

HITTING TIMES FOR TRANSIENT STABLE PROCESSES

S. C. PORT

In this paper we explicitly find the asymptotic behavior, for large t , of the probability that a transient d -dimensional stable process first (last) hits a bounded Borel set during the time interval (t, ∞) .

Assume that $X(t)$ is a stable process on R^d (d -dimensional Euclidean space) having exponent $\alpha < d$ and normalized so that the paths are right continuous with left-hand limits at every point. Assume further that $[X(t) - X(0)]t^{-1/\alpha}$ is distributed like $X(1) - X(0)$, and moreover, that $X(1) - X(0)$ has a genuinely d -dimensional distribution on R^d . [In particular, every symmetric stable process on R^d with 0 mean (when it exists) satisfies these conditions.]

From these assumptions it follows that $X(t) - X(0)$ has a bounded, continuous density, $f(t, x)$, which satisfies the well-known scaling property

$$(1.1) \quad f(t, x) = t^{-d/\alpha} f(1, t^{-1/\alpha} x).$$

For a Borel (more generally, analytic) set $B \subset R^d$, let

$$V_B = \inf \{t > 0: X(t) \in B\}$$

denote the first hitting time of B . As usual we set $V_B = \infty$ if

$$X(t) \notin B$$

for all $t > 0$. Our main purpose in this note is to establish the following.

THEOREM 1. *Let B be a bounded Borel (or analytic) subset of R^d . Then under the above assumptions on $X(t)$,*

$$(1.2) \quad \lim_{t \rightarrow \infty} t^{(d/\alpha)-1} P_x(t < V_B < \infty) = P_x(V_B = \infty) C(B) \left[\frac{d}{\alpha} - 1 \right]^{-1} f(1, 0),$$

where $C(B)$ is the natural capacity of B .

Previously, (by using a different method) Joffe [2] established this result for symmetric processes with $(d/2) < \alpha < 1$ when B has a non-empty interior, and Spitzer [4] (Lemma, p. 114) established this result for arbitrary compact B in the case of 3-dimensional Brownian motion. In the case of recurrent stable processes the analogue of Theorem 1 can be found in [3].

It is interesting to compare Theorem 1 with the following, much easier

THEOREM 2. *Let*

$$T_B = \inf\{t \geq 0: X(s) \notin B, \text{ all } s > t\}$$

be the last hitting time of B. Then under the same conditions as Theorem 1,

$$(1.3) \quad \lim_{t \rightarrow \infty} t^{d/\alpha-1} P_x(T_B > t) = C(B) \left[\frac{d}{\alpha} - 1 \right]^{-1} f(1, 0).$$

2 Proofs.

Proof of Theorem 1. A first passage decomposition yields

$$(2.1) \quad P_x(t < V_B < \infty) = \int_{\mathbb{R}^d} P_x(V_B > t, X(t) \in dy) P_y(V_B < \infty) \\ = \int_{\mathbb{R}^d} \left[f(t, y - x) - \int_0^t \int_{\bar{B}} H_B(x, ds, dz) f(t - s, y - z) \right] P_y(V_B < \infty) dy,$$

where here and in the following,

$$H_B(x, ds, dz) = P_x(V_B \in ds, X(s) \in dz),$$

and \bar{B} is the closure of B . But it is a known fact ([1] Prop. 18.4) that there is a measure, $e_B(dy)$, with support contained in \bar{B} (the capacity measure of B) and finite total mass $C(B)$ (the capacity of B), such that

$$(2.2) \quad P_y(V_B < \infty) = \int_{\bar{B}} g(u - y) e_B(du),$$

where

$$g(x) = \int_0^\infty f(t, x) dt$$

is the potential kernel density for the process $X(t)$. Setting

$$R(t, x) = \int_t^\infty f(s, x) ds$$

and using the fact that

$$(2.3) \quad \int_{\mathbb{R}^d} f(t, y - x) g(u - y) dy = R(t, u - x),$$

we obtain from (2.1) that

$$(2.4) \quad P_x(t < V_B < \infty) \\ \int_{\bar{B}} \left[R(t, y - x) - \int_{\bar{B}} \int_0^t H_B(x, ds, dz) R(t - s, y - z) \right] e_B(dy) .$$

From the scaling property (1.1) and the fact that $f(1, x)$ is continuous, we see that $\lim_{t \rightarrow \infty} t^{d/\alpha} f(t, x) = f(1, 0)$, uniformly in x on compacts, and thus

$$(2.5) \quad \lim_{t \rightarrow \infty} t^{(d/\alpha)-1} R(t, x) = f(1, 0) \left[\frac{d}{\alpha} - 1 \right]^{-1} ,$$

uniformly in x on compacts. Set

$$R(t) = t^{-(d/\alpha)+1} \left[\frac{d}{\alpha} - 1 \right]^{-1} .$$

Then from (2.5),

$$(2.6) \quad \lim_{t \rightarrow \infty} \int_{\bar{B}} \frac{R(t, y - x)}{R(t)} e_B(dy) = f(1, 0) C(B) ,$$

and

$$(2.7) \quad \lim_{T \rightarrow \infty} \lim_{t \rightarrow \infty} \int_0^T \left[\int_{\bar{B}} \int_{\bar{B}} H_B(x, ds, dz) R(t - s, y - z) e_B(dy) \right] R(t)^{-1} \\ = \lim_{T \rightarrow \infty} \int_0^T H_B(x, ds, \bar{B}) C(B) f(1, 0) = P_x(V_B < \infty) C(B) f(1, 0) .$$

From (2.4), we see that in order to complete the proof it suffices to show

$$(2.8) \quad \lim_{T \rightarrow \infty} \limsup_{t \rightarrow \infty} R(t)^{-1} \int_T^t \int_{\bar{B}} H_B(x, ds, dz) R(t - s, y - z) e_B(dy) = 0 .$$

To accomplish this, decompose \int_T^t as $\int_T^{t/2} + \int_{t/2}^{t-T} + \int_{t-T}^t$. Since

$$\sup_x f(1, x) = K < \infty ,$$

it follows from the scaling property that $R(t, x) \leq KR(t)$ for all $t > 0$. Setting $A = KC(B)$, we obtain

$$\int_T^{t/2} \leq A \int_T^{t/2} P_x(V_B \in ds) R(t - s) \leq AR(t/2) P_x(T < V_B < \infty) ,$$

and thus

$$\lim_{T \rightarrow \infty} \limsup_t R(t)^{-1} \int_T^{t/2} = 0 .$$

Next observe that

$$\int_{t/2}^{t-T} \leq A \int_{t/2}^{t-T} P_x(V_B \in ds)R(t-s) \leq AR(T)P_x(t/2 < V_B < \infty) .$$

By (2.4) this last term is dominated by $A^2R(T)R(t/2)$, and thus

$$\lim_T \limsup_t R(t)^{-1} \int_{t/2}^{t-T} = 0 .$$

Finally, from (2.2) we see that

$$\int_{t-T}^t \leq \int_{t-T}^t \int_{\bar{B}} H_B(x, ds, dz) \int_{\bar{B}} g(y-z)e_B(dy) \leq \int_{t-T}^t P_x(V_B \in ds) .$$

But

$$\begin{aligned} P_x(t-T < V_B \leq t) &= \int_{R^d} P_x(V_B > t-T, X(t-T) \in dy)P_y(V_B \leq T) \\ &\leq \int_{R^d} f(t-T, y-x)P_y(V_B \leq T)dy \leq K(t-T)^{-d/\alpha} \int_{R^d} P_y(V_B \leq T)dy . \end{aligned}$$

Since the paths $X(t)$ are bounded a.s. on $[0, T]$, we see that for each T there is a sphere $S_T \supset \bar{B}$, such that $P_y(X(t) \in S_T) \geq 1/2$ for all $t \leq T$ and $y \in \bar{B}$. But then

$$\begin{aligned} |S_T| &= \int_{R^d} P_x(X(T) \in S_T)dx \geq \int_{R^d} dx \int_0^T \int_{\bar{B}} H_B(x, ds, dy)P_y(X(T-s) \in S_T) \\ &\geq \frac{1}{2} \int_{R^d} P_x(V_B \leq T)dx . \end{aligned}$$

Thus

$$\lim_{t \rightarrow \infty} R(t)^{-1} \int_{t-T}^t = 0 .$$

This completes the proof.

Proof of Theorem 2. Clearly

$$P_x(T_B > t) = \int_{R^d} f(t, y-x)P_y(V_B < \infty)dy .$$

Using (2.2) and (2.3) we see that

$$P_x(T_B > t) = \int_{\bar{B}} R(t, y-x)e_B(dy) ,$$

from which the theorem follows.

REMARK. When $d/2 < \alpha < d$, it is possible to establish Theorem 1 by a much simpler argument. Set

$$Q_B^\lambda(x) = \int_0^\infty e^{-\lambda t} P_x(t < V_B < \infty) dt ,$$

$$H_B^\lambda(x, dy) = \int_0^\infty e^{-\lambda t} P_x(V_B \in dt, x(t) \in dy)$$

and

$$R^\lambda(x) = \int_0^\infty R(t, x) e^{-\lambda t} dt .$$

Then from (2.4) we obtain

$$(2.9) \quad Q_B^\lambda(x) = \int_{\bar{B}} \left[R^\lambda(y - x) - \int_{\bar{B}} H_B^\lambda(x, dz) R^\lambda(y - z) \right] e_B(dy) .$$

It follows from (2.5) that uniformly in x on compacts,

$$\lim_{\lambda \downarrow 0} R^\lambda(x) \lambda^{2-d/\alpha} = f(1, 0) \left[\frac{d}{\alpha} - 1 \right]^{-1} \Gamma(2 - d/\alpha) .$$

Consequently, from (2.9), we see that

$$\lim_{\lambda \downarrow 0} Q_B^\lambda(x) \lambda^{2-d/\alpha} = f(1, 0) C(B) P_x(V_B = \infty) \left[\frac{d}{\alpha} - 1 \right]^{-1} \Gamma(2 - d/\alpha) .$$

An appeal to Karamata's theorem, and the fact that $P_x(t < V_B < \infty)$ is monotone in t , then yields (1.2).

The above argument breaks down when $\alpha < d/2$ since

$$\lim_{\lambda \downarrow 0} R^\lambda(x) < \infty ,$$

and the more complicated proof given previously is needed.

REFERENCES

1. G. A. Hunt, *Markoff processes and potentials III*, Illinois J. Math. **2** (1958), 151-213.
2. A. Joffe, *Sojourn Time for Stable Processes*, Thesis, Cornell U., 1959.
3. S. C. Port, *Hitting times and potentials for recurrent stable processes*, J. D'Analyse Mathématique (to appear)
4. F. Spitzer, *Electrostatic capacity, heat flow, and Brownian motion*, Z. Warscheinlichk **3**, (1964).

Received June 30, 1966. This research is sponsored by the United States Air Force under Project RAND—Contract No. AF 49 (638)–1700 monitored by the Directorate of Development Plans, Deputy Chief of Staff, Research and Development, Hq USAF.

THE RAND CORPORATION

UNIVERSITY OF CALIFORNIA, LOS ANGELES

