# Seidel's long exact sequence on Calabi-Yau manifolds 

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#### Abstract

In this paper, we generalize construction of Seidel's long exact sequence of Lagrangian Floer cohomology to that of compact Lagrangian submanifolds with vanishing Malsov class on general Calabi-Yau manifolds. We use the framework of anchored Lagrangian submanifolds and some compactness theorem of smooth $J$-holomorphic sections of Lefschetz Hamiltonian fibration for a generic choice of $J$. The proof of the latter compactness theorem involves a study of proper pseudoholomorphic curves in the setting of noncompact symplectic manifolds with cylindrical ends.


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## 1. Introduction

To put the content of this paper in perspective, we first recall a long exact sequence for symplectic Floer cohomology of Lagrangian submanifolds, which was constructed by Seidel [Se3] originally for the category of exact Lagrangian submanifolds on (noncompact) exact symplectic manifolds.

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### 1.1. Dehn twists and Seidel's long exact sequence

Let ( $M, \omega, \alpha$ ) be an exact symplectic manifold with contact-type boundary: $\alpha$ is a contact 1-form on $\partial M$ which satisfies $d \alpha=\left.\omega\right|_{\partial M}$ and makes $\partial M$ convex. Assume that $[\omega, \alpha] \in H^{2}(M, \partial M ; \mathbb{R})$ is zero so that $\alpha$ can be extended to a 1 -form $\theta$ on $M$ satisfying $d \theta=\omega$.

THEOREM 1.1 (SEIDEL [Se3])
Let $L$ be an exact Lagrangian sphere in $M$ together with a preferred diffeomorphism $f: S^{2} \rightarrow L$. Denote by $\tau_{L}=\tau_{(L,[f])}$ the Dehn twist associated to ( $L,[f]$ ). For any two compact exact Lagrangian submanifolds $L_{0}, L_{1} \subset M$, there is a long exact sequence of Floer cohomology groups

$$
\begin{equation*}
\longrightarrow H F\left(\tau_{L}\left(L_{0}\right), L_{1}\right) \longrightarrow H F\left(L_{0}, L_{1}\right) \longrightarrow H F\left(L, L_{1}\right) \otimes H F\left(L_{0}, L\right) \longrightarrow . \tag{1.1}
\end{equation*}
$$

Due to the well-known difficulties in the construction of Lagrangian intersection Floer cohomology and a new nontrivial compactness issue arising from the singularities of the Lefschetz fibration used for the construction, Seidel used the exactness assumption to avoid these difficulties and work entirely with the exact Lagrangian category, and he left the extension to the more general situation, such as that for closed Calabi-Yau manifolds, as an open problem (see [Se3]). While the limitation to the exact Lagrangian category simplifies the analysis of holomorphic disc bubbles, it also forces him to work entirely with the language of exact Lefschetz fibrations and to make sure that he does not go out of this domain largely for the consistency of his exposition, as Seidel himself indicated. Because of this, [Se3] develops a fair amount of geometry of exact Lefschetz fibrations, some of which are not directly relevant to the construction of the long exact sequence. Partly due to this digression, it took some effort and time for the author of this paper to get to the main point of Seidel's construction in [Se3].

The cases of closed Calabi-Yau manifolds or Fano toric manifolds are the ones that are physically most relevant to mirror symmetry. According to Kontsevich [K2] and Seidel [Se3], the symplectic Dehn twists correspond to a particular class of autoequivalences, "twist functors along spherical objects," of derived categories of coherent sheaves, and this long exact sequence corresponds to an exact sequence of the same form in the mirror Calabi-Yau. Therefore it is important to establish the long exact sequence for a class of Lagrangian branes that is closed under the action of symplectic Dehn twists. The class of exact Lagrangian submanifolds in exact symplectic manifolds is one such class, which Seidel considered in [Se3].

One of the points Seidel tried to ensure by working with the exact Lagrangian category is to have not only single-valuedness of the action functional on the path space but coherence of the definition of the action functional between different exact Lagrangian submanifolds: this then allows one to have the energy estimate for the Floer trajectories and, more importantly, to have tiny-big decomposition of the Floer moduli spaces entering in the construction. This decomposition then enables him to apply the spectral sequence argument and derive the desired
conclusion based on the contribution coming from the tiny part of the Floer moduli spaces which can be explicitly analyzed.

### 1.2. Calabi-Yau Lagrangian branes

In regard to extending Seidel's construction to closed Calabi-Yau manifolds, we highlight two points that we take in this paper.

The first point is our restriction to the class of Lagrangian submanifolds with zero Maslov class. This class is closed under the action by symplectic Dehn twists and enables one to consider the involved cohomology as a $\mathbb{Z}$-graded group which is essential in the point of view of mirror symmetry. It appears to the author that for this kind of long exact sequence to exist some condition of this sort of Calabi-Yau property is needed (see the remark at the end of Section 7.4 for the reason).

The second point is the usage of the notion of anchors and anchored Lagrangian submanifolds introduced in $[\mathrm{FO}+2]$. We recall the definition of anchored Lagrangian submanifolds introduced in [FO+2].

## DEFINITION 1.1

Fix a base point $y$ of an ambient symplectic manifold $(M, \omega)$. Let $L$ be a Lagrangian submanifold of $(M, \omega)$. We define an anchor of $L$ to $y$ as a path $\gamma:[0,1] \rightarrow M$ such that

$$
\gamma(0)=y, \quad \gamma(1) \in L .
$$

We call a pair $(L, \gamma)$ an anchored Lagrangian submanifold.
This notion has its origin in the preprint $[\mathrm{FO}+3]$ when the authors take the based point of view of Lagrangian submanifolds in relation to the coherence of the definitions of various Maslov-type indices and of action functionals when one considers several Lagrangian submanifolds altogether as one studies the Fukaya category. From a technical point of view, consideration of anchored Lagrangian submanifolds enables one to maintain the consistency of the definitions both of action functionals and of the absolute gradings on the Calabi-Yau Lagrangian branes. Most importantly, this also enables us to provide a coherent filtration in the relevant Floer complexes and to have tiny-big decomposition of the relevant Floer moduli spaces of the kind Seidel considered in [Se3].

Now we introduce a class of decorated Lagrangian submanifolds on CalabiYau manifold $(M, \omega)$ which we call Calabi-Yau Lagrangian branes. It is expected that this class of Lagrangian submanifolds generates the Fukaya category of a Calabi-Yau manifold that is mirror to the derived category of coherent sheaves on the mirror Calabi-Yau. We refer readers to the main part of the paper for various undefined terms in the statement. We also omit the important datum of flat line bundles on $L$ in this definition because it does not play much of a role in our proof but can be easily incorporated in the construction.

## DEFINITION 1.2

Let $y \in M$ be a base point, and let $\Lambda_{y} \subset T_{y} M$ be a fixed Lagrangian subspace. Suppose that $\Theta$ is a quadratic complex volume form on $(M, \omega, J)$ with $\left\langle\Theta(y), \Lambda_{y}\right\rangle=1$. We consider the quadruple $((L, \gamma), s,[b])$, which we call an (anchored) Calabi-Yau Lagrangian brane, that satisfies the following data:
(1) $L$ is a Lagrangian submanifold of $M$ such that the Maslov index of $L$ is zero and $[\omega] \in H^{2}(M, L ; \mathbb{Z})$; we also enhance $L$ with a flat complex line bundle on it;
(2) $\gamma$ is an anchor of $L$ relative to $y$;
(3) $s$ is a spin structure of $L$;
(4) $[b] \in \mathcal{M}(L)$ is a bounding cochain described in Section 8.2.

We denote by $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$ the collection of Calabi-Yau Lagrangian branes and define $\operatorname{Fuk}\left(\mathcal{E}_{\text {brane }}^{\mathrm{CY}}\right)$ to be the Fukaya category generated by $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$.

We remark that the notion of an anchor to $L$ is introduced to solve the problems of grading and filtration on the Floer complex in a uniform way in $[\mathrm{FO}+2]$. In particular, it provides a canonical filtration on the associated Floer complex of anchored Lagrangian submanifolds which is needed to apply some spectral sequence argument in the proof (see the end of Section 10 in particular).

### 1.3. Statement of the main result and compactness issue

The main purpose of this paper is to construct an exact sequence for the CalabiYau Lagrangian branes on Calabi-Yau manifolds which is the analogue to that of Seidel [Se3].

We first note that each Dehn twist $\tau_{L}$ along a given Lagrangian sphere $L \subset M$ acts on $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$. We denote this action by

$$
\left(\tau_{L}\right)_{*}: \mathcal{E}_{\text {brane }}^{\mathrm{CY}} \rightarrow \mathcal{E}_{\text {brane }}^{\mathrm{CY}}
$$

and the image of $\mathcal{L}$ under this action by $\tau_{L}(\mathcal{L})=\left(\tau_{L}\right)_{*} \mathcal{L}$. This action defines an autoequivalence on $\operatorname{Fuk}\left(\mathcal{E}_{\text {brane }}^{\mathrm{CY}}\right)$, whose nonanchored versions should correspond to twist functors along spherical objects of derived categories of coherent sheaves alluded to in the beginning of this introduction.

## THEOREM 1.2

Let $(M, \omega)$ be a compact (symplectic) Calabi-Yau manifold, and let $y \in M$ be a base point. Let $L \subset M$ be a Lagrangian sphere together with a preferred diffeomorphism $f: S^{n} \rightarrow L$, and let $\mathcal{L}=\left((L, \gamma), s_{s t}, 0\right)$ be the associated Calabi-Yau Lagrangian brane. Denote by $\tau_{L}=\tau_{(L,[f])}$ the Dehn twist associated to $(L,[f])$.

Consider any Calabi-Yau Lagrangian branes $\mathcal{L}_{0}, \mathcal{L}_{1}$. Then there is a long exact sequence of $\mathbb{Z}$-graded Floer cohomologies

$$
\begin{equation*}
\longrightarrow H F\left(\left(\tau_{L}\right)_{*} \mathcal{L}_{0}, \mathcal{L}_{1}\right) \longrightarrow H F\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \longrightarrow H F\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes H F\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \longrightarrow \tag{1.2}
\end{equation*}
$$

as a $\Lambda_{\text {nov }}$-module, where the Floer cohomologies involved are the deformed Floer cohomologies constructed in $[F O+1]$.

We also have the long exact sequence for the nonanchored version of the Floer cohomology (see Section 11.2).

Once we use these frameworks, construction of the long exact sequence largely follows Seidel's strategy: we use the framework of Lefschetz fibration with Lagrangian boundary conditions for the construction of various operators appearing in the Floer theory, and we use the spectral sequence for the $\mathbb{R}$-filtered groups based on the tiny-big decomposition of the Floer moduli spaces. However, unlike the exact Lagrangian case, the definition of Lagrangian Floer cohomology for Calabi-Yau Lagrangian branes meets obstruction, as described in $[F O+1]$. Because of this we have to use the Maurer-Cartan elements $b_{i}$ and use the associated deformed Floer cohomology appearing in the statement of the main theorem (Theorem 1.2). (Since Lagrangian submanifolds with zero Maslov class in Calabi-Yau manifolds are semipositive, the related transversality issue is relatively standard, which is one of the advantages of considering this class of Lagrangian submanifolds.) For the reader's convenience and the readability of the paper, we borrow a fair amount of material from $[\mathrm{FO}+1]$ in our exposition. For the same reason, we also borrow much exposition from [Se3] and refer to the two for further details. In a way, most of the materials used in this paper are not new but have already been present in the literature in one way or another. We organize them in a coherent way to be able to construct the required long exact sequence. Familiarity with the scheme in [Se3] would be useful for the readers to follow the stream of the arguments used in this paper, especially those presented in Sections 9-11.

However, there is one nontrivial analytical issue that needs to be overcome. This concerns the issue of compactification of the smooth pseudoholomorphic section of Lefschetz (Hamiltonian) fibration when the fibration has nonempty critical fibers. By the definition of Lefschetz Hamiltonian fibration given in Definition 5.1, any smooth section avoids critical points of the fibration. However, a priori a sequence of smooth sections may approach critical points if the derivatives of the sections in the sequence blow up. When applied to a sequence of pseudoholomorphic sections, the bubble could touch the critical points. Therefore to define the relative Gromov-Witten-type invariants in the Lefschetz fibration setting, one should study the behavior of pseudoholomorphic sections approaching the critical points. This compactness result may be mathematically the most novel part of this paper and is performed in Section 7.

In this regard, we prove the following.

THEOREM 1.3 (THEOREM 7.6)
Let $\pi: E \rightarrow \Sigma$ be a Lefschetz Hamiltonian fibration with Lagrangian boundary $Q \subset$ $E_{\partial \Sigma}$ such that $E$ is fiberwise Calabi-Yau and $Q$ has vanishing fiberwise Maslov class. Then there exists a dense subset of $j$-compatible $J$ 's for which we have a constant $C>0$ depending only on $(E, J, j)$ and the section class $B \in \pi_{2}^{\sec }(E, Q)$, but independent of $s$, such that we have

$$
\begin{equation*}
\operatorname{dist}\left(\operatorname{Im} s, E^{\mathrm{crit}}\right) \geq C \tag{1.3}
\end{equation*}
$$

for any smooth section $s: \Sigma \rightarrow E$ with $[s, \partial s] \in B \in \pi_{2}^{\sec }(E, Q)$.
The proof of this theorem turns out to involve the compactification and the Fredholm theory in the setting of symplectic field theory (see [EGH]) in which we regard a bubble touching a critical point $x_{0}$ as a proper pseudoholomorphic curve on $\mathbb{C}$ in a punctured fiber $E_{z_{0}} \backslash\left\{x_{0}\right\}$ (see [FO+4], [OZ] for relevant studies of such compactification and Fredholm theory). This theorem may not hold without assuming that $E$ is fiberwise Calabi-Yau and that $Q$ has vanishing relative Maslov class (see the end of Section 7.4 for the reason).

Once this theorem is established, study of compactification of smooth pseudoholomorphic sections in the current case is essentially the same as the case of smooth Hamiltonian fibrations as studied in [En] and [MS].

The result in this paper was first announced at Eliashberg's 60th birthday conference "New Challenges and Perspectives in Symplectic Field Theory" held at Stanford University June 25-29, 2007 and then presented in various seminars and in conferences afterward. We apologize to readers for the long delay in appearance.

## 2. Basic facts on symplectic Dehn twists

In this section, we summarize basic facts on the Dehn twists in the symplectic point of view which Seidel [Se1], [Se3] extensively studied in a series of papers. We borrow the basic facts on the symplectic Dehn twists from them with a slight variation of the exposition that is necessary for the purpose of this paper.

Assume that $L \subset(M, \omega)$ is an embedded Lagrangian sphere together with an equivalence class $[f]$ of diffeomorphisms $f: S^{n} \rightarrow L$ : two $f_{1}, f_{2}$ are equivalent if and only if $f_{2}^{-1} f_{1}$ can be deformed inside $\operatorname{Diff}\left(S^{n}\right)$ to an element of $O(n+1)$. To any such $(L,[f])$ Seidel associates a Dehn twist $\tau_{L}=\tau_{(L,[f])} \in \operatorname{Symp}(M)$ using a model Dehn twist on the cotangent bundle $T:=T^{*} S^{n}$. Let $f: S^{n} \rightarrow L \subset M$ be a representative of the equivalence class [ $f$ ]. Denote by $T(r) \subset T^{*} S^{n}$ the disc bundle of radius $r$ in terms of the standard metric on the unit sphere $S^{n}=S^{n}(1) \subset \mathbb{R}^{n+1}$. Identifying $T=T S^{n}$ with respect to the standard metric, one considers the map

$$
\begin{equation*}
\sigma_{t}(u, v)=\left(\cos (t) u-\sin (t)\|u\| v, \cos (t) v+\sin (t) \frac{u}{\|u\|}\right) \subset \mathbb{R}^{n+1} \times S^{n} \tag{2.1}
\end{equation*}
$$

for $0<t<\pi ; \sigma_{\pi}$ is the antipodal involution $A(u, v)=(-u,-v)$. Next we fix a function $R \in C^{\infty}(\mathbb{R}, \mathbb{R})$ such that

$$
\begin{gather*}
\operatorname{supp} R \subset T(1) \\
R(-t)=R(t)-t \quad \text { for }|t| \leq \frac{1}{2} . \tag{2.2}
\end{gather*}
$$

Then we consider the rescaled function

$$
\begin{equation*}
R_{\lambda}(t)=\lambda R\left(\frac{t}{\lambda}\right) \tag{2.3}
\end{equation*}
$$

for all $0<\lambda \leq 1$. Then $R_{\lambda}$ is supported in $T(\lambda)$ and satisfies

$$
\begin{equation*}
R_{\lambda}(-t)=R_{\lambda}(t)-t \quad \text { for }|t| \leq \frac{\lambda}{2} . \tag{2.4}
\end{equation*}
$$

Insertion of 1-parameter $\lambda$ in our choice of $R$ is deliberate and is later explicitly related to the parameter that enters in the Lagrangian surgery.

The following lemma is a slight variation of [Se3, Lemma 1.8], whose proof is referred thereto.

## LEMMA 2.1

Let $\mu: T \backslash T(0) \rightarrow \mathbb{R}$ be the length function $\mu(u, v)=\|v\|$, and let $H_{\lambda}=R_{\lambda} \circ \mu$ on $T \backslash T(0)$. Then $\phi_{H_{\lambda}}^{2 \pi}$ extends smoothly over $T(0)$ to a symplectic diffeomorphism $\phi_{\lambda}$ of $T$. The function $K_{\lambda}=2 \pi\left(R_{\lambda}^{\prime} \circ \mu-R \circ \mu\right)$ also extends smoothly over $T(0)$ and satisfies

$$
\phi_{\lambda}^{*} \theta_{T}-\theta_{T}=d K_{\lambda} .
$$

These $\phi_{\lambda}$ are called model Dehn twists.
The model Dehn twists, denoted by $\tau_{\lambda}$, have the explicit formula

$$
\tau_{\lambda}(y)= \begin{cases}\sigma_{2 \pi R_{\lambda}^{\prime}(\mu(y))}(y), & y \in T(\lambda) \backslash T(0)  \tag{2.5}\\ A(y), & y \in T(0)\end{cases}
$$

where the angle of rotation goes from $2 \pi R_{\lambda}^{\prime}(0)=\pi$ to $2 \pi R_{\lambda}^{\prime}(\lambda)=0$. Note that as $\lambda \rightarrow 0$, we have

$$
2 \pi R_{\lambda}^{\prime}(\mu(y))=2 \pi R^{\prime}\left(\frac{\mu(y)}{\lambda}\right)
$$

and it changes from $2 \pi R^{\prime}(0)=\pi$ to $2 \pi R^{\prime}(1)=0$.
Now we take a Darboux-Weinstein chart, or a symplectic embedding $\iota$ : $T(\lambda) \rightarrow M$ such that

$$
\left.\iota\right|_{o_{T^{*} S^{n}}}=f, \quad \iota^{*} \omega=\omega_{T}\left(=-d \theta_{T}\right),
$$

for a representative of the framed Lagrangian sphere ( $L,[f]$ ).
Take a model Dehn twist $\tau$ supported in the interior of $T(1) \subset T^{*} S^{n}$.
We denote $U=\operatorname{im} \iota$ and fix the Darboux neighborhood once and for all, and we consider the 1-parameter family of Dehn twists $\tau_{r}$, any of which we denote by $\tau_{L}$ :

$$
\tau_{L}=\tau_{(L,[f] ; r)}= \begin{cases}\iota \circ \tau_{r} \circ \iota^{-1} & \text { on } \operatorname{im}(\iota)=U  \tag{2.6}\\ \mathrm{id} & \text { elsewhere }\end{cases}
$$

We quote the following basic fact on the Dehn twist $\tau_{(L,[f])}$ from [Se3] with a slight variation of the statements.

PROPOSITION 2.2 ([Se3, PROPOSITION 1.11])
Let $(L,[f])$ be a framed Lagrangian sphere in $M$. There is a 1-parameter family of

Lefschetz fibrations $\left(E_{\lambda}^{L}, \pi_{\lambda}^{L}\right) \rightarrow D(\lambda)$ together with an isomorphism $\phi_{\lambda}^{L}: E_{\lambda}^{L} \rightarrow M$ of symplectic manifolds, such that we have the following.
(1) Consider the rescaling $\operatorname{map} R_{\lambda}: D(\lambda) \rightarrow D(1)$ defined by $z \mapsto z / \lambda$. Then

$$
\left(R_{\lambda}\right)^{*} E_{1}^{L}=E_{\lambda}^{L}
$$

(2) If $\rho_{\lambda}^{L}$ is the symplectic monodromy around $\partial \bar{D}(\lambda)$, then $\phi_{\lambda}^{L} \circ \rho^{L} \circ\left(\phi_{\lambda}^{L}\right)^{-1}$ is a Dehn twist along $(L,[f])$.
(3) There exists a decomposition

$$
E^{L}=E \cup \bar{D}(\lambda) \times(M \backslash \iota(T(\lambda)) \backslash V)
$$

such that $E$ is the standard fibration $q: \mathbb{C}^{n+1} \rightarrow \mathbb{C}$ defined by

$$
q\left(z_{1}, \ldots, z_{n+1}\right)=z_{1}^{2}+\cdots+z_{n+1}^{2}
$$

We denote any of these maps by $\tau_{L}$.

An important point on this Dehn twist $\tau_{(L,[f] ; \lambda)}$ is that its support can be put into a Darboux neighborhood of the given Lagrangian sphere $L$ which can be made as close as $L$ by choosing $\lambda>0$ small, whose derivative can be controlled. One can choose $R$ such that for some $\delta>0$ we have

$$
\begin{align*}
& R^{\prime}(t) \geq 0 \quad \text { for all } t \geq 0  \tag{2.7}\\
& R^{\prime \prime}(t)<0 \quad \text { for all } t \geq 0 \text { such that } R^{\prime}(t) \geq \delta \tag{2.8}
\end{align*}
$$

and then consider $R_{\lambda}$ for any sufficiently small $\lambda>0$. According to Seidel's terminology, the corresponding Dehn twist is $\delta$-wobbly.

## 3. Action, grading, and anchored Lagrangian submanifolds

In this section, we consider the general Lagrangian submanifolds treated as in $[\mathrm{FO}+1]$. For the fine chain level analysis of the Floer complex, it is essential to analyze the $\mathbb{R}$-filtration on $C F\left(L_{0}, L_{1}\right)$ that is provided by the action functional $\mathcal{A}$ on $\Omega\left(L_{0}, L_{1}\right)$. This action functional is not single valued on $\Omega\left(L_{0}, L_{1}\right)$ itself, even for the pair $\left(L_{0}, L_{1}\right)$ of Calabi-Yau Lagrangian branes, but single valued only on some covering space. For the purpose of studying the Fukaya category and performing various constructions in the Floer homology in a coherent manner, we need to consider a whole collection of Lagrangian submanifolds and assign these auxiliary data to each pair of the given collection in a consistent way. For this purpose, $[\mathrm{FO}+1]$ and $[\mathrm{FO}+2]$ use the notion of anchored Lagrangian submanifolds. This auxiliary data is important later for consistency of definition of action functionals and in turn for the analysis of tiny-big decomposition of the various Floer moduli spaces entering in the construction of boundary map, chain map, chain homotopy, and pants products.

### 3.1. Novikov covering space

We consider the space of paths

$$
\Omega=\Omega\left(L_{0}, L_{1}\right)=\left\{\ell:[0,1] \rightarrow P \mid \ell(0) \in L_{0}, \ell(1) \in L_{1}\right\} .
$$

On this space, we are given the action 1 -form $\alpha$ defined by

$$
\alpha(\ell)(\xi)=\int_{0}^{1} \omega(\dot{\ell}(t), \xi(t)) d t
$$

for each tangent vector $\xi \in T_{\ell} \Omega$. From this expression, it follows that

$$
\operatorname{Zero}(\alpha)=\left\{\ell_{p}:[0,1] \rightarrow M \mid p \in L_{0} \cap L_{1}, \ell_{p} \equiv p\right\} .
$$

Using the Lagrangian property of ( $L_{0}, L_{1}$ ), a straightforward calculation shows that this form is closed. Note that $\Omega\left(L_{0}, L_{1}\right)$ is not connected but has countably many connected components. We work on a particular fixed connected component of $\Omega\left(L_{0}, L_{1}\right)$. We pick up a based path $\ell_{0} \in \Omega\left(L_{0}, L_{1}\right)$ and consider the corresponding component $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$; then we define a covering space

$$
\pi: \widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \Omega\left(L_{0}, L_{1} ; \ell_{0}\right)
$$

on which we have a single-valued action functional such that $d \mathcal{A}=-\pi^{*} \alpha$.
The base path $\ell_{0}$ automatically picks out a connected component from each of $L_{0}$ and $L_{1}$ as its initial and final points $x_{0}=\ell_{0}(0) \in L_{0}, x_{1}=\ell_{0}(1) \in L_{1}$. Then $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ is a subspace of the space of paths between the corresponding connected components of $L_{0}$ and $L_{1}$, respectively. Because of this we always assume that $L_{0}, L_{1}$ are connected from now on, unless otherwise stated.

Next we describe the Novikov covering of the component $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ of $\Omega\left(L_{0}, L_{1}\right)$. We first start by describing the universal covering space of $\Omega\left(L_{0}, L_{1}\right.$; $\left.\ell_{0}\right)$. Consider the set of all pairs $(\ell, w)$ such that $\ell \in \Omega\left(L_{0}, L_{1}\right)$ and such that $w:[0,1]^{2} \rightarrow M$ satisfies the boundary condition

$$
\left\{\begin{array}{l}
w(0, \cdot)=\ell_{0}, w(1, \cdot)=\ell \\
w(s, 0) \in L_{0}, w(s, 1) \in L_{1} \quad \text { for all } s \in[0,1] .
\end{array}\right.
$$

Considering $w$ as a continuous path $s \mapsto w(s, \cdot)$ in $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ from $\ell_{0}$ and $\ell$, the fiber at $\ell$ of the universal covering space of $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ can be represented by the set of path homotopy classes of $w$ relative to its end $s=0,1$.

### 3.2. The $\Gamma$-equivalence

We define a covering space of $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ by modding out the space of paths in $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ by another equivalence relation that is weaker than the homotopy. The deck transformation group of this covering space is shown to be abelian by construction.

Note that when we are given two pairs $(\ell, w)$ and $\left(\ell, w^{\prime}\right)$ from $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$, the concatenation

$$
\bar{w} \# w^{\prime}:[0,1] \times[0,1] \rightarrow M
$$

defines a loop $c: S^{1} \rightarrow \Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$. One may regard this loop as a map

$$
\begin{equation*}
C: S^{1} \times[0,1] \rightarrow M \tag{3.1}
\end{equation*}
$$

satisfying the boundary condition $C(s, 0) \in L_{0}, C(s, 1) \in L_{1}$. Obviously the symplectic area of $C$, denoted by

$$
I_{\omega}(c)=\int_{C} \omega,
$$

depends only on the homotopy class of $C$ satisfying (3.1) and so defines a homomorphism on $\pi_{1}\left(\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)\right)$, which we also denote by

$$
I_{\omega}: \pi_{1}\left(\Omega\left(L_{0}, L_{1}\right), \ell_{0}\right) \rightarrow \mathbb{R}
$$

Next we note that for the map $C: S^{1} \times[0,1] \rightarrow M$ satisfying (3.1), it associates a symplectic bundle pair $(\mathcal{V}, \lambda)$ defined by

$$
\mathcal{V}_{C}=C^{*} T M, \lambda_{C}=c_{0}^{*} T L_{0} \sqcup c_{1}^{*} T L_{1},
$$

where $c_{i}: S^{1} \rightarrow L_{i}$ is the map given by $c_{i}(s)=C(s, i)$ for $i=0,1$. This allows us to define another homomorphism

$$
I_{\mu}: \pi_{1}\left(\Omega\left(L_{0}, L_{1}\right), \ell_{0}\right) \rightarrow \mathbb{Z}, \quad I_{\mu}(c)=\mu\left(\mathcal{V}_{C}, \lambda_{C}\right)
$$

where $\mu\left(\mathcal{V}_{C}, \lambda_{C}\right)$ is the Maslov index of the bundle pair $\left(\mathcal{V}_{C}, \lambda_{C}\right)$.
Using the homomorphisms $I_{\mu}$ and $I_{\omega}$, we define an equivalence relation $\sim$ on the set of all pairs $(\ell, w)$ satisfying (3.1). For given such pair $w, w^{\prime}$, we denote by $\bar{w} \# w^{\prime}$ the concatenation of $\bar{w}$ and $w^{\prime}$ along $\ell$, which defines a loop in $\Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ based at $\ell_{0}$.

## DEFINITION 3.1

We say that $(\ell, w)$ is $\Gamma$-equivalent to $\left(\ell, w^{\prime}\right)$ and write $(\ell, w) \sim\left(\ell, w^{\prime}\right)$ if the following conditions are satisfied: $I_{\omega}\left(\bar{w} \# w^{\prime}\right)=0=I_{\mu}\left(\bar{w} \# w^{\prime}\right)$. We denote the set of equivalence classes $[\ell, w]$ by $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$ and call it the Novikov covering space.

There is a canonical lifting of $\ell_{0} \in \Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ to $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$ : this is just $\left[\ell_{0}, \widetilde{\ell}_{0}\right] \in \widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$, where $\widetilde{\ell}_{0}$ is the map $\widetilde{\ell}_{0}:[0,1]^{2} \rightarrow M$ with $\widetilde{\ell}_{0}(s, t)=\ell_{0}(t)$. In this way, $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$ also has a natural base point which we suppress from the notation.

We denote by $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$ the group of deck transformations of the covering space $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$. It is easy to see that the isomorphism class of $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$ depends only on the connected component containing $\ell_{0}$.

The two homomorphisms $I_{\omega}$ and $I_{\mu}$ push down to homomorphisms

$$
E: \Pi\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \mathbb{R}, \quad \mu: \Pi\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \mathbb{Z}
$$

defined by

$$
E(g)=I_{\omega}[C], \quad \mu(g)=I_{\mu}[C]
$$

for any map $C: S^{1} \times[0,1] \rightarrow M$ representing the class $g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$. By construction, it follows that the group $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$ is an abelian group. We
now define the Novikov ring $\Lambda\left(L_{0}, L_{1} ; \ell_{0}\right)$ associated to the abelian covering $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \Omega\left(L_{0}, L_{1} ; \ell_{0}\right)$ as a completion of the group ring $R\left[\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)\right]$. Here $R$ is a commutative ring with unit.

DEFINITION 3.2
$\Lambda_{k}^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)$ denotes the set of all (infinite) sums

$$
\sum_{\substack{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right) \\ \mu(g)=k}} a_{g}[g]
$$

such that $a_{g} \in R$ and, for each $C$, the set $\left\{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right) \mid E(g) \leq C, a_{g} \neq 0\right\}$ is of finite order. We put

$$
\Lambda^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)=\bigoplus_{k} \Lambda_{k}^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)
$$

The ring structure on $\Lambda^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)$ is defined by the convolution product

$$
\left(\sum_{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)} a_{g}[g]\right) \cdot\left(\sum_{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)} b_{g}[g]\right)=\sum_{g_{1}, g_{2} \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)} a_{g_{1}} b_{g_{2}}\left[g_{1} g_{2}\right] .
$$

It is easy to see that the term in the right-hand side is indeed an element in $\Lambda^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)$. Thus $\Lambda^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)=\bigoplus_{k} \Lambda_{k}^{R}\left(L_{0}, L_{1} ; \ell_{0}\right)$ becomes a graded ring under this multiplication. We call this graded ring the Novikov ring associated to the pair $\left(L_{0}, L_{1}\right)$ and the connected component containing $\ell_{0}$.

We also use the universal Novikov ring $\Lambda_{\text {nov }}$ in this paper. We recall its definition here. An element of $\Lambda_{\text {nov }}$ is a formal sum $\sum a_{i} T^{\lambda_{i}} e^{\mu_{i}}$ with $a_{i} \in \mathbb{C}$, $\lambda_{i} \in \mathbb{R}, \mu_{i} \in \mathbb{Z}$ such that $\lambda_{i} \leq \lambda_{i+1}$ and $\lim _{i \rightarrow \infty} \lambda_{i}=\infty . T$ and $e$ are formal parameters. We define a valuation $\mathfrak{v}: \Lambda_{\text {nov }} \rightarrow \mathbb{R}_{\geq 0}$ defined by

$$
\mathfrak{v}\left(\sum_{i=1}^{\infty} a_{i} T^{\lambda_{i}} e^{\mu_{i}}\right)=\lambda_{1} .
$$

We denote the corresponding valuation ring by

$$
\Lambda_{0, \text { nov }}=\left\{\sum_{i=1}^{\infty} a_{i} T^{\lambda_{i}} e^{\mu_{i}} \in \Lambda_{\mathrm{nov}} \mid \lambda_{i} \geq 0\right\} .
$$

It carries a unique maximal ideal consisting of $\sum a_{i} T^{\lambda_{i}} e^{\mu_{i}}$ with $\lambda_{i}>0$ for all $i$ which we denote by $\Lambda_{0, \text { nov }}^{+}$. We have a natural embedding

$$
\Lambda^{R}\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \Lambda_{\mathrm{nov}}
$$

given by

$$
\begin{equation*}
\sum_{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)} b_{g}[g] \mapsto \sum_{g \in \Pi\left(L_{0}, L_{1} ; \ell_{0}\right)} b_{g} T^{\omega(g)} e^{\mu(g) / 2} . \tag{3.2}
\end{equation*}
$$

Now for a given pair $(\ell, w)$, we define the action functional

$$
\mathcal{A}: \widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right) \rightarrow \mathbb{R}
$$

by the formula

$$
\mathcal{A}(\ell, w)=\int w^{*} \omega .
$$

It follows from the definition of $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$ that the integral depends only on the $\Gamma$-equivalence class $[\ell, w]$ and so pushes down to a well-defined functional on the covering space $\widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$.

We call an intersection point admissible to $\ell_{01}$ when the associated constant path lies in the same connected component $\Omega\left(L_{0}, L_{1} ; \ell_{01}\right)$ as $\ell_{01}$ in $\Omega\left(L_{0}, L_{1}\right)$.

## LEMMA 3.1

The set $\operatorname{Cr}\left(L_{0}, L_{1} ; \ell_{0}\right)$ of admissible critical points of $\mathcal{A}$ consists of the pairs of the type $\left[\ell_{p}, w\right]$, where $\ell_{p}$ is the constant path with $p \in L_{0} \cap L_{1}$ and $w$ is as in (2.2). $\operatorname{Cr}\left(L_{0}, L_{1} ; \ell_{0}\right)$ is invariant under the action of $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$ and so forms a principal bundle over a subset of $L_{0} \cap L_{1}$ with its fiber isomorphic to $\Pi\left(L_{0}, L_{1} ; \ell_{0}\right)$.

We put

$$
\operatorname{Cr}\left(L_{0}, L_{1}\right)=\bigcup_{\ell_{0, i}} \operatorname{Cr}\left(L_{0}, L_{1} ; \ell_{0, i}\right),
$$

where $\ell_{0, i}$ runs over the set of base points of connected components of $\Omega\left(L_{0}, L_{1}\right)$.
Next we assign an absolute Morse index to each critical point of $\mathcal{A}$. In general, assigning such an absolute index is not a trivial matter because the obvious Morse index of $\mathcal{A}$ at any critical point is infinite. For this purpose, we use the Maslov index of a certain bundle pair naturally associated to the critical point $\left[\ell_{p}, w\right] \in \operatorname{Cr}\left(L_{0}, L_{1} ; \ell_{0}\right)$.

We call this Morse index of $\left[\ell_{p}, w\right]$ the Maslov-Morse index (relative to the base path $\ell_{0}$ ) of the critical point. The definition of the index somewhat resembles that of $\mathcal{A}$. However, to define this we also need to fix a section $\lambda^{0}$ of $\ell_{0}^{*} \Lambda(M)$ such that

$$
\lambda^{0}(0)=T_{\ell_{0}(0)} L_{0}, \quad \lambda^{0}(1)=T_{\ell_{0}(1)} L_{1}
$$

Here $\Lambda(M)$ is the bundle of Lagrangian Grassmanians of $T M$,

$$
\Lambda(M)=\bigcup_{p \in M} \Lambda\left(T_{p} M\right)
$$

where $\Lambda\left(T_{p} M\right)$ is the set of Lagrangian subspaces of the symplectic vector space $\left(T_{p} M, \omega_{p}\right)$.

Let $\left[\ell_{p}, w\right] \in \operatorname{Cr}\left(L_{0}, L_{1} ; \ell_{0}\right) \subset \widetilde{\Omega}\left(L_{0}, L_{1} ; \ell_{0}\right)$ be an element whose projection corresponds to the intersection point $p \in L_{0} \cap L_{1}$.

As before, we associate a symplectic bundle pair $\left(\mathcal{V}_{w}, \lambda_{w}\right)$ over the square $[0,1]^{2}$, which is defined uniquely up to the homotopy. We first choose $\mathcal{V}_{w}=$
$w^{*} T M$. To define $\lambda_{w}$, let us choose a path $\alpha^{p}:[0,1] \rightarrow \Lambda\left(T_{p} M, \omega_{p}\right)$ satisfying

$$
\left\{\begin{array}{l}
\alpha^{p}(0)=T_{p} L_{0}, \alpha^{p}(1)=T_{p} L_{1} \subset T_{p} M, \\
\left(\alpha^{p}\right)(t) \oplus T_{p} L_{0}=T_{p} M, \\
\alpha^{p}(t) \in U_{0}\left(T_{p} L_{0}\right) \quad \text { for small } t,
\end{array}\right.
$$

where $U_{0}\left(T_{p} L_{0}\right)$ is as above.
Then we consider a continuous Lagrangian subbundle $\lambda_{w} \rightarrow \partial[0,1]^{2}$ of $\left.\mathcal{V}\right|_{\partial[0,1]^{2}}$ by the following formula: the fiber at each point of $\partial[0,1]^{2}$ is given as

$$
\left\{\begin{array}{l}
\lambda_{w}(s, 0)=T_{w(s, 0)} L_{0}, \lambda_{w}(1, t)=\alpha^{p}(t), \\
\lambda_{w}(s, 1)=T_{w(s, 1)} L_{1}, \lambda_{w}(0, t)=\lambda^{0}(0, t) .
\end{array}\right.
$$

It follows that the homotopy type of the bundle pair constructed as above does not depend on the choice of $\alpha^{p}$ either.

## DEFINITION 3.3

We define the Maslov-Morse index of $\left[\ell_{p}, w\right]$ (relative to $\lambda^{0}$ ) by

$$
\mu\left(\left[\ell_{p}, w\right] ; \lambda^{0}\right)=\mu\left(\mathcal{V}_{w}, \lambda_{w}\right) .
$$

### 3.3. Anchored Lagrangian submanifolds

Now we generalize this construction for a chain

$$
\mathcal{L}=\left(L_{0}, \ldots, L_{k}\right)
$$

of more than two Lagrangian submanifolds, that is, with $k \geq 2$.
To realize this purpose, we need the notion of anchors of Lagrangian submanifolds.

## DEFINITION 3.4

Fix a base point $y$ of an ambient symplectic manifold $(M, \omega)$. Let $L$ be a Lagrangian submanifold of $(M, \omega)$. We define an anchor of $L$ to $y$ as a path $\gamma:[0,1] \rightarrow M$ such that

$$
\gamma(0)=y, \quad \gamma(1) \in L .
$$

We call a pair $(L, \gamma)$ an anchored Lagrangian submanifold.

It is easy to see that any homotopy class of path in $\Omega\left(L, L^{\prime}\right)$ can be realized by a path that passes through the given point $y$. We denote the set of homotopy classes of the anchors $\gamma$ to $y \in M$ by $\pi_{1}(y, L)$.

The following lemma is easy to check.

LEMMA 3.2
Suppose that $L$ is connected. Then the set $\pi_{1}(y, L)$ of homotopy classes relative to the ends is a principal homogeneous space of $\pi_{1}(M, L)$; that is, it is a $\pi_{1}(M, L)-$ torsor. We call an element of $\pi_{1}(y, L)$ an anchor class of $L$ relative to $y$.

Proof
Since $L$ is connected, the natural map $\pi_{1}(M) \rightarrow \pi_{1}(M, L)$ is surjective, and so $\pi_{1}(M, L) \cong \pi_{1}(M) / \operatorname{im}\left(\pi_{1}(L) \rightarrow \pi_{1}(M)\right)$ forms a group. It is obvious to see that $\pi_{1}(M, L)$ acts on $\pi_{1}(y, L)$ by concatenation of paths on the right. By definition, this action is free. Transitivity is obvious by definition.

For a given pair $\left(\mathcal{L}, \mathcal{L}^{\prime}\right)$ of anchored Lagrangians $\mathcal{L}=(L, \gamma), \mathcal{L}^{\prime}=\left(L^{\prime}, \gamma^{\prime}\right)$, we denote

$$
\Omega\left(\mathcal{L}, \mathcal{L}^{\prime}\right):=\Omega\left(L, L^{\prime} ; \bar{\gamma} \# \gamma^{\prime}\right),
$$

where the latter is the path component of $\Omega\left(L, L^{\prime}\right)$ containing $\bar{\gamma} \# \gamma^{\prime}$. We also denote

$$
\begin{equation*}
\mathcal{L} \cap \mathcal{L}^{\prime}=\left\{p \in L \cap L^{\prime} \mid \widehat{p} \in \Omega\left(\mathcal{L}, \mathcal{L}^{\prime}\right)\right\} \tag{3.3}
\end{equation*}
$$

Here $\widehat{p}$ is the constant path $\widehat{p}(t) \equiv p$.
When we are given a Lagrangian chain

$$
\mathfrak{L}=\left(L_{0}, L_{1}, \ldots, L_{k}\right),
$$

we also consider a chain of anchors $\gamma_{i}:[0,1] \rightarrow M$ of $L_{i}$ to $y$ for $i=0, \ldots, k$. These anchors give a systematic choice of a base path $\ell_{i j} \in \Omega\left(L_{i}, L_{j}\right)$ by concatenating $\gamma_{i}$ and $\gamma_{j}$ as

$$
\ell_{i j}=\bar{\gamma}_{i} * \gamma_{j},
$$

where $\bar{\gamma}$ is the time reversal of $\gamma$ given by $\bar{\gamma}(t)=\gamma(1-t)$. The upshot of this construction is the following overlapping property:

$$
\begin{align*}
& \ell_{i j}(t)=\ell_{i \ell}(t) \quad \text { for } 0 \leq t \leq \frac{1}{2}  \tag{3.4}\\
& \ell_{i j}(t)=\ell_{\ell j}(t) \quad \text { for } \frac{1}{2} \leq t \leq 1
\end{align*}
$$

for all $j, \ell$.
We write $\vec{p}=\left(p_{10}, \ldots, p_{k(k-1)}\right)$. Let $\chi_{i}=\exp (-2 \pi i \sqrt{-1} / k)$. We consider the set of homotopy class of maps $v: D^{2} \rightarrow M$ such that $v\left(\overline{\chi_{i+1} \chi_{i}}\right) \subset L_{i}$ and $v\left(\chi_{i}\right)=$ $p_{i(i+1)}$. We denote it by $\pi_{2}(\mathfrak{L} ; \vec{p})$. If $\mathcal{E}$ is an anchored Lagrangian chain and $\mathfrak{L}$ is its underlying Lagrangian chain, we write $\pi_{2}(\mathcal{E} ; \vec{p})$ in place of $\pi_{2}(\mathfrak{L} ; \vec{p})$, sometimes by abuse of notation.

## DEFINITION 3.5

Let $\mathcal{E}=\left\{\left(L_{i}, \gamma_{i}\right)\right\}_{0 \leq i \leq k}$ be a chain of anchored Lagrangian submanifolds. A homotopy class $B \in \pi_{2}(\overline{\mathcal{L}} ; \vec{p})$ is called admissible to $\mathcal{E}$ if it can be obtained by a polygon that is a gluing of $k$ bounding strips $w_{i(i+1)}:[0,1] \times[0,1] \rightarrow M$ satisfying

$$
\begin{align*}
& w_{i(i+1)}(0, t)= \begin{cases}\gamma_{i}(2 t-1), & 0 \leq t \leq \frac{1}{2} \\
\gamma_{i+1}(2 t-1), & \frac{1}{2} \leq t \leq 1\end{cases}  \tag{3.5}\\
& w_{i(i+1)}(s, 0) \in L_{i}, \quad w_{i(i+1)}(s, 1) \in L_{i+1}, \tag{3.6}
\end{align*}
$$

$$
\begin{equation*}
w_{i(i+1)}(1, t)=p_{i(i+1)} \tag{3.7}
\end{equation*}
$$

When this is the case, we denote the homotopy class $B$ as

$$
B=\left[w_{01}\right] \#\left[w_{12}\right] \# \cdots \#\left[w_{k 0}\right]
$$

and the set of admissible homotopy classes by $\pi_{2}^{a d}(\mathcal{E} ; \vec{p})$.
We call such a tuple $\mathcal{E}$ an anchored Lagrangian chain.

## REMARK 3.6

We remark that we denote by $\mathfrak{L}$ a chain $\left(L_{0}, \ldots, L_{k}\right)$ of Lagrangian submanifolds and by $\mathcal{E}$ that of anchored Lagrangian submanifolds.

When the collection $\mathcal{E}=\left\{\left(L_{i}, \gamma_{i}\right)\right\}_{0 \leq i \leq k}$ is given, we note that not all homotopy classes in $\pi_{2}(\mathcal{L} ; \vec{p})$ are admissible. But we have the following basic lemma, which is enough for the construction of a Fukaya category, whose proof is easy and so is omitted.

LEMMA 3.3
Let $w_{i(i+1)}$ be given for $i=1, \ldots, k$, and let $B \in \pi_{2}(\mathcal{L} ; \vec{p})$. Then they canonically define a class $\left[w_{k 0}\right] \in \pi_{1}\left(\ell_{k 0}, p_{k 0}\right)$ by

$$
\left[w_{k 0}\right]:=\left[\bar{w}_{01}\right] \# \cdots \#\left[\bar{w}_{(k-1) k}\right] \# B .
$$

The following basic identity immediately follows from the definitions.

## PROPOSITION 3.4

Suppose that $B \in \pi_{2}(\mathcal{E}, \vec{p})$ is given as in Lemma 3.3, and provide the analytic coordinates at the puncture $z_{j} \in \partial D^{2} \subset D^{2}$ so that all $z_{j}$ are outgoing. Then we have $\omega(B)=\sum_{i=0}^{k} \mathcal{A}\left(\left[p_{i}, w_{i}\right]\right)$ and

$$
\begin{align*}
\omega(B) & =\sum_{i=0}^{k} \mathcal{A}_{\ell_{i(i+1)}}\left(\left[p_{i}, w_{i}\right]\right),  \tag{3.8}\\
\mu(\mathcal{E}, \vec{v} ; B) & =\sum_{i=0}^{k} \mu\left(\left[p_{i}, w_{i}\right] ; \lambda_{i(i+1)}\right) . \tag{3.9}
\end{align*}
$$

In particular, the sums in the right-hand sides do not depend on the choice of $\lambda_{i} \subset \gamma_{i}^{*} T M$.

Here the index $\mu\left(\left[p_{i}, w_{i}\right] ; \lambda_{i(i+1)}\right)$ is the Maslov-Morse index relative to the path $\lambda_{i(i+1)}$ of Lagrangian planes as defined in [FO+1], which provides a coherent grading $\mu: \operatorname{Crit} \mathcal{A} \rightarrow \mathbb{Z}$.

The action functional provides a canonical $\mathbb{R}$-filtration on the set

$$
\mathcal{A}: \operatorname{Crit} \mathcal{A} \rightarrow \mathbb{R} .
$$

In addition, inside the collection of anchored Lagrangian submanifolds $(L, \gamma)$ we are given a coherent system of single-valued action functionals

$$
\mathcal{A}: \widetilde{\Omega_{0}}\left(L_{i}, L_{j} ; \bar{\gamma}_{i} \# \gamma_{j}\right) \rightarrow \mathbb{R}
$$

at one stroke for any pair from the given collection $\mathcal{E}$ of anchored Lagrangian submanifolds.

### 3.4. Relation to the graded Lagrangian submanifolds

Now we go back to the collection of Lagrangian submanifolds with vanishing Maslov class on a Calabi-Yau manifold. In this case, we are able to obtain a canonical Lagrangian path $\lambda$ along a given anchor $\gamma$ of $L$ to $y$. We call a pair $(\gamma, \lambda)$ a graded anchor of $L$.

Let $J$ be a compatible almost complex structure. The assumption that $2 c_{1}(M)=0$ implies that the bundle $\Delta=\Lambda^{n}(T M, J)^{\otimes 2}$ is trivial. Choose a section $\Theta$ of $\Delta^{*}$ that has length one everywhere in terms of the metric $g=\omega(\cdot, J \cdot)$. This determines a map $\operatorname{det}_{\Theta}^{2}: \Lambda(M, \omega) \rightarrow S^{1}$ and then an $\infty$-fold Maslov covering by

$$
\begin{equation*}
\Lambda_{\infty}=\left\{(\Lambda, t) \in \Lambda \times \mathbb{R} \mid \operatorname{det}_{\Theta}^{2}(\Lambda)=e^{2 \pi i t}\right\} \tag{3.10}
\end{equation*}
$$

where $s_{L}:\left.L \rightarrow \Lambda(M, \omega)\right|_{L}$ is the natural section defined by the Gauss map $s_{L}(x)=T_{x} L$. An $\Lambda_{\infty}$-grading of a Lagrangian submanifold $L \subset(M, \omega)$ according to Seidel $[\mathrm{Se} 2]$ is just a lift to $\mathbb{R}$ of the map

$$
\operatorname{det}_{\Theta}^{2} \circ s_{L}: L \rightarrow S^{1} .
$$

First of all, the condition on $\mu_{L}=0$ implies that there is such a lifting to $\Lambda_{\infty}$. A grading of $L$ is a lift $\widetilde{s}_{L}$ of $s_{L}$.

We now explain how we give a coherent grading to anchored Lagrangian submanifolds with vanishing Maslov class.

First we go from a graded anchored Lagrangian submanifold to a graded Lagrangian submanifold. We first pick a Lagrangian subspace $V_{y} \in T_{y} M$ such that

$$
\operatorname{det}_{\Theta}^{2}(y)\left(V_{y}\right)=1
$$

We consider a graded anchored Lagrangian submanifold $(L, \gamma, \lambda)$ with $\lambda(0)=V_{y}$. We lift $\lambda$ to a section of $\gamma^{*} \widetilde{\operatorname{Lag}}(T M)$ so that $\widetilde{\lambda}(0)=\left(V_{y}, 0\right)$. Then $\widetilde{\lambda}(1)$ is a lifting of $\lambda(1)$ in $\widetilde{\operatorname{Lag}}\left(T_{\gamma(1)} M\right)$. Since the lifting of $\widetilde{\lambda}$ of $\lambda$ is homotopically unique, $\widetilde{\lambda}(1)$ depends only on $(L, \gamma)$ and the fixed $V_{y}$. Therefore if $\mu_{L}=0$, then this determines a unique grading $\widetilde{s}$ of $L$ with $\widetilde{s}(\gamma(1))=\widetilde{\lambda}(1)$.

DEFINITION 3.7
Let $(M, \omega)$ be such that $2 c_{1}(M)=0$, fix a base point $y \in M$, and let $\Lambda_{y} \in$ $\Lambda(M, \omega)_{y}$ be a Lagrangian subspace. Let $(L, \gamma)$ be a Calabi-Yau anchored Lagrangian. We denote the above common lifting by $\alpha_{\left(L,\left(x_{L}, \gamma_{L}\right)\right)}: L \rightarrow \mathbb{R}$ and call it the canonical grading of $(L, \gamma)$ relative to $\left(y, \Lambda_{y}\right)$.

Next we go from a lifting $\widetilde{s}_{L}$ to $\lambda$. Let $\widetilde{s}_{L}$ be a grading of $L$. Consider any anchored Lagrangian submanifold $(L, \gamma)$ with $\mu_{L}=0$. We take a section $\widetilde{\lambda}$ of the pullback $\gamma^{*}(\widetilde{\operatorname{Lag}}(T M)) \rightarrow[0,1]$ such that

$$
\widetilde{\lambda}(0)=\left(V_{y}, 0\right), \quad \widetilde{\lambda}(1)=\widetilde{s}_{L}\left(\lambda_{i}(1)\right) .
$$

Such a path is unique up to homotopy because $[0,1]$ is contractible, and so $\gamma_{i}^{*} \widetilde{\operatorname{Lag}}(T M)$ is simply connected. We push it out and obtain a section $\lambda$ in $\gamma^{*} \operatorname{Lag}(T M)$. In this way, a graded Lagrangian submanifold ( $L, \widetilde{s}$ ) canonically determines a grading $\lambda$ of an anchored Lagrangian submanifold ( $L, \gamma$ ).

We remark that the path $\lambda_{01}$ induced by these graded anchors lifts to $\widetilde{\lambda}_{01}$ joining $s_{0}\left(\ell_{01}(0)\right)$ to $s_{1}\left(\ell_{01}(1)\right)$.

We then define $\mu\left([p, w] ; \ell_{01}\right)$ using this path $\lambda_{01}$ as in Section 3.2.

LEMMA 3.5
$\mu\left([p, w] ; \ell_{01}\right)$ is independent of $w$.
Proof
Independence of the degree of $w$ is a consequence of our assumption that Maslov indices of $L_{0}, L_{1}$ are zero. We omit the detail.

We refer to [FO+2, Section 9.2] for a more detailed explanation on the above discussion.

Because of this presence of canonical grading associated to $(L, \gamma)$, we drop $\lambda$ from our notation $\mathcal{L}=(L, \gamma, \lambda)$ when we consider Calabi-Yau Lagrangian branes later in this paper.

Adapting to the convention from $[\mathrm{K} 1]$ and $[\mathrm{Se} 2]$, we denote $\widetilde{L}[0]=\left(L, \alpha_{(L, \gamma)}\right)$ and

$$
\widetilde{L}[k]=\left(L, \alpha_{(L, \gamma)}-k\right) .
$$

## 4. Calabi-Yau Lagrangian branes and Dehn twists

We restrict ourselves to the case of $(M, \omega)$ with $2 c_{1}(M, \omega)=0$ and $L \subset M$ whose Maslov class vanishes from now on. We give a precise definition of Calabi-Yau Lagrangian branes in this section. This is the case that is most relevant to mirror symmetry and to the extension of Seidel's long exact sequence of $\mathbb{Z}$-graded symplectic Floer cohomology.

Now we introduce a class of decorated Lagrangian submanifolds on CalabiYau manifold $(M, \omega)$ which we call Calabi-Yau Lagrangian branes.

## DEFINITION 4.1

Let $y \in M$ be a base point, and let $\Lambda_{y} \subset T_{y} M$ be a fixed Lagrangian subspace. Suppose that $\Theta$ is a quadratic complex volume form on $(M, \omega, J)$. Let $\mathcal{E}^{\mathrm{CY}}$ be the Calabi-Yau Lagrangian collection of $(M, \omega)$. We consider the triple $(\mathcal{L}, s,[b])$,
$\mathcal{L}=(L, \gamma)$, which we call an anchored Calabi-Yau Lagrangian brane, that satisfies the following data:
(1) $L$ is a Lagrangian submanifold of $M$ such that the Maslov index of $L$ is zero and $[\omega] \in H^{2}(M, L ; \mathbb{Z})$; we also enhance $L$ with a flat complex line bundle;
(2) $\gamma$ is an anchor of $L$ to $y$;
(3) $s$ is a spin structure of $L$;
(4) $[b] \in \mathcal{M}(L)$ is a bounding cochain described in Section 8.2.

We denote the collection of Calabi-Yau (CY) Lagrangian branes by $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$ and denote the Fukaya category generated by them by $\operatorname{Fuk}\left(\mathcal{E}_{\text {brane }}^{\mathrm{CY}}\right)$.

The first simplification arising from considering the CY Lagrangian collection is that we have only to use the Novikov ring of the form

$$
\begin{equation*}
\Lambda_{\mathrm{nov}}^{(0)}=\left\{\sum_{i=1}^{\infty} a_{i} T^{\lambda_{i}} \mid a_{i} \in \mathbb{Q}, \lambda_{i} \in \mathbb{R}, \lambda_{i} \leq \lambda_{i+1}, \lim _{i \rightarrow \infty} \lambda_{i}=\infty\right\}, \tag{4.1}
\end{equation*}
$$

which becomes a field. We also consider the subring

$$
\begin{equation*}
\Lambda_{0, \mathrm{nov}}^{(0)}=\left\{\sum_{i=1}^{\infty} a_{i} T^{\lambda_{i}} \in \Lambda_{\mathrm{nov}} \mid \lambda_{i} \geq 0\right\} . \tag{4.2}
\end{equation*}
$$

This is because the Maslov index satisfies $\mu(w)=\mu_{L}(\partial w)=0$ for any disc map $w:\left(D^{2}, \partial D^{2}\right) \rightarrow(M, L)$, where $\mu_{L} \in H^{1}(L ; \mathbb{Z})$ is the Maslov class of $L$.

REMARK 4.2
Furthermore, as we mentioned in Section 3.4, the anchor provides a canonical graded structure on a CY Lagrangian brane. Therefore it provides a canonical $\mathbb{R}$-filtration and a $\mathbb{Z}$-grading on

$$
C F\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)=C F\left(\left(L_{0}, \gamma_{0}\right),\left(L_{1}, \gamma_{1}\right)\right):=C F\left(L_{0}, L_{1} ; \bar{\gamma}_{0} \# \gamma_{1}\right)
$$

for any pair $\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ of CY Lagrangian branes, and hence on its cohomology $\operatorname{HF}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$. See Section 8 for a related discussion.

Now we examine the effect of Dehn twists on the CY Lagrangian collection.

## PROPOSITION 4.1

Let $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$ be the associated collection of anchored CY Lagrangian branes. Then $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$ is closed under the action of $\tau_{L}$ 's for all framed Lagrangian spheres $(L,[f])$ and so induces an autoequivalence of $\operatorname{Fuk}\left(\mathcal{E}_{\text {brane }}^{\mathrm{CY}}\right)$.

Proof
We note that the Dehn twist $\tau_{L}$ is a symplectic automorphism. Therefore it pushes the spin structure of $L_{0}$ to the image $\tau_{L}\left(L_{0}\right)$ and pulls back the Maslov class. Therefore the Maslov class of $\tau_{L}\left(L_{0}\right)$ is also zero. Similarly we can push forward the anchor of $L$ to $\tau_{L}\left(L_{0}\right)$. Finally, [FO+1, Theorem B (B.2)] states that
the bounding cochain can also be canonically pushed out under the symplectic automorphism and hence under the action of $\tau_{L}$. This finishes the proof.

Therefore we can ask the question about how the Floer cohomology changes under the Dehn twist along a Lagrangian sphere. The answer is supposed to come from a long exact sequence that Seidel introduced in [Se3] for the context of exact Lagrangian submanifolds. The rest of the paper is occupied with the construction of this long exact sequence for the CY Lagrangian branes on CalabiYau manifolds.

## 5. Lefschetz Hamiltonian fibration and coupling form

In this section, we first recall the basics on smooth Hamiltonian fibrations presented in [GLS] and [En] and extend our discussion to fibrations with Lefschetztype singular fibers. Especially we generalize the notion of coupling form to the current singular fibration and prove the uniqueness of the coupling form on a given Lefschetz Hamiltonian fibration when the fibration $\pi: E \rightarrow \Sigma$ is proper, that is, when the fiber of $E$ is compact.

The notion of Hamiltonian fibrations introduced by Guillemin, Lerman, and Sternberg [GLS] is the family of symplectic manifolds of a fixed isomorphism type, which could be twisted on the parameter space $\Sigma$. On the other hand, Seidel introduced the notion of exact Lefschetz fibrations which could have a finite number of singular fibers of type $A_{1}$-singularity.

Combining [GLS] and [Se3], we give the following definition.

## DEFINITION 5.1 (LEFSCHETZ HAMILTONIAN FIBRATION)

A Lefschetz Hamiltonian fibration over a compact surface $\Sigma$ with boundary $\partial \Sigma$ consists of the data ( $E, \pi, \Omega, J_{0}, j_{0}$ ) as follows.
(1) We have $\partial E=\pi^{-1}(\partial \Sigma)$, and $\left.\pi\right|_{\partial E} \rightarrow \partial \Sigma$ forms a smooth fiber bundle.
(2) The projection $\pi: E \rightarrow \Sigma$ can have at most finitely many critical points, and no two may lie on the same fiber. Denote $E^{\text {crit }} \subset E$ and $\Sigma^{\text {crit }} \subset \Sigma$.
(3) $J_{0}$ is a complex structure on a neighborhood of $E^{\text {crit }}, j_{0}$ is a positively oriented complex structure on a neighborhood of $\Sigma^{\text {crit }}$, and $\pi$ is $\left(J_{0}, j_{0}\right)$-holomorphic near $E^{\text {crit }}$. And the Hessian $D^{2} \pi$ at any critical point is nondegenerate as a complex quadratic form.
(4) $\Omega$ is a closed 2 -form on $E$ which must be nondegenerate on $\operatorname{ker} D \pi_{x}$ for each $x \in E$, and a Kähler form for $J_{0}$ in some neighborhood of $E^{\text {crit }}$.

We say that the fibration $\pi: E \rightarrow \Sigma$ is (symplectically) Calabi-Yau if $c_{1}\left(E_{z}^{v}\right)=0$ at all $z \notin \Sigma^{\text {crit }}$.

We remark that one may allow more than one critical point in the same fiber, which could be useful for the study of a family of Lefschetz Hamiltonian fibrations.

REMARK 5.2
We wish highlight that at a critical point $x \in E$ of $\pi$, Definition 5.1(4) implies that the 2 -form $\Omega_{x}$ is required to be nondegenerate on the whole tangent space $T_{e} E$ since ker $D \pi_{x}=T_{x} E$, while at a regular point $\Omega_{x}$ it is required to be so only on the vertical tangent space $T_{e}^{v} E$ as ker $D \pi_{x}=T_{x}^{v} E$.

When a generic fiber $E_{z}$ with $z \in \Sigma \backslash \Sigma^{\text {crit }}$ is compact, it is proved in [GLS] for a smooth Hamiltonian fibration that the choice of $\Omega$ is uniquely determined by the following additional requirement:

$$
\begin{equation*}
\pi_{*} \Omega^{n+1}=0, \tag{5.1}
\end{equation*}
$$

where $\pi_{*}$ is the integration along the fiber. Now we prove the following analogue to this result for the case of Lefschetz Hamiltonian fibrations.

## THEOREM 5.1

Let $\left(E, \pi, \Omega, J_{0}, j_{0}\right)$ be a Lefschetz Hamiltonian fibration as in Definition 5.1. Then there exists a closed 2 -form $\Omega^{\prime}$ smooth on $E \backslash E^{\text {crit }}$ that satisfies the following:
(1) $\left.\Omega^{\prime}\right|_{T^{v} E}=\left.\Omega\right|_{T^{v} E}$ at all $e \in E \backslash E^{\text {crit }}$,
(2) it satisfies (5.1) on $E \backslash \pi^{-1}\left(\Sigma^{\text {crit }}\right)$.

Furthermore, such a form $\Omega^{\prime}$ is unique. We call such a form the coupling form of $\left(E, \pi, \Omega, J_{0}, j_{0}\right)$.

Proof
We first consider the subset $E \backslash E^{\text {crit }}$. Then the form $\Omega$ induces a splitting

$$
\Gamma: T_{x} E=T_{x} E^{v} \oplus T_{x}^{h} E
$$

at any regular point $x \in E \backslash E^{\text {crit }}$, where the horizontal space is given by

$$
T_{x}^{h} E=\left\{\eta \in T_{x} E \mid \Omega(\eta, \xi)=0, \forall \xi \in T_{x}^{v} E\right\}
$$

and hence induces a natural (Ehresman) connection on $E \backslash E^{\text {crit }}$ whose monodromy is symplectic.

Because of the closedness of $\Omega$, the connection is Hamiltonian (see [GLS]) in that its curvature $\operatorname{curv}(\Gamma)$ has its values contained in $\operatorname{Ham}\left(E_{\pi(e)}\right)$, the set of Hamiltonian vector fields of the fiber $E_{\pi(e)}$, and hence the restriction $\pi$ : $E \backslash E^{\text {crit }} \rightarrow \Sigma \backslash \Sigma^{\text {crit }}$ is a smooth Hamiltonian fibration in the sense of [GLS]. In particular, if we restrict to $E \backslash \pi^{-1}\left(\Sigma^{\text {crit }}\right) \rightarrow \Sigma \backslash \Sigma^{\text {crit }}$, its fibers are all compact, and so we can construct a closed 2 -form $\Omega^{\prime}=\Omega+\pi^{*} \alpha$ on $E \backslash \pi^{-1}$ ( $\left.\Sigma^{\text {crit }}\right)$ for some closed 2 -form $\alpha$ on $\Sigma \backslash \Sigma^{\text {crit }}$ that satisfies (5.1) thereon. In fact, $\Omega^{\prime}$ (and so $d \beta$ ) can be explicitly constructed by requiring

$$
\begin{equation*}
\Omega^{\prime}\left(\eta_{1}^{\sharp}, \eta_{2}^{\sharp}\right)=H_{\eta_{1}, \eta_{2}}, \tag{5.2}
\end{equation*}
$$

where $H_{\eta_{1}, \eta_{2}}$ is the smooth function whose restriction to each fiber over a point in $\Sigma \backslash \Sigma^{\text {crit }}$ is uniquely determined by the following two requirements:
(1) $H_{\eta_{1}, \eta_{2}}$ generates the Lie algebra element $\operatorname{curv}_{\Gamma}\left(\eta_{1}, \eta_{2}\right)$ of $\operatorname{Ham}\left(E_{z}, \omega_{z}\right)$ which is a Hamiltonian vector field,
(2) it satisfies the normalization condition

$$
\int_{E_{z}} H_{\eta_{1}, \eta_{2}} \omega_{z}^{n}=0, \quad \omega_{z}=\left.\Omega\right|_{E_{z}},
$$

for all $z \in \Sigma \backslash \Sigma^{\text {crit }}$.
This finishes the proof.

## DEFINITION 5.3 (COUPLING FORM)

We call the above unique closed 2-form constructed in Theorem 5.1 the coupling form of the Lefschetz Hamiltonian fibration $E \rightarrow \Sigma$.

The following result was essentially proved by Seidel in [Se3, Lemma 1.6]. Seidel proved this for the context of exact Lefschetz fibrations, but the same proof applies if one ignores his consideration of generating functions of $Q$ therein.

## LEMMA 5.2 (SEE [Se3, LEMMA 1.6])

Let $\left(E, \pi, \Omega, J_{0}, j_{0}\right)$ be a Lefschetz Hamiltonian fibration, and let $x_{0}$ be a critical point of $\pi$. Then there are smooth families $\Omega^{\mu} \in \Omega^{2}(E), 0 \leq \mu \leq 1$, such that
(1) $\Omega^{0}=\Omega$;
(2) for all $\mu, \Omega^{\mu}=\Omega^{0}$ outside a small neighborhood of $x_{0}$;
(3) each $\left(E, \pi, \Omega^{\mu}, J_{0}, j_{0}\right)$ is a Lefschetz Hamiltonian fibration;
(4) there is a holomorphic Morse chart $(\xi, \Xi)$ around $x_{0}$ with $\Xi: V \subset \mathbb{C}^{n+1} \rightarrow$ $E$ such that $\Xi^{*} \Omega^{1}$ agree near the origin with the standard forms $\omega_{\mathbb{C}^{n+1}}=(i / 2) \times$ $\sum d x_{k} \wedge d \bar{x}_{k}$.

In fact, if near $E^{\text {crit }}$ we are given a 1-form $\Theta$ with $\Omega=d \Theta$ as in the exact cases, we can also deform the 1-form to $\Theta^{\mu}$ so that $\Xi^{*} \Theta^{1}$ becomes the standard 1-form

$$
\theta_{\mathbb{C}^{n+1}}=\frac{i}{4}\left(\sum x_{k} d \bar{x}_{k}-\bar{x}_{k} d x_{k}\right)
$$

(see [Se3, Lemma 1.6]).

## 6. Exact Lagrangian boundary condition

Now we consider a subbundle $i_{Q}: Q \rightarrow \partial \Sigma$ of the symplectic fiber bundle $\left(\left.E\right|_{\partial \Sigma},\left.\Omega\right|_{\partial \Sigma}\right)$ whose fiber $Q_{z}$ is a Lagrangian submanifold of $\Omega_{z}$ for each $z \in \partial \Sigma$. We call such a $Q$ a fiberwise Lagrangian submanifold of $\left(\left.E\right|_{\partial \Sigma},\left.\Omega\right|_{\partial \Sigma}\right)$.

We start with the notion of exact Lagrangian boundary over $\partial \Sigma$.

DEFINITION 6.1
We call a fiberwise Lagrangian submanifold $Q \subset \partial E$ an exact Lagrangian boundary over $\partial \Sigma$ if there exists a 1 -form $\kappa_{Q}$ on $Q$ such that

$$
\left.\kappa_{Q}\right|_{T(\partial E)^{v}} \equiv 0 \quad \text { and } \quad i_{Q}^{*} \Omega=d \kappa_{Q},
$$

where $i_{Q}: Q \rightarrow \partial E$ is the inclusion map and $T^{v}(\partial E)$ is the vertical tangent space of $\partial E$.

We note that when $\Sigma$ is oriented and the boundary orientation on $\partial \Sigma$ is provided by an orientation 1 -form, denoted by $d \theta$ with $\theta \in \partial \Sigma$, the connection induced by the form $\left.\Omega\right|_{\partial E}$ enables us to express any such 1 -form $\kappa_{Q}$ as

$$
\kappa_{Q}(\theta, x)=h_{i}(\theta, x) d \theta
$$

for $(\theta, x) \in \partial E$ with $\theta \in \partial_{i} \Sigma$ and $h_{i}: \partial E \rightarrow \mathbb{R}$, where $\partial \Sigma=\amalg \partial_{i} \Sigma$. The function $h_{i}$ is unique up to the addition of the function $c_{i}: \partial \Sigma_{i} \rightarrow \mathbb{R}$.

## DEFINITION 6.2

We define $\left\|\kappa_{Q}\right\|_{(1, \infty)}$ by

$$
\left\|\kappa_{Q}\right\|_{(1, \infty)}=\int_{\partial \Sigma} \operatorname{osc}\left(h_{\theta}\right) d \theta
$$

with $\operatorname{osc}\left(h_{\theta}\right):=\max _{x \in E_{x}} h(\theta, x)-\min _{x \in E_{x}} h(\theta, x)$ and call it the $L^{(1, \infty)}$-norm of $\kappa_{Q}$.

We remark that $\left\|\kappa_{Q}\right\|_{(1, \infty)}$ does not depend on the choice of the function $h$.
To give readers some insight on these definitions, we compare this with the classical notion of exact Lagrangian isotopy (see [Gr]).

EXAMPLE 6.3
Let $E=(\mathbb{R} \times[0,1]) \times(M, \omega)$ with $\pi:(\mathbb{R} \times[0,1]) \times(M, \omega) \rightarrow(M, \omega)$ be the projection. Consider the 2 -form $\Omega=\pi^{*} \omega$. Let $L_{i} \subset(M, \omega)$ be a Lagrangian submanifold, and let $\psi_{i}:[0,1] \times L \rightarrow M$ be a Lagrangian isotopy for $i=1,2$. The isotopy $\psi_{i}$ is an exact Lagrangian isotopy if there is a smooth function $h_{i}:[0,1] \times L \rightarrow \mathbb{R}$ such that $\psi_{i}^{*} \omega=d h_{i} \wedge d t=d\left(h_{i} d t\right)$ (see [Gr]). This definition is a special case of Definition 6.1: just consider the embedding

$$
\left(h_{i}, \psi_{i}\right): \mathbb{R} \times L_{i} \rightarrow \mathbb{R} \times[0,1] \times M
$$

set $Q_{i}=\operatorname{im}\left(h_{i}, \psi_{i}\right)$ as the Lagrangian suspension, and let $\kappa_{Q_{i}}(t, x)=h_{i}(t, x) d t$.

We now study the structure of $\pi_{2}(E, Q)$. We start with the following relative version of section class. A class $B \in \pi_{2}(E, Q)$ is a section class if $\pi_{*} B \in \pi_{2}(\Sigma, \partial \Sigma)$ is the positive generator with respect to the given orientation of $\Sigma$. We say that $B$ is a fiber class if it is in the image of $\pi_{2}(M, L) \rightarrow \pi_{2}(E, Q)$ induced from the inclusion of the fiber. The following is proved in [HL, Lemma 2.2] for the smooth

Hamiltonian fibration but the same proof applies to the current case with singular fibers. We refer readers to [HL] for its proof.

LEMMA 6.1
The following sequence of homotopy groups is exact at the middle term:

$$
\pi_{2}(M, L) \rightarrow \pi_{2}(E, Q) \rightarrow \pi_{2}(\Sigma, \partial \Sigma)
$$

The following proposition is the reason why the notion of exact Lagrangian boundary is relevant to the study of pseudoholomorphic curves with boundary later. Similar estimates were previously obtained in [En], [Oh3], and [Se3] in somewhat different contexts.

## PROPOSITION 6.2

Suppose that $\Sigma$ is oriented, and denote by dO a given orientation 1-form on $\partial \Sigma$. Let $Q \subset \partial E$ over $\partial \Sigma$ be an exact Lagrangian boundary of $E$, and let $\kappa_{Q}$ be a corresponding Hamiltonian 1-form. Consider a section s: $\Sigma \rightarrow E \backslash E^{\text {crit }}$ with $s(\partial \Sigma) \subset Q \subset \partial E$. Then for each given section class $[s, \partial s] \in \pi_{2}(E, Q ; \mathbb{Z})$, there exists a constant $C=C\left(\kappa_{Q},[s, \partial s]\right)>0$ such that the integral bound

$$
\left|\int_{\Sigma} s^{*} \Omega\right| \leq C=C(B)
$$

holds for any section s in a fixed class $B=[s, \partial s] \in H_{2}(E, Q ; \mathbb{Z})$. In fact, we have

$$
\begin{equation*}
\left|\int_{\Sigma} s_{2}^{*} \Omega-\int_{\Sigma} s_{1}^{*} \Omega\right| \leq\left\|\kappa_{Q}\right\|_{(1, \infty)} \tag{6.1}
\end{equation*}
$$

for any two such sections with $\left[s_{1}, \partial s_{1}\right]=\left[s_{2}, \partial s_{2}\right]$.
Proof
Recall that $i_{Q}^{*} \Omega=d \kappa_{Q}$ for a 1-form $\kappa_{Q}$ which exists by definition of exact Lagrangian boundary $Q$. Let $s_{i}, i=1,2$, be two sections of $E$ with $s_{i}(\partial \Sigma) \subset Q$ and $\left[s_{1}, \partial s_{1}\right]=\left[s_{2}, \partial s_{2}\right]$. Then we have a geometric chain $(S, C)$ with

$$
\partial S=s_{1} \coprod s_{2} \coprod C,
$$

and we have $\partial C=\partial s_{2}-\partial s_{1}$ as a chain in $Q$.
By Stokes's formula and closedness of $\Omega$, we have

$$
0=\int S^{*}(d \Omega)=\int_{\partial S} \Omega=\int_{\Sigma} s_{2}^{*} \Omega-\int_{\Sigma} s_{1}^{*} \Omega-\int C^{*} \Omega
$$

and hence

$$
\int_{\Sigma} s_{2}^{*} \Omega-\int_{\Sigma} s_{1}^{*} \Omega=\int_{C} \Omega
$$

But by the exactness of the fiberwise Lagrangian subbundle $Q$ and since $C$ has its image in $Q$, we obtain

$$
\int_{C} \Omega=\int_{C} d \kappa_{Q}=\int_{C} \kappa_{Q}=\int_{\partial s_{2}} \kappa_{Q}-\int_{\partial s_{1}} \kappa_{Q} .
$$

Therefore we have obtained

$$
\int_{\Sigma} s_{2}^{*} \Omega-\int_{\Sigma} s_{1}^{*} \Omega=\int_{\partial s_{2}} \kappa_{Q}-\int_{\partial s_{1}} \kappa_{Q}=\int_{\partial \Sigma}\left(h \circ s_{2}-h \circ s_{1}\right) d \theta,
$$

and so

$$
\left|\int_{\Sigma} s_{2}^{*} \Omega-\int_{\Sigma} s_{1}^{*} \Omega\right| \leq \int_{\partial \Sigma}\left(\max _{x \in E_{x}} h_{\theta}(x)-\min _{x \in E_{x}} h_{\theta}(x)\right) d \theta \leq\left\|\kappa_{Q}\right\|_{(1, \infty)} .
$$

Since $\left\|\kappa_{Q}\right\|_{(1, \infty)}$ does not depend on $s$, this finishes the proof.
Next we consider the topological index associated to the section $(s, \partial s)$ for $s$ which does not pass through critical points $E^{\text {crit }}$. By the pullback $s^{*}\left(T E^{v}\right)$, it defines a symplectic bundle pair $\left(s^{*} T E^{v},(\partial s)^{*} T Q^{v}\right)$, where $T Q^{v}=T Q^{v}=\left.T Q \cap T E^{v}\right|_{\partial \Sigma}$. Therefore we can associate the Maslov index, which we denote by $\mu([s, \partial s])$ (see [KL], $[\mathrm{FO}+1]$ ).

Now we examine the topological dependence of $\mu([s, \partial s])$. Note that each section $(s, \partial s)$ defines an element in $\pi_{2}(E, Q)$. We denote the corresponding class by $s_{*}([\Sigma, \partial \Sigma])$, where $[\Sigma, \partial \Sigma]$ is the fundamental class which is a generator of $H_{2}(\Sigma, \partial \Sigma ; \mathbb{Z}) \cong \mathbb{Z}$. The following lemma immediately follows from the definition of the Maslov index for the bundle pair.

LEMMA 6.3
Suppose that $\left(s_{1}\right)_{*}([\Sigma, \partial \Sigma])=\left(s_{2}\right)_{*}([\Sigma, \partial \Sigma])$. Then we have

$$
\mu\left(\left[s_{1}, \partial s_{1}\right]\right)=\mu\left(\left[s_{1}, \partial s_{2}\right]\right) .
$$

## DEFINITION 6.4

We denote by $\pi_{2}^{\sec }(E, Q) \subset \pi_{2}(E, Q)$ the subset of section classes $[s, \partial s]$ in $\pi_{2}(E$, $Q)$. We say that two section classes $B_{1}, B_{2}$ are $\Gamma$-equivalent if they satisfy

$$
\Omega\left(B_{1}\right)=\Omega\left(B_{2}\right), \quad \mu\left(B_{1}\right)=\mu\left(B_{2}\right)
$$

and denote by $\Pi(E, Q)$ the quotient group

$$
\Pi(E, Q)=\pi_{2}^{\sec }(E, Q) / \sim
$$

For the Calabi-Yau Lefschetz fibrations, one can proceed to the study of Maslov indices following the exposition given in [Se4]. Consider the bundle of relative quadratic volume forms

$$
\mathcal{K}_{E / \Sigma}^{2}=\pi^{*} \operatorname{det}_{\mathbb{C}}^{2}(T \Sigma) \otimes \operatorname{det}_{\mathbb{C}}^{\otimes-2} .
$$

By definition, if $E \rightarrow \Sigma$ is Calabi-Yau, we have nowhere a zero section $\eta_{E / B}^{2}$ of this on $E \backslash E^{\text {crit }}$. Furthermore, we can require $\eta_{E / B}^{2}$ to satisfy

$$
\eta_{E / B}^{2}=\frac{\left(d z_{1} \wedge \cdots \wedge d z_{n+1}\right)^{2}}{\left(2 z_{1} d z_{1}+\cdots+2 z_{n+1} d z_{n+1}\right)^{\wedge 2}}
$$

in a neighborhood $U \subset E$ of each critical point under the given identification $\pi: U \backslash\{x\} \rightarrow \pi(U \backslash\{x\})$ with $q: \mathbb{C}^{n+1} \backslash\{0\} \rightarrow \mathbb{C} \backslash 0$.

## DEFINITION 6.5

We say that a fiberwise Lagrangian submanifold $Q \subset E$ is (relatively) graded if there exists a function

$$
\alpha: Q \rightarrow \mathbb{R}
$$

such that

$$
\exp (2 \pi \sqrt{-1} \alpha(x))=\eta_{E / B}^{2}\left(T_{x}^{v} Q\right), \quad x \in Q
$$

We call $\alpha$ an $\mathcal{L}^{\infty}$-grading of $Q$.
For a given pair of a fiberwise Lagrangian submanifold $Q_{1}, Q_{2} \subset E$ intersecting transversely, we can associate a natural $\mathbb{Z}$-grading on the intersection $Q_{1} \cap Q_{2}$ in the following way.

We consider the 2 -form

$$
\omega_{E, \lambda}=\Omega+\lambda \pi^{*} \omega_{\Sigma}
$$

which is nondegenerate on $E \backslash E^{\text {crit }}$. The existence of such $\lambda>0$ is easy to check. The next lemma also immediately follows. We leave the proof to the readers.

LEMMA 6.4
Any fiberwise Lagrangian submanifold $Q \subset E$ is Lagrangian for $\omega_{E, \lambda}$.
Finally, we define the bundle analogue to the relative spin structure introduced in $[\mathrm{FO}+1]$.

## DEFINITION 6.6

Let $E \rightarrow \Sigma$ be a Lefschetz Hamiltonian fibration, and let $Q \subset \partial E \rightarrow \partial \Sigma$ with $Q \subset$ $\partial E$ be a Lagrangian boundary condition such that the vertical tangent bundle $T^{v} Q \rightarrow \partial \Sigma$ is orientable. A fiberwise relative spin structure of $Q \rightarrow \partial \Sigma$ is a pair of a vector bundle $V \rightarrow \Sigma$ and a spin structure on $T^{v} Q \oplus i_{\Sigma}^{*} V$ for the inclusion map $i_{\Sigma}: \partial \Sigma \rightarrow \Sigma$.

Any relative spin structure on $Q$ can be used to induce an orientation on the moduli space $\mathcal{M}(E, Q ; \vec{p} ; B)$ of pseudoholomorphic sections whose explanation is now in order. But in the present paper, we restrict our attention to the case of spin structures.

## 7. Pseudoholomorphic sections

In this section, we perform various studies of geometry and analysis of pseudoholomorphic sections with Lagrangian boundary condition in the setting of Lefschetz Hamiltonian fibrations.

Let $\left(E, \pi, \Omega, J_{0}, j_{0}\right)$ be a Lefschetz Hamiltonian fibration, and let $x_{0} \in E^{\text {crit }}$. Denote by $(\xi, \Xi)$ a holomorphic Morse chart at $x_{0}$, that is, a $j_{0}$-holomorphic coordinates $\xi: U \rightarrow S$ with $\xi(0)=z_{0}=\pi\left(x_{0}\right)$, where $U \subset \mathbb{C}$ is a neighborhood of the
origin and $\Xi: W \rightarrow E$ is a $J_{0}$-holomorphic chart with $W \subset \mathbb{C}^{n+1}$ a neighborhood of the origin in $\mathbb{C}^{n+1}$ with $\Xi(0)=x_{0}$ such that

$$
\begin{equation*}
\left(\xi^{-1} \circ \pi \circ \Xi\right)(x)=x_{1}^{2}+\cdots+x_{n+1}^{2} . \tag{7.1}
\end{equation*}
$$

With $(\xi, \Xi)$ fixed, we denote the model Lefschetz fibration by $q: \mathbb{C}^{n+1} \rightarrow \mathbb{C}$ defined by

$$
q\left(z_{1}, \ldots, z_{n}\right)=z_{1}^{2}+\cdots+z_{n+1}^{2}
$$

Now for the given $\left(E, \pi, \Omega, J_{0}, j_{0}\right)$ and the holomorphic Morse charts $(\xi, \Xi)$ at each of the critical points of $E$, we consider an almost complex structure $J$ on $E$ that satisfies the following:
(1) $J=J_{0}$ in a neighborhood of $E^{\text {crit }}$;
(2) $d \pi \circ J=j \circ d \pi$;
(3) $\left.\Omega(\cdot, J \cdot)\right|_{T E_{x}^{v}}$ is symmetric and positive definite for any $x \in E$.

Following Seidel [Se3], we call such $J$ compatible relative to $j$. An immediate consequence of the definition is the following lemma.

LEMMA 7.1
Let $J$ be compatible relative to $j$. Then for any given area form $\omega_{\Sigma}$ on $\Sigma$ with $\int_{\Sigma} \omega_{\Sigma}=1$, the 2-form $\Omega+\lambda \pi^{*} \omega_{\Sigma}$ tames $J$ for all sufficiently large $\lambda$.

We refer to [Se3] for a more complete explanation of the structure of $J$ 's compatible to $j$.

### 7.1. Energy estimates and Hamiltonian curvature

Next we study the energy estimates of pseudoholomorphic sections in terms of the topological action $\int u^{*} \Omega$ and the contribution coming from the curvature integral of $\int_{\Sigma} u^{*} \Omega$ of a canonical symplectic connection of the Hamiltonian fibration $E \rightarrow \Sigma$ associated to the coupling form $\Omega$ defined in Definition 5.3. This kind of estimate has been studied in [Se3], [Oh3], and [MS].

Using the connection associated to $\Omega$, we decompose $D u=(D u)^{v}+(D u)^{h}$ into vertical and horizontal components. Now we consider the symplectic form

$$
\omega_{E}=\Omega+\lambda \pi^{*} \omega_{\Sigma}
$$

with $\omega_{\Sigma}$ an area form on $\Sigma$ with $\int_{\Sigma} \omega_{\Sigma}=1$. We like to remark that an almost complex structure $J$ compatible to $j$ it is not compatible in the usual sense in that the bilinear form

$$
\Omega(\cdot, J \cdot)
$$

may not be symmetric. However, if $\lambda$ is sufficiently large, it is tame to $\omega_{E}$ (see [Se3, Lemma 2.1]). Therefore we can symmetrize this bilinear form and define
the associate metric $g_{J}$ by

$$
\begin{equation*}
\langle V, W\rangle=g_{J}(V, W):=\frac{1}{2}(\Omega(V, J W)+\Omega(W, J V)) \tag{7.2}
\end{equation*}
$$

We call (7.2) the metric associated to $J$ and denote

$$
|V|^{2}=|V|_{J}^{2}=g_{J}(V, V) .
$$

With respect to this metric, we still have the following basic identity, whose proof we omit.

## LEMMA 7.2

Let $s: \Sigma \rightarrow E$ be any smooth section $v$. Then we have

$$
\begin{equation*}
\frac{1}{2} \int|D s|^{2}=\int s^{*} \omega_{E}+\int\left|\bar{\partial}_{J} s\right|^{2} \tag{7.3}
\end{equation*}
$$

In particular, if $s$ is $J$-holomorphic, then

$$
\frac{1}{2} \int|D s|^{2}=\int s^{*} \omega_{E}
$$

To examine some positivity property of $J$-holomorphic sections, we now decompose $D s=(D s)^{v}+(D s)^{h}$ into vertical and horizontal parts and write

$$
|D s|^{2}=\left|(D s)^{v}\right|^{2}+\left|(D s)^{h}\right|^{2}+2\left\langle(D s)^{v},(D s)^{h}\right\rangle .
$$

Then it is straightforward to prove

$$
\begin{equation*}
\left|(D s)^{h}\right|^{2} \omega_{\Sigma}=2\left(s^{*} \Omega+\lambda \omega_{\Sigma}\right) \tag{7.4}
\end{equation*}
$$

by the identity

$$
\begin{aligned}
\sum_{i=1}^{2}\left|(D s)^{h}\left(e_{i}\right)\right|^{2} & =\sum_{i=1}^{2} \omega_{E, \lambda}\left((D s)^{h}\left(e_{i}\right), J(D s)^{h}\left(e_{i}\right)\right) \\
& =\sum_{i=1}^{2}\left(\Omega+\lambda \pi^{*} \omega_{\Sigma}\right)\left((D s)^{h}\left(e_{i}\right), J(D s)^{h}\left(e_{i}\right)\right) \\
& =2\left(\omega_{E}\left((D s)^{h}\left(e_{1}\right),(D s)^{h}\left(e_{2}\right)\right)+\lambda \omega_{\Sigma}\left(e_{1}, e_{2}\right)\right)
\end{aligned}
$$

for an orthonormal frame $\left\{e_{1}, e_{2}\right\}$ and then apply the curvature identity

$$
\left.d\left(\Omega\left(e_{1}^{\#}, e_{2}^{\#}\right)\right)=-\left[e_{1}^{\#}, e_{2}^{\#}\right]\right] \Omega: \text { fiberwise }
$$

Here $e_{i}^{\#}$ is the horizontal lift of $e_{i}$, and we have

$$
\operatorname{curv}_{\Gamma}\left(e_{1}, e_{2}\right)=\Omega\left(e_{1}^{\#}, e_{2}^{\#}\right)
$$

by the definition of coupling form (see [GLS, (1.12)] but with caution on the sign convention). The following is an immediate corollary of (7.4).

## PROPOSITION 7.3

Let $\Omega$ be the coupling form of $E$. Suppose that $\Omega+\lambda \pi^{*} \omega_{\Sigma}$ for $\lambda>0$ is positive,
that is, symplectic for an area form $\omega_{\Sigma}$ on $\Sigma$. Then we have the inequality as a 2-form

$$
\begin{equation*}
s^{*} \Omega+\lambda \omega_{\Sigma} \geq 0 \tag{7.5}
\end{equation*}
$$

for any J-holomorphic section s. In other words, if we write $s^{*} \Omega=f(s) \omega_{\Sigma}$ for a function $f: \Sigma \rightarrow \mathbb{R}$, then we have $f(s)+\lambda \geq 0$.

Proof
Just choose a complex structure $j$ on $\Sigma$ and a $j$-compatible $J$ on $E$. Positivity (7.5) immediately follows from (7.4).

### 7.2. Gromov-Floer moduli space of $J$-holomorphic sections

We first translate the anchor introduced in Section 3 in the setting of pointed Lagrangian boundary conditions in Hamiltonian fibrations over the surface $\Sigma$ with boundary $\partial \Sigma \neq \emptyset$. This is needed to study morphisms between two Floer chain modules constructed via the moduli space of pseudoholomorphic sections of Lefschetz fibrations over $\Sigma$.

Let $y \in M$ be a base point, and let $(L, \gamma),\left(L^{\prime}, \gamma^{\prime}\right)$ be two anchored Lagrangian submanifolds of $(M, \omega)$ which intersect transversely. Let $\gamma$ and $\gamma^{\prime}$ be the paths $\gamma(0)=\gamma^{\prime}(0)=y$ and $\gamma(1) \in L, \gamma^{\prime}(1)=L^{\prime}$ given as in the anchor data. We denote by $\widetilde{\gamma}$ the time reversal of $\gamma$, that is, the path defined by

$$
\widetilde{\gamma}(t)=\gamma(1-t) .
$$

Now to each intersection $p \in L \cap L^{\prime}$ we associate a Hamiltonian fibration over $[0,1]^{2}$. The paths $\gamma$ and $\gamma^{\prime}$ provide a path in $M$ along $\{0\} \times[0,1]$ via the obvious concatenation of $\widetilde{\gamma}$ and $\gamma^{\prime}$ with the midpoint given by $y$.

We take the trivial fibrations $E=[0,1]^{2} \times(M, \omega) \rightarrow[0,1]^{2}$. Then for each given pair $[p, w]$ with $p \in L \cap L^{\prime}$ and with a bounding strip $w:[0,1]^{2} \rightarrow M$ such that

$$
\begin{align*}
& w(0, t)=\widetilde{\gamma} \# \gamma^{\prime}(t), \quad w(1, t) \equiv p,  \tag{7.6}\\
& w(s, 0) \in L_{0}, \quad w(s, 1) \in L_{1},
\end{align*}
$$

we can associate a section of $s_{[p, w]}:[0,1]^{2} \rightarrow E$ by

$$
(s, t) \mapsto((s, t), w(s, t))[0,1]^{2} \rightarrow[0,1]^{2} \times(M, \omega) .
$$

We call this fibration over $[0,1]^{2}$ with a section $s_{[p, w]}$ an anchor cap associated to $[p, w]$ relative to the given anchor. For notational convenience, we denote the corresponding fibration with the fiberwise Lagrangian submanifolds by

$$
\left(E_{[p, w]} ;[0,1] \times\{0\} \times L_{0},[0,1] \times\{1\} \times L_{1} ; s_{[p, w]}\right)
$$

or simply as $\left(E_{[p, w]} ; s_{[p, w]}\right)$.
Note that the set of homotopy classes $[p, w]$ of $w$ relative to the boundary condition (7.6) has one-to-one correspondence with the homotopy class of sections with the obvious corresponding boundary condition on $\partial[0,1]^{2}$.

For a given compact surface $\Sigma$ with marked points $\vec{\zeta}=\left\{\zeta_{0}, \ldots, \zeta_{k}\right\}$, we consider the corresponding surface $\dot{\Sigma}=\Sigma \backslash \vec{\zeta}$ with punctures. We denote the given preferred holomorphic chart $\varphi_{\zeta}: D_{\zeta} \subset \Sigma \rightarrow D^{+}$of the half-disc $D^{+}=$ $D \cap\{\operatorname{im}(z) \geq 0\}$ with $\varphi_{\zeta}(\zeta)=0$. We also have a local trivialization

$$
\Phi_{\zeta}:\left.E\right|_{D_{\zeta} \backslash\{\zeta\}} \rightarrow D^{+} \backslash\{0\} \times M
$$

lying over $\varphi_{\zeta}$. When $Q \subset E$ is a Lagrangian boundary condition, we have a unique pair $L_{\zeta, \pm}$ of Lagrangian submanifolds of $M$ such that

$$
\Phi_{\zeta}(Q)=[-1,0) \times L_{\zeta,-} \cup(0,1] \times L_{\zeta,+} .
$$

For the given ordered chain of Lagrangian boundaries $\mathcal{Q}=\left(Q_{0}, Q_{1}, \ldots, Q_{k}\right)$, denote $Q=\bigcup_{i} Q_{i}$. Then we require the unique pair $L_{\zeta_{i}, \pm}$ at $\zeta_{i}$ to be

$$
\Phi_{\zeta_{i}}(Q)=[-1,0) \times L_{i} \cup(0,1] \times L_{i+1}
$$

at each $\zeta_{i}$. In this way, for each given $(\pi: E \rightarrow \Sigma ; \mathcal{Q})$ and a chain of intersection points $\vec{p}=\left(p_{0}, \ldots, p_{k}\right)$ with $p_{i} \in L_{i} \cap L_{i+1}$, we consider the moduli space

$$
\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)
$$

of smooth $J$-holomorphic sections for each section class $B \in \pi_{2}(E, \mathcal{Q} ; \vec{p})$.
The following lemma immediately follows from the definition.

LEMMA 7.4
For $\left[p_{i}, w_{i}\right]$ with $i=0, \ldots, k$ realizing $B=\left(\left[\widetilde{w}_{0}\right]\right) \#\left(\left[w_{1}\right] \# \cdots \#\left[w_{k}\right]\right)$, we have the identity

$$
\begin{equation*}
\int w_{0}^{*} \omega=\sum_{i=1}^{k} \int w_{i}^{*} \omega-\Omega(B) \tag{7.7}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
\Omega(B)=\sum_{i=1}^{k} \int w_{i}^{*} \omega-\int w_{0}^{*} \omega . \tag{7.8}
\end{equation*}
$$

Since any section $s$ in class $[s, \partial s]=B$ that satisfies the asymptotic condition $s\left(z_{i}\right)=x_{i}$, with $x_{i} \in Q$ with $x_{i}=\left(p_{i}, u\left(p_{i}\right)\right)$ in the trivialization, satisfies $\int s^{*} \Omega=$ $\Omega(B)$, Lemmas 7.2 and 7.4 imply

$$
\begin{equation*}
\frac{1}{2} \int\|D s\|^{2}=\sum_{i=1}^{k} \int w_{i}^{*} \omega-\int w_{0}^{*} \omega+\lambda<\infty . \tag{7.9}
\end{equation*}
$$

Once we have this energy estimate, it follows by a standard compactness argument that there are only finitely many section classes $B$ (and so $\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$ is nonempty) such that $\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$.

Finally, we state the Gromov-Floer type compactness of $\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$ in a precise way for later use. For this we introduce additional marked points in the interior of $\Sigma$ and on the boundary $\partial \Sigma$ besides the punctures $\vec{\zeta}$.

## DEFINITION 7.1

A configuration on $\dot{\Sigma}$ is the set of finite points consisting of

$$
\begin{array}{ll}
\left(x_{j}\right) \in \operatorname{Int} \Sigma & \text { for } j=1, \ldots, m \\
\left(y_{j}^{(i)}\right) \in \partial_{i} \dot{\Sigma} & \text { for } j=1, \ldots, n_{i} \text { and for } i=0, \ldots, k
\end{array}
$$

We denote by $\widetilde{C}_{m ; \vec{n}}$ the set of such configurations.
We denote such a configuration by

$$
C=\left(\left\{\left(x_{j}\right)\right\}_{1 \leq j \leq m} ;\left\{\left(y_{j}^{(i)}\right)\right\}_{1 \leq j \leq n_{i}}, 0 \leq i \leq k\right)
$$

in general. We note that there are no nontrivial holomorphic automorphisms of $C$ except the following cases:

- $\#(\vec{\zeta})=0$ and $2 m+\sum_{i=0}^{k} n_{i} \leq 2$,
- $\#(\vec{\zeta})=1, m=0$ and $k=1, n_{1}=1$.

We consider the pairs

$$
(s, C) \in \widetilde{\mathcal{M}}_{J}(E, \mathcal{Q} ; \vec{p} ; B) \times \widetilde{C}_{(m ; \vec{n})}, \quad \vec{n}=\left(n_{0}, n_{1}, \ldots, n_{k}\right) .
$$

There is a natural $\operatorname{Aut}(\dot{\Sigma})$-action on the product $\widetilde{\mathcal{M}}_{J}(E, \mathcal{Q} ; \vec{p} ; B) \times \widetilde{C}_{(m ; \vec{n})}$ defined by

$$
(s, C) \mapsto\left(s \circ \phi^{-1}, \phi(C)\right),
$$

where $\phi \in \operatorname{Aut}(\dot{\Sigma})$. We define $\mathcal{M}_{J,(m ; \vec{n})}(E, \mathcal{Q} ; \vec{p} ; B)$ to be the quotient

$$
\mathcal{M}_{J,(m ; \vec{n})}(E, \mathcal{Q} ; \vec{p} ; B)=\widetilde{\mathcal{M}}_{J}(E, \mathcal{Q} ; \vec{p} ; B) \times \widetilde{C}_{(m ; \vec{n})} / \operatorname{Aut}(\dot{\Sigma})
$$

We note that when $C \neq \emptyset$, we have the natural evaluation maps

$$
\mathrm{ev}: \mathcal{M}_{(m ; \vec{n})}(E, \mathcal{Q} ; \vec{p} ; B) \rightarrow E^{m} \times \prod_{i=0}^{k} Q_{i}^{n_{i}}
$$

which respect the above-mentioned $\operatorname{Aut}(\dot{\Sigma})$-action and so are well defined.
We denote by $\overline{\mathcal{M}}_{1}\left(E_{z}, J_{z} ; \alpha_{z}\right)$ the stable maps of genus zero with one marked point and by $\overline{\mathcal{M}}_{1}\left(E_{z}, Q_{z}, J_{z} ; \beta_{z}\right)$ the set of bordered stable maps with one marked point at a boundary of the disc. We consider the fiber product

$$
\begin{align*}
& \widetilde{\mathcal{M}}_{(m ; \vec{n})}\left(E, \mathcal{Q} ; \vec{p} ; B_{0} ;\left\{\alpha_{i}\right\},\left\{\beta_{j}^{i}\right\}\right):=\widetilde{\mathcal{M}}_{(m ; \vec{n})}\left(E, \mathcal{Q} ; \vec{p} ; B_{0}\right)  \tag{7.10}\\
& \mathrm{ev} \times_{\mathrm{ev}}\left(\prod_{i} \mathcal{M}\left(E_{z_{i}}, J_{z_{i}} ; \alpha_{i}\right) \times \prod_{i=0}^{k} \prod_{j=1}^{n_{i}} \mathcal{M}\left(E_{z_{i}}, Q_{z_{i}}, J_{z_{i}} ; \beta_{j}^{(i)}\right)\right)
\end{align*}
$$

with respect to the obvious evaluation maps.

### 7.3. Bubble may hit critical points

In this subsection, we analyze the failure of convergence of a sequence of smooth pseudoholomorphic sections $\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$ with

$$
\begin{equation*}
\operatorname{vir} \operatorname{dim} \mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)=0 \tag{7.11}
\end{equation*}
$$

of symplectic Lefschetz fibrations $E \subset \Sigma$ with the asymptotic condition provided by $\vec{p}$. We are especially interested in bubble components passing through critical points.

By definition, the pullback $\Xi^{*} J$ is the standard complex structure on $\mathbb{C}^{n+1}$, $\xi_{*} j$ is the standard one on $\mathbb{C}$, and $\xi^{-1} \circ \pi \circ \Xi=q$ on the neighborhoods $W \subset \mathbb{C}^{n+1}$, $U \subset \mathbb{C}$ of the origins, corresponding, respectively, to the given holomorphic Morse chart $(\xi, \Xi)$ at each $x_{0} \in E^{\text {crit }}$. We denote by $W_{x_{0}} \subset E, U_{x_{0}} \subset \Sigma$ the corresponding neighborhoods of $x_{0}$ and $\pi\left(x_{0}\right)$, respectively. We also denote by $B^{2 n+2}(r)$ the ball of radius $r>0$ in $\mathbb{C}^{n+1}$ with its center at the origin, and by $B_{x_{0}}(r)$ its image under $\Xi$ at $x_{0} \in E^{\text {crit }}$, and similarly for $B^{2}(\varepsilon)$ and $B_{z_{0}}(\varepsilon)$.

Since we assume that there are finitely many critical points, we can choose constants $\varepsilon, r>0$ and the above-mentioned balls $B_{x_{0}}(r)$ and $B^{2}(\varepsilon)$ so that

$$
\begin{equation*}
B_{z}(\varepsilon) \cap B_{z^{\prime}}(\varepsilon)=\emptyset \quad \text { for } z \neq z^{\prime} \text { with } z=\pi(x), z^{\prime}=\pi\left(x^{\prime}\right) \tag{7.12}
\end{equation*}
$$

for $x, x^{\prime} \in E^{\text {crit }}$ and

$$
\begin{equation*}
\pi\left(B_{x_{0}}(r)\right) \supset B_{z_{0}}(\varepsilon) \quad \text { for all } x_{0} \in E^{\text {crit }} \cap E_{z_{0}} . \tag{7.13}
\end{equation*}
$$

LEMMA 7.5
The graph of any differentiable section does not intersect $E^{\text {crit }}$.

## Proof

Since $s$ is a section, we have $\pi \circ s=\mathrm{id}$. By differentiating this, we obtain $d \pi \circ d s=$ Id. In particular, $d \pi$ is surjective at any point $s(z)$; that is, $s(z)$ must be a regular point of $\pi$, and hence $s(z) \in E \backslash E^{\text {crit }}$.

Obviously we have the inequality

$$
\begin{equation*}
\operatorname{dist}\left(s(z), E^{\text {crit }}\right) \geq C>0 \quad \text { for } z \in \Sigma \backslash \bigcup_{x \in E^{\text {crit }}} B_{\pi(x)}(\varepsilon) \tag{7.14}
\end{equation*}
$$

where $C=C(\varepsilon,(\xi, \Xi))$ is a constant depending only on $\varepsilon$ and the holomorphic Morse chart $(\xi, \Xi)$ independent of $s$. Now we consider the restriction of $s$ on $\bigcup_{x \in E^{\text {crit }}} B_{\pi(x)}(\varepsilon)$. On this neighborhood, we can identify the section $s$ to the holomorphic map

$$
f: B^{2}(\varepsilon) \subset \mathbb{C} \rightarrow B^{2 n+2}(r) \subset \mathbb{C}^{n+1}
$$

satisfying $q \circ f(z)=z$ for all $z \in B^{2}(\varepsilon)$; that is,

$$
f_{1}^{2}(z)+\cdots+f_{n+1}^{2}(z)=z
$$

In particular, we have

$$
|f(z)|^{2}=\sum_{j=1}^{n+1}\left|f_{j}^{2}(z)\right| \geq\left|f_{1}^{2}(z)+\cdots+f_{n+1}^{2}(z)\right|=|z|
$$

for all $z \in B^{2}(\varepsilon)$. Therefore we obtain

$$
\begin{equation*}
|f(z)| \geq \sqrt{\varepsilon} \quad \text { for } z \in \partial B^{2}(\varepsilon) . \tag{7.15}
\end{equation*}
$$

The following theorem is the main theorem proved in this section.

THEOREM 7.6
Suppose that the Lefschetz Hamiltonian fibration with Lagrangian boundary $Q \subset$ $\left.E\right|_{\partial \Sigma}$ such that $E$ is relative Calabi-Yau and $Q$ has vanishing fiberwise Maslov class. Then there exists a dense subset of $j$-compatible $J$ 's such that for any such $J$, there exists a constant $C>0$ depending only on $(E, Q, J, j)$, the section class $[s]$, and $\varepsilon>0$ such that we have

$$
\begin{equation*}
\operatorname{dist}\left(\operatorname{Im} s, E^{\text {crit }}\right) \geq C \tag{7.16}
\end{equation*}
$$

for any smooth section s: $\Sigma \rightarrow E$.
We prove Theorem 7.6 by contradiction. Let $J$ be $j$-compatible and suppose that there is a sequence $s_{i}$ of smooth $J$-holomorphic sections such that

$$
\begin{equation*}
\min _{z \in \Sigma} \operatorname{dist}\left(s_{i}(z), E^{\text {crit }}\right) \rightarrow 0 \tag{7.17}
\end{equation*}
$$

Since $E^{\text {crit }}$ is a finite set and by (7.12), we may choose a critical point $x_{0} \in E^{\text {crit }}$ and a sequence $z_{i} \in \Sigma$ such that

$$
\begin{equation*}
\operatorname{dist}\left(s_{i}\left(z_{i}\right), x_{0}\right) \rightarrow 0 \tag{7.18}
\end{equation*}
$$

By choosing a subsequence of $z_{i}$ if necessary, we may assume that $z_{i} \rightarrow z_{0}$, and so

$$
\begin{equation*}
z_{i} \in B_{z_{0}}(\varepsilon) \tag{7.19}
\end{equation*}
$$

for all $i$. By the Gromov-Floer convergence applied to $J$-holomorphic curves $s_{i}$ : $(\Sigma, j) \rightarrow(E, J)$, which are also $J$-holomorphic sections, there exists a subsequence that converges to

$$
s_{\infty}=s_{0}+\text { bubble components },
$$

where $s_{0}$ is a smooth section of $E \rightarrow \Sigma$ and where each bubble must be either a fiberwise pseudoholomorphic sphere or a fiberwise pseudoholomorphic disc. And each disc bubble has its boundary lying in the given Lagrangian boundary condition.

Due to property (7.19), at least one bubble must pass through the critical point $x_{0}$ whose image is contained in $E_{z_{0}}$. By the connectedness of the image of the limit, this bubble is contained in a bubble tree rooted at a point $z_{1} \in \Sigma$ in the principal component $\left(s_{0}, \Sigma\right)$. The image of this bubble tree itself must be contained in the same fiber $E_{z_{0}}$. Denote this bubble tree by $(v,(C, z)), z \in C$, which is a stable map in $E$ such that

$$
v(z)=s_{0}\left(z_{1}\right) \in E .
$$

However, since $E_{z_{0}}$ contains a singularity $x_{0}$ and so is not a smooth manifold, we need further clarification on the bubble component passing through $x_{0}$. Since $\pi: E \rightarrow \Sigma$ is isomorphic to the standard Lefschetz fibration

$$
q\left(z_{1}, \ldots, z_{n+1}\right)=z_{1}^{2}+\cdots+z_{n+1}^{2}
$$

near $x_{0}$, there is a well-defined multiplicity of the component at the critical point $x_{0}$. We make this statement precise in the next subsection.

### 7.4. Proper holomorphic curves in $E_{z_{0}} \backslash\left\{x_{0}\right\}$

Using the holomorphic Morse chart $(\xi, \Xi)$ at the critical point $x_{0} \in E_{z_{0}} \subset E$, we consider the decomposition

$$
E=B_{x_{0}}(\delta) \coprod\left(E \backslash B_{x_{0}}(\delta)\right),
$$

where $B_{x_{0}}(\delta)=\Xi^{-1}\left(B^{2(n+1)}(\delta)\right) \cap E_{z_{0}}$ for $0<\delta<\varepsilon$.
Now we consider the hypersurface

$$
q^{-1}(0)=\left\{\left(x_{1}, \ldots, x_{n+1}\right) \mid x_{1}^{2}+\cdots+x_{n+1}^{2}=0\right\} .
$$

The only singularity of this hypersurface is $0 \in \mathbb{C}^{n+1}$, and so $q^{-1}(0) \backslash\{0\}$ is a smooth complex hypersurface of $\mathbb{C}^{n+1} \backslash 0$. We denote by $\theta_{\mathbb{C}^{n+1}}$ the 1 -form

$$
\theta_{\mathbb{C}^{n+1}}=\frac{i}{4}\left(\sum x_{k} d \bar{x}_{k}-\bar{x}_{k} d x_{k}\right)
$$

and the standard Kähler form

$$
\omega_{\mathbb{C}^{n+1}}=-d \theta_{\mathbb{C}^{n+1}}\left(=\sum d q_{k} \wedge d p_{k}\right)
$$

where $x_{k}=q_{k}+i p_{k}$. We denote by $\theta$ and $\omega$ the restriction of these to $q^{-1}(0) \backslash\{0\}$.
Following [Se3], we denote $T=T^{*} S^{n}$, and $T(0)$ equals the zero section of $T$, and by $\theta_{T}$ and $\omega_{T}=-d \theta_{T}$, respectively, we denote the standard Liouville 1-form and the standard symplectic form on the cotangent bundle $T$. We identify $T$ with the subset

$$
\left\{(u, v) \in \mathbb{R}^{n+1} \times \mathbb{R}^{n+1} \mid\langle u, v\rangle=0,\|v\|=1\right\}
$$

and then consider the map

$$
\Phi: q^{-1}(0) \backslash\{0\} \rightarrow T \backslash T(0)
$$

defined by

$$
\begin{equation*}
\Phi(x)=\left(\operatorname{im}(\widehat{x})\|\operatorname{re}(\widehat{x})\|, \operatorname{re}(\widehat{x})\|\operatorname{re}(\widehat{x})\|^{-1}\right), \tag{7.20}
\end{equation*}
$$

where $s e^{i \alpha}$ are polar coordinates on the base of $q: \mathbb{C}^{n+1} \rightarrow \mathbb{C}$, and $\widehat{x}=e^{-i \alpha / 2}$. We note that this map is equivariant with respect to the canonical $O(n+1)$-actions on $q^{-1} \backslash\{0\} \subset \mathbb{C}^{n+1} \backslash\{0\}$ and $T=T^{*} S^{n}$.

The following is a consequence of a straightforward calculation, which is a restriction of the identity in [Se3, p. 1014].

## LEMMA 7.7

$\Phi$ is a diffeomorphism such that

$$
\Phi^{*} \theta_{T}=\theta,
$$

and so $\Phi^{*} \omega_{T}=\omega$. In particular, the symplectic manifold $\left(q^{-1}(0) \backslash\{0\}, \omega\right)$ is symplectomorphic to the cotangent bundle $T=T^{*} S^{n}$.

This lemma shows that $E_{z_{0}} \backslash\left\{x_{0}\right\}$ is a symplectic manifold with negative cylindrical end whose asymptotic boundary is symplectomorphic to the unit cosphere bundle $S^{1}\left(T^{*} S^{n}\right)$. Furthermore, complex structure on $\Xi^{-1}(W)$ is required to be induced from the standard complex structure from $q^{-1}(0) \backslash\{0\} \subset \mathbb{C}^{n+1} \backslash\{0\}$. Therefore any $j$-compatible $J$ provides a translational invariant almost complex structure with respect to the cylindrical structure, and any $J$-holomorphic curve is genuinely holomorphic near the end of $q^{-1}(0) \backslash\{0\}$.

In particular, any such curve converges to a Reeb orbit of $S^{1}\left(T^{*} S^{n}\right)$ with finite multiplicity. This motivates us to study the moduli problem of proper $J$-holomorphic curves from $\mathbb{C} \cong \mathbb{C} P^{1} \backslash\{N\}$ with the asymptotic boundary condition given by $(\gamma, k)$, where $\gamma$ is a simple Reeb orbit of $S^{1}\left(T^{*} S^{n}\right)$ and a multiplicity $k \in \mathbb{Z}_{+}$. We denote by $\widetilde{\mathcal{R}}_{1}\left(S^{1}\left(T^{*} S^{n}\right)\right)$ the set of parameterized Reeb orbits on $S^{1}\left(T^{*} S^{n}\right)$ with period $2 \pi$ and by $\mathcal{R}_{1}\left(S^{1}\left(T^{*} S^{n}\right)\right)$ the quotient by the natural $S^{1}$-action, that is, the set of unparameterized Reeb orbits.

We denote by $(s, \Theta)$ the cylindrical coordinates of $T \backslash T(0)$, where $s$ and $\Theta$ are defined by

$$
s(q, p)=\|p\|, \quad \Theta(q, p)=\left(q, \frac{p}{|p|}\right) .
$$

Note that all geodesics on $S^{n}$ have the same period, which implies that the contact manifold $S^{1}\left(T^{*} S^{n}\right)$ is foliated by Reeb orbits, all of which have the same period. We denote the corresponding cylindrical coordinates on $B_{x_{0}}(\varepsilon) \backslash\left\{x_{0}\right\} \subset$ $E \backslash\left\{x_{0}\right\}$ by the same letters $(s, \Theta)$. By a suitable translation of $s$-coordinates, we may assume the identification

$$
(s, \Theta): B_{x_{0}}(\varepsilon) \backslash\left\{x_{0}\right\} \rightarrow(-\infty, 0] \times S^{1}\left(T^{*} S^{n}\right)
$$

Now we are ready to define the moduli space of our interest. For simplicity of notation, we denote

$$
E_{z_{0}} \backslash\left\{x_{0}\right\}=E_{z_{0}}^{*},
$$

and $\dot{S}=S \backslash\left\{z_{0}\right\}$ is either an open Riemann surface isomorphic to $\mathbb{C}$ or an open Riemann surface with boundary isomorphic to $\mathbb{C} \backslash D^{1}(1)$. We fix an analytic coordinate $z=e^{\tau+i t}$ near $z_{0} \in S$. Let $u: \dot{S} \rightarrow E_{z_{0}} \backslash\left\{x_{0}\right\}$ be a pseudoholomorphic curve with Lagrangian boundary condition

$$
u(\partial S) \subset Q_{z_{0}} \subset E_{z_{0}} \backslash\left\{x_{0}\right\}
$$

Since the treatment of the latter is essentially similar to the former, we focus on the former case in the following exposition. We briefly mention the latter case at the end of our discussion.

By the properness and exponential convergence property of $u$ as $\tau \rightarrow-\infty$, we have

$$
\left.\operatorname{Im} u\right|_{(-\infty,-R] \times S^{1}} \subset B_{x_{0}}(\varepsilon)
$$

for a sufficiently large $R>0$. It is proved in $[\mathrm{H}]$ that

$$
T=\lim _{\tau \rightarrow-\infty} \int\left(\Theta \circ u_{\tau}\right)^{*} \lambda
$$

with $T=2 \pi k$ for some integer $k \geq 1$, where $u_{\tau}(t):=u(\tau, t)$. Then it is proved in $[\mathrm{H}]$ and $[\mathrm{Bo}]$ that there exist constants $C, \delta>0$ depending only on $\left(S^{1}\left(T^{*} S^{n}\right), \lambda\right)$ such that

$$
\lim _{\tau-\infty} \operatorname{dist}\left(u(\tau / T, t / T), u_{\gamma}\left(\tau+\tau_{0}, t+t_{0}\right)\right) \leq C e^{-\delta|\tau|}
$$

for some simple Reeb orbit $\gamma$ of $S^{1}\left(T^{*} S^{n}\right)$ and $\tau_{0} \in \mathbb{R}$ and $t_{0} \in S^{1}$. (See also $[\mathrm{FO}+4]$ for a proof of a similar exponential convergence result in the relative context.) Here $u_{\gamma}:[0, \infty) \times S^{1} \rightarrow(-\infty, 0] \times S^{1}\left(T^{*} S^{n}\right)$ denotes the trivial cylinder map $u_{\gamma}(t)=(\tau, \gamma(t))$.

By the above discussion, we can now define the following moduli spaces for each given integer $k \geq 1$ and a homotopy class $A$.

## DEFINITION 7.2

Let $\gamma \in \mathcal{R}_{1}\left(S^{1}\left(T^{*} S^{n}\right)\right)$ and $k \in \mathbb{Z}_{+}$. For each given $(\gamma, k)$, we define

$$
\begin{aligned}
& \mathcal{M}_{z_{0}}^{\mathrm{SFT}}\left(E^{*}, J_{0}, \gamma ; A, k\right)=\left\{u: \dot{\Sigma} \rightarrow E_{z_{0}}^{*} \mid \bar{\partial}_{J_{0}} u=0, \int u^{*} \Omega<\infty\right. \\
&\left.\lim _{\tau \rightarrow-\infty} u(\tau / 2 \pi k, t / 2 \pi k)=u_{\gamma}(t),[u]=A\right\} .
\end{aligned}
$$

We then define

$$
\mathcal{M}_{z_{0}}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)=\bigcup_{\gamma \in \mathcal{R}_{1}\left(S^{1}\left(T^{*} S^{n}\right)\right)} \mathcal{M}_{z_{0}}^{\mathrm{SFT}}\left(E^{*}, J_{0}, \gamma ; A, k\right) .
$$

The following general index formula can be derived from [Bo, Corollary 5.4]. In this regard, we note that the dimension of the space $\mathcal{R}_{1}\left(S^{1}\left(T^{*} S^{n}\right)\right)$ of simple Reeb orbits of $S^{1}\left(T^{*} S^{n}\right)$ is $n$.

## PROPOSITION 7.8

Let $u \in \mathcal{M}_{z_{0}}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)$. Then we have

$$
\begin{equation*}
\operatorname{Index} D_{u} \bar{\partial}_{J}=\left(-\mu_{C Z}(\gamma)+\frac{n}{2}\right)+(n-3)+2 c_{1}\left(u ; \phi_{\gamma}\right) \tag{7.21}
\end{equation*}
$$

where $\mu_{C Z}$ is the generalized Conley-Zehnder index defined by Robbin and Salamon [RS].

For the reader's convenience, we provide the precise definitions of $\mu_{C Z}(\gamma)$ and $c_{1}\left(u ; \phi_{\gamma}\right)$ in the appendix. Once the definitions are made precise, its proof follows from that of [Bo].

The following transversality result is an easy consequence of a standard argument whose proof is omitted. We first note that we have

$$
\mu_{C Z}(\gamma)=\operatorname{Morse}(\gamma)+\frac{\operatorname{dim} \mathcal{R}_{\operatorname{sim}}}{2}
$$

for the Reeb orbit at the negative end (see, e.g., [Mo], [EGH, Corollary 1.7.4] for such a formula). We denote by

$$
\mathcal{M}_{\mathrm{inj}}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right) \subset \mathcal{M}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)
$$

the open subset consisting of somewhere injective curves.

PROPOSITION 7.9
Let $x_{0} \in E^{\text {crit }}$. There exists a dense subset $\mathcal{J}^{\operatorname{tr}}\left(x_{0}\right)$ of the set $\mathcal{J}(j)$ of $\Omega$-compatible almost complex structures such that $\mathcal{M}_{\mathrm{inj}}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)$ is Fredholm-regular and so becomes a smooth manifold of dimension

$$
\begin{equation*}
-\operatorname{Morse}(\gamma)+(n-3)+2 c_{1}\left(u ; \phi_{\gamma}\right) \tag{7.22}
\end{equation*}
$$

for any $k \geq 1$.
An immediate corollary of this proposition is the following vanishing result.

COROLLARY 7.10
Let $x_{0} \in E^{\text {crit }}$ and $E^{*}=E \backslash\left\{x_{0}\right\}$. Suppose that the relative Maslov class of $E \rightarrow \Sigma$ is zero. Then for any $J_{0} \in \mathcal{J}^{\operatorname{tr}}\left(x_{0}\right), \mathcal{M}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)=\emptyset$ for all $A$ and $k$.

Proof
Since any element of $\mathcal{M}^{\mathrm{SFT}}\left(E_{z_{0}}^{*}, J_{0} ; A, k\right)$ is a composition of a somewhere injective curve and a branched covering of the domain, it is enough to prove that

$$
\mathcal{M}_{\mathrm{inj}}^{\mathrm{SFT}}\left(E^{*}, J_{0} ; A, k\right)=\emptyset
$$

for a generic choice of $J_{0}$.
By the assumption, we have $c_{1}(A)=0$ for all $A \in \pi_{2}\left(E_{z_{0}}^{*} ; z_{0}\right)$. Furthermore, the Morse index for the simple closed geodesic of $S^{n}$ is given by $n-1$ and is greater than $n-1$ for multiple geodesics. Therefore we derive

$$
\operatorname{dim} \mathcal{M}_{\mathrm{inj}}^{\mathrm{SFT}}\left(E_{z_{0}}^{*}, J_{0} ; A, k\right) \leq-(n-1)+n-3=-2
$$

for all $k \geq 1$ from (7.22) and hence the proof.
Finally, we briefly mention the case in which $\dot{S} \cong \mathbb{C} \backslash D^{2}(1)$ and $u(\partial S) \subset Q_{z_{0}}$. Since we assume that $Q$ has vanishing fiberwise Maslov class, the corresponding moduli space

$$
\mathcal{M}^{\mathrm{SFT}}\left(E_{z_{0}}^{*}, Q_{z_{0}}, J_{0} ; A, k\right)
$$

has its dimension exactly the same as that of $\mathcal{M}^{\mathrm{SFT}}\left(E_{z_{0}}^{*}, J_{0} ; A, k\right)$. Therefore the same dimension-counting argument applies in exactly the same way (see, e.g., [Mo] for a discussion on the dimension formula in the context of Lagrangian boundary).

We note that the above dimension-counting argument strongly relies on the fact that $E$ is (fiberwise) Calabi-Yau so that $c_{1}\left(u ; \phi_{\gamma}\right)=0$ and $Q$ has vanishing fiberwise Maslov class. It may be interesting to investigate how the long exact sequence will be transformed in other contexts, such as in the Fano case.

### 7.5. Bubble does not hit critical points

In this subsection, we restrict ourselves to the case of vanishing relative first Chern class $c_{1}(E)=0$.

We prove Theorem 7.6 in this subsection. In fact, it is enough to take

$$
\mathcal{J}^{\operatorname{tr}}:=\bigcap_{x_{0} \in E^{\text {crit }}} \mathcal{J}^{\operatorname{tr}}\left(x_{0}\right)
$$

for the dense subset of $j$-compatible $J$ 's.
Proof of Theorem 7.6
Let $J \in \mathcal{J}^{\text {tr }}$ be defined as above.
We have derived before by the Gromov-Floer compactness applied to $\mathcal{M}_{J}(E$, $\mathcal{Q} ; \vec{p} ; B)$ in $E$ that there exists a subsequence which converges to

$$
s_{\infty}=s_{0}+\text { bubble components },
$$

where $s_{0}$ is a smooth section of $E \rightarrow \Sigma$ and each bubble tree is contained in a fiber which consists of either fiberwise pseudoholomorphic spheres or discs. By Corollary 7.10, there cannot be any bubble tree passing through a critical point of $E$ and so contained in $E \backslash E^{\text {crit }}$. Since the principal component $s_{0}$, which is smooth, cannot pass through a critical point by Lemma 7.5, we have proved that the Gromov-Floer limit of $s_{i}: \Sigma \rightarrow E \backslash E^{\text {crit }}$ does not pass any critical point of $E$. Therefore the compactification $\overline{\mathcal{M}}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$ of $\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)$ has image contained in $E \backslash E^{\text {crit }}$.

Once we have achieved this, the rest of the proof follows, in the same way as in the case of smooth Hamiltonian fibration, by the standard dimension-counting argument from the fact that $E \rightarrow \Sigma$ is a Hamiltonian Lefschetz fibration with vanishing relative Maslov class. We consider the evaluation maps

$$
\mathrm{ev}: \mathcal{M}_{(m ; \vec{n})}(E, \mathcal{Q} ; \vec{p} ; B) \rightarrow E^{m} \times \prod_{i=0}^{k} Q_{i}^{n_{i}}
$$

and

$$
\mathrm{ev}_{1}: \mathcal{M}_{1}\left(E_{z}, J_{z} ; \alpha_{z}\right) \rightarrow E_{z}, \quad \mathrm{ev}_{1}: \mathcal{M}_{1}\left(E_{z}, Q_{z}, J_{z} ; \beta_{z}\right) \rightarrow E_{z}
$$

for $z \in \Sigma \backslash \Sigma^{\text {crit }}$ and consider the fiber product

$$
\begin{align*}
\widetilde{\mathcal{M}}_{(m ; \vec{n})}\left(E, \mathcal{Q} ; \vec{p} ; B_{0} ;\left\{\alpha_{i}\right\},\left\{\beta_{j}^{i}\right\}\right) & :=\widetilde{\mathcal{M}}_{(m ; \vec{n})}\left(E, \mathcal{Q} ; \vec{p} ; B_{0}\right) \\
\quad{ }_{\mathrm{ev}} \times_{\mathrm{ev}_{1}}\left(\prod_{i} \mathcal{M}\left(E_{z_{i}}, J_{z_{i}} ; \alpha_{i}\right)\right. & \left.\times \prod_{i=0}^{k} \prod_{j=1}^{n_{i}} \mathcal{M}\left(E_{z_{i}}, Q_{z_{i}}, J_{z_{i}} ; \beta_{j}^{(i)}\right)\right) \tag{7.23}
\end{align*}
$$

with respect to the obvious evaluation maps.
Note that the dimension of the moduli space of a holomorphic sphere in any class in a fiber $E_{z} \backslash E^{\text {crit }}$ has virtual dimension given by $2 n-6$, and so

$$
\operatorname{vir} . \operatorname{dim} \bigcup_{z \in \Sigma} \mathcal{M}_{1}\left(E_{z}, J_{z} ; \alpha_{z}\right)=2 n-4
$$

In particular, for a generic choice of $J$ all somewhere injective holomorphic spheres are regular, and hence we have the condition $c_{1}=0$, which implies semipositivity; the standard argument from [RT] implies that the evaluation image of the moduli space in each fiber has at least codimension 4, which obviously avoids the images of pseudoholomorphic sections in the classes whose associated moduli space has dimension zero. This proves that there cannot be any bubble passing through critical points of $E$ in the limit, and hence the above fiber product, for a generic choice of $J$, becomes empty. The relevant Fredholm theory needed to perform this kind of dimension-counting argument is by now standard. We refer to [Se3, Section 2] for an elegant exposition on this Fredholm theory in the context of exact Lefschetz fibrations, which applies to the current context of Lefschetz Hamiltonian fibrations without change. This proves that there exist a dense subset $\mathcal{J}^{\text {reg,tr }} \subset \mathcal{J}^{\text {tr }}$ such that for any given $(\vec{p} ; B)$ the moduli spaces $\mathcal{M}_{(m ; \vec{n})}(E, \mathcal{Q}, J ; \vec{p} ; B)$ is compact for $J \in \mathcal{J}^{\text {reg,tr }}$.

In particular, there exists a constant $C=C(E, Q, \vec{p} ; B)>0$ :

$$
\min _{z \in \Sigma} \operatorname{dist}\left(s_{\infty}(z), E^{\text {crit }}\right)>C
$$

But this together with Hausdorff convergence of the image of $s_{i}$ to that of $s_{\infty}$ contradicts the assumption that $s_{i}\left(z_{i}\right) \rightarrow x_{0} \in E^{\text {crit }}$. This finishes the proof.

This proposition shows that as far as compactness property of the set of smooth pseudoholomorphic sections is concerned, we can ignore the presence of critical points of the fibration $E \rightarrow \Sigma$.

### 7.6. Gromov-Floer moduli space of $J$-holomorphic trajectories

With Theorem 7.6 at our disposal, we can safely ignore the critical points in the study of smooth $J$-holomorphic sections and their degenerations for the Lefschetz Hamiltonian fibrations $\pi: E \rightarrow \Sigma$. This, together with the energy estimates derived from Section 6 , makes the study essentially the same as for the case of usual smooth Hamiltonian fibrations (without critical points) as studied in [En] and $[\mathrm{MS}]$.

We formulate the definition of Fukaya, Oh, Ohta, and Ono's $\mathfrak{m}_{k}$-maps in the setting of fibrations. Obviously this discussion can be extended to the Lefschetz Hamiltonian fibration $\pi: E \rightarrow \Sigma$ with punctures $\vec{\zeta}=\left\{\zeta_{0}, \ldots, \zeta_{k}\right\} \subset \partial \dot{\Sigma}$ and with a chain of Lagrangian boundary conditions

$$
\mathcal{Q}=\left(Q_{0}, \ldots, Q_{k}\right)
$$

by considering the anchor caps attached to $\left[p_{i j}, w_{i j}\right]$ with $p \in L_{i} \cap L_{j}$. Here $Q_{j}$ is the parallel fiberwise Lagrangian submanifold corresponding to $L_{j}$. We decompose

$$
\partial \dot{\Sigma}=\coprod_{j=0}^{k} \partial_{j} \Sigma
$$

where $\partial_{j} \Sigma$ is the $j$ th connected component of $\partial \dot{\Sigma}$. We briefly add some necessary modification from $[\mathrm{FO}+1]$ to accommodate the possible critical point in the fibration $\pi: E \rightarrow \Sigma$ following the notation from [Se3].

For a given compact surface $\Sigma$ with marked points $\vec{\zeta}$, we consider the corresponding surface $\dot{\Sigma}=\Sigma \backslash \vec{\zeta}$ with punctures. We fix a given preferred holomorphic chart $\varphi_{\zeta}: D_{\zeta} \subset \Sigma \rightarrow D^{+}$of the half-disc $D^{+}=D \cap\{\operatorname{im}(z) \geq 0\}$ with $\varphi_{\zeta}(\zeta)=0$. We consider the moduli space

$$
\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)
$$

as in Section 7.2 for all the section class $B \in \pi_{2}(E, \mathcal{Q} ; \vec{p})$ with

$$
\text { vit.dim } \mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)=0
$$

More specifically, the map is supposed to be given by

$$
\Phi_{0}^{\mathrm{rel}}(E, \pi ; \mathcal{Q})\left(\bigotimes_{j=1}^{k}\left[p_{j}, w_{j}\right]\right)=\sum_{k} \#\left(\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)\right)\left[p_{k}, w_{k}\right]
$$

with a suitable definition of the coefficient $\#\left(\mathcal{M}_{J}(E, \mathcal{Q} ; \vec{p} ; B)\right)$.
Finally, we recall the notion of broken Floer trajectory moduli spaces; that is, the case corresponds to $k+1=2$ (see [FOn] for the corresponding definition for the closed case).

## DEFINITION 7.3

Let $J=\left\{J_{t}\right\}_{0 \leq t \leq 1}$, and let $x, y \in L_{+} \cap L_{-}$. A stable broken Floer trajectory from $p$ to $q$ is a triple

$$
u=\left(\left(u_{1}, \ldots, u_{a}\right) ;\left(\sigma_{1}, \ldots, \sigma_{m}\right),\left(\gamma_{1}^{0}, \ldots, \gamma_{n_{0}}^{0}\right),\left(\gamma_{1}^{1}, \ldots, \gamma_{n_{1}}^{1}\right) ; o\right)
$$

that satisfies the following.
(1) For $i=1, \ldots, a-1, u_{i} \in \mathcal{M}\left(x_{i}, x_{i+1}\right)$ and satisfies

$$
\begin{align*}
u_{1}(-\infty) & =p, u_{a}(\infty)=q \\
u_{i}(\infty) & =u_{i+1}(-\infty) \quad \text { for } i=1, \ldots, a-1 \tag{7.24}
\end{align*}
$$

We call (7.24) the matching condition, and we say that a pair ( $u, u^{\prime}$ ) of Floer trajectories is gluable if it satisfies the matching condition.
(2) We have $\sigma_{i} \in \overline{\mathcal{M}}_{1}\left(J_{t_{i}} ; \alpha_{i}\right)$ for $i=1, \ldots, m$.
(3) We have $\gamma_{j}^{0} \in \overline{\mathcal{M}}_{1}\left(L_{0}, J_{0} ; \beta_{j}^{0}\right)$ for $j=1, \ldots, n_{0}$ and $\gamma_{k}^{1} \in \overline{\mathcal{M}}_{1}\left(L_{1}, J_{1} ; \beta_{k}^{1}\right)$ for $k=1, \ldots, n_{1}$.
(4) For each $\ell=1, \ldots, a$, either the map $u_{\ell}$ is nonstationary or $\Theta_{\ell} \cap \operatorname{Imo} \neq \emptyset$.

We denote the domain of $u$ simply by $\Theta_{u}$, which is the product of a broken configuration. This is the union of finite copies of

$$
\mathbb{R} \times[0,1]
$$

the principal components, and the prestable curves of closed or bordered Riemann surfaces of genus zero are the bubble components with their roots attached to the
principal components of $\Theta_{u}$. Each broken Floer trajectory $u: \Theta_{u} \rightarrow M$ can be regarded as a broken Floer trajectory into the fiber $M_{\zeta}$ for a $\zeta \in \vec{\zeta}$. We denote by $\widetilde{\mathcal{M}}\left(p, q ; B_{0} ;\left\{\alpha_{i}\right\},\left\{\beta_{0, j}\right\},\left\{\beta_{1, k}\right\}\right)$ the set of stable broken Floer trajectories in the prescribed topological types. The group $\operatorname{Aut}(\dot{\Sigma})$ acts on the moduli space $\widetilde{\mathcal{M}}\left(p, q ; B_{0} ;\left\{\alpha_{i}\right\},\left\{\beta_{0, j}\right\},\left\{\beta_{1, k}\right\}\right)$ by the simultaneous translation of the roots of the bubbles attached to the principal component. Then we denote by $\mathcal{M}(p, q ; B)$ the set of all stable broken trajectories in class $B \in \pi_{2}(p, q)$.

We then define

$$
\overline{\mathcal{M}}(E, \mathcal{Q} ; \vec{p} ; B)=\coprod \mathcal{M}_{(m ; \vec{n})}\left(E, \mathcal{Q} ; \vec{p}^{\prime} ; B^{\prime} ;\left\{\alpha_{i}\right\},\left\{\beta_{j}^{i}\right\}\right) \#\left(\prod_{i} \mathcal{M}\left(p_{i}^{\prime}, q_{i} ; B_{i}\right)\right)
$$

for all choices of $B_{0},\left\{\alpha_{i}\right\},\left\{\beta_{j}^{i}\right\}$, and $B_{i}$ 's satisfying

$$
B=B_{0}+\sum_{i} \alpha_{i}+\sum_{i=0}^{k}\left(\sum_{j=1}^{n_{i}} \beta_{j}^{(i)}\right)+\sum B_{i}
$$

and provide it with a topology of stable maps. We denote the corresponding decomposition of maps by

$$
s=s^{0} \#\left(\prod_{i} v_{i}\right) \#\left(\#_{i=0}^{k}\left(\prod_{j} w_{j}^{(i)}\right)\right) \#\left(u_{i}\right),
$$

and we call $s^{0}$ the principal component and other fiberwise curves the bubble components.

At this stage, we emphasize that this compactification is defined as a topological space for any choice of $(j, J)$ for a transversal chain $\left(L_{0}, \ldots, L_{k}\right)$ of Lagrangian submanifolds. The topological space $\overline{\mathcal{M}}(E, \mathcal{Q} ; \vec{p} ; B)$ will not be a smooth "manifold" even for a generic choice of $J$, but will be a space with Kuranishi structure (see [FOn]).

### 7.7. Orientation and signs

Recall that our main object of concern is the collection $\mathcal{E}_{\text {brane }}^{\mathrm{CY}}$ consisting of anchored CY Lagrangian branes and are the CY Lagrangian branes, that is, $((L, \gamma), s,[b])$. We also recall that the action of the Dehn twist $\tau_{L}$ preserves this collection.

In particular, we assume that $L$ is orientable and spin. Because of this we ignore the background class st in the definition of relative spin structures; that is, we restrict ourselves to $s t=0$. The existence of compatible orientations on the ordinary Floer moduli spaces

$$
\mathcal{M}_{k+1}(\mathcal{L}, \vec{p} ; B),
$$

together with some enhancement on their fiber products with chains in Lagrangian submanifolds, is already established in $[\mathrm{FO}+2$, Section 6] based on the analysis of compatible orientations given in [FO+1, Chapter 6]. This establishes all the necessary ingredients needed for the construction of $A_{\infty}$-structures on the Fukaya category $\mathfrak{F u k}\left(\mathcal{E}_{\text {brane }}^{\mathrm{CY}}\right)$.

For the purpose of constructing Seidel-type exact sequences over $\mathbb{Z}$-coefficients, not just $\mathbb{Z}_{2}$-coefficients, in addition to the study performed in $[\mathrm{FO}+1]$ and $[\mathrm{FO}+2]$, we need to study orientation of the moduli space of pseudoholomorphic sections

$$
\mathcal{M}(E, \mathcal{Q}, J ; \vec{p} ; B)
$$

and establish the morphism $\Phi_{0}^{\mathrm{rel}}(E, \pi ; \mathcal{Q})$ over $\mathbb{Z}$. Recall that Seidel considered this morphism over $\mathbb{Z}_{2}$ in [Se3] but extended his construction over $\mathbb{Z}$ by incorporating the study of orientations in [Se4] in the exact framework.

For the reader's convenience, we first borrow the exposition of orientations from $[F O+2]$ in this subsection. We examine the orientation on $\mathcal{M}(E, \mathcal{Q}, J ; \vec{p} ; B)$ later in Section 9.3 in relation to the construction of Seidel's map $\mathfrak{c}$ which is induced by the Lefschetz Hamiltonian fibration associated to the Dehn twist.

Let $p, q \in L_{0} \cap L_{1}$ and $B \in \pi_{2}(p, q)$. We consider $u: \mathbb{R} \times[0,1] \rightarrow M$ such that

$$
\begin{align*}
& \frac{d u}{d \tau}+J \frac{d u}{d t}=0,  \tag{7.25a}\\
& u(\tau, 0) \in L_{0}, \quad u(\tau, 1) \in L_{1}, \quad \int u^{*} \omega<\infty  \tag{7.25b}\\
& u(-\infty, \cdot) \equiv p, \quad u(\infty, \cdot) \equiv q \tag{7.25c}
\end{align*}
$$

It induces a continuous map $\bar{u}:[0,1]^{2} \rightarrow M$ with $\bar{u}(0, t) \equiv p, u(1, t) \equiv q$ in an obvious way. With an abuse of notation, we denote by $[u$ ] the homotopy class of the map $\bar{u}$ in $\pi_{2}(p, q)$. We denote by $\widetilde{\mathcal{M}}^{\circ}(p, q ; B)$ the moduli space consisting of the maps $u$ satisfying (7.25) and compactify $\widetilde{\mathcal{M}}^{\circ}(p, q ; B) / \mathbb{R}$ its quotient by the $\tau$-translations by using an appropriate notion of stable maps as in $[\mathrm{FO}+1$, Section 2.1].

If $\left(L_{0}, L_{1}\right)$ is a relative spin pair, then $\mathcal{M}(p, q ; B)$ is orientable. Furthermore, a choice of relative spin structures gives rise to a compatible system of orientations for $\mathcal{M}(p, q ; B)$ for all pairs $p, q \in L_{0} \cap L_{1}$ and $B \in \pi_{2}(p, q)$. For the sake of completeness, we now recall from [FO+1, Section 8.1]) how the relative spin structure gives rise to a system of coherent orientations.

Let $p \in L_{0} \cap L_{1}$ and $[w] \in \pi_{2}\left(\ell_{01} ; p\right)$ for a relevant map $w:[0,1]^{2} \rightarrow M$ as before. We denote by $\operatorname{Map}\left(\ell_{01} ; p ; L_{0}, L_{1} ; \alpha\right)$ the set of such maps $[0,1]^{2} \rightarrow M$ and denote its homotopy class by $[w]=\alpha$ in $\pi_{2}\left(\ell_{01} ; p\right)$. Let $w \in \operatorname{Map}\left(\ell_{01} ; p ; L_{0}, L_{1} ; \alpha\right)$. Let $\Phi: w^{*} T M \rightarrow[0,1]^{2} \times T_{p} M$ be a (homotopically unique) symplectic trivialization as before. The trivialization $\Phi$, together with the boundary condition $w(0, t)=\ell_{01}(t)$ and the Lagrangian path $\lambda_{01}$ along $\ell_{01}$, defines a Lagrangian path

$$
\lambda^{\Phi}=\lambda_{\left([p, w] ; \lambda_{01}\right)}^{\Phi}:[0,1] \rightarrow T_{p} M
$$

satisfying $\lambda^{\Phi}(0)=T_{p} L_{0}, \lambda^{\Phi}(1)=T_{p} L_{1}$. The homotopy class of this path does not depend on the trivialization $\Phi$ but depends only on $[p, w]$ and the homotopy class of $\lambda_{01}$. Hereafter we omit $\Phi$ from the notation.

We remark that relative spin structure determines a trivialization of $V_{\lambda_{01}(0)} \oplus$ $T_{\lambda_{01}(0)} L_{0}=V_{\lambda_{01}(0)} \oplus \lambda_{01}(0)$ and $V_{\lambda_{01}(1)} \oplus T_{p} L_{1}=V_{\lambda_{01}(1)} \oplus \lambda_{01}(1)$. We take and fix away to extend this trivialization to the family $\ell_{01}^{*} V \oplus \lambda_{01}$ on $[0,1]$.

We consider the following boundary-valued problem for the section $\xi$ of $w^{*} T M$ on $\mathbb{R}_{\geq 0} \times[0,1]$ of a $W^{1, p}$-class such that

$$
\begin{align*}
& D_{w} \bar{\partial}(\xi)=0  \tag{7.26a}\\
& \xi(0, t) \in \lambda_{01}(t), \quad \xi(\tau, 0) \in T_{p} L_{0}, \quad \xi(\tau, 1) \in T_{p} L_{1} \tag{7.26b}
\end{align*}
$$

Here $D_{w} \bar{\partial}$ is the linearization operator of the Cauchy-Riemann equation.
We define $W^{1, p}\left(\mathbb{R}_{\geq 0} \times[0,1], T_{p} M ; \lambda_{01}\right)$ to be the set of sections $\xi$ of $w^{*} T M$ on $\mathbb{R}_{\geq 0} \times[0,1]$ of a $W^{1, p}$-class satisfying (7.26b). Then (7.26a) induces an operator

$$
D_{w} \bar{\partial}: W^{1, p}\left(\mathbb{R}_{\geq 0} \times[0,1], T_{p} M ; \lambda\right) \rightarrow L^{p}\left(\mathbb{R}_{\geq 0} \times[0,1], T_{p} M \otimes \Lambda^{0,1}\right)
$$

which we denote by $\bar{\partial}_{\left([p, w] ; \lambda_{01}\right)}$. The following proposition was proved in $[\mathrm{FO}+1$, Lemma 3.7.69].

## PROPOSITION 7.11

We have

$$
\begin{equation*}
\operatorname{Index} \bar{\partial}_{\left([p, w] ; \lambda_{01}\right)}=\mu\left([p, w] ; \lambda_{01}\right) \tag{7.27}
\end{equation*}
$$

We denote its determinant line by

$$
\operatorname{det} \bar{\partial}_{\left([p, w] ; \lambda_{01}\right)}
$$

By varying $w$ in its homotopy class $\alpha \in \pi_{2}\left(\ell_{01} ; p\right)=\pi_{2}\left(\ell_{01} ; p ; L_{0}, L_{1}\right)$, these lines define a line bundle

$$
\begin{equation*}
\operatorname{det} \bar{\partial}_{\left([p, w] ; \lambda_{01}\right)} \rightarrow \operatorname{Map}\left(\ell_{01} ; p ; L_{0}, L_{1} ; \alpha\right) . \tag{7.28}
\end{equation*}
$$

The bundle (7.28) is trivial if $\left(L_{0}, L_{1}\right)$ is a relatively spin pair (see [FO+1, Section 8.1]).

We need to find a systematic way to orient (7.28) for various $\alpha \in \pi_{2}\left(\ell_{01} ; p\right)$ simultaneously. Following [FO+1, Section 8.1.3], we proceed as follows. Let $\lambda_{p}$ : $[0,1] \rightarrow T_{p} M$ be a path connecting from $T_{p} L_{0}$ to $T_{p} L_{1}$ in $\operatorname{Lag}^{+}\left(T_{p} M, \omega\right)$. The relative spin structure determines a trivialization of $V_{p} \oplus T_{p} L_{0}=V_{p} \oplus \lambda_{p}(0)$ and of $V_{p} \oplus T_{p} L_{1}=V_{p} \oplus \lambda_{p}(1)$. We fix an extension of this trivialization of the $[0,1]-$ parameterized family of vector spaces $V_{p} \oplus \lambda_{p}$. We define

$$
\begin{equation*}
Z_{+}=\left\{(\tau, t) \in \mathbb{R}^{2} \mid \tau \leq 0,0 \leq t \leq 1\right\} \cup\left\{(\tau, t) \mid \tau^{2}+(t-1 / 2)^{2} \leq 1 / 4\right\} \tag{7.29}
\end{equation*}
$$

We consider maps $\xi: Z_{+} \rightarrow T_{p} M$ of the $W^{1, p_{-} \text {-class and study the linear }}$ differential equation
(7.30a) $\bar{\partial} \xi=0$,
(7.30b) $\quad \xi\left(e^{\pi i(t-1 / 2)} / 2+i / 2\right) \in \lambda_{p}(t), \quad \xi(\tau, 0) \in T_{p} L_{0}, \quad \xi(\tau, 1) \in T_{p} L_{1}$.

It defines an operator

$$
W^{1, p}\left(Z_{+}, T_{p} M ; \lambda_{p}\right) \rightarrow L^{p}\left(Z_{+} ; T_{p} M \otimes \Lambda^{0,1}\right)
$$

which we denote by $\bar{\partial}_{\lambda_{p}}$. Let Index $\bar{\partial}_{\lambda_{p}}$ be its index, which is a virtual vector space. The following theorem is proved in [FO+2, Theorem 6.5] whose proof in turn follows that of $[\mathrm{FO}+1$, Chapter 6].

THEOREM 7.12
Let $\left(L_{0}, L_{1}\right)$ be a relatively spin pair of oriented Lagrangian submanifolds. Then for each fixed $\alpha$, the bundle (7.28) is trivial.

If we fix a choice of systems of orientations $o_{p}$ on $\operatorname{Index} \bar{\partial}_{\lambda_{p}}$ for each $p$, then it determines orientations on (7.28), which we denote by $o_{[p, w]}$.

Moreover $o_{p}, o_{[p, w]}$ determine an orientation of $\mathcal{M}(p, q ; B)$ denoted by o $(p, q$; B) by the gluing rule

$$
\begin{equation*}
o_{[q, w \# B]}=o_{[p, w]} \# o(p, q ; B) \tag{7.31}
\end{equation*}
$$

for all $p, q \in L_{0} \cap L_{1}$ and $B \in \pi_{2}(p, q)$ so that they satisfy the gluing formulae

$$
\partial o(p, r ; B)=o\left(p, q ; B_{1}\right) \# o\left(q, r ; B_{2}\right)
$$

whenever the virtual dimension of $\mathcal{M}(p, r ; B)$ is 1 . Here $\partial o(p, r ; B)$ is the induced boundary orientation of the boundary $\partial \mathcal{M}(p, r ; B)$ and $B=B_{1} \# B_{2}$, and $\mathcal{M}(p, q$; $\left.B_{1}\right) \# \mathcal{M}\left(q, r ; B_{2}\right)$ appears as a component of the boundary $\partial \mathcal{M}(p, r ; B)$.

REMARK 7.4
For the definition of the orientation of the moduli spaces for the filtered bimodule structure, see [FO+1, Sections 8.7, 8.8, Definition 8.8.11].

One can generalize the above discussion to the moduli space of pseudoholomorphic polygons, which we describe below.

Consider a disc $D^{2}$ with $k+1$ marked points $z_{0 k}, z_{k(k-1)}, \ldots, z_{10} \subset \partial D^{2}$ respecting the counterclockwise cyclic order of $\partial D^{2}$. We take a neighborhood $U_{i}$ of $z_{i(i-1)}$ and a conformal diffeomorphism $\varphi_{i}: U_{i} \backslash\left\{z_{i(i-1)}\right\} \subset D^{2} \cong(-\infty, 0] \times$ $[0,1]$ of each $z_{i(i-1)}$. For any smooth map

$$
w: D^{2} \rightarrow M ; w\left(z_{i(i-1)}\right)=p_{i(i-1)}, w\left(\overline{z_{(i+1) i} z_{i(i-1)}}\right) \subset L_{i},
$$

we deform $w$ so that it becomes constant on $\varphi_{i}^{-1}((-\infty,-1] \times[0,1]) \subset U_{i}$; that is, $w(z) \equiv p_{i(i-1)}$ for all $z \in \varphi^{-1}((-\infty,-1] \times[0,1])$. So assume that this holds for $w$ from now on. We now consider the Cauchy-Riemann equation

$$
\begin{align*}
& D_{w} \bar{\partial}(\xi)=0  \tag{7.32a}\\
& \xi(\theta) \in T_{w(t)} L_{i} \quad \text { for } \theta \in \overline{z_{(i+1) i} z_{i(i-1)}} \subset \partial D^{2} \tag{7.32b}
\end{align*}
$$

We remark that on $U_{i}=(-\infty, 0] \times[0,1]$ the boundary condition (7.32b) becomes

$$
\begin{equation*}
\xi(s, 0) \in L_{i-1}, \quad \xi(s, 1) \in L_{i} \tag{7.33}
\end{equation*}
$$

Equation (7.32) induces a Fredholm operator, which we denote by

$$
\begin{equation*}
\bar{\partial}_{w ; \mathfrak{L}}: W^{1, p}\left(D^{2} ; w^{*} T M ; \mathfrak{L}\right) \rightarrow L^{p}\left(D^{2} ; w^{*} \otimes \Lambda^{0,1}\right) \tag{7.34}
\end{equation*}
$$

Moving $w$, we obtain a family of Fredholm operators $\bar{\partial}_{(\mathfrak{L} ; \vec{p} ; B)}$ parameterized by a suitable completion of $\mathcal{F}(\vec{p} ; \mathfrak{L} ; B)$ for $B \in \pi_{2}(\vec{p} ; \mathfrak{L})$. Therefore we have a welldefined determinant line bundle

$$
\begin{equation*}
\operatorname{det} \bar{\partial}_{(\mathfrak{L} ; \vec{p} ; B)} \rightarrow \mathcal{F}(\mathfrak{L} ; \vec{p} ; B) \tag{7.35}
\end{equation*}
$$

The following theorem is an extension of Theorem 7.12, which is $[\mathrm{FO}+2$, Theorem 6.7], whose proof we refer readers thereto.

THEOREM 7.13
Suppose that $\mathfrak{L}=\left(L_{0}, \ldots, L_{k}\right)$ is a relatively spin Lagrangian chain. Then each $\operatorname{det} \bar{\partial}_{(\vec{p}, \mathfrak{L} ; B)}$ is trivial.

Moreover, we have the following. If we fix orientations o $o_{p_{i j}}$ on Index $\bar{\partial}_{\lambda_{p_{i j}}}$ as in Theorem 7.12 for all $p_{i j} \in L_{i} \cap L_{j}$, with $L_{i}$ transversal to $L_{j}$, then we have a system of orientations, denoted by $o_{k+1}(\vec{p} ; \mathfrak{L} ; B)$, on the bundles (7.35), so that it is compatible with the gluing map in an obvious sense.

We can prove that the orientation of $\bar{\partial}_{(\mathfrak{L} ; \vec{p} ; B)}$ depends on the choice of $o_{p_{(i+1) i}}$ (and so on $\lambda_{p}$ ) with $i=0, \ldots, k$, but is independent of the choice of $w_{(i+1) i}^{+}$, and so on. This is a consequence of the proof of Theorem 7.12. Therefore the orientation in Theorem 7.13 is independent of the choice of anchors.

## REMARK 7.5

To give an orientation of $\mathcal{M}(\mathfrak{L} ; \vec{p} ; B)$, we have to take the moduli parameters of marked points and the action of the automorphism group into account. We also treat the intersection point $p_{i(i-1)}$ as if it is a chain of codimension $\mu\left(\left[p_{i(i-1)}, w_{i(i-1)}^{+}\right] ; \lambda_{(i-1) i}\right)$ in a way similar to [FO+1, Chapter 8, Section 8.5].

## 8. Anchored Floer cohomology: Review

In the first subsection, we recall the exposition from $[\mathrm{FO}+2]$ on the Lagrangian Floer theory of anchored Lagrangian submanifolds.

### 8.1. Floer chain complex

Let $\left(L_{0}, L_{1}\right)$ be a pair with $L_{0}$ intersecting $L_{1}$ transversely.
Let $\left(L_{i}, \gamma_{i}\right), i=0,1$, be anchored Lagrangian submanifolds. Let $p, q \in L_{0} \cap L_{1}$ be admissible intersection points. We defined the set $\pi_{2}(p, q)=\pi_{2}\left(\left(L_{0}, L_{1}\right),(p, q)\right)$ in Section 3. We also defined $\pi_{2}\left(\ell_{01} ; p\right)$ there. We now define the following.

## DEFINITION 8.1

Let $R$ be the underlying coefficient field. We define $\operatorname{CF}\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)=C F\left(\left(L_{1}, \gamma_{1}\right)\right.$, $\left.\left(L_{0}, \gamma_{0}\right)\right)$ to be a free $R$-module over the basis $[p, w]$, where $p \in \mathcal{L}_{0} \cap \mathcal{L}_{1}$ is an admissible intersection point and $w$ is a map from $[0,1]^{2} \rightarrow M$ connecting $\ell_{01}$ and $\widehat{p}$.

Here $R$ is a ground ring such as $\mathbb{Q}, \mathbb{C}$, or $\mathbb{R}$.

## REMARK 8.2

We remark that the set $[p, w]$, where $p$ is the admissible intersection point, is identified with the set of the critical points of the action functional $\mathcal{A}$ defined on the Novikov covering space of $\Omega\left(L_{0}, L_{1} ; \ell_{01}\right)$. The group $\Pi\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ defined in Section 3.2 acts freely on it so that the quotient space is the set of admissible intersection points.

We next take a grading $\lambda_{i}$ to $\left(L_{i}, \gamma_{i}\right)$ as in Section 3.3. It induces a grading of $[p, w]$ given by $\mu\left([p, w] ; \lambda_{01}\right)$, which gives the graded structure on $C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$

$$
C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)=\bigoplus_{k} C F^{k}\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \lambda_{01}\right)
$$

where $C F^{k}\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \lambda_{01}\right)=\operatorname{span}_{R}\left\{[p, w] \in C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \mid \mu\left([p, w] ; \lambda_{01}\right)=k\right\}$.
For given $B \in \pi_{2}(p, q)$, we denote by $\operatorname{Map}(p, q ; B)$ the set of such $w$ 's in class $B$.

We summarize the extra structures added in the discussion of Floer homology for the anchored Lagrangian submanifolds in the following.
(1) We assume that $\left(L_{0}, L_{1}\right)$ is a relatively spin pair. We consider a pair $\left(L_{0}, \gamma_{0}\right),\left(L_{1}, \gamma_{1}\right)$ of anchored Lagrangian submanifolds and the base path $\ell_{01}=$ $\bar{\gamma}_{0} * \gamma_{1}$.
(2) We fix a grading $\lambda_{i}$ of $\gamma_{i}$ for $i=0,1$, which in turn induce a grading of $\ell_{01}, \lambda_{01}=\overline{\lambda_{0}} * \lambda_{1}$.
(3) We fix an orientation $o_{p}$ of Index $\bar{\partial}_{\lambda_{p}}$ for each $p \in L_{0} \cap L_{1}$ as in [FO+2].

Under these conditions, orientation of the Floer moduli space $\mathcal{M}(p, q ; B)$ is induced. Using the virtual fundamental chain technique (see [FOn], $[\mathrm{FO}+1$, Appendix A.1]), we can take a system of multisections and obtain a system of rational numbers $n(p, q ; B)=\#(\mathcal{M}(p, q ; B))$ whenever the virtual dimension of $\mathcal{M}(p, q ; B)$ is zero. Finally, we define the Floer "boundary" map $\partial: C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \rightarrow$ $C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$, defined in [Fl] and [Oh1], by the sum

$$
\begin{equation*}
\partial([p, w])=\sum_{q \in L_{0} \cap L_{1}} \sum_{B \in \pi_{2}(p, q)} n(p, q ; B)[q, w \# B] . \tag{8.1}
\end{equation*}
$$

By Remark 8.2, $\operatorname{CF}\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ carries a natural $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$-module structure and $C F^{k}\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \lambda_{01}\right)$ carries a $\Lambda^{(0)}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$-module structure, where

$$
\Lambda^{(0)}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)=\left\{\sum a_{g}[g] \in \Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \mid \mu([g])=0\right\} .
$$

We define

$$
\begin{equation*}
C\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)=C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \otimes_{\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)} \Lambda_{\mathrm{nov}}, \tag{8.2}
\end{equation*}
$$

where we use the embedding $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \rightarrow \Lambda_{\text {nov }}$ given in (3.2).
We write the $\Lambda_{\text {nov }}$-module (8.2) also as

$$
C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{\text {nov }}\right) .
$$

## DEFINITION 8.3

We define the energy filtration $F^{\lambda} C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ of the Floer chain complex $C F\left(L_{1}\right.$, $\left.\gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)$ ) (here $\left.\lambda \in \mathbb{R}\right)$ such that $[p, w]$ is in $F^{\lambda} C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ if and only if $\mathcal{A}([p, w]) \geq \lambda$.

This filtration also induces a filtration on (8.2).

REMARK 8.4
We remark that this filtration not only depends on the homotopy class but also on $\gamma_{i}$ 's themselves. A different choice of $\gamma_{i}$ 's induces a global uniform shift of filtration levels.

It is easy to see the following from the definition of $\partial$ above.

LEMMA 8.1
We have

$$
\partial\left(F^{\lambda} C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right) \subseteq F^{\lambda} C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)\right) .
$$

According to definition (8.1) of the map $\partial$, we have the formula for its matrix coefficients

$$
\begin{equation*}
\langle\partial \partial[p, w],[r, w \# B]\rangle=\sum_{q \in L_{0} \cap L_{1}} \sum_{B=B_{1} \# B_{2} \in \pi_{2}(p, r)} n\left(p, q ; B_{1}\right) n\left(q, r ; B_{2}\right) T^{\omega(B)}, \tag{8.3}
\end{equation*}
$$

where $B_{1} \in \pi_{2}(p, q)$ and $B_{2} \in \pi_{2}(q, r)$.
To prove that $\partial \partial=0$, one needs to prove that $\langle\partial \partial[p, w],[r, w \# B]\rangle=0$ for all pairs $[p, w],[r, w \# B]$. On the other hand, it follows from definition that each summand

$$
n\left(p, q ; B_{1}\right) n\left(q, r ; B_{2}\right) T^{\omega(B)}=n\left(p, q ; B_{1}\right) T^{\omega\left(B_{1}\right)} n\left(q, r ; B_{2}\right) T^{\omega\left(B_{2}\right)}
$$

and the coefficient $n\left(p, q ; B_{1}\right) n\left(q, r ; B_{2}\right)$ is nothing but the number of broken trajectories lying in $\mathcal{M}\left(p, q ; B_{1}\right) \# \mathcal{M}\left(q, r ; B_{2}\right)$. This number is nonzero in the general situation in which we work.

To handle the problem of obstruction to $\partial \circ \partial=0$ and of bubbling-off discs in general, a structure of a filtered $A_{\infty}$-algebra ( $C, \mathfrak{m}$ ) with nonzero $\mathfrak{m}_{0}$-term is associated to each Lagrangian submanifold $L$ (see $[\mathrm{FO}+3]$, $[\mathrm{FO}+1]$ ).

## 8.2. $A_{\infty}$-algebra

In this subsection, we review the notion and construction of a filtered $A_{\infty}$-algebra associated to a Lagrangian submanifold. To make the construction consistent with the one in Section 8.1, where $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ is used for the coefficient ring rather than the universal Novikov ring, we rewrite them using the smaller Novikov ring $\Lambda(L)$ which we define below. Let $L$ be a relatively spin Lagrangian submanifold. We have a homomorphism

$$
(E, \mu): H_{2}(M, L ; \mathbb{Z}) \rightarrow \mathbb{R} \times \mathbb{Z},
$$

where $E(\beta)=\beta \cap[\omega]$ and where $\mu$ is the Maslov index homomorphism. We put $g \sim g^{\prime}$ for $g, g^{\prime} \in H_{2}(M, L ;: \mathbb{Z})$ if $E(g)=E\left(g^{\prime}\right)$ and $\mu(g)=\mu\left(g^{\prime}\right)$. We write $\Pi(L)$ as the quotient with respect to this equivalence relation. It is a subgroup of $\mathbb{R} \times \mathbb{Z}$. We define

$$
\begin{aligned}
\Lambda(L)=\left\{\sum c_{g}[g]\right. & \mid g \in \Pi(L), c_{g} \in R, E(g) \geq 0, \\
& \left.\forall E_{0} \#\left\{g \mid c_{g} \neq 0, E(g) \leq E_{0}\right\}<\infty\right\} .
\end{aligned}
$$

We have the natural embedding $\Lambda(L) \rightarrow \Lambda_{0, \text { nov }}$ similarly as in (3.2).
Let $\bar{C}$ be a graded $R$-module, and let $C F=\bar{C} \widehat{\otimes}_{R} \Lambda(L)$. From now on we use the symbol $C F$ for the modules over $\Lambda(L)$ or $\Lambda\left(L_{0}, L_{1}\right)$ and $C$ for the modules over the universal Novikov ring.

We denote by $C F[1]$ its suspension defined by $C F[1]^{k}=C F^{k+1}$. We denote by $\operatorname{deg}(x)=|x|$ the degree of $x \in C$ before the shift and by $\operatorname{deg}^{\prime}(x)=|x|^{\prime}$ that after the degree shift, that is, $|x|^{\prime}=|x|-1$. Define the bar complex $B(C F[1])$ by

$$
B_{k}(C F[1])=(C F[1])^{k \otimes}, \quad B(C F[1])=\bigoplus_{k=0}^{\infty} B_{k}(C F[1]) .
$$

Here $B_{0}(C F[1])=R$ by definition. The tensor product is taken over $\Lambda(L)$. We provide the degree of elements of $B(C F[1])$ by the rule

$$
\begin{equation*}
\left|x_{1} \otimes \cdots \otimes x_{k}\right|^{\prime}:=\sum_{i=1}^{k}\left|x_{i}\right|^{\prime}=\sum_{i=1}^{k}\left|x_{i}\right|-k, \tag{8.4}
\end{equation*}
$$

where $|\cdot|^{\prime}$ is the shifted degree. The ring $B(C F[1])$ has the structure of a graded coalgebra.

## DEFINITION 8.5

A filtered $A_{\infty}$-algebra over $\Lambda(L)$ is a sequence of $\Lambda(L)$ module homomorphisms

$$
\mathfrak{m}_{k}: B_{k}(C F[1]) \rightarrow C F[1], \quad k=0,1,2, \ldots
$$

of degree +1 such that the coderivation $d=\sum_{k=0}^{\infty} \widehat{\mathfrak{m}}_{k}$ satisfies $d d=0$, which is called the $A_{\infty}$-relation. Here we denote by $\widehat{\mathfrak{m}}_{k}: B(C F[1]) \rightarrow B(C F[1])$ the unique extension of $\mathfrak{m}_{k}$ as a coderivation on $B(C F[1])$. A filtered $A_{\infty}$-algebra is an $A_{\infty^{-}}$ algebra with a filtration for which $\mathfrak{m}_{k}$ are continuous with respect to the induced non-Archimedean topology.

If we have $\mathfrak{m}_{1} \mathfrak{m}_{1}=0$, it defines a complex $\left(C F, \mathfrak{m}_{1}\right)$. We define the $\mathfrak{m}_{1}$-cohomology by

$$
\begin{equation*}
H\left(C F, \mathfrak{m}_{1}\right)=\operatorname{Ker} \mathfrak{m}_{1} / \operatorname{Im} \mathfrak{m}_{1} . \tag{8.5}
\end{equation*}
$$

The first two terms of the $A_{\infty}$-relation for a $A_{\infty}$-algebra are given as

$$
\begin{align*}
\mathfrak{m}_{1}\left(\mathfrak{m}_{0}(1)\right) & =0,  \tag{8.6}\\
\mathfrak{m}_{1} \mathfrak{m}_{1}(x)+(-1)^{|x|^{\prime}} \mathfrak{m}_{2}\left(x, \mathfrak{m}_{0}(1)\right)+\mathfrak{m}_{2}\left(\mathfrak{m}_{0}(1), x\right) & =0 \tag{8.7}
\end{align*}
$$

In particular, for the case when $\mathfrak{m}_{0}(1)$ is nonzero, $\mathfrak{m}_{1}$ does not necessarily satisfy the boundary property; that is, $\mathfrak{m}_{1} \mathfrak{m}_{1} \neq 0$ in general.

We now describe the $A_{\infty}$-operators $\mathfrak{m}_{k}$ in the context of $A_{\infty}$-algebra of Lagrangian submanifolds. For a given compatible almost complex structure $J$, consider the moduli space of stable maps of genus zero,

$$
\mathcal{M}_{k+1}(\beta ; L)=\left\{\left(w,\left(z_{0}, z_{1}, \ldots, z_{k}\right)\right) \mid \bar{\partial} w=0, z_{i} \in \partial D^{2},[w]=\beta \text { in } \pi_{2}(M, L)\right\} / \sim,
$$

where $\sim$ is the conformal reparameterization of the disc $D^{2}$. We require that $z_{0}, \ldots, z_{k}$ respect the counterclockwise cyclic order of $S^{1}$. (We wrote this moduli space $\mathcal{M}_{k+1}^{\text {main }}(\beta ; L)$ in $[F O+1$, Section 2.1]. The symbol "main" indicates the compatibility of $z_{0}, \ldots, z_{k}$ with counterclockwise cyclic order. We omit the symbol in this paper since we always assume it.)
$\mathcal{M}_{k+1}(\beta ; L)$ has a Kuranishi structure, and its dimension is given by

$$
\begin{equation*}
n+\mu(\beta)-3+(k+1)=n+\mu(\beta)+k-2 . \tag{8.8}
\end{equation*}
$$

Now let $\left[P_{1}, f_{1}\right], \ldots,\left[P_{k}, f_{k}\right] \in C_{*}(L ; \mathbb{Q})$ be $k$-smooth singular simplices of $L$. (Here we denote by $C(L ; \mathbb{Q})$ a suitably chosen countably generated cochain complex of smooth singular chains of $L$.) We put the cohomological grading $\operatorname{deg} P_{i}=$ $n-\operatorname{dim} P_{i}$ and consider the fiber product

$$
\mathrm{ev}_{0}: \mathcal{M}_{k+1}(\beta ; L) \times_{\left(\mathrm{ev}_{1}, \ldots, \mathrm{ev}_{k}\right)}\left(P_{1} \times \cdots \times P_{k}\right) \rightarrow L
$$

A simple calculation shows that the expected dimension of this chain is given by $n+\mu(\beta)-2+\sum_{j=1}^{k}\left(\operatorname{dim} P_{j}+1-n\right)$, or equivalently, we have the degree

$$
\operatorname{deg}\left[\mathcal{M}_{k+1}(\beta ; L) \times{ }_{\left(\mathrm{ev}_{1}, \ldots, \mathrm{ev}_{k}\right)}\left(P_{1} \times \cdots \times P_{k}\right), \mathrm{ev}_{0}\right]=\sum_{j=1}^{n}\left(\operatorname{deg} P_{j}-1\right)+2-\mu(\beta)
$$

For each given $\beta \in \pi_{2}(M, L)$ and $k=0, \ldots$, we define $\mathfrak{m}_{1,0}(P)= \pm \partial P$ and

$$
\begin{align*}
\mathfrak{m}_{k, \beta}\left(P_{1}, \ldots, P_{k}\right) & =\left[\mathcal{M}_{k+1}(\beta ; L) \times \times_{\left(\mathrm{ev}_{1}, \ldots, \mathrm{ev}_{k}\right)}\left(P_{1} \times \cdots \times P_{k}\right), \mathrm{ev}_{0}\right]  \tag{8.9}\\
& \in C(L ; \mathbb{Q}) .
\end{align*}
$$

(More precisely, we regard the right-hand side of (8.9) as a smooth singular chain by taking the appropriate multivalued perturbation (multisection) and choosing a simplicial decomposition of its zero set.)

We put

$$
C F(L)=C(L ; \mathbb{Q}) \widehat{\otimes}_{\mathbb{Q}} \Lambda(L) .
$$

We define $\mathfrak{m}_{k}: B_{k} C F(L)[1] \rightarrow B_{k} C F[1]$ by

$$
\mathfrak{m}_{k}=\sum_{\beta \in \pi_{2}(M, L)} \mathfrak{m}_{k, \beta} \otimes[\beta] .
$$

Then it follows that the map $\mathfrak{m}_{k}: B_{k} C F(L)[1] \rightarrow C F(L)[1]$ is well defined, has degree 1 , and is continuous with respect to non-Archimedean topology. We extend $\mathfrak{m}_{k}$ as a coderivation $\widehat{\mathfrak{m}}_{k}: B C F[1] \rightarrow B C F[1]$, where $B C F(L)[1]$ is the completion of the direct sum $\bigoplus_{k=0}^{\infty} B_{k} C F(L)[1]$, where $B_{k} C F(L)[1]$ itself is the
completion of $C F(L)[1]^{\otimes k}$. $B C F(L)[1]$ has a natural filtration defined similarly as in Definition 8.3. Finally, we take the sum

$$
\widehat{d}=\sum_{k=0}^{\infty} \widehat{\mathfrak{m}}_{k}: B C F(L)[1] \rightarrow B C F(L)[1] .
$$

We then have the following coboundary property.

## THEOREM 8.2

Let $L$ be an arbitrary compact relatively spin Lagrangian submanifold of an arbitrary tame symplectic manifold $(M, \omega)$. The coderivation $\widehat{d}$ is a continuous map that satisfies the $A_{\infty}$-relation $\widehat{d d}=0$, and so $(C F(L), \mathfrak{m})$ is a filtered $A_{\infty}$-algebra over $\Lambda(L)$.

We put

$$
C\left(L ; \Lambda_{0, \mathrm{nov}}\right)=C F(L) \widehat{\otimes}_{\Lambda(L)} \Lambda_{0, \mathrm{nov}}
$$

on which a filtered $A_{\infty}$-structure on $C\left(L ; \Lambda_{0, \text { nov }}\right)$ (over the ring $\left.\Lambda_{0, \text { nov }}\right)$ is induced. This is the filtered $A_{\infty}$-structure given in [FO+1, Theorem A].

In the presence of $\mathfrak{m}_{0}, \widehat{\mathfrak{m}}_{1} \widehat{\mathfrak{m}}_{1}=0$ no longer holds in general. This leads us to consider deforming Floer's original definition by a bounding cochain of the obstruction cycle arising from bubbling-off discs. One can always deform the given (filtered) $A_{\infty}$-algebra $(C F(L), \mathfrak{m})$ by an element $b \in C F(L)[1]^{0}$ by redefining the $A_{\infty}$-operators as

$$
\mathfrak{m}_{k}^{b}\left(x_{1}, \ldots, x_{k}\right)=\mathfrak{m}\left(e^{b}, x_{1}, e^{b}, x_{2}, e^{b}, x_{3}, \ldots, x_{k}, e^{b}\right)
$$

and taking the sum $\widehat{d}^{b}=\sum_{k=0}^{\infty} \widehat{\mathfrak{m}}_{k}^{b}$. This defines a new filtered $A_{\infty}$-algebra in general. Here we simplify notation by writing

$$
e^{b}=1+b+b \otimes b+\cdots+b \otimes \cdots \otimes b+\cdots
$$

Note that each summand in this infinite sum has degree zero in $C F(L)[1]$ and converges in the non-Archimedean topology if $b$ has positive valuation, that is, $\mathfrak{v}(b)>0($ see Section 3.2 for the definition of $\mathfrak{v})$.

## PROPOSITION 8.3

For the $A_{\infty}$-algebra $\left(C F(L), \mathfrak{m}_{k}^{b}\right), \mathfrak{m}_{0}^{b}=0$ if and only if $b$ satisfies

$$
\begin{equation*}
\sum_{k=0}^{\infty} \mathfrak{m}_{k}(b, \ldots, b)=0 \tag{8.10}
\end{equation*}
$$

This equation is a version of the Maurer-Cartan equation for the filtered $A_{\infty}$-algebra.

## DEFINITION 8.6

Let $(C F(L), \mathfrak{m})$ be a filtered $A_{\infty}$-algebra in general, and let $B C F(L)[1]$ be its bar complex. An element $b \in C F(L)[1]^{0}=C F(L)^{1}$ is called a bounding cochain if
it satisfies (8.10) and $v(b)>\mathfrak{v}(b)$. We denote by $\widetilde{\mathcal{M}}(L ; \Lambda(L))$ the set of bounding cochains.

In general, a given $A_{\infty}$-algebra may or may not have a solution to (8.10). In our case we define the following.

## DEFINITION 8.7

A filtered $A_{\infty}$-algebra $(C F(L), \mathfrak{m})$ is called unobstructed over $\Lambda(L)$ if (8.10) has a solution $b \in C F(L)[1]^{0}=C F(L)^{1}$ with $v(b)>\mathfrak{v}(b)$.

One can define the notion of homotopy equivalence between two bounding cochains as described in $[\mathrm{FO}+1$, Chapter 4]. We denote by $\mathcal{M}(L ; \Lambda(L))$ the set of equivalence classes of bounding cochains of $L$.

REMARK 8.8
In Definition 8.6, we consider the bounding cochain contained in $C F(L) \subset C(L$; $\Lambda_{0}$ ) only. This is the reason why we write $\widetilde{\mathcal{M}}(L ; \Lambda(L))$ in place of $\widetilde{\mathcal{M}}(L)$. (The latter is used in $[\mathrm{FO}+1]$.)

## 8.3. $A_{\infty}$-bimodule

Once the $A_{\infty}$-algebra is attached to each Lagrangian submanifold $L$, we then construct a structure of filtered $A_{\infty}$-bimodules on the module $\operatorname{CF}\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)=$ $C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)$, which was introduced in Section 8.1 as follows. This filtered $A_{\infty}$-bimodule structure is by definition a family of operators

$$
\begin{aligned}
\mathfrak{n}_{k_{1}, k_{0}}: B_{k_{1}}\left(C F\left(L_{1}\right)[1]\right) \hat{\otimes}_{\Lambda\left(L_{1}\right)} C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right) & \widehat{\otimes}_{\Lambda\left(L_{0}\right)} B_{k_{0}}\left(C F\left(L^{\prime}\right)[1]\right) \\
& \rightarrow C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)
\end{aligned}
$$

for $k_{0}, k_{1} \geq 0$. Here the left-hand side is defined as follows. It is easy to see that there are embeddings $\Lambda\left(L_{0}\right) \rightarrow \Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right), \Lambda\left(L_{1}\right) \rightarrow \Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$. Therefore a $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$-module $C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)$ can be regarded both as a $\Lambda\left(L_{0}\right)$-module and a $\Lambda\left(L_{1}\right)$-module. Hence we can take tensor product in the left-hand side ( $\widehat{\bigotimes}_{\Lambda\left(L_{i}\right)}$ is the completion of this algebraic tensor product). The left-hand side then becomes a $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$-module since the rings involved are all commutative.

We briefly describe the definition of $\mathfrak{n}_{k_{1}, k_{0}}$. A typical element of the tensor product

$$
B_{k_{1}}\left(C F\left(L_{1}\right)[1]\right) \hat{\otimes}_{\Lambda\left(L_{1}\right)} C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right) \hat{\otimes}_{\Lambda\left(L_{0}\right)} B_{k_{0}}\left(C F\left(L_{0}\right)[1]\right)
$$

has the form

$$
P_{1,1} \otimes \cdots \otimes P_{1, k_{1}} \otimes[p, w] \otimes P_{0,1} \otimes \cdots \otimes P_{0, k_{0}}
$$

with $p \in L_{0} \cap L_{1}$ being an admissible intersection point. Then the image $\mathfrak{n}_{k_{0}, k_{1}}$ thereof is given by

$$
\sum_{q, B} T^{\omega(B)} e^{\mu(B) / 2} \#\left(\mathcal{M}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,0}, \ldots, P_{0, k_{0}}\right)\right)[q, B \# w] .
$$

Here $B$ denotes homotopy class of Floer trajectories connecting $p$ and $q$, the summation is taken over all $[q, B]$ with

$$
\operatorname{dim} \mathcal{M}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)=0
$$

and $\#\left(\mathcal{M}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)\right)$ is the "number" of elements in the "zero"-dimensional moduli space $\mathcal{M}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)$. Here the moduli space $\mathcal{M}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)$ is the Floer moduli space $\mathcal{M}(p, q ; B)$ cut down by intersecting with the given chains $P_{1, i} \subset L_{1}$ and $P_{0, j} \subset L_{0}($ see $[F O+1$, Section 3.7]). An orientation on this moduli space can be given in $[\mathrm{FO}+1]$ and $[\mathrm{FO}+2]$.

## THEOREM 8.4

Let $\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ be a pair of anchored Lagrangian submanifolds. Then the family $\left\{\mathfrak{n}_{k_{1}, k_{0}}\right\}$ defines a left $\left(C F\left(L_{1}\right), \mathfrak{m}\right)$ and right $\left(C F\left(L_{0}\right), \mathfrak{m}\right)$ filtered $A_{\infty}$-bimodule structure on $C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$.

See $\left[\mathrm{FO}+1\right.$, Section 3.7] and $[\mathrm{FO}+2]$ for the definition of filtered $A_{\infty}$-bimodules. (In $[\mathrm{FO}+1]$ the case of the universal Novikov ring as a coefficient is considered. It is easy to modify it to our case of a $\Lambda\left(L_{0}, L_{1}\right)$-coefficient.)

In the case where both $L_{0}, L_{1}$ are unobstructed, we can perform this deformation of $\mathfrak{n}$ using bounding cochains $b_{0}$ and $b_{1}$ of $C F\left(L_{0}\right)$ and $C F\left(L_{1}\right)$, respectively, in a way similar to $\mathfrak{m}^{b}$. Namely, we define $\delta_{b_{1}, b_{0}}: C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ by

$$
\delta_{b_{1}, b_{0}}(x)=\sum_{k_{1}, k_{0}} \mathfrak{n}_{k_{1}, k_{0}}\left(b_{1}^{\otimes k_{1}} \otimes x \otimes b_{0}^{\otimes k_{0}}\right)=\widehat{\mathfrak{n}}\left(e^{b_{1}}, x, e^{b_{0}}\right) .
$$

We can generalize the story to the case where $L_{0}$ has clean intersection with $L_{1}$, especially to the case $L_{0}=L_{1}$. In the case $L_{0}=L_{1}$ we have $\mathfrak{n}_{k_{1}, k_{0}}=\mathfrak{m}_{k_{0}+k_{1}+1}$. So in this case, we have $\delta_{b_{1}, b_{0}}(x)=\mathfrak{m}\left(e^{b_{1}}, x, e^{b_{0}}\right)$.

We define the Floer cohomology of the pair $\mathcal{L}_{0}=\left(L_{0}, \gamma_{0}, \lambda_{0}\right), \mathcal{L}_{1}=\left(L_{1}, \gamma_{1}, \lambda_{1}\right)$ by

$$
H F\left(\left(\mathcal{L}_{1}, b_{1}\right),\left(\mathcal{L}_{0}, b_{0}\right)\right)=\operatorname{Ker} \delta_{b_{1}, b_{0}} / \operatorname{Im} \delta_{b_{1}, b_{0}} .
$$

This is a module over $\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$.

## THEOREM 8.5

$\operatorname{HF}\left(\left(\mathcal{L}_{1}, b_{1}\right),\left(\mathcal{L}_{0}, b_{0}\right)\right) \otimes_{\Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)} \Lambda_{\text {nov }}$ is invariant under the Hamiltonian isotopies of $\mathcal{L}_{0}$ and $\mathcal{L}_{1}$ and under the gauge equivalence of bounding cochains $b_{0}, b_{1}$.

We refer to $[\mathrm{FO}+1$, Section 4.3] for the definition of gauge equivalence and to [FO+1, Theorem 4.1.5] for the proof of this theorem.

## 9. Definitions of Seidel's maps $\mathfrak{b}, \mathfrak{c}$, and $\mathfrak{h}$

In this section, we recall the definition of Seidel's cochain maps $\mathfrak{b}, \mathfrak{c}$, and the homotopy $\mathfrak{h}$ and give the definition of the analogues thereof in our general setting.

They are the maps

$$
\begin{aligned}
\mathfrak{b}: C F\left(L, L_{1}\right) & \otimes C F\left(\tau_{L}\left(L_{0}\right), L\right) \\
\mathfrak{c}: C F\left(\tau_{L}\left(L_{0}\right), L_{1}\right) & \rightarrow C F\left(\tau_{L}\left(L_{0}\right), L_{1}\right),
\end{aligned}
$$

and the homotopy $\mathfrak{h}: C F\left(L, L_{1}\right) \otimes C F\left(\tau_{L}\left(L_{0}\right), L\right) \rightarrow C F\left(L_{0}, L_{1}\right)$ between the composition $\mathfrak{c} \circ \mathfrak{b}$ and the zero map.

We generalize these maps to our nonexact case and describe all the necessary properties of the maps in this section. We consider a quadruple of anchored Lagrangian submanifolds

$$
\begin{align*}
\mathcal{L} & =(L, \gamma), \quad \mathcal{L}_{0}=\left(L_{0}, \gamma_{0}\right), \quad \mathcal{L}_{1}=\left(L_{1}, \gamma_{1}\right), \\
\tau_{*} \mathcal{L} & =\left(\tau_{L}\left(L_{0}\right), \tau_{L}\left(\gamma_{0}\right)\right) . \tag{9.1}
\end{align*}
$$

For simplicity of notation in this section, we denote the action functional associated to the pair $(L, \gamma)$ and $\left(L^{\prime}, \gamma^{\prime}\right)$ just by $\mathcal{A}_{\mathcal{L} \mathcal{L}^{\prime}}$.

### 9.1. A simple invariant and its vanishing theorem

Let $\Phi_{1}(E, \pi, Q)$ be the invariant

$$
\Phi_{1}(E, \pi, Q)=\left(\mathrm{ev}_{\zeta}\right)_{*}\left[\mathcal{M}_{J}\right] \in H_{*}\left(Q_{\zeta} ; \mathbb{Z}\right)
$$

for the Calabi-Yau Lefschetz fibration. The following proposition replaces a similar proposition, [Se3, Proposition 2.13], for the present Calabi-Yau Lefschetz fibration setting.

For each given section class $A \in \pi_{2}^{\text {sec }}(E, Q)$, we define the moduli space $\mathcal{M}(E, Q, J ; A)$ of $J$-holomorphic section $s: D \rightarrow E$ with $[s, \partial s]=A$ and define an invariant

$$
\Phi_{1}(E, \pi, Q)=\sum_{A \in \pi_{2}^{\mathrm{sec}}(E, Q)}\left(\mathrm{ev}_{\zeta}\right)_{*}\left(\mathcal{M}_{J}(A)\right) T^{\Omega(A)} \in H_{*}\left(Q_{\zeta} ; \Lambda_{0, \mathrm{nov}}\right)
$$

We start with the following slight generalization of Proposition 2.2 in the general context of Lagrangian spheres in general symplectic manifolds.

PROPOSITION 9.1
Let $(L,[f])$ be a framed Lagrangian sphere in $M$. There is a 1-parameter family of Lefschetz Hamiltonian fibrations $\left(E_{r}^{L}, \pi_{r}^{L}\right) \rightarrow D(r)$ together with an isomorphism $\phi_{r}^{L}: E_{r}^{L} \rightarrow M$ of symplectic manifolds, such that we have the following.
(1) Consider the rescaling map $\lambda_{r}: D(r) \rightarrow D(1)$ defined by $z \mapsto z / r$. Then

$$
\left(\lambda_{r}\right)^{*} E_{1}^{L}=E_{r}^{L}
$$

(2) If $\rho_{r}^{L}$ is the symplectic monodromy around $\partial \bar{D}(r)$, then $\phi_{r}^{L} \circ \rho^{L} \circ\left(\phi_{r}^{L}\right)^{-1}$ is a Dehn twist along ( $L,[f]$ ).

We denote any of these maps by $\tau_{L}$, as before.

## Proof

The proof follows from Seidel's proof of [Se3, Proposition 1.11] stripping all the things related to the exactness requirement in the proof. In fact, the proof is easier because we do not have to concern ourselves with the exactness requirement in the construction.

We also recall that each fiber $E_{z}^{L}, z \neq 0$, of the fibration $E^{L}$ contains a distinguished Lagrangian sphere $\Sigma_{z}^{L}$. We call this fibration a standard Calabi-Yau Lefschetz fibration. The following is a crucial proposition needed in Seidel's construction of the long exact sequence in [Se3], whose proof goes through in our current context with orientation consideration incorporated.

## PROPOSITION 9.2

Let $L \subset M$ be any CY Lagrangian brane, and let $\pi^{L}: E^{L} \rightarrow D(r)$ be the associated standard Calabi-Yau Lefschetz fibration. Then we have

$$
\Phi_{1}\left(E^{L}, \pi^{L}, Q^{L}\right)=0 .
$$

## Proof

The proof of this proposition follows verbatim that of [Se3, Proposition 2.13] except for the consideration of orientations. The relevant consideration of orientation of the moduli space $\mathcal{M}(E, Q, J ; A)$ is given in [Se4, Section 17i] and so is omitted.

### 9.2. The map $\mathfrak{b}$

Let $\Sigma$ be a compact surface with boundary marked points $\vec{\zeta}=\left\{\zeta_{0}, \zeta_{1}, \zeta_{2}\right\}$. We denote $\dot{\Sigma}=\Sigma \backslash \vec{\zeta}$ and $\partial \dot{\Sigma}=\bigcup_{i=0}^{2} \partial_{i} \dot{\Sigma}$. We consider the three anchored Lagrangian submanifolds $\left(L_{0}, \gamma_{0}\right),(L, \gamma)$, and ( $L_{1}, \gamma_{1}$ ). Take the trivial Hamiltonian fibration $\pi: E=\dot{\Sigma} \times M \rightarrow \dot{\Sigma}$ with the 2 -form $\Omega$ equal to the 2 -form pulled back from $\omega$ in $M$. Equip this with the Lagrangian boundary condition

$$
Q=\left(\partial_{0} \dot{\Sigma} \times \tau_{L}\left(L_{0}\right)\right) \cup\left(\partial_{1} \dot{\Sigma} \times L\right) \cup\left(\partial_{3} \dot{\Sigma} \times L_{1}\right)
$$

$Q$ is an exact Lagrangian boundary with $\kappa_{Q}=0$. We note that $E$ has the trivial connection given by $K \equiv 0$ and hence has zero curvature.

LEMMA 9.3
Suppose that $\mathcal{L}=(L, \gamma), \mathcal{L}_{0}=\left(L_{0}, \gamma_{0}\right), \mathcal{L}_{1}=\left(L_{1}, \gamma_{1}\right)$ are given anchors of the type (9.1), and suppose that $\tau_{*} \mathcal{L}=\left(\tau_{L}\left(L_{0}\right), \tau_{L}\left(\gamma_{0}\right)\right)$. Let $s: \dot{\Sigma} \rightarrow \dot{\Sigma} \times M$ be a section with the given exact Lagrangian boundary condition $Q$ as above. Let $\left[p_{0}, w_{0}\right] \in$ $\operatorname{Crit} \mathcal{A}_{\tau_{L}\left(\mathcal{L}_{0}\right) \mathcal{L}_{1}},[p, w] \in \operatorname{Crit} \mathcal{A}_{\mathcal{L} \mathcal{L}_{1}}$, and $\left[p_{1}, w_{1}\right] \in \operatorname{Crit} \mathcal{A}_{\tau_{L}\left(\mathcal{L}_{0}\right) \mathcal{L}}$. Suppose that the homotopy class $[s, \partial s]$ is admissible and satisfies

$$
\begin{equation*}
\left[w_{0}\right]=[s, \partial s] *[w] *\left[w_{1}\right] \quad \text { in } \pi_{2}\left(\tau_{L}\left(L_{0}\right), L_{1} ; \tau_{L}\left(\bar{\gamma}_{0}\right) * \gamma_{1}\right) . \tag{9.2}
\end{equation*}
$$

Then we have

$$
\begin{equation*}
\int s^{*} \Omega=\mathcal{A}_{\tau_{L}\left(\mathcal{L}_{0}\right) \mathcal{L}_{1}}\left(\left[p_{0}, w_{0}\right]\right)-\mathcal{A}_{\mathcal{L} \mathcal{L}_{1}}([p, w])-\mathcal{A}_{\tau_{L}\left(\mathcal{L}_{0}\right) \mathcal{L}}\left(\left[p_{1}, w_{1}\right]\right) . \tag{9.3}
\end{equation*}
$$

Proof
This immediately follows from Proposition 3.4.
DEFINITION 9.1
Let $L, L^{\prime} \subset M$ be a pair of Lagrangian submanifolds. We say that $J$ lies in $\mathcal{J}^{\mathrm{reg}}\left(M ; L, L^{\prime}\right)$ if all Floer trajectories associated to $J$ for the pair $\left(L, L^{\prime}\right)$ are Fredholm regular.

At this point, we fix

$$
\begin{aligned}
J^{(1)} & \in \mathcal{J}^{\mathrm{reg}}\left(M ; \tau_{L}\left(L_{0}\right), L\right), \\
J^{(2)} & \in \mathcal{J}^{\mathrm{reg}}\left(M ; L, L_{1}\right), \\
J^{(3)} & \in \mathcal{J}^{\mathrm{reg}}\left(M ; \tau_{L}\left(L_{0}\right), L_{1}\right) .
\end{aligned}
$$

By choosing a horizontal $J \in \mathcal{J}\left(E, \pi, j, J^{(1)}, J^{(2)}, J^{(3)}\right)$, that is, $J$ satisfying $J_{x}\left(T E_{x}^{h}\right)=T E_{x}^{h}$ for all $x \in E \backslash E^{\text {crit }}$ with the above trivial fibration $E=\dot{\Sigma} \times M$, this gives rise to the standard pants product map

$$
\begin{equation*}
C F\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \rightarrow C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right) \tag{9.4}
\end{equation*}
$$

in the cochain level. Seidel's map $\mathfrak{b}$ in the cochain level is nothing but this pants product map.

By choosing bounding cochains $b_{0}, b_{1}$ of $L_{0}, L_{1}$, respectively, and considering $\left(\tau_{L}\right)_{*} b_{0}$ on $\tau_{L}\left(L_{0}\right)$, we define the deformed $\mathfrak{m}_{2}$ :

$$
\mathfrak{m}_{2}^{\vec{b}}: C F\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \rightarrow C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right),
$$

where $\vec{b}=\left(0,\left(\tau_{L}\right)_{*} b_{0}, b_{1}\right)$ and $\mathfrak{m}_{i}^{\vec{b}}$ is defined as in $[\mathrm{FO}+2,(8.15)]$. Here we use $\mathfrak{m}_{2}^{\vec{b}}$ instead of $\mathfrak{b}$ to be consistent with the notation from $[\mathrm{FO}+1]$, $[\mathrm{Fu}]$, and $[\mathrm{FO}+2]$ and to highlight the point that Seidel's map $\mathfrak{b}$ is nothing but the special case of the $\mathfrak{m}_{2}$-map therefrom.

According to $[\mathrm{FO}+1],[\mathrm{Fu}]$, and $[\mathrm{FO}+2]$, this induces a cochain map up to the higher homotopy map $\mathfrak{m}_{3}^{\vec{b}}$ and induces a homomorphism in cohomology. For the reader's convenience, we summarize the construction of this product map $\mathfrak{m}_{k}^{\vec{b}}$ in the appendix.

### 9.3. The map c

Let $\left(E^{L}, \pi^{L}\right)$ be the standard fibration over a disc $D(1 / 2)$ whose monodromy around $\partial D(1 / 2)$ is $\tau_{L}$, as defined in Section 2, by choosing $r$ small. Let $0<r<1 / 2$ be given, and choose a function $g$ with $g(t)=t$ for small $t, g(t) \equiv r$ for $t \geq r$, and $g^{\prime}(t) \geq 0$ everywhere. We consider the map $p: D(1 / 2) \rightarrow D(1 / 2)$ defined by
$p(z)=g(|z|) z /|z|$, and consider the pullback fibration

$$
\left(E^{p}, \pi^{p}\right)=p^{*}\left(E^{L}, \pi^{L}\right)
$$

This is flat on the annulus $D(1 / 2) \backslash \operatorname{Int}(D(r))$.
Now take the surface

$$
\Sigma^{f}=\mathbb{R} \times[-1,1] \backslash \operatorname{Int} D(1 / 2) \subset \mathbb{R}^{2}
$$

with coordinates $(s, t)$, and divide it into two parts $\Sigma^{f, \pm}=\Sigma^{f} \cap\left\{t \in \mathbb{R}_{ \pm}\right\}$so that

$$
\Sigma^{f,+} \cap \Sigma^{f,-}=\left(\left(-\infty,-\frac{1}{2}\right] \cup\left[\frac{1}{2}, \infty\right)\right) \times\{0\} .
$$

Consider trivial fibrations

$$
\pi^{f, \pm}: E^{f, \pm}=\Sigma^{f, \pm} \times M \rightarrow \Sigma^{f, \pm}
$$

over the two parts, and equip them with 2 -forms $\Omega^{f, \pm}$ the pullback of $\omega$. We define a fibration $\left(E^{f}, \pi^{f}\right)$ over $\Sigma^{f}$ by identifying the fibers $E_{(s, 0)}^{f,+} \rightarrow E_{(s, 0)}^{f,-}$ via $\operatorname{id}_{M}$ for $s \geq 1 / 2$ and via $\tau_{L}$ for $s \leq-1 / 2$. Because $\tau_{L}$ is symplectic, this defines a flat Hamiltonian fibration.

Using the fact that two fibrations $E^{p}$ and $E^{f}$ are flat close to the curve $|z|=1 / 2$, we now paste them along the curve. Denote the resulting fibration over

$$
\Sigma^{c}:=\Sigma^{p} \cup \Sigma^{f}=\mathbb{R} \times[-1,1]
$$

by $\left(E^{c}, \pi\right)$. Equip $\left(E^{c}, \pi\right)$ with the Lagrangian boundary condition

$$
Q^{c}=\left\{\begin{array}{l}
\mathbb{R} \times\{1\} \times L_{1} \subset E^{f,+}  \tag{9.5}\\
\mathbb{R} \times\{-1\} \times \tau_{L}\left(L_{0}\right) \subset E^{f,-}
\end{array}\right.
$$

This defines an exact Lagrangian boundary with the 1-form $\kappa_{Q^{c}}=0$. As explained in $\left[\operatorname{Se} 3\right.$, Section 3.3], $\left(E^{c}, \pi\right)$ is modeled on $\left(\tau_{L}\left(L_{0}\right), L_{1}\right)$ over the positive end of $\Sigma$, and over the negative end, it is modeled on $\left(L_{0}, L_{1}\right)$ due to the monodromy effect around the critical value $(0,0) \in \Sigma$.

Let $j$ be some complex structure on $\Sigma$, standard over the ends. Take some $J^{(3)} \in \mathcal{J}^{\mathrm{reg}}\left(M, \tau_{L}\left(L_{0}\right), L_{1}\right)$ as in Section 9.2, and choose an additional $J^{(5)} \in$ $\mathcal{J}^{\text {reg }}\left(M ; L_{0}, L_{1}\right)$. Using a regular $J^{(6)} \in \mathcal{J}\left(E^{c}, \pi, Q^{c}, j, J^{(3)}, J^{(5)}\right)$, we define a map

$$
\mathfrak{c}=C \Phi_{0}^{\mathrm{rel}}\left(E^{c}, \pi, Q^{c}, J^{(6)}\right): C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right) \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)
$$

as defined in [Se3, Section 3.3]. However, since we need to establish this morphism in $\mathbb{Z}$-coefficients, we need to incorporate the study of orientation in its definition. Furthermore, we also need to involve the bounding cochains since we need to deform Seidel's definition just as we do for the Floer coboundary map.

### 9.3.1. Seidel map $\mathfrak{c}$ over $\mathbb{Z}$

For each given $p \in L_{1} \cap \tau_{L}\left(L_{0}\right)$ and $q \in L_{1} \cap L_{0}$ and a section class $B \in \pi_{2}\left(E^{c}, Q^{c}\right.$; $p, q)$, we consider the moduli space $\mathcal{M}^{c}(p, q ; B)$ consisting of the $J^{(6)}$-holomorphic
sections $s: \Sigma^{c} \rightarrow E^{c}$ satisfying the conditions

$$
[s]=B, \quad s(-\infty)=\widetilde{p}, s(\infty)=\widetilde{q} .
$$

Here $\widetilde{p}, \widetilde{q}$ are horizontal sections of $\left(\left.E^{c}\right|_{ \pm \infty}, \partial E_{ \pm \infty}^{c}\right)$ which are trivialized to

$$
\left(\left.E^{c}\right|_{-\infty} ;\left.\partial E^{c}\right|_{-\infty}\right) \cong\left([-1,1] \times M ;\left\{\{-1\} \times \tau_{L}\left(L_{0}\right)\right\} \cup\left\{\{1\} \times L_{1}\right\}\right)
$$

and

$$
\left(\left.E^{c}\right|_{+\infty} ;\left.\partial E^{c}\right|_{+\infty}\right) \cong\left([-1,1] \times M ;\left\{\{-1\} \times L_{1}\right\} \cup\left\{\{1\} \times L_{0}\right\}\right)
$$

over $([0,1],\{0,1\})$ from the definition of the fibration $E^{c} \rightarrow \Sigma^{c}$ above.
We need to equip $\mathcal{M}^{c}(p, q ; B)$ with an orientation denoted by $o^{c}(p, q ; B)$ that is compatible with the orientations $o_{[p, w]}$ provided at $[p, w]$ defined in Theorem 7.12. This amounts to saying that we establish the existence of a relative spin structure on $(E, \mathcal{Q}) \rightarrow(\Sigma, \partial \Sigma)$ such that $\Sigma$ has one incoming end and one outgoing end of the type (9.5) and satisfies the following requirement.

Let $Z_{ \pm}$be the space given in (7.29), and consider the bundle pairs

$$
E_{ \pm}=Z_{ \pm} \times \mathbb{C}^{n}, \quad \lambda_{ \pm}
$$

defined right above (7.29). Then we consider the gluing bundle pair

$$
\left(E_{-}, \lambda_{-}\right) \#\left(s^{*} T^{v} E,(\partial s)^{*} T^{v} Q\right) \#\left(E_{+}, \lambda_{+}\right) .
$$

This is a bundle pair over a disc. By construction, this bundle pair is homotopic to the model bundle pair around the critical value $(0,0)$ induced by the model fibration $\left(E_{r}^{L}, \pi_{r}^{L}\right)$ with model Lagrangian boundary condition $Q_{r}^{L} \subset E_{r}^{L}$ given in Proposition 9.1.

We denote by $o(\cdot)$ the orientation given in the relevant real determinant bundle $(\cdot)$. Then the requirement for the choice of

$$
o^{c}(p, q ; B):=o\left(\operatorname{det}\left(s^{*} T^{v} E,(\partial s)^{*} T^{v} Q\right)\right)
$$

is the gluing rule

$$
\begin{align*}
& o\left(\operatorname{det} D \bar{\partial}_{\left(E_{r}^{L}, \pi_{r}^{L}\right)}\right)  \tag{9.6}\\
& \quad=o\left(\operatorname{det} D \bar{\partial}_{\left(E_{-}, \lambda_{-}\right)}\right) \# o^{c}(p, q ; B) \# o\left(\operatorname{det} D \bar{\partial}_{\left(E_{+}, \lambda_{+}\right)}\right) .
\end{align*}
$$

This coherent choice of orientations is already discussed in [Se4, Section 17] in a more general context of relative spin structure and so is omitted by referring readers thereto. The outcome then gives rise to the Seidel map

$$
\mathfrak{c}=C \Phi_{0}^{\mathrm{rel}}\left(E^{c}, \pi, Q^{c}, J^{(6)}\right): C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right) \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)
$$

over $\mathbb{Z}$ coefficients and so over arbitrary coefficients. However, this map may not satisfy the chain property for the same reason that the Floer map $\delta$ might not satisfy the coboundary property, which leads us to consider the deformed version of $\mathfrak{c}$ whose explanation is now in order.

### 9.3.2. Deformed Seidel map c

Since we use the deformed Floer complex by bounding cochains, we need to construct the deformed version of the map $\mathfrak{c}$. The construction will resemble that of an $A_{\infty}$-bimodule structure on $C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ performed in Section 8.3. Some details of construction are in order.

We first define a family of operators

$$
\begin{aligned}
\mathfrak{c}_{k_{1}, k_{0}}: B_{k_{1}}\left(C F\left(L_{1}\right)[1]\right) \widehat{\otimes}_{\Lambda\left(L_{1}\right)} C F\left(\mathcal{L}_{1}, \tau_{L}\left(\mathcal{L}_{0}\right)\right) & \widehat{\otimes}_{\Lambda\left(L_{0}\right)} B_{k_{0}}\left(C F\left(L_{0}\right)[1]\right) \\
& \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)
\end{aligned}
$$

for $k_{0}, k_{1} \geq 0$. A typical element of the tensor product

$$
B_{k_{1}}\left(C F\left(L_{1}\right)[1]\right) \widehat{\otimes}_{\Lambda\left(L_{1}\right)} C F\left(\mathcal{L}_{1}, \tau_{L}\left(\mathcal{L}_{0}\right)\right) \widehat{\otimes}_{\Lambda\left(L_{0}\right)} B_{k_{0}}\left(C F\left(L_{0}\right)[1]\right)
$$

has the form

$$
P_{1,1} \otimes \cdots \otimes P_{1, k_{1}} \otimes[p, w] \otimes P_{0,1} \otimes \cdots \otimes P_{0, k_{0}}
$$

with $p \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}_{1}$ being an admissible intersection point. Then the image $\mathfrak{c}_{k_{0}, k_{1}}$ thereof is given by

$$
\sum_{q, B} T^{\omega(B)} e^{\mu(B) / 2} \#\left(\mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,0}, \ldots, P_{0, k_{0}}\right)\right)[q, B \# w] .
$$

Here $B$ denotes the section class of $J$-holomorphic sections connecting $p$ and $q$, the summation is taken over all $[q, B]$ with

$$
\operatorname{dim} \mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)=0
$$

and $\#\left(\mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)\right)$ is the "number" of elements in the "zero"-dimensional moduli space $\mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)$. Here the moduli space $\mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)$ is the moduli space $\mathcal{M}^{c}(p, q ; B)$ of $J$-holomorphic sections of $\pi: E^{c} \rightarrow \Sigma$ cut down by intersecting with the given chains $P_{1, i} \subset L_{1}$ and $P_{0, j} \subset L_{0}$.

THEOREM 9.4
Let $\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)=\left(\left(L_{0}, \gamma_{0}\right),\left(L_{1}, \gamma_{1}\right)\right)$ be a pair of anchored Lagrangian submanifolds. Then the family $\left\{\mathfrak{c}_{k_{1}, k_{0}}\right\}$ defines a left $\left(C F\left(L_{1}\right), \mathfrak{m}\right)$ and right $\left(C F\left(L_{0}\right), \mathfrak{m}\right)$ filtered $A_{\infty}$-bimodule homomorphism from $\operatorname{CF}\left(\mathcal{L}_{1}, \tau_{L}\left(\mathcal{L}_{0}\right)\right)$ to $\operatorname{CF}\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$.

The proof of Theorem 9.4 is similar to that of Theorem 8.4 given in the proof of [FO+1, Theorem 3.7.21].

When both $L_{0}, L_{1}$ are unobstructed, we can perform this deformation of $\mathfrak{c}$ using bounding cochains $b_{0}$ and $b_{1}$ of $C F\left(L_{0}\right)$ and $C F\left(L_{1}\right)$, respectively, in a way similar to $\mathfrak{n}^{b_{0}, b_{1}}$. Namely, we define $\left.\mathfrak{c}^{b_{1}, b_{0}}: C F\left(\mathcal{L}_{1}, \tau_{L}\left(\mathcal{L}_{0}\right)\right) \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)\right)$ by

$$
\mathfrak{c}^{b_{1}, b_{0}}(x)=\sum_{k_{1}, k_{0}} \mathfrak{c}_{k_{1}, k_{0}}\left(b_{1}^{\otimes k_{1}} \otimes x \otimes b_{0}^{\otimes k_{0}}\right)=\widehat{\mathfrak{c}}\left(e^{b_{1}}, x, e^{b_{0}}\right)
$$

The following proposition is all we need to add to our context for the construction of Seidel's map $\boldsymbol{c}^{b_{1}, b_{0}}$ in [Se3].

PROPOSITION 9.5
Let $b_{0}, b_{1}$ be bounding cochains of $L_{0}, L_{1}$, respectively. Then $\mathfrak{c}$ defines a chain map

$$
\mathfrak{c}_{1}^{b_{1}, b_{0}}=C \Phi_{0}^{\mathrm{rel}}\left(E, \pi, Q, J^{(6)}\right):\left(C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right), \mathfrak{m}_{1}^{\left(\tau_{L}\right)_{*}\left(b_{0}\right)}\right) \rightarrow\left(C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right), \mathfrak{m}_{1}^{b_{1}}\right)
$$

and hence induces a homomorphism

$$
\mathfrak{c}^{b_{1}, b_{0}}: H F\left(\tau_{L}\left(\mathcal{L}_{0}\right),\left(\tau_{L}\right)_{*} b_{0},\left(\mathcal{L}_{1}, b_{1}\right)\right) \rightarrow H F\left(\left(\mathcal{L}_{1}, b_{1}\right),\left(\mathcal{L}_{0}, b_{0}\right)\right) .
$$

Proof
The proof proceeds in the same way as that of Theorem 8.4. The only difference from the latter is that the moduli space $\mathcal{M}^{c}\left(p, q ; B ; P_{1,1}, \ldots, P_{1, k_{1}} ; P_{0,1}, \ldots, P_{0, k_{0}}\right)$ does not have $\mathbb{R}$-action anymore, and so we consider the moduli space of sections without considering the quotient.

### 9.4. The gluing $\mathfrak{b} \#_{\rho} \mathfrak{c}$, the composition $\mathfrak{c} \circ \mathfrak{b}$, and the homotopy $\mathfrak{h}$

We denote by $\left(E^{b}, \pi^{b}\right), \Sigma^{b}$, and $Q^{b}$ the fibration, surface, and boundary condition associated to the map $b$ and by $\left(E^{c}, \pi^{c}\right), \Sigma^{c}$, and $Q^{c}$ those associated to the map $c$ constructed in Sections 9.1-9.3. We glue the positive end of $E^{c}$ and the negative end of $E^{b}$ to obtain $\left(E_{\rho}^{b c}, \pi_{\rho}^{b c}\right), \Sigma_{\rho}^{b c}$ and $Q_{\rho}^{b c}$, for a sufficiently large gluing parameter $\rho$ in the glued ends. This fibration provides a cochain map

$$
\mathfrak{b} \#_{\rho} \mathfrak{c}: C F\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \rightarrow C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)
$$

From now on we denote

$$
\mathfrak{b}=\mathfrak{m}_{2}^{\vec{b}}, \quad \mathfrak{c}=\mathfrak{c}^{b_{1}, b_{0}}
$$

for notational simplicity.

## LEMMA 9.6

Let $(j, J)$ be such that $j$ is a complex structure on $\Sigma_{\rho}^{b c}$ which is standard on its ends, and let $J \in \mathcal{J}\left(E_{\rho}^{b c}, \pi_{\rho}^{b c}, Q_{\rho}^{b c}, j, J^{(0)}, J^{(2)}, J^{(5)}\right)$. The cochain map $\mathfrak{b} \#_{\rho} \mathfrak{c}$ coincides with $\mathfrak{c} \circ \mathfrak{b}$ for any $\rho \geq \rho_{0}$ with $\rho_{0}$ sufficiently large.

Proof
We note that the composition $\mathfrak{c} \circ \mathfrak{b}$ is defined by counting the elements of the fiber product

$$
\mathcal{M}\left(E^{b}, \mathcal{Q}^{b} ; \vec{p}^{b} ; B^{b}\right)_{\mathrm{ev}_{+}} \#_{\mathrm{ev}_{-}} \mathcal{M}\left(E^{c}, \mathcal{Q}^{c} ; \vec{p}^{c} ; B^{c}\right)
$$

while the map $\mathfrak{b} \#_{\rho} \mathfrak{c}$ is defined by counting those in the moduli space

$$
\mathcal{M}\left(E_{\rho}^{b c}, \mathcal{Q}_{\rho}^{b c} ; \vec{p}_{\rho}^{b c} ; B_{\rho}^{b c}\right)
$$

where $B_{\rho}^{b c}=B^{b} \#{ }_{\rho} B^{c}$ is the obvious glued homotopy class. By a gluing theorem in the Floer complex (see, e.g., $[\mathrm{FO}+1]$ ), the two moduli spaces are diffeomorphic to each other, and hence we have the proof.

Finally, we can construct a homotopy from $\mathfrak{b} \#{ }_{\rho} \mathfrak{c}$ to the zero map following Seidel's argument from [Se3] verbatim (see [Se3, Figures 9, 12] in particular). We omit details of the construction.

## 10. Dichotomy of tiny and big pseudoholomorphic polygons

In this section, we study the decomposition of contributions in the cochain maps $\mathfrak{b}, \mathfrak{c}$, and the homotopy $\mathfrak{h}$ arising from pseudoholomorphic polygons of very small energy and of not small energy. We call the polygons of the first type tiny and the other big.

For this purpose, we first note that the Dehn twist $\tau_{L}^{-1}: M \rightarrow M$ acts by $\left(L_{i}, \gamma_{i}\right) \rightarrow\left(\tau_{L}^{-1}\left(L_{i}\right), \tau_{L}^{-1}\left(\gamma_{i}\right)\right)$ and induces a one-to-one correspondence

$$
\tau_{L}\left(L_{0}\right) \cap L \rightarrow L_{0} \cap L, \quad x \mapsto \tau_{L}^{-1}(x)
$$

Since $\left.\tau_{L}\right|_{L}=\left.\mathrm{id}\right|_{L}$, this lifts to a diffeomorphism

$$
\widetilde{\Omega}\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \rightarrow \widetilde{\Omega}\left(\mathcal{L}_{0}, \mathcal{L}\right), \quad[p, w] \mapsto\left[\tau_{L}^{-1}(p), \tau_{L}^{-1}(w)\right] .
$$

This latter diffeomorphism induces a filtration-preserving isomorphism

$$
\left(\tau_{L}^{-1}\right)_{*}: C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \rightarrow C F\left(\mathcal{L}_{0}, \mathcal{L}\right)
$$

that is, satisfies

$$
\mathcal{A}_{\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}}([p, w])=\mathcal{A}_{\mathcal{L}_{0}, \mathcal{L}}\left(\left[\tau_{L}^{-1}(p), \tau_{L}^{-1}(w)\right]\right) .
$$

By perturbing $L_{0}$ and $L_{1}$ if necessary and choosing $\varepsilon>0$ sufficiently small, we may assume that
(1) $L \cap L_{0}, L \cap L_{1}$, and $L_{0} \cap L_{1}$ are transverse intersections, and $L \cap L_{0} \cap$ $L_{1}=\emptyset ;$
(2) each $\iota^{-1}\left(L_{k}\right) \subset T(r)$ is a union of fibers; one can write this as

$$
\iota^{-1}\left(L_{k}\right)=\bigcup_{y \in \iota^{-1}\left(L \cap L_{k}\right)} T(r)_{y} ;
$$

(3) $R$ satisfies $0 \geq 2 \pi R(0)>-\varepsilon$ and is such that $\tau_{L}$ is $\delta$-wobbly.

The following lemma is a part of [Se3, Lemma 3.2]. For the reader's convenience, we provide its proof.

LEMMA 10.1
Suppose that $L_{0} \cap L \cap L_{1}=\emptyset$. Then we can choose $\operatorname{supp} \tau_{L}$ so close to $L$ that $L_{0} \cap L_{1} \subset M \backslash \operatorname{im} \tau_{L}$ and $\tau_{L}\left(L_{0}\right), L_{1}$ intersect transversally, and there are injective maps

$$
\begin{align*}
& p:\left(\tau_{L}\left(L_{0}\right) \cap L\right) \times\left(L \cap L_{1}\right) \rightarrow \tau_{L}\left(L_{0}\right) \cap L_{1},  \tag{10.1}\\
& q: L_{0} \cap L_{1} \rightarrow \tau_{L}\left(L_{0}\right) \cap L_{1}, \tag{10.2}
\end{align*}
$$

such that $\tau_{L}\left(L_{0}\right) \cap L_{1}$ is the disjoint union of their images.

## Proof

Conditions (1) and (2) right before Lemma 10.1 imply that $L_{0} \cap L_{1} \cap U=\emptyset$. Since $\tau_{L}$ is the identity outside $U$, one has $L_{0} \cap L_{1}=\left(\tau_{L}\left(L_{0}\right) \cap L_{1}\right) \backslash U$, so that $q$ can indeed be defined to be the inclusion. There is a bijective correspondence between pairs $\left(\widetilde{x}_{0}, x_{1}\right) \in\left(\tau_{L}\left(L_{0}\right) \cap L\right) \times\left(L \cap L_{1}\right)$ and $\left(y_{0}, y_{1}\right) \in \iota^{-1}\left(L_{0} \cap L\right) \times \iota^{-1}\left(L \cap L_{1}\right)$, given by setting $y_{0}=\iota^{-1}\left(\tau_{L}^{-1}\left(\widetilde{x}_{0}\right)\right), y_{1}=\iota^{-1}\left(x_{1}\right)$. As a consequence of condition (3) above,

$$
\begin{equation*}
\iota^{-1}\left(\tau_{L}\left(L_{0}\right) \cap L_{1}\right)=\bigcup_{y_{0}, y_{1}} \tau\left(T(\lambda)_{y_{0}}\right) \cap T(\lambda)_{y_{1}} . \tag{10.3}
\end{equation*}
$$

It is clear from their definitions that $p, q$ are injective. A point of $\tau_{L}\left(L_{0}\right) \cap L_{1}$ falls into $\operatorname{im}(q)$ or $\operatorname{im}(p)$ depending on whether it lies inside or outside $\operatorname{im}(\iota)$; hence the two images are disjoint and cover $\tau_{L}\left(L_{0}\right) \cap L_{1}$. The transversality follows from the definition of $\tau_{L}$ for $\operatorname{im}(p)$ and from that of $L_{0} \cap L_{1}$ for $\operatorname{im}(q)$.

We consider the triple

$$
\mathcal{L}=(L, \gamma), \quad \mathcal{L}_{0}=\left(L_{0}, \tau_{L}^{-1} \circ \gamma\right), \quad \mathcal{L}_{1}=\left(L_{1}, \gamma_{1}\right)
$$

We note that $\tau_{L}\left(\mathcal{L}_{0}\right)=\left(\tau_{L}\left(L_{0}\right), \gamma_{0}\right)=\left(\tau_{L}\left(L_{0}\right), \gamma\right)$. To make our discussion nontrivial, we may assume

$$
\begin{equation*}
\mathcal{L} \cap \mathcal{L}_{0} \neq \emptyset, \quad \mathcal{L} \cap \mathcal{L}_{1} \neq \emptyset \tag{10.4}
\end{equation*}
$$

We fix elements $x_{1} \in \mathcal{L} \cap \mathcal{L}_{1}$ and $\widetilde{x}_{0} \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}$ and $\widetilde{z}_{0}=p\left(\widetilde{x}_{0}, x_{1}\right)$, where $p$ is the injective map given in Lemma 10.1.

LEMMA 10.2
Suppose that $x_{1} \in \mathcal{L} \cap \mathcal{L}_{1}$ and $\widetilde{x}_{0} \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}$ and $\widetilde{z}_{0}=p\left(\widetilde{x}_{0}, x_{1}\right)$. Then we have $\widetilde{z}_{0} \in \mathcal{L}_{1} \cap \tau_{L}\left(\mathcal{L}_{0}\right)$.

Proof
We first note that there is a canonical homotopy class $B_{\text {can }}=B\left(\widetilde{x}_{0}, x_{1}, \widetilde{z}\right)$ spanned by a tiny triangle contained in $U=\operatorname{im} \iota$. Choose a path $w_{1}$ from $\bar{\gamma} * \gamma_{1}$ and $w_{0}$ from $\bar{\gamma}_{0} * \gamma$. Then it follows that we can choose a path $w$ from $\bar{\gamma}_{1} * \gamma_{0}$ defined by $u \# w_{1} \# w_{0}$, where $u$ is the above tiny triangle, that is, any $w$ such that

$$
[w]=B_{\mathrm{can}} \#\left[w_{1}\right] \#\left[w_{0}\right] .
$$

This finishes the proof.

Now we state the following lemma which is a variation of [Se3, Lemma 3.2] in our context.

## PROPOSITION 10.3

Let $L_{0}, L_{1}$, and $L$ be as in Lemma 10.1, and consider the maps $p, q$ defined therein. Then we can choose $\operatorname{im} \tau_{L}$ close to $L$ so that $\tau_{L}\left(L_{0}\right) \cap L_{1}$ satisfy the following properties in addition.
(1) The map $q$ is the inclusion $q(x)=x$. Moreover, for any $z \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}_{1}$ and $z \in \mathcal{L}_{0} \cap \mathcal{L}_{1}$ with $z \neq q(x)$, one has

$$
\mathcal{A}_{\mathcal{L}_{0} \mathcal{L}_{1}}([x, w])-\mathcal{A}_{\mathcal{L}_{0} \mathcal{L}_{1}}\left(\left[z, w^{\prime}\right]\right) \notin[0 ; 3 \varepsilon)
$$

whenever the corresponding Floer moduli space $\mathcal{M}\left(x, z ;\left[\bar{w} \# w^{\prime}\right]\right)$ is nonempty.
(2) Set $\widetilde{z}=p\left(\widetilde{x}_{0}, x_{1}\right)$. Then there is a canonical homotopy class $B_{\text {can }}=$ $B\left(\widetilde{x}_{0}, x_{1}, \widetilde{z}\right)$ spanned by a tiny triangle contained in $U=\operatorname{im} \iota$. And we have

$$
\begin{equation*}
\left|\Omega\left(B_{\text {can }}\right)\right|<\varepsilon . \tag{10.5}
\end{equation*}
$$

(3) For any $z \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}_{1}$ and $\left(\widetilde{x}_{0}, x_{1}\right) \in\left(\tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L}\right) \times\left(\mathcal{L} \cap \mathcal{L}_{1}\right)$ with $z \neq p\left(\widetilde{x}_{0}, x_{1}\right)$, or for $z=\widetilde{x}=p\left(\widetilde{x}_{0}, x_{1}\right)$ with $B \neq B_{\text {can }}$, we have

$$
\begin{equation*}
\Omega(B) \geq C=C(\mathcal{E} ; J) \tag{10.6}
\end{equation*}
$$

independent of $\varepsilon>0$ whenever $\mathcal{M}_{J}\left(\widetilde{x}_{0}, x_{1}, z ; B\right) \neq \emptyset$ for some $J$.
(4) Suppose that there are $x_{k} \in \mathcal{L} \cap \mathcal{L}_{k}, k=0,1$ whose preimages $y_{k}=\iota^{-1}\left(x_{k}\right)$ are antipodes on $S^{n}$. Since $\left.\tau\right|_{S^{n}}$ is the antipodal map, $\widetilde{x}_{0}=\tau_{L}\left(x_{0}\right)$ is equal to $x_{1}$. (Hence $x_{1} \in \tau_{L}\left(\mathcal{L}_{0}\right) \cap \mathcal{L} \cap \mathcal{L}_{1}$, and these are all such triple intersection points.) In that case $p\left(\widetilde{x}_{0}, x_{1}\right)=\widetilde{x}_{0}=x_{1}$.

Proof
The proof is a slight modification of that of [Se3, Lemma 3.2]. Since we need to strip all the exact Lagrangian setting away from Seidel's proof thereof and incorporate contributions coming from different choices of homotopy classes $B$, we give a complete proof of the proposition.

Now let $s(z)=(z, u(z))$ be the section of $E \rightarrow \dot{\Sigma}$ in class $B_{\text {can }}$ satisfying the Lagrangian boundary condition. By definition of $B_{\text {can }}$, we can choose $u$ so that its image is contained in $\operatorname{im} \iota$. Then (10.5) follows from the identity

$$
\Omega\left(B_{\text {can }}\right)=\pi^{*} \omega\left(B_{\text {can }}\right)=\omega_{0}\left(\iota^{-1} \circ u\right)
$$

which can be made as small as we want by choosing $r$ small in the definition of the Dehn twist $\tau_{L}$.

We now turn to (3). First, we recall that since $E$ is trivial (and so of zero curvature), we have

$$
\Omega(B)=\frac{1}{2} \int_{\dot{\Sigma}}\left\|(D u)^{v}\right\|_{J}^{2}
$$

and hence whenever $\mathcal{M}_{J}\left(\widetilde{x}_{0}, x_{1}, z ; B\right) \neq \emptyset$ for some $J$, we have $\Omega(B) \geq 0$. Define

$$
\begin{gathered}
C(\mathcal{E} ; J)=\inf _{u}\left\{\omega(u) \mid u \not \equiv \text { const }, u \in \mathcal{M}_{J}\left(\widetilde{x}_{0}, x_{1}, z ; B\right),\right. \\
\left.B \neq B_{\text {can }} \text { in } \Pi(E, Q)\right\} .
\end{gathered}
$$

## PROPOSITION 10.4

Let $B$ be an admissible class, and let $u \in \mathcal{M}_{J}\left(\widetilde{x}_{0}, x_{1}, z ; B\right)$ with $\mu(B)=0$ such that

Then we have $B=B_{\text {can }}$ in $\Pi(E, Q)$. In particular, we have

$$
\begin{equation*}
C(\mathcal{E} ; J)>0 \tag{10.7}
\end{equation*}
$$

## Proof

Since we know that $\mu\left(B_{\text {can }}\right)=0$, it is enough to prove that $\omega(B)=\omega\left(B_{\text {can }}\right)$ by definition of $\Pi(E, Q)$. For this purpose, we compare the action for the paths whose image is contained in the Darboux neighborhood $U=\iota(V)$ with $\iota:\left(V, \omega_{0}\right) \hookrightarrow$ $(M, \omega)$ and whose end points lie either on $L$, on the fibers $F$ of the cotangent bundle $T^{*} L \cap V$, or in the Dehn twists $\tau_{L}(F)$. We recall that the model Dehn twist $\tau$ is a Hamiltonian diffeomorphism that satisfies

$$
\tau^{*} \theta_{T}-\theta_{T}=d K_{\tau}
$$

for $K_{\tau}=2 \pi\left(\mu R^{\prime}(\mu)-R(\mu)\right)$.
On the cotangent bundle $T^{*} T$, the action functionals $\mathcal{A}_{o_{T} F_{1}}, \mathcal{A}_{\tau\left(F_{0}\right) o_{T}}$, and $\mathcal{A}_{\tau\left(F_{0}\right) F_{1}}$ are defined by

$$
\begin{aligned}
\mathcal{A}_{o_{T} F_{1}}(z) & =\int z^{*} \theta_{T} \\
\mathcal{A}_{\tau\left(F_{0}\right) o_{T}}(z) & =\int z^{*} \theta_{T}+K_{\tau} \circ \tau^{-1}(z(0)), \\
\mathcal{A}_{\tau\left(F_{0}\right) F_{1}}(z) & =\int z^{*} \theta_{T}+K_{\tau} \circ \tau^{-1}(z(0)),
\end{aligned}
$$

for a path $z:[0,1] \rightarrow T^{*} L$. Here we use the fact that

$$
\left.\left.\theta_{T}\right|_{F} \equiv 0 \equiv \theta_{T}\right|_{o_{T}},\left.\quad \theta_{T}\right|_{\tau(F)}=d\left(\left.K_{\tau} \circ \tau^{-1}\right|_{\tau(F)}\right)
$$

(see $[\mathrm{Oh} 2,(2.28)]$ or $[\mathrm{Se} 3,(1.1)])$. Since it is easy to realize such $B=[u]$ as the gluing $-\left[w_{0}\right] \#[w] \#\left[w_{1}\right]$ so that all $w_{i}, w$ have their images contained in $U$, we can write

$$
\omega(B)=\mathcal{A}_{o_{T} F_{1}}\left(\widehat{p}_{0}\right)-\mathcal{A}_{\tau\left(F_{0}\right) o_{T}}(\widehat{p})-\mathcal{A}_{\tau_{L}\left(F_{0}\right) F_{1}}\left(\widehat{p}_{1}\right)=\omega\left(B_{\text {can }}\right) .
$$

Here the "hat" denotes the constant path; for example, $\widehat{p}$ is the constant path $\widehat{p}(t) \equiv p$. This proves $B=B_{\text {can }}$ in $\Pi(E, Q)$.

For the proof of (10.7), it follows from the first part of the proof that for any $u \in \mathcal{M}_{J}\left(\widetilde{x}_{0}, x_{1}, z ; B\right)$ with $B \neq B_{\text {can }}$, the image of $u$ must go out of $U$. Now a simple application of the Gromov-type compactness theorem implies (10.7). This finishes the proof.

This finishes the proof of (3). The statements in (4) are obvious from definition.

Now we consider a Darboux neighborhood $U=\operatorname{im} \iota$ of $L$ such that $\operatorname{supp} \tau_{L} \subset U$ and Lemma 2.1 and Proposition 2.2 hold. We denote by $\mathcal{M}_{J}\left(\tau_{L}\left(L_{0}\right), L, L_{1}\right)$ the moduli space of $J$-holomorphic sections of the fibration $\left(E^{b}, \pi^{b}, Q^{b}\right)$, where $J$ is chosen as before. Then the arguments in [Se3, Section 3.2] give rise to the
decomposition

$$
\begin{equation*}
\mathfrak{b}=\beta+(\mathfrak{b}-\beta) \tag{10.8}
\end{equation*}
$$

that satisfies $\beta$ of order $[0 ; \varepsilon)$, while $(\mathfrak{b}-\beta)$ has order $[3 \varepsilon ; \infty)$. Furthermore, $b$ is precisely the class induced by the canonical "small" map continued from the constant map $p$.

Similar consideration of [Se3, Section 3.3] also gives rise to the decomposition

$$
\begin{equation*}
\mathfrak{c}=\gamma+(\mathfrak{c}-\gamma) \tag{10.9}
\end{equation*}
$$

such that $\gamma$ has order zero while $(\mathfrak{c}-\gamma)$ has order $[3 \varepsilon ; \infty)$. We refer to $[\mathrm{Se} 3]$ for the detailed proofs of the decomposition results.

Finally, we can construct a homotopy from $\mathfrak{b} \# \rho \mathfrak{c}$ to the zero map following Seidel's argument from [Se3]. For this purpose we use Proposition 9.2, which is a slight generalization of [Se3, Proposition 2.2] in the general context of Lagrangian spheres in general symplectic manifolds. We omit the details of this construction, referring readers to [Se3] for the details.

## 11. Construction of long exact sequence

In this section, we combine all the results obtained in Section 10 to construct the required long exact sequence. We first recall two basic lemmas that Seidel used in his construction of long exact sequences for the exact case.

## LEMMA 11.1

Let $D$ be an $\mathbb{R}$-graded vector space with a differential $d_{D}$ of order $[0 ; \infty)$. Suppose that $D$ has gap $[\varepsilon ; 2 \varepsilon)$ for some $\varepsilon>0$. One can write $d_{D}=\delta+\left(d_{D}-\delta\right)$ with $\delta$ of order $[0 ; \varepsilon)$, satisfying $\delta^{2}=0$, and $\left(d_{D}-\delta\right)$ of order $[2 \varepsilon ; \infty)$. Suppose in addition that $H(D, \delta)=0$; then $H\left(D, d_{D}\right)=0$.

Seidel then applied this lemma to the direct sum

$$
D=C^{\prime} \oplus C \oplus C^{\prime \prime}
$$

with the differentials given by

$$
d_{D}=\left(\begin{array}{ccc}
d_{C^{\prime}} & 0 & 0 \\
b & d_{C} & 0 \\
h & c & d_{C^{\prime \prime}}
\end{array}\right), \quad \delta=\left(\begin{array}{ccc}
0 & 0 & 0 \\
\beta & 0 & 0 \\
0 & \gamma & 0
\end{array}\right),
$$

where

$$
\begin{aligned}
C^{\prime} & =C F\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \\
C & =C F\left(\tau_{L}\left(L_{0}\right), L_{1}\right), \quad C^{\prime \prime}=C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right),
\end{aligned}
$$

and the entries of the matrices are given as stated in the lemma below.

LEMMA 11.2 ([Se3, LEMMA 2.32])
Take three $\mathbb{R}$-graded vector spaces $C^{\prime}, C, C^{\prime \prime}$, each of them with a differential of
order $(0 ; \infty)$. Suppose that we have the differential maps $b: C^{\prime} \rightarrow C, c: C \rightarrow C^{\prime \prime}$ and a homotopy $h: C^{\prime} \rightarrow C^{\prime \prime}$ between $c \circ b$ and the zero map such that the following conditions are satisfied for some $\varepsilon>0$.
(1) $C^{\prime}, C^{\prime \prime}$ have gap $(0,3 \varepsilon)$, and $C$ has gap $(0,2 \varepsilon)$.
(2) For all $r \in \operatorname{supp}\left(C^{\prime}\right)$ and $s \in \operatorname{supp}\left(C^{\prime \prime}\right),|r-s| \geq 4 \varepsilon$.
(3) One can write

$$
b=\beta+(b-\beta), \quad c=\gamma+(c-\gamma)
$$

with $\beta$ of order $[0 ; \varepsilon)$ and $(b-\beta)$ of order $[2 \varepsilon ; \infty)$ and with the same properties for $\gamma$ and $(c-\gamma)$. The lower order parts (which do not need to be differential maps) fit into a short exact sequence of modules

$$
\begin{equation*}
0 \rightarrow C^{\prime} \xrightarrow{\beta} C \xrightarrow{\gamma} C^{\prime \prime} \rightarrow 0 \tag{11.1}
\end{equation*}
$$

(4) The map $h$ is of order $[0 ; \infty)$.

Then the maps on cohomology induced by b, cfit into a long exact sequence

$$
\cdots \rightarrow H\left(C^{\prime} ; d_{C^{\prime}}\right) \xrightarrow{b_{*}} H\left(C ; d_{C}\right) \xrightarrow{c_{*}} H\left(C^{\prime \prime} ; d_{C^{\prime \prime}}\right) \xrightarrow{\delta} H\left(C^{\prime} ; d_{C^{\prime}}\right) \rightarrow \cdots .
$$

The proofs of both lemmas rely on an argument involving spectral sequences. For the exact case, all the complexes involved are finite-dimensional vector spaces with bounded filtration and gap and so existence of spectral sequences for such a complex is easy.

On the other hand, for the case of our current interest, the Floer complex as a $\mathbb{Q}$-vector space is infinite-dimensional with unbounded filtration and without gap on the vector space itself in general. The existence of a spectral sequence in this case is much more nontrivial (this has been studied by Fukaya, Oh, Ohta, and Ono in $[\mathrm{FO}+3],[\mathrm{FO}+1])$. A crucial algebraic model of such a spectral sequence is the notion of a differential graded completed free filtered $\Lambda_{0, \text { nov }}^{(0)}-$ module, abbreviated as d.g.c.f.z. The upshot of this definition is that the $\Lambda_{0, \text { nov }}$-module $C\left(L, L^{\prime} ; \Lambda_{0, \text { nov }}\right)$ naturally carries this structure. We summarize this construction of spectral sequences in Section A.2.

In the meantime, we remark that the proof of exactness of (11.1) is exactly the same as that of $[\mathrm{Se} 3]$ based on Lemmas $10.1,10.2$, and on some uniqueness result on a small pseudoholomorphic triangle. (See Section 3.2 [Se3] for details.)
11.1. $C F\left(\mathcal{L}, \mathcal{L}^{\prime}\right)$ versus $C\left(\mathcal{L}, \mathcal{L}^{\prime} ; \Lambda_{0, \text { nov }}\right)$

We first recall basic results on the structure of the Floer cochain group $C F\left(\mathcal{L}, \mathcal{L}^{\prime}\right)$ as a module of the Novikov ring $\Lambda\left(\mathcal{L}, \mathcal{L}^{\prime}\right)$. We note that in the current study of CY Lagrangian branes we can use the Novikov ring

$$
\Lambda_{0, \text { nov }}^{(0)}
$$

as the coefficient ring, where $\Lambda_{0, \text { nov }}^{(0)}$ is the degree zero part of $\Lambda_{0, \text { nov }}$ which is a field. Recall the definition from (4.2).

Having this in mind, we first recall the basic construction on the spectral sequence of the $\Lambda_{0, \text { nov }}^{(0)}$-module $C\left(L, L^{\prime} ; \Lambda_{0, \text { nov }}\right)$ from $[\mathrm{FO}+1$, Chapter 6$]$ (or $[\mathrm{FO}+3$, Appendix] $)$ restricting to the finitely generated case.

Under the given condition on $L_{0}, L_{1}$ and the given embedding $f: S^{n} \rightarrow$ $L \subset M, C F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right), C F\left(\mathcal{L}, \mathcal{L}_{1}\right)$ and $C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)$ are all finitely generated over $\Lambda_{\text {nov }}$. Here we note that there is a natural injective homomorphism

$$
I_{\mathcal{L}, \mathcal{L}^{\prime}}: \Lambda\left(\mathcal{L}, \mathcal{L}^{\prime}\right) \rightarrow \Lambda_{0, \text { nov }}
$$

and so we naturally extend their coefficient rings to $\Lambda_{0, \text { nov }}$.
The tiny-big decompositions of the maps $\mathfrak{b}$ and $\mathfrak{c}$ given in Section 10 imply that both maps are gapped in the sense of Definitions A. 3 and A.5. However, these vector spaces, as they are, do not quite manifest the structure of d.g.c.f.z. yet. Because of this, we follow the procedure given in $[\mathrm{FO}+1$, Section 12.4] turning these into a d.g.c.f.z. For the reader's convenience, we collect the definition of d.g.c.f.z. and the construction of spectral sequence given in $[\mathrm{FO}+1$, Appendix].

Let $\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ be a general relatively spin pair of anchored Lagrangian submanifolds of $M$. We first construct a $\Lambda_{0, \text { nov }}$-module $C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ which has a filtered $A_{\infty}$-bimodule structure over $\left(C\left(L_{0} ; \Lambda_{0, \text { nov }}\right), \mathfrak{m}^{(0)}\right)$ and $\left(C\left(L_{1} ; \Lambda_{0, \text { nov }}\right), \mathfrak{m}^{(1)}\right)$, where the latter are the filtered $A_{\infty}$-algebras defined in $[\mathrm{FO}+1]$. We consider the intersection $\mathcal{L}_{1} \cap \mathcal{L}_{0}$ and the $\mathbb{R}$-filtered set

$$
\widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right):=\left\{T^{\lambda} e^{\mu}[p, w] \mid p \in \mathcal{L}_{0} \cap \mathcal{L}_{1},[p, w] \in \widetilde{\Omega}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right), \lambda \in \mathbb{R}, \mu \in \mathbb{Z}\right\}
$$

This is an $\mathbb{R} \times \mathbb{Z}$ principal bundle over $\mathcal{L}_{0} \cap \mathcal{L}_{1}$.
We define an equivalence relation $\sim$ on $\widehat{I}\left(L_{0}, L_{1}\right)$ as follows. We say that

$$
T^{\lambda} e^{\mu}[p, w] \sim T^{\lambda^{\prime}} e^{\mu^{\prime}}\left[p^{\prime}, w^{\prime}\right]
$$

for $T^{\lambda} e^{\mu}[p, w], T^{\lambda^{\prime}} e^{\mu^{\prime}}\left[p^{\prime}, w^{\prime}\right] \in \widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ if and only if the following conditions are satisfied:

$$
\begin{aligned}
p & =p^{\prime} \\
\lambda+\int w^{*} \omega & =\lambda^{\prime}+\int\left(w^{\prime}\right)^{*} \omega \\
2 \mu+\mu\left([p, w] ; \lambda_{01}\right) & =2 \mu^{\prime}+\mu\left(\left[p^{\prime}, w^{\prime} ; \lambda_{01}\right]\right)
\end{aligned}
$$

Here $\mu\left([p, w] ; \lambda_{01}\right)$ is the Maslov-Morse index. It is easy to see that this relation is compatible with the conditions of the $\Gamma$-equivalence given, and so $\sim$ defines an equivalence relation on $\widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$. Furthermore, we define the action level on $E: \widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \rightarrow \mathbb{R}$ by

$$
E\left(T^{\lambda} e^{\mu}[p, w]\right)=\lambda+\int w^{*} \omega
$$

and the associated filtration on the set by setting

$$
T^{\lambda} e^{\mu}[p, w] \in F^{\lambda^{\prime}}\left(\widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)\right)
$$

if

$$
\lambda+\int w^{*} \omega \geq \lambda^{\prime}
$$

We now define

$$
I\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)=\widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) / \sim
$$

and somewhat ambiguously denote an element thereof still by $T^{\lambda} e^{\mu / 2}[p, w]$ as long as no danger of confusion arises. The above-mentioned filtration on $\widehat{I}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$ obviously induces on the quotient $I\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)$. We now define

$$
I_{\geq 0}\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right):=\left\{T^{\lambda} e^{\mu / 2}[p, w] \in I\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \mid \lambda+\int w^{*} \omega \geq 0\right\} .
$$

Consider the formal sum

$$
\alpha=\sum_{\lambda, \mu,[p, w]} a_{\lambda, \mu,[p, w]} T^{\lambda} e^{\mu / 2}[p, w]
$$

for $\lambda \in \mathbb{R}, \mu \in \mathbb{Z}$ and $[p, w] \in \operatorname{Crit} \mathcal{A}$, and define $\operatorname{supp} \alpha$ to be

$$
\operatorname{supp} \alpha=\left\{T^{\lambda} e^{\mu / 2}[p, w] \in I\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \mid a_{\lambda, \mu,[p, w]} \neq 0\right\}
$$

DEFINITION 11.1
We define by $C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ the $\Lambda_{0, \text { nov }}$-module

$$
\begin{aligned}
& C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right) \\
& \quad:=\left\{\alpha \mid E(\alpha) \geq 0, \#\left(\operatorname{supp} \alpha \cap E^{-1}((-\infty, \lambda])\right)<\infty \text { for all } \lambda \in \mathbb{R}\right\} .
\end{aligned}
$$

Obviously $C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ has the structure of a $\Lambda_{0, \text { nov }}$-module. In addition, we have the following.

## PROPOSITION 11.3

$C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ is a d.g.c.f.z.
Given the grading of an element $T^{\lambda} e^{\mu}[p, w] 2 \mu+\mu\left([p, w] ; \lambda_{01}\right)$, it becomes a filtered graded free $\Lambda_{0, \text { nov }}$-module. Following [FO+1, Section 5.1.3], we write

$$
\langle p\rangle= \begin{cases}T^{-\mathcal{A}_{\mathcal{L}_{0} \mathcal{L}_{1}}([p, w])} e^{-\mu\left([p, w] ; \lambda_{01}\right)}[p, w] & \text { if } \mu\left([p, w] ; \lambda_{01}\right) \text { is even, } \\ T^{-\mathcal{A}_{\mathcal{L}_{0} \mathcal{L}_{1}}([p, w])} e^{-\left(\mu\left([p, w] ; \lambda_{01}\right)-1\right) / 2}[p, w] & \text { if } \mu\left([p, w] ; \lambda_{01}\right) \text { is odd. }\end{cases}
$$

Thus we have $E(\langle p\rangle)=0$, and $\operatorname{deg}(\langle p\rangle)$ is either zero or 1 depending on the parity of $\mu\left([p, w] ; \lambda_{01}\right)$.

It is easy to see that $C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ is isomorphic to the completion (with respect to the filtration on $\Lambda_{0, \text { nov }}$ ) of the free $\Lambda_{0, \text { nov }}$-module generated by $\langle p\rangle$ for the intersection points $p \in \mathcal{L}_{0} \cap \mathcal{L}_{1}$. Namely, we have a canonical isomorphism

$$
C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right) \cong \widehat{\bigoplus}_{p \in \mathcal{L}_{0} \cap \mathcal{L}_{1}} \Lambda_{0, \text { nov }}\langle p\rangle
$$

as a ( $\mathbb{Z}_{2}$-graded) $\Lambda_{0, \text { nov }}$-module.

Recall the definition of the Floer cochain module. By definition, we have an inclusion

$$
C F^{*}\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \rightarrow C^{*}\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{\text {nov }}\right)
$$

defined by

$$
\begin{equation*}
[p, w] \mapsto e^{(\mu([p, w])-\mu(\langle p\rangle)) / 2} T^{\rho A_{\ell_{0}}([p, w])}\langle p\rangle . \tag{11.2}
\end{equation*}
$$

It is compatible with the obvious inclusion $I_{\mathcal{L}_{0}, \mathcal{L}_{1}}: \Lambda\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right) \rightarrow \Lambda_{\text {nov }}$.
We take the coefficient $R=\mathbb{Q}$ and recall that Floer cohomology

$$
H F\left(\mathcal{L}^{(1)}, \mathcal{L}^{(0)}\right)=\operatorname{Ker} \delta^{\vec{b}} / \operatorname{Im} \delta^{\vec{b}}, \quad \vec{b}=\left(b_{1}, b_{0}\right),
$$

is defined as a $\Lambda\left(\mathcal{L}^{(0)}, \mathcal{L}^{(1)}\right)$-module.
We remark that

$$
C F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \cong \Lambda\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)^{\#\left(\mathcal{L}_{0} \cap \mathcal{L}_{1}\right)}
$$

where $\# \mathcal{L}_{0} \cap \mathcal{L}_{1}$ is finite by the transversality hypothesis. Therefore we have the isomorphism

$$
C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{\text {nov }}\right) \cong C F\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \ell_{01}\right) \otimes_{\Lambda\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)} \Lambda_{\text {nov }}
$$

On the other hand, the Novikov ring $\Lambda\left(\mathcal{L}_{1}, \mathcal{L}^{(0)}\right)$ is a field if the ground ring is $\mathbb{Q}$. Therefore this leads to the isomorphism

$$
\begin{equation*}
H F\left(\left(\mathcal{L}_{1}, b_{1}\right),\left(\mathcal{L}_{0}, b_{0}\right) ; \Lambda_{\text {nov }}\right) \cong H F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \otimes_{\Lambda\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right)} \Lambda_{\text {nov }} . \tag{11.3}
\end{equation*}
$$

Finally, we explain how we combine the above-discussed anchored versions into a single nonanchored version of Floer cohomology following [FO+1, Section 5.1.3].

We first note that the filtered $\Lambda_{0, \text { nov }}$-module structure of $C\left(\mathcal{L}_{1}, \mathcal{L}_{0} ; \Lambda_{0, \text { nov }}\right)$ depends only on the homotopy class $\ell_{01}$. So we form the completed direct sum

$$
C\left(L_{1}, L_{0} ; \Lambda_{0, \text { nov }}\right)=\widehat{\bigoplus}_{\left[\ell_{0}\right] \in \pi_{0}\left(\Omega\left(L_{0}, L_{1}\right)\right)} C\left(L_{1}, L_{0} ; \ell_{0} ; \Lambda_{0, \text { nov }}\right) .
$$

We note that we have the natural inclusion map

$$
C\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \rightarrow C\left(L_{1}, L_{0} ; \bar{\gamma}_{0} * \gamma_{1} ; \Lambda_{0, \text { nov }}\right) \subset C\left(L_{1}, L_{0} ; \Lambda_{0, \text { nov }}\right)
$$

defined as (11.2). We define the corresponding Floer cohomology by

$$
H F\left(L_{1}, L_{0} ; \Lambda_{0, \text { nov }}\right):=\operatorname{Ker} \delta^{\vec{b}} / \operatorname{Im} \delta^{\vec{b}} .
$$

Then we have

$$
H F\left(L_{1}, L_{0} ; \Lambda_{0, \text { nov }}\right) \cong \bigoplus_{\left[\ell_{0}\right] \in \pi_{0}\left(\Omega\left(L_{0}, L_{1}\right)\right)} H F\left(\mathcal{L}_{1}, \mathcal{L}_{0}\right) \otimes_{\Lambda\left(L_{1}, L_{0} ; \ell_{0}\right)} \Lambda_{\mathrm{nov}}
$$

(see [FO+1, Proposition 5.1.17]).

### 11.2. Wrapping it up

Now combining all the discussions in the previous sections, we are ready to prove the main theorem, which we now restate here.

## THEOREM 11.4

Let $(M, \omega)$ be compact (symplectically) Calabi-Yau. Let L be a Lagrangian sphere in $M$ together with a preferred diffeomorphism $f: S^{2} \rightarrow L$. Denote by $\tau_{L}=\tau_{(L,[f])}$ the Dehn twist associated to $(L,[f])$.

Consider any CY Lagrangian branes $L_{0}, L_{1}$. Then for any pair $\left(b_{0}, b_{1}\right)$ of the Maurer-Cartan-solutions $b_{0} \in \mathcal{M}\left(L_{0} ; \Lambda\left(L_{0}\right)\right)$, $b_{1} \in \mathcal{M}\left(L_{1} ; \Lambda\left(L_{1}\right)\right)$, there is a long exact sequence of $\mathbb{Z}$-graded Floer cohomologies

$$
\begin{align*}
& \longrightarrow H F\left(\left(\tau_{L}\left(\mathcal{L}_{0}\right),\left(\tau_{L}\right)_{*}\left(b_{0}\right)\right),\left(\mathcal{L}_{1}, b_{1}\right)\right) \longrightarrow H F\left(\left(\mathcal{L}_{0}, b_{0}\right),\left(\mathcal{L}_{1}, b_{1}\right)\right)  \tag{11.4}\\
& \longrightarrow \operatorname{HF}\left((\mathcal{L}, 0),\left(\mathcal{L}_{1}, b_{1}\right)\right) \otimes \operatorname{HF}\left(\left(\mathcal{L}_{0}, b_{0}\right),\left(\mathcal{L}_{1}, b_{1}\right)\right) \longrightarrow
\end{align*}
$$

as a $\Lambda_{\text {nov }}$-module where the Floer cohomologies involved are the deformed Floer cohomologies constructed in $[F O+3]$ and $[F O+1]$. We also have the nonanchored version of the exact sequence.

The same exact sequence still holds for any orientable relatively spin pair ( $L_{0}, L_{1}$ ) if they are just unobstructed, whose Maslov classes do not necessarily vanish.

To highlight the main points of the construction, let us first assume that $b_{0}=$ $b_{1}=0$ are Maurer-Cartan-solutions. In this case, the Floer cohomology is the standard one which uses the Floer boundary map $\delta$. We first state the following lemma, which is a consequence of Corollary A. 8 and a variation of Lemma 11.2

## LEMMA 11.5

Let $D$ be an $[0 ; \infty)$-graded vector space with a differential $d_{D}$ of order $[0 ; \infty)$, which is not necessarily finite-dimensional but forms a d.g.c.f.z. in the sense of Definition A.2. Suppose that one can then write $d_{D}=\delta+\left(d_{D}-\delta\right)$ with $\delta$ of order $[0 ; \varepsilon)$, satisfying $\delta^{2}=0$, and $\left(d_{D}-\delta\right)$ of order $[2 \varepsilon ; \infty)$. Suppose in addition that $H(D, \delta)=0$; then $H\left(D, d_{D}\right)=0$.

Now we apply this lemma to the direct sum

$$
D=C^{\prime} \oplus C \oplus C^{\prime \prime}
$$

with the differentials given by

$$
d_{D}=\left(\begin{array}{ccc}
d_{C^{\prime}} & 0 & 0 \\
b & d_{C} & 0 \\
h & c & d_{C^{\prime \prime}}
\end{array}\right), \quad \delta=\left(\begin{array}{ccc}
0 & 0 & 0 \\
\beta & 0 & 0 \\
0 & \gamma & 0
\end{array}\right),
$$

where this time we consider the $\Lambda_{0, \text { nov }}$-modules

$$
\begin{aligned}
C^{\prime} & =C\left(\mathcal{L}, \mathcal{L}_{1}\right) \otimes C\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}\right) \\
C & =C\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1}\right), \quad C^{\prime \prime}=C\left(\mathcal{L}_{0}, \mathcal{L}_{1}\right)
\end{aligned}
$$

and the entries of the matrices are given as stated before.

Now Lemma 10.3 and the tiny-big decomposition results in Section 10 give rise to a long exact sequence

$$
\begin{align*}
& \longrightarrow H F\left(\tau_{L}\left(\mathcal{L}_{0}\right), \mathcal{L}_{1} ; \Lambda_{0, \text { nov }}\right) \longrightarrow H F\left(\mathcal{L}_{0}, \mathcal{L}_{1} ; \Lambda_{0, \text { nov }}\right) \\
& \longrightarrow H F\left(\mathcal{L}, \mathcal{L}_{1} ; \Lambda_{0, \text { nov }}\right) \otimes H F\left(\mathcal{L}_{0}, \mathcal{L}_{1} ; \Lambda_{0, \text { nov }}\right) \longrightarrow \tag{11.5}
\end{align*}
$$

for $\Lambda_{0, \text { nov }}^{(0)}$-modules. Since we have

$$
H F\left(\mathcal{L}, \mathcal{L}^{\prime} ; \Lambda_{\text {nov }}\right) \cong H F\left(\mathcal{L}, \mathcal{L}^{\prime} ; \Lambda_{0, \text { nov }}\right) \otimes_{\Lambda_{0, \text { nov }}} \Lambda_{\text {nov }}
$$

from (11.3) and since $\Lambda_{\text {nov }}$ is a field, tensoring (11.5) with $\Lambda_{\text {nov }}$ produces the exact sequence (11.4) for the case $b_{0}=b_{1}=0$.

Now the same reasoning as for the case $b_{i}=0$ induces the long exact sequence (11.4). This finishes the proof of Theorem 11.4.

## Appendix

## A.1. Index formula for $E_{z_{0}} \backslash\left\{x_{0}\right\}$

In this section, we prove the index formula (7.27). There is an index formula stated in the various literature in terms of the "capping surfaces" stated as in [EGH] and [Bo], which, however, does not fit our need. For this reason, we give a complete proof of (7.27).

In fact, we consider the following general setup. Consider a symplectic manifold $W$ with a contact-type boundary of the type

$$
\partial W \cong S^{1}\left(T^{*} N\right)
$$

with negative end for an oriented compact manifold $N$. We attach the cylinder $\mathbb{R}_{+} \times \partial W$ and also denote by $W$ the completed manifold. We denote $(r(x), \Theta(x))$ for a point $x \in \mathbb{R}^{+} \times S^{1}\left(T^{*} N\right)$. Composing this with the diffeomorphism

$$
(s, \Theta) \mapsto\left(e^{s}, \Theta\right), \quad \mathbb{R} \times \partial W \rightarrow \mathbb{R}_{+} \times \partial W
$$

we put a translational invariant almost complex structure $J$ on the end.
Next let $\gamma$ be a Reeb orbit of $S^{1}\left(T^{*} N\right)$ with period $T$. We note that the symplectic vector bundle $\gamma^{*} T\left(T^{*} N\right)$ carries a splitting

$$
\gamma^{*} T\left(T^{*} N\right)=\mathbb{C} \oplus \gamma^{*} \xi_{N},
$$

where $\xi_{N}$ is the contact distribution of $S^{1}\left(T^{*} N\right)$. Furthermore, we fix a Riemannian metric $g$ on $N$ and consider the canonical almost complex structure $J_{g}$ on $T^{*} N$. The projection of $\gamma$ to $N$ is nothing but a geodesic on $N$ with respect to $g$. Denote by $c=c_{\gamma}$ the associated geodesic on $N$. Since we assume that $N$ is oriented, we can take a trivialization $\gamma^{*} T\left(S^{1}\left(T^{*} N\right)\right)$ which is tangent to the vertical fibers of $T\left(T^{*} N\right)$. Using this, we can define the Conley-Zehnder index of $\gamma$ when $\gamma$ is nondegenerate, which we denote by $\mu_{C Z}(\gamma)$. For the Bott-Morse case, one uses the generalized Conley-Zehnder index defined by Robbin and Salamon [RS].

Next, this choice of trivialization of $\left.\gamma^{*}\left(T\left(S^{1}\left(T^{*} N\right)\right)\right)=\gamma^{*} T(\partial W)\right)$ also allows one to define a relative Chern number of a map $u: \dot{\Sigma} \rightarrow W$ with the asymptotic
condition

$$
\begin{equation*}
\lim _{\tau \rightarrow \infty} \Theta \circ u(\tau, t)=\gamma(t), \quad \lim _{\tau \rightarrow \infty} s \circ u(\tau, t)=-\infty . \tag{A.1}
\end{equation*}
$$

Denote by $\bar{u}:(\hat{\dot{\Sigma}}, \partial \widehat{\dot{\Sigma}}) \rightarrow(W, \gamma)$ the obvious compactified map.
Then $\bar{u}^{*}(T W)$ is a symplectic vector bundle with a trivialization $\phi_{\gamma}$ : $\gamma^{*}(T(\partial W)) \rightarrow S^{1} \times \mathbb{C}^{n-1}$ constructed above.

This gives rise to the main definition.

## DEFINITION A. 1

We define the relative Chern number, denoted by $c_{1}(u ; \gamma)$, by

$$
c_{1}(u ; \gamma)=c_{1}\left(u^{*} T W ; \phi_{\gamma}\right) .
$$

Once we have made the notions of relative Chern number and Conley-Zehnder index precise, the following index formula can be derived from the formula in [Bo, Corollary 5.4].

## THEOREM A. 1

The expected dimension of $\mathcal{M}(W, J ; \gamma ; A)$ is given by

$$
-\mu_{C Z}(\gamma)+(n-3)+2 c_{1}(u ; \gamma), \quad[u]=A
$$

for a nondegenerate geodesic.
For the Morse-Bott case in which $\mathcal{R}_{\text {sim }}$ forms a smooth manifold, the expected dimension of the moduli space $\mathcal{M}(W, J ; A ; 1)$ consisting of $J$-holomorphic u's with asymptotics

$$
\lim _{\tau \rightarrow \mathrm{inf}} \Theta \circ u\left(e^{2 \pi(\tau+i t)}\right)=\gamma(t),
$$

where $\gamma$ is a simple Reeb orbit, is given by

$$
-\mu_{C Z}(\gamma)+\frac{\operatorname{dim} \mathcal{R}_{\mathrm{sim}}}{2}+(n-3)+2 c_{1}(u ; \gamma)
$$

where $\mu_{C Z}$ is the generalized Conley-Zehnder index of $\gamma$.

## A.2. d.g.c.f.z. and spectral sequence

We first start from the following situation. Let $V=\left(\Lambda_{0, \text { nov }}^{(0)}\right)^{\oplus I}$ be a free $\Lambda_{0, \text { nov }}^{(0)}{ }^{-}$ module with $\#(I)$ finite. We define a filtration on $V$ in the obvious way, which induces a topology on $V$. Let $\widehat{V}$ be the completion of $V$. We call such $\widehat{V}$ a completed free filtered $\Lambda_{0, \text { nov }}^{(0)}$-module generated by energy zero elements or, in short, c.f.z. If $V$ is finitely generated (as a $\Lambda_{0, \text { nov }}^{(0)}$-module) in addition, we say that it is a finite c.f.z. We define a function, which we call the (action) level,

$$
E: \widehat{V} \backslash\{0\} \rightarrow \mathbb{R}_{\geq 0}
$$

such that

$$
\mathfrak{v} \in F^{E(\mathfrak{v})} V, \quad \mathfrak{v} \notin F^{\lambda} V \quad \text { if } \lambda>E(\mathfrak{v}) .
$$

Let $\bar{V}=V / \Lambda_{0, \text { nov }}^{+,(0)} V \cong R^{I}$. We always take an embedding (splitting)

$$
\bar{V} \subset V
$$

as the energy zero part of $V$ so that its composition with the projection $V \rightarrow \bar{V}$ is the identity map.

Let $\mathfrak{v} \in V$. We put

$$
\mathfrak{v}=\sum T^{\lambda_{i}} v_{i},
$$

where $v_{i} \in \bar{V}, \lambda_{i}<\lambda_{i+1}, \lim _{i \rightarrow \infty} \lambda_{i}=\infty$, and $v_{i} \neq 0$. We call $T^{\lambda_{i}} v_{i}$ the components of $\mathfrak{v}$, with $T^{\lambda_{1}} v_{1}$ the leading component and $v_{1}$ the leading coefficient of $\mathfrak{v}$. We denote the leading coefficient $v_{1}$ of $\mathfrak{v}$ by $\sigma(\mathfrak{v})$. We also define the leading component and the leading coefficient of an element of $\Lambda_{0, \text { nov }}^{(0)}$ in the same way.

Now we consider the case of graded $\Lambda_{0, \text { nov }}$-modules.

## DEFINITION A. 2 ([FO+1, DEFINITION 6.3.8])

Let $\widehat{C}$ be a graded $\Lambda_{0, \text { nov }}$-module. We assume that $\widehat{C}^{k}$ is a c.f.z. for each $k$. A differential graded c.f.z. (abbreviated as d.g.c.f.z.) is a pair $(\widehat{C}, \delta)$ with a degree 1 operator $\delta: \widehat{C} \rightarrow \widehat{C}$ such that

$$
\delta \circ \delta=0, \quad \delta\left(F^{\lambda} \widehat{C}\right) \subseteq F^{\lambda} \widehat{C} .
$$

We call the pair a finite d.g.c.f.z. if each $\widehat{C}^{k}$ is a finite c.f.z.
The following proposition is essential for the proof of some convergence properties of the spectral sequence.

PROPOSITION A. 2 ([FO+1, PROPOSITION 6.3.9])
Let $W$ be a finitely generated $\Lambda_{0, \text { nov }}^{(0)}$-submodule of $\widehat{C}^{k}$. Then there exists a constant $c$ depending only on $W$ but independent of $\lambda$ such that

$$
\delta(W) \cap F^{\lambda} \widehat{C}^{k+1} \subset \delta\left(W \cap F^{\lambda-c} \widehat{C}^{k}\right)
$$

Now let $(\widehat{C}, \delta)$ be a d.g.c.f.z., and let $\widehat{C}^{k}$ be a completion of $C^{k}$. We assume that $C^{k}$ is free over $\Lambda_{0, \text { nov }}^{(0)}$. We put

$$
\bar{C}=C / \Lambda_{0, \text { nov }}^{+(0)} C \cong \widehat{C} / \Lambda_{0, \text { nov }}^{+(0)} \widehat{C}
$$

and let $\bar{\delta}$ be the induced derivation on $\bar{C}$. We again embed $\bar{C} \subseteq C \subseteq \widehat{C}$ as the energy zero part. In general $\bar{C}$ is not a differential graded subalgebra of $\widehat{C}$. Let $\left\{\mathfrak{e}_{i}\right\}$ be a basis of $C$ (over $\Lambda_{0, \text { nov }}^{(0)}$ ), and let $\overline{\mathfrak{e}}_{i}$ be the corresponding basis of $\bar{C}$ over $R=\Lambda_{0, \text { nov }}^{(0)} / \Lambda_{0, \text { nov }}^{+(0)}$. We put

$$
\bar{\delta}\left(\overline{\mathfrak{e}}_{i}\right)=\sum \delta_{0, i j} \overline{\mathfrak{e}}_{j}
$$

and define $\delta_{0}: \widehat{C} \rightarrow \widehat{C}$ by $\delta_{0}=\bar{\delta} \otimes 1$, that is, by

$$
\delta_{0} \mathfrak{e}_{i}=\sum \delta_{0, i j} \mathfrak{e}_{j} .
$$

## DEFINITION A. 3

We say that $(\widehat{C}, \delta)$ satisfies the gapped condition if $\left(\delta-\delta_{0}\right)$ has order $\left[\lambda^{\prime \prime}, \infty\right)$, that is, if there exists $\lambda^{\prime \prime}>0$ such that for any $\lambda$ we have

$$
\delta \mathfrak{v}-\delta_{0} \mathfrak{v} \in F^{\lambda+\lambda^{\prime \prime}} \widehat{C}
$$

for all $\mathfrak{v} \in F^{\lambda} \widehat{C}$.
Under the gapped condition, we take a constant $\lambda_{0}$ with $0<\lambda_{0}<\lambda^{\prime \prime}$ and define a filtration on $\widehat{C}$ by

$$
F^{n} \widehat{C}=F^{n \lambda_{0}} \widehat{C}
$$

$[\mathrm{FO}+1]$ then uses this filtration to define a spectral sequence.

LEMMA A. 3 ([FO+1, LEMMA 6.3.20])
Denote

$$
\Lambda^{(0)}\left(\lambda_{0}\right)=\Lambda_{0, \text { nov }}^{(0)} / F^{\lambda_{0}} \Lambda_{0, \text { nov }}^{(0)} .
$$

Then there exists a $\Lambda^{(0)}\left(\lambda_{0}\right)$-module homomorphism

$$
\delta_{r}^{p, q}: E_{r}^{p, q}(\widehat{C}) \rightarrow E_{r}^{p+1, q+r-1}(\widehat{C})
$$

such that
(1) $\delta_{r}^{p+1, q+r-1} \circ \delta_{r}^{p, q}=0$,
(2) $\operatorname{Ker}\left(\delta_{r}^{p, q}\right) / \operatorname{Im}\left(\delta_{r}^{p-1, q-r+1}\right) \cong E_{r+1}^{p, q}(\widehat{C})$,
(3) $e^{ \pm 1} \circ \delta_{r}^{p, q}=\delta_{r}^{p \pm 2, q} \circ e^{ \pm 1}$.

Of course, the construction of $E_{r}^{p, q}(\widehat{C})$ is quite standard. One difference from the standard case is that the filtration used here is not bounded. Namely, we do not have $F^{n} \widehat{C}=0$ for large $n$. Hence the convergence property of our spectral sequence is far from being trivial in general. However, it is stable from below in that $F^{0} \widehat{C}=\widehat{C}$. As a consequence, we have the following.

LEMMA A. 4 ([FO+1, LEMMA 6.3.22])
There exists an injection

$$
E_{r+1}^{p, q}(\widehat{C}) \rightarrow E_{r}^{p, q}(\widehat{C})
$$

if $q-r+2 \leq 0$.

An immediate consequence of Lemma A. 4 is the following convergence result.

PROPOSITION A. 5
The projective limit

$$
E_{\infty}^{p, q}(\widehat{C}):=\lim _{\leftarrow} E_{r}^{p, q}(\widehat{C})
$$

exists.

Furthermore, from the construction, we have the description of the $E_{2}$-term of the associated spectral sequence.

LEMMA A. 6 ([FO+1, LEMMA 6.3.24])
We have an isomorphism

$$
E_{2}^{*, *}(\widehat{C}) \cong H(\bar{C} ; \bar{\delta}) \otimes_{R} \operatorname{gr}_{*}\left(F \Lambda_{0, \text { nov }}\right)
$$

$a s \operatorname{gr}_{*}\left(F \Lambda_{0, \text { nov }}\right)$-modules.

## Proof

By definition, we have

$$
E_{1}^{*, *}(\widehat{C}) \cong \bar{C} \otimes_{R} \operatorname{gr}_{*}\left(F \Lambda_{0, \text { nov }}\right)
$$

It follows from the gapped condition $\delta_{1}=\bar{\delta}$. Hence it finishes the proof.

## DEFINITION A. 4

We define $F^{q} H(\widehat{C}, \delta)$ to be the image of $H\left(F^{q} \widehat{C}, \delta\right)$ in $H(\widehat{C}, \delta)$.
To relate the limit $E_{\infty}^{p, q}$ of the spectral sequence and $F^{q} H(\widehat{C}, \delta)$, we need some finiteness assumptions which we now describe. Let $(C, \delta)$ and $\left(C^{\prime}, \delta^{\prime}\right)$ be d.g.c.f.z.'s satisfying the gap condition. Let $\varphi: C \rightarrow C^{\prime}$ be a map such that $\varphi \delta=\delta^{\prime} \varphi$, and let $\bar{\varphi}: \bar{C} \rightarrow \bar{C}^{\prime}$ be the map induced on $\bar{C}=C / \Lambda_{0, \text { nov }}^{+(0)} C$ and $\bar{C}^{\prime}=C^{\prime} / \Lambda_{0, \text { nov }}^{+(0)} C^{\prime}$, respectively. The induced map $\bar{\varphi}$ lifts to $\varphi_{0}=\bar{\varphi} \otimes 1: C \rightarrow C^{\prime}$.

## DEFINITION A. 5 ([FO+1, DEFINITION 6.3.26])

Under the situation above, we say that $\varphi: C \rightarrow C^{\prime}$ satisfies a gapped condition, or is a gapped cochain map, if there exists $\lambda^{\prime \prime}$ such that

$$
\left(\varphi-\varphi_{0}\right)\left(F^{\lambda} \widehat{C}\right) \subset F^{\lambda+\lambda^{\prime \prime}} \widehat{C}
$$

Using these definitions, $[\mathrm{FO}+1]$ proves the following.
THEOREM A. 7 ([FO +1 , THEOREM 6.3.28])
If $C$ is finite, then there exists $r_{0}$ such that

$$
E_{r_{0}}^{p, q}(\widehat{C}) \cong E_{r_{0}+1}^{p, q}(\widehat{C}) \cong \cdots \cong E_{\infty}^{p, q}(\widehat{C}) \cong F^{q} H^{p}(\widehat{C}, \delta) / F^{q+1} H^{p}(\widehat{C}, \delta)
$$

as $\Lambda^{(0)}\left(\lambda_{0}\right)=\Lambda_{0, \text { nov }}^{(0)} / F^{\lambda_{0}} \Lambda_{0, \text { nov }}^{(0)}$-modules.
We summarize the above discussion into the following vanishing result which is crucial in our spectral sequence arguments.

## COROLLARY A. 8

Let $C$ be a finite d.g.c.f.z. If $H(\bar{C}, \bar{\delta})=0$, then $H(C, \delta)=0$.

## A.3. Products

In this subsection, we recall the description of the deformed products $\mathfrak{m}_{k}^{\vec{b}}$ from $[\mathrm{FO}+2]$. We refer to $[\mathrm{Fu}]$ and $[\mathrm{FO}+2]$ for the relevant proofs of the statements we make without proofs here.

Let $\mathfrak{L}=\left(L_{0}, L_{1}, \ldots, L_{k}\right)$ be a chain of compact Lagrangian submanifolds in $(M, \omega)$ that intersect pairwise transversely without triple intersections.

Let $\vec{z}=\left(z_{0 k}, z_{k(k-1)}, \ldots, z_{10}\right)$ be a set of distinct points on $\partial D^{2}=\{z \in \mathbb{C} \mid$ $|z|=1\}$. We assume that they respect the counterclockwise cyclic order of $\partial D^{2}$. The group $\operatorname{PSL}(2 ; \mathbb{R}) \cong \operatorname{Aut}\left(D^{2}\right)$ acts on the set in an obvious way. We denote by $\mathcal{M}_{k+1}^{\text {main, }}$ the set of $\operatorname{PSL}(2 ; \mathbb{R})$-orbits of $\left(D^{2}, \vec{z}\right)$.

In this subsection, we consider only the case $k \geq 2$ since the case $k=1$ is already discussed Section A.2. In this case there is no automorphism on the domain $\left(D^{2}, \vec{z}\right)$; that is, $\operatorname{PSL}(2 ; \mathbb{R})$ acts freely on the set of such $\left(D^{2}, \vec{z}\right)$ 's.

Let $p_{j(j-1)} \in L_{j} \cap L_{j-1}(j=0, \ldots, k)$, be a set of intersection points.
We consider the pair $(w ; \vec{z})$, where $w: D^{2} \rightarrow M$ is a pseudoholomorphic map that satisfies the boundary condition

$$
\begin{equation*}
w\left(\overline{z_{j(j-1)} z_{(j+1) j}}\right) \subset L_{j} \tag{A.2a}
\end{equation*}
$$

We denote by $\widetilde{\mathcal{M}}^{\circ}(\mathfrak{L}, \vec{p})$ the set of such $\left(\left(D^{2}, \vec{z}\right), w\right)$.
We identify two elements $\left(\left(D^{2}, \vec{z}\right), w\right),\left(\left(D^{2}, \vec{z}^{\prime}\right), w^{\prime}\right)$ if there exists $\psi \in \operatorname{PSL}(2 ;$ $\mathbb{R})$ such that $w \circ \psi=w^{\prime}$ and $\psi\left(z_{j(j-1)}^{\prime}\right)=z_{j(j-1)}$. Let $\mathcal{M}^{\circ}(\mathfrak{L}, \vec{p})$ be the set of equivalence classes. We compactify it by including the configurations with disc or sphere bubbles attached and denote it by $\mathcal{M}(\mathfrak{L}, \vec{p})$. Its element is denoted by $((\Sigma, \vec{z}), w)$, where $\Sigma$ is a genus zero bordered Riemann surface with one boundary component, $\vec{z}$ are boundary marked points, and $w:(\Sigma, \partial \Sigma) \rightarrow(M, L)$ is a bordered stable map.

We can decompose $\mathcal{M}(\mathfrak{L}, \vec{p})$ according to the homotopy class $B \in \pi_{2}(\mathfrak{L}, \vec{p})$ of continuous maps satisfying (A.2a), (A.2b) into the union

$$
\mathcal{M}(\mathfrak{L}, \vec{p})=\bigcup_{B \in \pi_{2}(\mathfrak{L} ; \vec{p})} \mathcal{M}(\mathfrak{L}, \vec{p} ; B) .
$$

In the case when we fix an anchor $\gamma_{i}$ to each $L_{i}$ and put $\mathcal{E}=\left(\left(L_{0}, \gamma_{0}\right), \ldots\right.$, $\left(L_{k}, \gamma_{k}\right)$ ), we consider only admissible classes $B$ and put

$$
\mathcal{M}(\mathcal{E}, \vec{p})=\bigcup_{B \in \pi_{2}^{a d}(\mathcal{E} ; \vec{p})} \mathcal{M}(\mathcal{E}, \vec{p} ; B) .
$$

## THEOREM A. 9

Let $\mathfrak{L}=\left(L_{0}, \ldots, L_{k}\right)$ be a chain of Lagrangian submanifolds, and let $B \in \pi_{2}(\mathfrak{L} ; \vec{p})$. Then $\mathcal{M}(\mathfrak{L}, \vec{p} ; B)$ has an oriented Kuranishi structure (with boundary and corners). Its (virtual) dimension satisfies

$$
\begin{equation*}
\operatorname{dim} \mathcal{M}(\mathfrak{L}, \vec{p} ; B)=\mu(\mathfrak{L}, \vec{p} ; B)+n+k-2 \tag{A.3}
\end{equation*}
$$

where $\mu(\mathfrak{L}, \vec{p} ; B)$ is the polygonal Maslov index of $B$.

We next take graded anchors $\left(\gamma_{i}, \lambda_{i}\right)$ to each $L_{i}$. We assume that $B$ is admissible and write $B=\left[w_{01}^{-}\right] \#\left[w_{12}^{-}\right] \# \cdots \#\left[w_{k 0}^{-}\right]$as in Definition 3.5. We put $w_{(i+1) i}^{+}(s, t)=$ $w_{i(i+1)}^{-}(1-s, t)$. We also put $w_{k 0}^{+}(s, t)=w_{0 k}^{+}(s, 1-t)\left(\left[w_{k 0}^{+}\right] \in \pi_{1}\left(\ell_{k 0} ; p_{k 0}\right)\right)$. We also put $\lambda_{k 0}(t)=\lambda_{0 k}(1-t)$.

LEMMA A. 10 ([FO+2, LEMMA 8.11])
If $\operatorname{dim} \mathcal{M}(\mathfrak{L}, \vec{p} ; B)=0$, then we have

$$
\begin{equation*}
\left(\mu\left(\left[p_{k 0}, w_{k 0}^{+}\right] ; \lambda_{0 k}\right)-1\right)=1+\sum_{i=1}^{k}\left(\mu\left(\left[p_{i(i-1)}, w_{i(i-1)}^{+}\right] ; \lambda_{(i-1) i}\right)-1\right) . \tag{A.4}
\end{equation*}
$$

Using the case $\operatorname{dim} \mathcal{M}(\mathfrak{L}, \vec{p} ; B)=0$, we define the $k$-linear operator

$$
\begin{aligned}
\mathfrak{m}_{k} & : C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{k-1}, \gamma_{k-1}\right)\right) \otimes \cdots \otimes C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right) \\
& \rightarrow C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{0}, \gamma_{0}\right)\right)
\end{aligned}
$$

as follows:

$$
\begin{align*}
& \mathfrak{m}_{k}\left(\left[p_{k(k-1)}, w_{k(k-1)}^{+}\right],\left[p_{(k-1)(k-2)}, w_{(k-1)(k-2)}^{+}\right], \ldots,\left[p_{10}, w_{10}^{+}\right]\right)  \tag{A.5}\\
& \quad=\sum \#\left(\mathcal{M}_{k+1}(\mathfrak{L} ; \vec{p} ; B)\right)\left[p_{k 0}, w_{k 0}^{+}\right] .
\end{align*}
$$

Here the sum is over the basis $\left[p_{k 0}, w_{k 0}^{+}\right]$of $C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{0}, \gamma_{0}\right)\right)$, where $\vec{p}=$ $\left(p_{0 k}, p_{k(k-1)}, \ldots, p_{10}\right), B$ is as in Definition 3.5, and $w_{(i+1) i}^{+}(s, t)=w_{i(i+1)}^{-}(1-s, t)$.

Formula (A.4) implies that $\mathfrak{m}_{k}$ above has degree 1.
In general, the operator $\mathfrak{m}_{k}$ above does not satisfy the $A_{\infty}$-relation for the same reason as that of the case of boundary operators (see Section 8.1). We need to use bounding cochains $b_{i}$ of $L_{i}$ to deform $\mathfrak{m}_{k}$ in the same way as in the case of $A_{\infty}$-bimodules (see Section 8.3), whose explanation is now in order.

Let $m_{0}, \ldots, m_{k} \in \mathbb{Z}_{\geq 0}$ and $\mathcal{M}_{m_{0}, \ldots, m_{k}}(\mathfrak{L}, \vec{p} ; B)$ be the moduli space obtained from the set of $\left.\left(\left(D^{2}, \vec{z}\right),\left(\vec{z}^{(0)}, \ldots, \vec{z}^{(k)}\right), w\right)\right)$ by taking the quotient by $\operatorname{PSL}(2, \mathbb{R})-$ action and then by taking the stable map compactification as before. Here $z^{(i)}=$ $\left(z_{1}^{(i)}, \ldots, z_{k_{i}}^{(i)}\right)$ and $z_{j}^{(i)} \in \overline{z_{(i+1) i} z_{i(i-1)}}$ such that $z_{(i+1) i}, z_{1}^{(i)}, \ldots, z_{k_{i}}^{(i)}, z_{i(i-1)}$ respects the counterclockwise cyclic ordering;

$$
\left(\left(D^{2}, \vec{z}\right),\left(\vec{z}^{(0)}, \ldots, \vec{z}^{(k)}\right), w\right) \mapsto\left(w\left(z_{1}^{(0)}\right), \ldots, w\left(z_{m_{k}}^{(k)}\right)\right)
$$

induces an evaluation map

$$
\mathrm{ev}=\left(\mathrm{ev}^{(0)}, \ldots, \mathrm{ev}^{(k)}\right): \mathcal{M}_{m_{0}, \ldots, m_{k}}(\mathfrak{L}, \vec{p} ; B) \rightarrow \prod_{i=0}^{k} L_{i}^{m_{i}}
$$

Let $P_{j}^{(i)}$ be smooth singular chains of $L_{i}$, and put

$$
\vec{P}^{(i)}=\left(P_{1}^{(i)}, \ldots, P_{m_{i}}^{(i)}\right), \quad \overrightarrow{\vec{P}}=\left(\vec{P}^{(0)}, \ldots, \vec{P}^{(k)}\right)
$$

We then take the fiber product to obtain

$$
\mathcal{M}_{m_{0}, \ldots, m_{k}}(\mathfrak{L}, \vec{p} ; \vec{P} ; B)=\mathcal{M}_{m_{0}, \ldots, m_{k}}(\mathfrak{L}, \vec{p} ; B) \times_{\mathrm{ev}} \overrightarrow{\vec{P}}
$$

We use this to define

$$
\begin{aligned}
& \mathfrak{m}_{k ; m_{0}, \ldots, m_{k}}: B_{m_{k}}\left(C F\left(L_{k}\right)\right) \otimes C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{k-1}, \gamma_{k-1}\right)\right) \\
& \quad \otimes \cdots \otimes C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right) \otimes B_{m_{0}}\left(C F\left(L_{0}\right)\right) \rightarrow C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{0}, \gamma_{0}\right)\right)
\end{aligned}
$$

by

$$
\begin{aligned}
& \mathfrak{m}_{k ; m_{0}, \ldots, m_{k}}\left(\vec{P}^{(k)},\left[p_{k(k-1)}, w_{k(k-1)}^{+}\right], \ldots,\left[p_{10}, w_{10}^{+}\right], \vec{P}^{(0)}\right) \\
& \quad=\sum \#\left(\mathcal{M}_{k+1}(\mathfrak{L} ; \vec{p} ; \vec{P} ; B)\right)\left[p_{k 0}, w_{k 0}\right] .
\end{aligned}
$$

Finally, for each given $b_{i} \in C F\left(L_{i}\right)[1]^{0}\left(b_{i} \equiv 0 \bmod \Lambda_{+}\right), \vec{b}=\left(b_{0}, \ldots, b_{k}\right)$, and $x_{i} \in C F\left(\left(L_{i}, \gamma_{i}\right),\left(L_{i-1}, \gamma_{i-1}\right)\right)$, we put

$$
\begin{equation*}
\mathfrak{m}_{k}^{\vec{b}}\left(x_{k}, \ldots, x_{1}\right)=\sum_{m_{0}, \ldots, m_{k}} \mathfrak{m}_{k ; m_{0}, \ldots, m_{k}}\left(b_{k}^{m_{k}}, x_{k}, b_{k-1}^{m_{k-1}}, \ldots, x_{1}, b_{0}^{m_{0}}\right) \tag{A.6}
\end{equation*}
$$

THEOREM A. 11
If $b_{i}$ satisfies the Maurer-Cartan equation (8.10), then $\mathfrak{m}_{k}^{\vec{b}}$ in (A.6) satisfies the $A_{\infty}$-relation

$$
\begin{equation*}
\sum_{k_{1}, k_{2}, i}(-1)^{*} \mathfrak{m}_{k_{1}}^{\vec{b}}\left(x_{k}, \ldots, \mathfrak{m}_{k_{2}}^{\vec{b}}\left(x_{k-i-1}, \ldots, x_{k-i-k_{2}}\right), \ldots, x_{1}\right)=0 \tag{A.7}
\end{equation*}
$$

where we take the sum over $k_{1}+k_{2}=k+1, i=-1, \ldots, k-k_{2}$. (We write $\mathfrak{m}_{k}$ in place of $\mathfrak{m}_{k}^{\vec{b}}$ in (A.7).) The sign $*$ is $*=i+\operatorname{deg} x_{k}+\cdots+\operatorname{deg} x_{k-i}$.

We summarize the above discussion as follows.

## THEOREM A. 12

We can associate a filtered $A_{\infty}$-category to a symplectic manifold $(M, \omega)$ such that
(1) its object is $(\mathcal{L}, b, s p)$, where $\mathcal{L}=(L, \gamma, \lambda)$ is a graded anchored Lagrangian submanifold, $[b] \in \mathcal{M}(C F(L))$ is a bounding cochain, and sp is a spin structure of $L$;
(2) the set of morphisms is $\operatorname{CF}\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)$;
(3) $\mathfrak{m}_{k}^{\vec{b}}$ are the operations defined in (A.6).

## REMARK A. 6

Here we spell out the choice of orientations; $o_{p}$ of Index $\bar{\partial}_{\lambda_{p}}$ is included. This choice in fact does not affect the module structure $C F\left(\left(L_{1}, \gamma_{1}\right),\left(L_{0}, \gamma_{0}\right)\right)$ up to isomorphism: if we take an alternative choice $o_{p}^{\prime}$ at $p$, then all the signs appearing in the operations $\mathfrak{m}_{k}$ that involve $[p, w]$ for some $w$ are reversed. Therefore $[p, w] \mapsto-[p, w]$ gives the required isomorphism.

The operations $\mathfrak{m}_{k}$ are compatible with the filtration. Namely, we have the following.

PROPOSITION A. 13
If $x_{i} \in F^{\lambda_{i}} C F\left(\left(L_{i}, \gamma_{i}\right),\left(L_{i-1}, \gamma_{i-1}\right)\right)$, then

$$
\mathfrak{m}_{k}^{\vec{b}}\left(x_{k}, \ldots, x_{1}\right) \in F^{\lambda} C F\left(\left(L_{k}, \gamma_{k}\right),\left(L_{0}, \gamma_{0}\right)\right),
$$

where $\lambda=\sum_{i=1}^{k} \lambda_{i}$.

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