

Variational Characterization of the Speed of Propagation of Fronts for the Nonlinear Diffusion Equation

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Abstract: We give an integral variational characterization for the speed of fronts of the nonlinear diffusion equation $u_t = u_{xx} + f(u)$ with $f(0) = f(1) = 0$, and $f > 0$ in $(0, 1)$, which permits, in principle, the calculation of the exact speed for arbitrary f .

1. Introduction

The problem of the asymptotic speed of propagation of the interface between an unstable and stable state has received much attention in connection with different problems of population growth, chemical reactions, pattern formation and others. We refer to [1] for a recent review and references. The best understood of such problems is that of the nonlinear reaction diffusion equation

$$u_t = u_{xx} + f(u) \tag{1a}$$

with

$$f(0) = f(1) = 0, \quad f'(0) > 0 \quad \text{and} \quad f > 0 \quad \text{in} \quad (0, 1) \tag{1b}$$

for which Aronson and Weinberger [2] have shown that any positive sufficiently localized initial condition $u(x, 0)$ evolves into a front that joins the stable state $u = 1$ to $u = 0$. The asymptotic speed at which the front propagates is the minimal speed c^* for which there is a monotonic front joining $u = 1$ to $u = 0$. Moreover they show that the selected speed is bounded above and below by

$$2\sqrt{f'(0)} \leq c^* \leq 2 \sup \left\{ \sqrt{\frac{f(u)}{u}} \mid u \in (0, 1) \right\} \tag{2}$$

and that the asymptotic selected front approaches the fixed point $u = 0$ exponentially with slope

$$m = -\frac{1}{2}(c^* + \sqrt{c^{*2} - 4f'(0)}). \tag{3}$$

The lower bound $c_L \equiv 2\sqrt{f'(0)}$ is referred to in the literature as the linear value of the speed. For concave functions f , the upper and lower bounds coincide and the speed is exactly c_L [2,3]. However, the asymptotic speed of propagation can still be the linear value even when the upper and lower bounds do not coincide as explicit examples and a variational characterization [4] which provides improved upper bounds show. We have recently obtained an improved lower bound [5] on the speed of the front that enables one to decide when the selected speed is greater than the linear value case which is referred to as nonlinear marginal stability selection. There have been several reformulations of Aronson's and Weinberger's rigorous results for the nonlinear diffusion equation aiming to their heuristic extension to other higher order and pattern forming equations [6–9]. None of these approaches however provide the means to calculate a priori the velocity of the fronts.

The purpose of the present work is to extend our previous result [5] to show an integral variational characterization of the speed of the fronts of Eq. (1) which enables, in principle, its exact calculation for arbitrary f . Our main result is the following

Theorem (Variational Characterization of c^*). *Let $f \in C^1[0, 1]$ with $f(0) = f(1) = 0$, $f'(0) > 0$ and $f(u) > 0$ for $u \in (0, 1)$. Then*

$$c^* = J \equiv \sup\{I(g)|g \in \mathcal{D}\}. \tag{4}$$

Here

$$I(g) = 2 \frac{\int_0^1 \sqrt{fgh} du}{\int_0^1 g du}, \tag{5}$$

and \mathcal{D} is the space of functions in $C^1(0, 1)$ such that $g \geq 0$, $h \equiv -g' > 0$ in $(0, 1)$, and $\int_0^1 g(u) du < \infty$. Moreover, if $c^* \neq 2\sqrt{f'(0)}$, J is attained at some $\hat{g} \in \mathcal{D}$, and \hat{g} is unique up to a multiplicative constant.

Remark. The function \hat{g} , at which the supremum is attained when $c^* \neq 2\sqrt{f'(0)}$ actually satisfies $\hat{g}(1) = 0$ and $\hat{g}(q) \approx q^{1-\frac{c^*}{|m|}}$ near $q = 0$, where $|m|$ is given in terms of c^* and $f'(0)$ by Eq. (9a) below. In Sect. 2 we prove the theorem, and in Sect. 3 an example is given.

2. Proof of the Variational Characterization

We are interested in the calculation of the minimal speed for which Eq. (1) has a monotonic travelling front $u(x, t) = q(z)$ with $z = x - ct$ joining $u = 1$ to $u = 0$. Since the selected speed corresponds to that of a decreasing monotonic front, it is convenient to work in phase space. Calling $p(q) = -dq/dz$, where the minus sign is included so that p is positive, we find that the monotonic fronts are solutions of

$$p(q) \frac{dp}{dq} - cp(q) + f(q) = 0, \tag{7}$$

with

$$p(0) = 0, \quad p(1) = 0, \quad p > 0 \quad \text{in } (0, 1). \tag{8}$$

As shown by Aronson and Weinberger [2], the asymptotic speed of propagation of fronts of Eq. (1), c^* , is the minimum value c for which there is a solution of (7)

and (8). Aronson and Weinberger have proved that there is a (unique) p satisfying (7) and (8) for $c = c^*$ (see [2], Sect. 4). Moreover, the solution p is such that $p(q) \sim |m|q$ near $q = 0$, where $|m|$ is the largest root of the equation

$$x^2 - c^*x + f'(0) = 0, \quad (9)$$

i.e.,

$$|m| = \frac{1}{2}(c^* + \sqrt{c^{*2} - 4f'(0)}). \quad (9a)$$

We find it convenient to introduce the parameter λ defined as $\lambda = c^*/|m|$. In terms of λ one can write

$$c^* = \lambda \sqrt{\frac{f'(0)}{\lambda - 1}} \quad \text{and} \quad |m| = \sqrt{\frac{f'(0)}{\lambda - 1}}. \quad (10)$$

It is straightforward to verify that whenever $1 < \lambda < 2$ the value of $|m|$ given by (10) corresponds to the largest root of (9) and therefore to the asymptotic slope at the origin of the selected front [10]. At $\lambda = 2$ the speed c^* attains the linear value c_L .

Proof of the Theorem. First we show that

$$c^* \geq J, \quad (11)$$

and then we prove that the equality actually holds in (11). To prove (11) we need only to show that

$$c^* \geq I(g)$$

for all $g \in \mathcal{D}$. This latter fact has been proven by us in [5]. We repeat here the argument for completeness. Let g be any function in \mathcal{D} . Multiplying Eq. (7) by g/p and integrating with respect to q we find after integrating by parts,

$$c^* = \frac{\int_0^1 \left(hp + \frac{f(q)}{p} g \right) dq}{\int_0^1 g(q) dq}. \quad (12)$$

However since p, h, f and g are positive, for every fixed q

$$hp + \frac{f(q)g}{p} \geq 2\sqrt{fgh}, \quad (13)$$

hence

$$c^* \geq 2 \frac{\int_0^1 \sqrt{fgh} dq}{\int_0^1 g dq}, \quad (14)$$

which proves (11).

To finish the proof of the theorem we have to show that the equality holds in (11). From the results of Aronson and Weinberger [2] it follows that $c^* \geq 2\sqrt{f'(0)}$. We separate the proof that equality holds in (11) into two cases: Case i) $c^* = 2\sqrt{f'(0)}$ and Case ii) $c^* > 2\sqrt{f'(0)}$.

Case i) $c^* = 2\sqrt{f'(0)}$. Take $g_x(u) = \alpha u^{\alpha-1}$ for $0 < \alpha < 1$. Then $g_x \in \mathcal{D}$ and it is straightforward to verify that $I(g_x) \rightarrow 2\sqrt{f'(0)}$ as $\alpha \rightarrow 0$ so that $J \geq 2\sqrt{f'(0)}$. Hence, from (11) we have

$$c^* \geq J \geq 2\sqrt{f'(0)},$$

which implies $J = c^*$ in this case. (We thank the referee for bringing this point to our attention).

Case ii) $c^* > 2\sqrt{f'(0)}$. In this case we will not only prove that equality in (11) holds but also that there exists $\hat{g} \in \mathcal{D}$ such that $c^* = I(\hat{g})$.

Let $p(q)$ be the positive solution of (7) satisfying (8). The existence of such a solution has been established in [2]. Moreover, $p(q) \sim |m|q$ near $q = 0$.

A function \hat{g} will saturate the bound (14) if and only if (13) is satisfied as an equality, i.e., if \hat{g} is such that

$$hp = \frac{f\hat{g}}{p}, \tag{15}$$

where $h = -\hat{g}'$ and $p(q)$ is the solution of (7) mentioned above. From (7) and (15) we get

$$-\frac{\hat{g}'}{\hat{g}} = \frac{c^*}{p} - \frac{p'}{p},$$

which can be integrated to yield

$$\hat{g}(q) = \frac{p(q)}{c^*} \exp \left\{ \int_q^{q_0} \frac{c^*}{p} dq \right\}. \tag{16}$$

for some fixed $0 < q_0 < 1$. To complete the argument we need to show that $\hat{g} \in \mathcal{D}$. It follows from (15) and (16) that

$$h(q) = \frac{f(q)}{c^* p(q)} \exp \left\{ \int_q^{q_0} \frac{c^*}{p} dq \right\} \geq 0$$

in $(0,1)$. Moreover, since $p(q) > 0$ in $(0,1)$ and $p \in C^1(0,1)$, we have that $\hat{g} \in C^1(0,1)$. Thus \hat{g} is a continuous, positive and decreasing function in $(0,1)$. Hence \hat{g} is bounded away from the origin. In order to show that $\int_0^1 \hat{g}(q) dq$ is finite we have to determine the behavior of \hat{g} near $q = 0$. Since $p \sim |m|q$ near 0, we have from (16) that

$$\hat{g}(q) \sim \frac{1}{\lambda} \frac{1}{q^{\lambda-1}}$$

near 0. Therefore, if $\lambda < 2$ (i.e., if $c^* > 2\sqrt{f'(0)}$), we have $\int_0^1 \hat{g}(q) dq < \infty$ and $\hat{g} \in \mathcal{D}$. \square

Remark. It follows from (16) and the behavior of p near 1 that $\hat{g}(1) = 0$.

Note. The uniqueness of the maximizer follows from the fact that we need to choose \hat{g} in such a way that (13) is satisfied as an equality. Alternatively, we can prove uniqueness directly from the concavity of $I(g)$. From the homogeneity of $I(g)$ in g

we can restrict, without loss of generality, to g 's such that $\int_0^1 g(q) dq = 1$. Assume $c^* > c_L$ and that there are two g 's in \mathcal{D} , g_1 and g_2 say, such that $I(g_1) = I(g_2) = c^*$ (and $\int g_1 = \int g_2 = 1$). Then, let $g = \alpha g_1 + (1 - \alpha)g_2$, with $0 < \alpha < 1$. Clearly, $\int g = 1$, and $h \equiv -g' = \alpha h_1 + (1 - \alpha)h_2 > 0$. Therefore,

$$I(g) = \int_0^1 \sqrt{f(\alpha g_1 + (1 - \alpha)g_2)(\alpha h_1 + (1 - \alpha)h_2)} dq .$$

However, by the Cauchy–Schwarz inequality,

$$\sqrt{(\alpha g_1 + (1 - \alpha)g_2)}\sqrt{(\alpha h_1 + (1 - \alpha)h_2)} \geq \alpha\sqrt{h_1 g_1} + (1 - \alpha)\sqrt{h_2 g_2}$$

with equality if and only if $g_1 = \beta h_1$ and $g_2 = \beta h_2$. Thus,

$$c^* \geq I(g) \geq \alpha I(g_1) + (1 - \alpha)I(g_2) = c^* ,$$

which implies $h_1/g_1 = h_2/g_2$, and therefore $g_1 \equiv g_2$.

3. Example

In this section we illustrate the result by applying it to the exactly solvable case $f(u) = u(1 - u)(1 + au)$ for which it is known that

$$c^* = \begin{cases} \sqrt{\frac{2}{a}} + \sqrt{\frac{a}{2}} & \text{if } a > 2 \\ 2\sqrt{f'(0)} = 2 & \text{if } a \leq 2 \end{cases} .$$

The idea is to combine the variational characterization contained in Theorem 1, which given a trial function yields lower bounds for c^* , with the variational characterization of Haderer and Rothe [4], which provides upper bounds on c^* .

For a given $a > 0$, let us take

$$g(q) = \frac{(1 - q)^{\lambda+1}}{q^{\lambda-1}} \quad \text{with } \lambda = 1 + \frac{2}{a} .$$

Then

$$h = -g' = \frac{(1 - q)^\lambda}{q^\lambda} (\lambda - 1)(1 + aq)$$

and [11]

$$\int_0^1 g(q) dq = \frac{\Gamma(\lambda + 2)\Gamma(2 - \lambda)}{\Gamma(4)} \quad \text{if } \lambda < 2 ,$$

i.e., if $a > 2$.

We obtain

$$\int_0^1 \sqrt{fgh} dq = \sqrt{\lambda - 1} \left(\int_0^1 \frac{(1 - q)^{\lambda+1}}{q^{\lambda-1}} dq + a \int_0^1 (1 - q)^{\lambda+1} q^{2-\lambda} dq \right) ,$$

and therefore

$$I(g) = 2\sqrt{\lambda - 1} \left(1 + a \frac{\Gamma(3 - \lambda)\Gamma(4)}{\Gamma(2 - \lambda)\Gamma(5)} \right).$$

Using the definition of λ and $\Gamma(z + 1) = z\Gamma(z)$ we obtain

$$I(g) = \sqrt{\frac{2}{a}} + \sqrt{\frac{a}{2}}$$

for $a > 2$. Therefore, from Theorem 1 we have that

$$c^* \geq \begin{cases} \sqrt{\frac{2}{a}} + \sqrt{\frac{a}{2}} & \text{if } a > 2 \\ 2\sqrt{f'(0)} = 2 & \text{if } a \leq 2 \end{cases}.$$

To prove the reversed inequality we use Haderer and Rothe's variational characterization [4]:

$$c^* \leq \sup_u \left\{ \dot{\rho}(u) + \frac{f(u)}{\rho(u)} \right\}$$

for any $\rho \in C^1(0, 1)$ such that $\rho(0) = 0$ and $\rho(u) > 0$ in $(0, 1)$.

For $a \geq 2$ choose $\rho(u) = \sqrt{a/2}u(1 - u)$. Then $\dot{\rho} + (f/\rho) = \sqrt{a/2} + \sqrt{2/a}$. For $a \leq 2$ choose $\rho = u(1 - u)$. Hence, $\dot{\rho} + (f/\rho) = 2 + (a - 2)u$ and $\sup_{0 \leq u \leq 1} (\dot{\rho} + (f/\rho)) = 2$. Therefore, from Haderer's and Rothe's bound we get

$$c^* \leq \begin{cases} \sqrt{\frac{2}{a}} + \sqrt{\frac{a}{2}} & \text{if } a > 2 \\ 2\sqrt{f'(0)} = 2 & \text{if } a \leq 2 \end{cases},$$

which combined with our lower bound gives the desired result.

4. Conclusion

We have given a variational characterization of the minimal speed for which the nonlinear diffusion equation has monotonic fronts. As Aronson and Weinberger have shown this is the asymptotic speed of propagation of a sufficiently localized positive initial condition $u(x, 0)$.

The variational principle we have derived here can also be used to study the dependence of c on the parameters of f . Monotonicity properties can be immediately derived. Derivatives of c with respect to parameters of f can be obtained using the Feynman–Hellmann formula.

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