Critical scaling of stochastic epidemic models^{*}

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To Piet Groeneboom, on the occasion of his 39th birthday.

Abstract: In the simple mean-field *SIS* and *SIR* epidemic models, infection is transmitted from infectious to susceptible members of a finite population by independent p-coin tosses. Spatial variants of these models are proposed, in which finite populations of size N are situated at the sites of a lattice and infectious contacts are limited to individuals at neighboring sites. Scaling laws for both the mean-field and spatial models are given when the infection parameter p is such that the epidemics are *critical*. It is shown that in all cases there is a *critical threshold* for the numbers initially infected: below the threshold, the epidemic evolves in essentially the same manner as its *branching envelope*, but at the threshold evolves like a branching process with a sizedependent drift.

1. Stochastic epidemic models

1.1. Mean-field models

The simplest and most thoroughly studied stochastic models of epidemics are meanfield models, in which all individuals of a finite population interact in the same manner. In these models, a contagious disease is transmitted among individuals of a homogeneous population of size N. In the simple SIS epidemic, individuals are at any time either infected or susceptible; infected individuals remain infected for one unit of time and then become susceptible. In the simple SIR epidemic (more commonly known as the Reed-Frost model), individuals are either infected, susceptible, or recovered; infected individuals remain infected for one unit of time, after which they recover and acquire permanent immunity from future infection. In both models, the mechanism by which infection occurs is random: At each time, for any pair (i, s) of an infected and a susceptible individual, the disease is transmitted from i to s with probability $p = p_N$. These transmission events are mutually independent. Thus, in both the SIR and the SIS model, the number $J_{t+1} = J_t^N$ of infected individuals at time t + 1 is given by

(1)
$$J_{t+1} = \sum_{s=1}^{S_t} \xi_s.$$

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where $S_t = S_t^N$ is the number of susceptibles at time t and the random variables ξ_s are, conditional on the history of the epidemic to time t, independent, identically distributed Bernoulli-1 – $(1 - p)^{J_t}$. In the SIR model,

(2)
$$R_{t+1} = R_t + J_t \quad \text{and} \quad S_{t+1} = S_t - J_{t+1},$$

where R_t is the number of recovered individuals at time t, while in the SIS model,

(3)
$$S_{t+1} = S_t + J_t - J_{t+1}.$$

In either model, the epidemic ends at the first time T when $J_T = 0$. The most basic and interesting questions concerning these models have to do with the *duration* Tand size $\sum_{t \leq T} J_t$ of the epidemic and their dependence on the infection parameter p_N and the initial conditions.

1.2. Spatial SIR and SIS epidemics

In the simple SIS and SIR epidemics, no allowance is made for geographic or social stratifications of the population, nor for variability in susceptibility or degree of contagiousness. Following are descriptions of simple stochastic models that incorporate a geographic stratification of a population. We shall call these the (spatial) SIS-d and SIR-d epidemics, with d denoting the spatial dimension.

Assume that at each lattice point $x \in \mathbb{Z}^d$ is a homogeneous population of N_x individuals, each of whom may at any time be either susceptible or infected, or (in the *SIR* variants) recovered. These populations may be thought of as "villages". As in the mean-field models, infected individuals remain contagious for one unit of time, after which they recover with immunity from future infection (in the *SIR* variants) or once again become susceptible (in the *SIS* models). At each time $t = 0, 1, 2, \ldots$, for each pair (i_x, s_y) of an infected individual located at x and a susceptible individual at y, the disease spreads from i_x to s_y with probability $\alpha(x, y)$.

The simple Reed-Frost and stochastic logistic epidemics described in section 1.1 terminate with probability one, regardless of the value of the infection parameter p, because the population is finite. For the spatial *SIS* and *SIR* models this is no longer necessarily the case: If $\sum_{x \in \mathbb{Z}^d} N_x = \infty$ then, depending on the value of the parameter p and the dimension d, the epidemic may persist forever with positive probability. (For instance, if $N_x = 1$ for all x and $\alpha(x, y) = p$ for nearest neighbor pairs x, y but $\alpha(x, y) = 0$ otherwise, then the *SIS* -d epidemic is just oriented percolation on \mathbb{Z}^{d+1} , which is known to survive with positive probability if p exceeds a critical value $p_c < 1$ [6].) Obvious questions of interest center on how the epidemic spreads through space, and in cases where it eventually dies out, how far it spreads.

The figure below shows a simulation of an SIS-1 epidemic with village size $N_x = 20224$ and infection parameter 1/60672. At time 0 there were 2048 infected individuals at site 0; all other individuals were healthy. The epidemic lasted 713 generations (only the first 450 are shown).

1.3. Epidemic models and random graphs

All of the models described above have equivalent descriptions as structured random graphs, that is, percolation processes. Consider for definiteness the simple SIR

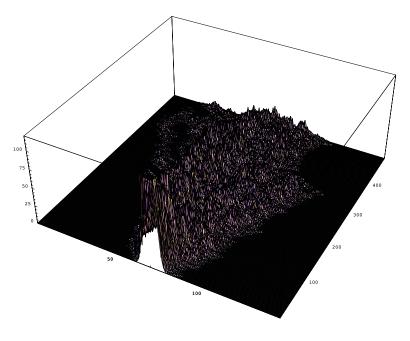


Fig 1.

(Reed-Frost) epidemic. In this model, no individual may be infected more than once; furthermore, for any pair x, y of individuals, there will be at most one opportunity for infection to pass from x to y or from y to x during the course of the epidemic. Thus, one could simulate the epidemic by first tossing a p-coin for every pair x, y, drawing an edge between x and y for each coin toss resulting in a Head, and then using the resulting (Erdös-Renyi) random graph determined by these edges to determine the course of infection in the epidemic. In detail: If Y_0 is the set of infected individuals at time 0, then the set Y_1 of individuals infected at time 1 consists of all $x \notin Y_0$ that are connected to individuals in Y_0 , and for any subsequent time n, the set Y_{n+1} of individuals infected at time n+1 consists of all $x \notin \bigcup_{j \leq n} Y_j$ who are connected to individuals in Y_n . Note that the set of individuals ultimately infected during the course of the epidemic is the union of those connected components of the random graph containing at least one vertex in Y_0 .

Similar random graph descriptions may be given for the simple *SIS* and the spatial *SIS* and *SIR* epidemic models.

1.4. Branching envelopes of epidemics

For each of the stochastic epidemic models discussed above there is an associated branching process that serves, in a certain sense, as a "tangent" to the epidemic. We shall refer to this branching process as the *branching envelope* of the epidemic. The branching envelopes of the simple mean-field epidemics are ordinary Galton-Watson processes; the envelopes of the spatial epidemics are branching random walks. There is a natural coupling of each epidemic with its branching envelope in which the set of infected individuals in the epidemic is at each time (and in the spatial models, at each location) dominated by the corresponding set of individuals in the branching envelope.

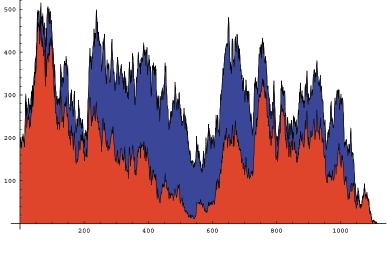


Fig 2.

Following is a detailed description of the natural coupling of the simple SIS epidemic with its branching envelope. The branching envelope is a Galton-Watson process Z_n with offspring distribution Binomial-(N, p), where p is the infection parameter of the epidemic, and whose initial generation Z_0 coincides with the set of individuals who are infected at time 0. Particles in the Galton-Watson process are marked *red* or *blue*: red particles represent infected individuals in the coupled epidemic, while blue offspring of red parents represent attempted infections that are not allowed because the attempt is made on an individual who is not susceptible, or has already been infected by another contagious individual. Colors are assigned as follows: (1) Offspring of blue particles are always blue. (2) Each red particle reproduces by tossing a p-coin N times, once for each individual i in the population. Each Head counts as an offspring, and each represents an *attempted* infection. If several red particles attempt to infect the same individual i, exactly one of these is marked as a success (red), and the others are marked as failures (blue). Also, if an attempt is made to infect an individual who is not susceptible, the corresponding particle is colored blue. Clearly, the collection of all particles (red and blue) evolves as a Galton-Watson process, while the collection of red particles evolves as the infected set in the SIS epidemic. See the figure below for a typical evolution of the coupling in a population of size N = 80,000 with p = 1/80000 and 200 individuals initially infected.

2. Critical behavior: mean-field case

When studying the behavior of the simple SIR and SIS epidemics in large populations, it is natural to consider the scaling $p = p_N = \lambda_N/N$ for the infection parameter p. In this scaling, $\lambda = \lambda_N$ is the mean of the offspring distribution in the branching envelope. If $\lambda < 1$ then the epidemic will end quickly, even if a large number of individuals are initially infected. On the other hand, if $\lambda > 1$ then with positive probability (approximately one minus the extinction probability for the associated branching envelope), even if only one individual is initially infected, the epidemic will be large, with a positive limiting fraction of the population eventually being infected. The large-N behavior of the size of the SIR epidemic in this case is well understood: see for example [12] and [14].

2.1. Critical scaling: size of the epidemic

The behavior of both the SIS and SIR epidemics is more interesting in the critical case $\lambda_N \approx 1$. When the set of individuals initially infected is sufficiently small relative to the population size, the epidemic can be expected to evolve in much the same manner as a critical Galton-Watson process with Poisson-1 offspring distribution. However, when the size of the initially infected set passes a certain critical threshold, then the epidemic will begin to deviate substantially from the branching envelope. For the SIR case, the critical threshold was (implicitly) shown by [11] and [1] (see also [12]) to be at $N^{1/3}$, and that the critical scaling window is of width $N^{-4/3}$:

Theorem 1 ([11], [1]). Assume that $p_N = 1/N + a/N^{4/3} + o(n^{-4/3})$, and that the number J_0^N of initially infected individuals is such that $J_0^N/N^{1/3} \to b$ as the population size $N \to \infty$. Then as $N \to \infty$, the size $U_N := \sum_t J_t$ obeys the asymptotic law

(4)
$$U_N/N^{2/3} \xrightarrow{\mathcal{D}} T_b$$

where T_b is the first passage time to the level b by $W_t + t^2/2 + at$, and W_t is a standard Wiener process.

The distribution of the first passage time T_b can be given in closed form: See [11], also [8], [13].

For the critical SIS epidemic, the critical threshold is at $N^{1/2}$, and the critical scaling window is of width $N^{-3/2}$:

Theorem 2 ([4]). Assume that $p_N = 1/N + a/N^{3/2} + o(n^{-3/2})$, and that the initial number of infected individuals satisfies $J_0^N \sim bN^{1/2}$ as $N \to \infty$. Then the total number of infections $U_N := \sum_t J_t$ during the course of the epidemic obeys

(5)
$$U_N/N \xrightarrow{\mathcal{D}} \tau(b-a;-a)$$

where $\tau(x; y)$ is the time of first passage to y by a standard Ornstein-Uhlenbeck process started at x.

2.2. Critical scaling: time evolution of the epidemic

For both the *SIR* and *SIS* epidemics, if the number of individuals initially infected is much below the critical threshold then the evolution of the epidemic will not differ noticeably from that of its branching envelope. It was observed by [7] (and proved by [9]) that a (near-) critical Galton-Watson process initiated by a large number M of individuals behaves, after appropriate rescaling, approximately as a *Feller* diffusion: In particular, if Z_n^M is the size of the *n*th generation of a Galton-Watson with $Z_0^M \sim bM$ with offspring distribution Poisson(1 + a/M)then as $M \to \infty$,

(6)
$$Z^M_{[Mt]}/M \xrightarrow{\mathcal{D}} Y_t$$

where Y_t satisfies the stochastic differential equation

(7)
$$dY_t = aY_t dt + \sqrt{Y_t} dW_t$$
$$Y_0 = b.$$

What happens at the critical threshold, in both the *SIR* and *SIS* epidemics, is that the deviation from the branching envelope exhibits itself as a size-dependent drift in the limiting diffusion:

Theorem 3 ([4]). Let $J^N(n) = J^N_{[n]}$ be the number infected in the nth generation of a simple SIS epidemic in a population of size N. Then under the hypotheses of Theorem 2,

(8)
$$J^N(\sqrt{Nt})/\sqrt{N} \xrightarrow{\mathcal{D}} Y_t$$

where $Y_0 = b$ and Y_t obeys the stochastic differential equation

(9)
$$dY_t = (aY_t - Y_t^2) dt + \sqrt{Y_t} dW_t$$

Note that the diffusion (9) has an entrance boundary at ∞ , so that it is possible to define a version Y_t of the process with initial condition $Y_0 = 0$. When the *SIS* epidemic is begun with $J_0^N \gg \sqrt{N}$ initially infected, the number J_t^N infected will rapidly drop (over the first $\varepsilon \sqrt{N}$ generations) until reaching a level of order \sqrt{N} , and then evolve as predicted by (8). The following figure depicts a typical evolution in a population of size N = 80,000, with infection parameter p = 1/N and $I_0 = 10,000$ initially infected.

Theorem 4 ([4]). Let $J^N(n) = J^N_{[n]}$ and $R^N(n) = R^N_{[n]}$ be the numbers of infected and recovered individuals in the nth generation of a simple SIR epidemic in a population of size N. Then under the hypotheses of Theorem 1,

(10)
$$\begin{pmatrix} N^{-1/3}J^N(N^{1/3}t) \\ N^{-2/3}R^N(N^{1/3}t) \end{pmatrix} \xrightarrow{\mathcal{D}} \begin{pmatrix} J(t) \\ R(t) \end{pmatrix}$$

where $J_0 = b$, $R_0 = 0$, and

(11)
$$dJ(t) = (aJ(t) - J(t)R(t)) dt + \sqrt{J(t)} dW_t,$$
$$dR(t) = J(t) dt.$$

Theorems 1–2 can be deduced from Theorems 3–4 by simple time-change arguments (see [4]).

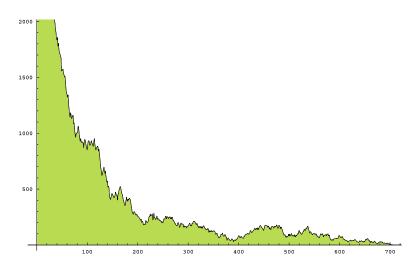


Fig 3.

2.3. Critical scaling: heuristics

The critical thresholds for the SIS-d and SIR-d epidemics can be guessed by simple comparison arguments using the standard couplings of the epidemics with their branching envelopes. Consider first the critical SIS epidemic in a population of size N. Recall (Section 1.4) that the branching envelope is a critical Galton-Watson process whose offspring distribution is Binomial-(N, 1/N). The particles of this Galton-Watson process are marked red or blue, in such a way that in each generation the number of red particles coincides with the number of infected individuals in the SIS epidemic. Offspring of blue particles are always blue, but offspring of red particles may be either red or blue; the blue offspring of red parents in each generation represent attempted infections that are suppressed.

Assume that initially there are N^{α} infected individuals and thus also N^{α} individuals in the zeroth generation of the branching envelope. By Feller's theorem, we may expect that the extinction time of the branching envelope will be on the order N^{α} , and that in each generation up to (shortly before) extinction the branching process will have order N^{α} individuals. If α is small enough that the SIS epidemic obeys the same rough asymptotics (that is, stays alive for $O(N^{\alpha})$ generations and has $O(N^{\alpha})$ infected individuals in each generation), then the number of blue off-spring of red parents in each generation will be on the order $N \times (N^{2\alpha}/N^2)$ (because for each of the N individuals of the population, the chance of a double infection is about $N^{2\alpha}/N^2$). Since the duration of the epidemic will be of the same rough order of magnitude as the size of the infected set in each generation, there should be at most O(1) blue offspring of red parents in any generation (if there were more, the red population would die out long before the blue). Thus, the critical threshold must be at $N^{1/2}$.

A similar argument applies for the SIR epidemic. The branching envelope of the critical SIR is once again a critical Galton-Watson process with offspring distribution Binomial-(N, 1/N), with constituent particles again labeled red or blue, red particles representing infected individuals in the epidemic. The rule by which red particles reproduce is as follows: Each red particle tosses a p-coin N times once for each individual i in the population. Each Head counts as an offspring, and represents an attempted infection. However, if a Head occurs on a toss at individual i where i was infected in an earlier generation, then the Head results in a *blue* offspring. Similarly, if more than one red particle tosses a Head at an individual i which has not been infected earlier, then one of these is labeled red and the excess are all labeled blue.

Assume that initially there are N^{α} infected individuals. As before, we may expect that the extinction time of the branching envelope will be on the order N^{α} , and that in each generation up to extinction the branching process will have order N^{α} individuals. If α is small enough, the extinction time and the size of the red population will also be $O(N^{\alpha})$. Consequently, the size of the recovered population will be (for all but the first few generations) on order $N^{2\alpha}$. Thus, in each generation, the number of blue offspring of red parents will be on order $(N^{2\alpha}/N) \times N^{\alpha}$ (because the chance that a recovered individual is chosen for attempted infection by an infected individual is $O(N^{\alpha}/N)$). Therefore, by similar reasoning as in the SIS case, the critical threshold is at $N^{1/3}$, as this is where the the number of blue offspring of red parents in each generation is O(1).

3. Critical behavior: SIS-1 and SIR-1 Spatial epidemics

Consider now the spatial SIS-d and SIR-d epidemic models on the d-dimensional integer lattice \mathbb{Z}^d . Assume that the village size $N_x = N$ is the same for all sites $x \in \mathbb{Z}^d$, and that the infection probabilities $\alpha(x, y)$ are nearest neighbor homogeneous, and uniform, that is,

(12)
$$\alpha(x,y) = \begin{cases} p = p_N, & \text{if } |x-y| \le 1; \\ 0, & \text{otherwise.} \end{cases}$$

3.1. Scaling limits of branching random walks

The branching envelope of a spatial SIS-d or SIR-d epidemic is a nearest neighbor branching random walk on the integer lattice \mathbb{Z}^d . This evolves as follows: Any particle located at site x at time t lives for one unit of time and then reproduces, placing random numbers ξ_y of offspring at the sites y such that $|y - x| \leq 1$. The random variables ξ_y are mutually independent, with Binomial- (N, p_N) distributions.

The analogue for branching random walks of Feller's theorem for Galton-Watson processes is *Watanabe's theorem*. This asserts that, after suitable rescaling, as the particle density increases, critical branching random walks converge to a limit, the *Dawson-Watanabe* process, also known as *super Brownian motion*. A precise statement follows: Consider a sequence of branching random walks, indexed by $M = 1, 2, \ldots$, with offspring distribution Binomial- (N, p_M) as above, and

(13)
$$p_M = p_{N,M} = \frac{1}{(2d+1)N} - \frac{a}{NM}$$

(Note: N may depend on M.) The rescaled measure-valued process X_t^M associated with the Mth branching random walk puts mass 1/M at location x/\sqrt{M} and time t for each particle of the branching random walk that is located at site x at time [Mt]. (Note: The branching random walk is a discrete-time process, but the associated measure-valued process runs in continuous time.)

Watanabe's theorem ([15]). Assume that the initial values X_0^M converge weakly (as finite Borel measures on \mathbb{R}^d) to a limit measure X_0 . Then under the hypothesis (13) the measure-valued processes X_t^M converge in law as $M \to \infty$ to a limit X_t :

(14)
$$X_t^M \Longrightarrow X_t$$
.

The limit process is the *Dawson-Watanabe* process with killing rate a and initial value X_0 . (The term *killing rate* is used because the process can be obtained from the "standard" Dawson-Watanabe process (a = 0) by elimination of mass at constant rate a.) The Dawson-Watanabe process X_t with killing rate a can be characterized by a martingale property: For each test function $\phi \in C_c^2(\mathbb{R}^d)$,

(15)
$$\langle X_t, \phi \rangle - \langle X_0, \phi \rangle - \frac{\sigma}{2} \int_0^t \langle X_s, \Delta \phi \rangle \, ds + a \int_0^t \langle X_s, \varphi \rangle \, ds$$

is a martingale. Here $\sigma^2 = 2d/(2d+1)$ is the variance of the random walk kernel naturally associated with the branching random walks. It is known [10] that in dimension d = 1 the random measure X_t is for each t absolutely continuous relative to Lebesgue measure, and the Radon-Nikodym derivative X(t, x) is jointly continuous in t, x (for t > 0). In dimensions $d \ge 2$ the measure X_t is almost surely singular, and is supported by a Borel set of Hausdorff dimension 2 [3].

3.2. Spatial SIS-1 and SIR-1 epidemics: critical scaling

As in the mean-field case, there are critical thresholds for the SIS-1 and SIR-1 epidemics at which they begin to deviate noticeably from their branching envelopes. These are at $N^{2/3}$ and $N^{2/5}$, respectively:

Theorem 5 ([5]). Fix $\alpha > 0$, and let X_t^N be the discrete-time measure-valued process obtained from an SIS-1 or an SIR-1 epidemic on a one-dimensional grid of size-N villages by attaching mass $1/N^{\alpha}$ to the point $(t, x/N^{\alpha/2})$ for each infected individual at site x at time $[tN^{\alpha}]$. Assume that X_0^N converges weakly to a limit measure X_0 as the village size $N \to \infty$. Then as $N \to \infty$,

(16)
$$X^N_{[N^{\alpha}t]} \xrightarrow{\mathcal{D}} X_t,$$

where X_t is a measure-valued process with initial value X_0 whose law depends on the value of α and the type of epidemic (SIS or SIR) as follows:

- (a) SIS: If $\alpha < \frac{2}{3}$ then X_t is a Dawson-Watanabe process with variance σ^2 .
- (b) SIS: If $\alpha = \frac{2}{3}$ then X_t is a Dawson-Watanabe process with variance σ^2 and killing rate

(17)
$$\theta(x,t) = X(x,t)/2.$$

- (c) SIR: If $\alpha < \frac{2}{5}$ then X_t is a Dawson-Watanabe process with variance σ^2 . (d) SIR: If $\alpha = \frac{2}{5}$ then X_t is a Dawson-Watanabe process with variance σ^2 and killing rate

(18)
$$\theta(x,t) = X(x,t) \int_0^t X(x,s) \, ds.$$

The Dawson-Watanabe process with variance σ^2 and (continuous, adapted) killing rate $\theta(t, x, \omega)$ is characterized [2] by a martingale problem similar to (15) above: for each test function $\phi \in C^2_c(\mathbb{R})$,

(19)
$$\langle X_t, \phi \rangle - \langle X_0, \phi \rangle - \frac{\sigma}{2} \int_0^t \langle X_s, \Delta \phi \rangle \, ds + \int_0^t \langle X_s, \theta \varphi \rangle \, ds$$

is a martingale. The law of this process is mutually absolutely continuous relative to that of the Dawson-Watanabe process with no killing, and there is an explicit formula for the Radon-Nikodym derivative – see [2].

3.3. Critical scaling for spatial epidemics: heuristics

Arguments similar to those given above for the mean-field SIS and SIR epidemics can be used to guess the critical thresholds for the spatial SIS-1 and SIR-1 epidemics. For the spatial epidemics, the associated branching envelopes are branching random walks. In the standard couplings, particles of the branching envelope are labeled either red or blue, with the red particles representing infected individuals in the epidemics. As in the mean-field cases, offspring of blue particles are always blue, but offspring of red particles may be either red or blue; blue offspring of red parents represent attempted infections that are suppressed. These may be viewed as an *attrition* of the red population (since blue particles created by red parents are potential red offspring that are not realized!).

Consider first the SIS-1 epidemic. Assume that initially there are N^{α} particles, distributed (say) uniformly among the $N^{\alpha/2}$ sites nearest the origin. Then by Feller's limit theorem (recall that the total population size in a branching random walk is a Galton-Watson process), the branching envelope can be expected to survive for $O_P(N^{\alpha})$ generations, and at any time prior to extinction the population will have $O_P(N^{\alpha})$ members. These will be distributed among the sites at distance $O_P(N^{\alpha/2})$ from the origin, and therefore in dimension d = 1 there should be about $O_P(N^{\alpha/2})$ particles per site. Consequently, for the SIS-1 epidemic, the rate of attrition per site per generation should be $O_P(N^{\alpha-1})$, and so the total attrition rate per generation should be $O_P(N^{3\alpha/2-1})$. If $\alpha = 2/3$, then the total attrition through the duration of the branching random walk envelope will be on the same order of magnitude as the population size N^{α} .

For the SIR-1 epidemic there is a similar heuristic calculation. As for the SIS-1 epidemic, the branching envelope will survive for $O_P(N^{\alpha})$ generations, and up to the time of extinction the population should have $O_P(N^{\alpha})$ individuals, about $O_P(N^{\alpha/2})$ per site. Therefore, through N^{α} generations, about $N^{\alpha} \times N^{\alpha/2}$ numbers j will be retired, and so the attrition rate per site per generation should be $O_P(N^{\alpha/2} \times N^{3\alpha/2})$, making the total attrition rate per generation $O_P(N^{5\alpha/2})$. Hence, if $\alpha = 2/5$ then the total attrition per generation should be $O_P(1)$, just enough so that the total attrition through the duration of the branching random walk envelope will be on the same order of magnitude as the population size.

3.4. Critical scaling in dimensions $d \geq 2$

In higher dimensions, the critical behavior of the SIS-d and SIR-d epidemics appears to be considerably different. We expect that there will be no analogous threshold effect, in particular, we expect that the epidemic will behave in the same manner as the branching envelope up to the point where the infected set is a positive fraction of the total population. This is because in dimensions $d \ge 2$, the particles of a critical branching random walk quickly spread out, so that (after a short initial period) there are only $O_P(1)$ (in dimensions $d \ge 3$) or $O_P(\log N)$ (in dimension d = 2) particles per site. (With N particles initially, a critical branching random walk typically lives O(N) generations, and particles are distributed among the sites at distance $O(\sqrt{N})$ from the origin; in dimensions $d \ge 2$, there are $O(N^{d/2})$ such sites, enough to accomodate the O(N) particles of the branching random walk without crowding.) Consequently, the rate at which "multiple" infections are attempted (that is, attempts by more than one contagious individual to simultaneously infect the same susceptible) is only of order $O_P(1/N)$ (or, in dimension d = 2, order $O_P(\log N/N)$).

The interesting questions regarding the evolution of critical epidemics in dimensions $d \ge 2$ center on the initial stages, in the relatively small amount of time (order o(N) generations) in which the particles spread out from their initial sites. These will be discussed in the forthcoming University of Chicago Ph. D. dissertation of Xinghua Zheng.

3.5. Weak convergence of densities

There is an obvious gap in the heuristic argument of Section 3.3 above: Even if the total number of infected individuals is, as expected, on the order N^{α} , and even if

these are concentrated in the sites at distance on the order $N^{\alpha/2}$ from the origin, it is by no means obvious that these will distribute themselves uniformly (or at least *locally* uniformly) among these sites. The key step in filling this gap in the argument is to show that the particles of the branching envelope distribute themselves more or less uniformly on scales smaller than $N^{\alpha/2}$.

Consider, as in Section 3.1, a sequence of branching random walks, indexed by $M = 1, 2, \ldots$, with offspring distribution Binomial- (N, p_M) as above, and p_M given by (13). Let $Y_t^M(x)$ be the number of particles at site x at time [t], and let $X^M(t, x)$ be the continuous function of $t \ge 0$ and $x \in \mathbb{R}$ obtained by linear interpolation from the values

(20)
$$X^{M}(t,x) = \frac{Y_{Mt}(\sqrt{M}x)}{\sqrt{M}} \quad \text{for} \quad Mt \in \mathbb{Z}_{+} \text{ and } \sqrt{M}x \in \mathbb{Z}.$$

Theorem 6 ([5]). Assume that d = 1. Assume also that the initial particle configuration is such that all particles are located in an interval $[-\kappa\sqrt{M}, \kappa\sqrt{M}]$ and such that the initial particle density satisfies

(21)
$$X^M(0,x) \Longrightarrow X(0,x)$$

as $M \to \infty$ for some continuous function X(0, x) with support $[-\kappa, \kappa]$. Then as $M \to \infty$,

(22)
$$X^M(t,x) \Longrightarrow X(t,x),$$

where X(t,x) is the density function of the Dawson-Watanabe process with killing rate a and initial value X(0,x). The convergence is relative to the topology of uniform convergence on compacts in the space $C(\mathbb{R}_+ \times \mathbb{R})$ of continuous functions.

Since the measure-valued processes associated with the densities $X^{M}(t, x)$ are known to converge to the Dawson-Watanabe process, by Watanabe's theorem, to prove Theorem 6 it suffices to establish tightness. This is done by a somewhat technical application of the Kolmogorov-Chentsov tightness criterion, based on a careful estimation of moments. See [5] for details.

It is also possible to show that convergence of the particle density processes holds in Theorem 5.

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