# Some New Stablity, Boundedness, and Square Integrability Conditions for Third-Order Neutral Delay Differential Equations 

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#### Abstract

In this paper, the authors establish some new sufficient conditions under which all solutions of a third order nonlinear neutral delay differential equation are stable, bounded, and square integrable. An example is also given to illustrate the results.


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## 1 Introduction

In this paper we consider third order neutral delay differential equations of the form

$$
\begin{equation*}
\left(q(t)\left(x^{\prime \prime}(t)+\beta(t) x^{\prime \prime}(t-r)\right)\right)^{\prime}+g\left(x(t), x^{\prime}(t)\right) x^{\prime \prime}(t)+f\left(x(t-r), x^{\prime}(t-r)\right)+h(x(t-r))=e(t) \tag{1.1}
\end{equation*}
$$

[^0]for all $t \geq t_{1} \geq t_{0}+r$, where $r>0$ and the functions $f\left(x(t-r), x^{\prime}(t-r)\right), g\left(x(t), x^{\prime}(t)\right), h(x(t-$ $r)$ ), and $q(t)$ are continuous on their respective domains, $q(t)>0$, and $h(0)=0$. In addition, it is also assumed that the derivatives $\left.\beta^{\prime}(t), g_{x}(x, y)=\frac{\partial}{\partial x} g(x, y)\right), f_{x}(x, y)=\frac{\partial}{\partial x} f(x, y), f_{y}(x, y)=$ $\frac{\partial}{\partial y} f(x, y)$, and $h^{\prime}(x)$ exist and are continuous.

Third order nonlinear equations in the form of (1.1) but with $q(t) \equiv 1, \beta(t) \equiv 0$, and $r=0$, have received a lot of attention over the years, and much of what was known prior to the 1970s can be found in the monograph by Reissig, Sansone, and Conti [15]. For example, the asymptotic stability of (1.1) with $e \equiv 0, \beta \equiv 0$, and $r=0$ has been previously discussed by several authors (see, for example, [13, 14, 12, 22]). Results on nonlinear third order equations with a delay can be found, for example, in $[1,8,16,17,18,19,20]$ where boundedness and stability properties of solutions are examined. Third order neutral equations are discussed in [3] for the case $g \equiv 0 \equiv f$; also see [25]. The oscillation of solutions of third order nonlinear equations has also proved to be a popular research topic as can be seen, for example, from the results in $[4,5,6,7,11,21,23,25]$. Physical applications of delay and neutral delay differential equations can be found in $[2,9,10]$. While it is clear that there are many possible special cases of equation (1.1), the types of results in this paper have not been previously obtained for an equation with such generality.

For convenience, we let

$$
Z(t)=x^{\prime \prime}(t)+\beta(t) x^{\prime \prime}(t-r) .
$$

By a solution of (1.1) we mean a nontrivial function $x \in C^{2}\left(\left[t_{x}, \infty\right), \mathbb{R}\right)$ such that $q(t) Z(t) \in$ $C^{1}\left(\left[t_{x}, \infty\right), \mathbb{R}\right)$ and which satisfies equation $(1.1)$ on $\left[t_{x}, \infty\right)$. Without further mention, we will assume throughout that every solution $x(t)$ of (1.1) under consideration here is continuable to the right and nontrivial, i.e., $x(t)$ is defined on some ray $\left[t_{x}, \infty\right)$. Moreover, we tacitly assume that (1.1) possesses such solutions.

This paper is organized as follows. In Section 2, we give stability results for (1.1) with $e \equiv 0$. In Section 3 the boundedness of solutions is discussed. Finally, in Section 4, sufficient conditions for the square integrability of solutions and their first and second derivatives are given.

## 2 Stability

We will assume that there are positive constants $a, b_{0}, b_{1}, q_{0}, q_{1}, q_{2}, \mu, \rho, \alpha, \beta_{1}, \delta_{0}, \delta_{1}, m$, $L$, and $K$ such that the following conditions on the functions in equation (1.1) are satisfied:
(i) $a+\mu \leq g(x, y) \leq a+\rho$, and $y g_{x}(x, y) \leq 0$;
(ii) $b_{1} \geq \frac{f(x, y)}{y} \geq b_{0},-K \leq f_{x}(x, y) \leq 0$, and $\left|f_{y}(x, y)\right| \leq L$;
(iii) $h^{\prime}(x) \leq \delta_{0}$ for all $x$ and $\frac{h(x)}{x} \geq \delta_{1}$ for $x \neq 0$;
(iv) $\left|\beta^{\prime}(t)\right| \leq \alpha, 0 \leq \beta(t) \leq \beta_{1}, q_{0} \leq q(t) \leq q_{1}$, and $\left|q^{\prime}(t)\right| \leq q_{2}$ for all $t \geq t_{0}$;
(v) $\int_{0}^{\infty}\left(\left|q^{\prime}(s)\right|+\left|\beta^{\prime}(s)\right|\right) d s<m<\infty$.

For ease of exposition, throughout this paper we will adopt the notation that

$$
A=-a b_{0}+q_{1} \delta_{0}+2+\frac{1}{2} q_{1} \beta_{1}\left(\delta_{0}+b_{1}\right), B=-\mu q_{0}+\beta_{1}\left(\rho q_{1}+1\right)+1, \text { and } M=\delta_{0}+K+L
$$

We begin with a stability result for equation (1.1) with $e \equiv 0$.
Theorem 2.1. In addition to conditions (i)-(v), assume that there exist positive constants $c_{1}, c_{2}$, and $\varepsilon$ such that
(vi) $A+\frac{1}{2}\left[\beta_{1}\left(q_{1} \delta_{0}+2+2 \beta_{1}+2 \alpha\right)+\delta_{0}\left(q_{1} \alpha+\beta_{1} q_{2}\right)+2 \alpha+4 \alpha \beta_{1}\right]=-c_{1}$,
(vii) $B+\frac{\beta_{1}}{2}\left[2+2 \beta_{1}+2 q_{1} \rho+q_{1} b_{1}\right]+\varepsilon=-c_{2}$,
(viii) $a-\frac{q_{1}^{2} \delta_{0}}{2}>0$.

Then zero solution of (1.1) is uniformly asymptotically stable provided

$$
r<\min \left\{\frac{2 c_{1}}{M a+2 \mu_{1}}, \frac{2 c_{2}}{M q_{1}+2 \mu_{2}}, \frac{2 \varepsilon}{M \beta_{1} q_{1}}\right\}
$$

where

$$
\begin{equation*}
\mu_{1}=\frac{1}{2}\left(q_{1}+a+\beta_{1} q_{1}\right)\left(\delta_{0}+K\right) \quad \text { and } \quad \mu_{2}=\frac{L}{2}\left(q_{1}+a+\beta_{1} q_{1}\right) \tag{2.1}
\end{equation*}
$$

Proof. We will write equation (1.1) as the equivalent differential system

$$
\left\{\begin{array}{l}
x^{\prime}(t)=y(t)  \tag{2.2}\\
y^{\prime}(t)=z(t) \\
{[q(t)(z(t)+\beta(t) z(t-r))]^{\prime}=-g(x(t), y(t)) z(t)} \\
\\
\quad-f(x(t), y(t))-h(x(t))+\Delta_{1}+\Delta_{2}+\Delta_{3}
\end{array}\right.
$$

where

$$
\begin{align*}
& \Delta_{1}(t)=\int_{t-r}^{t} f_{x}(x(s), y(s)) y(s) d s  \tag{2.3}\\
& \Delta_{2}(t)=\int_{t-r}^{t} f_{y}(x(s), y(s)) z(s) d s  \tag{2.4}\\
& \Delta_{3}(t)=\int_{t-r}^{t} h^{\prime}(x(s)) y(s) d s \tag{2.5}
\end{align*}
$$

By virtue of definition (2.2), we have

$$
\begin{align*}
x^{\prime}(t)+\beta(t) x^{\prime}(t-r) & =y(t)+\beta(t) y(t-r)=Y(t), \\
x^{\prime \prime}(t)+\beta(t) x^{\prime \prime}(t-r) & =z(t)+\beta(t) z(t-r)=Z(t), \\
Y^{\prime}(t) & =Z(t)+\beta^{\prime}(t) y(t-r) \tag{2.6}
\end{align*}
$$

The proof of this theorem depends on properties of the continuously differentiable function $W\left(t, x_{t}, y_{t}, z_{t}\right)=W$ defined by

$$
\begin{equation*}
W=V e^{-\frac{1}{\theta} \int_{0}^{t} \omega(s) d s} \tag{2.7}
\end{equation*}
$$

where

$$
\begin{equation*}
\omega(t)=\left|\beta^{\prime}(t)\right|+\left|q^{\prime}(t)\right| \tag{2.8}
\end{equation*}
$$

and

$$
\begin{align*}
V=V\left(t, x_{t}, y_{t}, z_{t}\right)= & a \int_{0}^{x} h(u) d u+q(t) h(x) Y+Y^{2}+\frac{1}{2}(a y+q(t) Z)^{2} \\
& +a \int_{0}^{y}[g(x, \xi)-a] \xi d \xi+q(t) \int_{0}^{y} f(x, \xi) d \xi+\lambda_{1} \int_{t-r}^{t} y^{2}(s) d s \\
& +\lambda_{2} \int_{t-r}^{t} z^{2}(s) d s+\hat{\mu_{1}} \int_{-r}^{0} \int_{t+s}^{t} y^{2}(\xi) d \xi d s+\hat{\mu_{2}} \int_{-r}^{0} \int_{t+s}^{t} z^{2}(\xi) d \xi d s . \tag{2.9}
\end{align*}
$$

Here, $\lambda_{1}, \lambda_{2}, \hat{\mu_{1}}, \hat{\mu_{2}}$, and $\theta$ are positive constants that will be determined later in the proof.
Noting that $h(0)=0$, it is easy to check that

$$
2 \int_{0}^{x} h^{\prime}(u) h(u) d u=h^{2}(x)
$$

and since $\frac{1}{2}(q(t) Z+a y)^{2} \geq 0$, this and (iii) imply that

$$
\begin{aligned}
a \int_{0}^{x} h(u) d u+q(t) h(x) Y & +Y^{2}+\frac{1}{2}(q(t) Z+a y)^{2} \\
& \geq a \int_{0}^{x} h(u) d u+\left(Y+\frac{q(t)}{2} h(x)\right)^{2}-\frac{q^{2}(t)}{4} h^{2}(x) \\
& \geq a \int_{0}^{x} h(u) d u-\frac{q^{2}(t)}{2} \int_{0}^{x} h^{\prime}(u) h(u) d u \\
& \geq \int_{0}^{x}\left(a-\frac{q_{1}^{2} \delta_{0}}{2}\right) h(u) d u \\
& \geq\left(a-\frac{q_{1}^{2} \delta_{0}}{2}\right) \frac{\delta_{1}}{2} x^{2}
\end{aligned}
$$

From (i)-(ii), it follows that

$$
a \int_{0}^{y}[g(x, \xi)-a] \xi d \xi \geq \frac{\mu a}{2} y^{2} \quad \text { and } \quad q(t) \int_{0}^{y} f(x, \xi) d \xi \geq \frac{q_{0} b_{0}}{2} y^{2}
$$

We see that $V=V\left(t, x_{t}, y_{t}, z_{t}\right) \geq 0$ and $V=V\left(t, x_{t}, y_{t}, z_{t}\right)=0$ if and only if $x=y=z=0$, so the functional V is positive definite. Thus, there exists a sufficiently small positive constant $K_{0}$ such that

$$
\begin{equation*}
V \geq K_{0}\left(x^{2}(t)+y^{2}(t)+Z^{2}(t)\right) \tag{2.10}
\end{equation*}
$$

By a straightforward calculation, the derivative of $V$ along the trajectories of system (2.2) is

$$
\begin{aligned}
V_{(2.2)}^{\prime}= & U_{1}+U_{2}+U_{3}+U_{4}-\hat{\mu_{1}} \int_{t-r}^{t} y^{2}(\theta) d \theta-\hat{\mu_{2}} \int_{t-r}^{t} z^{2}(\theta) d \theta \\
& +\left[2 \beta(t) \beta^{\prime}(t)-\lambda_{1}\right] y^{2}(t-r)-\lambda_{2} z^{2}(t-r) \\
& +a y(t) \int_{0}^{y} g_{x}(x, \xi) \xi d \xi+y(t) q(t) \int_{0}^{y} f_{x}(x, \xi) d \xi
\end{aligned}
$$

where

$$
\begin{aligned}
U_{1}= & q(t) h^{\prime}(x) y^{2}(t)-a f(x(t), y(t)) y(t)+\lambda_{1} y^{2}(t)+\hat{\mu_{1}} r y^{2}(t)+q^{\prime}(t) \int_{0}^{y} f(x, \xi) d \xi \\
& +\left(q(t)\left(a-g(x(t), y(t))+\lambda_{2}+\hat{\mu_{2}} r\right) z^{2}(t),\right. \\
U_{2}= & 2 y(t) z(t)+\left[q(t) \beta(t) h^{\prime}(x)+2 \beta^{\prime}(t)\right] y(t) y(t-r) \\
& +2 \beta(t) y(t) z(t-r)+2 \beta(t) y(t-r) z(t)+2 \beta^{2}(t) y(t-r) z(t-r) \\
& +\beta(t) q(t)[a-g(x(t), y(t))] z(t) z(t-r),
\end{aligned}
$$

and

$$
\begin{aligned}
U_{3}= & -\beta(t) q(t) f(x(t), y(t)) z(t-r)+\left[q(t) \beta^{\prime}(t)+\beta(t) q^{\prime}(t)\right] h(x) y(t-r) \\
& +q^{\prime}(t) h(x) y(t) \\
U_{4}= & (a y(t)+q(t) z(t)+\beta(t) q(t) z(t-r))\left(\Delta_{1}(t)+\Delta_{2}(t)+\Delta_{3}(t)\right)
\end{aligned}
$$

From conditions (i)-(iv), we see that

$$
a y(t) \int_{0}^{y} g_{x}(x, \xi) \xi d \xi+q(t) y(t) \int_{0}^{y} f_{x}(x, \xi) d \xi \leq 0
$$

and

$$
U_{1} \leq\left(q_{1} \delta_{0}-a b_{0}+\lambda_{1}+\frac{b_{1}}{2}\left|q^{\prime}(t)\right|+\hat{\mu_{1}} r\right) y^{2}(t)+\left(-\mu q_{0}+\lambda_{2}+\hat{\mu_{2}} r\right) z^{2}(t)
$$

Using Schwartz's inequality together with (i)-(iii), we obtain

$$
\begin{aligned}
U_{2} \leq & \frac{1}{2}\left[2+2 \beta_{1}+q_{1} \beta_{1} \delta_{0}+2\left|\beta^{\prime}(t)\right|\right] y^{2}(t)+\frac{1}{2}\left[q_{1} \beta_{1} \delta_{0}+2\left|\beta^{\prime}(t)\right|+2 \beta_{1}+2 \beta_{1}^{2}\right] y^{2}(t-r) \\
& +\frac{1}{2}\left[2+2 \beta_{1}+\beta_{1} q_{1} \rho\right] z^{2}(t)+\frac{1}{2}\left[2 \beta_{1}+2 \beta_{1}^{2}+2 \beta_{1} q_{1} \rho\right] z^{2}(t-r)
\end{aligned}
$$

On the other hand, by (2.3)-(2.5),

$$
\begin{aligned}
q(t) z(t)\left(\Delta_{1}(t)+\Delta_{2}(t)+\Delta_{3}(t)\right) \leq & \frac{M q_{1} r}{2} z^{2}(t)+\frac{q_{1}}{2}\left(\delta_{0}+K\right) \int_{t-r}^{t} y^{2}(s) d s \\
& +\frac{L q_{1}}{2} \int_{t-r}^{t} z^{2}(s) d s, \\
a y(t)\left(\Delta_{1}(t)+\Delta_{2}(t)+\Delta_{3}(t)\right) \leq & \frac{a M r}{2} y^{2}(t)+\frac{a}{2}\left(\delta_{0}+K\right) \int_{t-r}^{t} y^{2}(s) d s \\
& +\frac{a L}{2} \int_{t-r}^{t} z^{2}(s) d s,
\end{aligned}
$$

and

$$
\begin{aligned}
\beta(t) q(t) z(t-r)\left(\Delta_{1}(t)+\Delta_{2}(t)+\Delta_{3}(t)\right) \leq & \frac{\beta_{1} q_{1} M r}{2} z^{2}(t-r)+\frac{\beta_{1} q_{1}}{2}\left(\delta_{0}+K\right) \int_{t-r}^{t} y^{2}(s) d s \\
& +\frac{\beta_{1} q_{1} L}{2} \int_{t-r}^{t} z^{2}(s) d s
\end{aligned}
$$

Hence, we have the estimate

$$
\begin{aligned}
U_{4} \leq & \frac{M q_{1} r}{2} z^{2}(t)+\frac{M a r}{2} y^{2}(t)+\frac{\beta_{1} q_{1} M r}{2} z^{2}(t-r) \\
& +\frac{\delta_{0}+K}{2}\left(q_{1}+a+\beta_{1} q_{1}\right) \int_{t-r}^{t} y^{2}(s) d s \\
& +\frac{L}{2}\left(q_{1}+a+\beta_{1} q_{1}\right) \int_{t-r}^{t} z^{2}(s) d s
\end{aligned}
$$

Finally, condition (ii) and the facts that $h^{\prime}(x) \leq \delta_{0}$ and $h(0)=0$ imply

$$
\begin{aligned}
U_{3} \leq & \frac{1}{2}\left[\delta_{0}\left(q_{1}\left|\beta^{\prime}(t)\right|+\beta_{1}\left|q^{\prime}(t)\right|\right)+\left|q^{\prime}(t)\right| \delta_{0}\right] x^{2}(t) \\
& +\frac{1}{2}\left[\beta(t) q(t) b_{1}+\left|q^{\prime}(t)\right| \delta_{0}\right] y^{2}(t) \\
& +\frac{1}{2}\left(q_{1}\left|\beta^{\prime}(t)\right|+\beta_{1}\left|q^{\prime}(t)\right|\right) \delta_{0} y^{2}(t-r) \\
& +\frac{1}{2} \beta_{1} q_{1} b_{1} z^{2}(t-r) .
\end{aligned}
$$

With some rearrangement of terms and using the above estimates, we obtain

$$
\begin{aligned}
V^{\prime} \leq & \left(\frac{1}{2}\left[\delta_{0} q_{1}\left|\beta^{\prime}(t)\right|+\left(\beta_{1}+1\right) \delta_{0}\left|q^{\prime}(t)\right|\right]\right) x^{2}(t)+\left(\left|\beta^{\prime}(t)\right|+\left(\frac{\delta_{0}}{2}+\frac{3 b_{1}}{2}\right)\left|q^{\prime}(t)\right|\right) y^{2}(t) \\
& +\left(-a b_{0}+q_{1} \delta_{0}+2+\frac{1}{2} q_{1} \beta_{1}\left(\delta_{0}+b_{1}\right)+\lambda_{1}+\left(\frac{M a}{2}+\hat{\mu_{1}}\right) r\right) y^{2}(t) \\
& +\left[\frac{1}{2}\left[\beta_{1}\left(q_{1} \delta_{0}+2+2 \beta_{1}+2 \alpha\right)+\delta_{0}\left(q_{1} \alpha+\beta_{1} q_{2}\right)+2 \alpha+4 \alpha \beta_{1}\right]-\lambda_{1}\right] y^{2}(t-r) \\
& +\left[-\mu q_{0}+\rho \beta_{1} q_{1}+1+\beta_{1}+\lambda_{2}+\left(\hat{\mu_{2}}+\frac{M q_{1}}{2}\right) r\right] z^{2}(t) \\
& +\left[\frac{\beta_{1}}{2}\left[2+2 \beta_{1}+2 q_{1} \rho+q_{1} b_{1}\right]-\lambda_{2}+\frac{\beta_{1} q_{1} M}{2} r\right] z^{2}(t-r) \\
& +\left(\frac{1}{2}\left(q_{1}+a+\beta_{1} q_{1}\right)\left(\delta_{0}+K\right)-\hat{\mu}_{1}\right) \int_{t-r}^{t} y^{2}(s) d s \\
& +\left(\frac{L}{2}\left(q_{1}+a+\beta_{1} q_{1}\right)-\hat{\mu_{2}}\right) \int_{t-r}^{t} z^{2}(s) d s .
\end{aligned}
$$

If we now choose

$$
\begin{aligned}
\hat{\mu_{1}}=\frac{1}{2}\left(q_{1}+a+\beta_{1} q_{1}\right)\left(\delta_{0}+K\right) & =\mu_{1} \\
\hat{\mu_{2}}=\frac{L}{2}\left(q_{1}+a+\beta_{1} q_{1}\right) & =\mu_{2} \\
\frac{1}{2}\left[\beta_{1}\left(q_{1} \delta_{0}+2+2 \beta_{1}+2 \alpha\right)+\delta_{0}\left(q_{1} \alpha+\beta_{1} q_{2}\right)+2 \alpha+4 \alpha \beta_{1}\right] & =\lambda_{1} \\
\frac{\beta_{1}}{2}\left[2+2 \beta_{1}+2 q_{1} \rho+q_{1} b_{1}\right]+\varepsilon & =\lambda_{2}
\end{aligned}
$$

then from conditions (vi)-(vii), we see that

$$
\begin{aligned}
V^{\prime} \leq & K_{1}\left(\left|\beta^{\prime}(t)\right|+\left|q^{\prime}(t)\right|\right)\left(x^{2}(t)+y^{2}(t)\right)+\left(-c_{1}+\left(\frac{M a}{2}+\mu_{1}\right) r\right) y^{2}(t) \\
& +\left[-c_{2}+\left(\mu_{2}+\frac{M q_{1}}{2}\right) r\right] z^{2}(t)+\left[-\varepsilon+\frac{\beta_{1} q_{1} M}{2} r\right] z^{2}(t-r)
\end{aligned}
$$

where

$$
K_{1}=\frac{1}{2} \max \left\{\delta_{0} q_{1}, \delta_{0}\left(\beta_{1}+1\right), 2, \delta_{0}+3 b_{1}\right\} .
$$

Taking

$$
r<\min \left\{\frac{2 c_{1}}{M a+2 \mu_{1}}, \frac{2 c_{2}}{M q_{1}+2 \mu_{2}}, \frac{2 \varepsilon}{M \beta_{1} q_{1}}\right\}
$$

we obtain

$$
\begin{equation*}
V^{\prime} \leq K_{1} \omega(t)\left(x^{2}(t)+y^{2}(t)\right)-K_{2}\left(y^{2}(t)+z^{2}(t)\right) \tag{2.11}
\end{equation*}
$$

where

$$
K_{2}=\min \left\{c_{1}-\left(\frac{M a}{2}+\mu_{1}\right) r, c_{2}-\left(\mu_{2}+\frac{M q_{1}}{2}\right) r\right\}
$$

Hence, by (2.7), (2.10), and (2.11)

$$
\begin{aligned}
W^{\prime}= & \left(V^{\prime}-\frac{1}{\theta} \omega(t) V\right) e^{-\frac{1}{\theta} \int_{0}^{t} \omega(s) d s} \\
\leq & \left\{K_{1} \omega(t)\left[x^{2}(t)+y^{2}(t)\right]-K_{2}\left[y^{2}(t)+z^{2}(t)\right]\right. \\
& \left.-\frac{K_{0}}{\theta} \omega(t)\left[x^{2}(t)+y^{2}(t)+Z^{2}(t)\right]\right\} e^{-\frac{1}{\theta} \int_{0}^{t} \omega() d s}
\end{aligned}
$$

By taking

$$
\frac{K_{0}}{K_{1}}=\theta
$$

we obtain

$$
W^{\prime} \leq-K_{2}\left(y^{2}(t)+z^{2}(t)\right) e^{-\frac{1}{\theta} \int_{0}^{t} \omega(s) d s}
$$

Since

$$
\int_{t_{0}}^{t} \omega(s) d s<m, \text { for all } t \geq t_{0}
$$

we have

$$
\begin{equation*}
W^{\prime} \leq-K_{3}\left(y^{2}(t)+z^{2}(t)\right) \tag{2.12}
\end{equation*}
$$

where $K_{3}=K_{2} e^{-\frac{m K_{1}}{K_{0}}}$.
From (2.12) we see that $W^{\prime} \leq 0$ and $W^{\prime}=0$ on the set $M=\{(x, 0,0)\}$. Now the largest invariant set contained in $M$ is the origin, so by LaSalle's invariance principle, the zero solution of (2.2) is uniformly asymptotically stable.

## 3 BOUNDEDNESS

Our main theorem in this section is for the forced equation (1.1). In this case our system becomes

$$
\left\{\begin{array}{l}
x^{\prime}(t)=y(t)  \tag{3.1}\\
y^{\prime}(t)=z(t) \\
{[q(t)(z(t)+\beta(t) z(t-r))]^{\prime}=-g(x(t), y(t)) z(t)} \\
\quad-f(x(t), y(t))-h(x(t))+e(t)+\Delta_{1}+\Delta_{2}+\Delta_{3}
\end{array}\right.
$$

Theorem 3.1. In addition to the conditions of Theorem 2.1, assume that

$$
\left(I_{1}\right) \quad \int_{0}^{\infty} e(s) d s<\infty
$$

Then there exists a positive constant $D$ such that any solution of (3.1) satisfies

$$
\begin{equation*}
|x(t)| \leq D, \quad|y(t)| \leq D, \quad \text { and } \quad|Z(t)| \leq D \tag{3.2}
\end{equation*}
$$

Proof. On differentiating (2.9) along the solutions of system (3.1) we obtain

$$
V_{(3.1)}^{\prime} \leq-K_{2}\left(y^{2}(t)+z^{2}(t)\right)+a e(t) y+e(t) q(t) Z, \quad \text { for all } t \geq t_{1}
$$

Applying conditions ( $I_{1}$ ) and (iv) gives

$$
V_{(3.1)}^{\prime} \leq-K_{2}\left(y^{2}(t)+z^{2}(t)\right)+e(t)\left(a|y(t)|+q_{1}|Z(t)|\right)
$$

Now, the inequality $|u| \leq u^{2}+1$ leads to

$$
\begin{equation*}
V_{(3.1)}^{\prime} \leq-K_{2}\left(y^{2}(t)+z^{2}(t)\right)+K_{4} e(t)\left(y^{2}+Z^{2}+2\right), \text { for all } t \geq t_{1} \tag{3.3}
\end{equation*}
$$

where $K_{4}=\max \left\{a, q_{1}\right\}$.
In view of (2.10), the above estimates imply that

$$
\begin{equation*}
V_{(3.1)}^{\prime} \leq-K_{2}\left(y^{2}(t)+z^{2}(t)\right)+\frac{K_{4}}{K_{0}} e(t) V+2 K_{4} e(t) \tag{3.4}
\end{equation*}
$$

Integrating both sides (3.4) from $t_{1}$ to $t$, we easily obtain

$$
V(t)-V\left(t_{1}\right) \leq 2 K_{4} \int_{t_{1}}^{t} e(s) d s+\frac{K_{4}}{K_{0}} \int_{t_{1}}^{t} V(s) e(s) d s
$$

Condition ( $\mathrm{I}_{1}$ ) and an application of Gronwall's inequality shows that $V(t)$ is bounded, and the conclusion of the theorem follows immediately from (2.10).

Remark 3.2. The uniform asymptotic stability from Theorem 2.1 together with the boundedness of all solutions from Theorem 3.1 ensure the global uniform asymptotic stability of the zero solution of the unforced equation.

## 4 SQUARE INTEGRABILITY

Our next result concerns the square integrability of solutions of equation (1.1).
Theorem 4.1. If the conditions of Theorem 3.1 hold, then

$$
\int_{t_{0}}^{\infty}\left(x^{\prime \prime 2}(s)+x^{\prime 2}(s)+x^{2}(s)\right) d s<\infty .
$$

Proof. Define $H(t)$ by

$$
\begin{equation*}
H(t)=V(t)+\eta \int_{t_{1}}^{t}\left(z^{2}(s)+y^{2}(s)\right) d s \quad \text { for all } t \geq t_{1} \geq t_{0}+r \tag{4.1}
\end{equation*}
$$

where $\eta>0$ is a constant to be specified later. Differentiating $H$ and using (3.4) we obtain

$$
H^{\prime}(t) \leq\left(\eta-K_{2}\right)\left(y^{2}(t)+z^{2}(t)\right)+\left(\frac{K_{4}}{K_{0}} V+2 K_{4}\right) e(t) .
$$

If we Choose $\eta<K_{2}$, then from the boundedness of $V$, we see that

$$
\begin{equation*}
H^{\prime}(t) \leq K_{5} e(t), \tag{4.2}
\end{equation*}
$$

for some $K_{5}>0$. Integrating (4.2) from $t_{1}$ to $t$, and using condition $\left(I_{1}\right)$ we have that $H(t)$ is bounded. This implies the existence of positive constants $\kappa_{1}$ and $\kappa_{2}$ such that

$$
\int_{t_{1}}^{\infty} y^{2}(s) d s \leq \kappa_{1} \text { and } \int_{t_{1}}^{\infty} z^{2}(s) d s \leq \kappa_{2},
$$

so

$$
\begin{equation*}
\int_{t_{0}}^{\infty} x^{\prime 2}(s) d s<\infty \text { and } \int_{t_{0}}^{\infty} x^{\prime \prime 2}(s) d s<\infty . \tag{4.3}
\end{equation*}
$$

Next, we show that $\int_{t_{1}}^{\infty} x^{2}(s) d s<\infty$. Multiplying (1.1) by $x(t-r)$ gives

$$
\begin{align*}
\left(q ( t ) \left(x^{\prime \prime}(t)+\beta(t) x^{\prime \prime}( \right.\right. & t-r)))^{\prime} x(t-r)+g\left(x(t), x^{\prime}(t)\right) x^{\prime \prime}(t) x(t-r) \\
& +f\left(x(t-r), x^{\prime}(t-r)\right) x(t-r)+h(x(t-r)) x(t-r)=e(t) x(t-r), \tag{4.4}
\end{align*}
$$

and then integrating from $t_{1}$ to $t$, we have

$$
\begin{equation*}
\int_{t_{1}}^{t} h(x(s-r)) x(s-r) d s=L_{1}(t)+L_{2}(t)+L_{3}(t)+L_{4}(t) \tag{4.5}
\end{equation*}
$$

where

$$
\begin{aligned}
L_{1}(t) & =-\int_{t_{1}}^{t}\left(q(s)\left(x^{\prime \prime}(s)+\beta(s) x^{\prime \prime}(s-r)\right)\right)^{\prime} x(s-r) d s, \\
L_{2}(t) & =-\int_{t_{1}}^{t} g\left(x(s), x^{\prime}(s)\right) x^{\prime \prime}(s) x(s-r) d s, \\
L_{3}(t) & =-\int_{t_{1}}^{t} f\left(x(s-r), x^{\prime}(s-r)\right) x(s-r) d s, \\
L_{4}(t) & =\int_{t_{1}}^{t} e(s) x(s-r) d s .
\end{aligned}
$$

Integrating by parts,

$$
\begin{aligned}
L_{1}(t)= & -\left[\left(q(t)\left(x^{\prime \prime}(t)+\beta(t) x^{\prime \prime}(t-r)\right)\right) x(t-r)-\left(q\left(t_{1}\right)\left(x^{\prime \prime}\left(t_{1}\right)+\beta\left(t_{1}\right) x^{\prime \prime}\left(t_{1}-r\right)\right)\right) x\left(t_{1}-r\right)\right] \\
& +\int_{t_{1}}^{t}\left(q(s)\left(x^{\prime \prime}(s)+\beta(s) x^{\prime \prime}(s-r)\right)\right) x^{\prime}(s-r) d s \\
\leq & 2 q_{1}\left(1+\beta_{1}\right) D^{2}+\int_{t_{1}}^{t}\left(q(s) x^{\prime \prime}(s) x^{\prime}(s-r)+q(s) \beta(s) x^{\prime \prime}(s-r) x^{\prime}(s-r)\right) d s \\
= & 2 q_{1}\left(1+\beta_{1}\right) D^{2}+\frac{q_{1}}{2} \int_{t_{1}}^{t}\left[x^{\prime \prime 2}(s)+\beta_{1} x^{\prime \prime 2}(s-r)+2 x^{\prime 2}(t-s)\right] d s \\
\leq & l_{1}<\infty
\end{aligned}
$$

by (3.2) and (4.3).
In the same way, applying (i) and (4.3),

$$
\begin{aligned}
L_{2}(t) & \leq \int_{t_{1}}^{t}\left|g\left(x(s), x^{\prime}(s)\right) x^{\prime \prime}(s) x(s-r)\right| d s \\
& \leq\left\{\int_{t_{1}}^{t}\left[g\left(x(s), x^{\prime}(s)\right) x^{\prime \prime}(s)\right]^{2} d s\right\}^{\frac{1}{2}}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}} \\
& \leq\left\{(a+\rho)^{2} \int_{t_{1}}^{t}\left(x^{\prime \prime}(s)\right)^{2} d s\right\}^{\frac{1}{2}}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}} \\
& \leq l_{2}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}}
\end{aligned}
$$

Condition (ii) and (4.3) imply

$$
\begin{aligned}
L_{3}(t) & \leq \int_{t_{1}}^{t}\left|f\left(x(s-r), x^{\prime}(s-r)\right) x(s-r)\right| d s \\
& \leq\left\{b_{1}^{2} \int_{t_{1}}^{t}\left(x^{\prime}(s-r)\right)^{2} d s\right\}^{\frac{1}{2}}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}} \\
& \leq l_{3}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}}
\end{aligned}
$$

Finally,

$$
L_{4}(t) \leq \int_{t_{1}}^{t}|e(s) x(s-r)| d s \leq D \int_{t_{1}}^{t}|e(s)| d s \leq l_{4}<\infty
$$

On the other hand, from condition (iii), we have

$$
\int_{t_{1}}^{t} x(s-r) h(x(s-r)) d s \geq \delta_{1} \int_{t_{1}}^{t} x^{2}(s-r) d s
$$

Therefore,

$$
\delta_{1} \int_{t_{1}}^{t} x^{2}(s-r) d s \leq l_{1}+l_{2}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}}+l_{3}\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}}+l_{4}
$$

If

$$
\int_{t_{1}}^{t} x^{2}(s-r) d s \rightarrow \infty \text { as } t \rightarrow \infty
$$

then dividing both sides of by $\left\{\int_{t_{1}}^{t} x^{2}(s-r) d s\right\}^{\frac{1}{2}}$ we immediately obtain a contradiction. This completes the proof of the theorem.

We conclude our paper with an example to illustrate our theorems.
Example 4.2. Consider the third order nonlinear nonautonomous delay differential equation

$$
\begin{align*}
& \left(\ln (3+\cos t)\left(x^{\prime \prime}(t)+\frac{1}{54.93} \ln (2+\sin t) x^{\prime \prime}(t-r)\right)\right)^{\prime} \\
& +\left(46.4-\frac{1}{\pi} \arctan \left(x(t) x^{\prime}(t)\right)\right) x^{\prime \prime}(t)+13.5 x^{\prime}(t-r)-\frac{x^{\prime}(t-r)}{\pi} \arctan \left(e^{x(t-r))} x^{\prime}(t-r)\right) \\
&  \tag{4.6}\\
& \quad+x(t-r)+\frac{x(t-r)}{1+|x(t-r)|}=\frac{1}{1+t^{2}} .
\end{align*}
$$

Taking $a=2$, we see that

$$
\mu=43.9 \leq g(x, y)-a=45.5-\frac{1}{\pi} \arctan (x y) \leq 44.9=\rho,
$$

and

$$
y g_{x}(x, y)=-\frac{1}{\pi} \frac{y^{2}}{1+x^{2} y^{2}} \leq 0
$$

Now

$$
f(x, y)=13.5 y-\frac{y}{\pi} \arctan \left(e^{x} y\right),
$$

so we see that $f(x, 0)=f(0,0)=0$ for all $x$ and

$$
b_{0}=13<\frac{f(x, y)}{y}=13.5-\frac{1}{\pi} \arctan \left(e^{x} y\right)<14=b_{1} .
$$

Moreover,

$$
K=-\frac{1}{\pi} \leq f_{x}(x, y)=-\frac{1}{\pi} \frac{y^{2}}{1+\left(e^{x} y\right)^{2}} \leq 0
$$

and

$$
\left|f_{y}(x, y)\right|=\left|13.5-\frac{1}{\pi} \arctan \left(e^{x} y\right)-\frac{1}{\pi} \frac{e^{x} y}{1+\left(e^{x} y\right)^{2}}\right| \leq 14.5=L .
$$

Since

$$
h(x)=x+\frac{x}{1+|x|},
$$

it is clear that $h(0)=0$ and

$$
\left|h^{\prime}(x)\right|=\left|1+\frac{1}{(1+|x|)^{2}}\right| \leq 2=\delta_{0} .
$$

Since $0 \leq \frac{1}{1+|x|} \leq 1$ for all $x$, we have

$$
\frac{h(x)}{x} \geq 1=\delta_{1} \quad \text { for all } \quad x \neq 0 .
$$

Also, we have

$$
\begin{aligned}
q_{0} & =\ln 2 \leq q(t)=\ln (3+\cos t) \leq 2 \ln 2=q_{1}, \\
\left|q^{\prime}(t)\right| & =\left|\frac{\sin t}{3+\cos t}\right| \leq \frac{1}{2}=q_{2}, \\
\int_{-\infty}^{+\infty}\left|q^{\prime}(s)\right| d s & =\int_{-\infty}^{+\infty}\left|\frac{\sin s}{3+\cos s}\right| d s \leq \int_{-\infty}^{+\infty} \frac{1}{3+\cos s} d s \\
& =\int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \frac{1}{2+u^{2}} d u=\frac{2}{\sqrt{2}} \tan ^{-1}\left(\frac{\pi}{2 \sqrt{2}}\right), \\
0 & \leq \beta(t)=\frac{1}{54.93} \ln (2+\sin t) \leq \frac{\ln 3}{54.93}=\beta_{1}, \\
\left|\beta^{\prime}(t)\right| & =\frac{1}{54.93}\left|\frac{\cos t}{2+\sin t}\right| \leq \frac{1}{54.93}=\alpha,
\end{aligned}
$$

and

$$
\int_{0}^{t}\left(\left|q^{\prime}(s)\right|+\left|\beta^{\prime}(s)\right|\right) d s<m .
$$

for all $t \geq t_{0}$.
Choosing $\varepsilon=10^{-1}$, we obtain

$$
\begin{aligned}
-a b_{0}+q_{1} \delta_{0}+2+\frac{1}{2} q_{1} \beta_{1}\left(\delta_{0}+b_{1}\right)+\frac{1}{2} & {\left[\beta_{1}\left(q_{1} \delta_{0}+2+2 \beta_{1}+2 \alpha\right)\right.} \\
& \left.+\delta_{0}\left(q_{1} \alpha+\beta_{1} q_{2}\right)+2 \alpha+4 \alpha \beta_{1}\right]=-20.903=-c_{1}
\end{aligned},
$$

and

$$
a-\frac{q_{1}^{2} \delta_{0}}{2}=7.8188 \times 10^{-2}>0 .
$$

Hence, if $r<\min \left\{\frac{2 c_{1}}{M a+2 \mu_{1}}, \frac{2 c_{2}}{M q_{1}+2 \mu_{2}}, \frac{2 \varepsilon}{M \beta_{1} q_{1}}\right\}=\min \{1.0061,0.73072,0.42891\}=0.42891$, then all the conditions of Theorems 2.1,3.1, and 4.1 are satisfied, so for any solution $x(t)$ of equation (4.6), $x(t), x^{\prime}(t)$, and $x^{\prime \prime}(t)$ are bounded and square integrable. In addition, the zero solution of the unforced equation is globally uniformly asymptotically stable.

## Conclusion

A non-linear neutral delay differential equation of the third order is considered. Using Lyapunov's direct method, we have derived new sufficient conditions for the global uniform asymptotic stability of the zero solution as well as the boundedness and square integrability of all solutions. The results here will be of interest to other researchers working on qualitative behavior of solutions of differential equations.

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