

SUPERLINEAR INDEFINITE ELLIPTIC PROBLEMS AND NONLINEAR LIOUVILLE THEOREMS

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A Jean Leray, en témoignage de profonde admiration

1. Introduction

We study the following elliptic boundary value problem in a bounded domain Ω in \mathbb{R}^N , with smooth boundary:

$$(1.1) \quad \begin{aligned} u > 0, \quad Lu + a(x)g(u) = 0 & \quad \text{in } \Omega, \\ Bu = 0 & \quad \text{on } \partial\Omega. \end{aligned}$$

Here, L is a linear elliptic operator—we use summation convention—

$$L = a_{ij}(x) \frac{\partial^2}{\partial x_i \partial x_j} + b_i(x) \frac{\partial}{\partial x_i} + c(x),$$

with $a_{ij} \in C^2(\bar{\Omega})$, $b_i \in C^1(\bar{\Omega})$ and $c \in L^\infty$. We assume uniform ellipticity

$$c_0 |\xi|^2 \leq a_{ij}(x) \xi_i \xi_j \leq C_0 |\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \forall x \in \Omega,$$

with $c_0, C_0 > 0$. The boundary operator B is one of the following:

$$(1.2a) \quad Bu := u,$$

$$(1.2b) \quad Bu := \nu_j a_{jk} u_{x_k} + \alpha(x)u,$$

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where $\nu_i = (\nu_1, \dots, \nu_N)$ denotes the exterior unit normal on $\partial\Omega$; α is a given continuous nonnegative function on $\partial\Omega$.

We are interested in the case that the function a changes sign. The problem is then called one with indefinite nonlinearity. In case

$$L = \Delta - m(x), \quad m \in L^\infty,$$

several recent papers treat the problem for special functions g , by variational methods (see [4] and [1], when $g = u^p$, $1 < p < (N + 2)/(N - 2)$, [8] treats the problem on a compact manifold using bifurcation analysis). [1] also treats a wider class of functions g for the Dirichlet problem, i.e., for B given by (1.2a). In case g is odd they also obtain multiple solutions.

For general L , problem (1.1) does not admit a variational approach. We always assume that g is a C^1 function on \mathbb{R}^+ with

$$(1.3) \quad g(0) = g'(0) = 0 \quad \text{and} \quad g(s) > 0 \quad \text{for } s > s_1 > 0.$$

Our main result, which is proved with the aid of degree theory, is concerned with functions g which have precise power-like growth at infinity:

$$(1.4) \quad \lim_{s \rightarrow \infty} \frac{g(s)}{s^p} = l > 0, \quad \text{for some } p > 1.$$

Concerning the function a , we assume it belongs to $C^2(\overline{\Omega})$, that

$$\Omega^+ := \{x \in \Omega : a(x) > 0\} \quad \text{and} \quad \Omega^- := \{x \in \Omega : a(x) < 0\}$$

are nonempty, and that

$$(1.5) \quad \Gamma := \overline{\Omega^+} \cap \overline{\Omega^-} \subset \Omega, \quad \text{with } \nabla a(x) \neq 0 \quad \forall x \in \Gamma.$$

Our results depend on the sign of the principal eigenvalue $\lambda_1 = \lambda_1(-L)$ of the operator $-L$ in Ω under the boundary conditions $Bu = 0$. The eigenvalue λ_1 is such that there is a function φ satisfying

$$\begin{aligned} \varphi > 0, \quad (L + \lambda_1)\varphi &= 0 & \text{in } \Omega, \\ B\varphi &= 0 & \text{on } \partial\Omega. \end{aligned}$$

In case B is as in (1.2b), it follows easily with the aid of the Hopf lemma that

$$\varphi > 0 \quad \text{in } \overline{\Omega}.$$

Our main existence result is

THEOREM 1. *Assume (1.3), (1.4) and (1.5) and assume that $\lambda_1(-L) > 0$. Then problem (1.1) has a solution provided*

$$(1.6) \quad 1 < p < \frac{N+2}{N-1}.$$

The upper bound in (1.6) is less than the usual critical exponent $(N+2)/(N-2)$. The limitation (1.6) is due to that in our nonlinear Liouville Theorem 4 below. We believe that Theorem 1 should hold if $1 < p$ for $N \leq 2$, and if $N \geq 3$, also for $p < (N+2)/(N-2)$.

The next existence results refer to the problem involving a parameter $\tau \leq 0$:

$$(1.7) \quad \begin{aligned} u > 0, \quad (L - \tau)u + a(x)g(u) &= 0 && \text{in } \Omega, \\ Bu &= 0 && \text{on } \partial\Omega, \end{aligned}$$

in case $\lambda_1(-L) = 0$.

THEOREM 2. *Assume (1.3), (1.5) and that $\lambda_1(-L) = 0$. Assume further that*

$$(1.8) \quad \lim_{s \rightarrow 0} \frac{g(s)}{s^q} = \alpha \neq 0, \quad \text{for some } q > 1,$$

and that

$$(1.9) \quad \alpha \int_{\Omega} a(x)\varphi^q \psi < 0,$$

where $\psi > 0$ is the principal eigenfunction of the adjoint operator $L^* = \partial_{ij}a_{ij} - \partial_i b_i + c$, i.e. $L^*\psi = 0$ in Ω under the adjoint boundary condition $B^*\psi = 0$, with

$$B^*\psi = \psi \quad \text{in case (1.2a),}$$

$$B^*\psi = \nu_i(a_{ij}\psi)_j + (b_i\nu_i - \alpha)\psi \quad \text{in case (1.2b).}$$

Then there exists τ^* , $0 > \tau^* > -\infty$, such that for $0 > \tau > \tau^*$, problem (1.7) has a solution, but for $\tau < \tau^*$ it has no solution.

THEOREM 3. *Assume the conditions of Theorem 2 and, in addition, that (1.4) holds with $p < (N+2)/(N-1)$. Then (1.7) has a solution when $\tau = 0$.*

The proofs of the theorems involve several ingredients. Theorem 1 relies on Leray-Schauder degree theory, and for this purpose it is necessary to establish a priori estimates. These are derived in Section 3. There we obtain a bound for the L^∞ norm of solutions of a one-parameter family of problems

$$\begin{aligned} \rho \geq 0, \quad Lu + a(x)g(u^+) + \rho &= 0 && \text{in } \Omega, \\ Bu &= 0 && \text{on } \partial\Omega. \end{aligned}$$

The bound

$$(1.10) \quad u \leq \bar{C}$$

is established, with \bar{C} independent of ρ , but depending on L , the function a , and on Ω .

The derivation of (1.10) forms the heart of paper. The proof follows a line of argument similar to that used in Gidas and Spruck in [6]. It is indirect: we assume there is no such bound and then use blow up arguments to obtain a contradiction. A number of cases must be treated. To derive the contradiction we use two nonlinear forms of Liouville's theorem, Theorems 1.2 and 1.3 in [6]. In addition, we establish some new nonlinear Liouville theorems. In proving (1.10) we use the following one. (It is a consequence of a more general result of Liouville type, Theorem 2.1 in Section 2, which we present in the belief that it will prove useful in other problems.)

THEOREM 4. *In the half space*

$$\Sigma = \{x \in \mathbb{R}^N : x_N > 0\}$$

let u be a nonnegative function in $C^2(\Sigma)$ which is also bounded near the origin and satisfies

$$(1.11) \quad \Delta u + x_N u^p \leq 0 \quad \text{in } \Sigma.$$

Then

$$(1.12) \quad u \equiv 0 \quad \text{provided } p < \frac{N+2}{N-1}.$$

REMARKS. Note that no condition is assumed about the behaviour of u near infinity or near $\{x_N = 0\}$, except at the origin. If, in place of (1.11),

$$(1.13) \quad \Delta u + x_N u^p \equiv 0 \quad \text{in } \Sigma$$

holds, then we believe that the conclusion $u \equiv 0$ should be true for a larger range of p than given in (1.12).

For simplicity we carry out the proofs of Theorems 1–3 for L of the form

$$L = \Delta - m(x);$$

the arguments work also for our general L which is not self-adjoint.

It is clear from the arguments that more general functions a could be treated. For example, one might permit Γ to intersect $\partial\Omega$ transversally or Γ to have clean self-intersections. The more general Liouville theorem of Section 2 would then be called upon. We leave this for further consideration.

In Section 4, using the a priori estimate (1.10) we prove Theorem 1. Theorem 3 is then proved in Section 5, using Theorem 1 and a bifurcation analysis. Similar bifurcation analysis then yields small solutions of (1.7) in case $0 < -\tau$

is small. Finally, in Section 6, we use sub- and supersolutions to complete the proof of Theorem 2.

In the proof of Theorem 2 we make use of the following simple

LEMMA 1. *Suppose Ω^+ is nonempty and that g satisfies (1.3). If u is a solution of (1.7) for some $\tau < 0$, then*

$$(1.14) \quad -\tau \leq \tilde{C},$$

where \tilde{C} depends only on L , on a , and on a constant C which is such that

$$|g(s)| \leq Cs \quad \text{on } [0, s_1].$$

PROOF. Let B be an open ball in Ω^+ and let μ be the principal eigenvalue of $-L$ in B under Dirichlet boundary conditions, i.e. there is a function $\tilde{\varphi}$ in B satisfying

$$\begin{aligned} \tilde{\varphi} > 0, \quad (L + \mu)\tilde{\varphi} &= 0 & \text{in } B, \\ \tilde{\varphi} &= 0 & \text{on } \partial B. \end{aligned}$$

Since $a > 0$ in B , and $g(u) > 0$ for $u > s_1$, we see from (1.7) that in B ,

$$(L - \tau)u = -a(x)g(u) \leq C\|a\|_{L^\infty}u.$$

Hence,

$$u > 0, \quad Lu - (\tau + C\|a\|_{L^\infty})u \leq 0 \quad \text{in } B.$$

It follows (see [5]) that $-(\tau + C\|a\|_{L^\infty}) \leq \mu$, which yields (1.14).

The proof in Section 6 of Theorem 2, using sub- and supersolutions, is the same as one in [4].

It would be interesting to obtain some information about τ^* —even for the operator $\Delta - m(x)$ —and to determine if (1.7) has a solution when $\tau = \tau^*$.

A FINAL REMARK. The derivation in Section 3 of the a priori estimate (1.10) involves blow up arguments. When dealing with $L = \Delta - m(x)$, these arguments lead to equations of the form

$$(1.15) \quad v > 0, \quad \Delta v + v^p = 0$$

in all of \mathbb{R}^N or in a half space, with then

$$(1.16) \quad v = 0 \quad \text{or} \quad \nu \cdot \nabla v = 0 \quad \text{on the boundary.}$$

Or else, the blow up leads to

$$(1.17) \quad v > 0, \quad \Delta v + x_N v^p = 0 \quad \text{in } \mathbb{R}^N.$$

For general L , these equations would, instead, take the form

$$v > 0, \quad a_{jk}v_{x_j x_k} + v^p = 0$$

in \mathbb{R}^N or in a half space, with then

$$v = 0 \quad \text{or} \quad \nu_j a_{jk} v_{x_k} = 0 \quad \text{on the boundary.}$$

Or else we would find

$$v > 0, \quad a_{jk} v_{x_j x_k} + a_j x_j v^p = 0 \quad \text{in } \mathbb{R}^N$$

where the a_{jk} and a_j are constants, $\sum |a_j| > 0$, with $a_{jk} \partial_{jk}$ elliptic of course. After a suitable linear transformation of independent variables, and multiplication of v by a factor, we are easily reduced to the cases (1.15), (1.16) or (1.17).

Some of the results of this paper have been announced in [3].

2. Liouville theorems in cones

Let Σ be an open connected cone in \mathbb{R}^N , $N \geq 2$, with vertex at the origin and with $\bar{\Sigma} \neq \mathbb{R}^N$. Let u be a nonnegative function in $C^2(\Sigma)$, bounded near the origin, satisfying

$$(2.1) \quad \Delta u + h(x)u^p \leq 0,$$

where $0 \leq h \in C(\Sigma)$, bounded near the origin, and

$$(2.2) \quad h(x) = a|x|^\gamma \quad \text{for } |x| \text{ large, } \gamma > -2, a > 0.$$

THEOREM 2.1. *Let λ_1 be the principal eigenvalue for the Dirichlet problem on $\Sigma \cap S^{N-1}$ of $-\Delta_S$, the Laplacian on S^{N-1} . Define $\alpha > 0$ by the identity*

$$(2.3) \quad \lambda_1 = \alpha(N + \alpha - 2).$$

Assume that p satisfies

$$(2.4) \quad 1 < p \leq \frac{N + \alpha + \gamma}{N + \alpha - 2} := \sigma.$$

Let $u \geq 0$ satisfy (2.1), with h as above. Then $u \equiv 0$.

In the theorem no regularity of $\partial\Sigma$ is assumed and we always take $a = 1$.

COROLLARY 2.1. *Let Σ and u be as above. Assume that u satisfies (2.1), where h is a positive continuous function in Σ which is homogeneous of degree $\gamma > -2$ for $|x|$ large. Then $u \equiv 0$ if*

$$(2.5) \quad 1 < p < \frac{N + \alpha + \gamma}{N + \alpha - 2},$$

where α is defined as in Theorem 2.1.

PROOF OF COROLLARY 2.1. Observe that $h(x)|x|^{-\gamma}$, for large $|x|$, may vanish on $\partial\Sigma$. Then we simply apply Theorem 2.1 in a slightly smaller cone $\tilde{\Sigma}$. Corresponding to $\tilde{\Sigma}$ we may have $\tilde{\lambda}_1 > \lambda_1$, with $\tilde{\lambda}_1 - \lambda_1$ as small as we like (see

[5]). For the corresponding $\tilde{\alpha}$ in (2.3), we have $0 < \tilde{\alpha} - \alpha$ small, hence $0 < \sigma - \tilde{\sigma}$ small. By the theorem, $u \equiv 0$ in $\tilde{\Sigma}$ and hence in Σ .

PROOF OF THEOREM 2.1. Set $\Omega = \Sigma \cap S^{N-1}$ and let $\{\Omega_j\}$ be an increasing sequence of domains on S^{N-1} , with smooth boundaries, such that

$$\Omega_j \subset \bar{\Omega}_j \subset \Omega_{j+1} \subset \dots \subset \Omega, \quad \bigcup \Omega_j = \Omega.$$

Let φ_j be the principal eigenfunction for the Dirichlet problem of $-\Delta_S$ on Ω_j with principal eigenvalue λ_j , i.e.

$$\begin{aligned} \varphi_j > 0 \quad (\Delta_S + \lambda_j)\varphi_j &= 0, & \text{in } \Omega_j, \\ \varphi_j &= 0 & \text{on } \partial\Omega_j. \end{aligned}$$

We normalize the φ_j by requiring them to equal 1 at some fixed point $x_0 \in \Omega$. The functions φ_j are then uniformly bounded by some fixed constant C_1 (see Theorem 2.1 and its proof in [5]).

Let Σ_j be the cone generated by Ω_j , with vertex at 0. With $\alpha_j > 0$ chosen as in (2.3), the functions

$$g_j = |x|^{\alpha_j} \varphi_j(x/|x|)$$

are positive and harmonic in Σ_j and vanish on $\partial\Sigma_j$.

Let ζ be a C^∞ function defined on $[0, \infty)$ with $0 \leq \zeta \leq 1$, $\zeta(t) \equiv 1$ on $[0, 1/2]$, $\zeta(t) \equiv 0$ for $t \geq 1$. For $R > 0$, let $\zeta_R(x) = \zeta(|x|/R)$. With ε, R respectively small and large positive numbers and with $q = p/(p-1)$, set

$$I_{j,\varepsilon} = \int_{\Sigma_j} \zeta_R^q (1 - \zeta_\varepsilon) g_j h u^p.$$

By (2.1),

$$I_{j,\varepsilon} \leq - \int_{\Sigma_j} \zeta_R^q (1 - \zeta_\varepsilon) g_j \Delta u$$

so that

$$I_{j,\varepsilon} \leq \int_{\partial\Sigma_j} u \zeta_R^q (1 - \zeta_\varepsilon) \frac{\partial}{\partial \nu} g_j - \int_{\Sigma_j} u \Delta (\zeta_R^q (1 - \zeta_\varepsilon) g_j),$$

where ν is the exterior normal to $\partial\Sigma_j$. Since $\frac{\partial}{\partial \nu} g_j \leq 0$ on $\partial\Sigma_j$ we find

$$I_{j,\varepsilon} \leq - \int_{\Sigma_j} u \Delta (\zeta_R^q (1 - \zeta_\varepsilon) g_j).$$

We now let ε go zero. The only term which requires some care is

$$\int u \zeta_R^q \Delta (\zeta_\varepsilon g_j).$$

This is integrated over $x \in \Sigma_j, |x| < \varepsilon$. Since u is bounded near the origin this term is

$$O(\varepsilon^{\alpha_j} \varepsilon^{N-2}).$$

Here is the only place where we use $\alpha > 0$, in case $N = 2$. Consequently,

$$I_j := \int_{\Sigma_j} \zeta_R^q g_j h u^p \leq - \int_{\Sigma_j} u \Delta(\zeta_R^q g_j);$$

hence

$$I_j \leq - \int_{\Sigma_j} u g_j \Delta \zeta_R^q - 2 \int_{\Sigma_j} u \frac{\partial}{\partial r} \zeta_R^q \frac{\partial}{\partial r} g_j,$$

where $\partial/\partial r$ represents radial differentiation.

Using C to denote various constants independent of u , R and γ , we have

$$-\frac{C}{R} \leq \frac{\partial}{\partial r} \zeta_R^q \leq 0, \quad |\Delta \zeta_R| \leq \frac{C}{R^2} \quad \text{and} \quad -\Delta(\zeta_R^q) \leq -q \zeta_R^{q-1} \Delta \zeta.$$

Hence, since $\frac{\partial}{\partial r} g_j = \frac{\alpha_j}{r} g_j$,

$$\begin{aligned} I_j &\leq - \int_{\Sigma_j} u g_j q \zeta_R^{q-1} \Delta \zeta + \frac{C}{R^2} \int_{\Sigma_j} u \zeta_R^{q-1} g_j \\ &\leq \frac{C}{R^2} \int_{\Sigma_{j,R}} u \zeta_R^{q-1} g_j = \frac{C}{R^2} \int_{\Sigma_{j,R}} u \zeta_R^{q-1} h^{1/p} |x|^{-\gamma/p} g_j, \quad \text{for } R \text{ large,} \end{aligned}$$

where $\Sigma_{j,R} = \Sigma_j \cap \{x \in \mathbb{R}^N : R/2 < |x| < R\}$.

Thus, by Hölder's inequality,

$$\begin{aligned} I_j &\leq \frac{C}{R^2} \left[\int_{\Sigma_{j,R}} u^p \zeta_R^q h g_j \right]^{1/p} \left[\int_{\Sigma_{j,R}} |x|^{-\gamma/(p-1)} g_j \right]^{(p-1)/p} \\ &\leq \frac{C}{R^2} \left[\int_{\Sigma_{j,R}} u^p \zeta_R^q h g_j \right]^{1/p} R^{[N+\alpha_j-\gamma/(p-1)](p-1)/p}, \end{aligned}$$

with a different C . It follows that

$$I_j^{1-1/p} \leq C R^{[N+\alpha_j-\gamma/(p-1)](p-1)/p-2}.$$

Now let $j \rightarrow \infty$. The functions φ_j tend to φ , the principal eigenfunction of $-\Delta_S$ in Ω (see Section 4 in [5]). With $g = |x|^\alpha \varphi(x/|x|)$, we find that

$$I := \int_{\Sigma} u^p \zeta_R^q h g < \infty$$

and, furthermore

$$(2.6) \quad I_j \leq \frac{C}{R^2} \left[\int_{\Sigma_R} u^p \zeta_R^q h g \right]^{1/p} R^{[N+\alpha-\gamma/(p-1)](p-1)/p},$$

where $\Sigma_R := \Sigma \cap \{x \in \mathbb{R}^n : R/2 < |x| < R\}$. Consequently,

$$I^{1-1/p} \leq C R^{[N+\alpha-\gamma/(p-1)](p-1)/p-2}.$$

Condition (2.4) means that $\tau := [N + \alpha - \gamma/(p-1)]/(p-1)/p - 2 \leq 0$.

If $\tau < 0$, let $R \rightarrow \infty$. Then we conclude that

$$J := \int_{\Sigma} h g u^p = 0.$$

Thus, $u = 0$ for $|x|$ large. By the Maximum Principle it follows that $u \equiv 0$.

If $\tau = 0$ and we let $R \rightarrow \infty$, we conclude that $J < \infty$. But then, returning to (2.6) and letting $R \rightarrow \infty$, we find that the right-hand side in (2.6) tends to zero. Hence, $J = 0$ also in this case and we conclude as before that $u \equiv 0$.

PROOF OF THEOREM 4 OF INTRODUCTION. It is a special case of Corollary 2.1 when Σ is the half space $\{x_N > 0\}$ and $u \geq 0$ satisfies

$$\Delta u + x_N u^p \leq 0 \quad \text{in } \Sigma$$

and u is bounded near the origin. In this case, $g = x_N$ and so $\alpha = 1$, while $\gamma = 1$. From Corollary 2.1 it then follows that $u \equiv 0$ if $1 < p < (N + 2)/(N - 1)$.

The proof of Theorem 2.1 yields also the following simpler result:

THEOREM 2.2. *Let u be a nonnegative function in $\Sigma = \mathbb{R}^N - \{0\}$, $N \geq 3$, bounded near 0, and satisfying*

$$\Delta u + h(x)u^p \leq 0 \quad \text{in } \Sigma.$$

Here, $h \in C(\Sigma)$, $h \geq 0$ and $h(x) = a|x|^\gamma$, $a > 0$, $\gamma > -2$, for $|x|$ large. If

$$p \leq \frac{N + \gamma}{N - 2},$$

then $u \equiv 0$.

3. A priori estimates

In this section, for $\rho \geq 0$ we derive a priori estimates, independent of ρ , for the L^∞ norm of positive solutions of the problem

$$(3.1) \quad \begin{aligned} \rho \geq 0, \quad (\Delta - m(x))u + a(x)g(u) + \rho &= 0 && \text{in } \Omega, \\ Bu &= 0 && \text{on } \partial\Omega. \end{aligned}$$

Here, Ω is a bounded domain in \mathbb{R}^N with smooth boundary and B is one of the following boundary operators:

$$(3.2a) \quad Bu := u,$$

$$(3.2b) \quad Bu := \partial_\nu u + \alpha(x)u,$$

where ν is the outer normal to $\partial\Omega$ and $\alpha \geq 0$ is a given continuous function on $\partial\Omega$. We recall the assumption (1.4):

$$(3.3) \quad \lim_{s \rightarrow \infty} \frac{g(s)}{s^p} = l > 0 \quad \text{for some } p, \quad 1 < p < \frac{N + 2}{N - 1}.$$

By scaling we may suppose $l = 1$. Recall also (1.5):

$$(3.4) \quad \begin{aligned} a &\in C^2(\overline{\Omega}), \quad \Omega^+ := \{x \in \Omega : a(x) > 0\} \neq \emptyset, \\ \Omega^- &:= \Omega \setminus \overline{\Omega^+} \neq \emptyset, \\ \Gamma &:= \overline{\Omega^+} \cap \overline{\Omega^-} \subset \Omega \quad \text{and} \quad \nabla a \neq 0 \text{ on } \Gamma. \end{aligned}$$

Thus, a neighborhood in Ω of $\partial\Omega$ belongs either to Ω^+ or to Ω^- . We assume $m \in L^\infty$. In case of (3.2b) it follows with the aid of the Hopf lemma that $u > 0$ in $\overline{\Omega}$.

This section is devoted to the proof of

THEOREM 3.1. *Let $u \in W^{2,q}(\Omega)$, for all $q < \infty$, be a positive solution of (3.1). Suppose that (3.3) and (3.4) hold. Then, if $\|m\|_{L^\infty} \leq m_0$,*

$$(3.5) \quad 0 \leq u(x) \leq \overline{C}, \quad \forall x \in \overline{\Omega},$$

where \overline{C} is a constant depending only on m_0 , the function a , and Ω , and it is independent of $\rho \geq 0$.

From (3.5), with the aid of standard elliptic estimates we obtain the a priori estimate

$$(3.6) \quad \|u\|_{W^{2,q}} \leq \tilde{C} \quad \text{for } q < \infty,$$

where \tilde{C} depends only on q , m_0 , Ω , and the function a .

We need the following

LEMMA 3.1. *If u is a positive solution of (3.1) with $\|m\|_{L^\infty} \leq m_0$, then*

$$\rho \leq C \max u,$$

where C depends only on m_0 and Ω^+ .

PROOF. Since $g(u) > 0$ for u large, it follows from (3.1) that for some constant C_1 independent of ρ ,

$$\Delta u - m_0 u \leq C_1 - \rho \leq -\rho/2 \quad \text{in } \Omega^+$$

if $\rho > 2C_1$. Let σ be the solution of

$$\begin{aligned} (\Delta - m_0)\sigma &= -1 && \text{in } \Omega^+, \\ \sigma &= 0 && \text{on } \partial\Omega^+. \end{aligned}$$

Clearly, $\sigma > 0$ and for $\rho > 2C_1$, by the maximum principle,

$$w := u - \frac{\rho}{2}\sigma \geq 0 \quad \text{in } \Omega^+.$$

Since $C_2 := \max \sigma$ depends only on Ω^+ and on m_0 , it follows that $\rho \leq (2/C_2)u$ at a maximum point of σ and this yields the lemma.

PROOF OF THEOREM 3.1. We follow the approach in [6]. The proof is by contradiction and makes use of a blow up argument to reduce the problem of establishing the a priori estimate (3.5) to results of Liouville type.

Step 1. Suppose that the conclusion of Theorem 3.1 is false. Then there exist sequences m_j with $\|m_j\|_{L^\infty} \leq m_0$, $\rho_j \geq 0$ and a sequence $u_j \in W^{2,q}(\Omega)$ such that

$$(3.7) \quad \begin{aligned} u_j > 0, \quad (\Delta - m_j(x))u_j + a(x)g(u_j) + \rho_j &= 0 && \text{in } \Omega, \\ Bu_j &= 0 && \text{on } \partial\Omega \end{aligned}$$

and

$$(3.8) \quad M_j := \max_{\bar{\Omega}} u_j \rightarrow \infty \quad \text{as } j \rightarrow \infty.$$

We may assume that $M_j = u_j(x_j)$ for some $x_j \in \bar{\Omega}$, and that, for some $x_0 \in \bar{\Omega}$, $x_j \rightarrow x_0$ as $j \rightarrow \infty$. Let us introduce new scaled coordinates by setting

$$y = \frac{x - x_j}{\lambda_j}.$$

The positive scale factors λ_j will be chosen later with $\lambda_j \rightarrow 0$ as $j \rightarrow \infty$. Accordingly, we define a blow up function v_j by

$$(3.9) \quad v_j(y) = \frac{1}{M_j} u_j(x) = \frac{1}{M_j} u_j(\lambda_j y + x_j).$$

The function v_j is well defined for y in a suitable domain and

$$(3.10) \quad \max v_j = v_j(0) = 1, \quad j = 1, 2, \dots,$$

On the other hand, a direct computation shows that v_j satisfies

$$(3.11) \quad L_y v_j + \lambda_j^2 M_j^{p-1} \left(a(\lambda_j y + x_j) \frac{g(M_j v_j)}{M_j^p} + \frac{\rho_j}{M_j^p} \right) = 0,$$

where

$$L_y = \Delta_y - \lambda_j^2 m_j(\lambda_j y + x_j).$$

By Lemma 3.1, in every case

$$(3.12) \quad \rho_j / M_j^p \rightarrow 0 \quad \text{as } j \rightarrow \infty.$$

Observe also that for $\lambda_j \rightarrow 0$, $\lambda_j^2 m_j(\lambda_j y + x_j) \rightarrow 0$.

To proceed with the proof we consider several cases according to the location of the limit point x_0 , namely:

- (1) Case A: $x_0 \in \Gamma$,
- (2) Case B: $x_0 \in \Omega^+ \cup \Omega^-$,
- (3) Case C: $x_0 \in \partial\Omega$.

Step 2. Let us deal first with Case A. Set

$$\delta_j := \text{dist}(x_j, \Gamma) = |x_j - z_j|, \quad z_j \in \Gamma.$$

Since $\nabla a \neq 0$ on Γ , it follows that δ_j , which tends to zero, is given by

$$\delta_j = \pm \frac{\nabla a(z_j)}{|\nabla a(z_j)|} (x_j - z_j) \quad \text{for } j \text{ large,}$$

where the plus and the minus occur according as $x_j \in \Omega^+$ or $x_j \in \Omega^-$. Since $a(z_j) = 0$ and $a \in C^2(\Omega)$, we find by Taylor expansion,

$$a(\lambda_j y + x_j) = \pm |\nabla a(z_j)| \delta_j + \lambda_j \nabla a(z_j) \cdot y + O(\lambda_j^2 |y|^2 + \delta_j^2).$$

Substituting this expression into equation (3.11) we find

$$(3.13) \quad L_y v_j + \lambda_j^3 M_j^{p-1} \left[\nabla a(z_j) \cdot y \pm \frac{\delta_j}{\lambda_j} |\nabla a(z_j)| \right. \\ \left. + O\left(\lambda_j^2 |y|^2 + \frac{\delta_j^2}{\lambda_j}\right) \right] \frac{g(M_j v_j)}{M_j^p} + \frac{\lambda_j^2 \rho_j}{M_j} = 0$$

We now choose

$$(3.14) \quad \lambda_j = M_j^{(1-p)/3}.$$

Observe that equation (3.13) holds, for large j , in the ball

$$|y| < \frac{1}{2\lambda_j} \text{dist}(x_0, \partial\Omega).$$

We must now consider several cases:

Suppose

$$(i) \quad \delta_j/\lambda_j \rightarrow \infty \quad \text{possibly for a subsequence.}$$

Then set

$$\alpha_j := (\lambda_j/\delta_j)^{1/2} \quad \text{and} \quad \eta := y/\alpha_j.$$

Under this change of variables, equation (3.13) with λ_j specified by (3.14) becomes

$$(3.15) \quad L_\eta v_j + \alpha_j^2 \left[\alpha_j \nabla a(z_j) \cdot \eta \pm \frac{|\nabla a(z_j)|}{\alpha_j^2} \right. \\ \left. + O\left(\lambda_j \alpha_j^2 |\eta|^2 + \frac{\delta_j}{\alpha_j}\right) \right] \frac{g(M_j v_j)}{M_j^p} + \alpha_j^2 \lambda_j^2 \frac{\rho_j}{M_j} = 0,$$

where

$$L_\eta = \Delta_\eta - \alpha_j^2 \lambda_j^2 m_j (\lambda_j \alpha_j \eta + x_j).$$

By standard $W^{2,q}$ theory one may obtain estimates on the $v_j \leq 1$ ensuring that, for a subsequence, $v_j \rightarrow v$ locally uniformly, with v defined in all of \mathbb{R}^N

and $v \in W_{loc}^{2,q}$, for all $q < \infty$. Letting $j \rightarrow \infty$ in (3.15) we obtain, since $z_j \rightarrow x_0$, $\lambda_j, \delta_j, \alpha_j \rightarrow 0$ and using Lemma 3.1 and (3.3), that

$$(3.16) \quad \Delta_\eta v \pm |\nabla a(x_0)|v^p = 0 \quad \text{in } \mathbb{R}^N.$$

This should be explained in more detail. We have only to verify that

$$(3.17) \quad \sigma_j := \frac{g(M_j v_j)}{M_j^p} \rightarrow v^p.$$

At points x where $v(x) > 0$ this is clear. Suppose that $v(x) = 0$. Since the v_j are locally uniformly continuous, the set A where $v_j \rightarrow 0$ is closed and one easily verifies that $\sigma_j \rightarrow 0$ on A . Indeed, if $M_j v_j$ is bounded, then $\sigma_j \rightarrow 0$, while if, for a subsequence, $M_j v_j \rightarrow \infty$, then $\sigma_j \rightarrow v^p(x) = 0$. Hence, from (3.15) it follows that

$$\int_{\mathbb{R}^N} v \Delta \zeta \pm |\nabla a(x_0)|v^p \zeta = 0, \quad \forall \zeta \in C_0^\infty(\mathbb{R}^N),$$

which implies (3.16).

From (3.16) we find that $v \in C^2$. Furthermore, from (3.10),

$$v \geq 0, \quad \max v = v(0) = 1.$$

Because of the maximum principle, the minus sign cannot occur in (3.16), for $\nabla a(x_0) \neq 0$. Therefore, v satisfies

$$v \geq 0, \quad \Delta_\eta v + |\nabla a(x_0)|v^p = 0 \quad \text{in } \mathbb{R}^N.$$

By Theorem 1.2 in [6], this implies $v \equiv 0$, a contradiction with $v(0) = 1$.

Consider next the case

$$(ii) \quad \delta_j/\lambda_j \rightarrow 0 \quad \text{for a subsequence.}$$

Now let $j \rightarrow \infty$ in (3.13). As before, for a subsequence, $v_j \rightarrow v$. The limit equation is now

$$v \geq 0, \quad \Delta_y v + (\nabla a(x_0) \cdot y)v^p = 0 \quad \text{in } \mathbb{R}^N.$$

After a suitable rotation and rescaling, the equation reads

$$v \geq 0, \quad \Delta_\zeta v + \zeta_N v^p = 0 \quad \text{in } \mathbb{R}^N.$$

We now apply Theorem 4 and conclude that $v \equiv 0$, again a contradiction with $v(0) = 1$.

Finally, the case

$$(iii) \quad \delta_j/\lambda_j \rightarrow \delta_0 > 0 \quad \text{for some subsequence.}$$

The limit equation is now

$$(3.18) \quad v \geq 0, \quad \Delta_y v + (\pm \delta_0 |\nabla a(x_0)| + \nabla a(x_0) \cdot y)v^p = 0 \quad \text{in } \mathbb{R}^N.$$

The minus sign is impossible since at the maximum point $y = 0$, (3.18) would imply

$$\Delta_y v = \delta_0 |\nabla a(x_0)| > 0.$$

After a suitable change of variables and rescaling of v , (3.18) can be written as

$$v \geq 0, \quad \Delta_\zeta v + \zeta_N v^p = 0 \quad \text{in } \mathbb{R}^N,$$

with $v(0, \dots, 0, h) = 1$ for some $h \geq 0$.

Applying once more Theorem 4 we are led again to a contradiction.

Step 3. Consider now Case B, i.e. $x_0 \in \Omega^+ \cup \Omega^-$. This time choose

$$\lambda_j = M_j^{(1-p)/2}$$

in (3.11). Since $x_0 \in \Omega$, the same argument—Lemma 3.1 ensures that $\lambda_j^2 \rho_j / M_j \rightarrow 0$ —employed in the previous case leads to the limit equation

$$0 \leq v \leq 1, \quad \Delta_y v + a(x_0) v^p = 0 \quad \text{in all of } \mathbb{R}^N.$$

Since $v(0) = 1 = \max v$, the above and the maximum principle imply $a(x_0) > 0$.

Once more, by Theorem 1.2 in [6], $v \equiv 0$, contradicting $v(0) = 1$.

Step 4. We pass then to Case C, i.e. $x_0 \in \partial\Omega$. Here

$$d_j := \text{dist}(x_j, \partial\Omega) \rightarrow 0.$$

We need to consider two subcases:

- (a) $d_j M_j^{(p-1)/2} \rightarrow \infty$ for a subsequence,
- (b) $d_j M_j^{(p-1)/2} \rightarrow \delta_0 \geq 0$ for a subsequence.

Without loss of generality, we may assume that x_0 is the origin and that the exterior normal there is $-e_N = -(0, \dots, 0, 1)$. For all cases, we choose

$$\lambda_j = M_j^{(1-p)/2}.$$

In case (a), $d_j / \lambda_j \rightarrow \infty$. As before, letting $j \rightarrow \infty$, through a subsequence, in (3.11), we obtain a limit function v defined, then, in all of \mathbb{R}^N , and it satisfies

$$(3.19) \quad 0 \leq v \leq 1 = v(0), \quad \Delta_y v + a(x_0) v^p = 0 \quad \text{in } \mathbb{R}^N.$$

Since $a(x_0) \neq 0$, the desired contradiction follows as in Case B.

In case (b), for a subsequence, $d_j / \lambda_j \rightarrow \delta_0 \geq 0$. Going to the limit in (3.11) we now find that (3.19) holds in $\{y \in \mathbb{R}^N : y_N > -\delta_0\}$, and

$$(3.20a) \quad v = 0$$

or

$$(3.20b) \quad v_N = 0,$$

and, since now $\psi = \varphi$,

$$(5.3) \quad \alpha \int_{\Omega} a(x)\varphi^{q+1} < 0.$$

We have to prove that (5.1) has a solution when $\tau = 0$. For $\tau > 0$, (5.1) has a solution u_{τ} and $u_{\tau} \leq \bar{C}$. We have only to show that, as $\tau \rightarrow 0$, any positive solution stays away from the origin. Then, for a sequence $\tau_i \rightarrow 0$, the u_{τ_i} converge to a nonzero solution of (5.1) with $\tau = 0$. To show that u_{τ} stays away from zero we carry out a standard bifurcation analysis, using Lyapunov-Schmidt decomposition.

We show that for $0 < \tau$ small, there is no solution of (5.1) with small L^{∞} norm. Suppose there were such a solution u ; decompose it as a sum

$$(5.4) \quad u = t\varphi + v \quad \text{with} \quad \int_{\Omega} v\varphi = 0.$$

Then

$$(5.5) \quad (-\Delta + m)v = -\tau(t\varphi + v) + a(x)g(t\varphi + v)$$

and, necessarily, the right-hand side of (5.5) is L^2 -orthogonal to φ .

For the general operator L we would take v and the right-hand side of (5.5) orthogonal to ψ .

In the space of functions orthogonal to φ , the operator $-\Delta + m$ has a bounded inverse from L^p to $W^{2,p}$ for any p in $(1, \infty)$. It follows that

$$(5.6) \quad \|v\|_{W^{2,p}} \leq C\tau(|t| + \|v\|_{L^p}) + C\|g(t\varphi + v)\|_{L^{\infty}}.$$

If the L^{∞} norm of u is small, so are $|t|$ and the L^{∞} norm of v . Consequently, for τ small, by (5.6),

$$\|v\|_{W^{2,p}} \leq C\tau|t| + C(|t|^q + \|v\|_{L^{\infty}}^q).$$

Since $q > 1$ and $\|v\|_{L^{\infty}} \leq C\|v\|_{W^{2,p}}$, it follows that

$$(5.7) \quad \|v\|_{L^{\infty}} = |t|O(\tau + |t|^{q-1}).$$

From $u > 0$ it follows that $t \neq 0$, in fact $t > 0$. Next, we use the condition that the right-hand side of (5.5) is L^2 -orthogonal to φ ; we may now suppose $\int_{\Omega} \varphi^2 = 1$. Setting $v/t = w$ we find, on multiplying (5.5) by $\psi = \varphi$ and integrating,

$$(5.8) \quad \begin{aligned} t\tau &= \int_{\Omega} a(x)g(t\varphi + tw)\varphi \\ &= \alpha t^q \int_{\Omega} a(x)(\varphi + w)^q \varphi + o(t^q) \quad (\text{by (5.2)}) \\ &= \alpha t^q \int_{\Omega} a(x)\varphi^{q+1} + O(\tau t^q) + o(t^q) \quad (\text{by (5.7)}). \end{aligned}$$

Since $\tau, t > 0$, we see from (5.3) that this is impossible.

6. Proof of Theorem 2

Step 1. Using the same bifurcation analysis as above, we show first that for $0 < -\tau$ small, (5.1) has a solution $u = u_\tau$ which is close to 0 (we may assume $g(s) = 0$ for $s < 0$).

This follows from

LEMMA 6.1. *Under the conditions of Theorem 2, i.e. (1.3) and (5.2), (5.3), there is an interval $I = [0, t_0)$ and a continuous function $\delta(t)$ on I , with (recall $\varphi = \psi$)*

$$\delta(0) = \delta_0 := \alpha \int_{\Omega} a\varphi^{q+1} < 0,$$

and a C^1 function $u(t, x)$ for $t \in I$, $x \in \bar{\Omega}$, which for $0 < t < t_0$ is a solution of (5.1) with $\tau = t^{q-1}\delta(t)$.

PROOF. In (5.1) write u as in (5.4). We carry out the standard Lyapunov-Schmidt procedure. Consider the Banach spaces $X = \{u \in W^{2,p}(\Omega) : Bu = 0\}$ and $Y = L^p(\Omega)$, for some fixed $p > N$. Let P be the L^2 -orthogonal projection onto the subspace Y_1 of Y consisting of functions orthogonal to φ . We decompose (5.1) into two pieces which are to be solved for $v \in X$ and $\tau \leq 0$, depending on t :

$$(6.1) \quad (-\Delta + m)v = -\tau v + P[ag(t\varphi + v)]$$

$$(6.2) \quad t\tau = \int_{\Omega} ag(t\varphi + v)\varphi.$$

The right-hand side of (6.1) lies in $Y_1 = PY$ and it is of class C^1 in v, τ and t . With the aid of the Implicit Function Theorem we may solve (6.1) uniquely for $v = v(t, \tau)$ with $v(0, 0) = 0$. For t in some small interval I , and $|\tau|$ small, v belongs to C^1 . Furthermore, (5.6) holds and the derivatives v_t and v_τ vanish at $(0, 0)$. In fact, we find from (6.1), after taking the τ derivative,

$$(-\Delta + m)v_\tau = -v - \tau v_\tau + P[ag'(t\varphi + v)v_\tau].$$

Using (5.6) it follows that

$$(6.3) \quad \|v_\tau\|_{W^{2,p}} \leq C(t|\tau| + t^q).$$

We now substitute $v(t, \tau)$ in (6.2) and obtain

$$(6.3) \quad t\tau = \int_{\Omega} ag(t\varphi + v(t, \tau))\varphi.$$

This we solve for τ in the form $\tau = \delta t^{q-1}$, i.e. we solve for δ ,

$$(6.4) \quad \delta = \int_{\Omega} a\varphi \frac{g(t\varphi + v(t, \delta t^{q-1}))}{t^q} =: G(t, \delta).$$

As a function of $t \in I$ and of δ , $G(0, \delta) = \delta_0$. Furthermore, G is continuous in t and of class C^1 in δ . Indeed,

$$\frac{\partial}{\partial \delta} \frac{g(t\varphi + v(t, \delta t^{q-1}))}{t^q} = \frac{g'(t\varphi + v(t, \delta t^{q-1}))v_\tau}{t}.$$

By (6.3) we see that $G_\delta(t, \delta) = O(t^{q-1})$. We may therefore use the Implicit Function Theorem again and solve (6.4) for $\delta(t)$ with $\delta(0) = \delta_0$. The function $\delta(t)$ is continuous on I (possibly shortened). Then $u = t\varphi + v(t, t^{q-1}\delta(t))$ is a solution of (5.1) with $\tau = t^{q-1}\delta(t)$. The set of such τ covers a small interval $(\tau_0, 0)$.

Step 2. Completion of the proof of Theorem 2.

In Lemma 1 we showed that for τ large negative, problem (5.1) has no solution. The last assertion of Theorem 2 then follows from Lemma 6.1 and the following:

LEMMA 6.2. *If (5.1) has a solution u_{τ_1} for some $\tau_1 < 0$, then it has a solution for every τ in $(\tau_1, 0)$.*

PROOF. For every given τ in $(\tau_1, 0)$ consider problem (5.1). The function $\bar{u} = u_{\tau_1}$ is then a supersolution for (5.1). On the other hand, for $0 < \varepsilon$ small, the function $\underline{u} = \varepsilon\varphi$ is a subsolution. Furthermore, for ε small, $\underline{u} < \bar{u}$. This is clear for the Dirichlet problem, with the aid of the Hopf lemma. For B given by (1.2b), it holds because \bar{u} and φ are positive in $\bar{\Omega}$.

By the well known theory of sub- and supersolutions (see for example [2]), there is a solution $u = u_\tau$ of (5.1) with $\underline{u} < u < \bar{u}$.

Theorem 2 is proved.

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