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RENEWAL SERIES AND SQUARE-ROOT BOUNDARIES FOR BESSEL PROCESSES

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Abstract

We show how a description of Brownian exponential functionals as a renewal series gives access to the law of the hitting time of a square-root boundary by a Bessel process. This extends classical results by Breiman and Shepp, concerning Brownian motion, and recovers by different means, extensions for Bessel processes, obtained independently by Delong and Yor.

Let B_t be the standard real valued Brownian motion and for v > 0, introduce the geometric Brownian motion $\mathcal{E}_t^{(-v)}$ and its exponential functional $\mathcal{A}_t^{(-v)}$

$$\mathscr{E}_t^{(-v)} := exp(B_t - vt)$$

$$\mathscr{A}_t^{(-\nu)} := \int_0^t (\mathscr{E}_s^{(-\nu)})^2 ds.$$

Lamperti's representation theorem [5] applied to $\mathcal{E}_t^{(-\nu)}$ states

$$\mathcal{E}_{t}^{(-\nu)} = R_{op^{(-\nu)}}^{(-\nu)} \tag{0.1}$$

where $(R_u^{(-\nu)}, u \le T_0(R^{(-\nu)}))$ denotes the Bessel process of index $(-\nu)$ (equivalently of dimension $\delta = 2(1-\nu)$), starting at 1, which is an \mathbb{R}_+ -valued diffusion with infinitesimal generator $\mathcal{L}^{(-\nu)}$

given by

$$\mathscr{L}^{(-v)}f(x) = \frac{1}{2}f''(x) + \frac{1-2v}{2x}f'(x), \quad f \in C_b^2(\mathbb{R}_+^*).$$

Let us remark that, in the special case v=1/2, equation (0.1) is nothing else but the Dubins-Schwarz representation of the exponential martingale $\mathcal{E}_t^{(-1/2)}$ as Brownian motion time changed with $\mathcal{A}_t^{(-1/2)}$.

For a short summary of relations between Bessel processes and exponentials of Brownian motion, see e.g. Yor [10].

Let us consider now the following random variable Z, which is often called a perpetuity in the mathematical finance literature:

$$Z := \mathscr{A}_{\infty}^{(-\nu)} = \int_{0}^{\infty} (\mathscr{E}_{s}^{(-\nu)})^{2} ds$$

We deduce directly from (0.1) that

$$\mathscr{A}_{\infty}^{(-\nu)} = T_0(R^{(-\nu)})$$

where $T_0 := \inf\{u : X_u = 0\}$, and it is well-known (see [4], [11]), that

$$\mathscr{A}_{\infty}^{(-v)} \stackrel{(law)}{=} \frac{1}{2\gamma_{v}} \tag{0.2}$$

where γ_v is a gamma variable with parameter v (i.e. with density $\frac{1}{\Gamma(v)}x^{v-1}e^{-x}\mathbf{1}_{\mathbb{R}_+}$).

Our main result characterizes the law of the hitting time of a parabolic boundary by $R_u^{(-\nu)}$ which corresponds to a Bessel process of dimension d < 2.

Theorem 1. Let 0 < b < c, and $\sigma := \inf\{u : (R_u^{(-v)})^2 = \frac{1}{c}(b+u)\}$ with $R_0^{(-v)} = 1$.

$$E[(b+\sigma)^{-s}] = c^{-s} \frac{E[(1+2b\gamma_{v+s})^{-s}]}{E[(1+2c\gamma_{v+s})^{-s}]}, \quad \text{for any } s \ge 0$$
 (0.3)

Proof: using the strong Markov property and the stationarity of the increments of Brownian motion, we obtain that for any stopping time τ of the Brownian motion

$$\mathscr{A}_{\infty}^{(-\nu)} =: Z = \mathscr{A}_{\tau}^{(-\nu)} + (\mathscr{E}_{\tau}^{(-\nu)})^2 Z'$$

where Z' is independent of $(\mathscr{A}_{\tau}^{(-v)},\mathscr{E}_{\tau}^{(-v)})$ and $Z\stackrel{(law)}{=} Z'$.

This implies, by (0.1), that Z satisfies the following affine equation (see [8] for many results about these equations)

$$\mathscr{A}_{\infty}^{(-\nu)} =: Z = \mathscr{A}_{\tau}^{(-\nu)} + (R_{\mathscr{A}^{(-\nu)}}^{(-\nu)})^2 Z' \tag{0.4}$$

where Z' is independent of $(\mathscr{A}_{\tau}^{(-v)}, R_{\mathscr{A}_{\tau}^{(-v)}}^{(-v)})$ and $Z \stackrel{(law)}{=} Z'$.

Obviously, $\sigma < T_0(R^{(-\nu)})$. Taking now :

$$\tau = \inf\{t : (R_{\mathscr{A}_{t}^{(-\nu)}}^{(-\nu)})^{2} = \frac{1}{c}(b + \mathscr{A}_{t}^{(-\nu)})\}$$

we get $\mathcal{A}_{\tau}^{(-v)} = \sigma$, and the identity in law

$$b + Z \stackrel{(law)}{=} (b + \sigma)(1 + \frac{Z}{c}) \tag{0.5}$$

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where the variables σ and Z on the right-hand side are independent. As a result, we obtain the Mellin transform of $b + \sigma$ which is:

$$E[(b+\sigma)^{-s}] = c^{-s} \frac{E[(b+Z)^{-s}]}{E[(c+Z)^{-s}]}$$

But, from (0.2)

$$E[(b+\sigma)^{-s}] = c^{-s} \frac{E[(2\gamma_{v})^{s} \frac{1}{(1+2b\gamma_{v})^{s}}]}{E[(2\gamma_{v})^{s} \frac{1}{(1+2c\gamma_{v})^{s}}]}$$

which gives the result.

One can now use the duality between the laws of Bessel processes of dimension d and 4-d to get the analogous result of Theorem 1, and recover the result of Delong [2], [3], and Yor [9] which deals with the case $d \ge 2$.

Theorem 2. Let 0 < b < c, and $\sigma := \inf\{u : (R_u^{(v)})^2 = \frac{1}{c}(b+u)\}$ with $R_0^{(v)} = 1$.

$$E[(b+\sigma)^{-s}] = c^{-s} \frac{E[(1+2b\gamma_s)^{-s+\nu}]}{E[(1+2c\gamma_s)^{-s+\nu}]}, \quad \text{for any } s \ge 0.$$
 (0.6)

Proof: it is based on the following classical relation between the laws of the Bessel processes with indices v and -v:

$$\mathscr{P}_{x}^{(\nu)}|_{\mathscr{F}_{t}} = \frac{(X_{t \wedge T_{0}})^{2\nu}}{r^{2\nu}} \mathscr{P}_{x}^{(-\nu)}|_{\mathscr{F}_{t}}$$
(0.7)

which implies that

$$E_1^{(\nu)}[(b+\sigma)^{-s}] = E_1^{(-\nu)}[X_{\sigma}^{2\nu}(b+\sigma)^{-s}] = \frac{1}{c^{\nu}}E_1^{(-\nu)}[(b+\sigma)^{-s+\nu}]$$

Theorem 1 gives the result.

Finally, it is easily shown, thanks to the classical representations of the Whittaker functions (see Lebedev [6] page 279), that the right-hand sides of (0.3) and (0.6) are expressed in terms of ratios of Whittaker functions. Let us recall their integral representation:

$$W_{k,m}(z) = \frac{e^{-z/2}z^k}{\Gamma(\frac{1}{2} - k + m)} \int_0^\infty t^{-k - \frac{1}{2} + m} (1 + \frac{t}{z})^{k - \frac{1}{2} + m} e^{-t} dt$$

whenever $\Re(m-k+\frac{1}{2}) \ge 0$ and $\arg(z) < \pi$.

Using this identitity, the rhs of (0.3) and (0.6) take respectively the form

$$c^{-s} \frac{e^{\frac{1}{4b}} W_{\frac{1-\nu}{2} - s, \frac{\nu}{2}}(\frac{1}{2b})}{e^{\frac{1}{4c}} W_{\frac{1-\nu}{2} - s, \frac{\nu}{2}}(\frac{1}{2c})} \quad \text{and} \quad c^{-s} \frac{e^{\frac{1}{4b}} W_{\frac{1+\nu}{2} - s, \frac{\nu}{2}}(\frac{1}{2b})}{e^{\frac{1}{4c}} W_{\frac{1+\nu}{2} - s, \frac{\nu}{2}}(\frac{1}{2c})}.$$

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