

Generalized Variational Principles, Global Weak Solutions and Behavior with Random Initial Data for Systems of Conservation Laws Arising in Adhesion Particle Dynamics

Weinan E¹, Yu.G. Rykov², Ya.G. Sinai³

¹ Courant Institute of Mathematical Sciences, New York, NY 10012, USA

² Keldysh Institute of Applied Mathematics of Russian Academy of Sciences, Moscow, Russia

³ Mathematics Department, Princeton University, Princeton, NJ 08544, USA and Landau Institute of Theoretical Physics, Moscow, Russia

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Abstract: We study systems of conservation laws arising in two models of adhesion particle dynamics. The first is the system of free particles which stick under collision. The second is a system of gravitationally interacting particles which also stick under collision. In both cases, mass and momentum are conserved at the collisions, so the dynamics is described by 2×2 systems of conservation laws. We show that for these systems, global weak solutions can be constructed explicitly using the initial data by a procedure analogous to the Lax–Oleinik variational principle for scalar conservation laws. However, this weak solution is not unique among weak solutions satisfying the standard entropy condition. We also study a modified gravitational model in which, instead of momentum, some other weighted velocity is conserved at collisions. For this model, we prove both existence and uniqueness of global weak solutions. We then study the qualitative behavior of the solutions with random initial data. We show that for continuous but nowhere differentiable random initial velocities, all masses immediately concentrate on points even though they were continuously distributed initially, and the set of shock locations is dense.

1. Introduction

This paper has two main goals: The first is to give an explicit construction of weak solutions for the initial value problem of the systems of conservation laws:

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ (\rho u)_t + (\rho u^2)_x = 0 \end{cases} \quad (1.1)$$

and

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ (\rho u)_t + (\rho u^2)_x = -\rho g_x \\ g_{xx} = \rho \cdot \end{cases} \quad (1.2)$$

The second is to study the qualitative behavior of such weak solutions when initial data are random. We prove that for a wide class of probability distributions for the

initial data, almost every weak solution has the following structure: At *any* positive time $t > 0$, $\rho(\cdot, t)$ becomes a *purely singular measure* even though it may be continuous at $t = 0$. Moreover, this singular measure is supported on a dense set which can also be considered as the shock set of u . We will also study a variant of (1.2) in which a weighted velocity instead of momentum is conserved at the collisions [GMS, VFDN]:

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ u_t + (\frac{u^2}{2})_x = -g_x \\ g_{xx} = \rho. \end{cases} \tag{1.3}$$

Our construction of weak solutions for (1.1) is based on a connection between (1.1) and the “sticky particle model” of Zeldovich (see [Z] and also [CPY]). There is a similar connection between (1.2) and the gravitationally interacting sticky particles. Consider a system of particles on R^1 with initial velocities $\{v_j^0\}$, locations $\{x_j^0\}$ and masses $\{m_j^0\}$, $j \in \mathbf{Z}$. The particles move with constant velocities unless they collide. At collisions the colliding particles stick and form a new massive particle. The mass and velocity of this new particle are given by the laws of conservation of mass and momentum. This model was proposed by Zeldovich [SZ], and developed further by Kofman, Shandarin, et al. (see [GMS, KPS], and the survey paper [VDFN]) to explain the formation of large scale structures in the universe. In this context it is also referred to as the model of “adhesion dynamics.” One main result of this paper is that the adhesion dynamics of free particles is in a sense integrable, and this gives rise to weak solutions of (1.1).

A similar connection exists between (1.2) and the gravitationally interacting sticky particles. The Hamiltonian governing the dynamics between collisions is given by

$$H(p, x) = \sum_i \frac{p_i^2}{2m_i} + \sum_{i \neq j} m_i m_j |x_i - x_j|. \tag{1.4}$$

We will assume that $\sum_i m_i < \infty$. When particles collide, again they form a new particle with mass and velocity given by the conservation of mass and momentum. The gravitational force acting on a particle is proportional to the difference between the total masses from the right and from the left of the particle. This system is also integrable in the same sense and leads to weak solutions of (1.2).

For smooth solutions, (1.1) is equivalent to the Burgers equation

$$u_t + \left(\frac{u^2}{2}\right)_x = 0 \tag{1.5}$$

together with a scalar transport equation

$$\rho_t + (\rho u)_x = 0. \tag{1.6}$$

Given the initial data $\{\rho_0, u_0\}$, the solution of (1.5)–(1.6) can be easily found via the method of characteristics. Define the forward flow map $\varphi_t : R^1 \rightarrow R^1$ by

$$x = \varphi_t(y) = y + tu_0(y). \tag{1.7}$$

For small t , this map is usually invertible, and we have

$$u(x, t) = u_0(\varphi_t^{-1}(x)), \quad \rho(x, t) = \rho_0(\varphi_t^{-1}(x)) \left| \frac{\partial x}{\partial y} \right|^{-1}, \tag{1.8}$$

where $y = \varphi_t^{-1}(x)$ defines the backward flow map.

It is well-known [L] that this construction ceases to be valid after some critical time T^* at which the solution of (1.5) develops shocks. In general (1.1) and (1.5)–(1.6) also cease to be equivalent after T^* .

In analogy with fluid mechanics, we call y the Lagrangian coordinate and $\varphi_t(y)$ the Eulerian coordinate at time t . After T^* the mapping $y \rightarrow \varphi_t(y)$ is no longer one-to-one, and no longer defined by (1.7): a whole interval can be mapped to a single point which is the location of a shock.

However, in all cases φ_t defines a partition ξ_t of R^1 where elements of the partition are given by

$$\mathcal{D}_t(x) = \{\varphi_t^{-1}(x), x \in R^1\}. \tag{1.9}$$

We should stress that the solutions are assumed to be continuous from the right. The elements of ξ_t can be either single points, or segments. More importantly, knowing ξ_t , we can reconstruct φ_t and $u(\cdot, t)$ from the two conservation laws:

$$\varphi_t(y) = \frac{\int_{C_t(y)} (\eta + tu_0(\eta))\rho_0(\eta)d\eta}{\int_{C_t(y)} \rho_0(\eta)d\eta}, \quad u(x, t) = \frac{\int_{\mathcal{D}_t(x)} u_0(\eta)\rho_0(\eta)d\eta}{\int_{\mathcal{D}_t(x)} \rho_0(\eta)d\eta}, \tag{1.10}$$

where $C_t(y)$ denotes the element of the partition ξ_t containing y , and $\mathcal{D}_t(x) = \varphi_t^{-1}(x)$. In the more general case when the initial distribution of mass is given by a nonnegative Borel measure P_0 , (1.10) takes the form

$$\varphi_t(y) = \frac{\int_{C_t(y)} (\eta + tu_0(\eta))dP_0(\eta)}{\int_{C_t(y)} dP_0(\eta)}, \quad u(x, t) = \frac{\int_{\mathcal{D}_t(x)} u_0(\eta)dP_0(\eta)}{\int_{\mathcal{D}_t(x)} dP_0(\eta)}. \tag{1.11}$$

Both (1.10) and (1.11) state that $\varphi_t(y)$ is now the position of the center of mass of $C_t(y)$.

We are left with the key step of defining $\{\xi_t\}_{t>0}$. Let us first consider the simpler case of a finite number of particles with initial data $\{x_j^0, v_j^0, m_j^0\}$, $1 \leq j \leq N$. A crucial observation is that the necessary and sufficient condition for N particles to collide and form a single particle before, or at time t , is that

$$\frac{\sum_{j=1}^J (x_j^0 + tv_j^0)m_j^0}{\sum_{j=1}^J m_j^0} \geq \frac{\sum_{j=J+1}^N (x_j^0 + tv_j^0)m_j^0}{\sum_{j=J+1}^N m_j^0} \tag{1.12}$$

holds for all J , $1 \leq J \leq N - 1$. Indeed assume that (1.12) holds, yet there are more than one cluster of particles at time t . Without loss of generality, let us assume that there are two such clusters, $\{1, 2, \dots, J'\}$ and $\{J' + 1, \dots, N\}$ located at $X_1(t)$ and $X_2(t)$ with $X_1(t) < X_2(t)$. Then we have

$$X_1(t) = \frac{\sum_{j=1}^{J'} (x_j^0 + tv_j^0)m_j^0}{\sum_{j=1}^{J'} m_j^0} < X_2(t) = \frac{\sum_{j=J'+1}^N (x_j^0 + tv_j^0)m_j^0}{\sum_{j=J'+1}^N m_j^0}, \tag{1.13}$$

since the conservation of mass and momentum dictates that the cluster has to be located at the center of mass. Equation (1.13) contradicts the assumption that (1.12) holds for all j . On the other hand, assume that (1.12) is violated for some $J = J'$, then the group of particles $\{1, \dots, J'\}$ will never catch up with the group $\{J' + 1, \dots, N\}$ before time t . For details, see Sect. 3.

Before going into the continuous case, let us state the conditions we will impose on the initial data.

Let $P_0, I_0 \in M$: the space of Radon measures on $R^1, P_0 \geq 0$.

(A1) $P_0(\Delta) < \infty$ for any compact $\Delta \subset R^1$ and P_0 is either discrete or absolutely continuous with respect to the Lebesgue measure. In the latter case, we assume that density $\rho_0(x) > 0$, for $x \in \text{Supp}(P_0)$. If $\text{Supp}(P_0)$ is unbounded, we assume additionally

$$\int_0^x s dP_0(s) \rightarrow +\infty \text{ as } |x| \rightarrow +\infty.$$

(A2) The initial distribution of momentum I_0 is absolutely continuous with respect to P_0 . The Radon–Nikodym derivative $u(\cdot, 0) = \frac{dI_0}{dP_0}$ is the initial velocity. In the case when P_0 is absolutely continuous, we assume that $u(\cdot, 0)$ is also continuous.

(A3) For any $z > 0$

$$\sup_{|x| \leq z} |u_0(x)| \leq b_0(z) \text{ and } \lim_{|z| \rightarrow \infty} \frac{1}{z} b_0(z) = 0.$$

The first essential result of this paper is the following principle for the construction of ξ_t using the initial data.

Generalized Variational Principle (GVP): $y \in R^1$ is the left endpoint of an element of ξ_t iff for any $y^-, y^+ \in R^1$, such that $y^- < y < y^+$, the following holds:

$$\frac{\int_{[y^-, y]} (\eta + tu(\eta; 0)) dP_0(\eta)}{\int_{[y^-, y]} dP_0(\eta)} < \frac{\int_{[y, y^+]} (\eta + tu(\eta; 0)) dP_0(\eta)}{\int_{[y, y^+]} dP_0(\eta)}. \tag{1.14}$$

We can also formulate GVP for right endpoints of elements of ξ_t , but we will omit this since we do not need it.

Having $\{\xi_t\}_{t>0}$, we define φ_t via (1.11) and the density and momentum distributions at time t, P_t and I_t , by

$$P_t(\Delta) = P_0(\varphi_t^{-1}(\Delta)), \quad I_t(\Delta) = I_0(\varphi_t^{-1}(\Delta)) \tag{1.15}$$

for $\Delta \subset R^1$. In the case of continuous $u(x; 0)$ the mapping φ_t is also continuous. It is clear that I_t is absolutely continuous with respect to P_t , and we can introduce the Radon–Nikodym derivative

$$u(x, t) = \frac{dI_t}{dP_t}(x) \tag{1.16}$$

which is the velocity at (x, t) .

We will use the following definition of weak solutions.

Definition 1. Let (P_t, I_t) be a family of Borel measures, weakly continuous with respect to t , such that I_t is absolutely continuous with respect to P_t for each fixed t . Define u via (1.16). $(P_t, I_t, u)_{t \geq 0}$ is a weak solution of (1.1) if, for any $f, g \in C_0^1(R^1)$, the space of C^1 -functions on R^1 with compact support, and $0 < t_1 < t_2$,

$$(D1) \quad \int f(\eta) dP_{t_2}(\eta) - \int f(\eta) dP_{t_1}(\eta) = \int_{t_1}^{t_2} d\tau \int f'(\eta) dI_\tau(\eta),$$

$$(D2) \quad \int g(\eta) dI_{t_2}(\eta) - \int g(\eta) dI_{t_1}(\eta) = \int_{t_1}^{t_2} d\tau \int g'(\eta) u(\eta, \tau) dI_\tau(\eta).$$

Theorem 1. *Under the assumptions (A1–3), the family $(P_t, I_t, u)_{t \geq 0}$ constructed using GVP gives a weak solution of (1.1) with initial data (P_0, I_0) in the sense that*

$$\lim P_t = P_0, \quad \lim I_t = I_0$$

weakly as $t \rightarrow 0^+$.

Next we turn to (1.2). First of all, we remark that the third equation in (1.2) can be interpreted as

$$-g_x = \left(\int_x^{+\infty} \rho(\xi, t) d\xi - \int_{-\infty}^x \rho(\xi, t) d\xi \right), \tag{1.17}$$

i.e. the acceleration at a point is proportional to the difference between the total masses from the right and from the left of that point. For the initial data, in addition to (A1–A3) we will also assume

$$(A4) \quad P_0(R^1) < +\infty.$$

The characteristics of (1.2) are now given by quadratic functions of t :

$$x(t) = y + u_0(y)t + \frac{1}{2} a_0(y)t^2. \tag{1.18}$$

This has the effect of changing (1.11) to

$$\begin{aligned} \varphi_t(y) &= \frac{\int_{C_t(y)} (\eta + tu_0(\eta)) dP_0(\eta)}{\int_{C_t(y)} dP_0(\eta)} + a(C_t(y)) \frac{t^2}{2}, \\ u(x, t) &= \frac{\int_{\mathcal{D}_t(x)} u_0(\eta) dP_0(\eta)}{\int_{\mathcal{D}_t(x)} dP_0(\eta)} + a(\mathcal{D}_t(x))t, \end{aligned} \tag{1.19}$$

where

$$a(C_t(y)) = P_0(I^+) - P_0(I^-). \tag{1.20}$$

I^+ and I^- are respectively the right and left connected component of $R^1 \setminus C_t(y)$.

As before, the key issue is to construct the family of partitions of R^1 , $\{\xi_t\}_{t \geq 0}$. In this case we formulate an analogous Generalized Variational Principle (GVP): $y \in R^1$ is the left end point of an element of ξ_t iff for any $y^+, y^-, y^- < y < y^+$, the following holds:

$$\begin{aligned} & \frac{\int_{[y^-, y)} (\eta + tu_0(\eta)) dP_0(\eta)}{\int_{[y^-, y)} dP_0(\eta)} + \frac{t^2}{2} (P_0(y, +\infty) - P_0(-\infty, y^-)) \\ & < \frac{\int_{[y, y^+)} (\eta + tu_0(\eta)) dP_0(\eta)}{\int_{[y, y^+)} dP_0(\eta)} + \frac{t^2}{2} (P_0(y^+, \infty) - P_0(-\infty, y)). \end{aligned} \tag{1.21}$$

Having $\{\xi_t\}_{t \geq 0}$, we can construct φ_t, u using (1.19), P_t as before, and I_t from u and P_t by a simple integration.

Definition. *Let (P_t, I_t) be a family of Borel measures, weakly continuous with respect to t , such that I_t is absolutely continuous with respect to P_t for each fixed t . Define u via (1.16). $(P_t, I_t, u)_{t \geq 0}$ is a weak solution of (1.2) iff for any*

$f, g \in C_0^1(\mathbb{R}^1)$, and $0 < t_1 < t_2$,

$$(D1') \quad \int f(\eta)dP_{t_2}(\eta) - \int f(\eta)dP_{t_1}(\eta) = \int_{t_1}^{t_2} d\tau \int f'(\eta)dI_\tau(\eta),$$

(D2')

$$\begin{aligned} \int g(\eta)dI_{t_2}(\eta) - \int g(\eta)dI_{t_1}(\eta) &= \int_{t_1}^{t_2} d\tau \int g'(\eta)u(\eta, \tau)dI_\tau(\eta) \\ &+ \int_{t_1}^{t_2} d\tau \int g(\eta)(P_\tau(\eta, -\infty) - P_\tau(-\infty, \eta))dP_\tau(\eta). \end{aligned}$$

Theorem 2. *Under the assumptions (A1–4), the family $(P_t, I_t, u)_{t \geq 0}$ constructed using GVP gives a weak solution of (1.2) with initial data (P_0, I_0) in the sense that*

$$\lim P_t = P_0, \quad \lim I_t = I_0$$

weakly as $t \rightarrow 0^+$.

Before continuing, let us put these results in the perspective of general hyperbolic conservation laws. For obvious reasons, (1.1) is sometimes referred to as the pressureless gas dynamics equations. However, compared with the isentropic gas dynamics equation

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ (\rho u)_t + (\rho u^2 + p(\rho))_x = 0 \end{cases} \tag{1.22}$$

there are two important differences. First at a technical level, the natural space for (1.1) is M , the space of Radon measures, instead of BV or L^∞ . Secondly, the standard entropy condition, which in the present case, takes the form

$$(\rho S(\rho))_t + (u\rho S(\rho))_x \leq 0, \tag{1.23}$$

where S is convex, is not enough as a uniqueness criterion. Indeed in the context of particle systems, there is a whole family of inelastic collision rules that satisfy (1.1) and (1.23). The adhesion dynamics considered here is an extreme case of these collision rules. It is easy to see that the weak solutions of (1.1) constructed in this paper has the additional property:

$$u_x < \frac{1}{t}.$$

It is suggested by Jonathan Goodman that this might be sufficient as a uniqueness criterion. Moreover, it is natural to expect that the adhesion dynamics corresponds to a form of viscosity limit. But this also remains to be proven. Notice that (1.1) is an extreme case of a non-strictly hyperbolic system: the two characteristic fields coincide.

We will also consider a model closely related to (1.2) which is also discussed in the astrophysics literature [GMS, VDFN]:

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ u_t + (\frac{u^2}{2})_x = -g_x \\ g_{xx} = \rho. \end{cases} \tag{1.24}$$

The relation between (1.2) and (1.24) is analogous to the relation between (1.1) and (1.5–6). Let $h = g_x$, we can rewrite (1.24) as

$$\begin{cases} h_t + uh_x = 0 \\ u_t + (\frac{u^2}{2})_x = -h. \end{cases} \tag{1.25}$$

If P_0 is absolutely continuous, we can reduce (1.25) to a scalar conservation law of the form

$$u_t + \left(\frac{u^2}{2}\right)_x = -h_0(G(x - ut, t)), \tag{1.26}$$

which is amenable to standard methods. Using this, we are able to establish both existence and uniqueness of global weak solutions for (1.24). This is explained in Sect. 6.

In the second part of this paper we study the qualitative behavior of these models with continuous, but non-differentiable initial data, extending some results from [S]. We show that the solutions of these models share the common feature that almost surely, at any $t > 0$, all masses concentrate on points (i.e. the absolute continuous part vanishes), and the set of shock locations is dense. This behavior was to some extent anticipated by Zeldovich [Z] in his work on cosmology. In that context, these point masses are interpreted as the galaxies in a one-dimensional universe.

Before ending this introduction, let us mention that (1.1) and (1.2) also have an origin in kinetic equations. Consider

$$f_t + vf_x = 0. \tag{1.27}$$

If we look for solutions of the form

$$f(x, v, t) = \rho(x, t)\delta(v - u(x, t)), \tag{1.28}$$

we obtain (1.1) for (ρ, u) . Similarly consider the Vlasov–Poisson–Jeans equation [VDFN]

$$\begin{cases} f_t + vf_x - g_x f_v = 0 \\ g_{xx} = \int f(x, v, t) dv. \end{cases} \tag{1.29}$$

If we look for solutions of the form (1.28), we obtain (1.2) for (ρ, u) .

The paper has eight sections and one appendix. In Sect. 2, we compare our GVP with the Lax–Oleinik variational principle for scalar quasi-linear equations. In Sect. 3, we consider the discrete version of (1.1) and prove GVP for this case. In Sect. 4, we extend these results to the continuous case and complete the proof of Theorem 1. In Sect. 5, we explain the additional steps needed for the proof of Theorem 2. In Sect. 6, we consider the modified gravitational system (1.3).

Part II consists of two sections. In Sect. 7, we consider (1.1) and (1.2) with random initial data. In Sect. 8, we extend these results to (1.3).

After this paper was submitted for publication, we received a preprint [BG] by Brenier and Grenier in which existence of weak solutions of (1.1) was proved without resorting to GVP at the continuous level. [BG] also contains some very interesting ideas for the multi-dimensional version of (1.1). We thank Brenier and Grenier for timely communication of their results.

Part I. Generalized Variational Principles and Global Weak Solutions

2. Preliminaries and Comparison with the Lax–Oleinik Variational Principle

In the following we will concentrate on (1.1). The necessary changes for (1.2) will be summarized in Sect. 5.

Intuition from adhesion dynamics suggests that the masses cluster more and more, and the accumulated masses will never split apart again. We formulate this as:

Lemma 1. *The family of partitions $\{\xi_t\}_{t>0}$ determined with the help of GVP is decreasing. In other words, for $0 < t' < t$, each element of ξ_t is contained in an element of $\xi_{t'}$.*

Proof. Assume to the contrary that there exists $y \in \partial\xi_t$, but $y \notin \partial\xi_{t'}$, where $\partial\xi_t$ denotes the collection of points belonging to the boundary of some element of ξ_t . Without loss of generality, we can assume that y is the left end point of an element in ξ_t . Then for some $y^- < y < y^+$, we should have

$$\frac{\int_{[y^-,y)} (\eta + t'u_0(\eta))dP_0(\eta)}{\int_{[y^-,y)} dP_0(\eta)} \geq \frac{\int_{[y,y^+]} (\eta + t'u_0(\eta))dP_0(\eta)}{\int_{[y,y^+]} dP_0(\eta)}. \tag{2.1}$$

Consider two linear functions $l_1(s), l_2(s)$ of s defined by

$$l_1(s) = \frac{\int_{[y^-,y)} \eta dP_0(\eta)}{\int_{[y^-,y)} dP_0(\eta)} + s \frac{\int_{[y^-,y)} u_0(\eta) dP_0(\eta)}{\int_{[y^-,y)} dP_0(\eta)},$$

$$l_2(s) = \frac{\int_{[y,y^+]} \eta dP_0(\eta)}{\int_{[y,y^+]} dP_0(\eta)} + s \frac{\int_{[y,y^+]} u_0(\eta) dP_0(\eta)}{\int_{[y,y^+]} dP_0(\eta)}. \tag{2.2}$$

For sufficiently small s , we have $l_1(s) < l_2(s)$ while $l_1(t') \geq l_2(t')$. Since l_1 and l_2 are linear, we conclude that $l_1(t) > l_2(t)$. This contradicts the fact that $y \in \partial\xi_t$: i.e. y satisfies GVP at time t .

Now we will compare GVP with the classical Lax–Oleinik variational principle see [L, O]. We assume that P_0 has a density ρ_0 , and $0 < \text{const} \leq \rho_0(x) < \infty$. Introduce

$$\phi_1(y) = \int^y (\eta + tu_0(\eta))\rho_0(\eta)d\eta, \quad \phi_2(y) = \int^y \rho_0(\eta)d\eta \tag{2.3}$$

and

$$c(y', y'') = \frac{\phi_1(y'') - \phi_1(y')}{\phi_2(y'') - \phi_2(y')}. \tag{2.4}$$

ϕ_2 defines a C^1 -diffeomorphism of R^1 . Let $y = \phi_2^{-1}(z)$, and define

$$\tilde{c}(z', z'') = \frac{\phi_1 \circ \phi_2^{-1}(z'') - \phi_1 \circ \phi_2^{-1}(z')}{z'' - z'}. \tag{2.5}$$

In these notations, (1.14) becomes

$$\inf_{y^+ > y} c(y, y^+) \geq \sup_{y^- < y} c(y^-, y) \tag{2.6}$$

or, for $z = \phi_2(y)$,

$$\inf_{z^+ > z} \tilde{c}(z, z^+) \geq \sup_{z^- < z} \tilde{c}(z^-, z). \tag{2.7}$$

Let $\phi(z) = \phi_1 \circ \phi_2^{-1}(z)$. Geometrically (2.7) means that $(z, \phi(z))$ is a point of contact of the graph of ϕ and its convex hull constructed in the coordinate z .

In comparison, the Lax–Oleinik variational principle for the Burgers equation (1.5) can be formulated as

$$u(x, t) = \frac{\partial}{\partial x} \inf_z \left\{ \frac{(x - z)^2}{2t} + \int^z u_0(\eta) d\eta \right\}. \tag{2.8}$$

It is easy to see that finding the minimum in (2.8) is the same as constructing the convex hull of the function $F(y) = \int^y (\eta + t u_0(\eta)) d\eta$ and finding the points of contact between the graphs of F and its convex hull. In fact the set of points where the minimum in (2.8) is attained is exactly the same as the set of contact points between the graphs of F and its convex hull. Notice that F is a special case of ϕ_1 when $\rho_0 \equiv 1$. We see that the construction of GVP is the same as constructing the convex hull, but in a special coordinate z . It is in this sense that (1.14) and (1.21) generalize the variational principle of Lax and Oleinik.

3. The Discrete Case

It is instructive to consider first the case of a finite collection of particles $\{(x_i^0, v_i^0, m_i^0), i = 1, 2, \dots, N\}$, where x_i^0, v_i^0 and m_i^0 are respectively the location, velocity, and mass of the i^{th} particle. The particles undergo adhesion dynamics, defined in Sect. 1. The effect of the partitions $\{\xi_t\}_{t>0}$ is to divide the particles into ordered groups or clusters $G_1(t), G_2(t), \dots, G_k(t)$, so that each group of particles are glued to a single one before or at time t , and different groups are at different locations at time t . If G is one such group, say $G = \{i, i + 1, \dots, j\}$, we denote by $C_G(t)$ its location at time t . From the conservation of mass and momentum, we know that $C_G(t)$ has to be the center of mass of G :

$$C_G(t) = \frac{\sum_{l=i}^j m_l^0 (x_l^0 + t v_l^0)}{\sum_{l=i}^j m_l^0}. \tag{3.1}$$

More generally, we will denote the expression on the right-hand side of (3.1) by $C_{i,j}(t)$. It is a linear function of t .

Lemma 2. *Let G_1 and G_2 be two neighboring groups of particles such that $C_{G_1}(t) < C_{G_2}(t)$ for $t < t^*$, and $C_{G_1}(t^*) = C_{G_2}(t^*)$. Then for $t > t^*$,*

$$C_{G_1 \cup G_2}(t) < C_{G_1}(t).$$

Proof. Since both $C_{G_1}(t)$ and $C_{G_2}(t)$ are linear functions of t , we have for $t > t^*$,

$$C_{G_1}(t) > C_{G_2}(t).$$

Since $C_{G_1 \cup G_2}(t) = \alpha C_{G_1}(t) + (1 - \alpha)C_{G_2}(t)$ for some constant $\alpha \in (0, 1)$, we have for $t > t^*$,

$$C_{G_1 \cup G_2}(t) < C_{G_1}(t).$$

This proves the lemma.

Lemma 3. *Let $G \in \xi_t$, and $G = \{x_i^0, j' \leq i \leq j''\}$. Then for any $i, j' \leq i \leq j'' - 1$, we have*

$$C_{j',i}(t) \geq C_{i+1,j''}(t). \tag{3.2}$$

Proof. Assume to the contrary that there exists an i , such that (3.2) does not hold. Since $C_{j',j''}(t) = \alpha C_{j',i}(t) + (1 - \alpha)C_{i+1,j''}(t)$ for some $0 \leq \alpha \leq 1$, we have:

$$C_{j',i}(t) < C_{j',j''}(t). \tag{3.3}$$

Consider the evolution of the set of particles $I(0) = \{x_j^0, j' \leq j \leq i\}$. Each time the set is hit from the right by a particle or a cluster of particles, we add them to our set. In this way we obtain a growing family of sets $I(s) = \{x_j^0, j' \leq j \leq i(s)\}$. From Lemma 2, when new particles are added to $I(s)$, its center of mass is moved further to the left, i.e.

$$C_{j',i(s)}(s) < C_{j',i}(t). \tag{3.4}$$

From the assumption of the lemma we have $i(t) = j''$. Hence we have

$$C_{j',j''}(t) < C_{j',i}(t), \tag{3.5}$$

contradicting (3.3).

Theorem 1'. x_j^0 is the left endpoint of an element of the partition ξ_t iff

$$\max_{j' < j} C_{j',j-1}(t) < \min_{j'' \geq j} C_{j,j''}(t). \tag{3.6}$$

Proof. Assume that (3.6) holds, and x_j^0 is not the left endpoint of an element of ξ_t . Let $G \in \xi_t$ be the element of ξ_t containing x_j^0 , $G = \{x_i^0, i_0 \leq i \leq j_0\}$ and $i_0 < j$. From Lemma 3, we have

$$C_{i_0,j-1}(t) \geq C_{j,j_0}(t). \tag{3.7}$$

This clearly contradicts (3.6).

Assume now that x_j^0 is the left end point of an element of ξ_t . For any j', j'' , $j' < j < j''$, we want to show that $C_{j',j-1}(t) < C_{j,j''}(t)$. Let I_1, I_2, \dots, I_l be the consecutive elements of ξ_t to the left of x_j^0 , and $x_{j'}^0 \in I_1 = \{x_i^0, i_1 \leq i \leq i_2\}$. Let J_1, J_2, \dots, J_r be the consecutive elements of ξ_t to the right of x_j^0 (including the one containing x_j^0), and $x_{j''}^0 \in J_r = \{x_i^0, j_1 \leq i \leq j_2\}$. From Lemma 3, we have

$$C_{j',i_2}(t) \leq C_{I_1}(t), \quad C_{j_1,j''}(t) \geq C_{I_2}(t). \tag{3.8}$$

We also have

$$C_{I_1}(t) < \dots < C_{I_l}(t) < C_{J_1}(t) < \dots < C_{J_r}(t). \tag{3.9}$$

Hence, we must have

$$C_{j',j-1}(t) < C_{j,j''}(t). \tag{3.10}$$

We now show that the adhesion dynamics gives rise to weak solutions of (1.1) with initial data $P_0 = \sum_{i=1}^N m_i^0 \delta(x - x_i^0)$, $I_0 = \sum_{i=1}^N m_i^0 v_i^0 \delta(x - x_i^0)$.

Let $\{x_j(t + \Delta t), v_j(t + \Delta t), m_j(t + \Delta t), j = 1, \dots, N(t)\}$ be the locations, velocities, and masses of the particles at time $t + \Delta t$. Assume that the j^{th} particle is formed by gluing together some particles whose locations, velocities, and masses at time t were respectively $x_{ij}(t), v_{ij}(t)$ and $m_{ij}(t)$, $i \in I_j$ —a set of indices depending on j . We have

$$\begin{aligned} x_j(t + \Delta t) &= \frac{\sum_{i \in I_j} (x_{ij}(t) + \Delta t v_{ij}(t)) m_{ij}(t)}{m_j(t + \Delta t)}, \\ v_j(t + \Delta t) &= \frac{\sum_{i \in I_j} v_{ij}(t) m_{ij}(t)}{m_j(t + \Delta t)}, \\ m_j(t + \Delta t) &= \sum_{i \in I_j} m_{ij}(t). \end{aligned} \tag{3.11}$$

We also have $|x_{ij}(t) - x_j(t + \Delta t)| \leq \text{Const} \cdot \Delta t$ since all velocities are bounded. Let $f, g \in C_0^1(R^1)$. We first show that $\int f(\eta) dP_t(\eta)$ is a differentiable function of t , and

$$\frac{d}{dt} \int f(\eta) dP_t(\eta) = \int f'(\eta) dI_t(\eta). \tag{3.12}$$

This implies (D1). Indeed, we have

$$\begin{aligned} \int f(\eta) dP_t(\eta) &= \sum_{i,j} f(x_{ij}(t)) m_{ij}(t) \\ &= \sum_j \sum_{i \in I_j} f(x_j(t + \Delta t) + x_{ij}(t) - x_j(t + \Delta t)) m_{ij}(t) \\ &= \sum_j f(x_j(t + \Delta t)) m_j(t + \Delta t) \\ &\quad + \sum_j f'(x_j(t + \Delta t)) \sum_{i \in I_j} (x_{ij}(t) - x_j(t + \Delta t)) m_{ij}(t) + o(\Delta t). \end{aligned}$$

Since

$$\begin{aligned} \sum_{i \in I_j} (x_{ij}(t) - x_j(t + \Delta t)) m_{ij}(t) &= \sum_i x_{ij}(t) m_{ij}(t) - x_j(t + \Delta t) m_j(t + \Delta t) \\ &= -\Delta t \sum_i v_{ij}(t) m_{ij}(t) = -\Delta t v_j(t + \Delta t) m_j(t + \Delta t), \end{aligned}$$

we obtain

$$\begin{aligned} \int f(\eta) dP_t(\eta) &= \int f(\eta) dP_{t+\Delta t}(\eta) - \Delta t \sum_j f'(x_j(t + \Delta t)) v_j(t + \Delta t) m_j(t + \Delta t) + o(\Delta t). \end{aligned}$$

This implies (3.12).

In contrast $\int g(\eta)dI_t(\eta)$ is not C^1 in t because of the inelastic nature of the collisions. To see this, consider just two particles colliding at $t = \tau$. Before collision they are denoted by $(x_1(t), v_1(t), m_1(t)), (x_2(t), v_2(t), m_2(t)), t < \tau$, respectively. After collision they form a single particle $(x(t), v(t), m(t)), t \geq \tau$. We have $m(\tau) = m_1(\tau) + m_2(\tau)$, $m(\tau)v(\tau) = m_1(\tau)v_1(\tau) + m_2(\tau)v_2(\tau)$, $m(\tau)x(\tau) = m_1(\tau)x_1(\tau) + m_2(\tau)x_2(\tau)$. For $t < \tau$, we have

$$\begin{aligned} \int g(\eta)dI_t(\eta) &= g(x_1(t))v_1(t)m_1(t) + g(x_2(t))v_2(t)m_2(t) \\ &= g(x(\tau))(v_1(t)m_1(t) + v_2(t)m_2(t)) \\ &\quad + (t - \tau)g'(x(\tau))(v_1^2(t)m_1(t) + v_2^2(t)m_2(t)) + O((t - \tau)^2) \\ &= \int g(\eta)dI_\tau(\eta) + (t - \tau)g'(x(\tau))(v_1^2(\tau)m_1(\tau) + v_2^2(\tau)m_2(\tau)) \\ &\quad + O((t - \tau)^2). \end{aligned}$$

For $t > \tau$, we have

$$\begin{aligned} \int g(\eta)dI_t(\eta) &= g(x(t))v(t)m(t) \\ &= \int g(\eta)dI_\tau(\eta) + (t - \tau)g'(x(\tau))v^2(\tau)m(\tau) + O((t - \tau)^2). \end{aligned}$$

In general, energy decreases at collisions

$$v_1^2(\tau)m_1(\tau) + v_2^2(\tau)m_2(\tau) \neq v^2(\tau)m(\tau).$$

Hence $\int g(\eta)dI_t(\eta)$ is not C^1 in t . Nevertheless, we can prove that (D2) is still valid. We only have to prove this for the case when there is no collision in (t_1, t_2) and at $t = t_2$, a group of particles, say with indices $i_1, i_1 + 1, \dots, i_2$, are colliding to form a single particle at $x(t_2)$. We then have

$$\begin{aligned} \int_{t_1}^{t_2} d\tau \int g'(\eta)u(\eta, \tau)dI_\tau(\eta) &= \int_{t_1}^{t_2} d\tau \sum_{i=1}^N g'(x_i(\tau))v_i^2(\tau)m_i \\ &= \int_{t_1}^{t_2} d\tau \left[\sum_{i=1}^{i_1-1} m_i v_i \frac{d}{d\tau} g(x_i(\tau)) + \sum_{i=i_1}^{i_2} m_i v_i \frac{d}{d\tau} g(x_i(\tau)) + \sum_{i=i_2+1}^N m_i v_i \frac{d}{d\tau} g(x_i(\tau)) \right] \\ &= \sum_{i=1}^{i_1-1} m_i v_i g(x_i(t_2)) + g(x(t_2)) \sum_{i=i_1}^{i_2} m_i v_i + \sum_{i=i_2+1}^N m_i v_i g(x_i(t_2)) - \sum_{i=1}^{i_1-1} m_i v_i g(x_i(t_1)) \\ &= \int g(\eta)dI_{t_2}(\eta) - \int g(\eta)dI_{t_1}(\eta). \end{aligned}$$

4. Proof of Theorem 1

We will first give the proof for the case when P_0 is absolutely continuous, and then for the case when P_0 is discrete.

For initial data satisfying (A1–A3), we construct a decreasing family of partitions $\{\xi_t\}_{t>0}$ of R^1 according to GVP (1.14). Having $\{\xi_t\}_{t>0}$, we can define φ_t, P_t, I_t , and $u(x, t)$. Obviously φ_t is a non-decreasing function of y for any fixed t .

Furthermore, as a consequence of the assumption that $u(x; 0)$ is continuous, φ_t is also continuous, and we have

$$\begin{aligned} \int f(\eta)dP_t(\eta) &= \int f(\varphi_t(\eta))dP_0(\eta), \\ \int g(\eta)dI_t(\eta) &= \int g(\varphi_t(\eta))dI_0(\eta). \end{aligned} \tag{4.1}$$

We will prove Theorem 1 via discrete approximations. Take a sequence of measures $P_0^{(n)}$ concentrated on finite sets $\{x_i^{(n)}, i = 1, \dots, \}$ such that $P_0^{(n)} \rightarrow P_0$ weakly. Define $I_0^{(n)}$ to be a signed measure concentrated on the set $\{x_i^{(n)}, i = 1, \dots, \}$ such that $I_0^{(n)}(\{x_i^{(n)}\}) = u_0(x_i^{(n)})P_0^{(n)}(\{x_i^{(n)}\})$. Then $I_0^{(n)} \rightarrow I_0$ weakly. Using GVP, we construct the corresponding families of partitions $\xi_t^{(n)}$ and mappings $\varphi_t^{(n)}$. Moreover we already showed (Theorem 1') that for $f, g \in C_0(R^1)$,

$$\begin{aligned} \int f(\eta)dP_t^{(n)} &= \int f(\varphi_t^{(n)}(\eta))dP_0^{(n)}, \\ \int g(\eta)dI_t^{(n)} &= \int g(\varphi_t^{(n)}(\eta))dI_0^{(n)}. \end{aligned} \tag{4.2}$$

Here $P_t^{(n)}, I_t^{(n)}$ are constructed in Sect. 1. We also have for $0 < t_1 < t_2$,

$$\begin{aligned} \int f(\eta)dP_{t_2}^{(n)} - \int f(\eta)dP_{t_1}^{(n)} &= \int_{t_1}^{t_2} d\tau \int f'(\eta)dI_\tau^{(n)} = \int_{t_1}^{t_2} d\tau \int f'(\varphi_\tau^{(n)}(\eta))dI_0^{(n)}, \\ \int g(\eta)dI_{t_2}^{(n)} - \int g(\eta)dI_{t_1}^{(n)} &= \int_{t_1}^{t_2} d\tau \int g'(\varphi_\tau^{(n)}(\eta))u^{(n)}(\eta, \tau)dI_0^{(n)}. \end{aligned} \tag{4.3}$$

We can extend the definition of $\varphi_t^{(n)}$ to the whole line by putting $\varphi_t^{(n)}(y) = \varphi_t^{(n)}(x_i^{(n)})$ if $x_i^{(n)} \leq y \leq x_{i+1}^{(n)}$, $\varphi_t^{(n)}(y) = \varphi_t^{(n)}(x_1^{(n)})$ if $y < x_1^{(n)}$, $\varphi_t^{(n)}(y) = \varphi_t^{(n)}(x_N^{(n)})$ if $y \geq x_N^{(n)}$. Here $x_1^{(n)} = \min\{x_i^{(n)}\}$, $x_N^{(n)} = \max\{x_i^{(n)}\}$.

We will use superscript “n” to denote objects corresponding to $P_0^{(n)}, I_0^{(n)}$. If Δ is an interval in R^1 , we denote

$$C_\Delta(t) = \frac{\int_\Delta (\eta + tu(\eta; 0))dP_0(\eta)}{\int_\Delta dP_0(\eta)}, \quad u_0(\Delta) = \frac{\int_\Delta u(\eta; 0)dP_0(\eta)}{\int_\Delta dP_0(\eta)}. \tag{4.4}$$

The following lemma is a continuous version of Lemma 3. It can be proven in the same way as Lemma 3.

Lemma 3'. *Let $G = [a, b]$ be an element of ξ_t , $a < b$. Then for any $c \in (a, b)$, we have*

$$C_{[a,c]}(t) \geq C_{[c,b]}(t). \tag{4.5}$$

Lemma 4. $\varphi^{(n)} \rightarrow \varphi$ uniformly on compact subsets of $R^1 \times [0, \infty)$.

Proof. We first prove that for any fixed $t \geq 0$, the inverse images of $\varphi_t^{(n)}$ of a bounded interval $L = (c, d)$ are uniformly bounded. To that end, let $A^{(n)} = \{x_i^{(n)}, \varphi_t^{(n)}(x_i^{(n)}) \in L\}$, and $x_{\min}^{(n)} = \min\{x_i^{(n)}, x_i^{(n)} \in A^{(n)}\}$. Obviously $x_{\min}^{(n)}$ has to be the left end point of an element $G^{(n)}$ in the partition $\xi_t^{(n)}$. From Lemma 3, we have

$$x_{\min}^{(n)} + tu_0(x_{\min}^{(n)}) \geq C_{G^{(n)}}^{(n)}(t) \geq c. \tag{4.6}$$

Using (A3), we can write (4.6) as

$$x_{\min}^{(n)}(1 + o(1)t) \geq c. \tag{4.7}$$

Hence $\{x_{\min}^{(n)}\}$ is uniformly bounded from below. Similarly we can prove an upper bound.

Assume to the contrary that $\{\varphi_t^{(n)}\}$ do not converge uniformly on some bounded set, say $[c, d] \times [0, T]$. Then there exists $\varepsilon > 0$, and sequences $\{y_n, t_n\} \in [c, d] \times [0, T]$ such that

$$|\varphi_{t_n}^{(n)}(y_n) - \varphi_{t_n}(y_n)| > \varepsilon. \tag{4.8}$$

We can choose $\{y_n\}, \{t_n\}$ such that $\lim y_n = y^*, \lim t_n = t, \lim \varphi_t^{(n)}(y_n) = x^*$ exist. Since

$$|\varphi_{t_n}^{(n)}(y_n) - \varphi_t^{(n)}(y_n)| \leq |t_n - t| \max |u(x; 0)| \leq \text{Const} |t_n - t| \tag{4.9}$$

where the maximum is taken over a bounded set, for sufficiently large n we get from (4.8)

$$|\varphi_t^{(n)}(y_n) - \varphi_t(y_n)| > \frac{\varepsilon}{2}. \tag{4.10}$$

Therefore we have

$$|x^* - \varphi_t(y^*)| \geq \frac{\varepsilon}{2}. \tag{4.11}$$

Without loss of generality, let us assume $x^* < \varphi_t(y^*)$. Let $\Delta_n = (\varphi_t^{(n)})^{-1} \{\varphi_t^{(n)}(y_n)\} = \{x_{i_n}^{(n)} < \dots < x_{j_n}^{(n)}\}$ which is an element of $\xi_t^{(n)}$. By extracting subsequences, we can assume that Δ_n converges to an interval Δ , with $\bar{\Delta} = [a, b]$, in the sense that

$$\lim x_{i_n}^{(n)} = a, \quad \lim x_{j_n}^{(n)} = b. \tag{4.12}$$

Consequently, we also have $x^* = C_{\Delta}(t)$. Δ itself may be open, closed, or half-open. Since $y_n \in \Delta_n$, we must have $y^* \in \bar{\Delta}$. Let $\varphi_t^{-1} \{\varphi_t(y^*)\} = [c, d]$. By the continuity of φ_t , this has to be a closed interval which might consist of a single point. We will discuss the case when $a < b$. The reader can easily see that the case when $a = b$ also follows.

First we prove that $c, d \in (a, b)$. Assume to the contrary that $c \in (a, b)$. We can choose a sequence $\{x_{l_n}^{(n)}\}$, $x_{i_n}^{(n)} < x_{l_n}^{(n)} < x_{j_n}^{(n)}$, such that $\lim x_{l_n}^{(n)} = c$. Since $x_{l_n}^{(n)}$ is not the end point of any element in $\xi_t^{(n)}$, there exist j'_n, j''_n , such that

$$C_{j'_n, l_{n-1}}^{(n)}(t) \geq C_{l_n, j''_n}^{(n)}(t). \tag{4.13}$$

By extracting subsequences, we can further assume that $\{x_{j'_n}^{(n)}\}$ and $\{x_{j''_n}^{(n)}\}$ converge to c^* and d^* , respectively.

Taking the limit of (4.13), we obtain

$$\frac{\int_{c^*}^c (\eta + t u_0(\eta)) dP_0(\eta)}{\int_{c^*}^c dP_0(\eta)} \geq \frac{\int_c^{d^*} (\eta + t u_0(\eta)) dP_0(\eta)}{\int_c^{d^*} dP_0(\eta)}, \tag{4.14}$$

contradicting (1.14). Therefore $c \in (a, b)$. Similarly $d \in (a, b)$.

We are left with three different cases.

Case 1. $a < b = y^* = c < d$. From

$$C_{i_n, j_n}^{(n)}(t) \geq C_{j_n, j_n}^{(n)}(t)$$

a limiting argument similar to the one presented above gives

$$x^* \geq y^* + t u_0(y^*).$$

On the other hand, from Lemma 3', we have

$$y^* + t u_0(y^*) \geq \varphi_t(y^*).$$

Hence we have $x^* \geq \varphi_t(y^*)$ contradicting to the assumption that $x^* < \varphi_t(y^*)$.

Case 2. $c < a < b < d$. Choosing $\{x_{i_n}^{(n)}\}$ monotonically decreasing to c , and using $C_{i_n, i_n-1}^{(n)}(t) < \varphi_t^{(n)}(y_n)$ we obtain

$$C_{[c, a]}(t) \leq x^*.$$

From Lemma 3' we get

$$\varphi_t(y^*) \leq C_{[c, a]}(t).$$

Again we get $x^* \geq \varphi_t(y^*)$.

Case 3. $c < d = y^* = a < b$. The proof of Case 2 applies to this case also.

Proof of Theorem 1. To complete the proof of Theorem 1, we have to establish (D1) and (D2).

For any $t > 0$, we write

$$\begin{aligned} \int f(\eta) dP_t(\eta) &= \int [f(\varphi_t(\eta)) - f(\varphi_t^{(n)}(\eta))] dP_0^{(n)}(\eta) + \int f(\varphi_t(\eta)) (dP_0(\eta) \\ &\quad - dP_0^{(n)}(\eta)) + \int f(\varphi_t^{(n)}(\eta)) dP_0^{(n)}(\eta) \\ &= \int f(\varphi_t^{(n)}(\eta)) dP_0^{(n)}(\eta) + \delta_{i,n}^{(1)} + \delta_{i,n}^{(2)}, \end{aligned}$$

where $\delta_{i,n}^{(1)}, \delta_{i,n}^{(2)} \rightarrow 0$ as $n \rightarrow \infty$. Similarly we also have

$$\begin{aligned} \int_{t_1}^{t_2} dt \int f'(\eta) dI_t(\eta) &= \int_{t_1}^{t_2} dt \int f'(\varphi_t(\eta)) dI_0(\eta) \\ &= \int_{t_1}^{t_2} dt \int [f'(\varphi_t(\eta)) - f'(\varphi_t^{(n)}(\eta))] dI_0^{(n)}(\eta) + \int_{t_1}^{t_2} dt \int f'(\varphi_t(\eta)) [dI_0(\eta) \\ &\quad - dI_0^{(n)}(\eta)] + \int_{t_1}^{t_2} dt \int f'(\varphi_t^{(n)}(\eta)) dI_0^{(n)}(\eta) \\ &= \int_{t_1}^{t_2} dt \int f'(\varphi_t^{(n)}(\eta)) dI_0^{(n)}(\eta) + \beta_{i,n}^{(1)} + \beta_{i,n}^{(2)}, \end{aligned}$$

where $\beta_{i,n}^{(1)}, \beta_{i,n}^{(2)} \rightarrow 0$ as $n \rightarrow \infty$. Together with the results of Sect. 3, we obtain (D1).

To get (D2), let

$$\begin{aligned} \omega(y, t) &= \frac{\int_{C_t(y)} u_0(\eta) dP_0(\eta)}{\int_{C_t(y)} dP_0(\eta)} = u_0(C_t(y)), \\ \omega^{(n)}(y, t) &= \frac{\int_{C_t^{(n)}(y)} u_0^{(n)}(\eta) dP_0^{(n)}(\eta)}{\int_{C_t^{(n)}(y)} dP_0^{(n)}(\eta)} = u_0^{(n)}(C_t^{(n)}(y)). \end{aligned} \tag{4.17}$$

The difficulty lies in the fact that $\omega^{(n)}(y, t)$ may not converge to $\omega(y, t)$. Roughly speaking, at the point of shock interactions, $C_t(y)$ may differ from $C_t^{(n)}(y)$ by a set of finite measure. The main remaining task, therefore, is to show that such interactions occur rather rarely and give a small contribution to the integrals in (D2).

Let $\delta > 0$. By continuity of $u(\cdot, 0)$, we can choose $\varepsilon = \varepsilon(\delta) > 0$, such that if $|C_t(y) \Delta C_t^{(n)}(y)| < \varepsilon$ (where Δ denotes symmetric difference), then

$$|\omega(y, t) - \omega^{(n)}(y, t)| < \delta. \tag{4.18}$$

Let (y, τ) be such that

$$|\omega(y, \tau) - \omega^{(n)}(y, \tau)| > \delta, \tag{4.19}$$

and let $C_\tau(y) = [A, B]$, $C_\tau^{(n)}(y) = [C, D]$. Without loss of generality, we assume $A < C < B < D$. Then either, $D - B > \varepsilon$, or $C - A > \varepsilon$.

Let us assume that $D - B > \varepsilon$, and denote $t = \tau + \gamma$, $\varphi_t^{-1}(\varphi_t(y)) = [A', B']$. Assume $B' < D$, we want to estimate γ , i.e. the time needed for the atom at $\varphi_\tau(y)$ to absorb half of the masses in (B, D) . We have

$$\begin{aligned} C_{[B', D]}(t) &= C_{[B', D]}(\tau) + \gamma u_0([B', D]) = C_{[B', D]}^{(n)}(\tau) + \gamma u_0([B', D]) + o(1) \\ &\cong \varphi_\tau^{(n)}(y) + \gamma u_0([B', D]) + o(1). \end{aligned}$$

Here $o(1)$ denotes quantities that converge to 0 as n goes to infinity, with all the other parameters such as ε, δ held fixed. Since the center of mass is a linear function of t , and

$$C_{[A, B]}(\tau) \cong C_{[C, B]}(\tau),$$

we also have

$$\begin{aligned} C_{[A, B']}(t) &\cong C_{[C, B']}(t) = C_{[C, B']}^{(n)}(\tau) + \gamma u_0([C, B']) + o(1) \\ &\cong \varphi_\tau^{(n)}(y^*) + \gamma u_0([C, B']) + o(1). \end{aligned}$$

From the definition of B' , we have

$$C_{[A, B']}(t) \leq C_{[A', B']}(t) < C_{[B', D]}(t).$$

Therefore, we get

$$\gamma \{u_0([C, B']) - u_0([B', D])\} \leq o(1). \tag{4.20}$$

On the other hand, since $C_{[C, B']}^{(n)}(\tau) \geq C_{[B', D]}^{(n)}(\tau)$, we have

$$C_{[C, B']}(t) \geq C_{[B', D]}(\tau) + o(1).$$

Hence, we get

$$\begin{aligned} \tau\{u_0([C, B']) - u_0([B', D])\} &\geq \frac{\int_{B'}^D \eta dP_0(\eta)}{\int_{B'}^D dP_0(\eta)} - \frac{\int_C^{B'} \eta dP_0(\eta)}{\int_C^{B'} dP_0(\eta)} + o(1) \\ &\geq \beta(\varepsilon) + o(1). \end{aligned} \tag{4.21}$$

Here $\beta(\varepsilon)$ is some finite quantity depending on ε . Hence, we have

$$\gamma \leq \frac{o(1)}{\beta(\varepsilon)}. \tag{4.22}$$

Similarly, if $C - A > \varepsilon$, we can show that after time $\frac{o(1)}{\beta(\varepsilon)}$, at least half of the masses in (A, C) are absorbed by the atom at $\varphi_\tau^{(n)}(y)$.

Let $S = \{\tau \in [t_1, t_2], \text{ there exists } y \in L, L \text{ is a sufficiently large interval, such that (4.19) holds}\}$. Since either $P_0(A, C) \geq \text{Const} \cdot \varepsilon$, or $P_0(B, D) \geq \text{Const} \cdot \varepsilon$, and since $P_0(L) < \text{Const}$, we must have

$$|S| < \frac{o(1)}{\beta(\varepsilon)\varepsilon}. \tag{4.23}$$

Now let $g \in C_0^1(R^1)$. We choose L to contain the support of $g(\varphi_\tau(\cdot))$, $\tau \in [t_1, t_2]$,

$$\begin{aligned} &\left| \int_{t_1}^{t_2} d\tau \int g'(\eta) u_\tau^{(n)}(\eta)^2 dP_\tau^{(n)}(\eta) - \int_{t_1}^{t_2} d\tau \int g'(\eta) u(\eta, \tau)^2 dP_\tau(\eta) \right| \\ &= \left| \int_{t_1}^{t_2} d\tau \int g'(\varphi_\tau^{(n)}(\eta)) \omega^{(n)}(\eta, \tau)^2 dP_0^{(n)}(\eta) - \int_{t_1}^{t_2} d\tau \int g'(\varphi_\tau(\eta)) \omega(y, \tau)^2 dP_0(\eta) \right| \\ &\leq \left| \int_{t_1}^{t_2} d\tau \int (g'(\varphi_\tau^{(n)}(\eta)) - g'(\varphi_\tau(\eta))) \omega^{(n)}(\eta, \tau)^2 dP_0^{(n)}(\eta) \right| \\ &\quad + \left| \int_{t_1}^{t_2} d\tau \int g'(\varphi_\tau(\eta)) [\omega^{(n)}(\eta, \tau)^2 - \omega(\eta, \tau)^2] dP_0(\eta) \right| \\ &\leq o(1) + \left| \int_S d\tau \int g'(\varphi_\tau(\eta)) [\omega^{(n)}(\eta, \tau)^2 - \omega(\eta, \tau)^2] dP_0(\eta) \right| \\ &\quad + \left| \int_{[t_1, t_2] \setminus S} d\tau \int g'(\varphi_\tau(\eta)) [\omega^{(n)}(\eta, \tau)^2 - \omega(\eta, \tau)^2] dP_0(\eta) \right| \\ &\leq o(1) + \text{Const} \cdot |S| + \text{Const} \cdot \delta \leq o(1) + \frac{o(1)}{\beta(\varepsilon)\varepsilon} + \text{Const} \cdot \delta, \end{aligned}$$

where the constants depend on the support of g . Now we can choose n so large that $\frac{o(1)}{\beta(\varepsilon)\varepsilon} + o(1) < \delta$. Then we get

$$\left| \int_{t_1}^{t_2} d\tau \int g'(\eta) u_\tau^{(n)}(\eta)^2 dP_\tau^{(n)}(\eta) - \int_{t_1}^{t_2} d\tau \int g'(\eta) u(\eta, \tau)^2 dP_\tau(\eta) \right| \leq \text{Const} \cdot \delta.$$

Together with the argument which proved (D1), we get (D2).

Now we discuss in detail the case of discrete P_0 . We will use the same notations. We choose a natural approximating sequence $P^{(n)}, I^{(n)}$. Namely, we order all points by their masses and $P^{(n)}$ is concentrated on the set of all points whose masses are not less than $\frac{1}{n}$. $I^{(n)}$ is defined in an analogous way. We use the following lemma which will be proved later.

Lemma 4'. *For any $\beta > 0$ take $S_\beta = \{x: P_0(\{x\}) \geq \beta\}$. Then $\varphi_t^{(n)}(y)$ converges uniformly to $\varphi_t(y)$ on bounded subsets of $S_\beta \times R^+$ as $n \rightarrow \infty$.*

The proof of (D1) basically remains the same as in the continuous case. So the main problem is to prove (D2).

Fix β, ε . Assume that $(y, \tau), y \in S_\beta$ is such that for all large enough n ,

$$|u_0(C_\tau(y)) - u_0^{(n)}(C_\tau^{(n)}(y))| > \varepsilon .$$

We have

$$\begin{aligned} |u_0^{(n)}(C_\tau^{(n)}(y)) - u_0(C_\tau^{(n)}(y))| &= \left| \frac{\int_{C_\tau^{(n)}(y)} u(\eta; 0) dP_0^{(n)}(\eta)}{\int_{C_\tau^{(n)}(y)} dP_0^{(n)}(\eta)} - \frac{\int_{C_\tau^{(n)}(y)} u(\eta; 0) dP_0(\eta)}{\int_{C_\tau^{(n)}(y)} dP_0(\eta)} \right| \\ &\leq \frac{\text{const} \cdot \delta_n}{\beta^2} , \end{aligned}$$

where $\delta_n \rightarrow 0$ as $n \rightarrow \infty$. Choose n so large that $\frac{\text{const} \cdot \delta_n}{\beta^2} < \frac{\varepsilon}{2}$. Then

$$|u_0(C_\tau(y)) - u_0(C_\tau^{(n)}(y))| > \frac{\varepsilon}{2} . \tag{4.25}$$

Without any loss of generality we may assume that $C_\tau(y) = \Delta_1 \cup \Delta_2, C_\tau^{(n)}(y) = \Delta_2 \cup \Delta_3$ and $\Delta_1(\Delta_3)$ lies to the left (right) of Δ_2 . We can write

$$\begin{aligned} u_0(C_\tau(y)) &= \alpha u_0(\Delta_1) + (1 - \alpha)u_0(\Delta_2) , \\ u_0(C_\tau^{(n)}(y)) &= \xi u_0(\Delta_3) + (1 - \xi)u_0(\Delta_2) . \end{aligned}$$

We claim that there is a constant K , such that either $\alpha > K\varepsilon$ or $\xi > K\varepsilon$. Indeed if it were not true, we would have

$$|u_0(C_\tau(y)) - u_0(C_\tau^{(n)}(y))| \leq 4K(\max |u_0|) \cdot \varepsilon < \frac{\varepsilon}{2} ,$$

provided that $K = \frac{1}{8 \max |u_0|}$. The last inequality contradicts (4.25).

Consider for definiteness the case $\alpha > K\varepsilon$, i.e. $P_0(\Delta_1) > K\varepsilon P_0(\Delta_1 \cup \Delta_2) \geq K\beta\varepsilon$.

We will need an estimate for the time needed to absorb an essential part of Δ_1 by the atom of $\varphi_\tau^{(n)}$ in the dynamics with data (P_0^n, I_0^n) .

Let $\Delta_1 = J_1 \cup J_2$, where J_2 is the part which has been absorbed at $t = \tau + \gamma$, and $P_0(J_1) > \frac{K\beta\varepsilon}{2}$. Then

$$\begin{aligned} C_{J_1}(t) &> \varphi_\tau(y) + \gamma u_0(J_1) \geq \varphi_\tau(y) + \gamma u_0(\Delta_1) , \\ C_{J_1}^{(n)}(t) &\geq \varphi_\tau(y) + \gamma u_0(\Delta_1) + o(1) . \end{aligned}$$

Here $o(1)$ means quantities which tend to zero as $n \rightarrow \infty$ and β, ε remain fixed. We also have

$$C_{\mathcal{A}_2 \cup \mathcal{A}_3}^{(n)}(t) = \varphi_\tau^{(n)}(y) + \gamma u_0(C_\tau^{(n)}(y)) \leq \varphi_\tau^{(n)}(y) + \gamma u_0(C_\tau^{(n)}(y)) + o(1).$$

By assumption at time t , the set J_1 is not absorbed by the atom which at time τ were at $\varphi_\tau^{(n)}$, but J_2 is, we can write

$$\varphi_\tau(y) + \gamma u_0(\mathcal{A}_1) \leq \varphi_\tau^{(n)}(y) + \gamma u_0(C_\tau^{(n)}(y)) + o(1)$$

or

$$\gamma(u_0(\mathcal{A}_1) - u_0(C_\tau^{(n)}(y))) \leq o(1).$$

Since

$$u_0(\mathcal{A}_1) - u_0(C_\tau^{(n)}(y)) \geq u_0(C_\tau(y)) - u_0(C_\tau^{(n)}(y)) > \frac{\varepsilon}{2},$$

we conclude

$$\gamma \leq o(1).$$

Define $S^\beta = \{\tau \in [t_1, t_2] \text{ one can find } y \in L \cap S_\beta \text{ so that (4.24) holds}\}$. As in the continuous case, we have

$$|S^\beta| \leq o(1).$$

Now take $g \in C_0^1(R^1)$. We can write

$$\begin{aligned} & \left| \int_{t_1}^{t_2} dt \int g'(\eta)(u_\tau^{(n)}(\eta))^2 dP_0^{(n)}(\eta) - \int_{t_1}^{t_2} dt \int g'(\eta)(u_\tau(\eta))^2 dP_0(\eta) \right| \\ & \leq o(1) + \left| \int_{t_1}^{t_2} dt \int g'(\varphi_\tau(\eta))[(\omega^{(n)}(\eta; \tau))^2 - (\omega(\eta; \tau))^2] dP_0(\eta) \right| \\ & \leq o(1) + \text{const} \left| \int_{t_1}^{t_2} dt \int_{P_0(\eta) < \beta} dP_0(\eta) \right| \\ & \quad + \left| \int_{S^\beta} dt \int_{S_\beta} g'(\varphi_\tau(\eta))((\omega^{(n)}(\eta; \tau))^2 - (\omega(\eta; \tau))^2) dP_0(\eta) \right| \\ & \quad + \left| \int_{[t_1; t_2] \setminus S^\beta} \int_{S^\beta} g'(\varphi_\tau(\eta))(\omega^{(n)}(\eta; \tau))^2 - (\omega(\eta; \tau))^2 dP_0(\eta) \right| \\ & \leq o(1) + C(\beta) + \text{const} |S^\beta| + \text{const} \cdot \varepsilon = o(1) + C(\beta) + \text{const} \cdot \varepsilon. \end{aligned}$$

The rest of the proof is the same as in the continuous case.

Proof of Lemma 4'. We will follow the proof of Lemma 4. Assume that one can find a sequence $\{y^{(n)}\}$ such that $P_0(y^{(n)}) \geq \beta$ and

$$|\varphi_t^{(n)}(y^{(n)}) - \varphi_t(y^{(n)})| > \varepsilon$$

for some $\varepsilon > 0$ and an infinite sequence of values of n . For each of these n one can find three adjoining intervals $\Delta_1^{(n)}, \Delta_2^{(n)}, \Delta_3^{(n)}$ such that $C_{\xi^{(n)}}(y^{(n)}) = \Delta_1^{(n)} \cup \Delta_2^{(n)}, C_{\xi}(y^{(n)}) = \Delta_2^{(n)} \cup \Delta_3^{(n)}$. Let us assume that we have

$$\varphi_t^{(n)}(y^{(n)}) < \varphi_t(y^{(n)}). \tag{4.26}$$

Arguing as above we conclude that either $P_0(\Delta_1^{(n)}) > K_1\beta\varepsilon$, or $P_0(\Delta_3^{(n)}) > K_1\beta\varepsilon$ for some constant K_1 . For definiteness let us assume $P_0(\Delta_1^{(n)}) > K_1\beta\varepsilon$. Then GVP gives

$$C_{\Delta_1^{(n)}}(t) < C_{\Delta_2^{(n)} \cup \Delta_3^{(n)}}(t) = \varphi_t(y^{(n)}) \leq C_{\Delta_2^{(n)}}(t),$$

$$C_{\Delta_1^{(n)}}^{(n)}(t) \geq C_{\Delta_1^{(n)} \cup \Delta_2^{(n)}}^{(n)}(t) = \varphi_t^{(n)}(y^{(n)}) \geq C_{\Delta_2^{(n)}}^{(n)}(t),$$

or

$$C_{\Delta_1^{(n)}}(t) < \varphi_t(y^{(n)}) \leq C_{\Delta_2^{(n)}}(t),$$

$$C_{\Delta_1^{(n)}}^{(n)}(t) \geq \varphi_t^{(n)}(y^{(n)}) \geq C_{\Delta_2^{(n)}}^{(n)}(t).$$

Since $P_0(\Delta_1^{(n)}) > K\beta\varepsilon, P_0(\Delta_2^{(n)}) > \beta$, we have

$$|C_{\Delta_1^{(n)}}^{(n)}(t) - C_{\Delta_1^{(n)}}(t)| \leq o(1),$$

$$|C_{\Delta_2^{(n)}}^{(n)}(t) - C_{\Delta_2^{(n)}}(t)| \leq o(1).$$

This contradicts (4.26).

5. The Gravitationally Interacting Case

In this section we prove the GVP for (1.2) and Theorem 2. Since most of the argument is parallel to the ones in Sects. 2–4, we will only summarize the key steps.

Lemma 5. *The family of partitions $\{\xi_t\}_{t>0}$ is decreasing, i.e. for $0 < t' < t$, $\partial\xi_t \subseteq \partial\xi_{t'}$.*

Proof. Assume to the contrary that there exists $y \in \partial\xi_t$, but $y \notin \partial\xi_{t'}$. Then for some $y^-, y^+, y^- < y < y^+$, we should have

$$\frac{\int_{[y^-, y]}(\eta + t'u_0(\eta))dP_0(\eta)}{\int_{[y^-, y]}dP_0(\eta)} + \frac{(t')^2}{2}(P_0(y, \infty) - P_0(-\infty, y^-))$$

$$\geq \frac{\int_{[y, y^+]}(\eta + t'u_0(\eta))dP_0(\eta)}{\int_{[y, y^+]}dP_0(\eta)} + \frac{(t')^2}{2}(P_0(y^+, \infty) - P_0(-\infty, y)). \tag{5.1}$$

Consider two quadratic functions $q_1(s), q_2(s)$ defined by

$$\begin{aligned}
 q_1(s) &= \frac{\int_{[y^-, y]} (\eta + su_0(\eta)) dP_0(\eta)}{\int_{[y^-, y]} dP_0(\eta)} + \frac{s^2}{2} (P_0[y, \infty) - P_0(-\infty, y^-)) \\
 &= A^- + sB^- + \frac{s^2}{2} C^-, \\
 q_2(s) &= \frac{\int_{[y, y^+]} (\eta + su_0(\eta)) dP_0(\eta)}{\int_{[y, y^+]} dP_0(\eta)} + \frac{s^2}{2} (P_0[y^+, \infty) - P_0(-\infty, y)) \\
 &= A^+ + sB^+ + \frac{s^2}{2} C^+.
 \end{aligned}$$

For sufficiently small s , we have $A^- < A^+$ and $q_1(s) < q_2(s)$, while $q_1(t') \geq q_2(t')$. Since $C^- - C^+ = P_0[y^-, y^+] > 0, A^- - A^+ < 0$, one of the two roots of the equation $q_1(s) = q_2(s)$ is negative. Since there is one positive root in $(0, t')$, no roots exist in (t', ∞) . Therefore we have $q_1(s) > q_2(s)$ for $s > t'$, contradicting the assumption that $q_1(t) < q_2(t)$.

There is also an obvious analog of Lemma 2 for (1.2). We omit the details of that.

Next we study the discrete case, using the notations of Sect. 1. In the present situation, the center of masses of the particles $\{x_i^0, x_{i+1}^0, \dots, x_j^0\}$ becomes

$$C_{i,j}(t) = \frac{\sum_{i \leq l \leq j} m_l^0 (x_l^0 + tv_l^0)}{\sum_{i \leq l \leq j} m_l^0} + \frac{t^2}{2} \left(\sum_{l > j} m_l^0 - \sum_{l < i} m_l^0 \right). \tag{5.2}$$

Lemma 6. *Let $G \in \xi_t, G = \{x_i^0, j' \leq i \leq j''\}$. Then for any $i, j' \leq i \leq j'' - 1$, we have*

$$C_{j',i}(t) \geq C_{i+1,j''}(t). \tag{5.3}$$

Proof. We proceed as in the proof of Lemma 3. The key step is to show that if two groups of particles $\{(x_i^0, v_i^0, m_i^0), j_1 \leq i \leq j\}$ and $\{(x_i^0, v_i^0, m_i^0), j + 1 \leq i \leq j_2\}$ merge at time t' , then

$$C_{j_1, j_2}(s) \leq C_{j_1, j}(s), \quad s \geq t'. \tag{5.4}$$

Consider again two quadratic functions $q_1(s) = C_{j_1, j}(s) = A_1 + B_1s + C_1 \frac{s^2}{2}, q_2(s) = C_{j+1, j_2}(s) = A_2 + B_2s + C_2 \frac{s^2}{2}$. Obviously we have $A_1 - A_2 < 0, C_1 - C_2 = \sum_{i > j} m_i^0 - \sum_{i > j_1} m_i^0 - \sum_{i > j_2} m_i^0 + \sum_{i < j+1} m_i^0 = \sum_{j_1 \leq i \leq j_2} m_i^0 > 0$. Therefore the quadratic equation $q_1(s) - q_2(s) = 0$ has one negative root. Since there is already one root in $(0, t')$, there cannot be any more roots in (t', ∞) . Therefore, we have

$$C_{j_1, j}(s) \geq C_{j+1, j_2}(s), \quad \text{for } s > t'.$$

This implies (5.4).

Similarly we can formulate

Theorem 2'. x_j^0 is the left end point of an element of ξ_i iff

$$\max_{j' < j} C_{j', j-1}(t) < \min_{j \leq j''} C_{j, j''}(t). \tag{5.5}$$

We now prove that the adhesion dynamics of gravitationally interacting particles gives rise to weak solutions of (1.2). Since the proof of (D1') is entirely the same as the one for (D1), we will concentrate on (D2'). Again we only have to prove (D2') for the case when there is no collision in (t_1, t_2) , and at $t = t_2$, a group of particles, say indexed by i_1, \dots, i_2 , collide to form the new particle at $x(t_2)$. Let $h_i = \sum_{j>i} m_j - \sum_{j<i} m_j$, then we have

$$\begin{aligned} \frac{d}{d\tau} v_i(\tau) &= h_i, \\ \int_{t_1}^{t_2} d\tau \int g'(\eta) u(\eta, \tau) dI_\tau(\eta) &+ \int_{t_1}^{t_2} d\tau \int g(\eta) (P_\tau(\eta, \infty) - P_\tau(-\infty, \tau)) dP_\tau(\eta) \\ &= \int_{t_1}^{t_2} d\tau \sum_{i=1}^N g'(x_i(\tau)) v_i^2(\tau) m_i + \int_{t_1}^{t_2} d\tau \sum_{i=1}^N g(x_i(\tau)) h_i m_i \\ &= \int_{t_1}^{t_2} d\tau \sum_{i=1}^N m_i v_i(\tau) \frac{d}{d\tau} g(x_i(\tau)) + \int_{t_1}^{t_2} d\tau \sum_{i=1}^N g(x_i(\tau)) h_i m_i \\ &= \sum_{i=1}^N g(x_i(t_2)) m_i v_i(t_2) - \sum_{i=1}^N g(x_i(t_1)) m_i v_i(t_1) \\ &\quad - \int_{t_1}^{t_2} d\tau \sum_{i=1}^N m_i h_i g(x_i(\tau)) + \int_{t_1}^{t_2} d\tau \sum_{i=1}^N g(x_i(\tau)) h_i m_i \\ &= \int g(\eta) dI_{t_2}(\eta) - \int g(\eta) dI_{t_1}(\eta). \end{aligned}$$

To pass to the continuum limit, we follow the same steps as in the proof of Theorem 1. Some care has to be taken at the boundary of $\text{supp}\{P_0\}$ since $\rho'_0 > \text{const} > 0$ no longer holds. This can be done by making the set $\{x = \rho'_0(x) < \text{const}\}$ small by choosing the constant small.

6. The Modified Gravitational Model

In this section we study a variant of (1.2) for which some other weighted velocity, instead of momentum, is conserved.

$$\begin{cases} \rho_t + (\rho u)_x = 0 \\ u_t + (\frac{u^2}{2})_x = -g_x \\ g_{xx} = \rho. \end{cases} \tag{6.1}$$

This model, or more precisely its multidimensional version, is an approximation of the basic equations in astrophysics for the dynamic evolution of matter in which effects of Hubble expansion are neglected [P, GS]. For this system, we will prove

not only existence, but also *uniqueness* of solutions for the initial value problem. We will also obtain precise information on the characteristics of this system. This information will be used in the second part of this paper to study the behavior of solutions with random initial data.

As in Sect. 5, we assume that $\rho(x, 0) = \rho_0(x) \geq 0$, $u(x, 0) = u_0(x)$. Moreover, we will assume $\rho_0 \in L^1$, and (A3) holds. Without loss of generality, we let $\int \rho_0 dx = 1$.

Let $h(x, t) = \int_{-\infty}^x \rho(\eta, t) d\eta - \frac{1}{2}$. Obviously, for any reasonable solution of (6.1), we should have $h(\pm\infty, t) = \pm 1/2$.

In the following we will deal with solutions having the property that for any $t > 0$, $u(\cdot, t)$ is of bounded variation on any compact interval; $\rho(\cdot, t)$ is a non-negative Radon measure. We interpret the product $u\rho$ in the sense of Volpert [V]:

$$u\rho = \bar{u}\rho = \bar{u}h_x, \tag{6.2}$$

where \bar{u} is the symmetric mean of u : $\bar{u}(x, t) = u(x, t)$ if (x, t) is a point of approximate continuity of u . At jump points, $\bar{u}(x, t) = \frac{1}{2}(l_+u(x, t) + l_-u(x, t))$, where $l_+u(x, t)$ and $l_-u(x, t)$ are the limits of u at (x, t) from the two sides of the jump. Given this, we can rewrite (6.1) in an equivalent form,

$$\begin{cases} h_t + \bar{u}h_x = 0 \\ u_t + (\frac{u^2}{2})_x = -h, \end{cases} \tag{6.3}$$

where the first equation in (6.3) is understood in the sense of measures. We let $h_0(y) = h(y, 0)$. We assume $\rho_0, u_0 \in C^0$.

A special property of (6.3) is that for smooth solutions, the characteristics can be written down very easily:

$$\begin{cases} x = y + t u_0(y) - \frac{t^2}{2} h_0(y) \\ h(x, t) = h_0(y) \\ u(x, t) = u_0(y) - t h_0(y). \end{cases} \tag{6.4}$$

From (6.4), we get

$$x - u(x, t)t = y + \frac{t^2}{2} h_0(y). \tag{6.5}$$

Since the right-hand side of (6.5) is a strictly increasing function of y , we can invert it to get

$$y = G(x - u(x, t)t, t), \tag{6.6}$$

where $G(z, t)$ solves

$$z = G + \frac{t^2}{2} h_0(G). \tag{6.7}$$

Therefore, we have

$$h(x, t) = h_0(G(x - u(x, t)t, t)). \tag{6.8}$$

This suggests the following:

Lemma 7. *Let (h, u) be a solution of (6.3) such that u is locally BV on $R^1 \times (0, \infty)$. Then u is also a solution of*

$$u_t + \left(\frac{u^2}{2}\right)_x = -h_0(G(x - ut, t)). \tag{6.9}$$

Conversely, if u is a solution of (6.9) with locally bounded variations on $R^1 \times (0, \infty)$, and let $h(x, t) = h_0(G(x - ut, t))$, then (h, u) is also a solution of (6.3).

Proof. In the proof we will use $\hat{f}(u)$ to denote the functional composition of f and u : $\hat{f}(u) = \int_0^1 f(\alpha l_- u + (1 - \alpha)l_+ u) d\alpha$.

Let (h, u) be a solution of (6.3) such that u is locally BV on $R^1 \times (0, \infty)$. Using BV calculus [V], we see that u and h also satisfy

$$u_t + \bar{u}u_x = -h, \quad h_t + \bar{u}h_x = 0 \tag{6.10}$$

in the sense of measures.

Let $g(x, t) = h_0(G(x - ut, t))$. Again using BV calculus, we get

$$\begin{aligned} g_t + \bar{u}g_x &= \hat{h}'_0 \{G_t + \hat{G}_z(-u - u_t t) + \bar{u}\hat{G}_z(1 - u_x t)\} \\ &= \hat{h}'_0 \{G_t - t\hat{G}_z(u_t + \bar{u}u_x)\} = \hat{h}'_0 \{G_t + t\hat{G}_z h\}. \end{aligned}$$

The final expression only involves locally bounded functions. Therefore it can be simply written as $h'_0 \{G_t + tG_z h\}$. Since

$$G_t(z, t) = -\frac{th_0(G)}{1 + \frac{t^2}{2}h'_0(G)}, \quad G_z(z, t) = \frac{1}{1 + \frac{t^2}{2}h'_0(G)}, \tag{6.11}$$

we obtain

$$g_t + \bar{u}g_x = -\frac{h'_0(G)t}{1 + \frac{t^2}{2}h'_0(G)} \{g - h\}.$$

Let $e = g - h$. We have

$$e_t + \bar{u}e_x = -\frac{h'_0(G)t}{1 + \frac{t^2}{2}h'_0(G)} e$$

and $e(x, 0) = 0$.

This implies that $e \equiv 0$. Hence $h(x, t) = h_0(G(x - ut, t))$, i.e. u is a solution of (6.9).

Conversely assume u is solution of (6.9) with locally bounded variation. We let $h(x, t) = h_0(G(x - ut, t))$. Then a similar calculation shows that h satisfies

$$h_t + \bar{u}h_x = 0.$$

Hence (h, u) solves (6.3).

Now the problem is reduced to (6.9) which is a scalar equation. We are ready to prove:

Theorem 3. *Assume that ρ_0 and u_0 are continuous and satisfy (A3). Then there exists a weak solution (h, u) of (6.3) with initial data (h_0, u_0) such that*

$$u_x < \frac{1}{t} \tag{6.12}$$

in the distributional sense, and u is bounded locally and has locally bounded variations in $R^1 \times (0, \infty)$. Furthermore, such solutions are unique.

Proof. This theorem follows more or less from standard arguments on conservation laws [V], with some care on cutting off the growth at infinity. Here we will not repeat this argument. We will only indicate how (6.12) can be derived from the viscosity approximation.

Consider

$$u_t^\varepsilon + \left(\frac{(u^\varepsilon)^2}{2} \right)_x = -h_0(G(x - u^\varepsilon t, t)) + \varepsilon u_{xx}^\varepsilon. \tag{6.13}$$

Let $\xi = u_x^\varepsilon$. Differentiating both sides of (6.13), we get

$$\xi_t + u \xi_x + \xi^2 = -h_0' G_z(1 - \xi t) + \varepsilon \xi_{xx}. \tag{6.14}$$

The standard maximum principle implies that, if we let $f(t) = \sup_x \xi(x, t)$, then

$$\frac{d}{dt} f \leq h_0' G_z(f t - 1) - f^2.$$

Since $h_0' G_z \geq 0$, we obtain

$$f < \frac{1}{t}.$$

Taking the limit as $\varepsilon \rightarrow 0$, we get (6.12).

Next we summarize the main results on the generalized characteristics of (6.9), drawing mainly from the work of Dafermos [D]. A Lipschitz curve $\xi(\cdot) : [a, b] \rightarrow R^1$ is called a generalized characteristic if $\xi(t) \in [u(\xi(t)+, t), u(\xi(t)-, t)]$, for a.e. $t \in [a, b]$. It is called a genuine characteristic if $u(\xi(t)-, t) = u(\xi(t)+, t)$, for a.e. $t \in [a, b]$. The following holds if u is locally BV on $R^1 \times (0, \infty)$:

(1) If ξ is a generalized characteristic on $[a, b]$, then for a.e. $t \in [a, b]$

$$\dot{\xi}(t) = \begin{cases} u(\xi(t)\pm, t), & \text{if } u(\xi(t)+, t) = u(\xi(t)-, t) \\ \frac{1}{2}(u(\xi(t)+, t) + u(\xi(t)-, t)) = \bar{u}(\xi(t)), & \text{if not} \end{cases}$$

(2) Given $(x, t) \in R^1 \times (0, \infty)$, there exists a funnel of backward generalized characteristics from (x, t) to the initial time $t = 0$. Among them, there exists a unique minimal and a unique maximal one. They are both genuine.

(3) Let ξ be a genuine characteristic. Then by modifying on a measure zero set, one can extend $u(\xi(t)-, t) \stackrel{\text{a.e.}}{=} u(\xi(t)+, t)$ to a C^1 function $\bar{u}(t)$, such that

$$\frac{d\xi(t)}{dt} = \bar{u}(t), \quad \frac{d\bar{u}(t)}{dt} = h_0(G(\xi(t) - \bar{u}(t)t, t)). \tag{6.15}$$

In particular, if (\bar{x}, \bar{t}) is point of continuity, i.e. $u(\bar{x}+, \bar{t}) = u(\bar{x}-, \bar{t})$, then from (\bar{x}, \bar{t}) , there is a unique backward generalized characteristic. Moreover, (6.15) holds.

Recall that $h(x, t) = h_0(G(x - u(x, t)t, t))$. Along a genuine characteristic, we have

$$\begin{aligned} \frac{d}{dt} h(\xi(t), t) &= h_0' \left\{ G_t + G_z \left(\frac{d\xi}{dt} - \bar{u} - t \frac{d\bar{u}}{dt} \right) \right\} \\ &= h_0' \{ G_t + t h_0(G) G_z \} \\ &= 0. \end{aligned}$$

Part II. Behavior with Random Initial Data

7. Behavior of Solutions with Random Initial Data for Systems (1.1) and (1.2)

This part is devoted to the study of the qualitative behavior of solutions constructed above when the initial data are random. We prove that the systems (1.1), (1.2), and (1.3) share the following common feature: For continuous but nowhere differentiable random initial data, almost surely, the solution $u(x, t)$ becomes discontinuous for any $t > 0$, the set of discontinuities (shocks) is dense (on the appropriate intervals), and almost all masses are absorbed by the shocks. In particular for any $t > 0$, $\rho(\cdot, t)$ becomes a discrete measure even though it may have a smooth distribution at $t = 0$. Such behavior was, to some extent, anticipated by Zeldovich [Z], as part of his proposal for the formation of large scale structures in the universe. Technically, we extend some of the arguments in [S].

Let Q be the probability distribution of initial velocities, $u(x; 0) = u_0(x)$. We assume

(Q1) Q is defined on the Borel σ -algebra of the space of continuous functions on R^1 , and for a.e. u_0 (with respect to Q),

$$\lim_{|x| \rightarrow \infty} \frac{u_0(x)}{|x|} = 0.$$

(Q2) For a.e. u_0 , the following holds: for any $\eta_0 \in R^1$,

$$\begin{aligned} \limsup_{h \rightarrow 0+} \frac{\int_{\eta_0-h}^{\eta_0} (u(\eta; 0) - u(\eta_0; 0)) d\eta}{h^2} &= \infty, \\ \liminf_{h \rightarrow 0+} \frac{\int_{\eta_0}^{\eta_0+h} (u(\eta; 0) - u(\eta_0; 0)) d\eta}{h^2} &= -\infty. \end{aligned}$$

Concerning ρ_0 , the initial distribution of masses, we will assume either

$$(1) \quad \rho_0 \in C^1(R^1), \quad 0 < \text{const} \leq \rho_0 \leq \text{const} < \infty, \quad |\rho'| \leq \text{const} < \infty$$

or

$$(2) \quad \rho_0 \in C^1(R^1), \quad \rho_0(x) > 0 \text{ for } x \in (a, b), \text{ and } \rho_0(x) = 0, \text{ for } x \notin (a, b), \\ |\rho'_0(x)| \leq \text{const} < \infty. \text{ For (1.2) we will only consider the case (2).}$$

Theorem 4. For a.e. u_0 , the weak solutions of (1.1) and (1.2) constructed by GVP have the following properties: For any $t > 0$,

(a) The measure P_t is a pure point measure, i.e. $P_t(x) = \sum_i m_i \delta(x - x_i)$, ($m_i > 0$).

(b) The closure of the set $\{x_i\}$ is either R^1 in the case of (1), or a closed interval $I(t)$ in the case of (2).

Proof. We will only prove this theorem for (1.1). The modifications required for (1.2) are obvious.

We will show that for a.e. u_0 , the partition ξ_t given by GVP has the following properties:

- i) There exists a countable set of elements of ξ_t which are intervals of positive length.
- ii) The union of these intervals is a set of full measure
- iii) Between any two of these intervals there are infinitely many similar intervals.

Clearly (i)–(iii) implies the theorem.

Fix $y_0 \neq a, b$ and take $u(\cdot; 0)$ satisfying (Q1). Assume that y_0 is the left end point of an element of ξ_t , i.e. for any $h_1, h_2 > 0$,

$$\frac{\int_{y_0-h_1}^{y_0} (\eta + t u(\eta; 0)) \rho_0(\eta) d\eta}{\int_{y_0-h_1}^{y_0} \rho_0(\eta) d\eta} \leq \frac{\int_{y_0}^{y_0+h_2} (\eta + t u(\eta; 0)) \rho_0(\eta) d\eta}{\int_{y_0}^{y_0+h_2} \rho_0(\eta) d\eta}. \tag{7.1}$$

We will show that the Q-probability of such $u(\cdot; 0)$ is zero. Indeed we can rewrite (7.1) as

$$\begin{aligned} & \frac{\int_{y_0-h_1}^{y_0} (\eta - y_0) \rho_0(\eta) d\eta}{\int_{y_0-h_1}^{y_0} \rho_0(\eta) d\eta} + t \frac{\int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) \rho_0(\eta) d\eta}{\int_{y_0-h_1}^{y_0} \rho_0(\eta) d\eta} \\ & \leq \frac{\int_{y_0}^{y_0+h_2} (\eta - y_0) \rho_0(\eta) d\eta}{\int_{y_0}^{y_0+h_2} \rho_0(\eta) d\eta} + t \frac{\int_{y_0}^{y_0+h_2} (u(\eta; 0) - u(y_0; 0)) \rho_0(\eta) d\eta}{\int_{y_0}^{y_0+h_2} \rho_0(\eta) d\eta}. \end{aligned} \tag{7.2}$$

Since $\rho_0(y_0) > 0$, we have

$$\begin{aligned} & \frac{\int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) \rho_0(\eta) d\eta}{\int_{y_0-h_1}^{y_0} \rho_0(\eta) d\eta} = \frac{\rho_0(y_0) \int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) d\eta}{\rho_0(y_0) h_1 + \int_{y_0-h_1}^{y_0} (\rho_0(\eta) - \rho_0(y_0)) d\eta} \\ & + \frac{\int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) (\rho_0(\eta) - \rho_0(y_0)) d\eta}{\rho_0(y_0) h_1 + \int_{y_0-h_1}^{y_0} (\rho_0(\eta) - \rho_0(y_0)) d\eta} \\ & = \frac{\int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) d\eta}{h_1} (1 + o(1)) + O(h_1). \end{aligned} \tag{7.3}$$

Similar estimates hold for the integrals on $(y_0, y_0 + h_2)$. Therefore, we can write (7.1) as

$$\begin{aligned} & \frac{\int_{y_0-h_1}^{y_0} (u(\eta; 0) - u(y_0; 0)) d\eta}{h_1} (1 + o(1)) + O(h_1) \\ & \leq \frac{\int_{y_0}^{y_0+h_2} (u(\eta; 0) - u(y_0; 0)) d\eta}{h_2} (1 + o(1)) + O(h_2). \end{aligned} \tag{7.4}$$

From (Q2), this can only hold on a Q-measure 0 set of $u(\cdot; 0)$.

Next we consider the measurable space of $(u(\cdot; 0), y_0)$ equipped with the product measure of Q and the Lebesgue measure on R^1 . The Fubini theorem implies that for Q – a.e. $u(\cdot; 0)$ the set of y_0 satisfying (7.4) has Lebesgue measure zero. This proves (ii). Since two different intervals in ξ_t cannot have a common end point,

there have to be infinitely many such intervals in between. Finally, since φ_t is a continuous nondecreasing function, and φ_t is constant on each of these intervals, the range of φ_t from these intervals has to be dense on its full range.

8. Behavior of Solutions with Random Initial Data for (1.3)

In this section we prove an analogous result of Theorem 4 for the nonconservative model (1.3). We use mainly the method of characteristics, an approach potentially more general than the one used in Sect. 7. For the initial data we will assume that $\rho_0 \in C^1 \cap L^1(R^1)$, and $u(\cdot; 0) = u_0(\cdot)$ is random with distribution \mathcal{Q} , satisfying (Q1), whereas (Q2) is replaced by a more stringent condition:

(Q2') There exists a function $\omega(h)$, such that for almost a.e. u_0 ,

$$\begin{aligned} \limsup_{|h| \rightarrow 0} \frac{u_0(y+h) - u_0(y)}{\omega(h)} &= 1, \\ \liminf_{|h| \rightarrow 0} \frac{u_0(y+h) - u_0(y)}{\omega(h)} &= -1, \end{aligned} \tag{8.1}$$

holds for all $y \in R^1$, and the convergence is uniform on compact sets. Moreover, ω satisfies

(1) $\lim_{|h| \rightarrow 0} \omega(h) = 0, \quad \lim_{|h| \rightarrow 0} \frac{\omega(h)}{h} = +\infty.$

(2) There exists a constant $\varepsilon > 0$, such that for any $h > 0$, one can find $h_1 > 0$, such that $\omega(h_1) < \frac{1}{4}\omega(h)$ and $h_1 > \varepsilon h$.

(2) essentially says that ω has algebraic behavior near 0. Obviously, the Wiener process satisfies (Q2') with $\omega(h) = (2h \log \frac{1}{h})^{1/2}$.

Theorem 5. For a.e. u_0 , the weak solution of (1.3) constructed in Sect. 6 has the following properties: for any $t > 0$, let $\{x_i\}$ be the set of shock locations at time t . Then

(a) The measure $\rho(\cdot, t)$ is a pure point measure: $\rho(x, t) = \sum_i m_i \delta(x - x_i)$, $m_i \geq 0$.

(b) $\{x_i\}$ is dense in R^1 .

Proof. Take $u(x; 0)$ such that (Q1) and (Q2') are satisfied and fix any $t^* > 0$. We first show that the set of shock locations is dense in R^1 . Assume to the contrary that there exists an interval Δ on which $u(\cdot, t^*)$ is continuous. Then for any $x \in \Delta$ there is a unique backward characteristic $\xi : [0, t^*] \rightarrow R^1$ connecting x to a point y at $t = 0$. Furthermore, along ξ we have

$$\frac{d}{dt} \xi(t) = \bar{u}(t), \quad \frac{d}{dt} h(\xi(t), t) = 0, \quad \frac{d}{dt} \bar{u}(t) = -h(\xi(t), t). \tag{8.2}$$

Thus ξ is given by

$$\xi(t) = y + t u_0(y) - h_0(y) \frac{t^2}{2}, \quad \xi(t^*) = x. \tag{8.3}$$

The mapping from y to x is obviously continuous. It is also one-to-one since the characteristics do not intersect between 0 and t^* . Therefore, the inverse mapping

from x to y is also well-defined and continuous. Hence, there is a interval $J \in R^1$ at $t = 0$ on which the forward characteristics, also defined by (8.3), do not intersect between $t = 0$ and $t = t^*$, i.e. if we take $y_1, y_2 \in J, y_1 < y_2$, then we should also have for $0 \leq t \leq t^*$,

$$y_1 + tu_0(y_1) - \frac{t^2}{2}h_0(y_1) < y_2 + tu_0(y_2) - \frac{t^2}{2}h_0(y_2)$$

or

$$t(u_0(y_2) - u_0(y_1)) > y_1 - y_2 - \frac{t^2}{2}(h_0(y_1) - h_0(y_2)) \geq -C(y_2 - y_1) \tag{8.4}$$

for some constant C . This clearly violates (Q2'). Hence we have proved that the set of shock locations is dense in R^1 .

Let $x_0 \in \{x_i\}$. From x_0 we draw the minimal and maximal backward characteristics. Assume that they arrive at y_1 and y_2 respectively at $t = 0$. These characteristics are classical in the sense that (8.2) is satisfied along them. Therefore, we see that $h(\cdot, t^*)$ is discontinuous at x_0 , and

$$h(x_{0+}, t^*) = h_0(y_2), \quad h(x_{0-}, t^*) = h_0(y_1). \tag{8.5}$$

In other words, the masses initially distributed in (y_1, y_2) are absorbed into x_0 at $t = t^*$.

Let $S(u_0) = \{y : \text{There is a genuine (or classical) characteristic connecting } y \text{ at } t = 0 \text{ to some } x \text{ at } t = t^*\}$. We want to prove that the Lebesgue measure of $S(u_0)$ is zero.

Let $y_0 \in S(u_0)$. If $y \in S(u_0), y > y_0$, then we must have

$$y + tu(y; 0) - \frac{t^2}{2}h_0(y) > y_0 + tu(y_0; 0) - \frac{t^2}{2}h_0(y_0) \tag{8.6}$$

for $0 < t < t^*$. Equation (8.6) implies that

$$u(y; 0) - u(y_0; 0) > -C_0(y - y_0) \tag{8.7}$$

for some positive constant C_0 . Choose y_1 sufficiently close to $y_0, y_1 > y_0$ such that

$$(i) \quad u(y_1; 0) - u(y_0; 0) < -\frac{1}{2}\omega(y_1 - y_0),$$

$$(ii) \quad |u(y; 0) - u(y_1; 0)| < 2\omega(y_1 - y)$$

for $y \in (y_0, y_1)$. Then

$$u(y; 0) < 2\omega(y_1 - y) + u(y_1; 0) < 2\omega(y_1 - y) - \frac{1}{2}\omega(y_1 - y_0) + u(y_0; 0).$$

Next we choose $y_2 \in (y_0, y_1)$ such that

$$\omega(y_1 - y_2) < \frac{1}{4}\omega(y_1 - y_0)$$

and $y_1 - y_2 \geq \varepsilon(y_1 - y_0)$. For $y \in (y_2, y_1)$ we have

$$\begin{aligned} u(y; 0) - u(y_0; 0) &< 2\omega(y_1 - y_2) - \frac{1}{2}\omega(y_1 - y_0) \\ &< -\frac{1}{4}\omega(y_1 - y_0) < -C_0(y_1 - y_0) \leq -C_0(y - y_0) \end{aligned} \tag{8.8}$$

if y_1 is sufficiently close to y_0 . Comparing (8.8) with (8.7) we see that $(y_2, y_1) \notin S(u_0)$. Therefore, we have

$$\frac{|S(u_0) \cap [y_0, y_1]|}{|[y_0, y_1]|} \leq 1 - \varepsilon .$$

Here we used $|A|$ to denote the Lebesgue measure of the set A . Hence y_0 is not a point of density of $S(u_0)$. Since this is true for all $y_0 \in S(u_0)$, we have $|S(u_0)| = 0$.

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Appendix: Contact Points on the Convex Hull

In this appendix we summarize some properties of the contact points on the convex hull of $\phi_1 \circ \phi_2^{-1}$, extending the discussions in Sect. 2. We will assume that P_0 has a density $\rho_0(x), u_0$ is continuous.

Definition A1. *A point z is called special if*

$$\tilde{c}(z, z) = \min_{z^+ \geq z} \tilde{c}(z, z^+) = \max_{z^- \leq z} \tilde{c}(z^-, z) = \tilde{c}(z, z) .$$

Lemma A1. *The set of special points is closed.*

Proof. Let z_n be a sequence of special points, $z_n \rightarrow z$. We claim that z is a special point. Assume to the contrary, then for some $z^- < z < z^+$,

$$\tilde{c}(z^-, z) > \tilde{c}(z, z^+) .$$

But $\tilde{c}(z^-, z) = \lim_{n \rightarrow \infty} \tilde{c}(z^-, z_n)$, $\tilde{c}(z, z^+) = \lim_{n \rightarrow \infty} \tilde{c}(z_n, z^+)$. Therefore, for sufficiently large n ,

$$\tilde{c}(z^-, z_n) > \tilde{c}(z_n, z^+) ,$$

which contradicts our assumption that z_n is a special point.

Definition A2. *A special point z is called single if $\tilde{c}(z, z^+) > \tilde{c}(z^-, z)$ for any $z^- < z < z^+$. A special point z is called multiple if there exists either $z^+ > z$ for which $\tilde{c}(z^+, z) = \tilde{c}(z, z)$, or $z^- < z$ for which $\tilde{c}(z^-, z) = \tilde{c}(z, z)$.*

For every z consider a segment $[z_0^-, z_0^+]$, where $z_0^- = \min\{y \leq z : \tilde{c}(y, z) = \tilde{c}(z, z)\}$, $z_0^+ = \max\{y \geq z : \tilde{c}(z, y) = \tilde{c}(z, z)\}$. If a point z is multiple then $[z_0^-, z_0^+]$ has positive length.

Lemma A2. *Take $y \in (z, z_0^+)$ where $\tilde{c}(z, y) > \tilde{c}(z, z)$. Then y is not a special point.*

Proof. We have

$$\tilde{c}(z, z_0^+) = p_1 \tilde{c}(z, y) + p_2 \tilde{c}(y, z_0^+) ,$$

where

$$p_1 = \frac{\int_z^y \rho_0(\eta) d\eta}{\int_z^{z_0^+} \rho_0(\eta) d\eta}, \quad p_2 = \frac{\int_y^{z_0^+} \rho_0(\eta) d\eta}{\int_z^{z_0^+} \rho_0(\eta) d\eta}.$$

Since $\tilde{c}(z, y) > \tilde{c}(z, z_0^+)$, we conclude that $\tilde{c}(y, z_0^+) < \tilde{c}(z, y)$. This contradicts the definition of special point.

In the same way one proves the following lemma.

Lemma A2'. *If $y \in (z_0^-, z)$ and $\tilde{c}(y, z) < \tilde{c}(z, z)$ then y is not a special point.*

Lemma A3. *Take any point $z' \in [z, z_0^+]$ where $\tilde{c}(z, z') = \tilde{c}(z, z)$, then z' is a special point.*

Proof. We shall consider two cases.

1) $z < z^- < z', z' < z^+$. As above, we have

$$\tilde{c}(z, z) = \tilde{c}(z, z') = p_1 \tilde{c}(z, z^-) + p_2 \tilde{c}(z^-, z') = \min_{y \geq z} \tilde{c}(z, y)$$

for some $p_1, p_2 \geq 0, p_1 + p_2 = 1$, and $\tilde{c}(z, z') \leq \tilde{c}(z, z^-)$. We conclude that $\tilde{c}(z^-, z') \leq \tilde{c}(z, z')$. Also for some $p'_1, p'_2 \geq 0, p'_1 + p'_2 = 1$,

$$\tilde{c}(z, z^+) = p'_1 \tilde{c}(z, z') + p'_2 \tilde{c}(z', z^+) \geq \tilde{c}(z, z').$$

We conclude that $\tilde{c}(z', z^+) \geq \tilde{c}(z, z')$, i.e.

$$\tilde{c}(z^-, z') \leq \tilde{c}(z, z') \leq \tilde{c}(z', z^+).$$

2) $z^- < z, z' < z^+$.

We have for some $p''_1, p''_2 \geq 0, p''_1 + p''_2 = 1$,

$$\tilde{c}(z^-, z') = p''_1 \tilde{c}(z^-, z) + p''_2 \tilde{c}(z, z').$$

Since $\tilde{c}(z, z') = \tilde{c}(z, z)$, $\tilde{c}(z^-, z) \leq \tilde{c}(z, z')$, we find that $\tilde{c}(z^-, z') \leq \tilde{c}(z, z')$, while we have already shown that $\tilde{c}(z', z^+) \geq \tilde{c}(z, z')$.

The lemma is proven.

Lemma A3'. *Take any point $z' \in [z_0^-, z)$, where $\tilde{c}(z', z) = \tilde{c}(z, z)$. Then z' is a special point.*

Take the segment $[z_0^-, z_0^+]$ constructed for the point z . We shall denote the dependence on z as $[z_0^-, z_0^+]_z$.

Lemma A.4. *If z' is a special point $z' \in [z_0^-, z_0^+]_z$, then $[z_0^-, z_0^+]_{z'} = [z_0^-, z_0^+]_z$.*

Proof. We claim that $\tilde{c}(z, z) = \tilde{c}(z', z')$. Indeed, $\tilde{c}(z, z') = \tilde{c}(z, z)$ according to Lemmas 2, 2', 3, 3'. So the function $\tilde{c}(z, y)$ as a function of y attains its extremum on the semi-line at z' . z' is an interior point of the semi-line. So $\frac{\partial}{\partial y} \tilde{c}(z, y)|_{y=z'} = 0$. This gives

$$[z' + tu_0(z')] \int_z^{z'} \rho_0(\eta) d\eta = \int_z^{z'} (\eta + tu_0(\eta)) \rho_0(\eta) d\eta,$$

and so $\tilde{c}(z', z') = \tilde{c}(z', z) = \tilde{c}(z, z)$.

Suppose $z' < z$. We have, for some $p_1, p_2 \geq 0$, $p_1 + p_2 = 1$,

$$\tilde{c}(z, z) = \tilde{c}(z_0^-, z) = p_1 \tilde{c}(z_0^-, z') + p_2 \tilde{c}(z', z) = p_1 \tilde{c}(z_0^-, z') + p_2 \tilde{c}(z, z),$$

$$\tilde{c}(z_0^-, z') = \tilde{c}(z, z) = \tilde{c}(z', z').$$

If we have $z^- < z_0^-$ such that $\tilde{c}(z^-, z') = \tilde{c}(z', z')$, then for some $p'_1, p'_2 \geq 0$, $p'_1 + p'_2 = 1$,

$$\tilde{c}(z^-, z) = p'_1 \tilde{c}(z^-, z') + p'_2 \tilde{c}(z', z) = p'_1 \tilde{c}(z', z) + p'_2 \tilde{c}(z, z) = \tilde{c}(z, z).$$

That contradicts our assumption that z_0^- is the smallest with such a property. The analogous calculations can be done with z_0^+ .

The case $z < z'$ can be analyzed in a similar way.

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