ASYMPTOTIC BEHAVIOR OF OSCILLATING RADIAL SOLUTIONS TO CERTAIN NONLINEAR EQUATIONS*

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1. Introduction. Recently, in the investigation of thin film problems, a class of oscillating radial solutions has attracted the attention. The radial solutions of concern satisfy the following initial value problem:

(1.1)
$$u'' + \frac{n-1}{r}u' = f(u) \quad \text{in } \mathbb{R}_+, \quad u(0) = \alpha > 0, \quad u'(0) = 0.$$

where $f \in C^1(0, \infty)$ satisfies the following general conditions:

- (i) f has a single zero t_0 in $(0, \infty)$ satisfying $f'(t_0) < 0$,
- (ii) f is nonincreasing near 0 and $\lim_{t\to 0^+} f(t) = \infty$.

The radial solutions is a special case of more general thin film problem in a bounded domain Ω in \mathbb{R}^n with Neumann boundary condition

(1.2)
$$\Delta u = f(u), \quad x \in \Omega, \quad \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial \Omega.$$

A typical example is that $f(u) = u^{-p} - \mu_1 u^{-q} - \mu_2$ with constants $p > \max(q, 0)$. This kind of semilinear equation appears in several applications in mechanics and physics. In particular, it has been used to model the dynamics of thin films for viscous liquids. Some detailed physics background can be found in [1]-[3], [9] and [12]-[14]. Some recent mathematical analysis can be found in [4, 6, 7, 8, 10, 11, 15, 17] and the references therein.

It was proved in [8] (see also [11]) that in dimension $N \geq 3$, for each $\alpha \in (0,t_0)$, (1.1) has a unique positive solution u_{α} . Moreover, u_{α} oscillates around the constant t_0 , that is, there is an increasing positive sequence $\{r_{\alpha}^n\}$ such that $\{r \in (0,\infty) : u'_{\alpha}(r) = 0\} = \{r_{\alpha}^n\}$, and $\lim_{n \to \infty} r_{\alpha}^n = \infty$. Here r_{α}^{2i+1} are local maxima of u_{α} with $u_{\alpha}(r_{\alpha}^{2i+1}) > t_0$ for any $i \in \mathbb{N}$; while r_{α}^{2i} are local minima with $u_{\alpha}(r_{\alpha}^i) < t_0$. It is also proven that there exists a singular (or so-called a rupture) radial solution $u_0(r)$ of (1.1) such that $u_0 \in C(\mathbb{R}^N)$, $u_0(0) = 0$, $u_0(r) > 0$ for $r \in (0,\infty)$ and $f(u_0) \in L^1_{loc}(\mathbb{R}^N)$. Moreover, any singular radial solution of (1.1) is oscillatory around t_0 and converges to t_0 as $r \to \infty$.

In [11], Jiang and Ni have showed the existence and uniqueness of radial rupture solution for $f(u) = u^{-p} - \mu_2$ in \mathbb{R}^N with $p, \mu_2 > 0$ and $N \geq 2$. Moreover, they proved that $\lim_{n\to\infty} r_{\alpha}^{n+1} - r_{\alpha}^n$ exits, and hence obtained the asymptotic formula for the length of oscillating interval, which depends only on p and μ_2 .

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It it a natural question that whether more accurate asymptotic behaviors of the radial solutions can be obtained. Similar question also arises in the study of oscillating radial solutions of Allen Cahn equation

(1.3)
$$\Delta u + u - u^3 = 0, \quad x \in \mathbb{R}^n.$$

Generally speaking, it is important to understand radial entire solutions since they display one of the fundamental structures near interesting points such as singularities of solutions. The asymptotic behavior of an entire solution at infinity is an essential property since it reveals the condition away from the singularity. There has been a lot of literature on the asymptotic behavior of positive radial solutions, see, for example [16], [5], [18], etc. However, it seems that little has been done on oscillating radial solutions.

In this note, we shall show an accurate asymptotic behavior of radial solutions for a vast class of equations in term of Bessel functions. In Section 1, some basic information on Bessel functions are collected; In Section 2, we show, for small initial data, the existence of oscillating radial solutions which behaves like Bessel functions at infinity; In Section 3 and 4, we shall show that the asymptotic behavior of radial solutions in term of Bessel functions for rather general equations including Allen-Cahn equations and thin-film equations.

2. Preliminaries.

2.1. Bessel Function. We recall that the Bessel Function

(2.1)
$$x^2y'' + xy' + (x^2 - \alpha^2)y = 0 \qquad x > 0.$$

It has two independent (linearly) solutions: Bessel function of the first kind and second kind, $J_{\alpha}(x)$, $Y_{\alpha}(x)$ with asymptotic expansions

(2.2)
$$J_{\alpha}(x) \sim \sqrt{\frac{2}{\pi x}} \left[\cos\left(x - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) \sum_{j=0}^{\infty} \frac{(-1)^{j}(\alpha, 2j)}{(2x)^{2j}} - \sin\left(x - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) \sum_{j=0}^{\infty} \frac{(-1)^{j}(\alpha, 2j+1)}{(2x)^{2j+1}} \right]$$

and

(2.3)
$$Y_{\alpha}(x) \sim \sqrt{\frac{2}{\pi x}} \left[\sin\left(x - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) \sum_{j=0}^{\infty} \frac{(-1)^{j}(\alpha, 2j)}{(2x)^{2j}} + \cos\left(x - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) \sum_{j=0}^{\infty} \frac{(-1)^{j}(\alpha, 2j+1)}{(2x)^{2j+1}} \right]$$

where

(2.4)
$$(\alpha, k) = \begin{cases} \frac{(4\alpha^2 - 1)(4\alpha^2 - 3^2)\cdots(4\alpha^2 - (2k - 1)^2)}{2^k k!}, & k > 1\\ 1, & k = 0. \end{cases}$$

Moreover, the Wronskian determinant associated is

$$(2.5) W(J_{\alpha}(x), Y_{\alpha}(x)) = \frac{2}{\pi x}.$$

2.2. Radial solution for a linear equation. Now consider the linear Laplace equation

(2.6)
$$\Delta u + \beta^2 u = 0, \qquad x \in \mathbb{R}^n,$$

its radial solution u(r) satisfies

(2.7)
$$u'' + \frac{n-1}{r}u' + \beta^2 u = 0 \qquad r = |x| > 0.$$

Let $u = r^{\frac{2-n}{2}}y(\beta r)$, then y(x) satisfies (2.1) with $\alpha = \frac{n-2}{2}$. Hence (2.7) has two linearly independent solutions

(2.8)
$$u_1(r) = r^{\frac{2-n}{2}} J_{\alpha}(\beta r), \qquad u_2(r) = r^{\frac{2-n}{2}} Y_{\alpha}(\beta r)$$

and the Wronskian determinant is

(2.9)
$$W(u_1(r), u_2(r)) = u_1 u_2' - u_2 u_1' = \frac{2}{\pi} r^{1-n}.$$

In particular, we have

(2.10)
$$u_1(r) = \frac{2}{\sqrt{\pi\beta}} r^{\frac{1-n}{2}} \left[\cos\left(\beta r - \frac{n-1}{4}\pi\right) + O\left(\frac{1}{r}\right) \right]$$
$$u_2(r) = \frac{2}{\sqrt{\pi\beta}} r^{\frac{1-n}{2}} \left[\sin\left(\beta r - \frac{n-1}{4}\pi\right) + O\left(\frac{1}{r}\right) \right].$$

3. Existence of Radial Solutions for Nonlinear Equation. Let us consider the nonlinear problems

(3.1)
$$\begin{cases} u'' + \frac{n-1}{r}u' + \beta^2 u + f(u) = 0, & r > 0 \\ u(r) \to 0, & \text{as } r \to \infty \end{cases}$$

where f(0) = f'(0) = 0.

We suppose that $f(u) \in C^{1,\sigma}(-\delta_0, \delta_0)$, for some $\delta_0 > 0$, $\sigma > \frac{2}{n-1}$ if n > 3, and $f(u) \in C^{2,\sigma}(-\delta_0, \delta_0)$, if n = 3, $\sigma > 0$.

It is easy to see that (3.1) is equivalent to

(3.2)
$$u(r) = u_1(r) \int_R^r \frac{\pi}{2} s^{n-1} u_2(s) (-f(u(s))) ds - u_2(r) \int_R^r \frac{\pi}{2} s^{n-1} u_1(s) (-f(u(s))) ds + \gamma_1 u_1(r) + \gamma_2 u_2(r),$$

or we can write

(3.3)

$$u(r) = \frac{\pi}{2} \int_{R}^{r} s^{n-1} f(u(s)) [u_2(r)u_1(s) - u_1(r)u_2(s)] ds + \gamma_1 u_1(r) + \gamma_2 u_2(r) \qquad r \ge R.$$

Then we have the following

THEOREM 3.1. When $|(\gamma_1, \gamma_2)|$ is small enough, (3.3) has a unique solution u_{γ_1, γ_2} such that

$$u_{\gamma_1,\gamma_2} = (\tilde{\gamma}_1 + o(1))u_1 + (\tilde{\gamma}_2 + o(1))u_2$$

for some $(\tilde{\gamma}_1, \tilde{\gamma}_2)$, provided that one of the following conditions holds (3.4)

$$\begin{cases} \mathbf{a}): f(u) \in C^{1,\sigma}(-\delta_0, \delta_0), \ f(0) = f'(0) = 0, & \sigma > \frac{2}{n-1} \ if \ n > 3 \ or \\ \mathbf{b}): f(u) \in C^{2,\sigma}(-\delta_0, \delta_0), \ f(0) = f'(0) = f''(0) = 0, & \sigma > 0 \ if \ n = 3 \ or \\ \mathbf{c}): f(u) \in C^{3,\sigma}(-\delta_0, \delta_0), \ f(0) = f'(0) = f''(0) = f^{(3)}(0) = 0, \ \sigma > 0 \ if \ n = 2. \end{cases}$$

Proof. We define a Banach space H_R for any R > 0.

(3.5)
$$H_R = \left\{ u \in C(R, \infty); \sup_{r > R} \left| u(r) r^{\frac{n-1}{2}} \right| < \infty \right\}$$

with a weighted norm

$$||u||_{H_R} = \sup_{r>R} |u(r)r^{\frac{n-1}{2}}|.$$

It is easy to see that $\| \|_{H_R}$ is a norm in H_R . Define now an operator in H_R for any (γ_1, γ_2) by (3.6)

$$K_{\gamma_1,\gamma_2}(u)(r) = \frac{\pi}{2} \int_R^r s^{n-1} f(u(s)) [u_2(r)u_1(s) - u_1(r)u_2(s)] ds + \gamma_1 u_1(r) + \gamma_2 u_2(r).$$

Then there exists σ_0 , δ_0 such that if $|\gamma_1| + |\gamma_2| < \sigma_0$, $\delta < \delta_0$, there holds

(3.7)
$$||K_{\gamma_1,\gamma_2}(u)||_{H_R} \le \delta$$
 when $||u||_{H_R} < \delta$.

In fact,

$$\begin{split} \|K_{\gamma_{1},\gamma_{2}}(u)\|_{H_{R}} &\leq \frac{4}{\pi\beta} \frac{\pi}{2} \sup_{r \geq R} \int_{R}^{r} s^{n-1} |f(u(s))| s^{-\frac{n-1}{2}} \mathrm{d}s + 2\sqrt{\frac{2}{\pi\beta}} (|\gamma_{1}| + |\gamma_{2}|) \\ &\leq 2\sqrt{\frac{2}{\pi\beta}} (|\gamma_{1}| + |\gamma_{2}|) + \frac{4}{\beta} \sup_{r \geq R} \int_{R}^{r} \left| s^{\frac{n-1}{2}} u(s) \right| \left| \frac{f(u(s))}{u(s)} \right| \mathrm{d}s \\ &\leq 2\sqrt{\frac{2}{\pi\beta}} (|\gamma_{1}| + |\gamma_{2}|) + \frac{4M}{\beta} \sup_{r \geq R} \int_{R}^{r} \left| s^{\frac{n-1}{2}} u(s) \right| |u(s)|^{\alpha} \mathrm{d}s \\ &\leq 2\sqrt{\frac{2}{\pi\beta}} (|\gamma_{1}| + |\gamma_{2}|) + \frac{4M}{\beta} \|u\|_{H_{R}}^{1+\alpha} \left(\sup_{r \geq R} \int_{R}^{r} s^{-\frac{n-1}{2}\alpha} \mathrm{d}s \right) \\ &\leq 2\sqrt{\frac{2}{\pi\beta}} (|\gamma_{1}| + |\gamma_{2}|) + \frac{C}{\beta} \|u\|_{H_{R}}^{1+\alpha} R^{1-\frac{n-1}{2}\alpha}, \end{split}$$

where $\alpha = \sigma$ in case **a)** in (3.4); $\alpha = 1 + \sigma$ in case **b)** in (3.4), and $\alpha = 2 + \sigma$ in case **c)** in (3.4) and

$$C = 4M/(\frac{n-1}{2}\alpha - 1).$$

Here we have used

(3.8)
$$\left| \frac{f(u)}{u^{1+\alpha}} \right| \le M,$$

$$\left| \frac{f(u) - f(v)}{u - v} \right| \le M \max\{|u|^{\alpha}, |v|^{\alpha}\},$$

for u near u = 0 and $\alpha > \frac{2}{n-1}$.

Therefore, when

(3.9)
$$|\gamma_1| + |\gamma_2| < \gamma_0 := \sqrt{\frac{\pi\beta}{2}} \frac{1}{4} \delta,$$

$$\delta < \delta_0 := \left(\frac{\beta}{4C}\right)^{\frac{1}{\alpha}} R^{-\frac{1}{\alpha} + \frac{n-1}{2}},$$

we have

$$||K_{\gamma_1,\gamma_2}(u)||_{H_R} \le \frac{\delta}{2} + \delta^{1+\alpha} \times \frac{C}{\beta} R^{1-\frac{n-1}{2}\alpha} \le \frac{\delta}{2} + \frac{\delta}{4} < \delta, \quad \text{when } ||u||_{H_R} < \delta.$$

Furthermore, we know K_{γ_1,γ_2} is a contracting mapping in $B_{\delta}(0) \subset H_R$, since

$$||K_{\gamma_{1},\gamma_{2}}(u) - K_{\gamma_{1},\gamma_{2}}(v)||_{H_{R}} \leq \frac{4}{\beta} \sup_{r \geq R} \int_{R}^{r} s^{\frac{n-1}{2}} |f(u(s)) - f(v(s))| ds$$

$$\leq \frac{4M}{\beta} ||u - v||_{H_{R}} \delta^{\alpha} \left(\sup_{r \geq R} \int_{R}^{r} s^{-\frac{n-1}{2}\alpha} ds \right)$$

$$\leq \frac{1}{2} ||u - v||_{H_{R}}.$$

Here we have used (3.8).

It is easy to see that (3.8) and (3.11) hold if (3.4) is provided. Therefore $K_{\gamma_1,\gamma_2}(u)$ has a fixed point in $B_{\delta}(0) \subset H_R$ when (γ_1,γ_2) and δ satisfies (3.9).

Furthermore, the integral in (3.3) is integrable with r replaced by ∞ . Hence Theorem 3.1 is proven with $(\tilde{\gamma}_1, \tilde{\gamma}_2)$ sufficiently small when R is fixed. Note that if we allow R to be large, the size of γ_1, γ_2 can be indeed large as long (3.9) holds. \square

4. Application to Allen-Cahn Equation. Consider the radial solution to Allen-Cahn equation

(4.1)
$$\begin{cases} u'' + \frac{n-1}{r}u' - F'(u) = 0, & r = |x|, \ x \in \mathbb{R}^n \\ u(0) = u_0, & |u_0| < 1. \end{cases}$$

where $F(u) \in C^{2,\sigma}([-1,1])$, and satisfies

(4.2)
$$\begin{cases} F'(1) = F'(-1) = 0, & F(1) = F(-1) = 0 \\ F(u) > 0 & \text{if } |u| < 1 \\ F'(0) = 0, & F''(0) < 0, & F'(u) < 0 & \text{if } 0 < u < 1 \\ F'(u) > 0 & \text{if } -1 < u < 0. \end{cases}$$

We have then

(4.3)
$$\left[\left(\frac{u'(r)}{2} \right)^2 - F(u(r)) \right]' = -\frac{n-1}{r} u'^2 \le 0, \quad \forall r > 0,$$

and thus

(4.4)
$$F(u(r)) \ge F(u_0) + \left(\frac{u'(r)}{2}\right)^2 > 0 \quad \forall r > 0.$$

Hence

$$(4.5) |u(r)| < 1, \forall r > 0.$$

Lemma 4.1. We have

$$\lim_{r \to \infty} u(r) = 0.$$

Proof. If not, we assume that there exists $r_N \to \infty$ such that

$$u(r_N) \to \alpha \neq 0, \qquad u'(r_N) \to 0 \text{ as } n \to \infty.$$

Define

$$u_n(r) = u(r + r_N) \qquad r > 0.$$

We know by the standard elliptic theory that

$$(4.7) |\nabla u| \le C, |\nabla^2 u| \le C, ||u||_{C^{2,\sigma}(\mathbb{R}^2)} < C.$$

Then $u_n(r) \to u_\infty(r)$ in $C^2(0,\infty)$ up to a subsequence, and $u_0(r)$ satisfies $|u_\infty(r)| < |F^{-1}(F(u_0))| < 1, \forall r \geq 0$, and

(4.8)
$$\begin{cases} u_{\infty}''(r) - F'(u_{\infty}(r)) = 0, & r > 0 \\ u_{\infty}(0) = \alpha \neq 0, & u_{\infty}'(0) = 0, & |\alpha| < 1. \end{cases}$$

Equation (4.8) can be solved easily by integration, the only solution are periodic solutions.

(4.9)
$$u_{\infty}(r) = P_{\alpha}(r)$$
 with
$$\begin{cases} P_{\alpha}(r + T_{\alpha}) = P_{\alpha}(r), & \text{for some } T_{\alpha} > 0 \\ P_{\alpha}(0) = \alpha. \end{cases}$$

Without loss of generality, we assume $\alpha > 0$. Let r_k be the k-th local maximum point of u(r) in $(0, \infty)$. Then $u(r_k)$ is decreasing in k, since

$$(4.10) F(u(r_{k+1})) \ge F(u(r_k)), u(r_k) > 0.$$

Hence $\lim_{k\to\infty} u(r_k) = \alpha > 0$. We can then prove

(4.11)
$$||u(r_{k+1}+r) - P_{\alpha}(r)||_{C^{2,\alpha}([0,T_{\alpha}])} \to 0 \quad \text{as } k \to \infty.$$

On the other hand, (4.3) leads to

$$(4.12) F(u(r_k)) - F(u(r_N)) = \int_{r_N}^{r_k} \frac{n-1}{r} |u'(r)|^2 dr$$

$$= (n-1) \sum_{i=N}^{k-1} \int_{r_i}^{r_{i+1}} \frac{(u'(r))^2}{r} dr$$

$$\geq \frac{n-1}{2} \int_0^{T_\alpha} |P'_\alpha(r)|^2 dr \sum_{i=N}^{k-1} \frac{1}{r_{i+1}},$$

when N is sufficiently large. Letting $k \to \infty$, we have a contradiction. Hence Lemma 4.1 is proven. \square

Let $m_k = u(r_k)$. We know that m_k is decreasing in k and tends to 0 as k goes to infinity. Next we shall show

LEMMA 4.2. For any fixed $\epsilon > 0$, we have

$$(4.13) m_k \le Ck^{-\frac{n-1}{2} + \epsilon}.$$

Hence

$$(4.14) u(r) \le Cr^{-\frac{n-1}{2} + \epsilon}, \quad r > 0$$

Proof. Let $\beta^2 = -F''(0)$. First we note $u(r_k + r)/m_k$ converges to $\cos(\beta r)$ uniformly in any bounded interval as k tends to infinity. Indeed we have

$$||u(r) - m_k \cos(\beta(r - r_k))||_{C^1[r_k r_k + T_0]} = o(1)m_k$$

where $T_0 = \frac{2\pi}{\beta}$ and

$$r_{k+1} - r_k = T_0 + o(1).$$

Furthermore, we have

$$F(0) - F(m_k) = \frac{1}{2}(\beta^2 + o(1))m_k^2.$$

and

$$\begin{split} & \int_{r_k}^{r_{k+1}} \frac{n-1}{r} |u'(r)|^2 \mathrm{d}r \\ & = \int_0^{T_0} \frac{n-1}{k(T_0 + o(1))} m_k^2 \beta^2 \sin^2(\beta x) dx \\ & = (\frac{n-1}{2} + o(1)) \beta^2 \frac{m_k^2}{k}. \end{split}$$

Since

$$F(0) - F(m_k) = \int_{r_k}^{\infty} \frac{n-1}{r} |u'(r)|^2 dr.$$

Then, for any fixed small constant ϵ and k large enough we have

(4.15)
$$m_k^2 \ge (n - 1 - 2\epsilon) \sum_{i=k}^{\infty} \frac{m_i^2}{i}.$$

This leads to (4.13) by the following calculus lemma. \square

LEMMA 4.3. Let $\{a_k\}$ be sequence of nonincreasing positive numbers and satisfies

$$(4.16) a_k \ge p \sum_{i=k}^{\infty} \frac{a_i}{i}$$

for some positive constant p. Then for some positive constant C there holds

$$(4.17) a_k \le Ck^{-p}.$$

Proof. Define

$$A_k = \sum_{i=k}^{\infty} \frac{a_i}{i}, \quad k \ge 1.$$

Then

$$k(A_k - A_{k+1}) \ge pA_k, \quad k \ge 1.$$

Hence

$$\frac{A_k}{A_{k+1}} \ge \frac{k}{k-p}, \quad k > p.$$

Let $N_0 > p$ fixed. For any integer $N > N_0$, we have

$$\frac{A_{N_0}}{A_N} \ge \prod_{k=N_0}^{k=N-1} \frac{k}{k-p} \ge \prod_{k=N_0}^{k=N-1} (1 + \frac{p}{k}).$$

Then

$$\ln(A_{N-0}) - \ln(A_N) \ge \sum_{k=N_0}^{N-1} \ln(1 + \frac{p}{k}) \ge p \ln(N) - c, \quad N > N_0$$

where c is a constant independent of N. Therefore we obtain

$$A_N < CN^{-p}, \quad \forall N > 0$$

where C is a constant independent of N. Since a_i is nonincreasing, we obtain

$$A_k - A_{2k+1} = \sum_{i=1}^{2k} \frac{a_i}{i} \ge \frac{1}{2} a_{2k}, \quad \forall k.$$

Hence the lemma is proven. \square

Now we apply Theorem 3.1 to obtain

THEOREM 4.4. Assume f(u) = F''(0)u - F'(u) satisfies (3.4). Then when $|u_0| < 1$, the solution u(r) to (4.1) has asymptotic behavior

(4.18)
$$u(r) = (\tilde{\gamma}_1 + o(1))u_1(r) + (\tilde{\gamma}_2 + o(1))u_2(r)$$

where u_1 , u_2 are defined in (2.10).

Proof. Just note that we choose R large enough. Define γ_1, γ_2 by the following relation

(4.19)
$$\begin{cases} \gamma_1 u_1(R) + \gamma_2 u_2(R) = u(R) \\ \gamma_1 u_1'(R) + \gamma_2 u_2'(R) = u'(R) \end{cases}$$

Following Lemma 4.2 with $\epsilon < \frac{n-1}{2} - \frac{1}{\alpha}$, we see that (γ_1, γ_2) satisfies (3.9) when R is large enough. Hence (3.2) with (γ_1, γ_2) has a solution $\bar{u}(r)$ with asymptotic behavior (4.18). On the other hand, by the uniqueness of initial value problem for (4.1), we have $u(r) = \bar{u}(r)$, the theorem is proved. \square

Remark 4.5. In the typical Allen-Cahn equation, we have $f(u) = -u^3$, which satisfies (3.4) when $n \ge 3$. However, when n = 2 the condition (3.4) is not satisfied. It would be interesting to show that (4.18) still holds in this case. Similarly, if we consider $f(u) = |u|^{p-1}u$ with p > 1, it would be interesting to show that (4.18) still holds for $n \ge 2$. Note that when n = 1, the solution may not vanish, hence (4.18) does not hold in general.

5. Application to the thin film equation. We return to the study of the thin film equations. Consider

(5.1)
$$\begin{cases} u'' + \frac{n-1}{r}u' = g(u) & r > 0 \\ u(0) = u_0 > 0, \quad u'(0) = 0. \end{cases}$$

where the nonlinear term q(u) satisfies

$$(5.2) q'(1) < 0, q(1) = 0, q(u) > 0 \text{ for } 0 < u < 1, q(u) < 0 \text{ for } u > 1.$$

Let $G(u) = \int_1^u g(s) \mathrm{d}s$, then G(u) is non increasing for u > 1 and non decreasing for u < 1. Note that $g(u) = u^{-p} - u^{-q}$ with $p > q \ge 0$ is a typical example.

It is easy to see that

LEMMA 5.1. u is bounded above and below away from 0.

Proof. Note that

(5.3)
$$\left(\frac{u'^2}{2} - G(u)\right)' = -\frac{n-1}{r}(u')^2 \le 0, \qquad r > 0.$$

Then

(5.4)
$$G(u(r)) > G(u_0) + \frac{u'(r)^2}{2} \ge G(u_0), \qquad r \ge 0.$$

By (5.2), we know $0 < \delta < u(r) < C < \infty$, $\forall r$ for some $\delta, C > 0$. \square We can also obtain the limit of u easily.

Lemma 5.2. There holds

$$\lim_{r \to \infty} u(r) = 1.$$

Proof. Consider v(r) = u(r) - 1, then it satisfies

(5.6)
$$v'' + \frac{n-1}{r}v' - g'(1)v + f(v) = 0 \qquad r > 0,$$

where

(5.7)
$$f(v) = -g(1+v) + g'(1)v, \qquad -g'(1) = \beta^2 > 0.$$

Then the arguments in Lemma 4.1 leads to (5.5). \square

THEOREM 5.3. Assume (5.2) and (3.4) where f(v) is given by (5.7). Then the solution u of (5.1) has the following asymptotic behavior

(5.8)
$$u(r) = 1 + (\tilde{\gamma}_1 + o(1))u_1(r) + (\tilde{\gamma}_2 + o(1))u_2(r), \quad \text{as } r \to \infty.$$

Proof. The proof is similar to proof of Theorem 4.4. The detail is omit. \square

REMARK 5.4. In the typical thin film equation, we have $g(u) = u^{-p} - 1$ and $f(v) = -(1+v)^{-p} + 1 - pv$, which satisfies (3.4) when n > 3. However, when n = 2, 3 the condition (3.4) is not satisfied. It would be interesting to investigate if (4.18) still hold in this case.

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Note added at galley proof. Similar results for the cases n=2, 3 as in Remark 4.5 and Remark 5.4 have been proven by Changfeng Gui, Xue Luo and Feng Zhou in a forthcoming paper.

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