{T}}=C_{+}\left(\psi^{D_{-}}\right)^{\mathrm{T}} .
\end{array}
$$

Let us denote the set of Majorana spinors by

$$
E_{M}:=\left\{\psi \in E_{D} ; \psi^{c h_{-}}= \pm \psi\right\}
$$

and the set of pseudo-Majorana spinors by

$$
E_{p s M}:=\left\{\psi \in E_{D} ; \psi^{c h_{+}}= \pm \psi\right\} .
$$

Using definition of the charge conjugation and the properties of $B_{ \pm}$, it can be proved that Majorana spinors are realized only in the cases $p-q=0,6,7 \bmod 8$ and pseudo-Majorana spinors are realized only in the cases $p-q=0,1,2 \bmod 8$ (see, for example, [59]).

Let us denote a set of left Majorana-Weyl spinors by

$$
E_{L M W}:=\left\{\psi \in E_{L W} ; \psi^{c h_{-}}= \pm \psi\right\}=\left\{\psi \in E_{L W} ; \psi^{c h_{+}}= \pm \psi\right\}
$$

and a set of right Majorana-Weyl spinors by

$$
E_{R M W}:=\left\{\psi \in E_{R W} ; \psi^{c h_{-}}= \pm \psi\right\}=\left\{\psi \in E_{R W} ; \psi^{c h_{+}}= \pm \psi\right\}
$$

It can be proved that Majorana-Weyl spinors are realized only in the cases $p-q=0$ $\bmod 8$ (see, for example, [59]).
The question of existence of the spinors in the cases of different dimensions and signatures is related to the supersymmetry theory (see classic works on supersymmetry and supergravity [28], [21] and other papers and reviews [14], [70], [19], [73], [72], [41], [74], [71]).

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