# EIGENFUNCTION EXPANSIONS $\mathbb{N} L^{p}$ AND $C$ 

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In $L^{p}(1<p<\infty)$ [in $L^{1}$ or $C$ ] the Fourier series is (conditionally ${ }^{1}$ ) convergent $[(C, \theta)$ summable $(\theta>0)]$. Known equiconvergence theorems extend this result to the eigenfunction expansions of operators of the form

$$
T=D^{u} \quad(u \geqq 2),
$$

with $u$ suitable boundary conditions. It will be shown that in cases where the sequence of eigenvalues of $T$ also behaves reasonably, an operator $T+S$ will have the same type of spectral expansion as $T$. Here $S$ can be
(i) any bounded linear operator;
(ii) in very restricted cases, any operator

$$
B_{u-2} D^{u-2}+\cdots+B_{1} D+B_{0}
$$

where the $B_{i}$ are bounded linear operators.
The methods are those developed by Schwartz [4] and Kramer [2] for the $L^{2}$ case, but very much modified to deal with conditionally convergent or summable expansions.

## 1. Definitions and known results

$\mathfrak{B}$ will denote a Banach space with complex scalars. $\quad L^{p}$ will denote $L^{p}[0,1]$ and $C$ will denote $C[0,1]$ (or one of its principal subspaces; see $\S 5(1)$ ). If $T$ is a linear operator in $\mathfrak{B}$, the set of complex numbers $\lambda$ such that $(T-\lambda I)^{-1}$ exists and is a bounded linear operator on $\mathfrak{B}$ to $\mathfrak{B}$ will be written $\rho(T)$ and called the resolvent set of $T$. The complement of $\rho(T)$ in the complex plane is $\sigma(T)$, the spectrum of $T . \mathfrak{D}(T)$ will denote the domain, and $\Re(T)$ the range, of $T$.

If $G$ is a contour in $\rho(T)$, the following results are well known. (For our purposes it can be assumed that $G$ is a circle or a rectangle.)

Lemma 1.1. The operator

$$
P=\frac{-1}{2 \pi i} \int_{G}(T-\lambda I)^{-1} d \lambda
$$

and its complement $I-P$ are bounded projections onto subspaces invariant under $T$. $\Re(P) \subseteq \mathfrak{D}(T)$, and, as an operator in $\Re(P), T$ is bounded. The spectrum of $T$ in $\Re(P)$ is the part of the spectrum of $T$ lying inside $G$. If $\tau_{1}$, $\tau_{2}$ are disjoint components of $\sigma(T)$ and the curves $G_{1}, G_{2}, G_{3}$ contain respectively

[^0]the components $\tau_{1}, \tau_{2}, \tau_{1} \cup \tau_{2}$ of $\sigma(T)$, then for the corresponding projections we have
\[

$$
\begin{equation*}
P_{1} P_{2}=0 ; \quad P_{1}+P_{2}=P_{3} \tag{1.1}
\end{equation*}
$$

\]

Proof. See Taylor [6], especially Theorem 8.2 and the lines preceding the formulae (8.3).

If the contour $G_{i}$ contains only one point of $\sigma(T)$, say $\lambda_{i}$, the dimension of $\Re\left(P_{i}\right)$ will be called the generalised multiplicity of $\lambda_{i}$. It is easily shown by induction on $n$ that $\Re\left(P_{i}\right)$ contains the set

$$
\mathfrak{B}\left(\lambda_{i}\right)=\cup_{n=1}^{\infty}\left\{x:\left(T-\lambda_{i} I\right)^{n} x=0\right\}
$$

the set of proper vectors of $\lambda_{i}$. Thus the dimension of $\mathfrak{B}\left(\lambda_{i}\right)$ (which will be called the algebraic multiplicity of $\lambda_{i}$ ) is at most the generalised multiplicity of $\lambda_{i}$. The term eigenvector will be reserved for a (nonzero) solution of ( $T-\lambda I$ ) $x=0$, and the term eigenvalue for a scalar $\lambda$ for which eigenvectors exist. The multiplicity of $\lambda$ is the dimension of its space of eigenvectors. The point spectrum, written $\pi(T)$, is the set of eigenvalues of $T$.

Lemma 1.2. If $\Re(P)$ is finite-dimensional, then $G$ contains only a finite number of points of $\sigma(T)$, say $\lambda_{1}, \lambda_{2}, \cdots, \lambda_{n}$, and
(i) the generalised multiplicity of $\lambda_{i}, m_{i}$ say, is finite and equal to its algebraic multiplicity and is the same whether we consider the operator $T$ in $\mathfrak{B}$ or the operator $T_{\Re(P)}$ in $\Re(P)$;
(ii) $\sum_{1}^{n} m_{i}=\operatorname{dim} \Re(P)$;
(iii) $\sum_{1}^{n} P_{i}=P$.

Proof. The part of $\sigma(T)$ inside $G$ is the spectrum of $T$ as an operator in the finite-dimensional space $\Re(P)$, which must be a finite set of points. Enclosing each of these points in a contour $G_{i}$ we obtain (iii) from (1.1), and (ii) follows immediately. As $\Re\left(P_{i}\right)$ is finite-dimensional and the spectrum of $T$ in $\Re\left(P_{i}\right)$ is the one-point set $\left\{\lambda_{i}\right\}$, we have

$$
\mathfrak{R}\left(P_{i}\right)=\mathfrak{B}\left(\lambda_{i}\right),
$$

which proves (i).
A result proved by Nagy [3] is
Lemma 1.3. Let $E, F$ be bounded projections in a Banach space $\mathfrak{F}$, with $\|E-F\|<1$. Then $E \mathfrak{B}$ and $F \mathfrak{B}$ have the same dimension if the dimension of either is finite.

Proof. Let $\left(x_{i}\right)_{1 \leqq i \leqq m}$ be a basis for $E \mathfrak{B}$. We can write $E x=\sum_{1}^{m} y_{i}(\cdot) x_{i}$ with $y_{i} \in \mathfrak{B}^{*}(1 \leqq i \leqq m)$. If $\operatorname{dim} F \mathfrak{B}>m, F \mathfrak{B}$ contains a nonzero vector $x$ orthogonal to $\left(y_{i}\right)_{1 \leqq i \leqq m}$. Thus $F x=x, E x=0$, and

$$
\|E-F\| \geqq\|(E-F) x\| /\|x\|=1
$$

Thus $\operatorname{dim} F \mathfrak{B} \leqq m=\operatorname{dim} E \mathfrak{B}$, and the converse inequality is obtained by a similar argument.

We will say that an operator $T$ in a Banach space $\mathfrak{B}\left(=L^{p}\right.$ or $\left.C\right)$ is of the form $D^{u}$ with the boundary conditions

$$
\begin{equation*}
\sum_{j=0}^{u-1} \alpha_{i j} f^{(j)}(0)+\sum_{j=0}^{u-1} \beta_{i j} f^{(j)}(1)=0 \tag{1.2}
\end{equation*}
$$

where $i$ runs over a finite index set $I$ (possibly empty), if
(i) $\mathfrak{D}(T)$ is the set of functions $f$ on $[0,1]$ such that $f^{(j)}$ is absolutely continuous $(0 \leqq j \leqq u-1), f^{(u)} \in \mathfrak{B}$, and (1.2) is satisfied for all $i \in I$;
(ii) for $f \in \mathfrak{D}(T), T f=f^{(u)}$.

Lemma 1.4. An operator $T$ of the form $D^{u}$ with the boundary conditions (1.2) is a closed linear operator.

Proof. It is obvious from (i) and (ii) that $T$ is linear. For the operator $T_{1}$ of the form $D^{u}$ with boundary conditions

$$
f^{(j)}(0)=f^{(j)}(1)=0 \quad(0 \leqq j \leqq u-1)
$$

it is easily seen that the convergence in $\mathfrak{B}$ norm of a sequence $T_{1} f_{n}=f_{n}^{(u)}$ ( $n \geqq 1$ ) in $\Re\left(T_{1}\right)$ implies the uniform convergence of $f_{n}^{(j)}(0 \leqq j \leqq u-1)$ and the equation

$$
\left(\lim f_{n}\right)^{(u)}=\lim f_{n}^{(u)}
$$

It follows that $T_{1}$ is closed, and, as the graph of $T$ can be obtained from that of $T_{1}$ by removing a finite number of boundary conditions, $T$ is closed. (This proof was adapted from that given by Schwartz [4].)

Lemma 1.5. For $f_{1} \in C, g_{1} \in L^{1}$ we have

$$
\left\|f_{1}\right\|=\sup _{g \epsilon L_{1}}\left|\left(f_{1}, g\right)\right| /\|g\| ; \quad\left\|g_{1}\right\|=\sup _{f \epsilon C}\left|\left(f, g_{1}\right)\right| /\|f\|
$$

where

$$
(f, g)=\int_{0}^{1} f(t) g(t) d t
$$

Proof. After approximating to $g_{1}$ with a continuous function the proof is obvious.

Lemma 1.6. The Fourier series of a function $f$ in $\mathfrak{B}=L^{1}$ or $C\left[\mathfrak{B}=L^{p}\right.$ $(1<p<\infty)$ ] is $(C, \theta)$ summable $(\theta>0)$ [is convergent] to $f$ in the norm of $\mathfrak{B}$.

Proof. As regards $C$ and $L^{p}(1<p<\infty)$ see Zygmund [8], §§3.3 and 7.3. As regards $L^{1}$ : if $\theta>0$ and $\mathfrak{B}=C$ or $L^{1}$, let $F_{n}=F_{n}(\mathfrak{B})=F_{n}(\mathfrak{B}, \theta)$ be the operator in $\mathfrak{B}$ which maps each function in $\mathfrak{B}$ onto the $(C, \theta)$ mean of the first $n$ partial sums of its Fourier series. We have $F_{n}(C) x \rightarrow x(x \in C)$ so that by Lemma 1.5 and the Banach-Steinhaus theorem,

$$
\begin{aligned}
\left\|F_{n}\left(L_{1}\right)\right\| & =\sup \left\{\left|\left(F_{n} g, f\right)\right|: g \epsilon L^{1}, f \in C,\|f\|=\|g\|=1\right\} \\
& =\sup \left\{\left|\left(g, F_{n} f\right)\right|: g \in L^{1}, f \in C,\|f\|=\|g\|=1\right\} \\
& =\left\|F_{n}(C)\right\| \\
& =O(1) \text { as } n \rightarrow \infty .
\end{aligned}
$$

Since $F_{n} g \rightarrow g$ on the dense subspace of $L^{1}$ consisting of trigonometric polynomials, $F_{n} g \rightarrow g$ for all $g$ in $L^{1}$.

## 2. Axioms

All of the following axioms are satisfied by operators having the form $D^{u}$ and domains restricted by $u$ suitable boundary conditions. They have been separated into four groups needed for different purposes. Most of the axioms are relations between the distribution of the eigenvalues of $\bullet T$ and the norms of its spectral projections.

The first group of axioms comprises those needed in showing the existence of the resolvent of $T$.

Condition (a).
( $\mathrm{a}_{0}$ ) $T$ is closed.
( $\mathrm{a}_{1}$ ) $\pi(T)$ consists of a discrete sequence $\left(\lambda_{n}\right)_{n \geqq 1}$ of positive eigenvalues, each of multiplicity one.
( $a_{2}$ ) $\lambda_{n}$ is monotone increasing and

$$
\Delta \lambda_{n}=\lambda_{n+1}-\lambda_{n} \rightarrow \infty .
$$

$\left(a_{3}\right)$ There exist vectors $\left(\phi_{j}\right)_{j \geqq 1}$ in $\mathfrak{B}$ and $\left(\psi_{i}\right)_{i \geqq 1}$ in $\mathfrak{B}^{*}$ such that

$$
\left(T-\lambda_{j} I\right) \phi_{j} \equiv 0 \quad \text { and } \quad \psi_{i}\left(\phi_{j}\right) \equiv \delta_{i j} .
$$

( $a_{4}$ ) The eigenvectors $\left(\phi_{j}\right)_{j \geqq 1}$ form a total set.
$\left(\mathrm{a}_{5}\right)$ For the projections

$$
\begin{aligned}
E\left(\lambda_{j}\right) & =\psi_{j}(\cdot) \phi_{j} \\
E_{n}=\sum_{1}^{n} E\left(\lambda_{j}\right) & =\sum_{1}^{n} \psi_{j}(\cdot) \phi_{j}
\end{aligned}
$$

we have $\left\|E_{n}\right\|=o\left(\lambda_{n}\right)$ and

$$
\sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\lambda_{n+1} \cdot \lambda_{n}}=J<\infty .
$$

The main difficulty in discussing the perturbation of $T$ by a bounded linear operator $B$ is to estimate $\left\|(T-\lambda I)^{-1}\right\|$ on suitable contours. For this purpose we need the following axioms:

Condition (b). There exists a real number $\delta$ such that
( $\mathrm{b}_{1}$ ) $\quad 0<\delta<\frac{1}{2} ;$

$$
\begin{align*}
& \left(\mathrm{b}_{2}\right) \quad \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\lambda_{n+1} \cdot \lambda_{n}^{\delta}}=K<\infty ;  \tag{2}\\
& \left(\mathrm{b}_{3}\right) \quad \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\left|\mu_{r}-\lambda_{n+1}\right| \cdot\left|\mu_{r}-\lambda_{n}\right|^{\delta}}=H_{r}<H<\infty,
\end{align*}
$$

where $\mu_{r}=\frac{1}{2}\left(\lambda_{r}+\lambda_{r+1}\right)$ and $H$ is independent of $r$.
The following condition simplifies the discussion of the relation of the new eigenvalues to the old.

Condition (c). There exists a number $l>1$ such that, for $n>1$,

$$
\begin{equation*}
\Delta \lambda_{n-1}<l \Delta \lambda_{n} \tag{2.1}
\end{equation*}
$$

In discussing the perturbation of $T$ by an operator $S$ which is not necessarily bounded but whose domain includes that of $T^{\nu}$, for some $\nu$ with

$$
\begin{equation*}
0<\nu<1 \tag{2.2}
\end{equation*}
$$

we must estimate not only $\left\|R_{\lambda}\right\|$ but also $\left\|T^{\nu} R_{\lambda}\right\|$. We therefore require the following additional axioms:

Condition (d). For some choice of $\nu$ and $\delta$ (satisfying (b) and (2.2)) there exists a number $\tau$ such that
( $\mathrm{d}_{1}$ ) $\quad 0<\tau<1-\delta ;$
( $\left.\mathrm{d}_{2}\right) \quad \sum_{1}^{\infty}\left\|E_{n}\right\| \cdot\left|\Delta\left(\lambda_{n}^{-\nu}\right)\right|=P<\infty ; \quad\left\|E_{n}\right\|=o\left(\lambda_{n}^{\nu}\right) ;$
( $\mathrm{d}_{3}$ ) $\quad \sum_{1}^{\infty} \frac{\lambda_{n}^{\nu}\left\|E\left(\lambda_{n}\right)\right\|}{\lambda_{n}^{\tau}}=G<\infty$;
( $\left.\mathrm{d}_{4}\right) \quad \sum_{i}^{\infty} \frac{\lambda_{n}^{\nu}\left\|E\left(\lambda_{n}\right)\right\|}{\left|\mu_{r}-\lambda_{n}\right|^{\tau}}=F_{r}<F<\infty$,
where $F$ is independent of $r$.
By $\left(a_{1}\right)$ and ( $a_{2}$ ) the following contours pass through no point $\lambda_{n}$ :
The contour $\Gamma_{r}(r \geqq 1)$ is the square with centre the origin and sides parallel to the axes, whose right-hand vertical side (written $V_{r}$ ) passes through

$$
\begin{equation*}
\mu_{r}=\frac{1}{2}\left(\lambda_{r}+\lambda_{r+1}\right) . \tag{2.3}
\end{equation*}
$$

The contour formed by the upper, left-hand, and lower sides of $\Gamma_{r}$ will be written $L_{r}$.

The contour $\Omega_{r}$ (defined if (c) holds and $r>1$ ) is a circle with centre $\lambda_{r}$ and a radius $\rho_{r}$ satisfying the condition

$$
\begin{equation*}
\rho_{r} \leqq \Delta \lambda_{r-1} / 2 l \tag{2.4}
\end{equation*}
$$

Thus by (c), $\rho_{r} \leqq \inf \left\{\frac{1}{2} \Delta \lambda_{r-1}, \frac{1}{2} \Delta \lambda_{r}\right\}$.
Lemma 2.1. Let $\left(\mathrm{a}_{1}\right)$ and $\left(\mathrm{a}_{2}\right)$ be satisfied. Then
(i) there exists a number $k>0$ such that for all $r$, all $\lambda \in L_{r}$, and all $n$,

$$
\left|\lambda-\lambda_{n}\right|>k \lambda_{n}, \quad\left|\lambda-\lambda_{n}\right|>\frac{1}{2} \mu_{r}
$$

(ii) for all $r$, all $\lambda$ on $V_{r}\left(\right.$ say $\left.\lambda=\mu_{r}+i \eta\right)$ and all $n$,

$$
\left|\lambda-\lambda_{n}\right| \geqq\left|\mu_{r}-\lambda_{n}\right| \geqq \frac{1}{2} \Delta \lambda_{r}, \quad\left|\lambda-\lambda_{n}\right|>|\eta| .
$$

Proof. The proof is obvious from ( $a_{1}$ ).
Lemma 2.2. Let (a) and (c) be satisfied. Then for all $r>1$, all $\lambda$ on $\Omega_{r}$, and all $n \neq r$ we have

$$
\left|\lambda-\lambda_{n}\right| \geqq \Delta \lambda_{r-1} / 3 l, \quad\left|\lambda-\lambda_{n}\right| \geqq\left|\mu_{r-1}-\lambda_{n}\right| / 3 l .
$$

Proof. By (c) we have

$$
\begin{aligned}
(1 / 2 l) \Delta \lambda_{r-1} & \leqq \inf \left(\frac{1}{2} \Delta \lambda_{r-1}, \frac{1}{2} \Delta \lambda_{r}\right) \\
& \leqq \inf \left(\frac{1}{2}\left|\lambda-\lambda_{r-1}\right|, \frac{1}{2}\left|\lambda-\lambda_{r+1}\right|\right) \\
& \leqq\left|\lambda-\lambda_{n}\right|,
\end{aligned}
$$

by the definition of $\Omega_{r}$ and $\left(\mathrm{a}_{2}\right)$.
If $n<r$ we clearly have $\left|\lambda-\lambda_{n}\right|>\left|\mu_{r-1}-\lambda_{n}\right|$. If $n>r$,

$$
\begin{aligned}
\left|\mu_{r-1}-\lambda_{n}\right| & \leqq\left|\mu_{r-1}-\lambda_{r}\right|+\left|\lambda_{r}-\lambda\right|+\left|\lambda-\lambda_{n}\right| \\
& \leqq \frac{1}{2} \Delta \lambda_{r-1}+\frac{1}{2} \Delta \lambda_{r}+\left|\lambda-\lambda_{n}\right| \\
& \leqq \frac{1}{2}(l+1) \Delta \lambda_{r}+\left|\lambda-\lambda_{n}\right| \\
& \leqq 3 l\left|\lambda-\lambda_{n}\right| .
\end{aligned}
$$

## 3. Operational calculus

Let $\mathfrak{B}_{0}$ be the normed linear space consisting of finite linear combinations of eigenvectors of $T$. Let $T_{0}$ be the restriction of $T$ to $\mathfrak{B}_{0}$. We will set up a natural operational calculus for $T_{0}$ and extend it (subject to ( $a_{4}$ )) to an operational calculus for $T$.

We write the eigenvalues of $T$ as $\left(\lambda_{i}\right)_{i \geqq 1}$ and write $E_{0}\left(\lambda_{i}\right)$ for the projection of $\mathfrak{B}_{0}$ onto the set of eigenvectors of $\lambda_{i}$ parallel to the eigenvectors of all other eigenvalues. If $f(\lambda)$ is any function defined on the set $\left(\lambda_{i}\right)_{i \geqq 1}$ we may now define

$$
\begin{equation*}
f\left(T_{0}\right)=\sum_{1}^{\infty} f\left(\lambda_{i}\right) E_{0}\left(\lambda_{i}\right) . \tag{3.1}
\end{equation*}
$$

If the functions $f\left(\lambda_{i}\right)=1,1 /\left(\lambda_{i}-\mu\right), \lambda_{i}, \lambda_{i}^{\nu}, \lambda_{i}^{-\nu}, \lambda_{i}^{\nu}\left(\lambda_{i}-\mu\right)^{-1}$, are defined, we write the corresponding operators $f\left(T_{0}\right)=I_{0},\left(T_{0}-\mu I_{0}\right)^{-1}, T_{0}$, $T_{0}^{\nu}, T_{0}^{-\nu}, T_{0}^{\nu}\left(T_{0}-\mu I_{0}\right)^{-1}$.

Theorem 3.1. For the correspondence $f(\lambda) \rightarrow f\left(T_{0}\right)$ defined above,
(i) $\alpha f(\lambda)+\beta g(\lambda) \rightarrow \alpha f\left(T_{0}\right)+\beta g\left(T_{0}\right)$, for all scalars $\alpha, \beta$;
(ii) $f(\lambda) g(\lambda) \rightarrow f\left(T_{0}\right) g\left(T_{0}\right)$;
(iii) $f\left(T_{0}\right)$ has an inverse (not necessarily bounded) on $\mathfrak{B}_{0}$ to $\mathfrak{B}_{0}$ if and only if $f\left(\lambda_{i}\right)$ never vanishes, in which case $1 / f(\lambda) \rightarrow f\left(T_{0}\right)^{-1}$;
(iv) if $E_{0}\left(\lambda_{i}\right)$ is bounded for all $i$ and

$$
\begin{equation*}
\sum_{1}^{\infty}\left|f\left(\lambda_{i}\right)\right| \cdot\left\|E_{0}\left(\lambda_{i}\right)\right\|<\infty, \tag{3.2}
\end{equation*}
$$

the series (3.1) converges in operator norm to $f\left(T_{0}\right)$ (which is thus bounded);
(v) if $\left|f\left(\lambda_{i}\right)\right| \cdot\left\|E_{0 i}\right\| \rightarrow 0$ as $i \rightarrow \infty$ and $E_{0 i}$ is bounded for all $i$ and

$$
\begin{equation*}
\sum_{1}^{\infty}\left|\Delta f\left(\lambda_{n}\right)\right| \cdot\left\|E_{0 n}\right\|<\infty \tag{3.3}
\end{equation*}
$$

where $E_{0 n}=\sum_{1}^{n} E_{0}\left(\lambda_{i}\right)$, then the series

$$
\begin{equation*}
-\sum_{1}^{\infty} E_{0 n} \Delta f\left(\lambda_{n}\right) \tag{3.4}
\end{equation*}
$$

converges to $f\left(T_{0}\right)$ in operator norm.

Proof. Since eigenvectors of distinct eigenvalues are linearly independent, (i), (ii), and (iii) are obvious. If the conditions of (iv) are satisfied, we have

$$
\begin{aligned}
\left\|f\left(T_{0}\right)-\sum_{1}^{N} f\left(\lambda_{i}\right) E_{0}\left(\lambda_{i}\right)\right\| & =\sup _{\|x\|=1, x e ®_{0}}\left\|\sum_{N+1}^{\infty} f\left(\lambda_{i}\right) E_{0}\left(\lambda_{i}\right) x\right\| \\
& \leqq \sum_{N+1}^{\infty}\left|f\left(\lambda_{i}\right)\right| \cdot\left\|E_{0}\left(\lambda_{i}\right)\right\| \\
& \rightarrow 0 \quad \text { as } n \rightarrow \infty
\end{aligned}
$$

This proves (iv), and the proof of (v) is similar. For if we write $x=\sum_{1}^{m} E_{0}\left(\lambda_{i}\right) x$ with $m=m(x, N)>N$ we have

$$
\begin{aligned}
f\left(T_{0}\right) x & =\sum_{1}^{m} f\left(\lambda_{i}\right) E\left(\lambda_{i}\right) x \\
& =\sum_{1}^{N} E_{0 i} \Delta f\left(\lambda_{i}\right) x-\sum_{N+1}^{m-1} E_{0 i} \Delta f\left(\lambda_{i}\right) x+f\left(\lambda_{m}\right) E_{0 m} x .
\end{aligned}
$$

The following lemma contains our definition of $f(T)$.
Lemma 3.1. Let $\left(\mathrm{a}_{4}\right)$ be satisfied. Let $f\left(T_{0}\right)$ be bounded. Then the closure $\overline{f\left(T_{0}\right)}$ of $f\left(T_{0}\right)$ is the unique bounded linear operator $f(T)$ on $\mathfrak{B}$ to $\mathfrak{B}$ whose restriction to $\mathfrak{B}_{0}$ is $f\left(T_{0}\right)$.

Proof. The proof is obvious.
Lemma 3.2. Let (a) be satisfied. Then
(i) $\overline{E_{0}\left(\lambda_{i}\right)}=E\left(\lambda_{i}\right) ; \quad \overline{E_{0 i}}=E_{i} ;$
(ii) if $\lambda_{\ell}\left(\lambda_{n}\right)_{n \geqq 1}$, then

$$
(T-\lambda I)^{-1}=\sum_{1}^{\infty} \frac{E_{n} \Delta \lambda_{n}}{\left(\lambda-\lambda_{n+1}\right)\left(\lambda-\lambda_{n}\right)},
$$

with convergence in operator norm;
(iii) $\sigma(T)=\pi(T)=\left(\lambda_{n}\right)_{n \geqq 1}$;
(iv) $\overline{\left(T_{0}-\lambda I_{0}\right)^{-1}}=(T-\lambda I)^{-1}$;
(v) for $r>1$,

$$
E_{r}=\frac{-1}{2 \pi i} \int_{\Gamma_{r}}(T-\lambda I)^{-1} d \lambda, \quad E\left(\lambda_{r}\right)=\frac{-1}{2 \pi i} \int_{\Omega_{r}}(T-\lambda I)^{-1} d \lambda
$$

Proof. (i) It is obvious from the definitions of $E_{0}\left(\lambda_{i}\right)$ (see p. 87) and $E\left(\lambda_{i}\right)$ (see ( $a_{5}$ ) that $E\left(\lambda_{i}\right)$ is bounded and $E_{0}\left(\lambda_{i}\right)$ is the restriction of $E\left(\lambda_{i}\right)$ to $\mathfrak{B}_{0}$. The analogous facts about $E_{i}$ and $E_{0 i}$ are proved in the same way. Thus Lemma 3.1 gives both results.
(ii) By $\left(a_{5}\right)$ the series converges to a bounded linear operator, $S_{\lambda}$ say, on $\mathfrak{B}$ to $\mathfrak{B}$. By Theorem 3.1 (ii), (iii), and (iv), the restriction of $S_{\lambda}$ to $\mathfrak{B}_{0}$ is $\left(T_{0}-\lambda I_{0}\right)^{-1}$. Thus by Lemma 3.1,

$$
S_{\lambda}=\overline{\left(T_{0}-\lambda I_{0}\right)^{-1}}
$$

so that as $T$ is closed $S_{\lambda}$ is contained in $(T-\lambda I)^{-1}$. Since $(T-\lambda I)^{-1}$ is single-valued by $\left(a_{1}\right)$, and $\mathfrak{D}\left(S_{\lambda}\right)=\mathfrak{B}, \quad S_{\lambda}=(T-\lambda I)^{-1}$. This proves (ii), (iii), and (iv).
(v) By $\left(\mathrm{a}_{1}\right)$ and $\left(\mathrm{a}_{2}\right), \Gamma_{r}$ contains $\lambda_{1}, \lambda_{2}, \cdots, \lambda_{r}$, and the set $\left(\lambda_{n}\right)_{n>r}$ lies outside $\Gamma_{r}$. Since ( $\mathrm{a}_{2}$ ) and ( $\mathrm{a}_{5}$ ) imply that the series in (ii) converges uniformly on $\Gamma_{r}$, we have

$$
\begin{aligned}
\frac{-1}{2 \pi i} \int_{\Gamma_{r}}(T-\lambda I)^{-1} d \lambda & =\frac{-1}{2 \pi i} \int_{\Gamma_{r}} \sum_{1}^{\infty} \frac{E_{n} \Delta \lambda_{n}}{\left(\lambda-\lambda_{n+1}\right)\left(\lambda-\lambda_{n}\right)} d \lambda \\
& =\sum_{1}^{\infty} E_{n} \Delta \lambda_{n}\left(\frac{-1}{2 \pi i} \int_{\Gamma_{r}} \frac{d \lambda}{\left(\lambda-\lambda_{n+1}\right)\left(\lambda-\lambda_{n}\right)}\right) \\
& =E_{r}
\end{aligned}
$$

The discussion for $\Omega_{r}$ is similar.
Lemma 3.3. Let (a) and (b) be satisfied. Then
(i) if $\lambda$ lies on $L_{r}$,

$$
\left\|(T-\lambda I)^{-1}\right\|=O\left(\left|L_{r}\right|^{-1+\delta}\right)
$$

independently of $\lambda$;
(ii) if $\lambda$ lies on $V_{r}$, say $\lambda=\mu_{r}+i \eta$,

$$
\left\|(T-\lambda I)^{-1}\right\|=O\left(|\eta|^{-1+\delta}\right), \quad\left\|(T-\lambda I)^{-1}\right\|=O\left(\left(\Delta \lambda_{r}\right)^{-1+\delta}\right)
$$

independently of $\lambda$;
(iii) if (c) holds and $\lambda$ lies on the circle $\Omega_{r}$,

$$
\left\|(T-\lambda I)^{-1}\right\|=O\left(\left(\Delta \lambda_{r-1}\right)^{-1+\delta}+\frac{\left\|E_{r-1}\right\|+\left\|E_{r}\right\|}{\rho_{r}}\right)
$$

independently of $\lambda$ and $\rho_{r}$.
Proof. (i) By Lemmas 3.2 (ii) and 2.1,

$$
\begin{aligned}
\left\|(T-\lambda I)^{-1}\right\| & \leqq \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\left|\lambda-\lambda_{n+1}\right| \cdot\left|\lambda-\lambda_{n}\right|} \\
& \leqq \frac{1}{\inf \left|\lambda-\lambda_{n}\right|^{1-\delta}} \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\left|\lambda-\lambda_{n+1}\right| \cdot\left|\lambda-\lambda_{n}\right|^{\delta}} \\
& \leqq \frac{2^{1-\delta} k^{-1-\delta}}{\mu_{r}^{1-\delta}} \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\lambda_{n+1} \lambda_{n}^{\delta}} \\
& =O\left(\mu_{r}^{-1+\delta}\right),
\end{aligned}
$$

by ( $\mathrm{b}_{2}$ ). And $\left|L_{r}\right|=6 \mu_{r}$ by the definition of $L_{r}$.
(ii) To the second line of the proof of (i) apply Lemma 2.1 (ii) to obtain

$$
\begin{aligned}
\left\|(T-\lambda I)^{-1}\right\| & \leqq \frac{1}{\sup \left(|\eta|, \frac{1}{2} \Delta \lambda_{r}\right)^{1-\delta}} \sum_{1}^{\infty} \frac{\left\|E_{n}\right\| \Delta \lambda_{n}}{\left|\mu_{r}-\lambda_{n+1}\right| \cdot\left|\mu_{r}-\lambda_{n}\right|^{\delta}} \\
& \leqq \frac{H}{\sup \left(|\eta|, \frac{1}{2} \Delta \lambda_{r}\right)^{1-\delta}},
\end{aligned}
$$

by ( $\mathrm{b}_{3}$ ).
(iii) Lemma 3.2 (ii) gives the series

$$
\begin{aligned}
\sum_{1}^{r-2} \frac{E_{n} \Delta \lambda_{n}}{\left(\lambda-\lambda_{n+1}\right)\left(\lambda-\lambda_{n}\right)}- & \frac{E_{r-1}}{\lambda-\lambda_{r-1}} \\
& +\frac{E_{r-1}-E_{r}}{\lambda-\lambda_{r}}+\frac{E_{r}}{\lambda-\lambda_{r+1}}+\sum_{r+1}^{\infty} \frac{E_{n} \Delta \lambda_{n}}{\left(\lambda-\lambda_{n+1}\right)\left(\lambda-\lambda_{n}\right)}
\end{aligned}
$$

for $(T-\lambda I)^{-1}$. Thus by Lemma 2.2 and the relations $\left|\lambda-\lambda_{r}\right|=\rho_{r}$, $\left|\lambda-\lambda_{r \pm 1}\right| \geqq \rho_{r}$ (see the definition of $\Omega_{r}$ ), $\left\|(T-\lambda I)^{-1}\right\|$

$$
\begin{aligned}
& \leqq\left(\sum_{1}^{r-2}+\sum_{r+1}^{\infty}\right) \frac{\left\|E_{n}\right\| \Delta \lambda_{n}(3 l)^{2}}{\left|\mu_{r-1}-\lambda_{n+1}\right| \cdot\left|\mu_{r-1}-\lambda_{n}\right|^{\delta}\left(\Delta \lambda_{r-1}\right)^{1-\delta}}+2 \frac{\left\|E_{r-1}\right\|+\left\|E_{r}\right\|}{\rho_{r}} \\
& \leqq \frac{(3 l)^{2} H_{r-1}}{\left(\Delta \lambda_{r-1}\right)^{2}}+2 \frac{\left\|E_{r-1}\right\|+\left\|E_{r}\right\|}{\rho_{r}}
\end{aligned}
$$

by ( $b_{3}$ ).
Lemma 3.4. If (d) is satisfied, $T_{0}^{-\nu}$ is bounded.
Proof. Apply Theorem 3.1 (v) to ( $\mathrm{d}_{2}$ ).
We will write $T^{-\nu}$ for $\overline{T_{0}^{-\nu}}$ (cf. Lemma 3.1). Following Kramer [2] we prove
Lemma 3.5. If $S$ is a closed linear operator in $\mathfrak{B}$, if (d) is satisfied, and if $\mathfrak{D}(S) \supseteqq \Re\left(T^{-\nu}\right)$, then $S T^{-\nu}$ is bounded.

Proof. $S T^{-\nu}$ is closed and $\mathfrak{D}\left(S T^{-\nu}\right)=\mathfrak{B}$ (by the previous lemma) so that $S T^{-\nu}$ is bounded by the Closed Graph Theorem.

Lemma 3.6. Let (a), (b), and (d) be satisfied. Let $\mathfrak{D}(S) \supseteq \Re\left(T^{-\nu}\right)$. Then (i) if $\lambda$ lies on $L_{r}$,

$$
\left\|S(T-\lambda I)^{-1}\right\|=O\left(\left|L_{r}\right|^{-1+\tau}\right)
$$

independently of $\lambda$;
(ii) if $\lambda$ lies on $V_{r}$, say $\lambda=\mu_{r}+i \eta$

$$
\left\|S(T-\lambda I)^{-1}\right\|=O\left(\left(\Delta \lambda_{r}\right)^{-1+\tau}\right), \quad\left\|S(T-\lambda I)^{-1}\right\|=O\left(|\eta|^{-1+\tau}\right)
$$

independently of $\lambda$;
(iii) if (c) holds and $\lambda$ lies on $\Omega_{r}$,

$$
\left\|S(T-\lambda I)^{-1}\right\|=O\left(\left(\Delta \lambda_{r-1}\right)^{-1+\tau}+\rho_{r}^{-1}\left\|E\left(\lambda_{r}\right)\right\|\right)
$$

independently of $\lambda$ and $\rho_{r}$.
Proof. Since $S(T-\lambda I)^{-1}$ is closed, an argument on the lines of Lemma 3.1 shows that its order is the same as that of $S\left(T_{0}-\lambda I_{0}\right)^{-1}$, if the latter is bounded. By Theorem 3.1,

$$
S\left(T_{0}-\lambda I_{0}\right)^{-1}=S T_{0}^{-\nu} T_{0}^{\nu}\left(T_{0}-\lambda I_{0}\right)^{-1}
$$

so that by Lemma 3.5 it is sufficient to estimate the order of

$$
\left\|T_{0}^{\nu}\left(T_{0}-\lambda I_{0}\right)^{-1}\right\|=\left\|\sum_{1}^{\infty} \frac{\lambda_{n}^{\nu} E_{0}\left(\lambda_{n}\right)}{\lambda_{n}-\lambda}\right\|
$$

(i) Using Lemma 2.1 (i) we obtain

$$
\left\|T_{0}^{\nu}\left(T_{0}-\lambda I_{0}\right)^{-1}\right\| \leqq \frac{2^{1-\tau}}{\mu_{r}^{1-\tau}} \sum_{1}^{\infty} \frac{\lambda_{n}^{\nu}\left\|E_{0}\left(\lambda_{n}\right)\right\|}{k^{\tau} \lambda_{n}^{\tau}} \leqq \frac{G 2^{1-\tau}}{k^{\tau} \mu_{r}^{1-\tau}}
$$

by $\left(d_{3}\right)$.
(ii) Using Lemma 2.1 (ii) we obtain

$$
\begin{aligned}
\left\|T_{0}^{\nu}\left(T_{0}-\lambda I_{0}\right)^{-1}\right\| & \leqq \frac{1}{\sup \left(|\eta|, \frac{1}{2} \Delta \lambda_{r}\right)^{1-\tau}} \sum_{1}^{\infty} \frac{\lambda_{n}^{\nu}\left\|E\left(\lambda_{n}\right)\right\|}{\left|\mu_{r}-\lambda_{n}\right|^{\tau}} \\
& \leqq \frac{F}{\sup \left(|\eta|, \frac{1}{2} \Delta \lambda_{r}\right)^{1-\tau}}
\end{aligned}
$$

by $\left(d_{4}\right)$.
(iii) Using Lemma 2.2,

$$
\begin{align*}
\left\|T_{0}^{\nu}\left(T_{0}-\lambda I_{0}\right)^{-1}\right\| & \leqq \frac{3}{\left(\Delta \lambda_{r-1}\right)^{1-\tau}}\left(\sum_{1}^{r-1}+\sum_{r+1}^{\infty}\right) \frac{\lambda_{n}^{\nu}\left\|E\left(\lambda_{n}\right)\right\|}{\left|\mu_{r-1}-\lambda_{n}\right|^{\tau}}+\frac{\left\|E\left(\lambda_{r}\right)\right\|}{\rho_{r}} \\
& \leqq \frac{3 l F}{\left(\Delta \lambda_{r-1}\right)^{1-\tau}}+\frac{\left\|E\left(\lambda_{r}\right)\right\|}{\rho_{r}} \tag{r>1}
\end{align*}
$$

by $\left(d_{4}\right)$.

## 4. The perturbation theorems

Theorem 4.1. Let $T$ satisfy (a) and (b). Let $\Gamma_{r}(r \geqq 1)$ be the contour defined on $p$. 86. Let $S$ be any bounded linear operator on $\mathfrak{B}$ into $\mathfrak{B}$. Then
(i) $\sigma(T+S)$ consists of a sequence of eigenvalues $\nu_{i}(i \geqq k+1) o f^{2}$ generalised multiplicity $m_{i}<\infty$;
and for $r$ sufficiently large
(ii) $\sum_{k+1}^{r} m_{i}=r, m_{r}=1$ and $\nu_{r}$ lies between $\Gamma_{r-1}$ and $\Gamma_{r}$ (and if (c) is satisfied and $\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}=o\left(\Delta \lambda_{r-1}\right)$,

$$
\begin{equation*}
\left.\nu_{r}=\lambda_{r}+O\left(\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}\right)\right) \tag{4.1}
\end{equation*}
$$

(iii) the projection

$$
E_{r}^{\prime}=\frac{-1}{2 \pi i} \int_{\Gamma_{r}}(T+S-\lambda I)^{-1} d \lambda
$$

exists;
(iv) there are sets of linearly independent vectors $\left(\Phi_{p q}\right)_{p \geqq k+1,1 \leqq q \leqq m_{p}}$ in $\mathfrak{B}$

[^1]and $\left(\Psi_{i j}\right)_{i \geqq k+1,1 \leqq j \leqq m_{i}}$ in $\mathfrak{B}^{*}$ such that $\left(T+S-\nu_{p} I\right)^{m_{p}} \Phi_{p q} \equiv 0$ and $\Psi_{i j}\left(\Phi_{p q}\right) \equiv \delta_{i p} \delta_{j q}$, and
$$
E_{r}^{\prime}=\sum_{p \leqq r} \Psi_{p q}(\cdot) \Phi_{p q}
$$
(v) $\left\|E_{r}^{\prime}-E_{r}\right\| \rightarrow 0$.

Proof. (iii) Write $R_{\lambda}$ for $(T-\lambda I)^{-1}$. By ( $\mathrm{a}_{2}$ ) and Lemma 3.3, $\left\|S R_{\lambda}\right\|<\frac{1}{2}$ for $\lambda$ on $\Gamma_{r}=L_{r} \cup V_{r}$, for $r$ sufficiently large. Thus $\left(I+S R_{\lambda}\right)^{-1}$ exists and

$$
\begin{align*}
\left\|\left(I+S R_{\lambda}\right)^{-1}\right\| & =\left\|I-\left(S R_{\lambda}\right)+\left(S R_{\lambda}\right)^{2}-\cdots\right\| \\
& \leqq 1+\frac{1}{2}+\left(\frac{1}{2}\right)^{2}+\cdots=2 \tag{4.2}
\end{align*}
$$

With Lemma 3.2 (iii) this proves the existence on $\Gamma_{r}$ of the right-hand side of the equation

$$
\begin{equation*}
(T+S-\lambda I)^{-1}=R_{\lambda}-R_{\lambda}\left(I+S R_{\lambda}\right)^{-1} S R_{\lambda} \tag{4.3}
\end{equation*}
$$

which can now be verified by multiplying the right-hand side (on the left and on the right) by $T+S-\lambda I$. Thus $(T+S-\lambda I)^{-1}$ exists on $\Gamma_{r}$. This proves (iii).
(v) By (iii), Lemma 3.2 (v), equations (4.3) and (4.2), and Lemma 3.3,

$$
\begin{aligned}
\left\|E_{r}^{\prime}-E_{r}\right\| & =\frac{1}{2 \pi}\left\|\int_{\Gamma_{r}}\left((T+S-\lambda I)^{-1}-(T-\lambda I)^{-1}\right) d \lambda\right\| \\
& \leqq \frac{1}{2 \pi} \int_{\Gamma_{r}}\|S\| \cdot\left\|R_{\lambda}\right\|^{2} \cdot\left\|\left(I+S R_{\lambda}\right)^{-1}\right\| \cdot|d \lambda| \\
& \leqq \frac{\|S\|}{\pi} \int_{\Gamma_{r}}\left\|R_{\lambda}\right\|^{2}|d \lambda| \\
& =\frac{\|S\|}{\pi}\left\{\int_{L_{r}}\left\|R_{\lambda}\right\|^{2}|d \lambda|+\int_{-\Delta \lambda_{r}}^{\Delta \lambda_{r}}\left\|R_{\mu_{r}+i \eta}\right\|^{2} d \eta\right. \\
& \left.=O\left(\mid L_{r_{r}}^{-\Delta \lambda_{r}}+\int_{\Delta \lambda_{r}}^{\mu_{r}}\right)\left\|R_{\mu_{r}+i \eta}\right\|^{2} d \eta\right\} \\
& =O\left(\left(\Delta \lambda_{r}\right)^{-1+2 \delta}\right),
\end{aligned}
$$

since

$$
\int_{\Delta \lambda_{r}}^{\mu_{r}} \eta^{-2+2 \delta} d \eta=\left[(-1+2 \delta)^{-1} \eta^{-1+2 \delta}\right]_{\Delta \lambda_{r}}^{\mu_{r}}
$$

The result follows from ( $a_{2}$ ) and ( $b_{1}$ ).
(i), (ii), (iv). By (v) and Lemma 1.3, $E_{r}$ and $E_{r}^{\prime}$ project onto subspaces of the same dimension-which by $\left(a_{5}\right)$ is $r$-for all $r$ sufficiently large. Thus Lemma 1.2 proves (ii) (apart from the bracket), (i), and the existence of a linearly independent set $\left(\Phi_{p q}\right)_{k+1 \leq p, 1 \leqq q \leqq m_{p}}$ such that $\left(T+S-\nu_{p} I\right)^{m_{p}} \Phi_{p q} \equiv 0$ and the $\Phi_{p q}$ with $p \leqq r$ span $E_{r}^{\prime} \mathfrak{B}$. We can write any $x \in E_{r}^{\prime} \mathfrak{B}$ in the form
$\sum_{p \leqq r} x_{p q} \Phi_{p q}$ where the $x_{p q}$ are uniquely determined scalars. Since $\operatorname{dim} E_{r}^{\prime} \mathfrak{B}<\infty$, the mappings $x \rightarrow x_{p q}$ are bounded linear functionals on $E_{r}^{\prime} \mathfrak{B}$. Thus the linear functional $\Psi_{p q}$ on $\mathfrak{B}$ defined by

$$
\Psi_{p q}(x)=\left(E_{r}^{\prime} x\right)_{p q}
$$

for some chosen $r$ greater than $p$ has the required properties, for the following argument shows that the definition is independent of the choice of $r$. Take $p \leqq r \leqq r_{0}$; then

$$
E_{r}^{\prime} x=E_{r}^{\prime}\left(E_{r_{0}}^{\prime} x\right)=E_{r}^{\prime} \sum_{p \leqq r_{0}}\left(E_{r_{0}}^{\prime} x\right)_{p q} \Phi_{p q}=\sum_{p \leqq r}\left(E_{r_{0}}^{\prime} x\right)_{p q} \Phi_{p q}
$$

by Lemmas 1.1 and 1.2.
(ii) Let (c) be satisfied, and let $\rho_{r}$ (subject to (2.4)) have any order strictly higher than $\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}$. We can obtain (4.2) and (4.3) as before, but now with $\lambda$ on the circle $\Omega_{r}$ of centre $\lambda_{r}$ and radius $\rho_{r}$. Thus if we write

$$
E^{\prime}(r)=\frac{-1}{2 \pi i} \int_{\Omega_{r}}(T+S-\lambda I)^{-1} d \lambda
$$

we obtain, by the argument for (v), and (2.4), ( $a_{2}$ ) and ( $b_{1}$ ),

$$
\begin{aligned}
\left\|E^{\prime}(r)-E\left(\lambda_{r}\right)\right\| & \leqq \frac{\|S\|}{\pi} \int_{\Omega_{r}}\left\|R_{\lambda}\right\|^{2}|d \lambda| \\
& =O\left(\rho_{r}\left(\Delta \lambda_{r-1}\right)^{-2+2 \delta}+\frac{\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}}{\rho_{r}}\right) \\
& \rightarrow 0
\end{aligned}
$$

Thus by Lemma 1.3, $E^{\prime}(r)$ and $E\left(\lambda_{r}\right)$ project onto subspaces of the same di-mension-which by $\left(a_{5}\right)$ is one-for all $r$ sufficiently large. Thus $\Omega_{r}$ contains an eigenvalue of $T+S$, which must be $\nu_{r}$. In other words,

$$
\nu_{r}=\lambda_{r}+O\left(\rho_{r}\right)
$$

whenever $\rho_{r}$ has a strictly higher order than $\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}$. Thus

$$
\nu_{r}=\lambda_{r}+O\left(\left\|E_{r-1}\right\|^{2}+\left\|E_{r}\right\|^{2}\right)
$$

Corollary 4.1. Let $T$ and $S$ satisfy the conditions of the theorem, and let the vectors $\phi_{j}, \psi_{i}, \Phi_{p q}, \Psi_{p q}$ be those defined in ( $\mathrm{a}_{3}$ ) and part (iv) of the theorem. Let $x$ be a vector in $\mathfrak{B}$ such that

$$
\sum_{1}^{\infty} \psi_{i}(x) \phi_{i}=x
$$

in the sense of strong or weak convergence or strong or weak ( $C, \theta$ ) summability. Then

$$
\sum_{p=k+1}^{\infty} \sum_{q=1}^{m_{p}} \Psi_{p q}(x) \Phi_{p q}=x
$$

in the same sense.
Proof. The proof is obvious from part (v) of the theorem (with the definitions in ( $\mathrm{a}_{5}$ ) and part (iv) of the theorem).

Theorem 4.2. Let $T$ satisfy (a), (b) and (d) (for some $\nu, 0<\nu<1$ ). Let $S$ be a finite sum of operators of the form $B_{\alpha} S_{\alpha}$, where $B_{\alpha}$ is a bounded linear operator on $\mathfrak{B}$ into $\mathfrak{B}$ and $S_{\alpha}$ is a closed linear operator in $\mathfrak{B}$ such that $\mathfrak{D}\left(S_{\alpha}\right) \geqq \Re\left(T^{-\nu}\right) \cup \mathfrak{D}(T)$. Then the conclusions (i) to (v) of Theorem 4.1 are true if we insert before (4.1) the words,

$$
\text { "and }\left\|E_{r-1}\right\|+\left\|E_{r}\right\|=o\left(\Delta \lambda_{r-1}\right)^{1-\tau} "
$$

Proof. We can apply Lemma 3.6, with $S_{\alpha}$ replacing $S$.
(iii) Use the proof of Theorem 4.1 (iii), with Lemma 3.6 instead of Lemma 3.3.
(v) By (iii), Lemma 3.2 (v), equations (4.3) and (4.2), and Lemmas 3.3 and 3.6,

$$
\begin{aligned}
&\left\|E_{r}^{\prime}-E_{r}\right\|= \frac{1}{2 \pi}\left\|\int_{\Gamma_{r}}\left((T+S-\lambda I)^{-1}-(T-\lambda I)^{-1}\right) d \lambda\right\| \\
& \leqq \frac{1}{2 \pi} \int_{\Gamma_{r}}\left\|R_{\lambda}\right\| \cdot\left\|\left(I+S R_{\lambda}\right)^{-1}\right\| \cdot\left\|S R_{\lambda}\right\| \cdot|d \lambda| \\
& \leqq \frac{1}{\pi}\left\{\int_{L_{r}}\left\|R_{\lambda}\right\| \cdot\left\|S R_{\lambda}\right\| \cdot|d \lambda|+\int_{-\Delta \lambda_{r}}^{\Delta \lambda_{r}}\left\|R_{\mu_{r}+i \eta}\right\| \cdot\left\|S R_{\mu_{r}+i \eta}\right\| d \eta\right. \\
&\left.+\left(\int_{-\mu_{r}}^{-\Delta \lambda_{r}}+\int_{\Delta \lambda_{r}}^{\mu_{r}}\right)\left\|R_{\mu_{r}+i \eta}\right\| \cdot\left\|S R_{\mu_{r}+i \eta}\right\| d \eta\right\} \\
&= O\left\{\left|L_{r}\right| \cdot\left|L_{r}\right|^{-1+\delta}\left|L_{r}\right|^{-1+\tau}+\int_{-\Delta \lambda_{r}}^{\Delta \lambda_{r}}\left(\Delta \lambda_{r}\right)^{-1+\delta}\left(\Delta \lambda_{r}\right)^{-1+\tau} d \eta\right. \\
&\left.+\int_{\Delta \lambda_{r}}^{\mu_{r}} \eta^{-1+\delta} \eta^{-1+\tau} d \eta\right\} \\
&= O\left(\left|L_{r}\right|^{-1+\delta+\tau}\right)+O\left(\left(\Delta \lambda_{r}\right)^{-1+\delta+\tau}\right) \\
& \rightarrow 0
\end{aligned}
$$

by $\left(\mathrm{a}_{2}\right)$ and $\left(\mathrm{d}_{1}\right)$.
(i), (ii), (iv). Since the resolvent set of $T+S$ is not empty, $T+S$ is closed. Thus we may follow the discussion of the corresponding parts of Theorem 4.1.
(ii) Use the first four lines of the proof of this part of Theorem 4.1. Then we obtain, by (2.4), ( $\mathrm{a}_{2}$ ) and $\left(\mathrm{d}_{1}\right)$, and the argument for (v),

$$
\begin{aligned}
\| E^{\prime}(r)- & E\left(\lambda_{r}\right)\left\|\leqq \frac{1}{\pi} \int_{\Omega_{r}}\right\| R_{\lambda}\|\cdot\| S R_{\lambda} \| \cdot|d \lambda| \\
= & O\left(\left\{\left(\Delta \lambda_{r-1}\right)^{-1+\delta}+\frac{\left\|E_{r-1}\right\|+\left\|E_{r}\right\|}{\rho_{r}}\right\}\left\{\left(\Delta \lambda_{r-1}\right)^{-1+\tau}+\frac{\left\|E\left(\lambda_{r}\right)\right\|}{\rho_{r}}\right\} \rho_{r}\right) \\
= & O\left(\left(\Delta \lambda_{r-1}\right)^{-1+\delta+\tau}+\left(\Delta \lambda_{r-1}\right)^{-1+\tau}\left(\left\|E_{r-1}\right\|+\left\|E_{r}\right\|\right)\right. \\
& \left.\quad+\left(\Delta \lambda_{r-1}\right)^{-1+\delta}\left\|E\left(\lambda_{r}\right)\right\|+\left\|E\left(\lambda_{r}\right)\right\|\left(\left\|E_{r-1}\right\|+\left\|E_{r}\right\|\right) \rho_{r}^{-1}\right) \\
\rightarrow & 0 .
\end{aligned}
$$

We may now conclude with the last six lines of the proof of Theorem 4.1.
Corollary 4.2. The statement is verbally the same as that of Corollary 4.1.
Proof. The proof is the same as before.

## 5. Applications and extensions

(i) Let $T$ be a Sturm-Liouville operator of the simplest type, i.e., an operator of the form $-D^{2}$ with boundary conditions

$$
f(0) \cos \varphi+f^{\prime}(0) \sin \varphi=0, \quad f(1) \cos \psi+f^{\prime}(0) \sin \psi=0
$$

where $\varphi$ and $\psi$ are real numbers. The following facts are well known (see for example Titchmarsh [7]): $T$ has simple eigenvalues, and the eigenfunction expansion of any integrable function $f$ is uniformly equiconvergent with its Fourier series expansion. If the boundary conditions satisfy $\sin \varphi \sin \psi \neq 0$ $(\sin \varphi=0, \sin \psi \neq 0)(\sin \varphi \neq 0, \sin \psi=0)(\sin \varphi=\sin \psi=0)$, the term "Fourier series" means an expansion in terms of $\cos n \pi x\left(\sin \left(n+\frac{1}{2}\right) \pi x\right)$ $\left(\cos \left(n+\frac{1}{2}\right) \pi x\right)(\sin n \pi x)$. The space $C$ is the space of all continuous functions on $[0,1]$ (all which vanish at 0 ) (all which vanish at 1 ) (all which vanish at both 0 and 1). Thus if $f$ is in $\mathfrak{B}=C$ or $L^{1}\left[\mathfrak{B}=L^{p}(1<p<\infty)\right]$ its eigenfunction expansion is $(C, \theta)$ summable for all $\theta>0$ [its expansion is convergent] in the norm of $\mathfrak{B}$. Thus we have, by the Banach-Steinhaus theorem and the limitation theorem for Cesàro summability (Hardy [1], Theorem 4.6),

$$
\begin{equation*}
\left\|E_{n}\right\|=O\left(n^{\theta}\right) \quad(\theta>0) \quad\left[\left\|E_{n}\right\|=O(1)\right] \tag{5.1}
\end{equation*}
$$

where $E_{n}$ is the projection in $\mathfrak{B}$ mapping each function onto the $n^{\text {th }}$ partial sum of its Sturm-Liouville expansion. As we have also

$$
\lambda_{n} \sim \pi^{2} n^{2}, \quad \Delta \lambda_{n} \sim 2 \pi^{2} n
$$

we take $\theta$ as any number between 0 and $\frac{1}{2}$, and the only possible divergence from the axioms (a), (b), and (c) is that some eigenvalues may be negative. This is easily corrected by operating in the space $E_{N} \mathfrak{F}$, where $\lambda_{N}$ is the first positive eigenvalue. Let $T^{\prime}$ be any operator commuting with $E_{N}$ and such that $T^{\prime}$ is zero in $(I-E N) \mathfrak{B}$ and, as an operator in $E_{N} \mathfrak{B}, T^{\prime \prime}$ has a simple spectrum consisting of positive numbers less than $\lambda_{N}$. Then the operator $T-E_{N} T+T^{\prime}$ will satisfy axioms (a), (b), and (c) and so can be used instead of $T$. (Since $E_{N} \mathfrak{B}$ is finite-dimensional there exist suitable operators $T^{\prime}$, and these are bounded.)

Thus if $B$ is a bounded linear operator on $\mathfrak{B}$ to $\mathfrak{B}$ and $f$ is an arbitrary element of $\mathfrak{B}$, the expansion of $f$ in eigenvectors (and a finite number of other proper vectors, i.e., solutions of $\left.\left(T+B-\nu_{i} I\right)^{n} x=0\right)$ of $T+B$ is, by Corollary $4.1,(C, \theta)$ summable $(\theta>0)$ [convergent] in the norm of $\mathfrak{B}$. Also if we know that the $T$-expansion of some continuous function is uniformly convergent, the same will be true of its $(T+B)$-expansion.
(ii) Let $T$ be similar to a Sturm-Liouville operator, but with nonreal boundary conditions. Then by Schwartz [4], p. 439, the asymptotic behaviour of the eigenvalues will be satisfactory, but (a) the eigenvalues may be nonreal, and (b) $T$ may have some proper vectors which are not eigenvectors. Difficulty (a) could be met by adapting our proof, but the analysis would become more complicated. Difficulty (b) is easily dealt with by the device used in (i) for negative eigenvalues.
(iii) Let $T$ have the form $D^{u}$ with $u$ even $(u>2)$, and $u$ linearly independent boundary conditions. Judging by Tamarkin [5] various possibilities exist; but it seems that in many cases-for example, if $T$ is the $(u / 2)^{\text {th }}$ power of one of the Sturm-Liouville operators discussed in (i)-we will have simple real eigenvalues $\lambda_{n}$ such that

$$
\lambda_{n} \sim i^{u} \pi^{u} n^{u}, \quad \Delta \lambda_{n} \sim u i^{u} \pi^{u} n^{u-1}
$$

Thus, replacing $T$ (if $u / 2$ is odd) by $-T$, we may assume by the argument in (i) that all eigenvalues are positive. In the case of the $(u / 2)^{\text {th }}$ power of a Sturm-Liouville operator, the expansion is simply that of the SturmLiouville operator; in more general cases an equiconvergence theorem (Tamarkin [5], p. 45) reduces the expansion problem to the Fourier series case. Thus (5.1) is obtained as before, and taking

$$
\nu=(u-2) / u, \quad 1-1 / 4 u<\tau<1
$$

and $0<\delta<1-\tau$ we can verify (a), (b), (c), and (d) without difficulty. To show that $S$ can have the form $B_{u-2} D^{u-2}+\cdots+B_{1} D+B_{0}$ for arbitrary bounded linear operators $B_{i}$ on $\mathfrak{B}$ to $\mathfrak{B}$, it is thus sufficient by Lemma 1.4 to show that $\Re\left(\overline{T_{0}^{-\nu}}\right)$ consists of $(u-2)$-times differentiable functions $f$ with $f^{(u-2)}$ in $\mathfrak{B}$. Using the argument of Kramer [2], I can obtain this result but only where $T$ is the $(u / 2)^{\text {th }}$ power of a Sturm-Liouville operator $R$ (with eigenvalues $\Lambda_{i}=\lambda_{i}^{2 / u}$ ). Clearly $R$ and $T$ have the same normalised eigenvectors $\left(f_{i}\right)_{i \geqq 1}$. Thus

$$
\begin{aligned}
T_{0} & =\sum \lambda_{i} f_{i}(\cdot) f_{i} \\
T_{0}^{\nu} & =\sum \lambda_{i}^{\nu} f_{i}(\cdot) f_{i}=\sum \Lambda_{i}^{(u-2) / 2} f_{i}(\cdot) f_{i}=\sum R^{(u-2) / 2} f_{i}(\cdot) f_{i} \\
& \subseteq R^{(u-2) / 2}
\end{aligned}
$$

which is included in the operator of the form $D^{u-2}$ (or $-D^{u-2}$ ) and no boundary conditions, which is a closed operator by Lemma 1.4. Thus $\overline{T_{0}^{\nu}}$ exists, and its domain (which is $\Re\left(\overline{T_{0}^{-\nu}}\right)$ consists of $(u-2)$-times differentiable functions.
(iv) In dealing with operators of the form $D^{u}$ ( $u$ odd) I am handicapped by my ignorance of any equiconvergence theorem relating the expansions with different boundary conditions to one another. However, in the special case where the boundary conditions are

$$
f^{(i)}(0)=f^{(i)}(1) \quad(0 \leqq i \leqq u-1)
$$

the expansion is the Fourier series and the eigenvalues are simple. We have

$$
\lambda_{n}=(2 \pi i)^{u} n^{u}, \quad \Delta \lambda_{n}=(2 \pi i)^{u} u n^{u-1}+O\left(n^{u-2}\right)
$$

where $n$ runs through all the integers. Inspection of the proofs of our theorems shows that they can be adjusted to allow contours $\Gamma_{-m, n}$ which cut the spectrum of $T$ twice instead of once, as is necessary in this case. Thus, as in (iii), we can allow a perturbation by

$$
S=B_{u-2} D^{u-2}+\cdots+B_{1} D+B_{0}
$$

In this case we must take the space $C$ to be the space of continuous functions with period 1.
(v) In the cases where the boundary conditions allow multiple eigenvalues of $T$-for example, $D^{u}$ with $u$ even and the boundary conditions (5.2)—our arguments hold as far as line 27 on page 92 , at which point the fact that $\Delta \operatorname{dim} E_{r}$ is the (algebraic or generalised) multiplicity $m_{r}$ of $\lambda_{r}$ causes trouble. We can conclude that the annulus between $\Gamma_{n-1}$ and $\Gamma_{n}$ (or, at a later stage, the contour $\Omega_{n}$ ) contains a finite number of eigenvalues of $T+S$ whose generalised multiplicities add up to $m_{n}$, and obtain an expansion theorem, but we will know only the convergence of a suitable subsequence of the partial sums, those corresponding to the contours $\Gamma_{n}$. The cases where, say, $\lambda_{n} \sim K n^{u}$ but $\lim \inf \Delta \lambda_{n} / n^{u-1}=0$ are similar: we cannot put useful contours through too narrow gaps in $\sigma(T)$.

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    ${ }^{1}$ See the examples on p. 212 of Zygmund [8]. (If $p=2$, the convergence is unconditional.)

[^1]:    ${ }^{2} k$ should be regarded as the number of eigenvalues of $T$ which disappear, by amalgamating with other eigenvalues, when $S$ is added to $T$.

