## ON THE DETERMINATION OF A HILL'S EQUATION FROM ITS SPECTRUM

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A Hill's equation is an equation of the form:

(1) 
$$y'' + [\lambda - q(z)]y = 0, \quad q(z + \pi) = q(z),$$

where q(z) is assumed to be integrable over  $[0, \pi]$ . Without loss of generality, it is customary to assume that  $\int_0^{\pi} q(z) dz = 0$ . The discriminant of (1) is defined by

$$\Delta(\lambda) = y_1(\pi) + y'_2(\pi),$$

where  $y_1$  and  $y_2$  are solutions of (1) satisfying  $y_1(0)=y'_2(0)=1$  and  $y'_1(0)=y_2(0)=0$ .

The set of values of  $\lambda$  for which  $|\Delta| > 2$  consists of a finite or an infinite number of finite disjoint intervals and one infinite interval. These intervals are called instability intervals, since (1) has no solution which is bounded for all real z in these intervals. When  $|\Delta| < 2$ , all solutions of (1) are bounded for all real z and the corresponding intervals are called stability intervals. Pertinent information about stability and instability intervals of (1) can be found in Magnus and Winkler [1].

The following result has been proved:

THEOREM. If q(z) is real and integrable, and if precisely n finite instability intervals fail to vanish, then q(z) must satisfy a differential equation of the form

(2) 
$$q^{(2n)} + H(q, q', \cdots, q^{(2n-2)}) = 0, \quad a.e$$

where H is a polynomial of maximal degree n+2.

Borg [2], Hochstadt [3] and Ungar [4] proved this theorem for the case n=0, i.e. when all finite instability intervals vanish, and found that

(3) 
$$q(z) = 0$$
, a.e.

For the case n=1, Hochstadt [3] showed that q(z) is the elliptic function which satisfies

(4) 
$$q'' = 3q^2 + Aq + B$$
, a.e.

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where A and B are constants. (3) and (4) are equivalent to (2) for the cases n=0 and 1, respectively. In particular, for the case n=2, the explicit expression of (2) is

(5) 
$$q^{(4)} = 10qq'' + Aq'' + 5(q')^2 - 10q^3 + Bq^2 + Cq + D$$
, a.e.

where A, B, C and D are constants.

Erdélyi [5] investigated a Hill's equation where q(z) is a Lamé function and discovered situations where all but a finite number of finite instability intervals vanish. (4) provides a converse to some of his results.

Lax [6], through the study of partial differential operators, derived sufficient conditions for the vanishing of all but n finite instability intervals. These conditions coincide with (3), (4) and (5) for the cases n=0, 1 and 2, respectively. Whether there exist equivalent necessary conditions for higher values of n is still an open question.

The proof of the Theorem is accomplished by investigating the related problem

(6) 
$$u'' + [\lambda - q(z + \tau)]u = 0; \quad \tau \text{ real, arbitrary}$$

and by assuming the result [3] that q(z) is infinitely differentiable a.e. when *n* finite instability intervals fail to vanish.

In [3], Hochstadt showed that (6), when subject to  $u(0)=u(\pi)$ , has eigenvalues  $\mu_n(\tau)$ , where  $\mu_i(\tau)$  lies in the *i*th finite instability interval of (1). Furthermore, when precisely *n* finite instability intervals fail to vanish

$$u_{2}(\pi)\prod_{i=1}^{n} [\lambda - \mu_{i}(0)] = y_{2}(\pi)\prod_{i=1}^{n} [\lambda - \mu_{i}(\tau)],$$

where  $u_2(t)$  denotes the solution of (6) which satisfies  $u_2(0)=0$  and  $u'_2(0)=1$ . Suitable asymptotic expressions of  $u_2(\pi)$  and  $y_2(\pi)$  have been developed and our result follows upon their substitution into this equation. The details will appear in a later paper.

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