

Phase transition in one-dimensional Ising models with spatially inhomogeneous potentials

By

Munemi MIYAMOTO

(Received July 5, 1983)

Gibbs distributions in one-dimension have attracted interest of several authors [1-9]. In usual cases there can be no phase transition. But we have examples of phase transition in the following cases:

1. The set of values of random fields is $\{-1, +1\}$ and the interaction is of long range. Dyson [1].
 2. The random fields take values in a countably infinite set and the interaction is of nearest neighbour. Spitzer [7].
 3. The random fields take values in $\{-1, +1\}$ and the interaction is of nearest neighbour but it is spatially inhomogeneous. Sullivan [8].
- We consider Ising models in one-dimension with spatially inhomogeneous potentials. Let $\Omega = \{-1, +1\}^Z$. For a configuration $\sigma \in \Omega$, let σ_k be the spin of σ at $k \in Z$. The formal Hamiltonian $H(\sigma)$ is equal to

$$H(\sigma) = \sum_{k \in Z} J_k \sigma_k \sigma_{k+1} - \sum_{k \in Z} h_k \sigma_k \quad (\sigma \in \Omega).$$

We call $J = \{J_k\}$ and $h = \{h_k\}$ *interaction* and *external field*, respectively. In ordinary Ising models, J_k and h_k do not depend on k , but in our case they depend on k . For $n \leq m \in Z$ and $\sigma = (\sigma_n, \sigma_{n+1}, \dots, \sigma_m) \in \{-1, +1\}^{[n, m]}$, let

$$H^{[n, m]}(\sigma | \sigma_{n-1}, \sigma_{m+1}) = \sum_{k=n-1}^m J_k \sigma_k \sigma_{k+1} - \sum_{k=n}^m h_k \sigma_k,$$

$$q_{\sigma_{n-1}, \sigma_{m+1}}^{[n, m]}(\sigma) = \mathcal{E}^{[n, m]}(\sigma_{n-1}, \sigma_{m+1})^{-1} \exp\{-H^{[n, m]}(\sigma | \sigma_{n-1}, \sigma_{m+1})\},$$

where $\mathcal{E}^{[n, m]}(\sigma_{n-1}, \sigma_{m+1}) = \sum_{\sigma} \exp\{-H^{[n, m]}(\sigma | \sigma_{n-1}, \sigma_{m+1})\}$. A probability measure μ on Ω is called *Gibbs distribution with potential* (J, h) , if for each $\sigma \in \{-1, +1\}^{[n, m]}$

$$\mu(\sigma | \mathcal{B}_{[n, m]^c}) = q_{\sigma_{n-1}, \sigma_{m+1}}^{[n, m]}(\sigma) \quad a. e. (\mu).$$

The set of all Gibbs distributions with (J, h) is denoted by $\mathcal{G}(J, h)$. It is well known that if J_k and h_k do not depend on k , then $\mathcal{G}(J, h)$ consists

of a unique measure μ . The measure μ is a distribution of a time-homogeneous Markov chain.

Our aim is to clarify under what conditions the uniqueness of $\mathcal{G}(J, h)$ is broken. Put $\hat{J}_n = -|J_n|$ and $\hat{h}_n = s_n h_n$, where

$$s_n = \begin{cases} \prod_{k=0}^{n-1} \operatorname{sgn}(-J_k), & \text{if } n \geq 1, \\ 1, & \text{if } n = 0, \\ \prod_{k=n}^{-1} \operatorname{sgn}(-J_k), & \text{if } n \leq -1. \end{cases}$$

Define a transformation $\sigma \rightarrow \hat{\sigma}$ of Ω by

$$\hat{\sigma}_n = s_n \sigma_n.$$

Remark that the formal Hamiltonians induced by (J, h) and (\hat{J}, \hat{h}) coincide, i. e.,

$$\sum_{n \in \mathbb{Z}} J_n \sigma_n \sigma_{n+1} - \sum_{n \in \mathbb{Z}} h_n \sigma_n = \sum_{n \in \mathbb{Z}} \hat{J}_n \hat{\sigma}_n \hat{\sigma}_{n+1} - \sum_{n \in \mathbb{Z}} \hat{h}_n \hat{\sigma}_n.$$

Therefore, $\mathcal{G}(J, h)$ and $\mathcal{G}(\hat{J}, \hat{h})$ are isomorphic to each other. Let \hat{q} be the conditional Gibbs distribution induced by (\hat{J}, \hat{h}) . Since the potential (\hat{J}, \hat{h}) is ferromagnetic, we can see by the FKG inequality [10] that the limit

$$\hat{q}_{\tau', \tau} \equiv \lim_{\substack{n \rightarrow -\infty \\ m \rightarrow +\infty}} \hat{q}_{\tau', \tau}^{[n, m]}$$

exists for constant boundary conditions $\sigma_{n-1} = \tau'$ and $\sigma_{m+1} = \tau$ with $\tau', \tau = \pm 1$. Let $\mathcal{G}_{ex}(\hat{J}, \hat{h})$ be the set of all extremal measures of the convex set $\mathcal{G}(\hat{J}, \hat{h})$. It is easy to see that

$$\mathcal{G}_{ex}(\hat{J}, \hat{h}) = \{\hat{q}_{\tau', \tau}; \tau', \tau = \pm 1\}.$$

Returning to (J, h) , we have

$$\mathcal{G}_{ex}(J, h) = \{q_{\tau', \tau}; \tau', \tau = \pm 1\},$$

where $q_{\tau', \tau} = \lim_{s_{n-1}\tau', s_{m+1}\tau} q_{s_{n-1}\tau', s_{m+1}\tau}^{[n, m]}$.

For a 2×2 -matrix A , we denote the (σ', σ) -component of A by $A(\sigma', \sigma)$, i. e.,

$$A = \begin{pmatrix} A(-1, -1) & A(-1, +1) \\ A(+1, -1) & A(+1, +1) \end{pmatrix}.$$

Let $\mathbf{a}(\sigma)$ and $\mathbf{b}(\sigma)$ be the σ -th column and σ -th row of A , respectively. Put

$$\tilde{A} = (\mathbf{a}(-1)/\|\mathbf{a}(-1)\|, \mathbf{a}(+1)/\|\mathbf{a}(+1)\|),$$

$$\tilde{A}^* = \begin{pmatrix} \mathbf{b}(-1)/\|\mathbf{b}(-1)\| \\ \mathbf{b}(+1)/\|\mathbf{b}(+1)\| \end{pmatrix},$$

where $\| \cdot \|$ is the Euclidean norm. Put

$$Q_k = \begin{pmatrix} e^{-J_k - h_k} & e^{J_k - h_k} \\ e^{J_k + h_k} & e^{-J_k + h_k} \end{pmatrix},$$

i. e., $Q_k(\sigma', \sigma) = e^{-J_k \sigma' + h_k \sigma}$ and put

$$\Pi_c^d = \begin{cases} Q_c Q_{c+1} \cdots Q_d, & \text{if } d \geq c, \\ E, & \text{if } d = c - 1. \end{cases}$$

Let $n < l \leq r < m$ and let $\sigma = (\sigma_l, \sigma_{l+1}, \dots, \sigma_r) \in \{-1, +1\}^{[l, r]}$. We have

$$q_{\sigma_{n-1}, \sigma_{m+1}}^{[n, m]}(\sigma) = \frac{\Pi_{n-1}^{l-1}(\sigma_{n-1}, \sigma_l) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \Pi_r^m(\sigma_r, \sigma_{m+1})}{\Pi_{n-1}^m(\sigma_{n-1}, \sigma_{m+1})} \\ = \frac{\tilde{\Pi}_{n-1}^{*l-1}(\sigma_{n-1}, \sigma_l) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \tilde{\Pi}_r^m(\sigma_r, \sigma_{m+1})}{\tilde{\Pi}_{n-1}^{*l-1} \Pi_l^{r-1} \tilde{\Pi}_r^m(\sigma_{n-1}, \sigma_{m+1})}.$$

Since $\tilde{\Pi}_{n-1}^{*l-1}$ and $\tilde{\Pi}_r^m$ are bounded, we can extract subsequences $n_i \rightarrow -\infty$ and $m_j \rightarrow +\infty$ such that

$$\tilde{\Pi}_{-\infty}^{*l-1}(\tau, \sigma) \equiv \lim_i \tilde{\Pi}_{n_i-1}^{*l-1}(s_{n_i-1} \tau, \sigma),$$

$$\tilde{\Pi}_r^{+\infty}(\sigma, \tau) \equiv \lim_j \tilde{\Pi}_r^{m_j}(\sigma, s_{m_j+1} \tau)$$

exist for all l, r, τ and σ . Hence, we have

$$q_{\tau', \tau}(\sigma) = \frac{\tilde{\Pi}_{-\infty}^{*l-1}(\tau', \sigma_l) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau)}{\tilde{\Pi}_{-\infty}^{*l-1} \Pi_l^{r-1} \tilde{\Pi}_r^{+\infty}(\tau', \tau)}.$$

We can see by the following Lemma 1 that the matrices $\tilde{\Pi}_r^{+\infty}$ and $\tilde{\Pi}_{-\infty}^{*l-1}$ are positive.

Lemma 1. *There are sequences $\{\rho_r(\sigma)\}_r$ and $\{\rho_l^*(\sigma)\}_l$ of positive functions such that*

$$\tilde{\Pi}_r^{+\infty}(\sigma', \sigma) = Q_r \tilde{\Pi}_{r+1}^{+\infty}(\sigma', \sigma) \rho_r(\sigma), \\ \tilde{\Pi}_{-\infty}^{*l}(\sigma', \sigma) = \rho_l^*(\sigma') \tilde{\Pi}_{-\infty}^{*l-1} Q_l(\sigma', \sigma).$$

Proof is easy.

We say that $\{\tilde{\Pi}_r^{+\infty}\}$ is *asymptotically non-singular* if $\tilde{\Pi}_r^{+\infty}$ is non-singular for sufficiently large r . If not, Lemma 1 shows that $\tilde{\Pi}_r^{+\infty}$ is singular for all r . Asymptotic non-singularity of $\tilde{\Pi}_{-\infty}^{*l-1}$ is defined in the same way.

We prove the following Theorems.

Theorem 1. *Put*

$$\mathcal{M}_{+\infty}(J, h) = \begin{cases} \{-1, +1\}, & \text{if } \{\tilde{\Pi}_r^{+\infty}\} \text{ is asymptotically non-singular,} \\ \{+1\}, & \text{if otherwise.} \end{cases}$$

A set $\mathcal{M}_{-\infty}(J, h)$ is defined analogously. Put

$$\mathcal{M}(J, h) = \mathcal{M}_{-\infty}(J, h) \times \mathcal{M}_{+\infty}(J, h).$$

The set $\mathcal{G}_{ex}(J, h)$ is isomorphic to $\mathcal{M}(J, h)$. The mapping

$$q = q_{\tau', \tau}: \mathcal{M}(J, h) \rightarrow \mathcal{G}_{ex}(J, h)$$

is an isomorphism. If $\mathcal{M}_{+\infty}(J, h) = \{-1, +1\}$, then for $(\tau', \tau) \in \mathcal{M}(J, h)$,

$$q_{\tau', \tau}(\sigma_n = s_n \tau \text{ for all but finitely many } n \geq 0) = 1.$$

We call $\mathcal{M}(J, h)$ the Martin-Dynkin boundary of $\mathcal{G}(J, h)$.

Theorem 2. 1) Assume that $\mathcal{M}_{-\infty}(J, h)$ or $\mathcal{M}_{+\infty}(J, h)$ is equal to $\{+1\}$. Then, any $\mu \in \mathcal{G}(J, h)$ is a Markov chain.

2) Assume that $\mathcal{M}_{-\infty}(J, h) = \mathcal{M}_{+\infty}(J, h) = \{-1, +1\}$. Then, a measure

$$\mu = \sum \lambda_{\tau', \tau} q_{\tau', \tau} \in \mathcal{G}(J, h) \quad (\sum \lambda_{\tau', \tau} = 1, \lambda_{\tau', \tau} \geq 0)$$

is a Markov chain, if and only if

$$\det(\lambda_{\tau', \tau} / \tilde{\Pi}_{-\infty}^{*-1} \Pi_{\tau'}^{-1} \tilde{\Pi}_{\tau}^{+\infty}(\tau', \tau)) = 0.$$

This condition is consistent in the sense that it is equivalent to

$$\det(\lambda_{\tau', \tau} / \tilde{\Pi}_{-\infty}^{*l'-1} \Pi_{\tau'}^{l'-1} \tilde{\Pi}_{\tau}^{+\infty}(\tau', \tau)) = 0$$

for any $l' \leq r'$.

Theorem 3. If $\sum_{n=0}^{+\infty} e^{-2|J_n|} = +\infty$, then $\mathcal{M}_{+\infty}(J, h) = \{+1\}$ for any $h = \{h_n\}$.

Theorem 4. If

$$\sum_{n=0}^{+\infty} \exp 2(-|J_n| + |\sum_{k=0}^n s_k h_k|) < +\infty,$$

then $\mathcal{M}_{+\infty}(J, h) = \{-1, +1\}$. Conversely, if

$$\sum_{n=0}^{+\infty} \exp 2(-|J_n| + |\sum_{k=0}^n s_k h_k|) = +\infty$$

and if in addition $s_k h_k \geq 0$ (or ≤ 0) for all $k \geq 0$, then $\mathcal{M}_{+\infty}(J, h) = \{+1\}$.

As for homogeneous external fields, we have

Theorem 5. Assume that the external field h_n is equal to a constant h and $J_n \leq 0$ for all $n \geq 0$. Then, $\mathcal{M}_{+\infty}(J, h) = \{+1\}$ if and only if

$$\sum_{n=0}^{+\infty} e^{2(J_n + n|h|)} = +\infty.$$

Theorem 6. Assume that the external field h_n is equal to a constant and $J_n \geq 0$ for all $n \geq 0$. Then, $\mathcal{M}_{+\infty}(J, h) = \{+1\}$ if and only if

$$\sum_{n=0}^{+\infty} e^{-2J_n} = +\infty.$$

Theorems 5 and 6 are corollaries to Theorems 3 and 4.

We will discuss the one-dimensional Widom-Rowlison models in a forthcoming paper [5].

Proof of Theorem 1. The first statement in Theorem 1 follows immediately from Lemma 2 below. To prove the second statement, put

$$B_{\tau',\tau} = \{\sigma = (\sigma_m)_m; \lim_{n,m} q_{s_{n-1}\tau',\sigma_{m+1}}^{[n,m]} = q_{\tau',\tau}\}.$$

By Theorem 1 in [11], $q_{\tau',\tau}(B_{\tau',\tau}) = 1$. Assume $\tilde{\Pi}_r^{+\infty}(\sigma, -1) \neq \tilde{\Pi}_r^{+\infty}(\sigma, +1)$ for sufficiently large r . Then, $\sigma \in B_{\tau',\tau}$ if and only if $\sigma_m = s_m\tau$ for all but finitely many $m \geq 0$. Q. E. D.

Lemma 2. *It holds $q_{\tau'_1,\tau_1} = q_{\tau'_2,\tau_2}$, if and only if $(\tilde{\Pi}_r^{+\infty}(\sigma, \tau_i))_{\sigma,i}$ and $(\tilde{\Pi}_{-\infty}^{*-1}(\tau'_i, \sigma))_{i,\sigma}$ are singular for all l and r .*

Proof. Assume $q_{\tau'_1,\tau_1} = q_{\tau'_2,\tau_2}$. Take any $\sigma = (\sigma_i, \sigma_{i+1}, \dots, \sigma_r) \in \{-1, +1\}^{[l,r]}$. Put $P = \tilde{\Pi}_{-\infty}^{*-1} \Pi_l^{-1} \tilde{\Pi}_r^{+\infty}$. From

$$\frac{\tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma_r) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_1)}{P(\tau'_1, \tau_1)} = \frac{\tilde{\Pi}_{-\infty}^{*-1}(\tau'_2, \sigma_l) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_2)}{P(\tau'_2, \tau_2)},$$

it follows that

$$\begin{aligned} \tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma_l) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_1) / P(\tau'_1, \tau_1) \\ = \tilde{\Pi}_{-\infty}^{*-1}(\tau'_2, \sigma_l) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_2) / P(\tau'_2, \tau_2). \end{aligned}$$

Since $\sum_{\sigma_l} \tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma_l)^2 = \sum_{\sigma_r} \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_1)^2 = 1$, we have $P(\tau'_1, \tau_1) = P(\tau'_2, \tau_2)$.

Therefore,

$$\tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma_l) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_1) = \tilde{\Pi}_{-\infty}^{*-1}(\tau'_2, \sigma_l) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_2).$$

By the same argument as above, we have

$$\begin{aligned} \tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma_l) &= \tilde{\Pi}_{-\infty}^{*-1}(\tau'_2, \sigma_l), \\ \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_1) &= \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau_2). \end{aligned}$$

Conversely, assume that $(\tilde{\Pi}_{-\infty}^{*-1}(\tau'_i, \sigma))_{i,\sigma}$ and $(\tilde{\Pi}_r^{+\infty}(\sigma, \tau_i))_{\sigma,i}$ are singular. Then, $\tilde{\Pi}_{-\infty}^{*-1}(\tau'_1, \sigma) = \tilde{\Pi}_{-\infty}^{*-1}(\tau'_2, \sigma)$ and $\tilde{\Pi}_r^{+\infty}(\sigma, \tau_1) = \tilde{\Pi}_r^{+\infty}(\sigma, \tau_2)$, because the rows of $\tilde{\Pi}_{-\infty}^{*-1}$ and the columns of $\tilde{\Pi}_r^{+\infty}$ are non-negative and normalized, which implies $q_{\tau'_1,\tau_1} = q_{\tau'_2,\tau_2}$. Q. E. D.

Proof of Theorem 2. Take any $\mu = \sum_{\tau',\tau} \lambda_{\tau',\tau} q_{\tau',\tau} \in \mathcal{G}(J, h)$ and $\sigma = (\sigma_l, \sigma_{l+1}, \dots, \sigma_r) \in \{-1, +1\}^{[l,r]}$.

Put $\hat{\lambda}_{\tau',\tau} = \lambda_{\tau',\tau} / \tilde{\Pi}_{-\infty}^{*-1} \Pi_l^{-1} \tilde{\Pi}_r^{+\infty}(\tau', \tau)$. We have

$$\begin{aligned} \mu(\sigma) &= \sum_{\tau',\tau} \hat{\lambda}_{\tau',\tau} \tilde{\Pi}_{-\infty}^{*-1}(\tau', \sigma_l) \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) \tilde{\Pi}_r^{+\infty}(\sigma_r, \tau) \\ &= \prod_{k=l}^{r-1} Q_k(\sigma_k, \sigma_{k+1}) L(\sigma_r, \sigma_l), \end{aligned}$$

where $L = \tilde{\Pi}_r^{+\infty}(\hat{\lambda}_{\tau',\tau}) \tilde{\Pi}_{-\infty}^{*-1}$. Since

$$\mu(\sigma_l, \sigma_{l+1}, \dots, \sigma_{r-1}) = \prod_{k=l}^{r-2} Q_k(\sigma_k, \sigma_{k+1}) \sum_{\sigma} Q_{r-1}(\sigma_{r-1}, \sigma) L(\sigma, \sigma_l),$$

we have

$$\begin{aligned} \mu(\sigma_r | \sigma_l, \sigma_{l+1}, \dots, \sigma_{r-1}) &= Q_{r-1}(\sigma_{r-1}, \sigma_r) L(\sigma_r, \sigma_l) / \sum_{\sigma} Q_{r-1}(\sigma_{r-1}, \sigma) L(\sigma, \sigma_l) \\ &= \mu(\sigma_r | \sigma_l, \sigma_{r-1}). \end{aligned}$$

Consequently,

$$\begin{aligned} \mu(\sigma_r | \sigma_l = +1, \sigma_{r-1}) - \mu(\sigma_r | \sigma_l = -1, \sigma_{r-1}) &= \\ = \prod_{\sigma} \{Q_{r-1}(\sigma_{r-1}, \sigma) / Q_{r-1} L(\sigma_{r-1}, \sigma)\} \det L & \\ = \prod_{\sigma} \{Q_{r-1}(\sigma_{r-1}, \sigma) / Q_{r-1} L(\sigma_{r-1}, \sigma)\} \det \tilde{\Pi}_{-\infty}^{*-1} \det \tilde{\Pi}_r^{+\infty} \det(\hat{\lambda}_{\tau',\tau}). & \end{aligned}$$

If one of above three determinants vanishes, then

$$\mu(\sigma_r | \sigma_l = +1, \sigma_{r-1}) = \mu(\sigma_r | \sigma_l = -1, \sigma_{r-1}) = \mu(\sigma_r | \sigma_{r-1}),$$

from which follows $\mu(\sigma_r | \sigma_l, \sigma_{l+1}, \dots, \sigma_{r-1}) = \mu(\sigma_r | \sigma_{r-1})$.

Consistency in 2) is clear from Lemma 1.

Q. E. D.

Let A be a 2×2 -matrix whose columns are \mathbf{a} and \mathbf{a}' . Put

$$\Delta(A) = |\det \tilde{A}| = |\det A| / (\|\mathbf{a}\| \|\mathbf{a}'\|).$$

Denote the inner product of \mathbf{a} and \mathbf{a}' by $\langle \mathbf{a}, \mathbf{a}' \rangle$. It is easy to see

Lemma 3.

- 1) $\Delta(A) \leq |\det A| / \langle \mathbf{a}, \mathbf{a}' \rangle$.
- 2) Let $D = \begin{pmatrix} \alpha & 0 \\ 0 & \beta \end{pmatrix}$ with $\alpha, \beta \neq 0$. Then, $\Delta(AD) = \Delta(A)$.

Let us write $A \cong B$, if $A = \alpha B$ with a scalar $\alpha \neq 0$. Remark that $\Delta(A) = \Delta(B)$ if $A \cong B$. For $H(h) = \begin{pmatrix} e^{-h} & 0 \\ 0 & e^h \end{pmatrix}$ and $K(J) = \begin{pmatrix} e^{-J} & e^J \\ e^J & e^{-J} \end{pmatrix}$, we have

$$H(h) K(J) H(h)^{-1} \cong E + e^{2J} H(h)^2 F,$$

where $F = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. Put $H_k = H(\hat{h}_k)$, $K_k = K(\hat{J}_k)$, $Q_k = H_k K_k$ and

$$P_k = E + e^{2\hat{J}_k} \left(\prod_{j=r}^k H_j \right)^2 F.$$

Remark $\det P_k = 1 - e^{4f_k}$. It holds that

$$\begin{aligned} \prod_r^m &= \prod_{k=r}^m Q_k = \prod_{k=r}^m H_k K_k \\ &= \prod_{k=r}^m \left\{ \left(\prod_{j=r}^k H_j \right) K_k \left(\prod_{j=r}^k H_j \right)^{-1} \right\} \prod_{j=r}^m H_j \\ &\cong \prod_{k=r}^m \{ E + e^{2f_k} \left(\prod_{j=r}^k H_j \right)^2 F \} \prod_{j=r}^m H_j \\ &= \prod_{k=r}^m P_k \prod_{j=r}^m H_j. \end{aligned}$$

Therefore, by 2) in Lemma 3, we have

$$\Delta(\tilde{\Pi}_r^m) = \Delta(\Pi_r^m) = \Delta\left(\prod_{k=r}^m P_k\right).$$

Let $\mathbf{a}_m(\sigma)$ be the σ -th column of $\prod_{k=r}^m P_k$. It is easy to see by direct calculations

Lemma 4. Put $\hat{I}_m = \sum_{k=r}^m \hat{h}_k$ and $\mathbf{b}_m = e^{f_m} \mathbf{a}_{m-1}(+1) - e^{-f_m} \mathbf{a}_{m-1}(-1)$.

- 1) $\mathbf{a}_m(-1) = \mathbf{a}_{m-1}(-1) + e^{2(f_m + \hat{I}_m)} \mathbf{a}_{m-1}(+1)$,
 $\mathbf{a}_m(+1) = e^{2(f_m - \hat{I}_m)} \mathbf{a}_{m-1}(-1) + \mathbf{a}_{m-1}(+1)$.
- 2) $\langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle = e^{2f_m} \|\mathbf{b}_m\|^2 + (1 + e^{2f_m})^2 \langle \mathbf{a}_{m-1}(-1), \mathbf{a}_{m-1}(+1) \rangle$.

Proof of Theorem 3. We prove that $\mathcal{M}_{+\infty}(f, \hat{h}) = \{+1\}$ for any $\hat{h} = \{\hat{h}_n\}$ if $\sum_{n=r}^{+\infty} e^{2f_n} = +\infty$. By 2) in Lemma 4, we have

$$\langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle \geq \prod_{k=r+1}^m (1 + e^{2f_k})^2 \langle \mathbf{a}_r(-1), \mathbf{a}_r(+1) \rangle.$$

Therefore,

$$\begin{aligned} \Delta(\tilde{\Pi}_r^m) &= \Delta\left(\prod_{k=r}^m P_k\right) \leq \prod_{k=r}^m \det P_k / \langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle \\ &\leq \left\{ \prod_{k=r+1}^m (1 + e^{2f_k})^2 \langle \mathbf{a}_r(-1), \mathbf{a}_r(+1) \rangle \right\}^{-1}. \end{aligned}$$

The right-hand side diverges to 0 as $m \rightarrow +\infty$. Hence $\Delta(\tilde{\Pi}_r^{+\infty}) = 0$ for all r , which implies $\mathcal{M}_{+\infty}(f, \hat{h}) = \{+1\}$ by Theorem 1. Q. E. D.

Next, we prove the second statement in Theorem 4.

Lemma 5. If $\hat{h}_n \geq 0$ for all $n \geq 0$, then

$$\mathbf{b}_k \geq \prod_{j=r+2}^k e^{\hat{h}_j} \{1 - e^{2(f_{j-1} - \hat{h}_j)}\} \mathbf{b}_{r+1}.$$

Proof. We see by 1) in Lemma 4,

$$\begin{aligned} \mathbf{b}_k &= e^{f_k} \mathbf{a}_{k-1}(+1) - e^{-f_k} \mathbf{a}_{k-1}(-1) \\ &= (e^{\hat{h}_k} - e^{2f_{k-1} - \hat{h}_k}) e^{f_{k-1}} \mathbf{a}_{k-2}(+1) - (e^{-\hat{h}_k} - e^{2f_{k-1} + \hat{h}_k}) e^{-f_{k-1}} \mathbf{a}_{k-2}(-1). \end{aligned}$$

Since $e^{\hat{h}} - e^{2\hat{f} - \hat{h}} \geq e^{-\hat{h}} - e^{2\hat{f} + \hat{h}}$ if $\hat{h} \geq 0$, we have

$$\begin{aligned} \mathbf{b}_k &\geq (e^{\hat{h}_k} - e^{2\hat{f}_{k-1} - \hat{h}_k}) \mathbf{b}_{k-1} \\ &= e^{\hat{h}_k} \{1 - e^{2(\hat{f}_{k-1} - \hat{h}_k)}\} \mathbf{b}_{k-1}. \end{aligned} \quad \text{Q. E. D.}$$

Proof of the second statement in Theorem 4. It is enough to prove that $\mathcal{M}_{+\infty}(\hat{f}, \hat{h}) = \{+1\}$ if $\sum_{k=0}^{+\infty} \exp 2(\hat{f}_k + \sum_{k=0}^n \hat{h}_k) = +\infty$ and if $\hat{h}_n \geq 0$ for all $n \geq 0$. In case $\sum_{n=0}^{+\infty} e^{2\hat{f}_n} = +\infty$, $\mathcal{M}_{+\infty}(\hat{f}, \hat{h}) = \{+1\}$ by Theorem 3. So, we prove in case $\sum_{n=0}^{+\infty} e^{2\hat{f}_n} < +\infty$.

We replace \hat{f}_r by $\hat{f}'_r = -\hat{h}_{r+1}$. Then,

$$\begin{aligned} \mathbf{b}_{r+1} &= e^{\hat{h}_r + \hat{h}_{r+1}} \mathbf{a}_r(+1) - e^{-\hat{h}_r - \hat{h}_{r+1}} \mathbf{a}_r(-1) \\ &= \begin{pmatrix} e^{-\hat{h}_r} (e^{2\hat{f}'_r + \hat{h}_{r+1}} - e^{-\hat{h}_{r+1}}) \\ e^{\hat{h}_r} (e^{\hat{h}_{r+1}} - e^{2\hat{f}'_r - \hat{h}_{r+1}}) \end{pmatrix} \\ &= \begin{pmatrix} 0 \\ e^{\hat{h}_r} (e^{\hat{h}_{r+1}} - e^{-3\hat{h}_{r+1}}) \end{pmatrix} \geq \mathbf{0}. \end{aligned}$$

Therefore, we can see by Lemma 5,

$$\begin{aligned} \|\mathbf{b}_k\| &\geq \prod_{j=r+2}^k e^{\hat{h}_j} \{1 - e^{2(\hat{f}_{j-1} - \hat{h}_j)}\} \|\mathbf{b}_{r+1}\| \\ &\geq b \exp \sum_{j=r+2}^k \hat{h}_j, \end{aligned}$$

where $b = \|\mathbf{b}_{r+1}\| \prod_{j=r+2}^{+\infty} (1 - e^{2\hat{f}_{j-1}}) > 0$ by our assumption $\sum_{n=0}^{+\infty} e^{2\hat{f}_n} < +\infty$. On the other hand, by 2) in Lemma 4, we have

$$\begin{aligned} \langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle &= \sum_{k=r+1}^m \prod_{j=k+1}^m (1 + e^{2\hat{f}_j})^2 e^{2\hat{f}_k} \|\mathbf{b}_k\|^2 \\ &\quad + \prod_{j=r+1}^m (1 + e^{2\hat{f}_j})^2 \langle \mathbf{a}_r(-1), \mathbf{a}_r(+1) \rangle \\ &\geq \sum_{k=r+1}^m e^{2\hat{f}_k} \|\mathbf{b}_k\|^2 \\ &\geq b^2 \sum_{k=r+1}^m \exp 2(\hat{f}_k + \sum_{j=r+2}^k \hat{h}_j). \end{aligned}$$

The right-hand side diverges to $+\infty$ as $m \rightarrow +\infty$.

Put $P'_r = E + e^{2\hat{f}'_r} H_r^2 F$. Since

$$\begin{aligned} \Delta(P'_r \prod_{k=r+1}^m P_k) &\leq \det P'_r \prod_{k=r+1}^m \det P_k / \langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle \\ &\leq \langle \mathbf{a}_m(-1), \mathbf{a}_m(+1) \rangle^{-1}, \end{aligned}$$

we have $\lim_{m \rightarrow +\infty} \Delta(P'_r \prod_{k=r+1}^m P_k) = 0$. Therefore,

$$\begin{aligned} \Delta(\tilde{\Pi}_r^{+\infty}) &= \lim_{m \rightarrow +\infty} \Delta(\prod_{k=r}^m P_k) \\ &= \lim_{m \rightarrow +\infty} \Delta(P_r P_r^{-1} \{P'_r \prod_{k=r+1}^m P_k\}) = 0. \end{aligned}$$

Thus, $\mathcal{M}_{+\infty}(\hat{J}, \hat{h}) = \{+1\}$ by Theorem 1.

Q. E. D.

We proceed to the proof of the first statement in Theorem 4. Let $\|A\|$ be the operator norm of a matrix A , i. e.,

$$\|A\| = \sup_x \|A\mathbf{x}\| / \|\mathbf{x}\|.$$

Lemma 6. *Let $\{A_k\}_k$ be a sequence of matrices. If*

$$\sum_{k=1}^{+\infty} \|A_k\| < +\infty,$$

then $\prod_{k=1}^m (E + A_k)$ converges as $m \rightarrow +\infty$.

Proof. Let $n > m$. We have

$$\begin{aligned} \left\| \prod_{k=1}^n (E + A_k) - \prod_{k=1}^m (E + A_k) \right\| &\leq \left\| \prod_{k=1}^m (E + A_k) \right\| \left\| \prod_{k=m+1}^n (E + A_k) - E \right\| \\ &\leq \prod_{k=1}^m (1 + \|A_k\|) \left\{ \prod_{k=m+1}^n (1 + \|A_k\|) - 1 \right\}. \end{aligned}$$

The right-hand side converges to 0 as $n > m \rightarrow +\infty$.

Q. E. D.

Proof of the first statement in Theorem 4. We show that $\mathcal{M}_{+\infty}(\hat{J}, \hat{h}) = \{-1, +1\}$ if $\sum_{k=0}^{+\infty} \exp 2(\hat{J}_n + |\sum_{k=0}^n \hat{h}_k|) < +\infty$. Since

$$\begin{aligned} \sum_{k=r}^{+\infty} \|e^{2\hat{J}_k} (\prod_{j=r}^k H_j)^2 F\| &\leq \sum_{k=r}^{+\infty} e^{2\hat{J}_k} \|(\prod_{j=r}^k H_j)^2\| \|F\| \\ &= \sum_{k=r}^{+\infty} \exp 2(\hat{J}_k + |\sum_{j=r}^k \hat{h}_j|) < +\infty, \end{aligned}$$

$\prod_{k=r}^m P_k = \prod_{k=r}^m \{E + e^{2\hat{J}_k} (\prod_{j=r}^k H_j)^2 F\}$ converges as $m \rightarrow +\infty$ by Lemma 6.

On the other hand, for sufficiently large r ,

$$\det \prod_{k=r}^{+\infty} P_k = \prod_{k=r}^{+\infty} (1 - e^{4\hat{J}_k}) > 0.$$

Therefore, we have

$$\Delta(\tilde{\Pi}_r^{+\infty}) = \Delta(\prod_{k=r}^{+\infty} P_k) > 0,$$

which implies $\mathcal{M}_{+\infty}(\hat{J}, \hat{h}) = \{-1, +1\}$ by Theorem 1.

Q. E. D.

References

- [1] F. J. Dyson, Existence of a phase-transition in a one-dimensional Ising ferromagnet, *Comm. math. Phys.*, **12** (1969), 91-107.
- [2] F. J. Dyson, Non-existence of spontaneous magnetization in a one-dimensional Ising ferromagnet. *Comm. math. Phys.*, **12** (1969), 212-215.
- [3] E. Ising, Beitrag zur Theorie des Ferromagnetismus, *Z. Phys.*, **31** (1925), 253-258.
- [4] H. Kesten, Existence and uniqueness of countable one-dimensional Markov random fields, *Ann. Prob.*, **4** (1976), 557-569.
- [5] M. Miyamoto, Phase transition in one-dimensional Widom-Rowlinson models with spatially inhomogeneous potentials, to appear in *J. Math. Kyoto Univ.* **25** (1985).
- [6] J. K. Percus, Higher spin one-dimensional Ising lattice in arbitrary external field, *J. Math. Phys.*, **23** (1982), 1162-1167.
- [7] F. Spitzer, Phase transition in one-dimensional nearest-neighbor systems, *J. Func. Anal.*, **20** (1975), 240-255.
- [8] W. G. Sullivan, Markov processes for random fields, *Comm. Dublin Inst. for Advanced Studies* (1975).
- [9] W. G. Sullivan, P. Vanheuverzwijn, On the canonical Gibbs states associated with certain Markov chains, *Z. Wahr.*, **62** (1983), 171-183.
- [10] C. M. Fortuin, P. W. Kasteleyn, J. Ginibre, Correlation inequalities on some partially ordered sets, *Comm. math. Phys.*, **22** (1971), 89-103.
- [11] M. Miyamoto, Martin-Dynkin boundaries of random fields, *Comm. math. Phys.*, **36** (1974), 321-324.