

Blow-up for solution of a system of quasilinear hyperbolic equations involving the p -laplacian

Mamadou Sango

Abstract

We study the blow-up for the solution of a system of quasilinear hyperbolic equations involving the p -laplacian. We derive a differential inequality for a function involving some norms of the solution which yields the finite time blow-up.

1 Introduction

We are concerned with the blow-up of solutions of the initial boundary value problem for a class of quasilinear system of hyperbolic equations in a bounded domain $\Omega \subset \mathbf{R}^n$ ($n \geq 1$) with a sufficiently smooth boundary $\partial\Omega$:

$$u_{tt} - \sum_{i=1}^n \frac{\partial}{\partial x_i} \left(\left| \frac{\partial u}{\partial x} \right|^{p-2} \frac{\partial u}{\partial x_i} \right) - \Delta u_t + |v_t|^{\alpha_1} |u_t|^{\beta_1} \text{sign}(u_t) = |u|^{m_1-1} u \text{ in } \Omega \times (0, T), \quad (1)$$

$$v_{tt} - \sum_{i=1}^n \frac{\partial}{\partial x_i} \left(\left| \frac{\partial v}{\partial x} \right|^{p-2} \frac{\partial v}{\partial x_i} \right) - \Delta v_t + |v_t|^{\alpha_2} |u_t|^{\beta_2} \text{sign}(v_t) = |v|^{m_2-1} v \text{ in } \Omega \times (0, T), \quad (2)$$

$$u(x, t) = 0, v(x, t) = 0 \text{ on } \partial\Omega \times (0, T), \quad (3)$$

$$u(x, 0) = u_0(x), u_t(x, 0) = u_1(x) \text{ in } \Omega, \quad (4)$$

$$v(x, 0) = v_0(x), v_t(x, 0) = v_1(x) \text{ in } \Omega, \quad (5)$$

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where $0 < T \leq \infty$, p , m_i , α_i , β_i ($i = 1, 2$) are positive numbers subjected to some appropriate restrictions.

The problem is related to a class of nonlinear evolutions equations of the type

$$\begin{cases} \varphi_{tt} + A(t)\varphi - B(t)\varphi_t + D(t)\varphi_t = F(t)\varphi \\ \varphi(0) = \varphi_0 \\ \varphi_t(0) = \varphi_1 \end{cases} \quad (6)$$

where A , B , D and F are some nonlinear operators. Issues of global existence under various conditions were considered in [1], [14]; see the references in these papers. Equations of this type arise in several areas of physics. The most common of them being the case when

$$A\varphi = -\sum_{i=1}^n \frac{\partial}{\partial x_i} \vartheta_i \left(\frac{\partial \varphi}{\partial x} \right) \quad \text{and} \quad B\varphi_t = \Delta \varphi_t$$

which describes the longitudinal motion of a viscoelastic bar obeying the nonlinear Voight model. Physically the strong damping term $-\Delta \varphi_t$ and the nonlinear dissipative damping term $D(t)\varphi_t$ play a dissipative role in the energy accumulation in the configurations of viscoelastic materials, while the nonlinear source term $F(t)\varphi$ leads to the gathering of energy in the configurations. The interaction between these terms may lead to a lack of synchronization in the energy accumulation and as a result the configuration may break or burn out in finite time, this mathematically is expressed through the finite time blow-up of the solution.

Here we consider an initial boundary value problem involving a system of nonlinear hyperbolic equations with slightly more general nonlinear damping terms.

The case without sources terms which lead to global existence was considered in [1]. The study of finite time blow-up involving one equation (thus one of the parameters $\alpha_i = 0$ and $m_1 = m_2$) was considered in [14], [15].

The approach in the present paper follows closely that of [2], [7], [9]. We refer also to the important papers devoted to related questions such as [4], [3], [5], [6], [8], [10], [11], [12] (this paper treats hyperbolic systems with source terms without damping) and in the several references therein; the approach in some of these papers is mainly based on the potential well method which originated in the work of Sattinger [13]. We note that semilinear equations and systems (when $p = 2$) are the ones that have been widely studied. Nonlinear hyperbolic problems involving the p -Laplacian are becoming the object of increasing interest only in recent years.

The paper is organized as follows. In section 2, we state our main result. Section 3 is devoted to the proof of the main result through the derivation of a suitable differential inequality satisfied by a function involving some norms of the solution.

2 Preliminaries

We introduce some notations. By $L_p(\Omega)$ ($p \geq 1$) we denote the set of integrable functions u on Ω , such that the norm

$$\|u\|_{L_p(\Omega)} = \left(\int_{\Omega} |u|^p dx \right)^{1/p} < \infty.$$

Let X be a Banach space. By the symbol $L_p(0, T, Y)$ we mean the functions $u(x, t)$ that are L_p -integrable from $[0, T]$ into X and with the norm

$$\|u\|_{L_p(0,T,X)} = \left(\int_0^T \|u(t)\|_X^p dt \right)^{1/p}, \quad 1 \leq p < \infty$$

and

$$\|u\|_{L_\infty(0,T,X)} = \text{ess sup}_{t \in [0,T]} \|u(t)\|_X.$$

For $p > 1$, we consider the function space

$$\overset{\circ}{W}_p^1(\Omega) = \left\{ u \in L_p(\Omega) : u|_{\partial\Omega} = 0, \frac{\partial u}{\partial x_i} \in L_p(\Omega), i = 1, \dots, n \right\},$$

with the norm

$$\|u\|_{\overset{\circ}{W}_p^1(\Omega)} = \left(\sum_{i=1}^n \int_\Omega \left| \frac{\partial u}{\partial x_i} \right|^p dx \right)^{1/p};$$

when $p = 2$, we denote $\overset{\circ}{W}_p^1(\Omega)$ by $H_0^1(\Omega)$. We denote by X^2 the Cartesian product of a set X with itself. The letter C will stand for all constants depending only on the data.

We introduce the functions

$$H_u(t) = -\frac{1}{2} \int_\Omega u_t^2 dx - \frac{1}{p} \int_\Omega \left| \frac{\partial u}{\partial x} \right|^p dx + \frac{1}{m_1 + 1} \int_\Omega |u|^{m_1+1} dx, \tag{7}$$

$$H_v(t) = -\frac{1}{2} \int_\Omega v_t^2 dx - \frac{1}{p} \int_\Omega \left| \frac{\partial v}{\partial x} \right|^p dx + \frac{1}{m_2 + 1} \int_\Omega |v|^{m_2+1} dx, \tag{8}$$

$$H(t) = H_u(t) + H_v(t) \tag{9}$$

$$F(t) = \|u(t)\|_{L_2(\Omega)}^2 + \|v(t)\|_{L_2(\Omega)}^2 + \int_0^t \int_\Omega \left[\left| \frac{\partial u(\tau)}{\partial x} \right|^2 + \left| \frac{\partial v(\tau)}{\partial x} \right|^2 \right] dx d\tau; \tag{10}$$

for sake of simplicity we denote from now on the norm $\|\cdot\|_{L_p(\Omega)}$ by $\|\cdot\|_p$.

Our main result is

Theorem 1. *Let $U = (u, v)$ be a local weak solution of problem (1)-(5), in the sense that there exists a number $0 < T < \infty$ such that*

$$U \in \left[C \left(0, T, \overset{\circ}{W}_p^1(\Omega) \right) \right]^2 \cap C(0, T, L_{m_1+1}(\Omega)) \times C(0, T, L_{m_2+1}(\Omega)), \tag{11}$$

$$U_t = (u_t, v_t) \in [C(0, T, L_2(\Omega))]^2 \cap [L_2(0, T, H_0^1(\Omega))]^2 \tag{12}$$

$$\int_0^T \int_\Omega [|v_t|^{\alpha_1} |u_t|^{\beta_1+1} + |v_t|^{\alpha_2+1} |u_t|^{\beta_2}] dx dt \text{ is finite} \tag{13}$$

$$U_0 = (u_0, v_0) \in \left[\overset{\circ}{W}_p^1(\Omega) \right]^2, U_1 = (u_1, v_1) \in [L_2(\Omega)]^2 \tag{14}$$

and u satisfies (1)-(5) in the weak sense. Furthermore we assume that

$$H_u(0) \geq C_1 > 0, H_v(0) \geq C_2 > 0 F'(0) = \int_{\Omega} [u_0 u_1 + v_0 v_1] dx > 0 \tag{15}$$

$$\alpha_i > 0, \beta_i > 0, m_i > 1 (i = 1, 2), 0 < \alpha_1 - \alpha_2 < 1, 0 < \beta_2 - \beta_1 < 1 \tag{16}$$

$$\frac{\alpha_2 + 1}{\alpha_1} = \frac{\beta_2}{\beta_1} > \frac{m_1 + 1}{m_1}, \frac{\alpha_1}{\alpha_2} = \frac{\beta_1 + 1}{\beta_2} > \frac{m_2 + 1}{m_2}, \tag{17}$$

$$\max_i \left\{ 2, \frac{n(m_i + 1)}{n + m_i + 1} \right\} \leq p < \min_i \{m_i + 1, n\}. \tag{18}$$

Then u blows up in finite time, i.e., there exists a $T_0 > 0$ such that

$$\lim_{t \rightarrow T_0^-} \left[\|u(t)\|_{m_1+1}^{m_1+1} + \|v(t)\|_{m_2+1}^{m_2+1} + \|U_t(t)\|_2^2 \right] = \infty.$$

Remark 2. The constants C_1 and C_2 will be chosen later. Some few words about questions related to the problem (1)-(5) are in order. The global existence without the source terms was considered in [1]. In particular it was shown that if $p \geq 2$, $0 < \beta_1 < 1 - \alpha_1$, $0 < \alpha_2 < 1 - \beta_2$, $0 < \alpha_1 < 1$, $0 < \beta_2 < 1$ and the above conditions (14) are imposed on the initial data, then a global weak solution exists and decay estimates under further conditions were derived.

3 Proof of the theorem

The blow-up result will follow from a differential inequality satisfied by the function

$$W(t) = H^{1-\alpha}(t) + \varepsilon F'(t),$$

where α and ε are small parameters that will be chosen in the sequel. This idea goes back to Ball [2].

We start with the derivation of some useful informations on the function H which follow from a suitable identity.

Multiplying the equations (1) and (2) by u_t and v_t respectively and integrating over Ω , we get

$$\frac{d}{dt} \left[\frac{1}{2} \|u_t\|_2^2 + \frac{1}{p} \left\| \frac{\partial u}{\partial x} \right\|_p^p - \frac{1}{m_1 + 1} \|u\|_{m_1+1}^{m_1+1} \right] = - \left\| \frac{\partial u_t}{\partial x} \right\|_2^2 - \int_{\Omega} |v_t|^{\alpha_1} |u_t|^{\beta_1+1} dx$$

$$\frac{d}{dt} \left[\frac{1}{2} \|v_t\|_2^2 + \frac{1}{p} \left\| \frac{\partial v}{\partial x} \right\|_p^p - \frac{1}{m_2 + 1} \|v\|_{m_2+1}^{m_2+1} \right] = - \left\| \frac{\partial v_t}{\partial x} \right\|_2^2 - \int_{\Omega} |v_t|^{\alpha_2+1} |u_t|^{\beta_2} dx$$

This implies that

$$H'_u = \left\| \frac{\partial u_t}{\partial x} \right\|_2^2 + \int_{\Omega} |v_t|^{\alpha_1} |u_t|^{\beta_1+1} dx, \tag{19}$$

and

$$H'_v = \left\| \frac{\partial v_t}{\partial x} \right\|_2^2 + \int_{\Omega} |v_t|^{\alpha_2+1} |u_t|^{\beta_2} dx, \tag{20}$$

that is H_u and H_v are non decreasing. Thus from the assumption on $H_u(0)$ and $H_v(0)$, it follows that for all $t \geq 0$,

$$0 < H_u(0) \leq H_u(t) \leq \frac{1}{m_1 + 1} \|u\|_{m_1+1}^{m_1+1} \leq C \left\| \frac{\partial u}{\partial x} \right\|_p^{m_1+1}, \tag{21}$$

$$0 < H_v(0) \leq H_v(t) \leq \frac{1}{m_2 + 1} \|u\|_{m_2+1}^{m_2+1} \leq C \left\| \frac{\partial v}{\partial x} \right\|_p^{m_2+1}, \tag{22}$$

where we have used the fact that $W_p^1(\Omega)$ is embedded into $L_{m_i+1}(\Omega)$ since by the conditions (18) on p , $pn/(n-p) \geq \max_i \{m_i + 1\}$.

By approximating u with sufficiently smooth functions with respect to t , we can see that F'' satisfies

$$\begin{aligned} F''(t) &= 2 \int_{\Omega} (u_t^2 + v_t^2) dx + 2 \int_{\Omega} (uu_{tt} + vv_{tt}) dx \\ &\quad + 2 \int_{\Omega} \sum_{i=1}^n \left[\frac{\partial u}{\partial x_i} \frac{\partial u_t}{\partial x_i} + \frac{\partial v}{\partial x_i} \frac{\partial v_t}{\partial x_i} \right] dx \\ &= 2 \left(\|u_t\|_2^2 - \left\| \frac{\partial u}{\partial x} \right\|_p^p - \int_{\Omega} |v_t|^{\alpha_1} |u_t|^{\beta_1} u \operatorname{sign} u_t dx + \|u\|_{m_1+1}^{m_1+1} \right) \\ &\quad + 2 \left(\|v_t\|_2^2 - \left\| \frac{\partial v}{\partial x} \right\|_p^p - \int_{\Omega} |v_t|^{\alpha_2} |u_t|^{\beta_2} v \operatorname{sign} v_t dx + \|v\|_{m_2+1}^{m_2+1} \right). \end{aligned} \tag{23}$$

Thus substituting F'' as expressed in (23) into the relation

$$\frac{d}{dt} (H^{1-\alpha}(t) + \varepsilon F'(t)) = (1 - \alpha) H^{-\alpha} H + \varepsilon F''(t),$$

and using the definition of H , we get

$$\begin{aligned} &\frac{d}{dt} (H^{1-\alpha}(t) + \varepsilon F'(t)) \\ &= (1 - \alpha) H^{-\alpha} H' + \varepsilon (2 + p) \|U_t\|_2^2 + 2p\varepsilon H \\ &\quad + \varepsilon \left(2 - \frac{2p}{m_1 + 1} \right) \|u\|_{m_1+1}^{m_1+1} + \varepsilon \left(2 - \frac{2p}{m_2 + 1} \right) \|v\|_{m_2+1}^{m_2+1} \\ &\quad - 2\varepsilon \int_{\Omega} |v_t|^{\alpha_1} |u_t|^{\beta_1} u \operatorname{sign} u_t dx - 2\varepsilon \int_{\Omega} |v_t|^{\alpha_2} |u_t|^{\beta_2} v \operatorname{sign} v_t dx. \end{aligned} \tag{24}$$

Let us denote the two last integrals by I_1 and I_2 , respectively and estimate them.

By Hölder’s inequality, we have

$$\begin{aligned} I_1 &\leq \int_{\Omega} |v_t|^{\alpha_1} |u_t|^{\beta_1} |u| \, dx \\ &\leq \left(\int_{\Omega} (|v_t|^{\alpha_1} |u_t|^{\beta_1})^{(m_1+1)/m_1} \, dx \right)^{m_1/(m_1+1)} \left(\int_{\Omega} |u|^{m_1+1} \, dx \right)^{1/(m_1+1)} \\ &\leq \left(\int_{\Omega} |v_t|^{\alpha_2+1} |u_t|^{\beta_2} \, dx \right)^{\beta_1/\beta_2} \left(\int_{\Omega} |u|^{m_1+1} \, dx \right)^{1/(m_1+1)} \end{aligned}$$

where in the last inequality we used the restrictions (16) and (17).

Next writing

$$\frac{1}{m_1 + 1} = \frac{\beta_2 - \beta_1}{\beta_2} + \sigma_1; \quad \sigma_1 = \frac{1}{m_1 + 1} - \frac{\beta_2 - \beta_1}{\beta_2},$$

we have $\sigma_1 < 0$, and using Young’s inequality we get

$$\begin{aligned} I_1 &\leq \left(\int_{\Omega} |v_t|^{\alpha_2+1} |u_t|^{\beta_2} \, dx \right)^{\beta_1/\beta_2} \left(\int_{\Omega} |u|^{m_1+1} \, dx \right)^{\sigma_1} \left(\int_{\Omega} |u|^{m_1+1} \, dx \right)^{(\beta_2-\beta_1)/\beta_2} \\ &\leq C \left[\int_{\Omega} |v_t|^{\alpha_2+1} |u_t|^{\beta_2} \, dx + \int_{\Omega} |u|^{m_1+1} \, dx \right] \left(\int_{\Omega} |u|^{m_1+1} \, dx \right)^{\sigma_1}. \end{aligned}$$

By (21) and (20), we get

$$2\varepsilon I_1 \leq C\varepsilon (H_u(t))^{\sigma_1} \left[H'_v(t) + \int_{\Omega} |u|^{m_1+1} \, dx \right]. \tag{25}$$

Analogously we obtain

$$2\varepsilon I_2 \leq C\varepsilon (H_v(t))^{\sigma_2} \left[H'_u(t) + \int_{\Omega} |v|^{m_2+1} \, dx \right]; \quad \sigma_2 = \frac{1}{m_2 + 1} - \frac{\alpha_1 - \alpha_2}{\alpha_1}. \tag{26}$$

Combining these two last inequalities with (21)-(22), it follows that

$$2\varepsilon (I_1 + I_2) \leq 2\varepsilon [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] \left[H'(t) + \|u\|_{m_1+1}^{m_1+1} + \|v\|_{m_2+1}^{m_2+1} \right].$$

Taking account of (24), the following relation holds:

$$\begin{aligned} &\frac{d}{dt} \left(H^{1-\alpha}(t) + \varepsilon F'(t) \right) \\ &\geq \left\{ (1 - \alpha) H^{-\alpha}(0) - 2\varepsilon C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] \right\} H'(t) + \varepsilon (2 + p) \|U_t\|_2^2 + 2p\varepsilon H \\ &\quad + \varepsilon \left(2 - \frac{2p}{m_1 + 1} - C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] \right) \|u\|_{m_1+1}^{m_1+1} \\ &\quad + \varepsilon \left(2 - \frac{2p}{m_2 + 1} - C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] \right) \|v\|_{m_2+1}^{m_2+1}. \end{aligned} \tag{27}$$

We take

$$\alpha \in \left(0, \min \left\{ \frac{m_1 - 1}{2(m_1 + 1)}, \frac{m_2 - 1}{2(m_2 + 1)}, \frac{p - 2}{p} \right\} \right), \tag{28}$$

and choose $\varepsilon > 0$ such that

$$(1 - \alpha) H^{-\alpha}(0) - 2\varepsilon C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] > 0.$$

Also we choose the constants C_1 and C_2 in (15) in such a way that

$$C_1^{\sigma_1} + C_2^{\sigma_2} \leq \min_i \left\{ \frac{1}{C} \left(2 - \frac{2p}{m_i + 1} \right) \right\}.$$

Then the inequalities hold:

$$2 - \frac{2p}{m_1 + 1} - C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] > 0$$

$$2 - \frac{2p}{m_2 + 1} - C [H_u^{\sigma_1}(0) + H_v^{\sigma_2}(0)] > 0.$$

Thus from (19), (20) it follows that

$$\frac{d}{dt} (H^{1-\alpha}(t) + \varepsilon F'(t)) \geq C (\|U_t\|_2^2 + H(t) + \|u\|_{m_1+1}^{m_1+1} + \|u\|_{m_2+1}^{m_2+1}); \tag{29}$$

as a consequence we have that $W(t)$ is increasing since $H(t) > 0$ by (21)-(22). Therefore using the assumption that $F'(0) > 0$, we get

$$W(t) > 0, \forall t \geq 0.$$

We make a further restriction on α by requiring that $0 < \alpha < 1/2$. Then setting $\beta = 1/(1 - \alpha)$ (i.e., $2 > \beta > 1$) we claim the inequality

$$W'(t) \geq CW^\beta(t). \tag{30}$$

For the proof of (30), we consider two alternatives:

- If there exists a $t > 0$ such that $F'(t) < 0$, then

$$(H^{1-\alpha}(t) + \varepsilon F'(t))^\beta \leq H(t). \tag{31}$$

Thus (30) follows from (29).

- If there exists a $t > 0$ such that $F'(t) \geq 0$, then using Holder's and Young's inequalities we get

$$[F'(t)]^\beta = \left(2 \int_\Omega [u_t u + v_t v] dx + \left\| \frac{\partial u}{\partial x} \right\|_2^2 + \left\| \frac{\partial v}{\partial x} \right\|_2^2 \right)^\beta$$

$$\leq C \left[\|u\|_2^{\lambda_1 \beta} + \|u_t\|_2^{\mu_1 \beta} + \|v\|_2^{\lambda_2 \beta} + \|v_t\|_2^{\mu_2 \beta} + \left\| \frac{\partial u}{\partial x} \right\|_2^{2\beta} + \left\| \frac{\partial v}{\partial x} \right\|_2^{2\beta} \right] \tag{32}$$

where $\lambda_i^{-1} + \mu_i^{-1} = 1$. We take $\mu_i \beta = 2, i = 1, 2$. Thus $\mu_1 = \mu_2 = 2/\beta$ and

$$\lambda_1 = \lambda_2 = \lambda = \frac{2(1 - \alpha)}{1 - 2\alpha}.$$

By the restrictions on α , we have

$$\lambda\beta = \frac{2}{1 - 2\alpha} \leq \min \{m_i + 1\}, \quad 2\beta = \frac{2}{1 - \alpha} \leq p.$$

Thus from (32), using Hölder’s inequality we have

$$\begin{aligned}
 [F'(t)]^\beta &\leq C \left[\|u\|_{m_1+1}^{\lambda\beta} + \|u_t\|_2^2 + \|v\|_{m_2+1}^{\lambda\beta} + \|v_t\|_2^2 + \left\| \frac{\partial u}{\partial x} \right\|_p^{2\beta} + \left\| \frac{\partial v}{\partial x} \right\|_p^{2\beta} \right] \\
 &= C \left\{ \|u\|_{m_1+1}^{\lambda\beta-(m_1+1)} \|u\|_{m_1+1}^{m_1+1} + \|U_t\|_2^2 + \|v\|_{m_2+1}^{\lambda\beta-(m_2+1)} \|v\|_{m_2+1}^{m_2+1} \right. \\
 &\quad \left. + \left\| \frac{\partial u}{\partial x} \right\|_p^{2\beta-p} \left\| \frac{\partial u}{\partial x} \right\|_p^p + \left\| \frac{\partial v}{\partial x} \right\|_p^{2\beta-p} \left\| \frac{\partial v}{\partial x} \right\|_p^p \right\}.
 \end{aligned}$$

From the estimates (21) and (22) we deduce that

$$\begin{aligned}
 [F'(t)]^\beta &\leq C \left\{ [H_u(0)]^{[\lambda\beta-(m_1+1)]/[m_1+1]} \|u\|_{m_1+1}^{m_1+1} \right. \\
 &\quad \left. + \|U_t\|_2^2 + [H_v(0)]^{[\lambda\beta-(m_2+1)]/[m_2+1]} \|v\|_{m_2+1}^{m_2+1} \right. \\
 &\quad \left. + [H_u(0)]^{(2\beta-p)/(m_1+1)} \left\| \frac{\partial u}{\partial x} \right\|_p^p + [H_v(0)]^{(2\beta-p)/(m_2+1)} \left\| \frac{\partial v}{\partial x} \right\|_p^p \right\}.
 \end{aligned}$$

Thus

$$[F'(t)]^\beta \leq C \left[\|u\|_{m_1+1}^{m_1+1} + \|v\|_{m_2+1}^{m_2+1} + \|U_t\|_2^2 + \left\| \frac{\partial u}{\partial x} \right\|_p^p + \left\| \frac{\partial v}{\partial x} \right\|_p^p \right].$$

From the definition of H we have

$$H(t) + \frac{1}{p} \left[\left\| \frac{\partial u}{\partial x} \right\|_p^p + \left\| \frac{\partial v}{\partial x} \right\|_p^p \right] \leq \frac{1}{m_1+1} \|u\|_{m_1+1}^{m_1+1} + \frac{1}{m_2+1} \|v\|_{m_2+1}^{m_2+1}.$$

Thus

$$\begin{aligned}
 [F'(t)]^\beta &\leq C \left[\|u\|_{m_1+1}^{m_1+1} + \|v\|_{m_2+1}^{m_2+1} + \|U_t\|_2^2 + \left\| \frac{\partial u}{\partial x} \right\|_p^p + \left\| \frac{\partial v}{\partial x} \right\|_p^p + H(t) \right] \\
 &\leq C \left[\|u\|_{m_1+1}^{m_1+1} + \|v\|_{m_2+1}^{m_2+1} + \|U_t\|_2^2 \right],
 \end{aligned}$$

and hence

$$(H^{1-\alpha}(t) + \varepsilon F'(t))^\beta \leq C \left[\|u\|_{m_1+1}^{m_1+1} + \|v\|_{m_2+1}^{m_2+1} + \|U_t\|_2^2 \right]. \tag{33}$$

This inequality together with (29) imply (30).

Now integrating both sides of (30) over the interval $[0, t]$, it follows that there exists a $T_0 > 0$ such that

$$\lim_{t \rightarrow T_0^-} (H^{1-\alpha}(t) + \varepsilon F'(t)) = \infty.$$

This limit combined with (33), (31), (21) and (22) give

$$\lim_{t \rightarrow T_0^-} \left[\|u(t)\|_{m_1+1}^{m_1+1} + \|v(t)\|_{m_2+1}^{m_2+1} + \|U_t(t)\|_2^2 \right] = \infty.$$

The theorem is proved.

References

- [1] Biazutti, A. C. On a nonlinear evolution equation and its applications. *Nonlinear Anal.* 24 (1995), no. 8, 1221–1234.
- [2] Ball, J. M. Remarks on blow-up and nonexistence theorems for nonlinear evolution equations. *Quart. J. Math. Oxford Ser. (2)* 28 (1977), no. 112, 473–486.
- [3] Eloulaimi, R., Guedda, M. Nonexistence of global solutions of nonlinear wave equations. *Port. Math. (N.S.)* 58 (2001), no. 4, 449–460.
- [4] Galaktionov, V. A.; Pohozaev, S. I. Blow-up and critical exponents for nonlinear hyperbolic equations. *Nonlinear Anal.* 53 (2003), no. 3-4, 453–466.
- [5] Guedda, M., Labani, H. Nonexistence of global solutions to a class of nonlinear wave equations with dynamic boundary conditions. *Bull. Belg. Math. Soc. Simon Stevin* 9 (2002), no. 1, 39–46.
- [6] Glassey, R. T. Blow-up theorems for nonlinear wave equations. *Math. Z.* 132 (1973), 183–203.
- [7] Ikehata, R., Some remarks on the wave equations with nonlinear damping and source terms. *Nonlinear Anal.* 27 (1996), no. 10, 1165–1175.
- [8] Kirane, M., Messaoudi, S. Nonexistence results for the Cauchy problem of some systems of hyperbolic equations. *Ann. Polon. Math.* 78 (2002), no. 1, 39–47
- [9] Lions, J.-L. Contrôle des systèmes distribués singuliers. (French) [Control of singular distributed systems] *Méthodes Mathématiques de l'Informatique [Mathematical Methods of Information Science]*, 13. Gauthier-Villars, Montrouge, 1983.
- [10] Ono, K. On global solutions and blow-up solutions of nonlinear Kirchhoff strings with nonlinear dissipation. *J. Math. Anal. Appl.* 216 (1997), no. 1, 321–342.
- [11] Ono, K. Global existence, decay, and blowup of solutions for some mildly degenerate nonlinear Kirchhoff strings. *J. Differential Equations* 137 (1997), no. 2, 273–301.
- [12] Pohozaev, S.; Véron, L. Blow-up results for nonlinear hyperbolic inequalities. *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)* 29 (2000), no. 2, 393–420.
- [13] Sattinger, D. H. On global solution of nonlinear hyperbolic equations. *Arch. Rational Mech. Anal.* 30 1968 148–172.
- [14] Yang, Zhijian. Blowup of solutions for a class of non-linear evolution equations with non-linear damping and source terms. *Math. Methods Appl. Sci.* 25 (2002), no. 10, 825–833.

- [15] Yang, Zhijian. Existence and asymptotic behaviour of solutions for a class of quasi-linear evolution equations with non-linear damping and source terms. *Math. Methods Appl. Sci.* 25 (2002), no. 10, 795–814.

Department of Mathematics
University of Pretoria, Mamelodi Campus
Pretoria 0002, South Africa
Email: sango7777@yahoo.com, mamadou.sango@up.ac.za