# One Dimensional $p$-Laplacian with a Concentrated Nonlinear Source 

Yuanyuan Ke Jingxue Yin* Shuhe Wang


#### Abstract

We establish the existence of local solutions for the Cauchy problem for one-dimensional $p$-Laplacian equation with a concentrated nonlinear source.


## 1 Introduction

Consider the Cauchy problem for one-dimensional $p$-Laplacian with a concentrated nonlinear source

$$
\begin{align*}
\frac{\partial u}{\partial t}-D\left(|D u|^{p-2} D u\right) & =\delta(x) f(u), & & (x, t) \in Q_{T}  \tag{1.1}\\
u(x, 0) & =u_{0}(x), & & x \in \mathbb{R}, \tag{1.2}
\end{align*}
$$

where $\delta(x)$ is the Dirac measure, $p>2, D=\frac{\partial}{\partial x}, Q_{T}=\mathbb{R} \times(0, T)$.
During the past years, there are rich references concerning the partial differential equations with measure data as the sources, see for example [1]-[5]. In particular, the problems with the source of the form $\delta(x) f(u)$ have been considered by some authors in recent years, see for instance [6]-[10]. If $p=2$, Olmstead and Roberts [6] studied the initial boundary value problem, which is motivated by the model that a combustible medium be ignited by using a heated wire or a pair of small electrodes to supply a large amount of energy to a very confined area. The authors discussed a possible blow-up solution of the problem by using Green's function and analyzing its corresponding nonlinear Volterra equation of

[^0]the second kind at the site of the concentrated source. For the porous medium equation with this kind of source, Yin et al [10] studied the existence of generalized solutions for the Cauchy problem based on some a priori estimates on solutions. When the nonlinear source term in Eq. (1.1) is replaced by Dirac measure $\delta(x)$, the existence of generalized solutions was obtained by Li et al [11] for $p>2$; while when the nonlinear source term is replaced by a bounded Borel measure, in [12] the authors considered the existence of solutions and gave summability results for the gradients of solutions for the case $p>1$. And also there are a large amount of papers devoted to this type of degenerate parabolic problems when the nonlinear source is of the form $f(u)$, such as [13], in which the authors studied the local existence of solutions by the regularized methods.

The purpose of this paper is to investigate the local existence of generalized solutions for the problem (1.1)-(1.2). Due to the degeneracy, we have to consider solutions in some weak sense, namely

Definition 1.1. A nonnegative function $u: Q_{T} \rightarrow[0,+\infty)$ is called a generalized solution of the problem (1.1)-(1.2) in $Q_{T}$, if $u \in C((\mathbb{R} \backslash\{0\}) \times(0, T)) \cap$ $L^{\infty}\left(0, T ; L_{\text {loc }}^{\infty}(\mathbb{R})\right) \cap B V_{x}\left(Q_{T}\right)$, and $u$ satisfies

$$
\begin{aligned}
\int_{\mathbb{R}} u_{0}(x) \varphi(x, 0) d x & +\iint_{Q_{T}}\left(u \varphi_{t}-|D u|^{p-2} D u D \varphi\right) d x d t \\
& +\int_{0}^{T} \frac{1}{2}\left(f\left(u^{l}(0, t)\right)+f\left(u^{r}(0, t)\right)\right) \varphi(0, t) d t=0
\end{aligned}
$$

for any $\varphi \in C^{\infty}\left(Q_{T}\right)$, which vanishes for large $|x|$ and $t=T$, where $B V_{x}\left(Q_{T}\right)$ is a subset of $L_{l o c}^{1}\left(Q_{T}\right)$, in which the derivatives in $x$ of each function are Radon measures in $Q_{T}$, for fixed $t \in(0, T), u^{l}(0, t)$ and $u^{r}(0, t)$ are the left limit and the right limit of $u$ at $x=0$.

The main result of the paper is the following theorem.
Theorem 1.1. Let $f(s)$ be nonnegative, bounded and continuously differentiable, and $u_{0}(x)$ be nonnegative and Hölder continuous with compact support. Assume that $u_{0}(x) \in W^{1, p}(\mathbb{R})$. Then there exists a finite $T^{*}>0$, such that the Cauchy problem (1.1)-(1.2) admits a generalized solution in $Q_{T^{*}}$.

Since $p>2$, the appearance of the diffusion term makes the theory of Green's function used in [6] inapplicable. In addition, by noticing that both $\delta(x)$ and $f(u)$ appear in the source term, some estimates we needed cannot be obtained by using the methods in [11]-[13]. In our paper, just as done in [10], we should first make an approximation of the Dirac measure. Owing to the degeneracy of Eq. (1.1), we should also use parabolic regularization to approach the equation. Then we establish the existence of local solutions by some suitable estimates on the approximate solutions.

## 2 Proof of the Main Result

In this section, we give the proof of our main result. In order to prove the existence of generalized solutions, we first consider the following approximate problem

$$
\begin{align*}
\frac{\partial u}{\partial t}-D\left(\left(|D u|^{2}+\frac{1}{n}\right)^{(p-2) / 2} D u\right) & =\delta_{\varepsilon}(x) f(u), & & (x, t) \in Q_{R, T}  \tag{2.1}\\
u(x, 0) & =u_{0}(x), & & x \in(-R, R)  \tag{2.2}\\
u( \pm R, t) & =0, & & t \in(0, T) \tag{2.3}
\end{align*}
$$

where $Q_{R, T}=(-R, R) \times(0, T), R$ is a properly large positive constant, and

$$
\begin{aligned}
& \delta_{\varepsilon}(x)=\frac{1}{\varepsilon} j\left(\frac{x}{\varepsilon}\right), \quad 0<\varepsilon<1, \\
& j(x)= \begin{cases}\frac{1}{A} e^{1 /\left(|x|^{2}-1\right)}, & |x|<1, \\
0, & |x| \geq 1,\end{cases} \\
& A=\int_{-1}^{1} e^{1 /\left(|x|^{2}-1\right)} d x .
\end{aligned}
$$

Obviously, we have

$$
\begin{aligned}
& \int_{\mathbb{R}} \delta_{\varepsilon}(x) d x=1, \quad 0 \leq \delta_{\varepsilon}(x) \in C_{0}^{\infty}(\mathbb{R}) \\
& \quad \operatorname{supp} \delta_{\varepsilon}(x)=\{x \in \mathbb{R},|x| \leq \varepsilon\}
\end{aligned}
$$

and

$$
\lim _{\varepsilon \rightarrow 0^{+}} \int_{\mathbb{R}} \delta_{\varepsilon}(x) \phi(x) d x=\phi(0), \quad \forall \phi \in C(\mathbb{R})
$$

By classical theories for parabolic equations, the problem (2.1)-(2.3) admits a unique nonnegative solution $u_{\varepsilon, R, n}$.

Next we make some estimates on $u_{\varepsilon, R, n}$.
Lemma 2.1. There exists $T^{\prime} \in[0, T]$, such that

$$
\left\|u_{\varepsilon, R, n}\right\|_{L^{\infty}\left(Q_{R, T^{\prime}}\right)} \leq C_{1},
$$

where $C_{1}$ is a positive constant depending on $\varepsilon$ and $T^{\prime}$.
Proof. Let $w(t)$ be the solution of the ordinary differential equation

$$
\begin{align*}
& \frac{d w}{d t}=M_{\varepsilon} f(w)  \tag{2.4}\\
& w(0)=\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})} \tag{2.5}
\end{align*}
$$

where $M_{\varepsilon}=\max _{x \in \mathbb{R}} \delta_{\varepsilon}(x)$. Then there exists $T_{0} \in(0, T)$ such that the problem (2.4)(2.5) has a solution $w(t)$ on $\left[0, T_{0}\right]$ and $T_{0}$ depends only on $\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})}$, see [14].

Set $\phi=u_{\varepsilon, R, n}-w$. Then $\phi$ satisfies

$$
\begin{aligned}
& \frac{\partial \phi}{\partial t}-D\left(\left(|D \phi|^{2}+\frac{1}{n}\right)^{(p-2) / 2} D \phi\right) \\
= & \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right)-M_{\varepsilon} f(w) \\
\leq & M_{\varepsilon}\left(f\left(u_{\varepsilon, R, n}\right)-f(w)\right) \\
= & M_{\varepsilon}\left(u_{\varepsilon, R, n}-w\right) \int_{0}^{1} f^{\prime}\left(\theta u_{\varepsilon, R, n}+(1-\theta) w\right) d \theta \\
= & M_{\varepsilon} C_{\varepsilon, R, n}(x, t)\left(u_{\varepsilon, R, n}-w\right) \\
= & M_{\varepsilon} C_{\varepsilon, R, n}(x, t) \phi .
\end{aligned}
$$

Thus $\phi$ satisfies

$$
\frac{\partial \phi}{\partial t}-D\left(\left(|D \phi|^{2}+\frac{1}{n}\right)^{(p-2) / 2} D \phi\right)-M_{\varepsilon} C_{\varepsilon, R, n} \phi \leq 0
$$

and $\phi( \pm R, t) \leq 0$ on $t \in\left[0, T_{0}\right], \phi(x, 0) \leq 0$ on $x \in(-R, R)$. By the maximum principle, $\phi \leq 0$ on $Q_{R, T_{0}}$. Hence, there exists $T^{\prime} \in\left(0, T_{0}\right)$, such that

$$
\left\|u_{\varepsilon, R, n}\right\|_{L^{\infty}\left(Q_{R, T^{\prime}}\right)} \leq \max _{t \in\left(0, T^{\prime}\right)} w(t) .
$$

Taking $T^{\prime}=T_{0} / 2$ and $C_{1}=w\left(T^{\prime}\right)$, we obtain

$$
\left\|u_{\varepsilon, R, n}\right\|_{L^{\infty}\left(Q_{R, T^{\prime}}\right)} \leq C_{1} .
$$

The proof is complete.
Lemma 2.2. There exists a positive constant $C_{2}$ depending on $\varepsilon$ and $T^{\prime}$, such that

$$
\iint_{Q_{R, T^{\prime}}}\left|D u_{\varepsilon, R, n}\right|^{p} d x d t \leq C_{2} .
$$

Proof. Multiplying (2.1) by $u_{\varepsilon, R, n}$, and integrating it over $Q_{R, T^{\prime}}$, we get

$$
\begin{aligned}
& \frac{1}{2} \int_{-R}^{R} u_{\varepsilon, R, n}^{2}\left(x, T^{\prime}\right) d x-\frac{1}{2} \int_{-R}^{R} u_{0}(x) d x \\
& \quad+\iint_{Q_{R, T^{\prime}}}\left(\left|D u_{\varepsilon, R, n}\right|^{2}+\frac{1}{n}\right)^{(p-2) / 2}\left|D u_{\varepsilon, R, n}\right|^{2} d x d t \\
& =\iint_{Q_{R, T^{\prime}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right) u_{\varepsilon, R, n} d x d t .
\end{aligned}
$$

By virtue of Lemma 2.1, we further obtain

$$
\begin{aligned}
& \iint_{Q_{R, T^{\prime}}}\left|D u_{\varepsilon, R, n}\right|^{p} d x d t \\
\leq & \iint_{Q_{R, T^{\prime}}}\left(\left|D u_{\varepsilon, R, n}\right|^{2}+\frac{1}{n}\right)^{(p-2) / 2}\left|D u_{\varepsilon, R, n}\right|^{2} d x d t \\
\leq & \frac{1}{2} \int_{-R}^{R} u_{0}(x) d x-\frac{1}{2} \int_{-R}^{R} u_{\varepsilon, R, n}^{2}\left(x, T^{\prime}\right) d x \\
& +C_{1} \iint_{Q_{R, T^{\prime}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right) d x d t .
\end{aligned}
$$

According to the assumption on $f$ and $u_{0}$, there exists a positive constant $C_{2}$, such that

$$
\iint_{Q_{R, T^{\prime}}}\left|D u_{\varepsilon, R, n}\right|^{p} d x d t \leq C_{2} .
$$

The proof is complete.
Lemma 2.3. There exists a positive constant $C_{3}$ depending on $\varepsilon$ and $T^{\prime}$, such that

$$
\iint_{Q_{R, T^{\prime}}}\left|\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right|^{2} d x d t \leq C_{3} .
$$

Proof. Multiplying (2.1) by $\frac{\partial u_{\varepsilon, R, n}}{\partial t}$ and integrating on $Q_{R, T^{\prime}}$, we have

$$
\begin{aligned}
& \iint_{Q_{R, T^{\prime}}}\left|\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right|^{2} d x d t \\
& +\iint_{Q_{R, T^{\prime}}}\left(\left|D u_{\varepsilon, R, n}\right|^{2}+\frac{1}{n}\right)^{(p-2) / 2} D u_{\varepsilon, R, n} D\left(\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right) d x d t \\
= & \iint_{Q_{R, T^{\prime}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right) \frac{\partial u_{\varepsilon, R, n}}{\partial t} d x d t .
\end{aligned}
$$

Since

$$
\begin{aligned}
& \iint_{Q_{R, T^{\prime}}}\left(\left|D u_{\varepsilon, R, n}\right|^{2}+\frac{1}{n}\right)^{(p-2) / 2} D u_{\varepsilon, R, n} D\left(\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right) d x d t \\
= & \frac{1}{2} \int_{-R}^{R} \int_{0}^{T^{\prime}} \frac{\partial}{\partial t}\left(\int_{0}^{\left|D u_{\varepsilon, R, n}\right|^{2}}\left(s+\frac{1}{n}\right)^{(p-2) / 2} d s\right) d t d x \\
= & \frac{1}{2} \int_{-R}^{R}\left(\int_{0}^{\left|D u_{\varepsilon, R, n}\left(x, T^{\prime}\right)\right|^{2}}\left(s+\frac{1}{n}\right)^{(p-2) / 2} d s\right) d x \\
& -\frac{1}{2} \int_{-R}^{R}\left(\int_{0}^{\left|D u_{0}(x)\right|^{2}}\left(s+\frac{1}{n}\right)^{(p-2) / 2} d s\right) d x \\
= & \frac{1}{p} \int_{-R}^{R}\left(\left|D u_{\varepsilon, R, n}\left(x, T^{\prime}\right)\right|^{2}+\frac{1}{n}\right)^{p / 2} d x \\
& -\frac{1}{p} \int_{-R}^{R}\left(\left|D u_{0}(x)\right|^{2}+\frac{1}{n}\right)^{p / 2} d x,
\end{aligned}
$$

and

$$
\begin{aligned}
& \iint_{Q_{R, T^{\prime}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right) \frac{\partial u_{\varepsilon, R, n}}{\partial t} d x d t \\
\leq & C_{\varepsilon} \iint_{Q_{R, T^{\prime}}}\left|f\left(u_{\varepsilon, R, n}\right)\right|^{2} d x d t+\frac{1}{2} \iint_{Q_{R, T^{\prime}}}\left|\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right|^{2} d x d t
\end{aligned}
$$

it follows that

$$
\begin{aligned}
& \iint_{Q_{R, T^{\prime}}}\left|\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right|^{2} d x d t \\
\leq & \frac{2}{p} \int_{-R}^{R}\left(\left|D u_{0}(x)\right|^{2}+\frac{1}{n}\right)^{p / 2} d x+2 C_{\varepsilon} \iint_{Q_{R, T^{\prime}}}\left|f\left(u_{\varepsilon, R, n}\right)\right|^{2} d x d t .
\end{aligned}
$$

According to Lemma 2.1-2.2, there exists a positive constant $C_{3}$, such that

$$
\iint_{Q_{R, T^{\prime}}}\left|\frac{\partial u_{\varepsilon, R, n}}{\partial t}\right|^{2} d x d t \leq C_{3}
$$

The proof is complete.
Utilizing the results in [15], we can easily obtain the uniform Hölder norm estimates on $u_{\varepsilon, R, n}$. From Lemma 2.1-2.3, we conclude that there exists a subsequence of $u_{\varepsilon, R, n}$, denoted by $u_{\varepsilon, R, n}$ itself, and a function $u_{\varepsilon, R}$ in $Q_{R, T^{\prime}}$, such that

$$
\begin{aligned}
& u_{\varepsilon, R, n} \rightarrow u_{\varepsilon, R}, \\
& \frac{\partial u_{\varepsilon, R, n}}{\partial t} \rightharpoonup \frac{\partial u_{\varepsilon, R}}{\partial t}, \quad \text { a.e. in } Q_{R, T^{\prime}}, \\
& \left(\left|D u_{\varepsilon, R, n}\right|^{2}+\frac{1}{n}\right)^{(p-2) / 2} \quad D u_{\varepsilon, R, n} \rightharpoonup \zeta, \quad \text { in } L^{p /(p-1)}\left(Q_{R, T^{\prime}}\right),
\end{aligned}
$$

Using the methods in [16], we can get

$$
\iint_{Q_{R, T^{\prime}}}\left|D u_{\varepsilon, R}\right|^{p-2} D u_{\varepsilon, R} D \varphi d x d t=\iint_{Q_{R, T^{\prime}}} \zeta D \varphi d x d t
$$

where $\varphi \in C^{\infty}\left(\bar{Q}_{R, T^{\prime}}\right)$, which vanishes for $t=T^{\prime}, x= \pm R$.
Consider the following problem

$$
\begin{align*}
\frac{\partial u_{\varepsilon}}{\partial t}-D\left(\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon}\right) & =\delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right), & & (x, t) \in Q_{R, T^{\prime}}  \tag{2.6}\\
u_{\varepsilon}(x, 0) & =u_{0}(x), & & x \in(-R, R)  \tag{2.7}\\
u_{\varepsilon}( \pm R, t) & =0, & & t \in\left(0, T^{\prime}\right) . \tag{2.8}
\end{align*}
$$

Due to the degeneracy of Eq. (2.6), we should give the definition of generalized solutions of the problem (2.6)-(2.8).

Definition 2.1. A nonnegative function $u_{\varepsilon}: Q_{R, T^{\prime}} \rightarrow[0,+\infty)$ is called a generalized solution of problem (2.6)-(2.8) on $Q_{R, T^{\prime}}$, if $u_{\varepsilon} \in C\left(0, T^{\prime} ; L^{\infty}(-R, R)\right) \cap L^{p}\left(0, T^{\prime}\right.$; $\left.W_{0}^{1, p}(-R, R)\right), \frac{\partial u_{\varepsilon}}{\partial t} \in L^{2}\left(0, T^{\prime} ; L^{2}(-R, R)\right)$, and for all $0<t_{1}<t_{2}<T^{\prime}, u_{\varepsilon}$ satisfies

$$
\begin{align*}
& \int_{-R}^{R} u_{\varepsilon}\left(x, t_{2}\right) \varphi\left(x, t_{2}\right) d x-\int_{-R}^{R} u_{\varepsilon}\left(x, t_{1}\right) \varphi\left(x, t_{1}\right) d x=\int_{t_{1}}^{t_{2}} \int_{-R}^{R} u_{\varepsilon} \varphi_{t} d x d t \\
& -\int_{t_{1}}^{t_{2}} \int_{-R}^{R}\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon} D \varphi d x d t+\int_{t_{1}}^{t_{2}} \int_{-R}^{R} \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) \varphi(x, t) d x d t \tag{2.9}
\end{align*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} \int_{-R}^{R} u_{\varepsilon}(x, t) \psi(x) d x=\int_{-R}^{R} u_{0}(x) \psi(x) d x \tag{2.10}
\end{equation*}
$$

where $\varphi \in C^{\infty}\left(\bar{Q}_{R, T^{\prime}}\right)$, which vanishes for $t=T^{\prime}, x= \pm R$, and $\psi(x) \in C_{0}^{\infty}(-R, R)$.
Similarly, to define a lower(super) solution $\underline{u}_{\varepsilon}(x, t)\left(\bar{u}_{\varepsilon}(x, t)\right)$, we need only to ask $\varphi(x, t) \geq 0, \psi(x) \geq 0, \underline{u}_{\varepsilon}( \pm R, t) \leq 0\left(\bar{u}_{\varepsilon}( \pm R, t) \geq 0\right)$ in $\left(0, T^{\prime}\right)$ and the equality in (2.9) and (2.10) is replaced by $\leq(\geq)$.

Utilizing the methods in [15], we can derive the following result.
Lemma 2.4. Suppose that $\underline{u}_{\varepsilon}$ and $\bar{u}_{\varepsilon}$ are lower and super solutions of (2.6)-(2.8) respectively, then $\underline{u}_{\varepsilon} \leq \bar{u}_{\varepsilon}$ a.e. in $Q_{R, T^{\prime}}$.

By using a standard limiting process and Lemma 2.4, it is easy to see that $u_{\mathcal{\varepsilon}, R}$ is a unique generalized solution of the problem (2.6)-(2.8).

Lemma 2.5. There exist positive constants $\lambda_{1}$ and $\lambda_{2}$ depending on $R$, such that

$$
u_{\varepsilon, R}(x, t) \leq \lambda_{1}-\int_{-R}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s, \quad \forall(x, t) \in Q_{R, T^{\prime}}
$$

where $H(x)$ is the Heaviside function.
Proof. Let

$$
\bar{u}_{\varepsilon, R}(x, t)=\lambda_{1}-\int_{-R}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s, \quad(x, t) \in Q_{R, T^{\prime}}
$$

where $\lambda_{1}$ and $\lambda_{2}$ are two positive constants to be determined. Then
$\int_{t_{1}}^{t_{2}} \int_{-R}^{R} \frac{\partial \bar{u}_{\varepsilon, R}}{\partial t} \varphi d x d t-\int_{t_{1}}^{t_{2}} \int_{-R}^{R} D\left(\left|D \bar{u}_{\varepsilon, R}\right|^{p-2} D \bar{u}_{\varepsilon, R}\right) \varphi d x d t=\int_{t_{1}}^{t_{2}} \int_{-R}^{R} \lambda_{2} \delta_{\varepsilon}(x) \varphi d x d t$,
where $0<t_{1}<t_{2}<T^{\prime}, \varphi(x, t) \geq 0, \varphi \in C^{\infty}\left(\bar{Q}_{R, T^{\prime}}\right)$, which vanishes for $t=T^{\prime}, x= \pm R$.

It is easy to see that

$$
\int_{-R}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s=0, \quad \text { for } x \leq-\varepsilon
$$

and for $x>-\varepsilon$, we have

$$
\begin{aligned}
& \int_{-R}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s \\
& =\int_{-\varepsilon}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s \\
& \leq \int_{-\varepsilon}^{x}\left(\lambda_{2} \int_{\mathbb{R}} \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s \\
& \leq \lambda_{2}^{1 /(p-1)}(|x|+1)
\end{aligned}
$$

According to the assumption on $f$, there exists a positive constant $M_{0}$, such that

$$
f(s) \leq M_{0}, \quad \forall s \in \mathbb{R}
$$

Choosing a positive constant $\lambda_{2} \geq M_{0}$ and $\lambda_{1}-\lambda_{2}^{1 /(p-1)}(R+1) \geq\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})}$. Thus we can obtain

$$
\begin{aligned}
\int_{t_{1}}^{t_{2}} \int_{-R}^{R} \lambda_{2} \delta_{\varepsilon}(x) \varphi d x d t \geq \int_{t_{1}}^{t_{2}} \int_{-R}^{R} \delta_{\varepsilon}(x) f\left(\bar{u}_{\varepsilon, R}\right) \varphi d x d t \\
\forall(x, t) \in(-R, R) \times\left(t_{1}, t_{2}\right)
\end{aligned}
$$

Noticing that

$$
u_{0}(x) \leq\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})} \leq \lambda_{1}-\lambda_{2}^{1 /(p-1)}(R+1) \leq \bar{u}_{\varepsilon, R}(x, 0), \quad \forall x \in(-R, R)
$$

and

$$
\bar{u}_{\varepsilon, R}( \pm R, t) \geq 0, \quad \forall t \in\left(0, T^{\prime}\right)
$$

thus $\bar{u}_{\varepsilon, R}$ is a super solution of the problem (2.6)-(2.8). By Lemma 2.4, we get

$$
u_{\varepsilon, R}(x, t) \leq \lambda_{1}-\int_{-R}^{x}\left(\lambda_{2} \int_{\mathbb{R}} H(s-y) \delta_{\varepsilon}(y) d y\right)^{1 /(p-1)} d s, \quad \forall(x, t) \in Q_{R, T^{\prime}}
$$

The proof is complete.
Now, we prove the property of finite propagation of disturbances of $u_{\varepsilon, R}$.
Lemma 2.6. Let supp $u_{0} \subset I, r_{1}<0<r_{2}$, where $I=\left[r_{1}, r_{2}\right] \subset(-R, R)$ is a closed set in $\mathbb{R}$. Then there exists $0<T^{*}<T^{\prime}$, such that

$$
\operatorname{supp} u_{\varepsilon, R}(\cdot, t) \subset\left[R_{1}, R_{2}\right], \quad \text { a.e. } t \in\left(0, T^{*}\right)
$$

where

$$
R_{1}=r_{1}-C_{4} T^{* \mu}, \quad R_{2}=r_{2}+C_{5} T^{* \mu}
$$

with the positive constants $C_{4}, C_{5}, \mu$ depending on $p$ and $R$.
Proof. Set

$$
g_{m}(y)=\iint_{Q_{R, t}}(x-y)_{+}^{m}\left|D u_{\varepsilon, R}\right|^{p} d x d \tau, \quad m=1,2, \cdots, y \geq r_{2}, 0<t \leq T^{\prime}
$$

Using the techniques in [11], we can obtain

$$
g_{1}^{\prime}(y) \leq-C t^{-\lambda /(\theta+1)}\left[g_{1}(y)\right]^{1 /(\theta+1)},
$$

where

$$
\lambda=\frac{1-a}{\sigma(2 p+1)}, \quad \theta=\frac{p \gamma}{\sigma(2 p+1)^{2}}-\frac{1}{2 p+1},
$$

with the positive constants

$$
a=\frac{\frac{1}{2}+\frac{1}{p+2}-\frac{1}{p}}{\frac{1}{2}+\frac{2}{p+2}-\frac{1}{p}}, \quad \sigma=1-\frac{p+1}{2 p+1} \gamma, \quad \gamma=a+(1-a) \frac{p}{2} .
$$

If $g_{1}\left(r_{2}\right)=0$, then $D u_{\varepsilon, R}(x, t)=0$ for $x \in\left[r_{2}, R\right]$, and hence from the boundary value condition, we see that $u_{\varepsilon, R}(x, t)=0$ for $x \in\left[r_{2}, R\right]$, i.e., supp $u_{\varepsilon, R}(\cdot, t) \subset$ $\left[-R, r_{2}\right]$. If $g_{1}\left(r_{2}\right) \neq 0$, then there exists an interval $\left(r_{2}, R^{*}\right)$, such that $g_{1}(y)>0$ in $\left(r_{2}, R^{*}\right)$, but $g_{1}\left(R^{*}\right)=0$. Therefore, for $y \in\left(r_{2}, R^{*}\right)$, we have

$$
\left(g_{1}(y)^{\theta /(\theta+1)}\right)^{\prime}=\frac{\theta}{\theta+1} \frac{g_{1}^{\prime}(y)}{g_{1}(y)^{1 /(\theta+1)}} \leq-C t^{-\lambda /(\theta+1)}
$$

Integrating the above inequality on $\left(r_{2}, R^{*}\right)$, we obtain

$$
g_{1}\left(R^{*}\right)^{\theta /(\theta+1)}-g_{1}\left(r_{2}\right)^{\theta /(\theta+1)} \leq-C t^{-\lambda /(\theta+1)}\left(R^{*}-r_{2}\right)
$$

Therefore

$$
R^{*} \leq r_{2}+C t^{\lambda /(\theta+1)} g_{1}\left(r_{2}\right)^{\theta /(\theta+1)}
$$

According to Lemma 2.2, we have

$$
\iint_{Q_{R, t}}\left|D u_{\varepsilon, R, n}\right|^{p} d x d t \leq \frac{1}{2} \int_{-R}^{R} u_{0}(x) d x+\iint_{Q_{R, t}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R, n}\right) u_{\varepsilon, R, n} d x d t .
$$

Letting $n \rightarrow \infty$ and using Lemma 2.5, we further obtain

$$
\begin{aligned}
& \iint_{Q_{R, t}}\left|D u_{\varepsilon, R}\right|^{p} d x d t \\
& \leq \frac{1}{2} \int_{-R}^{R} u_{0}(x) d x+\iint_{Q_{R, t}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon, R}\right) u_{\varepsilon, R} d x d t \\
& \leq \frac{1}{2} \int_{-R}^{R} u_{0}(x) d x+\lambda_{1} M_{0} \iint_{Q_{R, t}} \delta_{\varepsilon}(x) d x d t \\
& \leq M_{1}+M_{2} t
\end{aligned}
$$

which implies that

$$
g_{1}\left(r_{2}\right)=\iint_{Q_{R, t}}\left(x-r_{2}\right)_{+}\left|D u_{\varepsilon, R}\right|^{p} d x d \tau \leq\left(R-r_{2}\right)\left(M_{1}+M_{2} t\right)
$$

and hence

$$
R^{*} \leq r_{2}+\left(M_{3}+M_{4} t\right) t^{\lambda /(\theta+1)}
$$

Obviously, there exists $0<T_{1}^{*}<T^{\prime}$, such that

$$
r_{2}+\left(M_{3}+M_{4} T_{1}^{*}\right) T_{1}^{* \lambda /(\theta+1)}=r_{2}+C_{5} T_{1}^{* \mu}<R
$$

which implies

$$
\operatorname{supp} u_{\varepsilon, R}(\cdot, t) \subset\left[-R, r_{2}+C_{5} T_{1}^{* \mu}\right], \quad \text { a.e. } t \in\left(0, T_{1}^{*}\right)
$$

Similarly, there exists $0<T_{2}^{*}<T^{\prime}$, such that

$$
\operatorname{supp} u_{\varepsilon, R}(\cdot, t) \subset\left[r_{1}-C_{4} T_{2}^{* \mu}, R\right], \quad \text { a.e. } t \in\left(0, T_{2}^{*}\right)
$$

Taking $T^{*}=\min \left\{T_{1}^{*}, T_{2}^{*}\right\}$, we obtain

$$
\operatorname{supp} u_{\varepsilon, R}(\cdot, t) \subset\left[r_{1}-C_{4} T^{* \mu}, r_{2}+C_{5} T^{* \mu}\right]=\left[R_{1}, R_{2}\right], \quad \text { a.e. } t \in\left(0, T^{*}\right)
$$

The proof is complete.
Considering the Cauchy problem

$$
\begin{array}{rlrl}
\frac{\partial u_{\varepsilon}}{\partial t}-D\left(\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon}\right) & =\delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right), & & (x, t) \in Q_{T^{*}} \\
u_{\varepsilon}(x, 0) & =u_{0}(x), & x \in \mathbb{R}, \tag{2.12}
\end{array}
$$

we can also give the definition of generalized solution of the problem (2.11)(2.12).

Definition 2.2. A nonnegative function $u_{\varepsilon}: Q_{T^{*}} \rightarrow[0,+\infty)$ is called a generalized solution of the problem (2.11)-(2.12) in $Q_{T^{*}}$, if $u_{\varepsilon} \in C\left(0, T^{*} ; L^{\infty}(\mathbb{R})\right) \cap L^{p}\left(0, T^{*}\right.$; $\left.W^{1, p}(\mathbb{R})\right), \frac{\partial u_{\varepsilon}}{\partial t} \in L^{2}\left(0, T^{*} ; L^{2}(\mathbb{R})\right)$, and $u_{\varepsilon}$ satisfies

$$
\begin{aligned}
\int_{\mathbb{R}} u_{0}(x) \varphi(x, 0) d x+ & \iint_{Q_{T^{*}}}\left(u_{\varepsilon} \varphi_{t}-\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon} D \varphi\right) d x d t \\
& +\iint_{Q_{T^{*}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) \varphi(x, t) d x d t=0
\end{aligned}
$$

for any $\varphi \in C^{\infty}\left(Q_{T^{*}}\right)$, which vanishes for large $|x|$ and $t=T^{*}$.
Let

$$
u_{\varepsilon}(x, t)= \begin{cases}u_{\varepsilon, R}(x, t), & x \in[-R, R], \\ 0, & x \in \mathbb{R} \backslash[-R, R] .\end{cases}
$$

Then by Lemma 2.6, we can see that $u_{\varepsilon}(x, t)$ is a generalized solution of problem (2.11)-(2.12) on $Q_{T^{*}}$.

We want to use Moser iteration to estimate the local boundedness on $u_{\varepsilon}$. Firstly, we need the following result.

Lemma 2.7. Let $x_{0} \neq 0,\left(x_{0}-2 \rho, x_{0}+2 \rho\right) \subset \mathbb{R} \backslash(-\varepsilon, \varepsilon)$. For any $\beta>0$, it has

$$
\int_{0}^{T^{*}} \int_{x_{0}-\rho}^{x_{0}+\rho} u_{\varepsilon}^{\beta+2(p-1)} d x d t \leq C_{6}
$$

where $C_{6}$ is a positive constant depending on $\beta, p, \rho, T^{*}$ and $\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})}$.
Proof. Let $\eta(x) \in C_{0}^{\infty}(\mathbb{R})$, and

$$
\begin{aligned}
& 0 \leq \eta(x) \leq 1 ; \quad \eta(x)=1, \quad \forall x \in\left(x_{0}-\rho, x_{0}+\rho\right) ; \\
& \eta(x)=0, \quad \forall x \in \mathbb{R} \backslash\left(x_{0}-2 \rho, x_{0}+2 \rho\right) ; \quad\left|\eta^{\prime}(x)\right| \leq \frac{C}{\rho}
\end{aligned}
$$

Multiplying (2.1) by $\varphi=\eta^{2 q} u_{\varepsilon}^{\beta}$ and integrating on $\left(x_{0}-2 \rho, x_{0}+2 \rho\right) \times(0, t)$, it is easy to obtain

$$
\int_{0}^{t} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \frac{\partial u_{\varepsilon}}{\partial \sigma}\left(\eta^{2 q} u_{\varepsilon}^{\beta}\right) d x d \sigma+\int_{0}^{t} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon} D\left(\eta^{2 q} u_{\varepsilon}^{\beta}\right) d x d \sigma=0
$$

where $0<t \leq T^{*}, q>p / 2$ is a constant to be determined. Then

$$
\begin{aligned}
& \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \frac{1}{\beta+1} \eta^{2 q} u_{\varepsilon}^{\beta+1} d x+\int_{0}^{t} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \beta \eta^{2 q} u_{\varepsilon}^{\beta-1}\left|D u_{\varepsilon}\right|^{p} d x d \sigma \\
\leq & C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{q}\left|\left(\eta^{q}\right)^{\prime}\right| \cdot u_{\varepsilon}^{\beta} \cdot\left|D u_{\varepsilon}\right|^{p-1} d x d t+\int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \frac{1}{\beta+1} \eta^{2 q}\left(u_{0}(x)\right)^{\beta+1} d x
\end{aligned}
$$

Taking the supremum with respect to $t$ and utilizing Young's inequality, we have

$$
\begin{aligned}
& \sup _{t \in\left(0, T^{*} *\right.} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \frac{1}{\beta+1} \eta^{2 q} u_{\varepsilon}^{\beta+1} d x+\int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \beta \eta^{2 q} u_{\varepsilon}^{\beta-1}\left|D u_{\varepsilon}\right|^{p} d x d t \\
& \leq C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{q}\left|\left(\eta^{q}\right)^{\prime}\right| \cdot u_{\varepsilon}^{\beta} \cdot\left|D u_{\varepsilon}\right|^{p-1} d x d t+M_{5} \\
& \leq C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q-p}\left|\eta^{\prime}\right|^{p} u_{\varepsilon}^{\beta+p-1} d x d t+\frac{\beta}{2} \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta-1}\left|D u_{\varepsilon}\right|^{p} d x d t+M_{5},
\end{aligned}
$$

where $M_{5}$ is a positive constant depending on $\beta, \rho$ and $\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})}$. Hence

$$
\begin{align*}
& \sup _{t \in\left(0, T^{*}\right)} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta+1} d x+\int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta-1}\left|D u_{\varepsilon}\right|^{p} d x d t \\
& \leq C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q-p}\left|\eta^{\prime}\right|^{p} u_{\varepsilon}^{\beta+p-1} d x d t+M_{6} . \tag{2.13}
\end{align*}
$$

For any $t \in\left(0, T^{*}\right)$, using Hölder's inequality and Lemma 2.5, we have

$$
\begin{aligned}
& \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta+2(p-1)} d x \\
\leq & \left(\int_{x_{0}-2 \rho}^{x_{0}+2 \rho} u_{\varepsilon}^{p} d x\right)^{(p-1) / p}\left(\int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left(\eta^{2 q} u_{\varepsilon}^{\beta+p-1}\right)^{p} d x\right)^{1 / p} \\
\leq & C\left(\int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left(\eta^{2 q / p} u_{\varepsilon}^{(\beta+p-1) / p}\right)^{p^{2}} d x\right)^{\frac{1}{p}} \\
\leq & C\left(\sup _{x \in\left(x_{0}-2 \rho, x_{0}+2 \rho\right)}\left(\eta^{2 q / p} u_{\varepsilon}^{(\beta+p-1) / p}\right)\right)^{p} .
\end{aligned}
$$

Utilizing the embedding theorem, it has

$$
\sup _{x \in\left(x_{0}-2 \rho, x_{0}+2 \rho\right)}\left(\eta^{2 q / p} u_{\varepsilon}^{(\beta+p-1) / p}\right) \leq C\left(\int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left|D\left(\eta^{2 q / p} u_{\varepsilon}^{(\beta+p-1) / p}\right)\right|^{p} d x\right)^{1 / p}
$$

Thus

$$
\int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta+2(p-1)} d x \leq C \int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left|D\left(\eta^{2 q / p} u_{\varepsilon}^{(\beta+p-1) / p}\right)\right|^{p} d x .
$$

Integrating the above inequality with respect to $t$ on $\left(0, T^{*}\right)$, and utilizing (2.13), we have

$$
\begin{aligned}
& \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta+2(p-1)} d x d t \\
\leq & C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left(\eta^{2 q} u_{\varepsilon}^{\beta-1}\left|D u_{\varepsilon}\right|^{p}+\eta^{2 q-p}\left|\eta^{\prime}\right|^{p} u_{\varepsilon}^{\beta+p-1}\right) d x d t \\
\leq & C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q-p}\left|\eta^{\prime}\right|^{p} u_{\varepsilon}^{\beta+p-1} d x d t+M_{7} .
\end{aligned}
$$

Applying Young's inequality, it follows that

$$
\begin{aligned}
& C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q-p}\left|\eta^{\prime}\right|^{p} u_{\varepsilon}^{\beta+p-1} d x d t \\
\leq & \frac{1}{2} \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{(2 q-p) \frac{\beta+2(p-1)}{\beta+p-1}} u_{\varepsilon}^{\beta+2(p-1)} d x d t+C \int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho}\left|\eta^{\prime}\right|^{p \cdot \frac{\beta+2(p-1)}{p-1}} d x d t .
\end{aligned}
$$

Let $(2 q-p) \frac{\beta+2(p-1)}{\beta+p-1}=2 q$. Then $q=\frac{p[\beta+2(p-1)]}{2(p-1)}>\frac{p}{2}$. Hence, there exists a positive constant $C_{6}$, such that

$$
\int_{0}^{T^{*}} \int_{x_{0}-2 \rho}^{x_{0}+2 \rho} \eta^{2 q} u_{\varepsilon}^{\beta+2(p-1)} d x d t \leq C_{6} .
$$

Therefore

$$
\int_{0}^{T^{*}} \int_{x_{0}-\rho}^{x_{0}+\rho} u_{\varepsilon}^{\beta+2(p-1)} d x d t \leq C_{6} .
$$

The proof is complete.
Next, we consider the local boundedness of $u_{\varepsilon}$ by using the technique of Moser's iteration.

Lemma 2.8. For $0<\tau<1$, there exists a positive constant $C_{\tau}$ depending on $\tau, p, \rho, T^{*}$ and $\left\|u_{0}(x)\right\|_{L^{\infty}(\mathbb{R})}$, such that

$$
u_{\varepsilon}(x, t) \leq C_{\tau}, \quad \forall(x, t) \in Q_{\tau \rho}
$$

where $Q_{\rho}=Q_{\rho}\left(x_{0}, t_{0}\right)=\left(x_{0}-\rho, x_{0}+\rho\right) \times\left(t_{0}-\rho^{p}, t_{0}+\rho^{p}\right), 0<t_{0}<T^{*}$, and $\left(t_{0}-\rho^{p}, t_{0}+\rho^{p}\right) \subset\left(0, T^{*}\right)$.

Proof. For any $\tau \rho \leq h<h^{\prime} \leq \rho$, take $\eta(x) \in C_{0}^{\infty}(\mathbb{R})$, and

$$
\begin{aligned}
& 0 \leq \eta(x) \leq 1 ; \quad \eta(x)=1, \quad \forall x \in\left(x_{0}-h, x_{0}+h\right) ; \\
& \eta(x)=0, \quad \forall x \in \mathbb{R} \backslash\left(x_{0}-h^{\prime}, x_{0}+h^{\prime}\right) ; \quad\left|\eta^{\prime}(x)\right| \leq \frac{C}{h^{\prime}-h} .
\end{aligned}
$$

And for any $s \in\left(t_{0}-h^{p}, t_{0}+h^{p}\right)$, we choose a function $\xi(t) \in C^{\infty}(-\infty, s]$, such that

$$
\begin{aligned}
& \xi(t)=1, \quad \forall t \in\left[t_{0}-h^{p}, s\right] ; \quad \xi(t)=0, \quad \forall t \in\left(-\infty, t_{0}-h^{\prime p}\right] \\
& 0 \leq \xi(t) \leq 1, \quad 0 \leq \xi^{\prime}(t) \leq \frac{C}{\left(h^{\prime}-h\right)^{p}}, \quad \forall t \leq s
\end{aligned}
$$

and extend it to be zero for $t>s$. Let $v_{\varepsilon}(x, t)=\max \left\{u_{\varepsilon}(x, t), 1\right\}$. Multiplying (2.11) by $\varphi=\tilde{\zeta}^{2} \eta^{p} v_{\varepsilon}^{\gamma}$ and integrating on $Q_{\rho}$, we can obtain

$$
\iint_{Q_{\rho}} \frac{\partial v_{\varepsilon}}{\partial t}\left(\tilde{\zeta}^{2} \eta^{p} v_{\varepsilon}^{\gamma}\right) d x d t+\iint_{Q_{\rho}}\left|D v_{\varepsilon}\right|^{p-2} D v_{\varepsilon} D\left(\tilde{\zeta}^{2} \eta^{p} v_{\varepsilon}^{\gamma}\right) d x d t=0 .
$$

Thus

$$
\begin{aligned}
& \frac{1}{\gamma+1} \iint_{Q_{h^{\prime}}^{s}} \frac{\partial}{\partial t}\left(\tilde{\xi}^{2} \eta^{p} v_{\varepsilon}^{\gamma+1}\right) d x d t-\frac{2}{\gamma+1} \iint_{Q_{h^{\prime}}^{s}} \xi \xi^{\prime} \eta^{p} v_{\varepsilon}^{\gamma+1} d x d t \\
& +\gamma \iint_{Q_{h^{\prime}}^{s}} \xi^{2} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t+p \iint_{Q_{h^{\prime}}^{s}} \xi^{2} \eta^{p-1} \eta^{\prime} v_{\varepsilon}^{\gamma}\left|D v_{\varepsilon}\right|^{p-2} D v_{\varepsilon} d x d t=0,
\end{aligned}
$$

where $Q_{h^{\prime}}^{S}=\left(x_{0}-h^{\prime}, x_{0}+h^{\prime}\right) \times\left(t_{0}-h^{\prime p}, s\right)$.
Note that for $t=t_{0}-h^{\prime 2}$, it has $\xi(t)=0$. Hence

$$
\iint_{Q_{h^{\prime}}^{s}} \frac{\partial}{\partial t}\left(\tilde{\zeta}^{2} \eta^{p} v_{\varepsilon}^{\gamma+1}\right) d x d t=\left.\int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \tilde{\zeta}^{2} \eta^{p} v_{\varepsilon}^{\gamma+1}\right|_{t=s} d x .
$$

Using Young's inequality, we have

$$
\begin{aligned}
& p \iint_{Q_{h^{\prime}}^{s}} \xi^{2} \eta^{p-1} \eta^{\prime} v_{\varepsilon}^{\gamma}\left|D v_{\varepsilon}\right|^{p-2} D v_{\varepsilon} d x d t \\
\leq & \frac{1}{2} \iint_{Q_{h^{\prime}}^{s}} \tilde{\xi}^{2} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t+C \iint_{Q_{h^{\prime}}^{s}} \tilde{\xi}^{2}\left|\eta^{\prime}\right|^{p} v_{\varepsilon}^{\gamma+p-1} d x d t .
\end{aligned}
$$

Therefore

$$
\begin{aligned}
& \left.\frac{1}{\gamma+1} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \xi^{2} \eta^{p} v_{\varepsilon}^{\gamma+1}\right|_{t=s} d x+\gamma \iint_{Q_{h^{\prime}}^{s}} \tilde{\xi}^{2} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t \\
\leq & \frac{1}{2} \iint_{Q_{h^{\prime}}^{s}} \xi^{2} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t+C \iint_{Q_{h^{\prime}}^{s}} \xi^{2}\left|\eta^{\prime}\right|^{p} v_{\varepsilon}^{\gamma+p-1} d x d t \\
& +\frac{2}{\gamma+1} \iint_{Q_{h^{\prime}}^{s}} \xi \xi^{\prime} \eta^{p} v_{\varepsilon}^{\gamma+1} d x d t,
\end{aligned}
$$

it follows that

$$
\begin{align*}
& \quad \sup _{t \in\left(t_{0}-h^{p}, t_{0}+h^{p}\right)} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma+1} d x+\int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t \\
& \leq \frac{C}{\left(h^{\prime}-h\right)^{p}} \iint_{Q_{h^{\prime}}} v_{\varepsilon}^{\gamma+p-1} d x d t . \tag{2.14}
\end{align*}
$$

Let $\chi(t)$ be the characteristic function of the interval $\left[t_{0}-h^{p}, t_{0}+h^{p}\right]$. Utilizing the embedding theorem, we have

$$
\begin{aligned}
&\left(\frac{1}{\rho^{p+1}} \iint_{Q_{\rho}}\left|\chi(t) \eta^{p / 2} v_{\varepsilon}^{(\gamma+1) / 2}\right|^{2 q} d x d t\right)^{1 / q} \\
& \leq \frac{C}{\rho}\left(\sup _{t_{0}-\rho^{p} \leq t \leq t_{0}+\rho^{p}} \int_{x_{0}-\rho}^{x_{0}+\rho}\left|\chi(t) \eta^{p / 2} v_{\varepsilon}^{(\gamma+1) / 2}\right|^{2} d x+\iint_{Q_{\rho}}\left|\chi(t) D\left(\eta^{p / 2} v_{\varepsilon}^{(\gamma+1) / 2}\right)\right|^{2} d x d t\right) \\
& \leq \frac{C}{\rho}\left(\sup _{t_{0}-h^{p} \leq t \leq t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma+1} d x\right. \\
&\left.+\int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}}\left|\eta^{p / 2} D\left(v_{\varepsilon}^{(\gamma+1) / 2}\right)+\frac{p}{2} \eta^{(p-2) / 2} \eta^{\prime} v_{\varepsilon}^{(\gamma+1) / 2}\right|^{2} d x d t\right) \\
& \leq \frac{C}{\rho}\left(\sup _{t_{0}-h^{p} \leq t \leq t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma+1} d x+\frac{(\gamma+1)^{2}}{2} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{2} d x d t\right. \\
&\left.+\frac{p^{2}}{2} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p-2}\left|\eta^{\prime}\right|^{2} v_{\varepsilon}^{\gamma+1} d x d t\right) \\
& \leq \frac{C}{\rho}\left(\sup _{t_{0}-h^{p} \leq t \leq t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma+1} d x+M_{8} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma-1} d x d t\right. \\
&+M_{9} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p} v_{\varepsilon}^{\gamma-1}\left|D v_{\varepsilon}\right|^{p} d x d t \\
&+M_{10} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p-2} v_{\varepsilon}^{\gamma+1} d x d t \\
&\left.+M_{11} \int_{t_{0}-h^{p}}^{t_{0}+h^{p}} \int_{x_{0}-h^{\prime}}^{x_{0}+h^{\prime}} \eta^{p-2}\left|\eta^{\prime}\right|^{p} v_{\varepsilon}^{\gamma+1} d x d t\right) .
\end{aligned}
$$

Using (2.14), it gives

$$
\begin{aligned}
& \left(\frac{1}{\rho^{p+1}} \iint_{Q_{h}}\left|v_{\varepsilon}^{(\gamma+1) / 2}\right|^{2 q} d x d t\right)^{1 / q} \\
& \leq\left(\frac{1}{\rho^{p+1}} \iint_{Q_{\rho}}\left|\chi(t) \eta^{p / 2} v_{\varepsilon}^{(\gamma+1) / 2}\right|^{2 q}\right)^{1 / q} \\
& \leq \frac{C}{\rho\left(h^{\prime}-h\right)^{p}} \iint_{Q_{h^{\prime}}} v_{\varepsilon}^{\gamma+p-1} d x d t .
\end{aligned}
$$

In fact, since the dimension is 1 , according to the embedding theorem, we have $q=\frac{5}{3}$. Hence

$$
\frac{1}{\rho^{p+1}} \iint_{Q_{h}} v_{\varepsilon}^{\frac{5}{3}(\gamma+1)} d x d t \leq\left[\frac{C}{\rho\left(h^{\prime}-h\right)^{p}} \iint_{Q_{h^{\prime}}} v_{\varepsilon}^{\gamma+p-1} d x d t\right]^{5 / 3}
$$

Let

$$
\begin{aligned}
& h_{k}=\tau \rho\left(1+\frac{1-\tau}{\tau 2^{k-1}}\right), \quad h=h_{k+1}, \quad h^{\prime}=h_{k} ; \\
& \mu=\frac{5}{3}, \quad \mu^{k}=\gamma-\frac{3}{2} p+4 .
\end{aligned}
$$

Then

$$
\frac{1}{\rho^{p+1}} \iint_{Q_{h_{k+1}}} v_{\varepsilon}^{\mu^{k+1}+\frac{5}{2}(p-2)} d x d t \leq\left[\frac{C}{\rho^{p+1}} \iint_{Q_{h_{k}}} v_{\varepsilon}^{\mu^{k}+\frac{5}{2}(p-2)} d x d t\right]^{\mu},
$$

namely,

$$
\left[\frac{1}{\rho^{p+1}} \iint_{Q_{h_{k+1}}} v_{\varepsilon}^{\mu^{k+1}+\frac{5}{2}(p-2)} d x d t\right]^{1 / \mu^{k+1}} \leq C\left[\frac{1}{\rho^{p+1}} \iint_{Q_{h_{k}}} v_{\varepsilon}^{\mu^{k}+\frac{5}{2}(p-2)} d x d t\right]^{1 / \mu^{k}}
$$

Utilizing Moser iteration, we have

$$
\begin{equation*}
\sup _{Q_{\tau \rho}} v_{\varepsilon} \leq C\left[\frac{1}{\rho^{p+1}} \iint_{Q_{\rho}} v_{\varepsilon}^{\mu+\frac{5}{2}(p-2)} d x d t\right]^{1 / \mu} . \tag{2.15}
\end{equation*}
$$

According to Lemma 2.7 and the definition of $v_{\mathcal{\varepsilon}}$, for $p>\frac{8}{3}$, we can choose $\beta>0$, such that

$$
\beta+2(p-1)=\mu+\frac{5}{2}(p-2)
$$

namely, $\beta=\frac{1}{2} p-\frac{4}{3}$. Then there exists a constant $M_{12}$, such that

$$
\iint_{Q_{\rho}} v_{\varepsilon}^{\mu+\frac{5}{2}(p-2)} d x d t \leq M_{12} .
$$

For $2<p \leq \frac{8}{3}$ and $\beta>0$, it has $\beta+2(p-1)>\mu+\frac{5}{2}(p-2)$. Then there exists a constant $M_{13}$, such that

$$
\iint_{Q_{\rho}} v_{\varepsilon}^{\mu+\frac{5}{2}(p-2)} d x d t \leq C \iint_{Q_{\rho}} v_{\varepsilon}^{\beta+2(p-1)} d x d t \leq M_{13} .
$$

According to (2.15), there exists a constant $C_{\tau}$, such that

$$
u_{\varepsilon} \leq C_{\tau}, \quad \forall(x, t) \in Q_{\tau \rho} .
$$

The proof is complete.
Utilizing the results in [15], we can obtain the Hölder estimate on $u_{\varepsilon}$.
Lemma 2.9. For $0<\theta<\tau$, there exist positive constants $\lambda$ and $\sigma \in(0,1)$, such that
$\left|u_{\varepsilon}\left(x_{1}, t_{1}\right)-u_{\varepsilon}\left(x_{2}, t_{2}\right)\right| \leq C_{7}\left(\left|x_{1}-x_{2}\right|+\lambda\left|t_{1}-t_{2}\right|^{1 / p}\right)^{\sigma}, \quad \forall\left(x_{i}, t_{i}\right) \in Q_{\theta \rho}, \quad i=1,2$, where $C_{7}$ is a positive constant depending on $p, \theta, \tau, \rho, T^{*}$ and $C_{\tau}$.

In addition, it is easy to obtain the following estimate.
Lemma 2.10. There exists a positive constant $C_{8}$ independent of $\varepsilon$, such that

$$
\iint_{Q_{\tau \rho}}\left|D u_{\varepsilon}\right|^{p} d x d t \leq C_{8}
$$

Now, we are in the position to prove the main result.
Proof of Theorem 1.1. According to Lemma 2.8-2.10, there are a subsequence of $\left\{u_{\varepsilon}\right\}$ (without loss of generality, we denote it by $\left\{u_{\varepsilon}\right\}$ itself) and a function $u$, such that for any compact set $K \subset(\mathbb{R} \backslash\{0\}) \times\left(0, T^{*}\right)$, we have

$$
\begin{gathered}
u_{\varepsilon} \rightarrow u, \quad \text { a.e. in } K, \\
\left(\left|D u_{\varepsilon}\right|\right)^{p-2} D u_{\varepsilon} \rightharpoonup \omega, \quad \text { in } L^{p /(p-1)}(K),
\end{gathered}
$$

and we can also prove that the following equality

$$
\iint_{K}|D u|^{p-2} D u D \varphi d x d t=\iint_{K} \omega D \varphi d x d t
$$

holds for any $\varphi \in C^{\infty}\left(Q_{T^{*}}\right)$, which vanishes for large $|x|$ and $t=T^{*}$.
In addition, according to Lemma 2.5, Lemma 2.8-2.10, we have

$$
u \in L^{\infty}\left(0, T^{*} ; L_{l o c}^{\infty}(\mathbb{R})\right) \cap C\left((\mathbb{R} \backslash\{0\}) \times\left(0, T^{*}\right)\right) \cap B V_{x}\left(Q_{T^{*}}\right)
$$

Then for almost all $t \in\left(0, T^{*}\right), u^{l}(0, t)$ and $u^{r}(0, t)$ both exist.
Now we show that $u$ satisfies the integral equation in Definition 1.1. For any $\varphi(x, t)$ given as before, we have

$$
\begin{align*}
& \int_{\mathbb{R}} u_{0}(x) \varphi(x, 0) d x+\iint_{Q_{T^{*}}} u_{\varepsilon} \varphi_{t} d x d t-\iint_{Q_{T^{*}}}\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon} D \varphi d x d t \\
&+\iint_{Q_{T^{*}}} \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) \varphi(x, t) d x d t=0 \tag{2.16}
\end{align*}
$$

The second term of (2.16) can be rewritten as

$$
\begin{equation*}
\iint_{Q_{T^{*}}} u_{\varepsilon} \varphi_{t} d x d t=\int_{0}^{T^{*}} \int_{\mathbb{R} \backslash(-\tau, \tau)} u_{\varepsilon} \varphi_{t} d x d t+\int_{0}^{T^{*}} \int_{-\tau}^{\tau} u_{\varepsilon} \varphi_{t} d x d t \tag{2.17}
\end{equation*}
$$

where $\tau>0$. According to Lemma 2.5, we have

$$
\int_{0}^{T^{*}} \int_{-\tau}^{\tau} u_{\varepsilon} \varphi_{t} d x d t \leq \int_{0}^{T^{*}} \int_{-\tau}^{\tau} u_{\varepsilon}\left|\varphi_{t}\right| d x d t \leq \int_{0}^{T^{*}} \int_{-\tau}^{\tau} \lambda_{1}\left|\varphi_{t}\right| d x d t \leq C \tau \lambda_{1} T^{*}
$$

Thus letting $\varepsilon \rightarrow 0^{+}, \tau \rightarrow 0^{+}$in (2.17), we obtain

$$
\iint_{Q_{T^{*}}} u_{\varepsilon} \varphi_{t} d x d t \rightarrow \iint_{Q_{T^{*}}} u \varphi_{t} d x d t .
$$

Similarly, we can prove that

$$
\lim _{\varepsilon \rightarrow 0^{+}} \iint_{Q_{T^{*}}}\left|D u_{\varepsilon}\right|^{p-2} D u_{\varepsilon} D \varphi d x d t=\iint_{Q_{T^{*}}}|D u|^{p-2} D u D \varphi d x d t .
$$

In addition,

$$
\begin{aligned}
& \iint_{Q_{T^{*}}} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x d t \\
= & \int_{0}^{T^{*}} \int_{-\varepsilon}^{\varepsilon} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x d t \\
= & \int_{0}^{T^{*}} \int_{-\varepsilon}^{0} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x d t+\int_{0}^{T^{*}} \int_{0}^{\varepsilon} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x d t .
\end{aligned}
$$

Next we show the following equality holds

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0^{+}} \int_{-\varepsilon}^{0} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x=\frac{1}{2} f\left(u^{l}(0, t)\right) \varphi(0, t), \tag{2.18}
\end{equation*}
$$

for any $t \in\left(0, T^{*}\right)$. In fact,

$$
\begin{aligned}
& \left|\int_{-\varepsilon}^{0} \delta_{\varepsilon}(x) \varphi(x, t) f\left(u_{\varepsilon}(x, t)\right) d x-\frac{1}{2} f\left(u^{l}(0, t)\right) \varphi(0, t)\right| \\
\leq & \left|\int_{-\varepsilon}^{0} \delta_{\varepsilon}(x)\left(\varphi(x, t) f\left(u_{\varepsilon}(x, t)\right)-\varphi(0, t) f\left(u^{l}(0, t)\right)\right) d x\right| \\
& +\left|\left(\int_{-\varepsilon}^{0} \delta_{\varepsilon}(x) d x-\frac{1}{2}\right) \varphi(0, t) f\left(u^{l}(0, t)\right)\right| \\
\leq & \omega(\varepsilon) \int_{-\varepsilon}^{0} \delta_{\varepsilon}(x) d x+\varphi(0, t) f\left(u^{l}(0, t)\right)\left|\int_{-\varepsilon}^{0} \delta_{\varepsilon}(x) d x-\frac{1}{2}\right|,
\end{aligned}
$$

where $\omega(\varepsilon)=\sup _{-\varepsilon \leq x \leq 0}\left|\varphi(x, t) f\left(u_{\varepsilon}(x, t)\right)-\varphi(0, t) f\left(u^{l}(0, t)\right)\right|$, and

$$
\lim _{\varepsilon \rightarrow 0^{+}} \omega(\varepsilon)=0, \quad \lim _{\varepsilon \rightarrow 0^{+}} \int_{-\varepsilon}^{0} \delta_{\varepsilon}(x) d x=\frac{1}{2}
$$

we see that (2.18) holds. We can use the similar method to prove that

$$
\lim _{\varepsilon \rightarrow 0^{+}} \int_{0}^{\varepsilon} \delta_{\varepsilon}(x) \varphi(x, t) f\left(u_{\varepsilon}(x, t)\right) d x=\frac{1}{2} f\left(u^{r}(0, t)\right) \varphi(0, t) .
$$

Hence we obtain

$$
\begin{align*}
& \lim _{\varepsilon \rightarrow 0^{+}} \iint_{Q_{T^{*}}} \varphi(x, t) \delta_{\varepsilon}(x) f\left(u_{\varepsilon}(x, t)\right) d x d t \\
= & \int_{0}^{T^{*}} \frac{1}{2}\left(f\left(u^{l}(0, t)\right)+f\left(u^{r}(0, t)\right)\right) \varphi(0, t) d t \tag{2.19}
\end{align*}
$$

Thus letting $\varepsilon \rightarrow 0^{+}$in (2.16), we get

$$
\begin{aligned}
\int_{\mathbb{R}} u_{0}(x) \varphi(x, 0) d x & +\iint_{Q_{T^{*}}}\left(u \varphi_{t}-|D u|^{p-2} D u D \varphi\right) d x d t \\
& +\int_{0}^{T^{*}} \frac{1}{2}\left(f\left(u^{l}(0, t)\right)+f\left(u^{r}(0, t)\right)\right) \varphi(0, t) d t=0 .
\end{aligned}
$$

The proof of Theorem 1.1 is complete.

## Acknowledgments

This work is supported by the Fundamental Research Funds for the Central Universities and the Research Funds of Renmin University of China(2010030171).

## References

[1] Zhiyuan Pang, Yaodong Wang and Lishang Jiang, Semilinear diffusion equations with Dirac measure on the right side and the related optimal control problems, Acta Math. Sinica, 24(5)(1981), 780-796.
[2] L. Boccardo and T. Gallouët, Nonlinear elliptic and parabolic equations involving measure data, J. Func. Anal., 87(1)(1989), 149-169.
[3] L. Boccardo and T. Gallouët, Nonlinear elliptic equations with right hand side measures, Comm. Partial Differential Equations, 17(3/4)(1992), 641-655.
[4] Fengquan Li, Nonlinear degenerate parabolic equations with measure data, Nonlinear Analysis TMA, 61(7)(2005), 1269-1282.
[5] T. Del Vecchio, Nonlinear elliptic equations with measure data, Potential Anal., 4(2)(1995), 185-203.
[6] W. E. Olmstead and C. A. Roberts, Explosion in a diffusive strip due to a concentrated nonlinear source, Methods Appl. Anal., 1(4)(1994), 435-445.
[7] W. E. Olmstead and C. A. Roberts, Explosion in a diffusive strip due to a source with local and nonlocal features, Methods Appl. Anal., 3(3)(1996), 345-357.
[8] C. Y. Chan and H. Y. Tian, Single-point blow-up for a degenerate parabolic problem due to a concentrated nonlinear source, Quart. Appl. Math., 61(2)(2003), 363-385.
[9] C. Y. Chan and R. Boonklurb, A blow-up criterion for a degenerate parabolic problem due to a concentrated nonlinear source, Quart. Appl. Math., 65(4)(2007), 781-787.
[10] Jingxue Yin, Shuhe Wang and Yuanyuan Ke, A class of degenerate parabolic equations with a concentrated nonlinear source, Nonlinear Analysis TMA, 72(1)(2010), 123-131.
[11] Yinghua Li, Yuanyuan Ke and Zejia Wang, Cauchy problem for onedimensional $p$-Laplacian equation with point source, Journal of Partial Differential Equations, 18(1)(2005), 22-34.
[12] L. Boccardo, A. Dall'Aglio, T. Gallouët and L. Orsina, Nonlinear parabolic equations with measure data, Journal of Functional Analysis, 147(1)(1997), 237-258.
[13] Jing Wang, Zejia Wang and Jingxue Yin, A class of degenerate diffusion equations with mixed boundary conditions, J. Math. Anal. Appl., 298(2)(2004), 589-603.
[14] M. Berger, Nonlinearity and Functional Analysis, Academic Press, New York, 1977.
[15] E. DiBenedetto, Degenerate Parabolic Equations, Springer Verlag, New York, 1993.
[16] L. C. Evans, Weak Convergence Methods for Nonlinear Partial Differential Equations, CBMS Regional Conference Series in Mathematics, 74., Published for the Conference Board of the Mathematical Sciences, Washington, DC; by the American Mathematical Society, Providence, RI, 1990.

School of Information, Renmin University of China, Beijing 100872, P.R. China

School of Mathematical Sciences, South China Normal University, Guangzhou 510631, P.R. China
Email: yjx@scnu.edu.cn
Department of Mathematics, Jilin University,
Changchun 130012, P.R. China
Email: wangshuhe0821@yahoo.com.cn


[^0]:    *Corresponding author.
    Received by the editors October 2010 - In revised form in April 2011. Communicated by P. Godin.
    Key words and phrases : p-Laplacian, concentrated nonlinear source, local solution.

