

BOUNDARY BEHAVIOR OF LIMITS OF DISCRETE SERIES REPRESENTATIONS OF REAL RANK ONE SEMISIMPLE GROUPS

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The decomposition of the reducible unitary principal series of a connected semisimple Lie group having real rank one and a simply connected complexification is exhibited on a global analytic level in such a way that it is seen to correspond to a phenomenon in classical Fourier analysis. This is done by embedding limits of discrete series representations via a group equivariant passage to boundary values analogous to the classical Hardy space inclusion used by Bargmann in the case of $SL(2, \mathbf{R})$. The boundary value map is shown to be a factor of the projection operator given by the Knapp-Stein intertwining operator. From a representation theoretic view, while these decompositions are already known, the method of computing the leading term of the asymptotic expansion of matrix coefficients is new and does not require a K -finiteness assumption.

1. Introduction and preliminaries. The decomposition of representations in the unitary principal series of a connected semisimple Lie group G having real rank one and a simply connected complexification is well understood [10], [11]. In particular, Knapp and Wallach having used Szegő kernels to decompose all reducible unitary principal series representations as sums of limits of discrete series representations [11, §12]. In this paper we exhibit these reducibility results on a global analytic level by explicitly embedding limits of discrete series representations in the reducible principal series. This is achieved by realizing the representations in question in suitable function spaces and providing a group equivariant passage to boundary values analogous to the Hardy space inclusion of $H^2(\mathbf{R})$ in $L^2(\mathbf{R})$ that was used by Bargmann in the case $G = SL(2, \mathbf{R})$ [1] and Knapp and Okamoto [9] more generally in the case of limits of holomorphic discrete series.

Throughout this paper we assume that G satisfies the properties listed above. Furthermore, from the point of view of exhibiting reducibility results, there is no loss of generality in assuming that G has a compact Cartan subgroup $T \subseteq K$ where K is a maximal compact subgroup of G corresponding to a Cartan involution θ [10, p. 543–544]. Then G has discrete series $\mathcal{E}^2(G)$ [5]. To each nonsingular integral form Λ on the Lie algebra \mathfrak{t} of T , Harish-Chandra associates an invariant eigendistribution

Θ_Λ [4, Theorem 2] and proves the existence of a discrete series representation (π_Λ, H^Λ) with character Θ_Λ [5] and that these representations exhaust $\mathcal{E}^2(G)$; we call Λ a Harish-Chandra parameter.

Let \mathfrak{g} and \mathfrak{k} denote the Lie algebras of G and K , and let Δ (respectively Δ_k, Δ_n) denote the roots (respectively compact roots, noncompact roots) of $(\mathfrak{g}^{\mathbb{C}}, \mathfrak{k}^{\mathbb{C}})$. Normalize root vectors E_α ($\alpha \in \Delta$) according to [6, 155–156]. If Λ is a Harish-Chandra parameter we order Δ so that Λ is Δ^+ -dominant; Δ^+ is thereby uniquely determined. If instead the integral parameter Λ is singular, but not orthogonal to any compact root, it is easy to see that there is a noncompact root α for which $\{\pm\alpha\}$ is precisely the set of roots orthogonal to Λ [11, Lemma 12.5]. For such a parameter, called here a limit Harish-Chandra parameter, there are two possible choices of positive roots Δ^+ for which Λ is Δ^+ -dominant. Whichever the choice, the unique positive root orthogonal to Λ is noncompact and simple [11, Lemma 12.5].

Let Λ be either a Harish-Chandra parameter or a limit Harish-Chandra parameter. Order Δ so that Λ is Δ^+ -dominant and put $\delta = \frac{1}{2}\sum_{\alpha \in \Delta^+} \alpha$, $\delta_k = \frac{1}{2}\sum_{\alpha \in \Delta_k^+} \alpha$, and $\delta_n = \delta - \delta_k$. Let α_0 be any simple noncompact root if Λ is nonsingular and the unique positive root orthogonal to Λ (hence also simple noncompact) if Λ is singular. Then α_0 is a fundamental sequence of positive noncompact roots in the sense of [11, §4] and α_0 determines an Iwasawa decomposition $G = ANK$ with the Lie algebra \mathfrak{a} of A given by $\mathfrak{a} = \mathbf{R} \cdot (E_{\alpha_0} + E_{-\alpha_0})$ and $E_{\alpha_0} + E_{-\alpha_0}$ in the positive chamber of \mathfrak{a} . Observe that if Λ is singular, the Iwasawa decomposition does not depend on which of the two possible systems of positive roots Δ^+ that is used. Let M (respectively M') denote the centralizer (respectively normalizer) of A in K and denote by P the minimal parabolic subgroup MAN of G . Let $\lambda = \Lambda - \delta_k + \delta_n$ be the Blattner parameter corresponding to (Λ, Δ^+) . Thus, when Λ is nonsingular, $\lambda = \lambda(\Lambda)$ is the lowest K -type in π_Λ . Even when Λ is singular, λ is integral and Δ_k^+ -dominant [11, p. 198]; the Blattner parameter λ' corresponding to $(\Lambda, \Delta^{+'})$, where $\Delta^{+'} = (\Delta^+ - \{\alpha_0\}) \cup \{-\alpha_0\}$ is the other possible positive root system, is given by $\lambda' = \lambda - \alpha_0$. For μ integral and Δ_k^+ -dominant let (τ_μ, V_μ) denote an irreducible unitary representation of K .

A convenient realization of the discrete series representation π_Λ (Λ nonsingular) is one the space of square integrable functions in

$$(1.1) \quad C^\infty(G, \tau_\Lambda) \\ = \{F \in C^\infty(G, V_\lambda) \mid F(kg) = \tau_\lambda(k)F(g), k \in K, g \in G\}$$

that are annihilated by a certain first order elliptic differential operator \mathcal{D}_λ [13, 14] ($\lambda = \lambda(\Lambda)$).

For $\nu \in \text{Hom}_{\mathbf{R}}(\mathfrak{a}, \mathbf{C}) = \mathfrak{a}'_{\mathbf{C}}$ and (σ, H) a irreducible unitary representation of M , let $U(\sigma : \nu)$ denote the nonunitary principal series representation realized in the compact picture on $L^2(K, \sigma)$ (cf. §2). In [11] Knapp and Wallach associate to the parameter Λ (and the ordering Δ^+ if Λ is nonsingular) an irreducible unitary representation $(\sigma_\lambda, H_\lambda)$ of M with highest weight λ and $H_\lambda \subset V_\lambda$ ($\lambda = \lambda(\Lambda, \Delta^+)$), a parameter $\nu(\lambda)$ in $\mathfrak{a}'_{\mathbf{C}}$, and an integral formula S_λ defined on the dense subspace

$$(1.2) \quad C^\infty(K, \sigma_\lambda) = \{f \in C^\infty(K, H_\lambda) \mid f(mk) = \sigma_\lambda(m)f(k), m \in M, k \in K\}$$

of $L^2(K, \sigma_\lambda)$ by

$$(1.3) \quad S_\lambda f(x) = \int_K \tau_\lambda(k)^{-1} f(kx) dk \quad (x \in G).$$

The dependence on $\nu(\lambda)$ is incorporated into the extension of f to G required for formula (1.3) (cf. §2). The point is that S_λ carries $C^\infty(K, \sigma_\lambda)$ G -equivariantly into the kernel of \mathcal{D}_λ in $C^\infty(G, \tau_\lambda)$ and thus provides a quotient map of $U(\sigma_\lambda : \nu(\lambda))$ onto π_Λ when Λ is nonsingular. When Λ is singular, the two Blattner parameters λ and λ' give rise to equivalent M -types σ_λ and $\sigma_{\lambda'}$ and the formulas for both $\nu(\lambda)$ and $\nu(\lambda')$ reduce to ρ (cf. §2). The unitary principal series representations $U(\sigma_\lambda : \rho)$ and $U(\sigma_{\lambda'} : \rho)$ are therefore equivalent. Nevertheless, their images under the Szegő maps S_λ and $S_{\lambda'}$ respectively have lowest K -types λ and $\lambda' = \lambda - \alpha_0$ respectively and so are independent limits of discrete series representations. Knapp and Wallach showed that the unitary principal series representation $U(\sigma_\lambda : \rho)$ is infinitesimally equivalent with the direct sum of the K -finite images of $U(\sigma_\lambda : \rho)$ and $U(\sigma_{\lambda'} : \rho)$ under S_λ and $S_{\lambda'}$ respectively [11, Theorem 12.6]. The completeness result that all reducibility of the unitary principal series may be so accounted for is Theorem 12.7 of [11].

In this paper we will establish this decomposition in a global analytic fashion by means of a boundary value embedding \mathcal{L} carried out in §3. The point is that, although $\lim_{a \rightarrow \infty} S_\lambda f(a) = 0$ for f in $C^\infty(K, \sigma_\lambda)$, we can write $S_\lambda f(a) = c(f)e^{-\rho \log a} + \text{lower order terms}$, with $c(f) \neq 0$ in general. The boundary map is then defined by the constant term in the expansion of $e^{\rho \log a} S_\lambda f(a)$ after projecting by E_λ from V_λ onto H_λ :

$$(1.4) \quad \mathcal{L}(S_\lambda f)(k) = E_\lambda c(U(\sigma_\lambda : \rho : w^{-1}k)f)$$

where w is a certain representative of the nontrivial Weyl group element. The bulk of §3 is devoted to establishing the finite, generally non-zero

limit in (1.4). The main tools in this analysis were developed in [10], which we quote frequently. Particularly important for our purposes are the mean value property [10, Proposition 20]

$$(1.5) \quad \int_{c < |v| < d} \frac{\sigma_\lambda(vw)^{-1}}{|v|} dv = 0$$

and another result of Knapp-Stein which we include here as Lemma 3.1. Some consequences of the proof of this lemma, such as Proposition 3.3 and 3.4, may be of independent interest. The final limit result needed to define \mathcal{L} is given in Theorem 3.14. In Theorem 3.16 it is shown that \mathcal{L} maps the limit of discrete series representation with lowest K -type λ G -equivariantly into $U(\sigma_\lambda : \rho)$.

Embedding theorems for limits of discrete series for the classical real rank one groups were given in [12]. In addition to the greater generality of the present paper, the results given here may be of interest through their relationship with the Knapp-Stein intertwining operators. These were given in [10] where it is shown [10, p. 517] that in the noncompact picture $\mathcal{U}(\sigma_\lambda : \rho)$ the intertwining operators consist of linear combinations of the identity and the convolution operator with kernel $|v|^{-1}\sigma_\lambda(vw)^{-1}$. We show in Theorem 3.16 that the composition

$$\mathcal{U}(\sigma_\lambda : \rho) \xrightarrow{W^{-1}} U(\sigma_\lambda : \rho) \xrightarrow{S_\lambda} \text{Limit of Discrete Series} \xrightarrow{\mathcal{L}} U(\sigma_\lambda : \rho) \xrightarrow{W} \mathcal{U}(\sigma_\lambda : \rho)$$

(cf. §2 for the definition of W) is indeed of the type described.

Some of these results were announced in [2]. It is a pleasure to thank Professors A. W. Knapp and N. R. Wallach for their valuable suggestions.

2. The Szegő integral. Let Λ be a limit Harish-Chandra parameter. Order Δ so that Λ is Δ^+ -dominant and let α_0 , A , and α be as defined in §1. Let Φ denote the restricted roots of \mathfrak{g} with respect to \mathfrak{a} ; for $\gamma \in \Phi$ let \mathfrak{g}^γ denote the corresponding restricted root space and set $\gamma \in \Phi^+$ if $\gamma(E_{\alpha_0} + E_{-\alpha_0}) > 0$. Our assumption that G has real rank one results in Φ^+ having the form $\Phi^+ = \{\alpha\}$ or $\Phi^+ = \{\alpha, 2\alpha\}$. In our notation α will denote the smallest positive root. Let $p = \dim_{\mathbf{R}} \mathfrak{g}^\alpha$, $q = \dim_{\mathbf{R}} \mathfrak{g}^{2\alpha}$ where $\mathfrak{g}^{2\alpha} = (0)$ if $2\alpha \notin \Phi$, and let ρ denote half the sum of the positive restricted roots with multiplicity so that

$$(2.1) \quad \rho = \frac{1}{2}(p + 2q)\alpha.$$

Let $\mathfrak{n} = \sum_{\beta \in \Phi} \mathfrak{g}^\beta$ and let N and \bar{N} denote the analytic subgroups of G corresponding to \mathfrak{n} and $\theta\mathfrak{n}$.

Since $\dim_{\mathbb{R}} \alpha = 1$, the Weyl group $w = M'/M$ has order two. Let M' act in each equivalence class $[\sigma]$ in \hat{M} , the set of equivalence classes of irreducible unitary representations of M , by

$$(2.2) \quad w\sigma(m) = \sigma(w^{-1}mw) \quad (w \in M'; m \in M: [\sigma] \in \hat{M}).$$

By [8] we can choose a representation w of the nontrivial Weyl group element that centralizes M so that

$$(2.3) \quad w\sigma = \sigma.$$

We denote the factors of an element g in the Iwasawa decomposition $G = ANK$ by

$$(2.4) \quad g = \exp H(g) \cdot n\kappa(g) \quad (H(g) \in \alpha, \kappa(g) \in K)$$

and write $\log a$ for $H(a)$ when $a \in A$. Every element g not in the lower dimensional set Pw where $P = MAN$ also has a unique Gelfand-Naimark decomposition

$$(2.5) \quad g = m(g)a(g)n\bar{n}(g)$$

with factors in M, A, N , and \bar{N} respectively. By means of this decomposition we extend representations σ of M and characters χ of A to functions defined almost everywhere on G with respect to Haar measure:

$$\sigma(man\bar{n}) = \sigma(m), \quad \chi(man\bar{n}) = \chi(a)$$

where we adopt without further reference the lower case convention for group elements with the exception that v will always denote an element of \bar{N} . The Bruhat decomposition shows that for each g in G there is at most one v in \bar{N} for which $\bar{n}(vg)$ is not defined. If \bar{N}_g is this exceptional set, then $\bar{N}_w = \bar{N} - \{1\}$.

Let λ be the Blattner parameter associated to (Λ, Δ^+) as described in §1 and let $(\tau_\lambda, V_\lambda)$ denote a K -type with highest weight λ . Let ϕ_λ be a highest weight vector of τ_λ of length one, let H_λ be the M -cyclic subspace of V_λ generated by ϕ_λ , and let σ_λ be the representation of M given by τ_λ operating on H_λ . The proof of Proposition 5.5 of [11] and Lemma 12.3 of [11] show that $(\sigma_\lambda, H_\lambda)$ is an irreducible representation of M .

We recall from [3] that

$$(2.6) \quad \int_{\bar{N}} e^{(1+z)\rho H(v)} dv < \infty \quad \text{if } \operatorname{Re} z > 0$$

where dv is unimodular Haar measure on \bar{N} . We normalize Haar measures on M, K , and \bar{N} so that

$$(2.7) \quad \int_M dm = \int_K dk = \int_{\bar{N}} e^{2\rho H(v)} dv = 1.$$

We arrange parameters so that induction of $\sigma_\lambda \otimes e^\rho \otimes 1$ from MAN to G gives rise to a unitary representation. In the compact picture of this unitary principal series $U = U(\sigma_\lambda; \rho)$ the representation space is the closed subspace $L^2(K, \sigma_\lambda)$ of $L^2(K, H_\lambda)$ consisting of functions f such that for every m

$$(2.8) \quad f(mk) = \sigma_\lambda(m)f(k)$$

dk -almost everywhere in K . The action of G on $L^2(K, \sigma_\lambda)$ is

$$(2.9) \quad U(g)f(k) = e^{\rho H(kg)}f(\kappa(kg)).$$

We let $C^\infty(K, \sigma_\lambda)$ denote the space of smooth functions in $L^2(K, \sigma_\lambda)$. Then $C^\infty(K, \sigma_\lambda)$ is dense and is the space of C^∞ -vectors for U . If l belongs to K , it will be convenient to denote its action on f under U by ${}^l f$.

In the unitarily equivalent noncompact picture \mathcal{U} of U , the Hilbert space is $L^2(\bar{N}, H_\lambda)$ and the group action is given by

$$(2.10) \quad \mathcal{U}(g)F(v) = e^{\rho \log v g} \sigma(vg)F(\bar{n}(vg)).$$

The intertwining operator W between these two pictures is

$$(2.11) \quad Wf(v) = e^{\rho H(v)}f(\kappa(v)) \quad (v \in \bar{N}; f \in L^2(K, \sigma_\lambda)).$$

For f in $C^\infty(K, \sigma_\lambda)$ Knapp and Wallach define the Szegő map S_λ with parameter λ by

$$(2.12) \quad S_\lambda f(x) = \int_K e^{\rho H(lx^{-1})} \tau_\lambda(\kappa(lx^{-1}))^{-1} f(l) dl \quad (x \in G).$$

Extending f in $C^\infty(K, \sigma_\lambda)$ to G by $\tilde{f}(g) = e^{\rho H(g)}f(\kappa(g))$ so that $\tilde{f}(manx) = e^{\rho \log a} \sigma_\lambda(m) \tilde{f}(x)$ and \tilde{f} is in the induced picture of U , we have [11, p. 178]

$$(2.13) \quad S_\lambda \tilde{f}(x) = \int_K \tau_\lambda(k)^{-1} \tilde{f}(kx) dk \quad (x \in G),$$

exhibiting the G -equivariance of the Szegő map into the space $C^\infty(G, \tau_\lambda)$. It is shown in [11] that the image of $C^\infty(K, \sigma_\lambda)$ under S_λ is in the kernel of \mathcal{D}_λ in $C^\infty(G, \tau_\lambda)$ and that infinitesimally the K -finite image of S_λ is a direct summand of $U(\sigma_\lambda; \rho)$.

We will need another integral formula for the operator S_λ , one that will be of use in conjunction with the noncompact picture $\mathcal{U}(\sigma_\lambda; \rho)$.

LEMMA 2.1. *Let f belong to $C^\infty(K, \sigma_\lambda)$ and let S_λ be defined by (2.12) or the equivalent formula (2.13). Then*

$$(2.14) \quad S_\lambda f(a) = \int_{\bar{N}} e^{\rho H(va)} \tau_\lambda(\kappa(va)w)^{-1} (W^w f)(v) dv.$$

Proof. We use the integral formula

$$(2.15) \quad \int_K \varphi(k) dk = \int_{\bar{N}} \int_M \varphi(m\kappa(v)) e^{2\rho H(v)} dm dv$$

of Harish-Chandra [3, p. 287]. Thus, since $wa^{-1}w^{-1} = a$, $H(m\kappa(v)a) = -H(v) + H(va)$, and $\kappa(m\kappa(v)a) = m\kappa(va)$, we have

$$\begin{aligned} S_\lambda f(a) &= \int_K e^{\rho H(la^{-1})} \tau_\lambda(\kappa(la^{-1}))^{-1} f(l) dl \\ &= \int_K e^{\rho H(la)} \tau_\lambda(\kappa(la)w)^{-1} f(lw) dl \\ &= \int_{\bar{N}} \int_M e^{\rho H(m\kappa(v)a)} \tau_\lambda(\kappa(m\kappa(v)a)w)^{-1} {}^w f(m\kappa(v)) e^{2\rho H(v)} dm dv \\ &= \int_{\bar{N}} \int_M e^{\rho H(v)} e^{\rho H(va)} \tau_\lambda(\kappa(va)w)^{-1} \tau_\lambda(m)^{-1} \sigma_\lambda(m) {}^w f(\kappa(v)) dm dv \\ &= \int_{\bar{N}} e^{\rho H(va)} \tau_\lambda(\kappa(va)w)^{-1} e^{\rho H(v)} {}^w f(\kappa(v)) dv. \quad \square \end{aligned}$$

In view of the G -equivariance of S_λ this formula can be used globally on G via the Cartan decomposition:

$$(2.16) \quad S_\lambda(f:kak) = \tau_\lambda(k) S_\lambda({}^k f:a).$$

Let $\mathcal{S}_\lambda: L^2(\bar{N}, H_\lambda) \rightarrow C^\infty(A, V_\lambda)$ be defined by

$$(2.17) \quad \mathcal{S}_\lambda F(a) = e^{\rho \log a} \int_{\bar{N}} e^{\rho H(va)} \tau_\lambda(\kappa(va)w)^{-1} F(v) dv.$$

By abuse of notation we define \mathcal{S}_λ on $C^\infty(K, \sigma_\lambda)$ by

$$(2.18) \quad \mathcal{S}_\lambda f(a) = e^{\rho \log a} S_\lambda f(a);$$

then we have by Lemma 2.1

$$(2.19) \quad \mathcal{S}_\lambda f = \mathcal{S}_\lambda(W {}^w f).$$

3. Boundary values of Szegő integrals. The group A acts on \bar{N} by the dilations δ_a where

$$(3.1) \quad \delta_a v = a^{-1}va \quad (v \in \bar{N})$$

with change of variables given by

$$(3.2) \quad d(\delta_a v) = e^{2\rho \log a} dv.$$

The homogeneous norm $|v|$ on \bar{N} [10, p. 512] given by

$$(3.3) \quad |v| = e^{-\rho \log(vw)} \quad (v \in \bar{N}_w = \bar{N} - \{1\})$$

is α -homogeneous of degree $p + 2q$ and invariant under conjugation by M , i.e.,

$$(3.4) \quad |\delta_a v| = e^{2\rho \log a} |v| \quad \text{and} \quad |mvm^{-1}| = |v| \quad (v \neq 1).$$

The function $v \rightarrow \sigma_\lambda(vw)$ is of class C^∞ away from $v = 1$ and has the homogeneity property

$$(3.5) \quad \sigma_\lambda(\delta_a v \cdot w) = \sigma_\lambda(vw).$$

These facts may be found in [10, §6 and §8]. The following lemma may also be found in [10] but we provide an outline of its proof because we will need several consequences of the proof not found in [10].

LEMMA 3.1. ([10, Lemma 29].) *The map $v \rightarrow |v|^2$ is a polynomial on \bar{N} that is α -homogeneous of degree less than $2(p + 2q)$ such that*

$$(3.6) \quad e^{-2\rho H(v)} = 1 + P_1(v) + \cdots + P_s(v) + |v|^2.$$

Consequently $e^{2\rho H(v)} \leq 1$ and

$$(3.7) \quad e^{2\rho H(v)} \leq \frac{1}{|v|^2} \quad (v \neq 1).$$

Proof (sketch). Let π be a finite dimensional irreducible representation with α -weights $\rho = \mu_0, \mu_1, \dots, \mu_{s+1}$ of which ρ is the highest and such that the compact real form $\mathfrak{k} \oplus i\mathfrak{p}$ of \mathfrak{g}^C acts by skew-Hermitian transformations. If ϕ_ρ is a highest α -weight vector of length one then $\|\pi(g)^{-1}\phi_\rho\|^2 = e^{-2\rho H(g)}$. Let E_{μ_i} denote the orthogonal projection onto the weight space belonging to μ_i and put $P_j(g) = \|E_{\mu_j}\pi(g)^{-1}\phi_\rho\|^2$. Then

$$(3.8) \quad e^{-2\rho H(g)} = P_0(g) + \cdots + P_{s+1}(g) \quad (g \in G).$$

Routine computation shows that when gw belongs to $MAN\bar{N}$

$$(3.9) \quad \|E_{w\rho}\pi(g)^{-1}\phi_\rho\|^2 = e^{-\rho \log a(gw)}$$

and in particular when $v \neq 1$

$$(3.10) \quad \|E_{w\rho}\pi(v)^{-1}\phi_\rho\|^2 = e^{-2\rho \log(vw)}.$$

Since $P_{s+1}(v) = e^{-2\rho \log(vw)} = |v|^2$, all statements follow from (3.8) and (3.10). \square

We define the kernel $K(v : a)$ on $\bar{N} \times A$ by

$$(3.11) \quad K(v : a) = e^{2\rho \log a} e^{\rho H(\delta_a v)}.$$

Then for f in $C^\infty(K, \sigma_\lambda)$

$$\mathcal{S}_\lambda f(a) = \int_{\bar{N}} K(v : a) \tau_\lambda(\kappa(\delta_a v)w)^{-1} W^w f(v) dv.$$

The notation $a \rightarrow \infty$ will signify $a = \exp t(E_{\alpha_0} + E_{-\alpha_0})$ with $t \rightarrow \infty$.

LEMMA 3.2. For v in \bar{N} different from 1 we have

$$(3.12) \quad K(v : a) \leq \frac{1}{|v|},$$

$$(3.13) \quad \lim_{a \rightarrow \infty} K(v : a) = \frac{1}{|v|}, \quad \text{and}$$

$$(3.14) \quad \int K(v : a) dv = \int K(\delta_a v) d(\delta_a v) \quad \text{where } K(v) = K(v : 1).$$

Proof. Statement (3.12) follows immediately from (3.7) and (3.4) as does (3.14) from (3.2). Statement (3.13) is a simple consequence of (3.6). \square

PROPOSITION 3.3. The map $\mu : g \rightarrow e^{-2\rho \log a(gw)}$ which is defined on $MAN\bar{N}w$ can be continuously extended to G by putting $\mu(man) = 0$.

Proof. By (3.9), $\mu(g) = \|E_{w\rho}\pi(g)^{-1}\phi_\rho\|^2$ when g belongs to $MAN\bar{N}w$. Since M preserves the highest α -weight space of π , A acts by scalars, and N acts trivially, $\|E_{w\rho}\pi(man)^{-1}\phi_\rho\|^2 = 0$, the result follows because $G = P \cup P\bar{N}w$ (essentially the Bruhat decomposition of G). \square

The significance of Proposition 3.3 is that $\mu|_K$ is zero precisely on M and so provides a means of testing when an element of K belongs to M .

PROPOSITION 3.4. For v in \bar{N} different from 1, $\lim_{a \rightarrow \infty} \kappa(\delta_a v)w = m(vw)$.

Proof. For every a in A and v in \bar{N} , writing $\delta_a v$ as $e^{H(\delta_a v)}n\kappa(\delta_a v)$ we get $\kappa(\delta_a v)w^2 = w^2 e^{-H(\delta_a v)}n'v'$ belongs to $MAN\bar{N}$, since w^2 is in M , and so, in particular,

$$\mu(\kappa(\delta_a v)w) = e^{-2\rho \log a(\kappa(\delta_a v)w^2)} = e^{2\rho H(\delta_a v)}.$$

Thus, for $v \neq 1$,

$$\lim_{a \rightarrow \infty} \mu(\kappa(\delta_a v)w) = \lim_{a \rightarrow \infty} e^{-2\rho \log a} K(v : a) = 0$$

by (3.12) and so $\kappa(\delta_a v)w$ tends to M as $a \rightarrow \infty$. Since the Gelfand-Naimark decomposition is continuous on $MAN\bar{N}$, the A, N , and \bar{N} components of $\kappa(\delta_a v)w$ for $v \neq 1$ each converge to 1 as $a \rightarrow \infty$. The result follows since $m(\kappa(\delta_a v)w) = m(\delta_a v w) = m(vw)$ [10, formula (6.12)]. \square

Some remarks are appropriate before proceeding to the next sequence of lemmas. If we define $\mathcal{S}_\lambda f(a)$ by (2.18), then by (2.14) and (3.11)

$$(3.15) \quad \mathcal{S}_\lambda f(a) = \int_{\bar{N}} K(v : a) \tau_\lambda(\kappa(\delta_a v)w)^{-1} W^w f(v) dv$$

or

$$(3.16) \quad \mathcal{S}_\lambda f(a) = \int_{\bar{N}} e^{\rho H(v)} K(v : a) \tau_\lambda(\pi(\delta_a v)w)^{-1} f(\kappa(v)w) dv.$$

We have now shown that the integrand converges pointwise, except for $v = 1$, to $(\sigma_\lambda(vw)^{-1}/|v|)W^w f(v)$, but since $|v|^{-1}$ just fails to be integrable, the dominated convergence theorem is not applicable. Instead, we obtain more precise information about the rate at which $\kappa(\delta_a v)w$ approaches $m(vw)$ in the case, essentially, of $SU(2, 1)$ to which the general solution can be reduced. Positive constants that appear in these lemmas depend in an essential way only on the subscripted objects and may change from line to line. Let B denote the Killing form on \mathfrak{g} and let B_θ denote the positive definite norm on \mathfrak{g} given by $B_\theta = B \circ (1 \times -\theta)$. The associated norm on \mathfrak{g} will be denoted by $\|\cdot\|$.

LEMMA 3.5. *Suppose Y and Z are nonzero elements of $\mathfrak{g}^{-\alpha}$ and $\mathfrak{g}^{-2\alpha}$ respectively. Then*

$$(3.17) \quad [Z, \theta Z] \in \mathfrak{a},$$

$$(3.18) \quad [Y, [Y, \theta Z]] \text{ is a nonzero element of } \mathfrak{m},$$

$$(3.19) \quad [Z, [Z, \theta Z]] = |2\alpha|^2 \|Z\|^2 Z, \text{ and}$$

$$(3.20) \quad (\text{ad } Y)^4 \theta Z = c_{Y,Z} Z \text{ where } c_0 \|Y\|^4 \leq c_{Y,Z} \leq c_1 \|Y\|^4$$

for two positive constants c_0 and c_1 .

Proof. The first statement is obvious, the second can be found in [7, Lemma 1.8], and the third is an immediate computation. For (3.20), observe that

$$\begin{aligned} B_\theta((\text{ad } Y)^4 \theta Z, Z) &= -B([Y, [Y, \theta Z]], [Y, [Y, \theta Z]]) \\ &= \|[Y, [Y, \theta Z]]\|^2, \end{aligned}$$

since by (3.18) $[Y, [Y, \theta Z]]$ belongs to \mathfrak{k} and so is θ -invariant. \square

Let $d(k, k')$ denote a translation invariant metric on K . For X in \mathfrak{n} let $k_X = \exp(X + \theta X)$.

LEMMA 3.6. *There exists a neighborhood I of 0 in \mathfrak{n} and a positive number ε_0 such that*

$$(3.21) \quad d(k_X, 1) \leq C_I e^{-\varepsilon_0 \rho \log a(k_X w)} \quad (X \in I).$$

Proof. We give the proof for the case where $\Phi^+ = \{\alpha, 2\alpha\}$, it being the more difficult. The modifications necessary when $\Phi^+ = \{\alpha\}$ are evident in the proof. For X in \mathfrak{n} write $\theta X = Y + Z$ with Y in $\mathbf{R} \cdot X_{-\alpha}$ and Z in $\mathbf{R} \cdot X_{-2\alpha}$ where $X_{-\alpha}$ and $X_{-2\alpha}$ are nonzero vectors in $\mathfrak{g}^{-\alpha}$ and $\mathfrak{g}^{-2\alpha}$ respectively. Let \mathfrak{g}_X be the Lie subalgebra of \mathfrak{g} generated by $X_{-\alpha}$, $X_{-2\alpha}$, $\theta X_{-\alpha}$, and $\theta X_{-2\alpha}$ and let G_X be the analytic subgroup of G corresponding to \mathfrak{g}_X . Then \mathfrak{g}_X is isomorphic to $\mathfrak{su}(2, 1)$ [7, p. 54] and direct computation shows that the analogue ρ_X of ρ for \mathfrak{g}_X is given by $\rho_X = 2\alpha$. Let π_X be the representation of G_X constructed in the proof of Lemma 3.1 so that in fact π_X acts on $(\mathfrak{g}_X, B_\theta)$ with $\pi_X(X + \theta X) = \text{ad}_{\mathfrak{g}_X}(X + \theta X)$. We carry over from Lemma 3.1 the notation for weight vectors and projections. In a neighborhood I of 0 in \mathfrak{n} we have

$$(3.22) \quad \left\| E_{-2\alpha} \left(\sum_{n=0}^4 \frac{1}{n!} \pi_X(X + \theta X)^n \right) \phi_{-2\alpha} \right\|^2 \leq c_I \|E_{-2\alpha} \exp \pi_X(X + \theta X) \phi_{2\alpha}\|^2.$$

But the left side of (3.22) is

$$\left\| \left[\frac{1}{2}(\text{ad } Z)^2 + \frac{1}{6} \text{ad } Y \text{ad } Z \text{ad } Y + \frac{1}{6}(\text{ad } Y)^2 \text{ad } Z + \frac{1}{6} \text{ad } Z(\text{ad } Y)^2 + \frac{1}{24}(\text{ad } Y)^4 \right] \phi_{2\alpha} \right\|^2$$

and since $[Y, Z] = 0$, hence $\text{ad } Z(\text{ad } Y)^2 = \text{ad } Y \text{ad } Z \text{ad } Y = (\text{ad } Y)^2 \text{ad } Z$, (3.22) for X in I simplifies to

$$\left\| \left[\frac{1}{2}(\text{ad } Z)^2 + \frac{1}{2}(\text{ad } Y)^2 \text{ad } Z + \frac{1}{24}(\text{ad } Y)^4 \right] \phi_{2\alpha} \right\|^2 \leq c_I \|E_{-2\alpha} \pi_X(k_X) \phi_{2\alpha}\|^2.$$

Now, if $Z \neq 0$,

$$(3.23) \quad \phi_{2\alpha} = \|Z\|^{-1} \theta Z,$$

hence by (3.17) $(\text{ad } Y)^2(\text{ad } Z) \phi_{2\alpha} = 0$, so that

$$(3.24) \quad \left\| \frac{1}{2}(\text{ad } Z)^2 \phi_{2\alpha} + \frac{1}{24}(\text{ad } Y)^4 \phi_{2\alpha} \right\|^2 \leq c_I \|E_{-2\alpha} \pi_X(k_X) \phi_{2\alpha}\|^2.$$

Substituting in (3.23) and using (3.19) and (3.20) we get

$$2|\alpha|^2 \|Z\|^2 + \frac{1}{24} c_0 \|Y\|^4 \leq c_I \|E_{-2\alpha} \pi_X(k_X) \phi_{2\alpha}\|.$$

By shrinking the neighborhood I if necessary so that if $X \in I$

$$d(k_X, 1) \leq c_I \|Y + Z\| \leq c_I \left(2|\alpha|^2 \|Z\|^2 + \frac{1}{24} c_0 \|Y\|^4 \right)^{1/4}.$$

we have in view of (3.9)

$$d(k_X, 1) \leq c_I \left(e^{\rho_X \log a(k_X^{-1}w)} \right)^{1/4} = c_I e^{-(p+2q)^{-1} \rho \log a(k_X^{-1}w)}.$$

Since $d(k_X^{-1}, 1) = d(k_X, 1)$, choosing I to be symmetric we get

$$d(k_X, 1) \leq c_I e^{-\varepsilon_0 \rho \log a(k_X w)}$$

with $\varepsilon_0 = (p + 2q)^{-1}$. □

PROPOSITION 3.7. *The operator valued map of \bar{N} into $\text{End}_{\mathbb{C}}(V_\lambda)$ defined by $v \rightarrow e^{\rho H(v)} [\tau_\lambda(\kappa(v)w)^{-1} - \sigma_\lambda(vw)^{-1}]$ is integrable.*

Proof. Because $m(vw) = m(\kappa(v)w)$, we have the Lipschitz inequality

$$\| \tau_\lambda(\kappa(v)w)^{-1} - \sigma_\lambda(vw)^{-1} \| \leq cd(\kappa(v)w, m(\kappa(v)w)).$$

Thus, in view of (2.6) it suffices to show that for some $\varepsilon > 0$

$$(3.25) \quad d(\kappa(v)w, m(\kappa(v)w)) \leq ce^{\varepsilon \rho H(v)}$$

for $|v|$ sufficiently large. But by Proposition 3.4, for $|v|$ sufficiently large $k = m(\kappa(v)w)^{-1} \kappa(v)w$ is sufficiently close to 1 so that the neighborhood I in Lemma 3.6 may be used as a chart via $k_X = \exp(X + \theta X)$. Thus, by (3.21)

$$d(\kappa(v)w, m(\kappa(v)w)) \leq c_I e^{-\varepsilon_0 \rho \log a(m(\kappa(v)w)^{-1} \kappa(v)w^2)}.$$

Since $a(m(\kappa(v)w)^{-1} \kappa(v)w^2) = a(\kappa(v)) = -H(v)$, (3.25) follows. □

We may now use Proposition 3.7 to deal with the singularity of $|v|^{-1}$ at $v = 1$. To do so we will construct a modification of a partition of unity found in [10]. Following [10, p. 521, 523] we identify \bar{N} and $\theta \mathfrak{n}$ and transfer the norm $\| \cdot \|$ to \bar{N} . The norm on V_λ will be denoted by $\| \cdot \|_{V_\lambda}$. Fix any positive number R_0 (with the intention of doing a Taylor expansion in $\{ \|v\| < R_0 \}$). Let $\varphi(s)$ be a nonincreasing element of $C_0^\infty([0, \infty), [0, 1])$ that is equal to 1 for $0 \leq s \leq d$ and to 0 for $b \leq s < \infty$ where $0 < d < b$ and b is chosen so that $\{ |v| < b \} \subset \{ \|v\| < R_0 \}$ (cf. [10, p. 529]). Define $\psi_1(k)$ by

$$(3.26) \quad \psi_1(k) = \begin{cases} \varphi(|v|) & \text{if } k = m\kappa(v)w \text{ for some } m \in M, v \in \bar{N} \\ 0 & \text{otherwise.} \end{cases}$$

LEMMA 3.8. *The function ψ_1 defined on K by (3.26) is a well-defined, left M -invariant, smooth separation of the two closed disjoint subsets M and Mw of K , and does not depend on the function f .*

Proof. That ψ_1 is well defined and smooth follows from the Gelfand-Naimark decomposition. It is clear that ψ_1 is left M -invariant and $\psi_1|_{Mw} \equiv 1$. The existence of an element m in M for which $\psi_1(m) \neq 0$ would imply the existence of an element v in \bar{N} for which $\kappa(v)$ belongs to Mw , contradicting the disjointness of \bar{N} and $MANw$. Thus $\psi_1|_M \equiv 0$. \square

Let $\psi_2 = 1 - \psi_1$, $\varphi_1 = \varphi$, and $\varphi_2 = 1 - \varphi_1$. Thus $\psi_i f$ is in $C^\infty(K, \sigma_\lambda)$ ($i = 1, 2$) and $\mathcal{S}_\lambda f = \mathcal{S}_\lambda(\psi_1 f) + \mathcal{S}_\lambda(\psi_2 f)$. Let $Yf(v) = W^w f(v) - f(w)e^{\rho H(v)}$ and let $Zf(v) = f(\kappa(v)w) - f(w)$, that is, $Yf(v) = e^{\rho H(v)}Zf(v)$. Then

$$(3.27) \quad \mathcal{S}_\lambda f(a) = \sum_{i=1}^2 \int_{\bar{N}} \varphi_i(|v|) K(v:a) \tau_\lambda(\kappa(\delta_a v)w)^{-1} Yf(v) dv \\ + \int_{\bar{N}} K(v:a) \tau_\lambda(\kappa(\delta_a v)w)^{-1} f(w) e^{\rho H(v)} dv.$$

We will deal with the integrals in (3.27) in the order in which they are written. With ν an integer to be determined by Lemma 3.9, the function $Zf(v)$ has a Taylor expansion in $\{\|v\| < R_0\}$ of the form (cf. [10, p. 523])

$$(3.28) \quad Zf(v) = \sum_{j=1}^{2\nu} f_j(v) + R_\nu(v)$$

where $f_j(v)$ is α -homogeneous of degree j and

$$(3.29) \quad \|R_\nu(v)\|_{V_\lambda} \leq c \|v\|^{\nu+1}.$$

LEMMA 3.9. *For ν sufficiently large*

$$\lim_{a \rightarrow \infty} \int_{\bar{N}} \varphi(|v|) K(v:a) \tau_\lambda(\kappa(\delta_a v)w)^{-1} e^{\rho H(v)} R_\nu(v) dv$$

exists and equals

$$\int_{\bar{N}} \varphi(|v|) \frac{\sigma_\lambda(vw)^{-1}}{|v|} e^{\rho H(v)} R_\nu(v) dv.$$

Proof. By (3.12), (3.13), and Proposition 3.4, the result will follow from the dominated convergence theorem if we show that $e^{\rho H(v)}|v|^{-1}\|v\|^{\nu+1}\chi_{\{|v|<b\}}(v)$ is integrable for ν sufficiently large. This follows from [10, p. 529 (10.2)]. \square

LEMMA 3.10. For each $j = 1, \dots, 2\nu$

$$\lim_{a \rightarrow \infty} \int_{\bar{N}} \varphi(|v|) K(v : a) \tau_\lambda(\kappa(\delta_a v) w)^{-1} e^{\rho H(v)} f_j(v) dv$$

exists and equals

$$\int_{\bar{N}} \varphi(|v|) \frac{\sigma_\lambda(vw)^{-1}}{|v|} e^{\rho H(v)} f_j(v) dv.$$

Proof. Here it suffices to exhibit the integrability of the dominating function

$$G_j(v) = |v|^{-1} \|f_j(v)\|_{V_\lambda \mathcal{X}_{\{|v| < b\}}(v)}.$$

Now

$$h_j(v) = |v|^{-j/(p+2q)} \|f_j(v)\|_{V_\lambda}$$

is α -homogeneous of degree 0. By [10, Proposition 3] there exists a real number $e(h_j)$ for which

$$\int_{\bar{N}} G_j(v) dv = e(h_j) \int_0^\infty \left(\frac{1}{r}\right)^{1-j/(p+2q)} \varphi(r) dr.$$

The right hand side is clearly finite for any $j > 0$. □

COROLLARY 3.11. For f in $C^\infty(K, \sigma_\lambda)$,

$$\lim_{a \rightarrow \infty} \int_{\bar{N}} \varphi(|v|) K(v : a) \tau_\lambda(\kappa(\delta_a v) w)^{-1} Yf(v) dv$$

exists and equals

$$\int_{\bar{N}} \varphi(|v|) \frac{\sigma_\lambda(vw)^{-1}}{|v|} Yf(v) dv.$$

LEMMA 3.12. For f in $C^\infty(K, \sigma_\lambda)$,

$$\lim_{a \rightarrow \infty} \int_{\bar{N}} \varphi_2(|v|) K(v : a) \tau_\lambda(\kappa(\delta_a v) w)^{-1} Yf(v) dv$$

exists and equals

$$\int_{\bar{N}} \varphi_2(|v|) \frac{\sigma_\lambda(vw)^{-1}}{|v|} Yf(v) dv.$$

Proof. Here we may take $2\|f\|_{L^\infty}|v|^{-1}e^{\rho H(v)}\chi_{\{d < |v|\}}(v)$ as the pointwise dominating function and its integrability follows from (2.6) and (3.6). □

Lemma 3.12 together with Corollary 3.11 show that

$$(3.30) \quad \lim_{a \rightarrow \infty} \mathcal{S}_\lambda(Yf)(a) = \int_{\bar{N}} \frac{\sigma_\lambda(vw)^{-1}}{|v|} Yf(v) dv.$$

Only the last integral in (3.27) remains. As a first step in this consideration we write

$$(3.31) \quad \int_{\bar{N}} K(v:a)\tau_\lambda(\kappa(\delta_a v)w)^{-1}f(w)e^{\rho H(v)}dv \\ = \int_{\bar{N}} \xi_a(v)f(w)e^{\rho H(v)}dv + \int_{\bar{N}} K(v:a)\sigma_\lambda(vw)^{-1}f(w)e^{\rho H(v)}dv$$

where $\xi_a: \bar{N} \rightarrow \text{End } V_\lambda$ is defined by

$$(3.32) \quad \xi_a(v) = K(v:a)\left[\tau_\lambda(\kappa(\delta_a v)w)^{-1} - \sigma_\lambda(vw)^{-1}\right].$$

Note that by Proposition 3.7 $\xi_1(v) = e^{\rho H(v)}[\tau_\lambda(\kappa(v)w)^{-1} - \sigma_\lambda(vw)^{-1}]$ is integrable. Furthermore, by the homogeneity property (3.5) and by (3.14),

$$(3.32) \quad \xi_a(v)dv = \xi_1(\delta_a v)d(\delta_a v).$$

Thus, if we let Ξ be the element of $\text{End}(V_\lambda)$ given by

$$(3.33) \quad \Xi = \int_{\bar{N}} \xi_1(v)dv,$$

a standard approximation to the identity argument gives

$$(3.34) \quad \lim_{a \rightarrow \infty} \int_{\bar{N}} \xi_a(v)f(w)e^{\rho H(v)}dv = \Xi f(w)$$

and so we may turn our attention to establishing the limit of the last term in (3.31). We introduce the truncated kernel

$$(3.35) \quad K'(v:a) = \begin{cases} |v|^{-1} & \text{if } e^{-2\rho \log a} \leq |v| \\ 0 & \text{otherwise} \end{cases}$$

and we let

$$(3.36) \quad \theta_a(v)[K(v:a) - K'(v:a)]\sigma_\lambda(vw)^{-1}.$$

It is a simple matter using (3.6) to verify that $\theta_1(v)$ belongs to $L^1(\bar{N}, \text{End}(H_\lambda))$. Furthermore, $\theta_a(v)dv = \theta_1(\delta_a v)d(\delta_a v)$ by (3.2) and

(3.4). If we let ϑ be the element of $\text{End}(H_\lambda)$ given by

$$(3.37) \quad \vartheta = \int_{\bar{N}} \theta_1(v) dv$$

then an approximation to the identity argument gives

$$(3.38) \quad \lim_{a \rightarrow \infty} \int_{\bar{N}} \theta_a(v) f(w) e^{\rho H(v)} dv = \vartheta f(w)$$

LEMMA 3.13.

$$\lim_{a \rightarrow \infty} \int_{\bar{N}} K(v:a) \sigma_\lambda(vw)^{-1} f(w) e^{\rho H(v)} dv$$

exists and equals

$$\begin{aligned} & \int_{|v| \leq 1} \frac{e^{\rho H(v)} - 1}{|v|} \sigma_\lambda(vw)^{-1} f(w) dv \\ & + \int_{1 < |v|} \frac{e^{\rho H(v)}}{|v|} \sigma_\lambda(vw)^{-1} f(w) dv + \vartheta f(w). \end{aligned}$$

Proof.

$$\begin{aligned} & \int_{\bar{N}} K'(v:a) \sigma_\lambda(vw)^{-1} f(w) e^{\rho H(v)} dv \\ & = \int_{e^{-2\rho \log a} \leq |v|} \frac{e^{\rho H(v)}}{|v|} \sigma_\lambda(vw)^{-1} f(w) dv \\ & = \int_{e^{-2\rho \log a} \leq |v| \leq 1} \frac{e^{\rho H(v)}}{|v|} \sigma_\lambda(vw)^{-1} f(w) dv \\ & \quad + \int_{1 < |v|} \frac{e^{\rho H(v)}}{|v|} \sigma_\lambda(vw)^{-1} f(w) dv \end{aligned}$$

Now

$$\int_{e^{-2\rho \log a} \leq |v| \leq 1} \frac{\sigma_\lambda(vw)^{-1}}{|v|} f(w) dv = 0$$

for every a by [10, Proposition 20] so we may subtract it from the first term on the right hand side above. Since $(e^{\rho H(v)} - 1)|v|^{-1}$ is continuous on the compact set $\{|v| \leq 1\}$, the lemma follows by decomposing $K(v:a)$ as $K(v:a) = \theta_a(v) + K'(v:a)$ and using (3.38). \square

We combine the results of (3.27), Corollary 3.11, Lemma 3.12, (3.31), (3.34), and Lemma 3.13 as

THEOREM 3.14. *Let f belong to $C^\infty(K, \sigma_\lambda)$ and let \mathcal{S}_λ , Ξ and ϑ be the operators defined by (2.18), (3.33), and (3.37) respectively. Then*

$$\begin{aligned} \lim_{a \rightarrow \infty} \mathcal{S}_\lambda f(a) &= \Xi f(w) + \vartheta f(w) \\ &+ \int_{\bar{N}} \frac{\sigma_\lambda(vw)^{-1}}{|v|} [W^w f(v) - e^{\rho H(v)} f(w)] dv \\ &+ \int_{|v| \leq 1} \frac{\sigma_\lambda(vw)^{-1}}{|v|} [e^{\rho H(v)} f(w) - f(w)] dv \\ &+ \int_{1 < |v|} \frac{\sigma_\lambda(vw)^{-1}}{|v|} e^{\rho H(v)} f(w) dv. \end{aligned}$$

We can write this result in a more convenient form through the use of principal value integrals. Thus, for F in $L^2(\bar{N}, H_\lambda)$ we make the following interpretation of the singular integral with kernel $|v|^{-1} \sigma_\lambda(vw)^{-1}$:

$$\int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} F(v) dv = \lim_{\epsilon \rightarrow 0} \int_{\epsilon < |v|} |v|^{-1} \sigma_\lambda(vw)^{-1} F(v) dv.$$

Then, by using the mean value zero property (1.5) of $|v|^{-1} \sigma_\lambda(vw)^{-1}$ on spherical shells, we can rewrite the conclusion of Theorem 3.14 as

$$(3.39) \quad \lim_{a \rightarrow \infty} \mathcal{S}_\lambda f(a) = \Xi f(w) + \vartheta f(w) + \int_{\bar{N}} \frac{\sigma_\lambda(vw)^{-1}}{|v|} W^w f(v) dv.$$

Furthermore, the same device shows that

$$\int_{\bar{N}} e^{\rho H(v)} \tau_\lambda(\kappa(v)w)^{-1} f(w) dv$$

has meaning as a principal value integral and in fact

$$\lim_{V \rightarrow \infty} \int_{|v| < V} e^{\rho H(v)} \tau_\lambda(\kappa(v)w)^{-1} f(w) dv = (\Xi + \vartheta) f(w).$$

Thus,

$$(3.40) \quad \begin{aligned} \lim_{a \rightarrow \infty} \mathcal{S}_\lambda f(a) &= \int_{\bar{N}} e^{\rho H(v)} \tau_\lambda(\kappa(v)w)^{-1} f(w) dv \\ &+ \int_{\bar{N}} \frac{\sigma_\lambda(vw)^{-1}}{|v|} W^w f(v) dv. \end{aligned}$$

Let E_λ be the projection of V_λ onto H_λ . Observe that the last term in (3.40) already has values in H_λ and is not affected by the projection E_λ .

LEMMA 3.15. *There exists a constant a_λ such that*

$$(3.41) \quad \lim_{a \rightarrow \infty} E_\lambda \mathcal{L}_\lambda f(a) = a_\lambda f(w) + \int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} W^w f(v) dv.$$

Proof. That

$$E_\lambda \int_{\bar{N}} e^{\rho H(v)} \tau_\lambda(\kappa(v)w)^{-1} dv = a_\lambda I$$

is a consequence of Schur’s lemma. Indeed, for m in M , we have $wm = mw$, $m^{-1}\kappa(v)m = \kappa(m^{-1}vm)$, and $d(mvm^{-1}) = dv$. Hence, $\sigma_\lambda(m)$ and

$$E_\lambda \int_{\bar{N}} e^{\rho H(v)} \tau_\lambda(\kappa(v)w)^{-1} dv$$

commute. □

To each element F in the image of S_λ in $C^\infty(G, \tau_\lambda)$ we now associate a function $\mathcal{L}F$ on K with values in H_λ as follows:

$$(3.42) \quad \mathcal{L}F(k) = \lim_{a \rightarrow \infty} E_\lambda \mathcal{L}_\lambda(w^{-1}kf)(a) \quad (F = S_\lambda f; f \in C^\infty(K, \sigma_\lambda)).$$

Using the form of S_λ given in (1.3) we have

$$(3.43) \quad \mathcal{L}F(k) = \lim_{a \rightarrow \infty} e^{\rho \log a} E_\lambda \int_K \tau_\lambda(l)^{-1} f(law^{-1}k) dl,$$

from which it follows that

$$\mathcal{L}F(mk) = \sigma_\lambda(m) \mathcal{L}F(k).$$

We extend the definition of $\mathcal{L}F$ to G by

$$(3.44) \quad \mathcal{L}F(g) = e^{\rho H(g)} \mathcal{L}F(\kappa(g)).$$

THEOREM 3.16. *The boundary value map \mathcal{L} defined by (3.42) maps $S_\lambda(C^\infty(K, \sigma_\lambda))$ into $L^2(K, \sigma_\lambda)$ in a G -equivariant manner. Furthermore, the intertwining operator that is the composite*

$$\mathcal{U}(\sigma_\lambda : \rho) \xrightarrow{W^{-1}} U(\sigma_\lambda : \rho) \xrightarrow{\mathcal{L} \circ S_\lambda} U(\sigma_\lambda : \rho) \xrightarrow{W} \mathcal{U}(\sigma_\lambda : \rho)$$

is the projection $a_\lambda I + \int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} F(v \cdot) dv$, i.e.

$$(3.45) \quad W\mathcal{L}(S_\lambda W^{-1}F)(u) = a_\lambda F(u) + \int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} F(vu) dv$$

for a smooth element F in $L^2(\bar{N}, H_\lambda)$.

Proof. Since S_λ is G -equivariant, to establish the G -equivariance of \mathcal{L} it must be shown that

$$(3.46) \quad \mathcal{L}(S_\lambda U(g)f)(x) = \mathcal{L}(S_\lambda f)(xg) \quad (x, g \in G).$$

It suffices to prove (3.45) when $x = 1$. For g in K this follows from the definition. For $g = a_0 \in A$ we have by (3.43)

$$\begin{aligned} \mathcal{L}(S_\lambda U(a_0)f)(1) &= \lim_{a \rightarrow \infty} e^{\rho \log a} E_\lambda \int_K \tau_\lambda(l)^{-1} f(law^{-1}a_0) dl \\ &= e^{\rho \log a_0} \lim_{a \rightarrow \infty} e^{\rho \log aa_0^{-1}} E_\lambda \int_K \tau_\lambda(l)^{-1} f(laa_0^{-1}w^{-1}) dl \\ &= e^{\rho \log a_0} \mathcal{L}(S_\lambda f)(1) = \mathcal{L}(S_\lambda f)(a_0). \end{aligned}$$

It follows from the Cartan decomposition $G = KAK$ that (3.46) holds for every g when $x = 1$. Finally, from the explicit limit formula given in Theorem 3.14, it is clear that $\mathcal{L}F(k)$ is continuous; hence, $\mathcal{L}F$ belongs to $L^2(K, \sigma_\lambda)$. This proves the first part.

Now let F be a smooth element of $L^2(\bar{N}, H_\lambda)$. Then

$$(W^{-1}F)(k) = e^{\rho \log a(k)} \sigma_\lambda(m(k)) F(\bar{n}(k)).$$

By (3.41) and (3.42) we get

$$\mathcal{L}(S_\lambda W^{-1}F)(k) = a_\lambda W^{-1}F(k) + \int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} W^k(W^{-1}F)(v) dv.$$

Since W , W^{-1} , and $\mathcal{L} \circ S_\lambda$ are equivariant, we have for u in \bar{N}

$$\begin{aligned} W(\mathcal{L}(S_\lambda W^{-1}F))(u) &= W(\mathcal{L}(S_\lambda W^{-1} \mathcal{Q}(u) F))(1) \\ &= e^{\rho H(1)} \mathcal{L}(S_\lambda W^{-1} \mathcal{Q}(u) F)(1) \\ &= a_\lambda F(u) + \int_{\bar{N}} |v|^{-1} \sigma_\lambda(vw)^{-1} F(vu) dv. \quad \square \end{aligned}$$

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