# GENERALIZED RADON TRANSFORM AND LÉVY'S BROWNIAN MOTION, I*) 

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## §1. Introduction

In connection with a Gaussian system $X=\{X(x) ; x \in M\}$ called Lévy's Brownian motion (Definition 1), we shall introduce two integral transformations of special type-one is a generalized Radon transform $R$ on a measure space $(M, m)$, and the other is a dual Radon transform $R^{*}$ on another measure space $(H, \nu)$ such that $H \subset 2^{M}$, the set of all subsets of $M$ (Definition 2). To each Lévy's Brownian motion $X$, there is attached a distance $d(x, y):=E\left[(X(x)-X(y))^{2}\right]$ on $M$ having a notable property named $L^{1}$-embeddability ([3]). The above measure $\nu$ on $H$ is then chosen to satisfy

$$
d(x, y)=\nu\left(B_{x} \triangle B_{y}\right) \quad \text { with } B_{x}:=\{h \in H ; x \in h\},
$$

where $\triangle$ stands for the symmetric difference.
It turns out that these transforms constitute a factorization of the covariance operator of $X$ (Theorem 3); a more explicit link between $X$ and $R^{*}$ can be noticed in the somewhat informal expression

$$
X(x)=\left(R^{*} W\right)(x)
$$

where $W=\{W(d h) ; h \in H\}$ is a Gaussian random measure with mean 0 and variance $\nu(d h)$. In view of the quite simple probabilistic structure of $W$, an idea comes to mind: The deep study of $R$ and $R^{*}$ will yield fruitful results on $X$. Thus, we shall investigate the transforms $R$ and $R^{*}$ as well as the Lévy's Brownian motion $X$ in the present and subsequent papers.

The main purpose of the present paper (I) is to obtain the singular value decomposition of $R^{*}$ (Theorem 5), which gives us the Karhunen-

[^0]Loève expansion of $X$ (Theorem 6). The second paper (II) will concentrate on the investigation of the null spaces of $R^{*}$ :

$$
N_{1}(A):=\left\{g \in L^{2}(H, \nu) ;\left(R^{*} g\right)(x) \equiv 0, x \in A\right\}, \quad A \subset M
$$

The structure of the closed linear span $[X(x) ; x \in A]$ in $L^{2}(\Omega, P)$ will be described in terms of $N_{1}(A)$ and $W$.

In order to give some interpretation to the representation of Chentsov type which is useful for our study, we begin with a familiar Brownian motion $X=\left\{X(x) ; x \in R^{n}\right\}$ with $n$-dimensional parameter. The variance of the increment $X(x)-X(y)$ is, by definition, equal to the Euclidean distance $|x-y|$ between $x$ and $y$. The idea of Chentsov [6] (cf. [24] and [26]) now leads us to take the following measure space ( $H, \nu$ ): $H$ is the set of all half-spaces $h_{t, \omega}:=\left\{x \in R^{n} ;(x, \omega)>t\right\}$ not containing the origin $O$; an element $h_{t, \omega} \in H$ is parametrized by the distance $t>0$ and the direction $\omega \in S^{n-1}:=\left\{\omega \in R^{n} ;|\omega|=1\right\}$. The measure $\nu$ is an invariant measure on $H$, explicitly given by

$$
\nu\left(d h_{t, \omega}\right)=\frac{n-1}{\left|S^{n-2}\right|} d t d \omega .
$$

Then it is easy to verify that $\nu\left(B_{x} \triangle B_{y}\right)=|x-y|$. We thus get at the conclusion that $X$ is expressed in the form

$$
\begin{equation*}
X(x)=\int_{B_{x}} W(d h)=W\left(B_{x}\right) . \tag{1}
\end{equation*}
$$

A general framework behind the representation (1) of Chentsov type consists of the following:
(i) A centered Gaussian system $X=\{X(x) ; x \in M\}$ with parameter space $M$; the variance of the increment is denoted by $d(x, y):=E[(X(x)$ $\left.-X(y))^{2}\right]$.
(ii) A Gaussian random measure $W=\{W(d h) ; h \in H\}$ based on a measure space $(H, \nu)$ such that $H \subset 2^{M}$ and $\nu\left(B_{x}\right)<\infty$ for all $x \in M$. It follows from (1) that

$$
\begin{equation*}
d(x, y)=\int_{H}\left|\chi_{B_{x}}(h)-\chi_{B_{y}}(h)\right| \nu(d h)=\nu\left(B_{x} \triangle B_{y}\right), \tag{2}
\end{equation*}
$$

where $\chi_{B}$ denotes the indicator function of a subset $B \subset H$. Conversely, this equation (2) guarantees the existence of such a representation (1). The variance of $X$ admitting a representation (1) of Chentsov type is
therefore a (semi-)metric on $M$ of the form $\left\|\chi_{B_{x}}-\chi_{B_{y}}\right\|_{L^{1}(H, \nu)}$; such a metric is said to be $L^{1}$-embeddable ([3]).

We are now in a position to introduce the following
Definition 1. Let $(M, d)$ be an $L^{1}$-embeddable metric space. Then a centered Gaussian system $X=\{X(x) ; x \in M\}$ with the variance $d(x, y)$ of the increment $X(x)-X(y)$ is called Lévy's Brownian motion with parameter space ( $M, d$ ).

With this terminology, our first conclusion (Theorem 1) is that every Lévy's Brownian motion admits of a representation of the form (1).

Another ingredient in our study is a pair of integral transformations associated with the expression (1).

Definition 2. Let $m(d x)$ be a reference measure on $M$. The integral transform

$$
\begin{equation*}
(R f)(h):=\int_{h} f(x) m(d x), \quad f \in L^{1}(M, m) \tag{3}
\end{equation*}
$$

(resp.

$$
\begin{equation*}
\left.\left(R^{*} g\right)(x):=\int_{B_{x}} g(h) \nu(d h), \quad g \in L^{2}(H, \nu),\right) \tag{4}
\end{equation*}
$$

is called a generalized (resp. dual) Radon transform.
The reason for using the symbol $R^{*}$ lies in the obvious relation of duality:

$$
(R f, g)_{L^{2}(H, \nu)}=\left(f, R^{*} g\right)_{L^{2}(M, m)}
$$

In case $X$ is a Brownian motion with $n$-dimensional parameter, the value $(R f)\left(h_{t, \omega}\right)$ is nothing but the integral of $f$ over the half-space $h_{t, \omega}$ and hence the classical Radon transform, the integral over the hyperplane $\delta h_{t, \omega}$ (Radon's celebrated paper [31]; see also [8], [15] and [23]) can be derived from the first variation of $R$ (cf. [19], p. 47). On the other hand, the dual Radon transform $R^{*}$ is closely related to the one studied by Cormack and Quinto [7], because the set $B_{x}$ is changed into the open ball $\tilde{B}_{x}$ with diameter $\overline{O x}$ by means of the mapping
$h_{t, \omega} \in H \longmapsto y=t \omega \in R^{n} \backslash\{O\}$, the foot of the perpendicular from
$O$ to the hyperplane $\delta h_{t, \alpha}$.
Another important example should be mentioned here; it is a Lévy's Brownian motion with parameter space $\left(S^{n}, d_{G}\right)$, $d_{G}$ being the geodesic
distance on $S^{n}$. Due to Lévy [21] (cf. also [18]), the corresponding measure space $(H, \nu)$ is chosen to be the set of all hemispheres endowed with an invariant measure $\nu$. In this case, the transforms $R$ and $R^{*}$ take the same form-the integral over a hemisphere. Since the integral over a great circle can also be derived from the first variation of $R$, the study of $R$ and $R^{*}$ has another origin in Funk [11] and [12].

In Section 2 we shall establish the representation (1) of Chentsov type, and give several examples of ( $M, d$ ) and ( $H, \nu$ ), except the case of $M=R^{n}$, the usual parameter space of random fields. A variety of $L^{1}$-embeddable metrics $d$ on $R^{n}$ will be described in the second paper (II).

Section 3 is devoted to the study of fundamental properties of $R$ and $R^{*}$. In particular, we shall obtain their singular value decompositions, which will be applied to show that $X$ admits of the Karhunen-Loève expansion in terms of an i.i.d. sequence of standard Gaussian random variables.

Section 4 will concern the $n$-sphere $M=S^{n}$ equipped with the uniform probability measure $\sigma$. The Karhunen-Loéve expansion will be explicitly calculated for a certain class of Lévy's Brownian motions $X=$ $\left\{X(x) ; x \in S^{n}\right\}$ including the one due to Lévy [21] mentioned above; all of them have probability laws invariant under every rotation on $S^{n}$.

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## §2. Representations of Chentsov type

The purpose of this section is two-fold: to prove the representation (1) of Chentsov type for each Lévy's Borwnian motion $X$, and to give several examples of $(H, \nu)$ combined with ( $M, d$ ) via the equality (2). Particular attention will be paid to the case of $M=S^{n}$.

Suppose that ( $M, d$ ) is an $L^{1}$-embeddable metric space; by definition, there exist a measure space ( $T, \mu$ ) and a mapping $x \in M \mapsto f_{x}(t) \in L^{1}(T, \mu)$ such that $d(x, y)=\left\|f_{x}(t)-f_{y}(t)\right\|_{L^{1}(T, \mu)}$. Then, as was shown by Assouad and Deza [3], we can find another measure space ( $H, \nu$ ) satisfying $H \subset 2^{M}$ and

$$
\begin{equation*}
d(x, y)=\nu\left(B_{x} \triangle B_{y}\right)=\int_{H} \pi_{h}(x, y) \nu(d h), \tag{2}
\end{equation*}
$$

where we have used the notation

$$
\pi_{h}(x, y):=\left|\chi_{h}(x)-\chi_{h}(y)\right|=\left|\chi_{B_{x}}(h)-\chi_{B_{y}}(h)\right| .
$$

Among various kinds of possible realizations of the distance $d$, the above one in terms of the indicator function $\chi_{B_{x}}(h)$ in $L^{1}(H, \nu)$ is most convenient for us to associate the transforms $R$ and $R^{*}$ with the expression (1). It was called a multiplicity realization in [3]. The correspondence $(M, d) \mapsto(H, \nu)$ has a tiny fault, however; it is not one to one (see Example 1b below).

Having found a multiplicity realization $\chi_{B_{x}}(h)$ in $L^{1}(H, \nu)$ of a given $L^{1}$-embeddable metric $d$ on $M$, our first conclusion follows immediately:

Theorem 1. A Lévy's Brownian motion $X$ with parameter space ( $M, d$ ) admits of the representation

$$
\begin{equation*}
X(x)=\int_{B_{x}} W(d h) \tag{1}
\end{equation*}
$$

in terms of a Gaussian random measure $W$ based on the measure space ( $H, \nu$ ).

Now choose and fix a point $O \in M$ as the origin. In view of a simple fact that $\pi_{h c} \equiv \pi_{h}$, we may change an element $h \in H$ with its complement $h^{c}$ if $O \in h$, so that $H \subset\left(2^{M}\right)_{o}:=\{h \subset M$; $O \oplus h\}$. This choice of $H$ implies that $B_{o}=\phi$, which leads to the assumption $X(O)=0$ often added in the definition of Lévy's Brownian motion.

Example 1. Let us mention a couple of examples in which $(M, d)$ is induced by a graph $G([14])$, i.e., $M$ is the set of all vertexes and $d(x, y)$ is the number of edges in a shortest path between $x$ and $y$.
(a) $G=T$, a tree. At each edge $e$ of $T, M$ is separated into the two complementary subsets $h_{e}$ and $h_{e}^{c}$; the root $O$ of $T$ always belongs to $h_{e}^{c}$. Define

$$
H=\left\{h_{e} \text { for all edges } e\right\} \subset\left(2^{M}\right)_{o} \text { with weight } \nu\left(h_{e}\right) \equiv 1,
$$

to get the desired distance $d$ on $M$. With this choice of $(H, \nu)$, the representation (1) of Chentsov type can be regarded as a simple extension of partial sums of a sequence of i.i.d. Gaussian random variables.
(b) $G=K_{m}$, the complete graph of $m$ vertexes. It is possible to find several different kinds of $(H, \nu)$. Indeed, for each $k, 1 \leq k \leq[m / 2]$, take

$$
H_{k}=\{\text { all subsets } h \text { of } k \text { vertexes }\} \text { with weight } \nu_{k}(h) \equiv\left\{2\binom{m-2}{k-1}\right\}^{-1}
$$

Then it is easy to show that

$$
d(x, y):=\sum_{h \in \vec{H}_{k}} \pi_{h}(x, y) \nu_{k}(h) \equiv 1 \quad \text { for any } x, y \in M
$$

We note that ( $M, d$ ) induced by a cyclic graph $C_{m}$ is a discrete analogue of ( $S^{1}, d_{G}$ ) and hence the corresponding measure space ( $H, \nu$ ) can be constructed after the manner of the one described in Section 1.

Example 2. In case $M$ is the set of all natural numbers, an easy way to get an $L^{1}$-embeddable metric $d$ on $M$ is as follows: Take

$$
H:=\left\{h_{m} ; m \geq 2\right\} \quad \text { with weight } \nu\left(h_{m}\right) \geq 0,
$$

and define

$$
d(x, y):=\sum_{m=2}^{\infty} \pi_{h_{m}}(x, y) \nu\left(h_{m}\right),
$$

where $h_{m}:=\{m k ; k=1,2, \cdots\}$ is the set of all multiples of $m$. The special choice of weight

$$
\nu\left(h_{m}\right)=\log p \text { if } m \text { has only one prime factor } p,=0 \text { otherwise }
$$

gives us the interesting distance $d(x, y)=\log (x \cup y / x \cap y)$ mentioned in [1] and [2], where $x \cup y$ (resp. $x \cap y$ ) denotes the L. C.M. (resp. G. C. M.) of $x$ and $y$.

A generalized Radon transform of the form

$$
(R f)\left(h_{m}\right):=\sum_{k=1}^{\infty} f(m k)
$$

was considered by Strichartz [34], who gave the inversion formula

$$
\begin{equation*}
f(x)=\sum_{k=1}^{\infty} \mu(k)(R f)\left(h_{x k}\right), \tag{5}
\end{equation*}
$$

where $\mu(k)$ is the Möbius function defined by

$$
\mu(k):=\left\{\begin{array}{cl}
(-1)^{l}, & \text { if } k \text { has } l \text { distinct prime factors }, \\
0, & \text { if } k \text { is divisible by the square of a prime. }
\end{array}\right.
$$

The representation (1) for a Lévy's Brownian motion $X$ with parameter space ( $M, d$ ) now takes the form

$$
X(x)=\sum_{m \mid x} W\left(h_{m}\right), \quad x \geq 2, \quad \text { and } \quad X(1)=0
$$

which is canonical ([16]) in the sense that

$$
[X(2), \cdots, X(m)]=\left[W\left(h_{2}\right), \cdots, W\left(h_{m}\right)\right] \quad \text { for every } m \geq 2
$$

To be more precise, we obtain the exact expression of $W$ in terms of $X$ :

$$
\begin{equation*}
W\left(h_{m}\right)=\sum_{x \mid m} \mu(m / x) X(x) . \tag{6}
\end{equation*}
$$

The proof of (5) consists of an application of the inversion formula (5) to a general relation

$$
\sum_{x=2}^{\infty} f(x) X(x)=\sum_{m=2}^{\infty}(R f)\left(h_{m}\right) W\left(h_{m}\right) .
$$

The rest of this section concentrates of the case of the $n$-sphere ( $M, m$ ) $=\left(S^{n}, \sigma\right)$. For each $\rho \in(0,2 \pi)$, set

$$
C_{\rho}(p):=\left\{x \in S^{n} ;(x, p)>\cos (\rho / 2)\right\} .
$$

This is an open cap with north pole $p \in S^{n}$ and in particular $C_{\pi}(p)$ is the hemisphere. Take $H_{\rho}:=\left\{C_{\rho}(p) ; p \in S^{n}\right\}$ with an invariant measure

$$
d \nu\left(C_{\rho}(p)\right)=c d \sigma(p), \quad c=\nu\left(H_{o}\right)>0 .
$$

Then the corresponding distance becomes

$$
\begin{equation*}
d_{\rho}(x, y):=c \int_{S^{n}} \pi_{C_{\rho}(p)}(x, y) \sigma(d p)=c \sigma\left(C_{\rho}(x) \triangle C_{\rho}(y)\right) \tag{7}
\end{equation*}
$$

which is rotation-invariant and hence of the form $c r_{\rho}\left(d_{G}(x, y)\right)$, where $d_{G}(x, y):=\arccos (x, y)$. Since, $\pi_{C_{2 \pi-\rho}(p)}=\pi_{C_{\rho}(-p)}$, we have $r_{2 \pi-\rho}(t) \equiv r_{\rho}(t)$. Furthermore, a straightforward computation $(\rho=\pi)$ yields the explicit form of $r_{\pi}: r_{\pi}(t)=t / \pi$ (cf. [18] and [21]).

A Lévy's Brownian motion $X$ with parameter space ( $S^{n}, d_{\rho}$ ) is then expressed in the form

$$
X(x)=\sqrt{c} \int_{C_{p}(x)} W_{0}(d y)
$$

where $W_{0}=\left\{W_{0}(d y) ; y \in S^{n}\right\}$ is a Gaussian random measure based on the uniform probability space ( $S^{n}, \sigma$ ). Instead of the pair of $R$ and $R^{*}$ associated with (1), it is more convenient to treat the following transform associated with ( $1^{\prime}$ ):

$$
\begin{equation*}
\left(R_{\rho} f\right)(x):=\int_{C_{\rho}(x)} f(y) \sigma(d y) \tag{8}
\end{equation*}
$$

which is a self-adjoint operator on $L^{2}\left(S^{n}, \sigma\right)$. The expression ( $1^{\prime}$ ) as well as the transform $R_{\rho}$ will be further discussed in Section 4.

In the one-dimensional case $n=1$, we can go further by making a superposition of $\left\{d_{\rho}: 0<\rho \leq \pi\right\}$ :

$$
\begin{equation*}
d(x, y):=\int_{(0, \pi]} d_{\rho}(x, y) \mu(d \rho) \tag{9}
\end{equation*}
$$

where $\mu$ is a probability measure on $(0, \pi]$. A measure space $(H, \nu)$ combined with this $d$ is obviously taken as follows:
$H:=\left\{h_{\rho, p}:=C_{\rho}(p) ; 0<\rho \leq \pi, p \in S^{1}\right\} \quad$ with $\quad \nu\left(d h_{\rho, p}\right)=c \mu(d \rho) \sigma(d p)$.
Observe that the rotation-invariant distance $d$ on $S^{1}$ takes the form $d(x, y)=r\left(d_{G}(x, y)\right)$, where

$$
r(t)=c \int_{(0, \pi]} r_{\rho}(t) \mu(d \rho)=2 c \int_{(0, \pi]} \min (t, \rho) \mu(d \rho) .
$$

The right derivative $r_{+}^{\prime}(t)$ is of the form $2 c \mu((t, \pi])$ and therefore nonincreasing in $0 \leq t<\pi$.

What we have just observed is summed up in the following
Proposition 2. Suppose that $r(t)$ is a continuous function on $[0, \pi]$, $r(0)=0$ and has the right derivative $r_{+}^{\prime}(t) \geq 0$, non-increasing on $[0, \pi)$. Then the distance $d(x, y):=r\left(d_{G}(x, y)\right)$ on $S^{1}$ is $L^{1}$-embeddable.

## §3. Generalized Radon transform and its dual

This section is devoted to the study of basic properties of the generalized Radon transform $R$ and the dual Radon transform $R^{*}$. The main fact we prove is the singular value decomposition of $R^{*}$ regarded as a Hilbert-Schmidt operator from $L^{2}(H, \nu)$ to $L^{2}(M, \alpha(x) m(d x))$, where the density $\alpha(x)$ is chosen from among positive functions in $L^{1}(M, m)$ satisfying

$$
\int_{M} \nu\left(B_{x}\right) \alpha(x) m(d x):=C<\infty .
$$

The decomposition of $R^{*}$ implies the Karhunen-Loève expansion of a Lévy's Brownian motion $X$ with parameter space ( $M, d$ ).

We shall begin by discussing the covariance operator of $X$. The representation (1) of $X$ implies that the covariance function $\Gamma(x, y):=$ $E[X(x) X(y)]$ is equal to $\nu\left(B_{x} \cap B_{y}\right)$. With a choice of $\alpha$ mentioned above, we consider the Hilbert space $L^{2}(M, \tilde{m}), \tilde{m}(d x):=\alpha(x) m(d x)$, instead of the usual $L^{2}(M, m)$. Then, the equation

$$
\begin{equation*}
(\Gamma f)(x)=\int_{M} \Gamma(x, y) f(y) \tilde{m}(d y), \tag{10}
\end{equation*}
$$

defines a positive, self-adjoint and trace class operator on $L^{2}(M, \tilde{m})$ (cf. [5], p. 294). The operator $\Gamma$ is called the covariance operator of $X$.

We next consider the generalized Radon transform $R$. Observe that multiplication by $\alpha$ is a well-defined operator from $L^{2}(M, \tilde{m})$ to $L^{1}(M, m)$ :

$$
\left(T_{\alpha} f\right)(x):=\alpha(x) f(x), \quad f \in L^{2}(M, \tilde{m})
$$

So we can form the composition $R \circ T_{\alpha}$ to infer that it is a bounded operator from $L^{2}(M, \tilde{m})$ to $L^{2}(H, \nu)$. The proof of this assertion is an easy computation:

$$
\begin{aligned}
& \left\|\left(R \circ T_{\alpha} f\right)(h)\right\|_{L^{2}(H, \nu)}^{2} \\
& \quad \leq \int_{H} \nu(d h)\left\{\iint_{M^{2}} \chi_{h}(x) \chi_{h}(y)|f(x) \| f(y)| \tilde{m}(d x) \tilde{m}(d y)\right\} \\
& \quad=\iint_{M^{2}} \nu\left(B_{x} \cap B_{y}\right)|f(x) \| f(y)| \tilde{m}(d x) \tilde{m}(d y) \\
& \quad \leq\left\{\int_{M} \sqrt{\left.\nu\left(B_{x}\right)|f(x)| \tilde{m}(d x)\right\}^{2} \leq C\|f\|_{L^{2}(M, \tilde{m})}^{2} .}\right.
\end{aligned}
$$

A similar argument implies that the dual Radon transform $R^{*}$ is bounded from $L^{2}(H, \nu)$ to $L^{2}(M, \tilde{m})$. We need one more step to get at the following

Theorem 3. We have a factorization of $\Gamma$ :

$$
\begin{equation*}
\Gamma=R^{*} \circ\left(R \circ T_{a}\right) . \quad L^{2}(M \circ \tilde{m}) \stackrel{\Gamma}{\longrightarrow} L^{2}(M, \tilde{m}) \tag{11}
\end{equation*}
$$

The proof of (11) is immediate:

$$
\begin{aligned}
\left(R^{*} \circ R \circ T_{a} f\right)(x) & =\int_{M} f(y) \tilde{m}(d y)\left\{\int_{H} \chi_{B_{x}}(h) \chi_{B_{y}}(h) \nu(d h)\right\} \\
& =\int_{M} \Gamma(x, y) f(y) \tilde{m}(d y)=(\Gamma f)(x), \quad f \in L^{2}(M, \tilde{m})
\end{aligned}
$$

We are now going to give the singular value decompositions of the two factors, $R \circ T_{\alpha}$ and $R^{*}$, in Theorem 3. Positive eigenvalues $\lambda_{i}^{2}$ of the covariance operator $\Gamma$ is enumerated by means of index $i \in I$, where $I$ is a finite or countable infinite set and $\left\{\lambda_{i}\right\} \in l^{2}(I)$. Set

$$
N_{0}:=\left\{f \in L^{2}(M, \tilde{m}) ;(\Gamma f)(x) \equiv 0, x \in M\right\}, \quad \text { the null space of } \Gamma .
$$

Then we can select in $N_{o}^{\perp}$ a CONS $\left\{f_{i}(x) ; i \in I\right\}$ consisting of eigenfunctions of $\Gamma$ :

$$
\left(\Gamma f_{i}\right)(x)=\lambda_{i}^{2} f_{i}(x), \quad i \in I
$$

Note that any non-negative function in $L^{2}(M, \tilde{m})$ cannot be in $N_{0}$ since $\Gamma(x, y) \geq 0$ for all $x, y \in M$.

Now, put

$$
g_{i}(h):=\left(R \circ T_{\alpha} f_{i}\right)(h) / \lambda_{i} \in L^{2}(H, \nu),
$$

and $N_{1}:=\left[g_{i} ; i \in I\right]^{\perp}$, where $\left[g_{i} ; i \in I\right]$ stands for the closed linear span of $\left\{g_{i} ; i \in I\right\}$ in $L^{2}(H, \nu)$. The functions $g_{i}(h), i \in I$, constitute a CONS in $N_{1}^{\perp}$; the proof of this assertion is carried out by using Theorem 3:

$$
\begin{aligned}
\left(g_{i}, g_{j}\right)_{L^{2}(H, \nu)} & =\left(\lambda_{i} \lambda_{j}\right)^{-1}\left(R^{*} \circ R \circ T_{\alpha} f_{i}, f_{j}\right)_{L^{2}(M, \tilde{m})} \\
& =\left(\lambda_{i} \lambda_{j}\right)^{-1}\left(\Gamma f_{i}, f_{j}\right)_{L^{2}(M, \tilde{m})}=\lambda_{i} \lambda_{j}^{-1}\left(f_{i}, f_{j}\right)_{L^{2}(M, \tilde{m})}=\delta_{i, j} .
\end{aligned}
$$

For our purpose we need the following
Lemma 4. We have an expansion

$$
\begin{equation*}
\chi_{h}(x)=\chi_{B_{x}}(h)=\sum_{i \in I} \lambda_{i} f_{i}(x) g_{i}(h), \quad x \in M \quad \text { and } \quad h \in H \tag{12}
\end{equation*}
$$

Proof. We write the Fourier series of $\chi_{B_{x}}(h)$ as an element of $L^{2}(H, \nu)$ :

$$
\chi_{B_{x}}(h)=\sum_{i \in I} c_{i}(x) g_{i}(h)+g^{0}(h),
$$

where $c_{i}(x):=\left(\chi_{B_{x}}(h), g_{i}(h)\right)_{L^{2}(H, \nu)}=\left(R^{*} g_{i}\right)(x)$ and $g^{0} \in N_{1}$. Since

$$
\left(R^{*} g^{0}, f_{i}\right)_{L^{2}(M, \tilde{m})}=\left(g^{0}, R \circ T_{\alpha} f_{i}\right)_{L^{2}(H, \nu)}=\lambda_{i}\left(g^{0}, g_{i}\right)_{L^{2}(H, \nu)}=0
$$

for any $i \in I$, we have $R^{*} g^{0} \in N_{0}$. Actually this function $\left(R^{*} g^{0}\right)(x)$ is constantly equal to 0 , because it is a non-negative function in $N_{0}$ :

$$
\left(R^{*} g^{0}\right)(x)=\left(\chi_{B_{x}}(h), g^{0}(h)\right)_{L^{2}(H, \nu)}=\left\|g^{0}\right\|_{L^{2}(H, \nu)}^{2} \geq 0 .
$$

We have thus proved that $g^{0}(h) \equiv 0$.
The next task is to calculate the Fourier coefficients $c_{i}(x)=\left(R^{*} g_{i}\right)(x)$, $i \in I$. Since

$$
\left(R^{*} g_{i}, f^{0}\right)_{L^{2}(M, \tilde{m})}=\left(g_{i}, R \circ T_{a} f^{0}\right)_{L^{2}(H, \nu)}=\left(f_{i}, \Gamma f^{0}\right)_{L^{2}(M, \tilde{m})} / \lambda_{i}=0
$$

for any $f^{0} \in N_{0}$, we have $R^{*} g_{i} \in N_{0}^{\perp}$. Furthermore, the equality

$$
\left(R^{*} g_{i}, f_{j}\right)_{L^{2}(M, \tilde{m})}=\left(g_{i}, R \circ T_{\alpha} f_{j}\right)_{L^{2}(H, \nu)}=\lambda_{j} \delta_{i, j}
$$

shows that $c_{i}(x)=\lambda_{i} f_{i}(x)$, which completes the proof. We note that (12) is also the Fourier series of $\chi_{h}(x)$ as an element of $L^{2}(M, \tilde{m})$.

In view of the expressions

$$
\left(R \circ T_{\alpha} f\right)(h)=\left(\chi_{h}(x), f(x)\right)_{L^{2}(M, \tilde{m})} \quad \text { and } \quad\left(R^{*} g\right)(x)=\left(\chi_{B_{x}}(h), g(h)\right)_{L^{2}(H, \nu)},
$$

Lemma 4 immediately gives us their singular value decompositions having common positive singular values $\left\{\lambda_{i} ; i \in I\right\} \in l^{2}(I)$.

Theorem 5. (i) The operator $R \circ T_{\alpha}$ is a Hilbert-Schmidt operator from $L^{2}(M, \tilde{m})$ to $L^{2}(H, \nu)$ and has the singular value decomposition

$$
\begin{equation*}
\left(R \circ T_{\alpha} f\right)(h)=\sum_{i \in I} \lambda_{i}\left(f, f_{i}\right)_{L^{2}(M, \tilde{m})} g_{i}(h) \tag{13}
\end{equation*}
$$

The null space of $R \circ T_{\alpha}$ is $N_{0}=\left[f_{i} ; i \in I\right]^{\perp}$.
(ii) The dual Radon transform $R^{*}$ is a Hilbert-Schmidt operator from $L^{2}(H, \nu)$ to $L^{2}(M, \tilde{m})$ and has the singular value decomposition

$$
\begin{equation*}
\left(R^{*} g\right)(x)=\sum_{i \in I} \lambda_{i}\left(g, g_{i}\right)_{L^{2}(H, \nu)} f_{i}(x) \tag{14}
\end{equation*}
$$

The null space of $R^{*}$ is $N_{1}=\left[g_{i} ; i \in I\right]^{\perp}$.
An application of Theorem 5 (ii) to the representation (1) is now in order. Let us define

$$
\xi_{i}:=\int_{H} g_{i}(h) W(d h),
$$

to get an i.i.d. sequence $\xi=\left\{\xi_{i} ; i \in I\right\}$ of standard Gaussian random variables. Since (1) is rewritten as $X(x)=\left(R^{*} W\right)(x)$, the decomposition (14) yields

$$
\begin{equation*}
X(x)=\sum_{\imath \in I} \lambda_{i} \xi_{i} f_{i}(x) \tag{15}
\end{equation*}
$$

which is nothing but the Karhunen-Loéve expansion usually derived from Mercer's theorem (cf. [5] and [17]):

$$
\begin{equation*}
\Gamma(x, y)=\sum_{i \in I} \lambda_{i}^{2} f(x) f_{i}(y) \tag{16}
\end{equation*}
$$

Moreover, orthonormality of the system $\left\{f_{i} ; i \in I\right\}$ in $L^{2}(M, \tilde{m})$ implies the inverse expression of $\xi$ in terms of $X$ :

$$
\xi_{i}=\left(X(x), f_{i}(x)\right)_{L^{2}(M, \tilde{m})} / \lambda_{i} .
$$

Summing up what we have just proved, we get
Theorem 6. Every Lévy's Brownian motion $X$ with parameter space ( $M, d$ ) admits of the Karhunen-Loève expansion (15) in terms of $\xi$, and moreover we have

$$
\begin{equation*}
[X(x) ; x \in M]=\left[\xi_{i} ; i \in I\right]=\left\{\int_{H} g(h) W(d h) ; g \in N_{\frac{1}{1}}^{\perp}\right\} \tag{17}
\end{equation*}
$$

As a direct consequence of (15), we obtain another useful expression of the $L^{1}$-embeddable metric $d(x, y)=\nu\left(B_{x} \triangle B_{y}\right)$ on $M$ :

$$
\begin{equation*}
d(x, y)=\sum_{\imath \in I} \lambda_{i}^{2}\left(f_{i}(x)-f_{i}(y)\right)^{2}, \tag{18}
\end{equation*}
$$

which is equivalent to (16).

## §4. Lévy's Brownian motion with parameter space ( $\boldsymbol{S}^{n}, d_{\rho}$ )

The final section concerns the concrete examples on $S^{n}$ discussed in Section 2. We shall calculate explicitly the eigenvalues and eigenfunctions of the self-adjoint operator $R_{\rho}$ on $L^{2}\left(S^{n}, \sigma\right)$; then an application of the decomposition of $R_{\rho}$ to the expression ( $1^{\prime}$ ) will yield a new representation for a Lévy's Brownian motion $X$ with parameter space ( $S^{n}, d_{\rho}$ ). In addition, we shall investigate the $M(t)$-process of $X$ introduced by Lévy [20].

We recall some known facts about spherical harmonics (cf. [10] and [33]). Let $S H_{m}$ denote the set of all spherical harmonics of degree $m$; then the dimension of $S H_{m}$ is

$$
h(m):=\frac{2 m+n-1}{m+n-1}\binom{m+n-1}{m} .
$$

We get the direct sum decomposition $L^{2}\left(S^{n}, \sigma\right)=\sum_{m=0}^{\infty} \oplus S H_{m}$, as well as a CONS $\left\{S_{m, k}(x) ;(m, k) \in \Delta\right\}$ consisting of spherical harmonics, where $\Delta:=$ $\{(m, k) ; m \geq 0$ and $1 \leq k \leq h(m)\}$. In the sequel we shall make use of the addition formula

$$
\frac{1}{h(m)} \sum_{k=1}^{h(m)} S_{m, k}(x) S_{m, k}(y)=C_{m}^{\lambda}((x, y)) / C_{m}^{\lambda}(1)
$$

where $C_{m}^{2}(u)$ is the Gegenbauer polynomial of degree $m$ with $\lambda:=(n-1) / 2$.
Let us proceed to prove the explicit form of (12) in the present situation where $d=d_{\rho}$ and $(M, m)=\left(S^{n}, \sigma\right) \sim\left(H_{\rho}, \nu\right)$ by the mapping $x \in S^{n}$ $\mapsto C_{\rho}(x) \in H_{\rho}$.

Lemma 7 (cf. [32]). We have an expansion

$$
\begin{align*}
\chi_{C_{\rho}(x)}(y) & =\sum_{(m, k) \in \Delta} \lambda_{m}(\rho) S_{m, k}(x) S_{m, k}(y)  \tag{19}\\
& =\sum_{m=0}^{\infty} \lambda_{m}(\rho) h(m) C_{m}^{2}((x, y)) / C_{m}^{\lambda}(1)
\end{align*}
$$

where

$$
\lambda_{m}(\rho)=\left\{\begin{array}{l}
\frac{\left|S^{n-1}\right|}{\left|S^{n}\right|} \int_{\cos (\rho / 2)}^{1}\left(1-u^{2}\right)^{\lambda-1 / 2} d u, \quad m=0 \\
\frac{\left|S^{n-1}\right|}{\left|S^{n}\right| n} \frac{C_{m+1}^{\lambda+1}(\cos (\rho / 2))}{C_{m-1}^{\lambda+1}(1)} \sin ^{2 \lambda+1}(\rho / 2), \quad m \geq 1
\end{array}\right.
$$

Proof. Appealing to the Funk-Hecke theorem ([10]), we have (19) with

$$
\lambda_{m}(\rho)=\frac{\left|S^{n-1}\right|}{\left|S^{n}\right|} \int_{\cos (\rho / 2)}^{1} \frac{C_{m}^{\lambda}(u)}{C_{m}^{2}(1)}\left(1-u^{2}\right)^{\lambda-1 / 2} d u
$$

To compute the integral $\int_{\cos (\rho / 2)}^{1} C_{m}^{\lambda}(u)\left(1-u^{2}\right)^{\lambda-1 / 2} d u$ for $m \geq 1$, we use the formula

$$
\begin{aligned}
& C_{m}^{2}(u)=b_{m}^{2}\left(1-u^{2}\right)^{-\lambda+1 / 2} \frac{d^{m}}{d u^{m}}\left(1-u^{2}\right)^{m+\lambda-1 / 2} \\
& b_{m}^{\lambda}=(-1)^{m}(2 \lambda)_{m} /(2 m)!!(\lambda+1 / 2)_{m}
\end{aligned}
$$

where $(a)_{m}:=\prod_{j=0}^{m-1}(a+j)$. Since

$$
\begin{aligned}
\lambda_{m}(\rho) & =\frac{\left|S^{n-1}\right|}{\left|S^{n}\right|} \frac{b_{m}^{2}}{C_{m}^{\lambda}(1)}\left[\frac{d^{m-1}}{d u^{m-1}}\left(1-u^{2}\right)^{m+\lambda-1 / 2}\right]_{\cos (\rho / 2)}^{1} \\
& =\frac{\left|S^{n-1}\right| 2 \lambda}{\left|S^{n}\right| m(2 \lambda+m) C_{m}^{\lambda}(1)} C_{m-1}^{\lambda+1}(\cos (\rho / 2)) \sin ^{2 \lambda+1}(\rho / 2) \\
& =\frac{\left|S^{n-1}\right|}{\left|S^{n}\right| n} \frac{C_{m-1}^{\lambda+1}(\cos (\rho / 2))}{C_{m-1}^{\lambda+1}(1)} \sin ^{2 \lambda+1}(\rho / 2),
\end{aligned}
$$

the proof is completed.
The generalized (or dual) Radon transform $R_{\rho}$ associated with (1') is a self-adjoint and Hilbert-Schmidt operator on $L^{2}\left(S^{n}, \sigma\right)$, and the factorization of the covariance operator $\Gamma$ (Theorem 3) takes the simpler form

$$
\Gamma=\left(\sqrt{c} R_{\rho}\right)^{2} . \quad \begin{aligned}
& L^{2}\left(S^{n}, \sigma\right) \stackrel{\Gamma}{\longmapsto} L^{2}\left(S^{n}, \sigma\right) \\
& \sqrt{c} R_{\rho} \bigvee_{L^{2}\left(S^{n}, \sigma\right)} \int^{\sqrt{c} R_{\rho}} \\
& \hline
\end{aligned}
$$

In order to state the decomposition of $R_{\rho}$, we set

$$
\Delta_{\rho}:=\left\{(m, k) \in \Delta ; \lambda_{m}(\rho)=0\right\}=\left\{(m, k) \in \Delta ; m \geq 2, C_{m-1}^{\lambda+1}(\cos (\rho / 2))=0\right\}
$$

which corresponds to the null space $N$ of $R_{\rho}$, and $I_{\rho}:=\Delta \backslash \Delta_{\rho}$. Recalling that $\left(R_{\rho} f\right)(x)=\left(\chi_{C_{\rho}(x)}(y), f(y)\right)_{L^{2}\left(s^{n}, \sigma\right)}$, Lemma 7 implies the following

Theorem 8. We have

$$
\begin{equation*}
\left(R_{\rho} f\right)(x)=\sum_{(m, k) \in I_{\rho}} \lambda_{m}(\rho)\left(f, S_{m, k}\right)_{L^{2}\left(S^{n}, o\right)} S_{m, k}(x), \tag{20}
\end{equation*}
$$

and the null space $N=\left[S_{m, k}(x) ;(m, k) \in \Delta_{\rho}\right]$.
By applying (20) to the expression $X(x)=\sqrt{c}\left(R_{\rho} W_{0}\right)(x)$, we obtain the Karhunen-Loève expansion of $X$ in terms of the i.i.d. sequence

$$
\xi=\left\{\xi_{m, k}:=\int_{S^{n}} S_{m, k}(y) W_{0}(d y) ;(m, k) \in I_{\rho}\right\}
$$

of standard Gaussian random variables.
Theorem 9. A Lévy's Brownian motion $X$ with parameter space ( $S^{n}, d_{\rho}$ ) admits of a representation

$$
\begin{equation*}
X(x)=\sqrt{c} \sum_{(m, k) \in I_{\rho}} \lambda_{m}(\rho) \xi_{m, k} S_{m, k}(x) . \tag{21}
\end{equation*}
$$

Moreover, we have

$$
\begin{equation*}
\left[X(x) ; x \in S^{n}\right]=\left[\xi_{m, k} ;(m, k) \in I_{\rho}\right]=\left\{\int_{S_{n}} g(y) W_{0}(d y) ; g \in N^{\perp}\right\} \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
d_{\rho}(x, y)=2 c \sum_{m=1}^{\infty} \lambda_{m}^{2}(\rho) h(m)\left\{1-C_{m}^{\lambda}((x, y)) / C_{m}^{\lambda}(1)\right\} \tag{23}
\end{equation*}
$$

We now focus our attention on the case $\rho=c=\pi$, i.e., $X$ is a Lévy's Brownian motion with the geodesic distance $d_{G}$. In this case,

$$
\Delta_{\pi}=\{(2 j, k) ; j=1,2, \cdots \text { and } 1 \leq k \leq h(2 j)\}
$$

and

$$
\lambda_{m}(\pi)=\left\{\begin{array}{l}
1 / 2, \quad m=0 \\
\frac{\Gamma((n+1) / 2)(n-2)!!}{\Gamma(n / 2) \sqrt{\pi}}(-1)^{j} \frac{(2 j-1)!!}{(2 j+n)!!}, \quad m=2 j+1
\end{array}\right.
$$

With the help of these values, one can compute the coefficients $2 \pi \lambda_{m}^{2}(\pi) h(m)$ in (23), to find the formula of $d_{G}$ due to Gangoli [13] and Molčan [25] (cf. also [27], p. 143) who proved it via an entirely different approach. In view of the special form of $\Delta_{\pi}$, it is natural to assume that $X(x)$ is odd, i.e., $X(x)+X(-x)=0$; the expression (21) then becomes

$$
\begin{align*}
X(x)= & \frac{\Gamma((n+1) / 2)(n-2)!!}{\Gamma(n / 2)} \sum_{j=0}^{\infty}(-1)^{j} \frac{(2 j-1)!!}{(2 j+n)!!} \\
& \times \sum_{k=1}^{h(2 j+1)} \xi_{2 j+1, k} S_{2 j+1, k}(x) .
\end{align*}
$$

In connection with the $M(t)$-process ([20]), we need another transform $\left(T_{\rho} f\right)(x)$ (cf. [15]) which is given by the mean value of $f$ over the small (or great in case $\rho=\pi$ ) circle $\delta C_{\rho}(x), 0<\rho<2 \pi$.

Definition 3. For $f \in C\left(S^{n}\right)$, the set of all continuous functions on $S^{n}$, the integral transformation defined by

$$
\begin{equation*}
\left(T_{\rho} f\right)(x):=\int_{\partial C_{\rho}(x)} f(y) s(d y) / s\left(\delta C_{\rho}(x)\right), \tag{24}
\end{equation*}
$$

is called the mean value operator over $\delta C_{\rho}(x)$, where $s$ denotes the ( $n-1$ )dimensional surface measure on $\delta C_{\rho}(x)$. For each fixed $x_{0} \in S^{n}$, the Gaussian process

$$
\begin{equation*}
M(t):=\left(T_{2 t} X\right)\left(x_{0}\right)-X\left(x_{0}\right), \quad 0<t<\pi \tag{25}
\end{equation*}
$$

is called the $M(t)$-process.
By appealing, again, to the Funk-Hecke theorem, we get the decomposition of $T_{\rho}$.

Proposition 10. The mean value operator $T_{\rho}$ on $C\left(S^{n}\right)$ is extended to be a self-adjoint, compact operator on $L^{2}\left(S^{n}, \sigma\right)$, and it has the decomposition

$$
\begin{equation*}
\left(T_{\rho} f\right)(x)=\sum_{\langle m, k) \in I_{\rho}} \frac{C_{m}^{\lambda}(\cos (\rho / 2))}{C_{m}^{2}(1)}\left(f, S_{m, k}\right)_{L^{2}\left(S^{n}, \rho\right)} S_{m, k}(x), \tag{26}
\end{equation*}
$$

where $\tilde{I}_{\rho}:=\Delta \mid \tilde{\Lambda}_{\rho}$ and $\tilde{\Lambda}_{\rho}:=\left\{(m, k) \in \Delta ; C_{m}^{\lambda}(\cos (\rho / 2))=0\right\}$. Moreover, the null space of $T_{\rho}$ is $\left[S_{m, k}(x) ;(m, k) \in \tilde{\mathcal{I}}_{\rho}\right]$.

By the combination of (21) and (26), we write

$$
\begin{aligned}
M(t) & =\sqrt{c} \sum_{\langle m, k) \in I_{\rho}}\left\{\begin{array}{c}
C_{m}^{\lambda}(\cos t) \\
C_{m}^{\lambda}(1)
\end{array}-1\right\} \lambda_{m}(\rho) \xi_{m, k} S_{m, k}\left(x_{0}\right) \\
& =\sqrt{c} \sum_{m \in J_{\rho}} \lambda_{m}(\rho) \sqrt{ } h(m) \eta_{m}\left\{1-C_{m}^{\lambda}(\cos t) / C_{m}^{\lambda}(1)\right\},
\end{aligned}
$$

where we have put

$$
\eta_{m}:=\frac{-1}{\sqrt{h(m)}} \sum_{k=1}^{h(m)} \xi_{m, k} S_{m, k}\left(x_{0}\right) \quad \text { for } m \in J_{\rho}:=\left\{m \geq 1 ; \lambda_{m}(\rho) \neq 0\right\}
$$

It is shown that the $\eta_{m}$ form an i.i.d. sequence of standard Gaussian random variables.

Proposition 11 (cf. [27] in case $\rho=\pi$ ). The $M(t)$-process of a Lévy's Brownian motion $X$ with parameter space $\left(S^{n}, d_{\rho}\right)$ is expressed in the form

$$
\begin{equation*}
M(t)=\sqrt{c} \sum_{m \in J_{\rho}} \lambda_{m}(\rho) \sqrt{h(m)} \eta_{m}\left\{1-C_{m}^{\lambda}(\cos t) / C_{m}^{\lambda}(1)\right\} \tag{27}
\end{equation*}
$$

## References

[1] P. Assouad, Produit tensoriel, distances extrémales et realisation de covariance, I et II, C. R. Acad. Sci. Paris Ser. A, 288 (1979), 649-652 et 675-677.
[2] P. Assouad et M. Deza, Espaces métriques plongeables dans un hypercube: Aspects combinatories, Ann. Discrete Math., 8 (1980), 197-210.
[3] P. Assouad and M. Deza, Metric subspaces of $L^{1}$, Publications math. d’Orsay, Université de Paris-Sud, 1982.
[4] S. Campi, On the reconstruction of a function on a sphere by its integrals over great circles, Boll. Un. Math. Ital. (5), 18-c (1981), 195-215.
[5] P. Cartier, Une étude des covariances measurables, Math. Analysis and Applications, Part A, Advances in Math. Supplementary Studies, 7A (1981), 267-316.
[6] N. N. Chentsov, Lévy Brownian motion for several parameters and generalized white noise, Theory Probab. Appl., 2 (1957), 265-266 (English translation).
[7] A. M. Cormack and E. T. Quinto, A Radon transform on spheres through the origin in $R^{n}$ and applications to the Darboux equation, Trans. Amer. Math. Soc., 260 (1980), 575-581.
[8] S. R. Deans, The Radon transform and some of its applications, John Wiley \& Sons, New York, 1983.
[9] P. Diaconis and R. L. Graham, The Radon transform on $\boldsymbol{Z}_{2}^{k}$, Pacific J. Math., 118 (1985), 323-345.
[10] A. Erdélyi, W. Magnus, F. Oberhettinger and F. G. Tricomi, Higher transcendental functions (Bateman manuscript project), Vol. II., McGraw-Hill, New York, 1953.
[11] P. Funk, Über Flächen mit lauter geschlossenen geodätischen Linien, Math. Ann., 74 (1913), 278-300.
[12] _- Über eine geometrische Anwendung der Abelschen Integralgleichung, Math. Ann., 77 (1916), 129-135.
[13] R. Gangolli, Positive definite kernels on homogeneous spaces and certain stochastic processes related to Lévy's Brownian motion of several parameters, Ann. Inst. H. Poincarè Sect. B, 3 (1967), 121-225.
[14] R. L. Graham, On isometric embeddings of graphs, in Progress in graph theory (J. A. Bondy and U.S.R. Murty, ed.), Academic Press, New York, 1984, 307-322.
[15] S. Helgason, The Radon transform, Birkhäuser, Boston, 1980.
[16] T. Hida and M. Hitsuda, Gaussian processes (in Japanese), Kinokuniya, Tokyo, 1976.
[17] K. Karhunen, Über lineare Methoden in der Wahrscheinlichkeitsrechnung, Ann. Acad. Sci. Fennicae Ser. A, I. Math. Phys., 37 (1947), 79 pp.
[18] J. B. Kelly, Hypermetric spaces, in The geometry of metric and linear spaces, Lecture Notes in Math., No. 490, Springer-Verlag, Berlin, 1975, 17-31.
[19] P. Lévy, Problèmes concrets d'analyse fonctionnelle, Gauthier-Villars, Paris, 1951.
[20] ——, Processus stochastiques et mouvement brownien, Gauthier-Villars, Paris, 1965.
[21] ——, Le mouvement brownien fonction d'un point de la sphère de Riemann, Rend. Circ. Mat. Palermo (2), 8 (1959), 1-14.
[22] M. A. Lifshits, On representation of Lévy's fields by indicators, Theory Probab. Appl., 24 (1979), 629-633 (English translation).
[23] D. Ludwig, The Radon transform on Euclidean space, Comm. Pure Appl. Math., 19 (1966), 49-81.
[24] H. P. McKean, Brownian motion with a several-dimensional time, Theory Probab. Appl., 8 (1963), 335-354.
[25] G. M. Molčan, The Markov property of Lévy fields on spaces of constant curvature, Soviet Math. Dokl., 16 (1975), 528-532.
[26] E. A. Morozova and N. N. Chentsov, P. Lévy's random fields, Theory Probab. Appl., 13 (1968), 153-156 (English translation).
[27] A. Noda, Lévy's Brownian motion; Total positivity structure of $M(t)$-process and deterministic character, Nagoya Math. J., 94 (1984), 137-164.
[28] E. T. Quinto, The dependence of the generalized Radon transform on defining measures, Trans. Amer. Math. Soc., 257 (1980), 331-346.
[29] -_, Null spaces and ranges for the classical and spherical Radon transforms, J. Math. Anal. Appl., 90 (1982), 408-420.
[30] -, Singular value decompositions and inversion methods for the exterior Radon transform and a spherical transform, J. Math. Anal. Appl., 95 (1983), 437-448.
[31] J. Radon, Über die Bestimmung von Funktionen durch ihre Integralwerte längs gewisser Mannigfaltigkeiten, Ber. Verh. Sächs. Akad. Wiss. Leipzig, Math.-Nat. kl., 69 (1917), 262-277.
[32] R. Schneider, Über eine Integralgleichung in der Theorie der konvexen Körper, Math. Nachr., 44 (1970), 55-75.
[33] R. T. Seeley, Spherical harmonics, Amer. Math. Monthly, 73 (1966), 115-121.
[34] R. S. Strichartz, Radon inversion-variations on a theme, Amer. Math. Monthly, 89 (1982), 377-384 and 420-423.
[35] S. Takenaka, I. Kubo and H. Urakawa, Brownian motion parametrized with metric spaces of constant curvature, Nagoya Math. J., 82 (1981), 131-140.

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