# RENEWAL SERIES AND SQUARE-ROOT BOUNDARIES FOR BESSEL PROCESSES 

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## Abstract

We show how a description of Brownian exponential functionals as a renewal series gives access to the law of the hitting time of a square-root boundary by a Bessel process. This extends classical results by Breiman and Shepp, concerning Brownian motion, and recovers by different means, extensions for Bessel processes, obtained independently by Delong and Yor.

Let $B_{t}$ be the standard real valued Brownian motion and for $v>0$, introduce the geometric Brownian motion $\mathscr{E}_{t}^{(-v)}$ and its exponential functional $\mathscr{A}_{t}^{(-v)}$

$$
\begin{aligned}
\mathscr{E}_{t}^{(-v)} & :=\exp \left(B_{t}-v t\right) \\
\mathscr{A}_{t}^{(-v)} & :=\int_{0}^{t}\left(\mathscr{E}_{s}^{(-v)}\right)^{2} d s .
\end{aligned}
$$

Lamperti's representation theorem [5] applied to $\mathscr{E}_{t}^{(-v)}$ states

$$
\begin{equation*}
\mathscr{E}_{t}^{(-v)}=R_{\mathscr{A}_{t}^{(-v)}}^{(-v)} \tag{0.1}
\end{equation*}
$$

where $\left(R_{u}^{(-v)}, u \leq T_{0}\left(R^{(-v)}\right)\right.$ ) denotes the Bessel process of index $(-v)$ (equivalently of dimension $\delta=2(1-v))$, starting at 1 , which is an $\mathbb{R}_{+}$-valued diffusion with infinitesimal generator $\mathscr{L}^{(-v)}$
given by

$$
\mathscr{L}^{(-v)} f(x)=\frac{1}{2} f^{\prime \prime}(x)+\frac{1-2 v}{2 x} f^{\prime}(x), \quad f \in C_{b}^{2}\left(\mathbb{R}_{+}^{\star}\right)
$$

Let us remark that, in the special case $v=1 / 2$, equation (0.1) is nothing else but the DubinsSchwarz representation of the exponential martingale $\mathscr{E}_{t}^{(-1 / 2)}$ as Brownian motion time changed with $\mathscr{A}_{t}^{(-1 / 2)}$.
For a short summary of relations between Bessel processes and exponentials of Brownian motion, see e.g. Yor [10].
Let us consider now the following random variable $Z$, which is often called a perpetuity in the mathematical finance literature:

$$
Z:=\mathscr{A}_{\infty}^{(-v)}=\int_{0}^{\infty}\left(\mathscr{E}_{s}^{(-v)}\right)^{2} d s
$$

We deduce directly from (0.1) that

$$
\mathscr{A}_{\infty}^{(-v)}=T_{0}\left(R^{(-v)}\right)
$$

where $T_{0}:=\inf \left\{u: X_{u}=0\right\}$, and it is well-known (see [4], [11]), that

$$
\begin{equation*}
\mathscr{A}_{\infty}^{(-v)} \stackrel{(\text { law })}{=} \frac{1}{2 \gamma_{v}} \tag{0.2}
\end{equation*}
$$

where $\gamma_{v}$ is a gamma variable with parameter $v$ (i.e. with density $\frac{1}{\Gamma(v)} x^{v-1} e^{-x} \mathbf{1}_{\mathbb{R}_{+}}$).
Our main result characterizes the law of the hitting time of a parabolic boundary by $R_{u}^{(-v)}$ which corresponds to a Bessel process of dimension $d<2$.
Theorem 1. Let $0<b<c$, and $\sigma:=\inf \left\{u:\left(R_{u}^{(-v)}\right)^{2}=\frac{1}{c}(b+u)\right\}$ with $R_{0}^{(-v)}=1$.

$$
\begin{equation*}
E\left[(b+\sigma)^{-s}\right]=c^{-s} \frac{E\left[\left(1+2 b \gamma_{v+s}\right)^{-s}\right]}{E\left[\left(1+2 c \gamma_{v+s}\right)^{-s}\right]}, \quad \text { for any } s \geq 0 \tag{0.3}
\end{equation*}
$$

Proof: using the strong Markov property and the stationarity of the increments of Brownian motion, we obtain that for any stopping time $\tau$ of the Brownian motion

$$
\mathscr{A}_{\infty}^{(-v)}=: Z=\mathscr{A}_{\tau}^{(-v)}+\left(\mathscr{E}_{\tau}^{(-v)}\right)^{2} Z^{\prime}
$$

where $Z^{\prime}$ is independent of $\left(\mathscr{A}_{\tau}^{(-v)}, \mathscr{E}_{\tau}^{(-v)}\right)$ and $Z \stackrel{(\text { law })}{=} Z^{\prime}$.
This implies, by (0.1), that $Z$ satisfies the following affine equation (see [8] for many results about these equations)

$$
\begin{equation*}
\mathscr{A}_{\infty}^{(-v)}=: Z=\mathscr{A}_{\tau}^{(-v)}+\left(R_{\mathscr{A} \tau}^{(-v)}(-v)\right)^{2} Z^{\prime} \tag{0.4}
\end{equation*}
$$

where $Z^{\prime}$ is independent of $\left(\mathscr{A}_{\tau}^{(-v)}, R_{\mathscr{A}_{\tau}^{(-v)}}^{(-v)}\right)$ and $Z \stackrel{(l a w)}{=} Z^{\prime}$.
Obviously, $\sigma<T_{0}\left(R^{(-v)}\right)$. Taking now :

$$
\tau=\inf \left\{t:\left(R_{\mathscr{A}_{t}^{(-v)}}^{(-v)}\right)^{2}=\frac{1}{c}\left(b+\mathscr{A}_{t}^{(-v)}\right)\right\}
$$

we get $\mathscr{A}_{\tau}^{(-v)}=\sigma$, and the identity in law

$$
\begin{equation*}
b+Z \stackrel{(\text { law })}{=}(b+\sigma)\left(1+\frac{Z}{c}\right) \tag{0.5}
\end{equation*}
$$

where the variables $\sigma$ and $Z$ on the right-hand side are independent.
As a result, we obtain the Mellin transform of $b+\sigma$ which is:

$$
E\left[(b+\sigma)^{-s}\right]=c^{-s} \frac{E\left[(b+Z)^{-s}\right]}{E\left[(c+Z)^{-s}\right]}
$$

But, from (0.2)

$$
E\left[(b+\sigma)^{-s}\right]=c^{-s} \frac{E\left[\left(2 \gamma_{v}\right)^{s} \frac{1}{\left(1+2 b \gamma_{v}\right)^{s}}\right]}{E\left[\left(2 \gamma_{v}\right)^{s} \frac{1}{\left(1+2 c \gamma_{v}\right)^{s}}\right]}
$$

which gives the result.
One can now use the duality between the laws of Bessel processes of dimension $d$ and $4-d$ to get the analogous result of Theorem 1, and recover the result of Delong [2], [3], and Yor [9] which deals with the case $d \geq 2$.

Theorem 2. Let $0<b<c$, and $\sigma:=\inf \left\{u:\left(R_{u}^{(v)}\right)^{2}=\frac{1}{c}(b+u)\right\}$ with $R_{0}^{(v)}=1$.

$$
\begin{equation*}
E\left[(b+\sigma)^{-s}\right]=c^{-s} \frac{E\left[\left(1+2 b \gamma_{s}\right)^{-s+v}\right]}{E\left[\left(1+2 c \gamma_{s}\right)^{-s+v}\right]}, \quad \text { for any } s \geq 0 \tag{0.6}
\end{equation*}
$$

Proof : it is based on the following classical relation between the laws of the Bessel processes with indices $v$ and $-v$ :

$$
\begin{equation*}
\mathscr{P}_{x}^{(v)}{ }_{\mid \mathscr{F}_{t}}=\frac{\left(X_{t \wedge T_{0}}\right)^{2 v}}{x^{2 v}} \cdot \mathscr{P}_{x}^{(-v)}{ }_{\mid \mathscr{F}_{t}} \tag{0.7}
\end{equation*}
$$

which implies that

$$
E_{1}^{(v)}\left[(b+\sigma)^{-s}\right]=E_{1}^{(-v)}\left[X_{\sigma}^{2 v}(b+\sigma)^{-s}\right]=\frac{1}{c^{v}} E_{1}^{(-v)}\left[(b+\sigma)^{-s+v}\right]
$$

Theorem 1 gives the result.
Finally, it is easily shown, thanks to the classical representations of the Whittaker functions (see Lebedev [6] page 279), that the right-hand sides of (0.3) and (0.6) are expressed in terms of ratios of Whittaker functions. Let us recall their integral representation:

$$
W_{k, m}(z)=\frac{e^{-z / 2} z^{k}}{\Gamma\left(\frac{1}{2}-k+m\right)} \int_{0}^{\infty} t^{-k-\frac{1}{2}+m}\left(1+\frac{t}{z}\right)^{k-\frac{1}{2}+m} e^{-t} d t
$$

whenever $\Re\left(m-k+\frac{1}{2}\right) \geq 0$ and $\arg (z)<\pi$. Using this identitity, the rhs of (0.3) and (0.6) take respectively the form

$$
c^{-s} \frac{e^{\frac{1}{4 b}} W_{\frac{1-v}{2}-s, \frac{v}{2}}\left(\frac{1}{2 b}\right)}{e^{\frac{1}{4 c}} W_{\frac{1-v}{2}-s, \frac{v}{2}}\left(\frac{1}{2 c}\right)} \quad \text { and } \quad c^{-s} \frac{e^{\frac{1}{4 b}} W_{\frac{1+v}{2}-s, \frac{v}{2}}\left(\frac{1}{2 b}\right)}{e^{\frac{1}{4 c}} W_{\frac{1+v}{2}-s, \frac{v}{2}}\left(\frac{1}{2 c}\right)} .
$$

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