# Square integrable normal differentials on Riemann surfaces* 

Dedicated to Professor Y. Akizuki on his 60th birthday

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Introduction. In this article we shall study some properties on periods of abelian differentials square integrable on open Riemann surfaces. In §1 we shall show the relation of canonical differentials defined in [4] and Ahlfors' reproducing differentials [2], which play a fundamental role in the theory of abelian differentials. In $\S 2$ we shall prove a general periods theorem, namely, let $\Gamma_{0}\left(\subset \Gamma_{a}\right)$ be a Hilbert space of square integrable analytic differentials and $\gamma$ be a set of cycles for which corresponding reproducing differentials in $\mathrm{I}_{0}$ are linearly independent, then Theorem 1 gives a necessary and sufficient condition in order that given complex numbers should be the periods of a differential of $\Gamma_{0}$ along $\gamma$. As its corollary we know the existence of three kinds of normal differentials, in particular, normal differentials of the first kind, for which in $\S 3$ some properties on corresponding Riemann matrix are shown. Theorem 3 is partly identical with Oikawa's results in his unpublished paper studying normal differentials under Ahlfors' and Accola's method.

## $\S 1$

Let $R$ be an arbitrary Riemann surface of genus $g(0 \leq g \leq \infty)$

[^0]and $\left\{A_{n}, B_{n}\right\}_{n=1, \cdots, g}$ and $\left\{C_{i}\right\}_{v=1, \cdots, p}(0 \leq p \leq \infty)$ be a canonical homology basis on $R$ such that (i) any cycle in $R$ is homologous to a finite sum $\sum\left(p_{n} A_{n}+q_{n} B_{n}\right)+\sum r_{\nu} C_{v}$ where $p_{n}, q_{n}$ and $r_{\nu}$ are integers; (ii) the intersection numbers are characterized by $A_{i} \times B_{j}=\delta_{i j}$, $A_{i} \times A_{j}=B_{i} \times B_{j}=0$ for $i, j=1, \cdots, g$; (iii) any dividing cycle is homologous to a finite sum $\sum r_{\nu} C_{\nu}$.

Let $\Gamma_{a}$ be the Hilbert space of analytic differentials square integrable over $R$ where the scalar product and norm are defined by

$$
\begin{equation*}
(\phi, \psi r)=\frac{i}{2} \int_{R} \phi \wedge \bar{\psi}, \quad\|\phi\|^{2}=(\phi, \phi) . \tag{1}
\end{equation*}
$$

For $\phi=d u+i^{*} d u, \psi=d u^{\prime}+i^{*} d u^{\prime}$ one can write

$$
(\phi, \psi)=D_{R}\left(d u, d u^{\prime}\right)-i D_{R}\left(d u, * d u^{\prime}\right),
$$

where $D_{R}\left(d u, d u^{\prime}\right)=\int_{R} d u \wedge * d u^{\prime}=\int_{R}\left(u_{x} u_{x}^{\prime}+u_{y} u_{y}^{\prime}\right) d x d y$ stands for the mixed Dirichlet integral. Let $\mathrm{I}_{a e}$ and $\mathrm{I}_{a s e}$ denote the subspaces of $\Gamma_{a}$ consisting of exact and semi-exact differentials respectively. Now Virtanen's decomposition theorem can be expressed ([4]) as

$$
\mathrm{\Gamma}_{a}=\mathrm{\Gamma}_{a s e} \oplus \mathrm{I}_{C}, \quad \mathrm{\Gamma}_{a s e}=\mathrm{\Gamma}_{a e} \oplus \mathrm{\Gamma}_{A B}^{\prime}
$$

where $\mathrm{I}_{c}$, orthogonal complement of $\mathrm{I}_{\text {ase }}$, is spanned by analytic differentials $\left\{\mathcal{P}_{c_{\nu}}\right\}$ derived from generalized harmonic measures associated with dividing cycles $\left\{C_{\nu}\right\}$ and $1_{A B}^{\prime}$ is spanned by semiexact canonical differentials $\left\{\mathscr{\rho}_{A_{n}}, \mathscr{P}_{B_{n}}\right\}$ such that $\operatorname{Re} \mathscr{\rho}_{A_{n}}, \operatorname{Re} \mathscr{P}_{B_{n}}$ have no periods except along $B_{n}$ resp. $A_{n}$ where

$$
\begin{equation*}
\operatorname{Re} \int_{B_{n}} \mathcal{P}_{A_{n}}=-\operatorname{Re} \int_{A_{n}} \mathcal{P}_{B_{n}}=1 \tag{2}
\end{equation*}
$$

Proposition 1. For any $\phi \in \mathrm{I}_{\text {ase }}$ we have

$$
\begin{equation*}
\left(\phi, \mathcal{P}_{A_{n}}\right)=i \int_{A_{n}} \phi, \quad\left(\phi, \mathscr{P}_{B_{n}}\right)=i \int_{B_{n}} \phi . \tag{3}
\end{equation*}
$$

Proof. Take an exhaustion of compact domains $\left\{R_{m}\right\}$ of $R$ such that each component of the boundary $\partial R_{m}$ consists of a dividing analytic curve. Then for $\phi=d u+i^{*} d u, \mathcal{P}_{A_{n}}=d u_{A_{n}}+i^{*} d u_{A_{n}}$ and $R_{m}$ containing the cycle $A_{n}$ we have by the period relation

$$
D_{R m}\left(d u_{A_{n}}, d u\right)=-\int_{A_{n}} * d u+\int_{\partial R m} u_{A_{n}}^{*} d u
$$

Since $u_{A_{n}}$ is single-valued outside $R_{m}$, moreover canonical, it follows that $\int_{\partial R m} u_{A_{n}}{ }^{*} d u \rightarrow 0$ for $m \rightarrow \infty$ (cf. Lemma 4 [4]). Hence

$$
D_{R}\left(d u_{A_{n}}, d u\right)=-\int_{A_{n}} * d u
$$

Analogously we have

$$
\begin{equation*}
D_{R}\left(d u_{A_{n}}, * d u\right)=-D_{R}\left(d u, * d u_{A_{n}}\right)=\int_{A_{n}} d u \tag{4}
\end{equation*}
$$

By ( $1^{\prime}$ ) we get therefore the first formula of (3), similarly the second one. Q. E. D.

We recall [4] that for any $\phi \in \Gamma_{a}$ and harmonic differential $d U$ with finite Dirichlet integral

$$
\begin{equation*}
\left(\phi, \rho_{C_{\nu}}\right)=-i \int_{C_{\nu}} \phi, \quad D_{R}\left(d U, * d \omega_{C_{\nu}}\right)=\int_{C_{\nu}} d U \tag{5}
\end{equation*}
$$

where $\mathcal{P}_{C_{\nu}}=d \omega_{C_{\nu}}+i^{*} d \omega_{C_{\nu}}$. (3), (4), (5) imply that our canonical differentials $\varphi_{A_{n}}, \mathscr{P}_{B_{n}}$ and $\mathcal{P}_{C_{\nu}}$ are identical, except constant factors, with reproducing analytic differentials in $\mathrm{I}_{a s e}$ resp. $\mathrm{I}_{a}{ }_{a}$ in the Ahlfors sense [2] and $d u_{A_{n}}, d u_{B_{n}}$ and $d \omega_{C_{\nu}}$ are also reproducing (real) harmonic differentials in $1_{h s e}^{*}$ resp. $\Gamma_{h}$ under the Dirichlet norm.

## § 2. General Periods Theorem

Let $I_{0}^{\prime} \subset \Gamma_{a}$ be any Hilbert space. Now let $P_{\gamma}=P_{\gamma}\left\{\Gamma_{0}\right\}$ be a linear mapping of $1_{0}$ into a complex vector space $E$ such that

$$
P_{\gamma}: \phi \rightarrow P_{\gamma} \phi=\left(\int_{\gamma_{1}} \phi, \int_{\gamma_{2}} \phi, \cdots\right), \quad \gamma=\left\{\gamma_{n}\right\}_{n=1,2, \cdots}
$$

We note that the kernel of $P_{\gamma}$ is a Hilbert space $\subset \Gamma_{0}$ while the image of $P_{\gamma}$, say $P_{\gamma}\left(\Gamma_{0}\right)$, is in general not an $l_{2}$-(Hilbert) space, namely we can easily find such an example among hyperelliptic surfaces.

We say a set of cycles $\gamma=\left\{\gamma_{j}\right\}_{j=1,2}, \ldots$ be admissible for $\Gamma_{0}$ if the corresponding reproducing differentials $\sigma_{\gamma_{j}}$ in $\Gamma_{0}$ defined by

$$
\left(\phi, \sigma_{\gamma_{j}}\right)=\int_{\gamma_{j}} \phi \quad \text { for all } \phi \in \mathrm{\Gamma}_{0}
$$

are linearly independent (for any finite pair). $\sigma_{\gamma_{j}} \in \Gamma_{0}$ exist, because $\int_{\gamma_{j}} \phi$ is a bounded linear functional on $\Gamma_{0}$. Let

$$
\Gamma_{0}^{\prime}=\left[\sigma_{\gamma_{j}}\right] \subset \Gamma_{0}
$$

be a subspace of $\Gamma_{0}$ spanned by $\sigma_{\gamma_{j}}$, and $\left\{\sigma_{\gamma_{j}}^{\prime}\right\}$ be an orthonormal system constructed by the Schmidt method:

$$
\begin{equation*}
\sigma_{\gamma_{j}}^{\prime}=\sum_{k=1}^{j} s_{j k} \sigma_{\gamma_{k}} \quad(j=1,2, \cdots) . \tag{6}
\end{equation*}
$$

We call the triangular matrix $S_{\gamma}=S_{\gamma}\left\{\Gamma_{0}\right\}=\left(s_{j k}\right)$ with $s_{j k}=0(j<k)$ the matrix (mapping) associated with $\Gamma_{0}$ and $\gamma$.

Proposition 2. $S_{\gamma}\left\{\mathrm{I}_{0}\right\}=\left(s_{m m}\right)$ has the properties:
(i) $s_{n n}>0$;
(ii) $s_{m n}$ are real if $\left(\sigma_{\gamma_{j}}, \sigma_{\gamma_{k}}\right)$ are real for $j, k=1,2, \cdots$;
(iii) $\sum_{m}\left|s_{m n}\right|^{2}=\left\|\phi_{n}\right\|^{2}<\infty$ if there exists a $\phi_{n} \in \Gamma_{0}$ such that $P_{\gamma} \phi_{n}=e_{n}$ (unit vector in $E$ with $n$-th component one).
(i), (ii) are immediate consequences from the orthogonalization process and (iii) follows easily from the following Theorem 1.

Now $S_{\gamma}$ defines a linear mapping of $E$ into itself such that

$$
S_{\gamma}: \alpha \rightarrow S_{\gamma}(\alpha)=\left(\bar{s}_{11} \alpha_{1}, \bar{s}_{21} \alpha_{1}+\bar{s}_{22} \alpha_{2} \cdots\right),
$$

that is, $S_{\gamma}(\alpha)$ is the product of matrix ( $\bar{s}_{j k}$ ) (bar stands for complex conjugate) and the transposed vector ${ }^{t} \alpha$ of $\alpha=\left(\alpha_{1}, \alpha_{2}, \cdots\right) \in E . \quad S_{\gamma}$ is one-to-one by (i).

Theorem 1. For given $\Gamma_{0}$ and $\gamma=\left\{\gamma_{j}\right\}$ admissible for $\Gamma_{0}$, (I) $S_{\gamma} \circ P_{\gamma}$ is a continuous linear mapping of $\Gamma_{0}$ onto the complex $l_{2}$ space ${ }^{1)}$ such that

$$
\left\|S_{\gamma} \circ P_{\gamma}(\phi)\right\|_{l_{2}} \leq\|\phi\|, \quad \phi \in \mathrm{I}_{0}^{\prime} .
$$

(II) $S_{\gamma} \circ P_{\gamma}$ is one-to-one on $\mathrm{I}_{0}^{\prime}$ and $P_{\gamma}\left(\mathrm{I}_{0}\right)=P_{\gamma}\left(\Gamma_{0}^{\prime}\right)=S_{\gamma}^{-1}\left(l_{2}\right)$.

Proof. Let $\phi \in \mathrm{I}_{0}^{\prime}$ and $\phi^{\prime}$, be the orthogonal projection of $\phi$ onto $\Gamma_{0}^{\prime}$. Then $\phi^{\prime}$ can be written as $\phi^{\prime}=\sum \beta_{n} \sigma_{\gamma_{n}}^{\prime}$ with the Fourier coefficients $\beta_{n}$, given by

1) This is just a complex vector space of dimenslon $q$ if $r$ consists of $q(<\infty)$ cycles,

$$
\beta_{n}=\left(\phi^{\prime}, \sigma_{\gamma_{n}}^{\prime}\right)=\left(\phi, \sigma_{\gamma_{n}}^{\prime}\right)=\sum_{k=1}^{n} \bar{s}_{n k}\left(\phi, \sigma_{\gamma_{k}}\right)=\sum_{k=1}^{n} \bar{s}_{n k} \int_{\gamma_{k}} \phi .
$$

Hence by Parseval's equality we have

$$
\sum\left|\beta_{n}\right|^{2}=\left\|S_{\gamma} \circ P_{\gamma}(\phi)\right\|_{i_{2}}^{2}=\left\|\phi^{\prime}\right\|^{2} \leq\|\phi\|^{2} .
$$

which shows (I) except "onto" and $P_{\gamma}\left(\mathrm{I}_{0}^{\prime}\right) \subset S_{\gamma}^{-1}\left(l_{2}\right)=\left\{\alpha ; S_{\gamma}(\alpha) \in l_{2}\right\}$. Now let $\beta^{*}=\left(\beta_{1}^{*}, \beta_{2}^{*} \cdots\right) \in l_{2}$, then there is an $\alpha$ such that $S_{\gamma}(\alpha)=\beta^{*}$;

$$
\begin{equation*}
\beta_{n}^{*}=\sum_{k=1}^{n} \bar{s}_{n k} \alpha_{k} \tag{7}
\end{equation*}
$$

Since $\sum\left|\beta_{n}^{*}\right|^{2}<\infty$,

$$
\begin{gather*}
\phi^{*} \equiv \sum \beta_{n}^{*} \sigma_{\gamma_{n}}^{\prime} \in \mathbf{I}_{0}^{\prime} \subset \mathbf{I}_{0}^{\prime},  \tag{8}\\
\beta_{n}^{*}=\left(\phi^{*}, \sigma_{\gamma_{n}}^{\prime}\right)=\sum_{k=1}^{n} \bar{s}_{n k} \int_{\gamma_{k}} \phi^{*} . \tag{9}
\end{gather*}
$$

Therefore by (7), (9) and (i) we have successively

$$
\int_{\gamma_{k}} \phi^{*}=\alpha_{k} \quad(k=1,2, \cdots), \text { i. e., } \alpha=P_{\gamma} \phi^{*},
$$

which implies $S_{\gamma}^{-1}\left(l_{2}\right) \subset P_{\gamma}\left(\mathrm{I}_{0}^{\prime}\right) \subset P_{\gamma}\left(\Gamma_{0}\right)$, hence they coincide. Since any differential of $\Gamma_{0}^{\prime}$ orthogonal to every $\sigma_{\gamma_{j}}$ is identically zero, $S_{\gamma} \circ P_{\gamma}$ is one-to-one on $\Gamma_{0}^{\prime}$.

Corollary. Suppose $s_{j k}$ are all, except finite numbers, real (or pure imaginary) and let $P_{\gamma} \phi=\alpha+i \beta, \phi \in \Gamma_{0}$ where $\alpha, \beta$ are real vectors then there exist $\phi_{1}, \phi_{2} \in \Gamma_{0}^{\prime}$ such that

$$
P_{\gamma} \phi_{1}=\alpha, \quad P_{\gamma} \phi_{2}=\beta .
$$

Here we refer to important special cases. For canonical basis $\left\{A_{n}, B_{n}\right\}$ and $\left\{C_{v}\right\}$ we know by (3), (5) that

$$
\begin{array}{ll}
\sigma_{A_{n}}=i \rho_{A_{n}}, \quad \sigma_{B_{n}}=i \varphi_{B_{n}} & \text { for } \mathrm{I}_{0}^{\prime}=\mathrm{\Gamma}_{a s e}  \tag{10}\\
\sigma_{C_{\nu}}=-i \varphi_{C_{\nu}} & \text { for } \mathrm{I}_{0}^{\prime}=\mathrm{I}_{a}^{\prime}
\end{array}
$$

and $\left(\sigma_{A_{j}}, \sigma_{A_{k}}\right),\left(\sigma_{B j}, \sigma_{B_{k}}\right)$ and $\left(\sigma_{C_{j}}, \sigma_{C_{k}}\right)$ are real, but $\left(\sigma_{A_{j}}, \sigma_{B_{k}}\right)$ are not necessarily real ; moreover, $\left\{C_{\nu}\right\}^{\prime}=\left\{C_{\nu} \in\left\{C_{\nu}\right\} ; \sigma_{C_{\nu}} \neq 0\right\}$ is admissible for $\mathrm{I}_{a}$, and $\left\{A_{n}\right\},\left\{B_{n}\right\}$ are admissible for $\Gamma_{a s e}$. In fact, let $\left\{\omega_{j}\right\}$ denotes any one of the systems $\left\{\sigma_{A j}\right\},\left\{\sigma_{B j}\right\}$ and $\left\{\sigma_{C j}\right\}$
$\left(C_{j} \in\left\{C_{\nu}\right\}^{\prime}\right)$, then $\left\{\omega_{j}\right\}$ are linearly independent in the real field, hence for finite number of real $x_{j}$ the quadratic form

$$
\sum_{i, j} x_{i} x_{j} D_{R}\left(R e \omega_{i}, \operatorname{Re} \omega_{j}\right)=D_{R}\left(\sum_{j} x_{j} \omega_{j}\right) \geq 0
$$

is positive definite, i.e.

$$
\operatorname{det}\left(D_{R}\left(\operatorname{Re} \omega_{i}, \operatorname{Re} \omega_{j}\right)\right)>0
$$

Noting that $\left(\omega_{j}, \omega_{j}\right)=D_{R}\left(R e \omega_{i}, R e \omega_{j}\right)$ (=real) we find easily that $\left\{\omega_{j}\right\}$ are linearly independent in the complex field.

Virtanen [7] observed first the case $\Gamma_{0}=\Gamma_{a}, \gamma=\left\{A_{n}\right\}$ on Riemann surfaces of parabolic type (where $\Gamma_{a}=\Gamma_{a s e}=\Gamma_{A B}$ ). For general surfaces, I [4] treated the case $\Gamma_{0}=\Gamma_{C}, \gamma=\left\{C_{\gamma}\right\}^{\prime}$ and Sainouchi did independently the case $\mathrm{I}_{0}^{\prime}=\Gamma_{a s e}, \gamma=\left\{A_{n}\right\}$ in his recent paper [6]. While in the case where $\gamma$ contains mixed basis it is generally difficult to decide whether $\gamma$ is admissible or not. For instance if $R$ is closed or parabolic surface of finite genus $g, \Gamma_{a}=\Gamma_{A B}=$ $\left[\sigma_{A_{n}}, \sigma_{B n}\right]_{n=1, \ldots, g}$ reduces $g$-dimensional. If $R$ is a bordered surface with $p$ contours, $\left\{A_{n}, B_{n}, C_{v}\right\}_{n=1, \cdots, g ; v=1, \cdots, p^{-1}}$ is admissible for $\Gamma_{0}=$ $\Gamma_{a S}=\Gamma_{A B} \oplus \Gamma_{C}$ and $P_{\gamma}$ gives a one-to-one mapping of $\Gamma_{0}=\Gamma_{0}^{\prime}{ }^{\prime}$ to ( $2 g+p-1$ )-dimensional vector space by (II) (Ahlfors [1]).

## § 3. Normal Differentials

Hereafter we fix the canonical basis $\left\{A_{n}, B_{n}\right\}$ once for all and define a Hilbert space $\Gamma_{A}^{0}$ by

$$
\Gamma_{A}^{0}=\left\{\phi \in \mathrm{I}_{a s e}^{\prime}, \int_{A_{n}} \phi=0, n=1,2, \cdots\right\}
$$

i. e., the kernel of $P_{\gamma}\left\{\mathrm{\Gamma}_{\text {ase }}\right\}, \gamma=\left\{A_{n}\right\}$. Let $\mathrm{I}_{A}$ be the orthogonal complement of $\Gamma_{A}^{0}$ in $\Gamma_{\text {ase }}$. Obviously

$$
\Gamma_{a e} \subset \Gamma_{A}^{0} \subset 1_{a s e}, \quad \Gamma_{A} \subset \Gamma_{A B}
$$

Every differential of $\Gamma_{A}$ is by definition uniquely determined by its $A$-periods. Moreover among differentials of $\Gamma_{\text {ase }}$ having the same $A$-periods with $\phi \in \Gamma_{\text {ase }}$ the orthogonal projection of $\phi$ onto $\Gamma_{A}$ has the minimum norm.

Proposition 3. $\Gamma_{A}=\left[\mathcal{P}_{A_{n}}\right]=\left[\sigma_{A_{n}}\right] \quad(n=1, \cdots, g)$,

This is an immediate consequence of Proposition 1.
Since for general surfaces $\Gamma_{A}$ does not necessarily coincide with $\Gamma_{A B}$, we must confine ourselves to the space $\Gamma_{A}$ to achieve the complex normalization about $A$-periods. Now we take $\Gamma_{0}=$ $\mathrm{I}_{\text {ase }}, \gamma=\left\{A_{n}\right\}$, then $\mathrm{I}_{0}^{\prime}=\mathrm{\Gamma}_{A}$ by Proposition 3 and $S_{\gamma}=\left(s_{m n}\right)$ is real. For $\mathscr{P}_{B_{n}} \in \Gamma_{0} \int_{A_{m}}-\mathscr{P}_{B_{n}}=\delta_{m n}+i()$ hence by Corollary of Theorem 1 we have the following theorem [8].

Theorem 2. There exist normal differentials $\omega_{A_{n}} \in \Gamma_{A}$ such that

$$
\int_{A_{m}} \omega_{A_{n}}=\delta_{m n} \quad(m, n=1, \cdots, g) .
$$

The $\omega_{A_{n}}$ can be written by (8) as

$$
\begin{equation*}
\omega_{A_{n}}=\sum_{m} s_{m n} \sigma_{A_{m}}^{\prime}=\sum_{m} s_{m n}\left(\sum_{l=1}^{m} s_{m l} \sigma_{A_{l}}\right), \tag{11}
\end{equation*}
$$

where $\left(s_{m n}\right)$ is a real matrix associated with $\mathrm{\Gamma}_{\text {ase }}$ and $\left\{A_{n}\right\}$ for which

$$
\begin{equation*}
\sum_{m}\left|s_{m n}\right|^{2}=\left\|\omega_{A_{n}}\right\|^{2}<\infty \tag{12}
\end{equation*}
$$

For the following purposes we note that

$$
\begin{equation*}
s_{m n}=-\operatorname{Re} b_{m n}, \quad \text { where } b_{m n}=-i \int_{B_{n}} \overline{\sigma_{A_{m}}^{\prime}} \tag{13}
\end{equation*}
$$

To see this, let $\mathscr{\varphi}_{B_{n}}^{0}$ be the orthogonal projection $\mathscr{P}_{B_{n}}$ to $\Gamma_{A}$, then it can be written as

$$
\begin{equation*}
\mathscr{P}_{B_{n}}^{0}=\sum b_{m n} \sigma_{A m}^{\prime}, \tag{14}
\end{equation*}
$$

where $\quad b_{m n}=\left(\mathscr{P}_{B_{n}}^{n}, \sigma_{A_{m}}^{\prime}\right)=\left(\mathcal{P}_{B_{n}}, \sigma_{A_{m}}^{\prime}\right)=-i \int_{B_{n}} \overline{\sigma_{A_{m}}^{\prime}}=-\int_{B_{n}} \sum_{k=1}^{m} s_{m k} \overline{\mathcal{P}}_{A_{k}}$, hence (13) holds. Now set

$$
\begin{array}{ll}
\int_{B j} \omega_{A k}=\sigma_{j k}+i \tau_{j k}, & (j, k=1, \cdots, g), \\
\sigma=\left(\sigma_{j k}\right), \quad \boldsymbol{\tau}=\left(\tau_{j k}\right) . &
\end{array}
$$

We say an (infinite) quadratic form $\sum \xi_{i} \xi_{j} a_{i j}$ is convergent if for given $\varepsilon>0$ there exists an integer $N$ such that $\left|\sum_{i, j=n}^{m} \xi_{i} \bar{\xi}_{j} a_{i j}\right|<\varepsilon$
for $m>n>N$. Now our Riemann matrix has the following properties.

Theorem 3. $1^{\circ}$ ) $\boldsymbol{\tau}$ is symmetric and positive definite; more precisely, for any finite or infinite numbers of complex numbers $\xi_{1}, \xi_{2}, \cdots$ not simultaneously zero

$$
\sum \xi_{i} \bar{\xi}_{j} \tau_{i j}>0
$$

provided that the quadratic form is convergent.
$2^{\circ}$ ) $\boldsymbol{\tau}={ }^{t} S_{\gamma} \cdot S_{\gamma}$ (multiplication in the obvious sense) where $S_{\gamma}$ is a triangular real matrix associated with $\Gamma_{\text {ase }}$ and $\left\{A_{n}\right\}$.
$3^{\circ}$ ) if $\Gamma_{A}^{0}=\mathrm{\Gamma}_{a e}$, equivalently $\Gamma_{A}=\Gamma_{A B}$, then $\sigma$ is symmetric.
Proof. From (10), (11)

$$
\begin{align*}
\tau_{j k}=\sum_{m} s_{m k} \int_{B j} I m \sigma_{A_{m}}^{\prime} & =\sum_{m} s_{m k} \sum_{l=1}^{m} s_{m l} \int_{B j} R e \mathcal{P}_{A_{l}}  \tag{15}\\
& =\sum_{m} s_{m k} s_{m j} .
\end{align*}
$$

The series on the right-hand side is absolutely convergent by (12) and Schwarz' inequality. This formula shows $2^{\circ}$ ) and the symmetry of $\tau$. On the other hand

$$
\left(\omega_{A j}, \omega_{A_{k}}\right)=\lim _{N \rightarrow \infty}\left(\sum_{m=1}^{N} s_{m j} \sigma_{A_{m}}^{\prime}, \sum_{n=1}^{N} s_{n k} \sigma_{A_{n}}^{\prime}\right)=\sum_{m} s_{m j} s_{m k} .
$$

Hence we have

$$
\left(\omega_{A j}, \omega_{A_{k}}\right)=\tau_{j k},
$$

consequently for any finite pair $\left(\xi_{1}, \cdots, \xi_{n}\right) \neq(0, \cdots, 0)$

$$
\left\|\xi_{1} \omega_{A_{\nu_{1}}}+\cdots+\xi_{n} \omega_{A_{\nu_{n}}}\right\|^{2}=\sum_{i, j=1}^{n} \xi_{i} \bar{\xi}_{j} \tau_{\nu_{i} \nu_{j}}>0
$$

Now let $\Omega_{n}=\sum_{i=1}^{n} \xi_{i} \omega_{A_{i}}(n=1,2, \cdots)$ and quadratic form $\sum \xi_{i} \bar{\xi}_{j} \tau_{i j}$ be convergent, then for given $\varepsilon>0$

$$
\left\|\Omega_{m}-\Omega_{n}\right\|^{2}=\sum_{i, j=n+1}^{m} \xi_{i} \bar{\xi}_{j} \tau_{i j}<\varepsilon \quad(m>n>N) .
$$

Therefore there is an $\Omega \in \mathrm{I}_{A}$ such that $\left\|\Omega_{n}-\Omega\right\| \rightarrow 0(n \rightarrow \infty)$ and

$$
\|\Omega\|^{2}=\lim _{n \rightarrow \infty}\left\|\Omega_{n}\right\|^{2}=\sum \xi_{i} \bar{\xi}_{j} \tau_{i j} \geq 0
$$

If the last quadratic form vanishes, $\Omega=0$ and since $\Omega_{n}$ converge to $\Omega$ uniformly on every compact set on $R$, we have

$$
0=\int_{A j} \Omega=\lim _{n \rightarrow \infty} \int_{A j} \Omega_{n}=\xi_{j} \quad(j=1,2, \cdots) .
$$

Next we prove $3^{\circ}$ ). By (13), (15)

$$
\begin{aligned}
\sigma_{j k} & =-\sum_{m}\left(\operatorname{Re} b_{m k}\right) \int_{B j} \operatorname{Re} \sigma_{A m}^{\prime}=-\sum_{m}\left(\operatorname{Re} b_{m k}\right)\left(\operatorname{Im} b_{m j}\right) \\
& =-\sum_{m}\left[b_{m k} b_{m j}-\bar{b}_{m k} \bar{b}_{m j}-2 i \operatorname{Im}\left(b_{m j} \bar{b}_{m k}\right)\right] / 4 i .
\end{aligned}
$$

Hence the symmetry of $\sigma$ is equivalent to the condition

$$
\sum_{m} \operatorname{Im}\left(b_{m j} \bar{b}_{m k}\right)=0
$$

On the other hand we have by (14)

$$
\left(\varphi_{B j}^{0}, \varphi_{B_{k}}^{n}\right)=\sum b_{m j} \bar{b}_{m k} .
$$

Therefore if $\mathrm{I}_{A}^{\prime}=\Gamma_{A B}$, then $\mathcal{P}_{B j}=\mathscr{P}_{B j}^{\cap}, \mathscr{P}_{B k}=\mathscr{P}_{B_{k}}^{\cap}$ and we know $\left(\mathcal{P}_{B j}^{0}, \mathscr{P}_{B_{k}}^{n}\right)=\left(\mathcal{P}_{B j}, \mathscr{P}_{B_{k}}\right)$ are real. Q. E. D.

REMARK. In case of $\Gamma_{A}=\Gamma_{A B}, \sigma_{j k}$ are given by

$$
\sigma_{j k}=-\frac{1}{2} \operatorname{Im} \sum_{m} b_{m j} b_{m k}=\frac{1}{2} D\left(d u_{B j}, * d \tilde{u}_{B k}\right) .
$$

where

$$
\sum_{n} \bar{b}_{m k} \sigma_{A_{m}}^{\prime}=d \tilde{u}_{B k}+i^{*} d \tilde{u}_{B_{k}}
$$

For the sake of convenience we reformulate a result contained in the above proof as follows.

Proposition 4. Let $\left\{\alpha_{n}\right\}$ be given complex numbers for which $\sum \alpha_{i} \bar{\alpha}_{j} \tau_{i j}$ is convergent, then there exists in $\Gamma_{A}$ a unique differential $\Omega$ with $A$-periods $\left\{\alpha_{n}\right\}$, which is expressed as

$$
\begin{equation*}
\Omega=\sum \alpha_{n} \omega_{A_{n}}, \quad\|\Omega\|^{2}=\sum \alpha_{i} \bar{\alpha}_{j} \tau_{i j}<\infty, \tag{16}
\end{equation*}
$$

and the convergence is uniform on every compact set on $R$.
It is desirable that every differential $\phi \in \mathrm{\Gamma}_{A}$ can be written as (16). Of course it is true in case of finite genus. As for general
case we know $\sum_{n}\left|\sum_{k=1}^{n} s_{n k} \alpha_{k}\right|^{2}<\infty$ for $\alpha=P_{\gamma}(\phi)$ while by (15) $\sum \alpha_{i} \bar{\alpha}_{j} \tau_{i j}$ means the interchange of above summation. I have no answer about this problem. Here we shall give only some criteria about the convergence of hermitian form $J=\sum \alpha_{i} \bar{\alpha}_{j} \tau_{i j}$.
a) Since $\tau_{i j}^{2}<\tau_{i i} \tau_{j j}, J$ is convergent if

$$
\sum\left|\alpha_{i}\right| \sqrt{\tau_{i i}}=\sum\left|\alpha_{i}\right| \| \omega_{A_{i}}| |<\infty
$$

(cf. [5], [6], [7]).
b) As $\tau_{i j}$ (or $s_{i j}$ ) are complicatedly dependent on the structure of $R$, the following criterion by geometric quantities seems to be useful.

Proposition 5. Let $\lambda(\gamma)$ denote the extremal length of curve family homologous to a cycle $\gamma$. Then $J$ is convergent if

$$
\begin{equation*}
\sum_{i}\left|\alpha_{i}\right| \sqrt{\lambda\left(B_{i}\right)}<\infty \tag{17}
\end{equation*}
$$

Proof. We note [3] that for $\omega \in \mathrm{I}_{a}$

$$
\begin{equation*}
\left|\int_{\gamma} \omega\right|^{2} \leq \lambda(\gamma)\|\omega\|^{2} \tag{18}
\end{equation*}
$$

On account of symmetry of $\tau_{i j}$ it suffices to treat the case of real $\alpha_{i}$. Applying (18) for $\omega=\Omega_{n}^{n}=\sum_{i=n}^{m} \alpha_{i} \omega_{A_{i}}, \gamma=B_{j}$ we have

$$
\left|\sum_{i=n}^{m} \alpha_{i} \tau_{j \boldsymbol{i}}\right|^{2} \leq\left|\int_{B j} \Omega_{n}^{m}\right|^{2} \leq \lambda\left(B_{j}\right)\left\|\Omega_{n}^{m}\right\|^{2}
$$

Multiplying by $\left|\alpha_{j}\right|$ and summing from $j=n$ to $m$, then

$$
\left|\sum_{i, j=n}^{m} \alpha_{i} \alpha_{j} \tau_{i j}\right| \leq \sum_{j}\left|\alpha_{j}\right|\left|\sum_{i} \alpha_{i} \tau_{j i}\right| \leq \| \Omega_{n}^{m}| | \sum_{j=n}^{m}\left|\alpha_{j}\right| \sqrt{\lambda\left(B_{j}\right)} .
$$

Since $\left\|\Omega_{n}^{m}\right\|=\left(\sum_{i, j=n}^{m} \alpha_{i} \alpha_{j} \tau_{i j}\right)^{1 / 2}$ we have our conclusion.
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## REFERENCES

[1] Ahlfors, L., Open Riemann surfaces and extremal problems on compact subregion. Comm. Math. Helv. 24 (1950) 100-134.
[2] Ahlfors, L. and Sario, L., Riemann surfaces. Princeton (1960).
[3] Kusunoki, Y., On Riemann's periods relations on open Riemann surfaces. Mem. Col. Sci., Univ. of Kyoto, Ser. A. Math. 30 (1956) 1-22.
[4] Kusunoki, Y., Theory of abelian integrals and its applications to conformal map. pings. Ibid. 32 (1959) 235-258.
[5] Myrberg, P. J., Über transzendente hyperelliptische Integrale erster Gattung. Ann. Acad. Sci. Fenn. 14 (1943).
\{6\} Sainouchi, Y., On the semiexact analytic differentials on an open Riemann surface (to appear).
[7] Virtanen, K. I., Über Abelsche Integrale auf nullberandeten Riemannschen Flächen von unendlichem Geschlecht. Ibid. 56 (1949).
[8] Virtanen, K. I., Bemerkungen zur Theorie der quadratisch integrierbaren analytischen Differentiale. Ibid. 78 (1950).


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