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## Note on Nonlinear Differential Equation of Catalysis<sup>1)</sup>

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1. Mr. Yoshiyuki Suehiro, a chemical engineer, has consulted the author about the solutions of his differential equation of catalysis by spherical tablets,

(1) 
$$\begin{cases} (m-1)\left(Da\frac{dy}{dx}-wa\frac{y}{t}\right)=wa,\\ \frac{dwa}{dx}dx=cy\frac{adx}{\gamma}; \end{cases}$$

or for the case m=2, eliminating w in (1), we have

(2) 
$$\frac{d}{dx}\left(\frac{Dat}{y+t},\frac{dy}{dx}\right) = -\frac{ca}{\gamma}y;$$

or fully written,

$$(3) \qquad \frac{d^2 y}{dx^2} + \frac{2}{x} \frac{dy}{dx} - \frac{1}{y+t} \left(\frac{dy}{dx}\right)^2 - n^2 \frac{y(y+t)}{t} = 0,$$

where  $n = \sqrt{C/Dr}$ . In these equations, x means the distance of any point of the tablet from its centre; y is the concentration of the reacting substance at x; w is the quantity of mass flow of the reacted substance through the spherical surface of radius x in the tablet; a is the total area occupied by the pores on the spherical surface of radius x, and hence proportional to  $x^2$ , while the remainings are chemical constants, positive. Solutions for x > 0, satisfying the conditions  $\frac{dy}{dx} = 0$  at x = 0, are required.

2. We may suppose t=1 by writing y instead of ty (also n=1 by writing x instead of nx).

<sup>1)</sup> Read before the last autumn meeting of Japanese Mathematical Society held in Kyoto.

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Putting t=1, and

(4) 
$$e^z = 1 + y$$
, we have  $\frac{d\left(x^2 \cdot \frac{dz}{dx}\right)}{dx} = n^2 x^2 y$ .

Further putting

(5) 
$$\frac{\chi}{x} = z$$
, we have  $\frac{d^2\chi}{dx^2} = n^2xy$ .

The differential equation is very near to Emden's one<sup>2)</sup>

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = \theta^n.$$

By its knowledge we may prove the following proposition:

The solution y(x) of (4) which is continuous in an interval in which 0 < x and  $y(+0) = a_0$  finite and determinate and greater than -1, is normal to y-axis, i.e., y'(+0)=0. We may define y(0) $\equiv y(+0)$ , then y(x) is continuous for  $0 \le x$  and  $y(0) = a_0$ .

Since  $z(x) = \log(1+y(x))$  is also continuous and y(0) > -1,  $z(0) (\equiv z(+0))$  is also finite and determinate. Hence we have

$$\lim_{x \to +0} \chi(x) = \lim_{x \to +0} xz(x) = 0, \quad \text{i. e.,} \quad \chi(0) (\equiv \chi(+0)) = 0.$$

Nextly we have

$$\frac{\left.\frac{dz}{dx}\right|_{x=+0} = \lim_{x \to +0} \frac{(\chi/x) - z(0)}{x} = \lim_{x \to +0} \frac{\chi - xz(0)}{x^2}$$
$$= \lim_{x \to +0} \frac{\frac{d\chi}{dx} - z(0)}{2x},$$

provided the last limit is determinate. On the other hand

$$\frac{d\chi}{dx}\Big|_{x=+0} = \lim_{x\to+0} \frac{\chi(x)-\chi(0)}{x} = \lim_{x\to+0} \frac{\chi}{x} = z(0).$$

From (5) we have

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$$\frac{d\chi}{dx} = z(0) + n^2 \int_0^x xy \, dx \, .$$

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<sup>2)</sup> Emden, Gaskugeln (1907); Chandrasekhar, Steller Structure (1938); by the author, Emden's differential Equation (in japanese), Sankaidō & Co., (1945).

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$$\lim_{x \to +0} \frac{\frac{d\chi}{dx} - z(0)}{2x} = \lim_{x \to +0} \frac{n^2 x y(x)}{2} = 0,$$

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since by hypothesis  $y(0) = a_0$  is finite. Hence we may conclude that

$$\frac{dz}{dx}\Big|_{x=+0}=0;$$

hence by (4), we have

$$\left. \frac{dy}{dx} \right|_{x=+0} = 0. \qquad \text{Q. E. D.}$$

In the following for simplicity we take  $a_0 \ge 0$ , since y is the concentration, not negative.

**3.** Our equation has an integral near  $0 \le x$  of the form:

 $y(x) = a_0 + a_2 x^2 + a_4 x^4 + \cdots$ (6)

where  $a_0 \ge 0$  is given, while the other coefficients can be found successively from the equation :

 $2a_{3}x + 4a_{4}x^{3} + \cdots$ (7)

$$=n^{2}\left(\frac{a_{0}}{3}x+\frac{a_{2}}{5}x^{3}+\cdots\right)(1+a_{0}+a_{2}x^{2}+a_{4}x^{4}+\cdots).$$

This power-series is dominated by

$$\psi(x) = A_0 + A_2 x^2 + A_4 x^4 + \cdots,$$

where  $A_0 = a_0$  and

(8) 
$$\frac{d\psi}{dx} = \frac{n^2}{3} x \psi (1+\psi),$$

which gives

$$\psi(x) = \frac{a_0}{1+a_0} e^{n^2 x^2/6} \left| \left( 1 - \frac{a_0}{1+a_0} e^{n^2 x^2/6} \right) \right|.$$

Hence the convergence abscissa  $\rho$  of y(x) is

$$\rho \geq \frac{\sqrt{6}}{n} \sqrt{\log\left(1 + \frac{1}{a_0}\right)}.$$

Our solution y(x) with the initial condition  $y(0) = a_0$  is unique on the right of y-axis and it is a power-series of x and  $a_0$ , with positive

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coefficients, so that our solution, written  $y(x, a_0)$  increases with  $a_0$  for fixed x.

To obtain (7), integrating (4) we have

$$x^2 \frac{dz}{dx} = n^2 \int_0^x x^2 y \, dx ,$$
$$x^2 \frac{dy}{dx} = n^2 (1+y) \int_0^x x^2 y \, dx$$

or

Putting (6), we have (7). From (7) we may find all the coefficients except  $a_0$ ; they are positive for  $a_0 > 0$ .

To find  $\psi(x)$ , we consider instead of (7), the relation:

$$2A_2x + 4A_4x^3 + \cdots$$
  
=  $\frac{n^2}{3}(A_0x + A_2x^3 + \cdots)(1 + A_0 + A_2x^2 + A_4x^4 + \cdots)$ 

We see easily that for  $a_0 = A_0 (>0)$ , we have  $a_2 = A_2$ ,  $a_4 < A_4$ ,... and that  $\psi(x)$  satisfies (8).

It is evident that the solution y(x) with the initial condition  $y(0)=a_0>0$  is unique and it is continuous with respect to  $a_0$ ; the first quadrant of the coordinate plane is swept by our integrals with  $0 \le a_0 < \infty$ .

We remark that above considerations may easily be extended for the cases  $a_0 > -1$  and x < 0.

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