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ASYMPTOTIC BEHAVIOR AND DEGENERACY OF BIHARMONIC FUNCTIONS ON RIEMANNIAN MANIFOLDS

BY LUNG OCK CHUNG

One of the most fascinating results in harmonic classification theory is the is the identity $O_{HD}^{N}=O_{HC}^{N}$, where H stands for the class of harmonic functions h, $\Delta h=0$, with $\Delta=d\delta+\delta d$ the Laplace-Beltrami operator, and HD, HC are the subclasses of functions which are Dirichlet finite, or bounded Dirichlet finite, respectively. For any class F of functions, O_{F} , \tilde{O}_{F} denote the classes of Riemannian manifolds on which $F \subset \mathbf{R}$ or $F \not\in \mathbf{R}$ respectively, and O_{F}^{N} , \tilde{O}_{F}^{N} are the corresponding subclasses of manifolds of dimension $N \ge 2$.

A striking phenomenon in biharmonic classification theory is that, in contrast with the harmonic case, the inclusion $O_{H^2D} \subset O_{H^2C}$ is strict, with H^2 the class of nonharmonic biharmonic functions. This has been, however, known only in the 2-dimensional case, in which it was established by undoubtedly the most intricate counterexample in all classification theory (Nakai-Sario [6]). The technique of complex analysis used therein is not available for an arbitrarily high dimension.

Combining certain recent results in the biharmonic classification of the Poincaré N-ball for the subclasses H^2D , H^2B of H^2 functions which are Dirichlet finite or bounded, respectively (Hada-Sario-Wang [2], [3]), one can draw the conclusion that $O_{H^2D}^N \subset O_{H^2C}^N$ is strict for $N \ge 5$. However, for N=3, 4, the reasoning fails and the question remains unsettled.

The first purpose of the present paper is to give a complete and unified solution to this problem by proving the strict inclusion

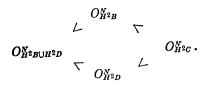
$$O_{H^{2}D}^{N} < O_{H^{2}C}^{N}$$

for any dimension $N \ge 2$. We shall, in fact, show more generally that $O_{H^2B}^N \subset O_{H^2D}^N$. On the other hand, from recent results on the Poincaré *N*-ball (Hada-Sario-Wang [2], [3]), we infer that $O_{H^2D}^N \subset O_{H^2B}^N$. In summary, we have the following string of strict inclusion relations:

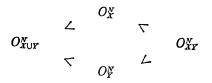
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Proceeding from the special to the general, we state our most general result which will be the content of our paper:



for and $N \ge 2$; $X = H^2B$, Γ , G, HP, HB, HD, HC; $Y = H^2D$, H^2L^p . Here $HF = H \cap F$, $H^2F = H^2 \cap F$; $1 \le p < \infty$; Γ is the class of biharmonic Green's functions (Sario [9]); G is the class of harmonic Green's functions; and P is the class of positive functions. Of these relations, the following cases, in addition to the aforementioned partial relations on H^2B and H^2D , have been previously known: $(X, Y) = (HP, H^2D)$, (Sario-Wang [11]); $(X, Y) = (HD, H^2D)$, (HB, H^2D) , (Sario-Wang [11]); $(X, Y) = (HD, H^2D)$, (HB, H^2D) , (Sario-Wang [13]); $(X, Y) = (G; H^2D)$, (Nakai-Sario [8], Sario-Wang [12]); $(X, Y) = (\Gamma, H^2D)$, (Wang [14]). The rest are new: in addition to the aforementioned unsettled relation between H^2B and H^2D , the cases (X, H^2L^p) , where X = G, HP, HB, HD, HC, Γ , and H^2B .

An essential aspect of our paper is that all the above inclusion relations, old and new, are obtained in a simple and unified manner. The N-cylinder, endowed with various simple metrics, is the only manifold we will need as a counterexample. This unification of approach is made possible by a systematic use of the asymptotic behavior of solutions of differential equations.

The proof of the above statement on the classes O_x and O_y will be presented in Lemmas 1-25 and §5.

1. Consider the N-cylinder.

$$M = \mathbf{R} \times S^{N-1} = \{ |x| < \infty, |y_i| \le \pi, i = 1, \dots, N-1 \}$$

with the faces $y_i = \pi$ and $y_i = -\pi$ identified, for each *i*, by a parellel translation perpendicular to the x-axis. Endow M with the metric

$$ds^{2} = \varphi^{2}(x)dx^{2} + \psi^{2}(x)dy_{1}^{2} + \sum_{i=2}^{N-1} dy_{i}^{2},$$

where φ , $\psi \in C^{\infty}(-\infty, \infty)$. The proof of our theorem will consist, in essence, of two parts. First we show that for a suitable choice of φ , ψ ,

$$M \in O_G^N \cap O_{HF}^N \cap O_F^N \cap O_{H^2B}^N \cap O_{H^2D}^N \cap O_{H^2L}^N p,$$

and then that for another choice of φ , ψ ,

$$M_1 \in \tilde{O}_G^N \cap \tilde{O}_{HF}^N \cap \tilde{O}_T^N \cap \tilde{O}_{H^2B}^N \cap O_{H^2D}^N \cap O_{H^2L}^N p$$
,

where M_1 is the manifold with the new metric, and F=P, B, D, C. This will establish our claims $O_X^v \cap \tilde{O}_Y^v \neq \phi$ and $\tilde{O}_X^v \cap O_Y^v \neq \phi$. The remaining relations $\tilde{O}_X^v \cap \tilde{O}_Y^v \neq \phi$ and $O_X^v \cap O_Y^v \neq \phi$ will then follow from other quite trivial choices of φ and φ .

2. To establish the first string of relations in §1, we choose $\varphi = \psi$ on $(-\infty, \infty)$, $\varphi(x) = |x|^{-3}$ for |x| > 1.

LEMMA 1. A harmonic function $h(x, y), y=(y_1, \dots, y_{N-1})$, has a representation $h(x, y)=f_0(x)+\sum_{n=1}^{\infty}f_n(x)G_n(y)$, where $G_n(y)=\prod_{i=1}^{N-1}G_n^i(y_i)$ with $G_n^i(y_i)=\pm \sin n_i y_i$ or $\pm \cos n_i y_i$ for some integer n_i . The series converges absolutely and uniformly on compact sets.

In fact, by a standard application of the Peter-Weyl theorem, we obtain for any x_0 , $h(x_0, y) = f_0(x_0) + \sum_{n=1}^{\infty} f_n(x_0)G_n(y)$. Here the G_n are invariant under varying x_0 by virtue of continuity. The convergence follows by a standard argument using differentiation with respect to y.

LEMMA 2. f(x) is harmonic if and only if f(x)=ax+b.

For the proof, solve the equation $\Delta f = -g^{-1/2}f'' = 0$, where $\sqrt{g} dxdy$ is the volume element.

LEMMA 3. $M \in O_X$ with $X = \Gamma$, G, HP, HB, HD, HC.

From the harmonic classification theory, we have the the inclusions $O_G < O_{HP} < O_{HB} < O_{HD} = O_{HC}$. Moreover, $O_G < O_T$ (Wang [15]). Thus it suffices to show that $M \in O_G$. The harmonic measure ω of $\{x=c>0\}$ on $\{0 < x < c\}$ is x/c in view of Lemma 2. As $c \to \infty$, $\omega \to 0$. Similarly, the harmonic measure of the boundary component at $x=-\infty$ vanishes. Therefore, $M \in O_G$.

3. Having discussed the spaces O_G , O_{HF} , O_{Γ} of the first string of relations in § 2, we turn to the spaces related to biharmonic functions. First we present some preparatory results.

LEMMA 4. If f(x)G(y) is harmonic, then f is strictly monotone.

Suppose the claim false. Then for $c_1 < c_2$, say, $f | \{c_1 < x < c_2\}$ is not strictly monotone, and f takes on its maximum or minimum on $\{c_1 < x < c_2\}$ at some point of this open interval. So does, a fortiori, fG, in violation of the maximum principle for harmonic functions.

LEMMA 5. If $f(x)G(y_1)$ is harmonic, $G(y_1) = \pm \sin n_1 y_1$ or $\pm \cos n_1 y_1$ with $n_1 \neq 0$, then $f(x) = ae^{-n_1 x} + be^{n_1 x}$.

We obtain successively

$$\begin{aligned} \mathcal{\Delta}(fG) = (\mathcal{\Delta}f)G + f\mathcal{\Delta}G = 0, \\ \mathcal{\Delta}f = -g^{-1/2}f'', \\ \mathcal{\Delta}G = g^{-1/2}(\varphi^{-2}g^{1/2}n_1^2G) = n_1^2g^{-1/2}G, \\ \mathcal{\Delta}(fG) = (-g^{-1/2}f'' + n_1^2g^{-1/2}f)G = 0, \end{aligned}$$

with the fundamental solutions $f_1(x) = e^{n_1 x}$ and $f_2(x) = e^{-n_1 x}$.

LEMMA 6. If $f(x)G(y_i)$ is harmonic with $G(y_i)$ not constant, $i \neq 1$, then

$$f(x) = ax(1+o(1))+b(1+o(1)), \quad a \neq 0$$

either as $x \rightarrow \infty$ or else as $x \rightarrow -\infty$.

This time we have

$$\Delta(fG) = \left(-\frac{1}{\sqrt{g}}f'' + n_i^2 f\right)G = 0,$$

hence

$$f''=n_i^2\sqrt{g}f.$$

We now make use of the following theorem of Haupt [4] and Hille [5]: A necessary and sufficient condition for the equation

$$f''(x) = p(x)f(x)$$

on $(0, \infty)$ to have the fundamental solutions

$$f_1(x) = x(1+o(1))$$

 $f_2(x) = 1+o(1)$

as $x \rightarrow \infty$ is that

 $xp(x) \in L^1(0,\infty)$.

Since $n_i^2 \sqrt{g} = n_i^2 |x|^{-3}$ for |x| > 1, the condition of the theorem of Haupt and Hille is satisfied, and we conclude that

$$f(x) = a_1 x(1+o(1)) + b_1(1+o(1)) \quad \text{as} \quad x \to \infty,$$

$$f(x) = a_2 x(1+o(1)) + b_2(1+o(1)) \quad \text{as} \quad x \to -\infty.$$

By Lemma 3, fG is not bounded and the same is true of f. Consequently $a_1 \neq 0$ or $a_2 \neq 0$.

LEMMA 7. If $f(x)G(y_2, y_3, \dots, y_{N-1})$ is harmonic, with G not constant, then

or

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$$f(x) = ax(1+o(1))+b(1+o(1))$$

either as $x \rightarrow \infty$ or else as $x \rightarrow -\infty$.

The proof is the same as for Lemma 6, the equation

$$f'' = \left(\sum_{i=2}^{N-1} n_i^2\right) \sqrt{g} f$$

again satisfying the Haupt-Hille condition.

LEMMA 8. If f(x)G(y) is harmonic with $G(y) = \prod_{i=1}^{N-1} G^i(y_i)$, $G^1(y_i)$ not constant, then

$$f(x) \sim a e^{n_1 |x|} \qquad with \quad a \neq 0$$

either as $x \rightarrow \infty$ or else as $x \rightarrow -\infty$.

The equation $\Delta(fG)=0$ gives

$$f'' = \sqrt{g} (n_1^2 \varphi^{-2} + \sum_{i=2}^{N-1} n_i^2) f,$$

which for |x| > 1 is reduced by the transformation $f(x) = f(n_1 x)$ to the form

$$f''(x) = (1+c|x|^{-3})f(x)$$
.

We now make use of the following theorem of Bellman [1]:

If $p(x)\to 0$ as $x\to\infty$ and if $\int_0^{\infty} p^2 dx < \infty$, then the equation f''=(1+p)f in $(0,\infty)$ has the fundamental solutions

$$f_{1}(x) \sim \exp\left[-\left(x + \frac{1}{2} \int_{x_{0}}^{x} p(x) dx + o(1)\right)\right],$$

$$f_{2}(x) \sim \exp\left[x + \frac{1}{2} \int_{x_{0}}^{x} p(x) dx + o(1)\right].$$

In the present case, $p(x)=c|x|^{-3}$ satisfies the condition of Bellman's theorem, and we obtain $f=a_1f_1+b_1f_2$ for x>1 and $f=a_2f_1+b_2f_2$ for x<-1 with

$$f_{1}(x) \sim \exp\left[-\left(|x| + \frac{1}{2} \int_{x_{0}}^{x} p dx + o(1)\right)\right]$$
$$f_{2}(x) \sim \exp\left[|x| + \frac{1}{2} \int_{x_{0}}^{x} p dx + o(1)\right].$$

If $b_1=b_2=0$, then $f(x)\to 0$ as $|x|\to\infty$, in violation of the maximum principle. Therefore either $b_1\neq 0$ or $b_2\neq 0$, and in view of the above transformation, we have the lemma.

LEMMA 9. A solution of $\Delta q = 1$ is $q_0(x) = \int_0^x \int_0^t \sqrt{g(s)} \, ds dt$. The general solution of $\Delta q = c$ is $cq_0(x) + h(x, y)$ where $h \in H$. Every q is unbounded.

The only part of the lemma that needs proving is the unboundedness of q. Suppose that there exists a bounded q. Then the transform $(Tq)(x) = \int_{y} q(x, y) dy = aq_0(x) + bx + c$ is bounded. Since $q_0 \to -\infty$ as $|x| \to \infty$, whereas bx changes its sign with x, we have a contradiction.

LEMMA 10. A solution of $\Delta^2 u(x) = 0$ is $u_0(x) = \int_0^x \int_{-\infty}^t s \sqrt{g(s)} \, ds dt$. It satisfies $u_0(x) \sim \pm a \log |x|$ for some constant a as $x \rightarrow \pm \infty$, respectively. The general solution $c_0 u_0(x) + c_1 q_0(x) + c_2 x + c_3$ is unbounded.

The proof is analogous to that of Lemma 9.

LEMMA 11. $M \in \tilde{O}_{H^2D}$.

The function $u_0(x)$ of Lemma 10 is Dirichlet finite:

$$D(u_0) = c \int_{-\infty}^{\infty} (u'_0)^2 \varphi^{-2} \varphi^2 dx$$

= $c_1 + c \Big(\int_{-\infty}^{-1} + \int_{1}^{\infty} \Big) |x|^{-2} dx < \infty$.

Lemma 12. $M \in \tilde{O}_{H^{2}L^{p}}$.

If fact,

$$\|u_0\|_p^p = c \int_{-\infty}^{\infty} |u_0|^p \sqrt{g} \, dx < \infty$$
 ,

since $|u_0(x)| \sim |a \log |x||$ but $\sqrt{g} \sim |x|^{-3}$ as $x \to \pm \infty$.

LEMMA 13. Let v(x) satisfy the equation $\Delta(v(x)G(y_1))=f(x)G(y_1)\in H$ with $fG\neq 0$ and $G(y_1)$ not constant. Then v is unbounded.

We have

$$\left(-\frac{1}{\sqrt{g}}v''+\frac{n_1^2}{\sqrt{g}}v\right)G=fG,$$

hence

 $v''=n_1^2v-\sqrt{g}f.$

By Lemma 5, $f(x) = ae^{n_1x} + be^{-n_1x}$ with $|a| + |b| \neq 0$. We may assume a < 0; the proof for the other case is analogous.

Suppose v is bounded. For sufficiently large x>0, $n_1^2v - \sqrt{g}f$ grows at the rate of $x^{-3}e^{n_1x}$. We thus have

$$v'(x) = v'(x_0) + \int_{x_0}^x (n_1^2 v - a x^{-3} e^{n_1 x}) dx$$

where we may choose $x_0 > 1$. It follows that

$$v(x) = v(x_0) + v'(x_0)(x - x_0) + \int_{x_0}^x \int_{x_0}^t (n_1^2 v(s) - as^{-3} e^{n_1 s}) ds dt$$

which is clearly not bounded.

LEMMA 14. Let v(x) satisfy the exuation $\Delta(v(x)G(y_i))=f(x)G(y_i)\in H$ with $fG\neq 0$, i>1, and $G(y_i)$ not constant. Then v is unbounded.

The proof is analogous to that of Lemma 13, with

$$v''=n_i^2\sqrt{g}\,v-\sqrt{g}\,f.$$

In applying Lemma 6, we may assume that f(x)=ax(1+o(1))+b(1+o(1)) as $x\to -\infty$ with a<0. If v is bounded, we have for x<-1,

$$v'(x) = \int_{-\infty}^{x} n_{1}^{2} |t|^{-3} v(t) dt - \int_{-\infty}^{x} |t|^{-3} f dt ,$$

$$v(x_{1}) - v(x_{2}) = \int_{x_{2}}^{x} \int_{-\infty}^{x} n_{1}^{2} |t|^{-3} v(t) dt dx - \int_{x_{2}}^{x_{1}} \int_{-\infty}^{x} |t|^{-3} f dt dx$$

for $x_2 < x < x_1 < -1$. As $x_2 \rightarrow -\infty$, the first integral converges but the second does not. Thus $v(x_2)$ cannot be bounded as $x_2 \rightarrow -\infty$, in violation of the assumption.

LEMMA 15. Let v(x) satisfy $\Delta(v(x)G(y)) = f(x)G(y) \in H$, with f(x)G(y) not constant. Then v(x) is not bounded.

We may assume $n_1 \neq 0$ and at least one $n_i \neq 0$, i > 1. We now have

$$v'' = (n_1^2 + \sum_{i=2}^{N-1} n_i^2 \sqrt{g})v - \sqrt{g}f.$$

Since fG is harmonic, $f \sim ae^{n_1|x|}$ for either $x \to \infty$ or else $x \to -\infty$. We may assume the former. Clearly $|\sqrt{g}f| \to \infty$ as $x \to \infty$. If v is bounded, then $\sqrt{g}f$ will dominate the right-hand side of the equation. On integrating as in the proof of Lemmas 13 and 14, we arrive at the contradiction that v is both bounded and unbounded.

LEMMA 16. $M \in O_{H^{2}B}$.

Suppose there exists a $u(x, y) \in H^2B$. Write $u(x, y) = v_0(x) + \sum_{n=1}^{\infty} v_n(x)G_n(y)$ with $G_n \neq G_m$ for $n \neq m$. Either $v_0(x)$ or some v_nG_n is not harmonic. Suppose this is true of $v_{n_0}G_{n_0}$. Then the thransform

$$(Tu)(x) = \int_{y} u G_{n_0} dy = c v_{n_0}(x)$$

is bounded, in violation of Lemma 15.

With Lemma 16, the proof of the first string of inclusion relations in §1 is complete.

4. We turn to the second string of relations in §1. We now choose $\varphi \equiv 1$ and ψ a positive symmetric C^{∞} function with $\psi(x) = \exp e^{|x|}$ for |x| > 1, and denote

the resulting manifold by M_1 .

The same proof as for Lemma 1 shows that every harmonic function h on M_1 has a representation

$$h(x, y) = f_0(x) + \sum_{n=1}^{\infty} f_n(x) G_n(y)$$
.

LEMMA 17. f(x) is harmonic if and only if $f(x) = a \int_0^x \phi^{-1} dx + b$.

This is seen by solving the harmonic equation $\Delta f(x) = -\phi^{-1}(\phi f')' = 0$.

LEMMA 18. $M_1 \in \tilde{O}_G \cap \tilde{O}_{HX}$, where X=P, B, D, C.

The function $f(x) = \int_0^x \phi^{-1} dx$ is bounded and its Dirichlet integral is

$$D(f) = \int_{M_1} (f')^2 \psi dx dy = c \int_{-\infty}^{\infty} \psi^{-1} dx < \infty .$$

LEMMA 19. The function

$$q_0(x) = -\int_0^x \psi^{-1}(t) \int_0^t \psi(s) ds dt$$

is quasiharmonic, that is, $\Delta q_0=1$. Every quasiharmonic function has the form $q=q_0+h$ with $h\in H$.

This is verified by direct computation of Δq_0 .

LEMMA 20. $M_1 \in \tilde{O}_{H^2B}$.

In fact, $q_0 \in H^2B$, since

$$\left|\psi^{-1}(t)\int_{0}^{t}\psi(s)ds\right|\sim e^{-|t|}$$
 as $|t|\rightarrow\infty$.

For verification, first apply l'Hospital's rule to the left-hand side to see that it goes to 0 as $|t| \to \infty$, and then show, again by l'Hospital's rule, that $\left| e^{it} \psi^{-1}(t) \int_{0}^{t} \psi(s) ds \right| \to 1$ as $|t| \to \infty$.

LEMMA 21. For $h(x) \in H$, the function

$$u_0(x) = \int_0^x \psi^{-1}(t) \int_0^t \psi(s) h(s) ds dt$$

is biharmonic. Every biharmonic function of the form u(x) can be written $u(x) = u_0(x) + c$.

The proof is again by direct computation.

LEMMA 22. Every nonharmonic biharmonic function of the form u(x) has infinite Dirichlet and L^p norms.

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An estimate similar to that in the proof of Lemma 20 shows that $|u'(t)| \sim e^{-|t|}$ either as $t \to \infty$ or else as $t \to -\infty$. The Dirichlet integral is

$$D(u)=c\int_{-\infty}^{\infty}(u')^2\psi dx=\infty.$$

Since u'(t) does not decrease faster than $e^{-|t|}$ at least in one direction, the same is true of u(t). Therefore,

$$\|u\|_p^p = c \int_{-\infty}^{\infty} |u|^p \psi dx = \infty.$$

LEMMA 23. If v(x)G(y) is a nonharmonic biharmonic function, with G(y) not constant, then $v \in L^p$.

Suppose $v \in L^p$ for some $1 \leq p < \infty$. Then $|v(x)|^p \phi(x)$ is integrable and decreases to 0. Let $\mathcal{A}(vG) = fG$. Since vG is nonharmonic biharmonic, f does not vanish in the neighborhood of at least one component of the ideal boundary, say $x = \infty$. As in Lemma 15, $\mathcal{A}(vG) = fG$ gives

$$(\psi v')' = (n_1^2 \psi^{-1} + \sum_{i=2}^{N-1} n_i^2 \psi) v - \psi f.$$

For large x>0, we may assume $f(x)<\varepsilon<0$, by changing the sign of G if necessary. Since $|v|\rightarrow 0$ rapidly, the dominating term on the right-hand side is $-\psi f$, and we obtain

$$(\psi v')' \ge c \psi > 0$$

for all sufficiently large x>0. On integrating from a sufficiently large x_0 to a larger x, we obtain

$$\psi v' \ge c \int_{x_0}^x \psi dx \, .$$

An estimation exactly as that in the proof of Lemma 20 yields

$$v' \ge ce^{-x}$$
.

Thus v can not be decreasing faster than ce^{-x} . This contradicts $|v|^p \phi(x) \rightarrow 0$ and completes the proof of the lemma.

LEMMA 24. The Dirichlet and L^p norms of the function vG of Lemma 23 are infinite.

By Lemma 23, $v \oplus L^p$. Therefore

$$\|vG\|_p^p = c \int_{-\infty}^{\infty} |v|^p \psi dx = \infty.$$

By the proof of Lemma 23, $|v'| \ge ce^{-|x|}$ either as $x \to \infty$ or else as $x \to -\infty$. Therefore,

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$$D(vG) = \int_{\mathcal{M}_1} (v'G)^2 \psi dx dy + \int_{\mathcal{M}_1} \sum_{i=1}^{N-1} \left(v \frac{\partial G}{\partial y^i} \right)^2 g^{ii} dx dy$$
$$\geq c \int_{-\infty}^{\infty} (v')^2 \psi dx = \infty.$$

Lemma 25. $M_1 \in O_{H^2L^p} \cap O_{H^2D}$.

Let u(x, y) be a nonharmonic biharmonic function. Write $u(x, y) = v_0(x) + \sum_{n=1}^{\infty} v_n(x)G_n(y)$. By Lemmas 22 and 24, neither v_0 nor any v_nG_n belongs to $D \cup L^p$ if it is nonharmonic. By the Dirichlet orthogonality of v_0 and the v_nG_n , we conclude that $v_0 + \sum_{n=1}^{\infty} v_nG_n$ is Dirichlet infinite.

Suppose $v_0 + \sum_{n=1}^{\infty} v_n G_n \in L^p$. Choose a nonharmonic term $v_{n_0} G_{n_0}$. Since $v_{n_0} G_{n_0} \notin L^p$, there exists an L^q function fG_{n_0} such that $(v_{n_0} G_{n_0}, fG_{n_0}) = \int_{M_1} v_{n_0} G_{n_0} fG_{n_0} dV$ = ∞ . On the other hand, $(v_{n_0} G_{n_0}, fG_{n_0}) = (v_0 + \sum_{n=1}^{\infty} v_n G_n, fG_{n_0}) < \infty$, a contradiction.

With Lemma 25, the proof of the second string of relations in §1 is complete.

5. It remains to show that $O_X \cap O_Y \neq \phi$ and $\tilde{O}_X \cap \tilde{O}_Y \neq \phi$. The metrics we shall choose will result in simple computations which also are completely analogous to those in §§ 2-4, and we can be brief.

To show that $\tilde{O}_x \cap \tilde{O}_r \neq \phi$, we choose $\psi = 1$ and $\varphi(x) = |x|^{-4}$ for |x| > 1. Then the solutions $\mathcal{\Delta}(f(x)) = 0$ and $\mathcal{\Delta}(q(x)) = 1$ turn out to belong to the desired function classes X, Y.

To prove $O_X \cap O_Y \neq \phi$, let $\varphi = \psi = 1$. It is easy to explicitly solve the equation $\mathcal{J}^2 u = 0$ in all cases and to show that the solutions do not belong to X or Y.

6. We have completed, by Lemmas 1-25, and \S 5, the proof of the following result:

THEOREM. The classification scheme

holds for X=G, HP, HB, HD, HC, Γ , H^2B ; $Y=H^2D$, H^2L^p .

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UNIVERSITY OF CALIFORNIA, LOS ANGELES NORTH CAROLINA STATE UNIVERSITY, RALEIGH