

Absence of Diffusion in the Anderson Tight Binding Model for Large Disorder or Low Energy

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Abstract. We prove that the Green's function of the Anderson tight binding Hamiltonian decays exponentially fast at long distances on \mathbb{Z}^v , with probability 1. We must assume that either the disorder is large or the energy is sufficiently low. Our proof is based on perturbation theory about an infinite sequence of block Hamiltonians and is related to KAM methods.

1. Introduction

1.1. General Background

In this paper we analyze the Schrödinger operator (Hamiltonian)

$$H = H(v) = -\Delta + v, \quad (1.1)$$

where Δ is the finite difference Laplacian on \mathbb{Z}^v , and $v = \{v(j)\}$ is a random potential. We shall consider the case in which the $v(j)$ are independent (e.g. Gaussian) random variables with mean 0 and variance γ . The Hamiltonian (1.1) was introduced by Anderson [1] to model the dynamics of a quantum mechanical particle – the electron – moving in a random medium. The random medium may be thought of as a crystal with impurities of random strength $v(j)$. The variance γ measures the overall strength of the impurities or of the disorder. In this model the electron interacts only with the impurities. Electron-electron interactions and thermal effects are neglected.

Let $\psi_t = e^{itH}\psi_0$ be the time evolution of a wave function ψ_0 supported near the origin, e.g. $\psi_0(j) = \text{const} e^{-|j|}$. In order to describe the long time behaviour of our particle, consider the spread of ψ_t as measured by

$$r^2(t) \equiv \sum_{x \in \mathbb{Z}^v} x^2 |\psi_t(x)|^2. \quad (1.2)$$

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The quantity $r^2(t)$ represents the mean square of the distance of the particle from the origin at time t . In three or more dimensions and for small disorder γ , the particle is expected to diffuse with probability 1. In other words, except for a set of potentials v of measure 0,

$$r^2(t) \cong Dt, \quad D > 0. \quad (1.3)$$

Here D represents the diffusion constant which depends on v and ψ_0 . In solid state theory, D is proportional to the conductivity. It is easy to show that if we set $v \equiv 0$ or equal to a fixed periodic potential, $r^2(t) \approx \text{const} t^2$, hence there is infinite conductivity for such a system. (A proof of this behaviour can be based on analyzing the Heisenberg equations of motion.)

The main purpose of this paper is to prove that if the disorder γ is sufficiently large, then the diffusion constant D and the conductivity vanish with probability 1. If the disorder is small we still prove the absence of diffusion for sufficiently low energy E . This means that if the wave function ψ_0 in (1.2) belongs to the spectral subspace of states characterized by $|H - E| \leq \kappa$, with $|E|$ large enough and $\kappa \leq 1$, then $D \equiv D(E) = 0$. See (1.15)–(1.18) for the precise definition of D and of $D(E)$ used in our paper.

Anderson [1] argued that, for large disorder, the particle is trapped or localized near the origin, i.e.

$$r^2(t) \leq \text{const}, \quad \text{uniformly in } t.$$

While our results do not imply localization, we hope that some variant of the methods described in this paper will yield this result.

In one dimension, localization has been established for all strictly positive values of γ (see [2–5]). Moreover it is known that H has dense pure point spectrum with eigenfunctions which decay exponentially about some point in space. The simplest example of this phenomenon is given by the degenerate case $H = v$: Clearly $r^2(t) \leq \text{const}$, and the eigenfunctions are simply (Kronecker) delta functions.

In three dimensions and for weak disorder, the spectrum of H is expected to consist of a band of absolutely continuous spectrum filling out an interval (m, m') , outside which the spectrum is dense pure point. The numbers m, m' are called *mobility edges*. Recently, Kunz and Souillard [6] have established the existence of absolutely continuous spectrum when \mathbb{Z}^v is replaced by the Cayley tree. They have also proven results on the typical decay rate of eigenfunctions for energies near the mobility edges.

The main estimates of this paper concern the decay of the Green's function

$$G(E + i\varepsilon, v; x, y) = [H - E - i\varepsilon]^{-1}(x, y), \quad (1.4)$$

for large γ , or for $|E - 2v| \gg 1$. Roughly speaking, we establish exponential decay of G in $|x - y|$, provided v is excluded from a set of measure $\alpha|x - y|^{-p}$. The constant p can be made arbitrarily large by increasing $|E - 2v|$ or γ . The diffusion constant $D(E)$ and conductivity have expressions in terms of the Green's function, known as the Kubo formulae. The methods we use to derive estimates on G are related to those of the KAM scheme for handling small divisors. In fact, small divisors appear in the resolvent because eigenvalues of H come arbitrarily close to E . We also rely heavily on techniques we have developed to prove phase transitions in the

$1/r^2$ Ising model [7] and in the two-dimensional Coulomb gas [8]. However, no knowledge of these papers is assumed.

Our proof is based on perturbation theory about an infinite sequence of “block” Hamiltonians. The “blocks” correspond to regions of \mathbb{Z}^v on which the potential is singular, in the sense that the “block” Hamiltonian has eigenvalues close to E . In order to decouple a “block” of singular sites, Dirichlet boundary conditions are imposed, thereby decoupling H along the boundary of the region. The regions are chosen depending on the particular configuration v . We inductively use information obtained for small blocks to obtain results for larger blocks. As the block size increases, the eigenvalues, E_p , of the block Hamiltonian are permitted to get closer to E . The divergence arising from small divisors $(E - E_i)^{-1}$ will be offset by small factors arising from the Green’s function over long distances proportional to the length of the block.

Finally we remark that our analysis applies not only to the Gaussian but to a wide class of distributions. What is important for our analysis is that the density of states $\varrho(E)$ defined by (1.11) be bounded and small.

1.2. Definitions and Notation

The state space for a quantum mechanical particle moving in \mathbb{Z}^v is given by $l_2(\mathbb{Z}^v)$ – the space of square summable functions $f(j), j \in \mathbb{Z}^v$. The kinetic energy operator of such a particle is the finite difference Laplacian which is the bounded operator on $l_2(\mathbb{Z}^v)$ defined by

$$(\Delta f)(j) \equiv \sum_{i:|i-j|=1} (f(i) - f(j)). \tag{1.5}$$

Thus, the Hamiltonian $H(v)$, given by (1.1), is a self-adjoint operator defined on the domain of the potential v . Hence the Green’s function (the matrix elements of the resolvent) of H , given by (1.4), is well defined for $\varepsilon \neq 0$.

Let $\text{spec} H \equiv \sigma(H)$ denote the spectrum of H , $\sigma_{pp}(H)$ the pure point spectrum, and $\sigma_c(H)$ the continuous spectrum of H .

The random potential v belongs to the probability space

$$\Omega = \prod_{j \in \mathbb{Z}^v} (\mathbb{R}, d\lambda(v(j))), \tag{1.6}$$

where $d\lambda(v)$ is a probability measure on the reals with a bounded density with respect to Lebesgue measure. We set

$$dP(v) \equiv \prod_{j \in \mathbb{Z}^v} d\lambda(v(j)). \tag{1.7}$$

Note that the group of translations acts ergodically on Ω and leaves the spectrum of H invariant. Consequently it can be shown, see [3], that with probability one

$$\sigma(H(v)) = [0, 4v] + \text{supp } d\lambda, \tag{1.8}$$

and that $\sigma_{pp}(H(v))$ and $\sigma_c(H(v))$ are independent of v , dP -almost everywhere.

Let A be some subset of \mathbb{Z}^v . We define Δ_A to be the finite difference Laplacian with zero Dirichlet data on the boundary ∂A (\equiv the set of links connecting sites in

A to sites in the complement, $\sim A$, of A). Let

$$H_A(v) \equiv -\Delta_A + v\chi_A, \quad (1.9)$$

where χ_A is the characteristic function of A . This operator acts on the Hilbert space $l_2(A)$. If A is finite $l_2(A) \cong \mathbb{C}^{|A|}$, where $|A|$ is the number of sites in A , is finite dimensional, and the spectrum of $H_A(v)$ is discrete. This permits us to define $N_A(E, v)$ as the number of eigenvalues of $H_A(v)$ less than, or equal to E .

By a simple argument involving the ergodic theorem (see e.g. [9])

$$\lim_{A \uparrow \mathbb{Z}^v} \frac{1}{|A|} N_A(E, v) \equiv N(E) \quad (1.10)$$

exists and is independent of v , almost surely. Under our assumptions, *the density of states*

$$\varrho(E) \equiv \frac{dN(E)}{dE} \quad (1.11)$$

exists as a bounded function. In fact, Wegner [10] has shown that (see Appendix C)

$$\varrho(E) \leq \sup_v \frac{d\lambda(v)}{dv}, \quad (1.12)$$

and

$$\varrho(E) > 0, \quad \text{for } E \in \sigma(H). \quad (1.13)$$

If $\frac{d\lambda}{dv}$ has certain analyticity properties, e.g.

$$\frac{d\lambda(v)}{dv} = (2\pi\gamma)^{-1/2} \exp[-v^2/2\gamma], \quad (1.14)$$

then $\varrho(E)$ is analytic in E when either $\gamma \gg 1$, or $|E - 2v| \gg 1$ (see [9]).

Next, we turn to the formal definition of the *diffusion constant* and the *conductivity*, as given by Kubo (see e.g. [11]):

$$\varrho(E) D(E) \equiv \lim_{\varepsilon \downarrow 0} \frac{2\varepsilon^2}{\pi v} \sum_x |x|^2 \cdot \int dP(v) |G(E + i\varepsilon, v; 0, x)|^2, \quad (1.15)$$

$$D \equiv \int \varrho(E) D(E) dE, \quad (1.16)$$

and the conductivity is given by

$$\sigma(E) = e^2 \varrho(E) D(E),$$

where e is the electric charge.

To relate the time-independent definition of $D(E)$, D , in terms of the Green's function of H , to the time-dependent definitions (1.3), we choose

$$\psi_0(j) = \delta(j, 0),$$

and note that if $f(t)/t \rightarrow D$, as $t \rightarrow \infty$, for some function $f(t)$ defined on $[0, \infty)$, then

$$D = \lim_{\varepsilon \downarrow 0} \varepsilon^2 \int_0^\infty e^{-\varepsilon t} f(t) dt, \quad (1.17)$$

and formally,

$$\varepsilon^2 \int_0^\infty e^{-\varepsilon t} |(e^{-iH\varepsilon} \psi_0)(x)|^2 dt = \frac{\varepsilon^2}{2\pi} \int dE \left| G\left(E + \frac{i\varepsilon}{2}, v; 0, x\right) \right|^2 dt. \tag{1.18}$$

It is not hard to show that if D , as given by (1.15), (1.16), vanishes, an ergodic mean of $r^2(t)/t$ vanishes [i.e. existence of the limit of $r^2(t)/t$, as $t \rightarrow \infty$, is not required in the above arguments].

Before summarizing our main results we recall, for completeness, a convenient sufficient condition for the existence of point spectrum, essentially due to Ruelle [12]¹:

If $f(t)$ is a non-negative function on $[0, \infty)$ and $\varepsilon = T^{-1}$ then

$$\begin{aligned} \frac{1}{T} \int_0^T f(t) dt &\leq \varepsilon \int_0^{\varepsilon^{-1}} e^{1-\varepsilon t} f(t) dt \\ &\leq e\varepsilon \int_0^\infty e^{-\varepsilon t} f(t) dt. \end{aligned} \tag{1.19}$$

Let \mathcal{E}_c denote the projection onto the spectral subspace of H corresponding to $\sigma_c(H)$, and let $P_{|x| \geq R}$ be the projection onto wave functions which vanish in $\{x \in \mathbb{Z}^v : |x| < R\}$. Let $h(R)$ be a positive function, with $h(R) \rightarrow +\infty$, as $R \rightarrow \infty$. Then, using (1.19) and results of [12, 3],

$$\begin{aligned} \|\mathcal{E}_c \psi\|^2 &= \lim_{R \rightarrow \infty} \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \|P_{|x| \geq R} e^{iHt} \psi\|^2 dt \\ &\leq e \lim_{R \rightarrow \infty} \overline{\lim}_{\varepsilon \downarrow 0} \varepsilon \int_0^\infty e^{-\varepsilon t} \|P_{|x| \geq R} e^{-iHt} \psi\|^2 dt \\ &\leq e \overline{\lim}_{R \rightarrow \infty} (1/h(R)) \overline{\lim}_{\varepsilon \downarrow 0} \varepsilon \int_0^\infty e^{-\varepsilon t} \|\sqrt{h(|x|)} e^{-iHt} \psi\|^2 dt. \end{aligned} \tag{1.20}$$

By arguments similar to the ones used to obtain (1.18) we conclude from (1.20) that in order to show that, for dP -almost every v ,

$$\sigma(H) \cap I = \sigma_{pp}(H), \quad \text{i.e.} \quad \sigma_c(H) \cap I = \emptyset,$$

for some interval I , it suffices to show that

$$\varepsilon \sum_x h(|x|) \int_I dE \int_\Omega dP(v) |G(E + i\varepsilon, v; 0, x)|^2$$

remains bounded, as $\varepsilon \downarrow 0$. This, however, will not follow from the results proven in this paper, no matter how I is chosen.

1.3. Main Theorems

We now summarize the main results of this paper. Let δ be a measure of the disorder given by

$$\delta^{-1} \equiv \sup(d\lambda(v)/dv). \tag{1.21}$$

¹ We thank B. Simon for explaining to us Ruelle's criterion in the form stated here which is essentially in [3]

Theorem 1.1. 1) For large disorder $\delta, D(E)=0$, for all E .

2) If $d\lambda(v)$ is given by the Gaussian distribution (1.14) ($\delta = \sqrt{2\pi\gamma}$) then there is a finite constant E_γ such that $D(E)=0$, for $|E - 2v| > E_\gamma$.

Remarks. (1) Theorem 1.1 holds in arbitrary dimension v . Our estimates for E_γ approach $2v$ as $\gamma \downarrow 0$ since the density of states $\varrho(E)$ goes to 0 for $|E - 2v| > 2v$ as $\gamma \downarrow 0$.

(2) A result analogous to Theorems 1.1 and 1.2 holds for a large class of distributions $d\lambda(v)$ briefly described in Appendix C.

Theorem 1.1 is a simple corollary of the following more general result.

Theorem 1.2. There are constants $\lambda(p, m)$ and $E(\gamma, p, m)$ such that, for every $m < \infty$ and every $p < \infty$,

$$\text{Prob}\{v : |G(E + i\varepsilon, v; 0, x)| \leq e^{m(N - |x|)}, \varepsilon \neq 0\} \geq 1 - C_p/N^p, \tag{1.22}$$

uniformly in $\varepsilon \neq 0$, for some finite constant C_p , provided $\delta^{-1} \leq \lambda(p, m)$ or $|E - 2v| \geq E(\gamma, p, m)$ (when $d\lambda$ is Gaussian), respectively.

Remark. We have used (and shall use) the convention that

$$\text{Prob } V \equiv \int_V dP(v) \equiv P(V), \tag{1.23}$$

where V is some measurable subset of Ω .

Next we show that Theorem 1.2 implies Theorem 1.1. We define

$$V_N \equiv \{v : |G(E + i\varepsilon, v; 0, x)| \leq e^{m(N - |x|)}, \varepsilon \neq 0\}. \tag{1.24}$$

By Eq. (1.15)

$$\varrho(E) D(E) = \lim_{\varepsilon \downarrow 0} \frac{2\varepsilon^2}{\pi v} \sum_x |x|^2 \int |G(E + i\varepsilon, v; 0, x)|^2 dP(v).$$

We decompose \mathbb{Z}^v into annuli, $A_j, j=0, 1, 2, \dots$, where, for some constant R to be chosen later,

$$A_0 \equiv \{x : |x| < R\},$$

$$A_j \equiv \{x : R2^{j-1} \leq |x| < R2^j\}, \quad j=1, 2, 3, \dots$$

Then, using the definition of V_N , we get

$$\varrho(E) D(E) \leq \lim_{\varepsilon \downarrow 0} \frac{2}{\pi v} \sum_{j=0}^j \left\{ \varepsilon^2 e^{mN_j} P(V_{N_j}) \left[\sum_{x \in A_j} |x|^2 e^{-m|x|} \right] + c_v (R2^j)^{v+2} P(\sim V_{N_j}) \right\}, \tag{1.25}$$

where c_v is a constant (\approx volume of the unit ball in \mathbb{R}^v), and in the second term on the right side of (1.25) we have used the trivial upper bound

$$\varepsilon^2 |G(E + i\varepsilon, v; 0, x)|^2 \leq 1;$$

the sequence $(N_j)_{j=0}^\infty$ can be chosen at convenience. Clearly

$$\sum_{x \in A_j} |x|^2 e^{-m|x|} \leq \text{const } e^{-(m/4)R2^j}, \quad \text{for } j \geq 1, \tag{1.26}$$

and, by Theorem 1.2,

$$P(\sim V_{N_j}) = 1 - P(V_{N_j}) \leq C_p/N_j^p, \tag{1.27}$$

for any $p < \infty$, provided $|E|$ or δ are large enough. We now choose

$$N_j = \frac{1}{8} R 2^j.$$

It follows from (1.25)–(1.27) that

$$\begin{aligned} \varrho(E) D(E) \leq \lim_{\varepsilon \downarrow 0} \text{const} \left\{ \varepsilon^2 \left(e^{(m/8)R} R^{v+2} + \sum_{j=1}^j e^{-(m/8)R 2^j} \right) \right. \\ \left. + \sum_{j=0}^{\infty} c_v (R 2^j)^{v+2-p} C_p 8^p \right\}. \end{aligned}$$

If we choose $p > v + 2$, let ε tend to 0 and R tend to ∞ we conclude that $\varrho(E) D(E) = 0$. Since $\varrho(E) > 0$, for $E \in \text{spec} H$, see (1.13), it follows that $D(E) = 0$.

In the proof of Theorem 1.2 we shall take advantage of the fact that it is really enough to prove inequality (1.22) for the Green’s function of the Hamiltonian in a finite volume. Let A be some finite sublattice, and let $G_A(z, v; x, y)$ be the Green’s function of the Hamiltonian $H_A = -\Delta_A + v\chi_A$. In Appendix B we show that, for all $\varepsilon \neq 0$,

$$|G_A(E + i\varepsilon, v; x, y)| \leq C_\varepsilon e^{-m(\varepsilon)|x-y|}, \tag{1.28}$$

for some constant C_ε which is finite for $\varepsilon \neq 0$, and some $m(\varepsilon)$ proportional to ε , both independent of v and A . From this and the resolvent equation it follows that, for $\varepsilon \neq 0$,

$$\lim_{A \uparrow \mathbb{Z}^v} G_A(E + i\varepsilon, v; x, y) = G(E + i\varepsilon, v; x, y), \tag{1.29}$$

uniformly for x in finite subsets of \mathbb{Z}^v .

Theorem 1.2 follows from (1.29) and the following result, (see Appendix B for some details):

Theorem 1.2_A. *There are constants $\lambda(p, m)$ and $E(\gamma, p, m)$ such that, for $m > 0$ and $p < \infty$,*

$$\text{Prob}\{v : |G_A(E + i\varepsilon, v; 0, x)| \leq e^{m(N-|x|)}, \varepsilon \neq 0\} \geq 1 - K_p/N^p, \tag{1.30}$$

uniformly in $\varepsilon \neq 0$ and in A , for some finite constant K_p , provided $\delta^{-1} \leq \lambda(p, m)$, or $|E - 2v| \geq E(\gamma, p, m)$.

The proof of Theorem 1.2_A is the main goal of our paper.

The organization of our paper is as follows:

In Sect. 2 we sketch the basic strategy of the proof of Theorem 1.2_A and introduce an inductive scheme, extending over a sequence of increasing length scales, on which that proof is based. We reduce the proof of Theorem 1.2_A to some estimates on the Green’s function of $H_A = -\Delta_A + v\chi_A$, where A is an arbitrary finite sublattice having certain properties. Those estimates hold under inductively defined conditions on v . Furthermore, we state estimates on the probability that a potential v satisfies those conditions.

In Sect. 3 we prove our basic (*deterministic*) estimates on finite volume Green’s functions. Sections 2 and 3 really form the core of this paper.

In Sect. 4 we prove some purely combinatorial results concerning subsets of \mathbb{Z}^v which we shall need in our probability estimates. The results and the ideas of the proofs are closely related to results and arguments which appeared in [7, 8].

The basic probability estimates are stated and proven in Sect. 5, and in Sect. 6 the proof of Theorem 1.2_A is completed.

Some auxiliary results are collected in several appendices.

2. An Outline of the Basic Strategy

In this section we develop an inductive method which permits us to prove results concerning the decay of the Green's function in a finite volume A , for a class of "nonsingular" potentials which will turn out to have dP -measure close to 1.

We choose an energy E , to be kept fixed in the following, and define a sequence of singular sets $S_j = S_j(E; v)$ contained in \mathbb{Z}^v which depend on E and on the choice of the potential v . The sets S_j have the property that

$$S_0 \supseteq S_1 \supseteq S_2 \supseteq \dots, \tag{2.1}$$

with $\bigcap_j S_j = S_\infty = \emptyset$ with probability one.

The set of singular sites S_0 is defined by

$$S_0 = \{x \in \mathbb{Z}^v : |v(x) - E + 2v| \leq N(E, \delta)\}, \tag{2.2}$$

where $N(E, \delta) = \frac{|E|}{2}$ when $d\lambda$ is gaussian and $|E|$ is large, and $N(E, \delta) = \delta^{1/2}$ whenever the disorder δ given by (1.21) is large. The sets S_k , $k = 1, 2, \dots$, will be defined inductively.

Let $G_A(E; x, y)$ denote the Green's function of $H_A = -\Delta_A + v\chi_A$, where Δ_A is the finite difference Laplacian with 0 Dirichlet boundary conditions on ∂A . We set

$$G_A(E; x, y) = 0, \quad \text{unless } x \text{ and } y \text{ belong to } A.$$

It is easy to show that if $A \cap S_0 = \emptyset$ and $N(E, \delta) > 2v$ then

$$|G_A(E + i\varepsilon; x, y)| \leq e^{-m(E)|x-y|}, \quad \forall \varepsilon \in \mathbb{R},$$

where $m(E) \geq \ln(N/2v)$, x, y in A , (see Appendix A). As i increases, S_i consists of increasingly singular points (sites) of v , in the sense that eigenvalues of "block" Hamiltonians, associated with certain subsets of S_i , are permitted to get closer to E . In Sect. 6 we prove that the probability that S_i intersect some fixed, finite region goes to 0 rapidly with i .

The main purpose of this section is to give a precise definition of S_i , $i = 0, 1, 2, \dots$, and to state decay estimates for the Green's function $G_A(E; x, y)$, assuming that $A \cap S_k = \emptyset$, for some $k < \infty$. We then present an estimate on the probability of choosing a potential v with the property that for a given subset A and an integer k , $A \cap S_k(E; v) = \emptyset$. These results suffice to prove our main result, Theorem 1.2_A.

We now present our inductive definition of S_k , $k \geq 1$: Assume that

$$S_0 \supseteq S_1 \supseteq S_2 \supseteq \dots \supseteq S_i$$

have already been defined, with S_0 given by (2.2). We define

$$S_{i+1} \equiv S_i \sim S_i^g, \quad (\text{with } A \sim B \equiv A \setminus B), \quad (2.3)$$

where

$$S_i^g \equiv \bigcup_{\alpha} C_i^{\alpha} \subseteq S_i$$

is the subset of “gentle singular sites of strength i ” of S_i . S_i^g is a *maximal* union of components C_i^{α} satisfying the following condition.

Condition A(i).

$$(a) \quad \text{diam } C_i^{\alpha} \leq d_i \equiv \exp[K_0(5/4)^i] \equiv r_{i-1}; \quad (2.4)$$

$$(b) \quad \text{dist}(C_i^{\alpha}, S_i \sim C_i^{\alpha}) \geq 2r_i \equiv 2d_i^{5/4}; \quad (2.5)$$

$$(c) \quad \text{dist}(\text{spec } H_{\overline{C}_i^{\alpha}}, E) \geq e^{-V d_i}, \quad (2.6)$$

for all α .

Here $\text{spec } H_D$ is the set of eigenvalues of the “block Hamiltonian” $H_D = -\Delta_D + v\chi_D$, (D some finite region), and the set \overline{C}_i^{α} is defined by

$$\overline{C}_i^{\alpha} \equiv \{x : \text{dist}(x, C_i^{\alpha}) \leq 4d_i\}. \quad (2.7)$$

In Condition A(i) and in the following, $\text{diam } C$ is the diameter of $C \subset \mathbb{Z}^v$, and

$$\text{dist}(A, B) \equiv \inf|a - b|, \quad a \in A, b \in B,$$

and $|\cdot|$ is the Euclidean distance.

Definition. A set A is k -admissible iff

$$\partial A \cap \overline{C}_j^{\alpha} = \emptyset, \quad j=0, 1, \dots, k. \quad (2.8)$$

A set is admissible iff it is k -admissible, for all k .

We now describe some ideas underlying our analysis of the Green’s function, $G_A(z; x, y)$, where A is some k -admissible subset of \mathbb{Z}^v , $k=0, 1, 2, \dots$. Our estimates are based upon quantum mechanical perturbation theory, where the perturbations are chosen to be the operators which couple components of sets of singular points to their complement in A . In order to give a more precise description of our perturbation expansion we establish some more notation: Let B be some subset of A , and let $\gamma \equiv \partial B$ be the set of links (nearest neighbor pairs) connecting sites in B to sites in $\sim B$. We assume that $\gamma \subset A$. We define Γ to be the operator corresponding to γ , whose matrix elements are given by

$$\Gamma(i, j) = \begin{cases} 1, & (i, j) \in \gamma \quad (|i - j| = 1) \\ 0, & \text{otherwise.} \end{cases} \quad (2.9)$$

Note that the Hilbert space, $l_2(A)$, associated with A , is a direct sum

$$l_2(A) = l_2(B) \oplus l_2(A \sim B),$$

and, correspondingly, the Hamiltonian H_A is given by

$$H_A = H_B \oplus H_{A \sim B} - \Gamma. \quad (2.10)$$

Thus Γ is the operator coupling the system in B to the system in $A \sim B$. We define

$$G_{B/A} \equiv G_B \oplus G_{A \sim B}. \tag{2.11}$$

The first resolvent identity says that

$$G_A = G_{B/A} + G_{B/A} \Gamma G_A, \tag{2.12}$$

or

$$G_A(z; x, y) = G_{B/A}(z; x, y) + \sum_{(j, j') \in \gamma} G_{B/A}(z; x, j) G_A(z; j', y). \tag{2.13}$$

Our perturbation expansion is generated by iterating this identity, using different decompositions of A which depend on v . It will be quite important to note that if $x \in B$,

$$G_{B/A}(z; x, j) = G_B(z; x, j),$$

and $G_B(z; x, j)$ vanishes unless $j \in B$; furthermore $G_{B/A}(z; x, y) = 0$, unless x and y both belong to B or to $A \sim B$, etc.

In the next section we prove the following decay estimate on the Green's function by iterated use of the resolvent identity.

Theorem 2.1. *For $|E - 2v| + \delta$ sufficiently large and arbitrary real ε there is a constant $m = m(E, \delta) > 0$ such that if A is a k -admissible subset of the lattice \mathbb{Z}^v , and $A \cap S_{k+1} = \emptyset$ then*

$$|G_A(E + i\varepsilon; x, y)| \leq \exp[-m|x - y|],$$

provided $|x - y| \geq \frac{1}{5} r_k \equiv \frac{1}{5} d_{k+1}$. The constant $m(E, \delta)$ is independent of k and A . Moreover

$$m(E, \delta) \geq \frac{1}{2} \ln \{N(E, \delta)\}$$

for $|E - 2v| + \delta$ sufficiently large. [Recall $N(E, \delta)$ is defined after Eq. (2.2).]

Note that this result which is proven in Sect. 3 is purely *deterministic*, i.e. it holds for each potential v with the property that A is k -admissible and $S_{k+1}(E; v) \cap A = \emptyset$.

Our next task will be to estimate the dP -measure of the set of potentials, $W_{A,k}$, defined by

$$W_{A,k} \equiv \{v : A \cap S_k(E; v) = \emptyset\}. \tag{2.14}$$

In (2.14) the potential v ought to be chosen in such a way that A is k -admissible, and we are really interested in estimating the *conditional probability*

$$\text{Prob}(W_{A,k} | A \text{ is } k\text{-admissible}).$$

This conditional probability can be estimated in terms of $P(W_{A,k})$, for arbitrary, but v -independent subsets A . Moreover, bounding $P(W_{A,k})$, for v -independent subsets A , can be reduced to estimating the probability that a component of $S_j^q(E; v)$ in the sense of Condition A(j) contains a given site x of the lattice, e.g. the origin; ($j \geq k$). Thus, let

$$V_{x,j} \equiv \{v : S_j^q(E; v) \ni x\}. \tag{2.15}$$

In Sects. 4 and 5 we shall prove combinatorial and probabilistic results which yield the following theorem, proven in Sect. 6.

Theorem 2.2. *For large disorder δ , or – if $d\lambda$ is the Gaussian, given by (1.14) – for large $|E - 2v|$, there exists a positive constant $k(E, \delta)$, with $k(E, \delta) \rightarrow +\infty$, as $\delta \rightarrow \infty$ or $|E| \rightarrow \infty$, such that*

$$\begin{aligned} \text{Prob}(V_{x,j}) &\equiv \int_{V_{x,j}} dP(v) \\ &\leq \exp[-k(E, \delta) \ln d_{j-1}] \\ &= \exp[-k(E, \delta) K_0(5/4)^{j-1}]. \end{aligned} \quad (2.16)$$

Now suppose that the diameter of the region A in Theorem 2.1 is bounded by

$$\text{diam } A \leq r_k = \exp(K_0(5/4)^{k+1}). \quad (2.17)$$

Let $i \geq k$. It follows from Condition A(i), (b) that at most one component of $S_i^g(E; v)$ can intersect A . It then follows from Theorem 2.2 by straightforward arguments that

$$\begin{aligned} 1 - \text{Prob}(W_{A,k}) &\leq \sum_{j=k}^{\infty} \sum_{|x| \leq r_k} \text{Prob}(V_{x,j}) \\ &\leq r_k^v \sum_{j=k}^{\infty} \exp[-k(E, \delta) K_0(5/4)^{j-1}] \\ &\leq \exp[-\bar{k}(E, \delta) K_0(5/4)^k] \equiv d_k^{-\bar{k}(E, \delta)}, \end{aligned} \quad (2.18)$$

by (2.16) and (2.17). Here $\bar{k}(E, \delta)$ is a constant which tends to $+\infty$, as δ , or $|E|$, tend to ∞ .

If we now combine Theorem 2.1 with Theorem 2.2 and (2.18) we arrive at the following basic probability estimate for the decay of the Green's function in a finite region. See Sects. 5 and 6 for a proof.

Theorem 2.3. *Given any finite, positive number p , there exist constants λ_p, E_p such that, for*

$$\delta^{-1} = \sup_v (d\lambda(v)/dv) \leq \lambda_p, \quad \text{or} \quad |E - 2v| \geq E_p,$$

the following event, \mathcal{E}_l , holds with probability at least $1 - l^{-p}$:

\mathcal{E}_l is the set of all potentials v with the property that there exists an admissible region A such that $0 \in A$ and

$$l/2 \leq \min_{b \in \partial A} \text{dist}(0, b) \leq \max_{b \in \partial A} \text{dist}(0, b) \leq l,$$

and

$$|G_A(E + i\varepsilon, v; x, y)| \leq e^{-m|x-y|},$$

provided x and y are in A , and $|x - y| \geq l^{3/4}$. The decay rate $m = m(E, \delta)$ is as in Theorem 2.1.

Remark. The constraint $|x - y| \geq l^{3/4}$ can be replaced by $|x - y| \geq l^\sigma$, for arbitrary $\sigma > 0$.

In the proof of Theorem 2.3 (see Sects. 5 and 6) and in the considerations below we also use the following elementary lemma which is proven by using bounds on the density of states; see Appendix C (and Appendix A for additional information).

Lemma 2.4. *Let $d\lambda(v)/dv$ be bounded. Then, for an arbitrary finite subset, Λ , of \mathbb{Z}^v and an arbitrary $\kappa \leq 1$,*

$$\text{Prob}\{v : \text{dist}(\text{spec}H_A(v), E) \leq \kappa\} \leq \sqrt{2\kappa} \bar{q}(E, \delta) |\Lambda|,$$

where $\bar{q}(E, \delta) \rightarrow 0$, as $|E|$, or δ tend to $+\infty$.

The end of this section is devoted to proving that Theorem 2.3 and Lemma 2.4 imply our main result, Theorem 1.2_A (Sect. 1):

First, we assume that

$$|x| \geq N/2, \tag{2.19}$$

where N is the positive number specified in Theorem 1.2_A. Let $(A_j)_{j=0,1,2,\dots}$ be a sequence of admissible sets for which Theorem 2.3 holds, with

$$l = l_j \equiv 2|x| \cdot 4^j. \tag{2.20}$$

We set $\gamma_j \equiv \partial A_j$, and let Γ_j be the corresponding boundary operator coupling the systems in A_j and in $\sim A_j$. To simplify notation we set

$$G \equiv G_A, \quad G_j \equiv G_{A_j}.$$

[Note that, for j large enough, $A_j \supseteq A$, and all sites in $\sim A$ belong to $\sim S_0$, since $v(j) = +\infty$, for $j \in \sim A$.] By iterating the resolvent identity (2.13) we have

$$G(E + i\varepsilon, v; 0, x) = [G_0 + G_0 \Gamma_0 G_1 + G_0 \Gamma_0 G_1 \Gamma_1 G_2 + \dots](E + i\varepsilon, v; 0, x). \tag{2.21}$$

Notice that the matrix elements of G_j in the series on the right hand side of (2.21) are evaluated at sites which are separated by a distance of at least

$$-2|x|4^{j-1} + \frac{1}{2}2|x|4^j \geq \frac{1}{2}|x|4^j. \tag{2.22}$$

By Theorem 2.3 we know that if $v \in \mathcal{E}_{l_j}$ then

$$|G_j(E + i\varepsilon; z, z')| \leq \exp[-m(E, \delta)|z - z'|], \tag{2.23}$$

provided $|z - z'| \geq l_j^{3/4} = (2|x|4^j)^{3/4}$. By (2.22), this condition is satisfied in all terms on the right hand side of (2.21). Thus, it is easy to see that the series on the right hand side of (2.21) converges, for every potential $v \in \bigcap_{j=0}^{\infty} \mathcal{E}_{l_j}$, and is bounded by

$$|G(E + i\varepsilon, v; 0, x)| \leq K e^{-m(E, \delta)|x|},$$

for $|x| \geq N/2$, where K is some constant independent of E , N and of $v \in \bigcap_{j=0}^{\infty} \mathcal{E}_{l_j}$.

We now estimate $\text{Prob}\left[\bigcap_{j=0}^{\infty} \mathcal{E}_{l_j}\right]$: We choose an arbitrary, positive integer p and require that

$$\delta^{-1} \leq \lambda_p, \quad \text{or} \quad |E - 2v| \geq E_p,$$

where λ_p and E_p are the constants introduced in Theorem 2.3. Then

$$\begin{aligned} \text{Prob} \left[\bigcap_{j=0}^{\infty} \mathcal{E}_{l_j} \right] &\equiv P \left[\bigcap_{j=0}^{\infty} \mathcal{E}_{l_j} \right] \\ &= 1 - P \left[\bigcup_j (\sim \mathcal{E}_{l_j}) \right] \\ &\geq 1 - \sum_j P[\sim \mathcal{E}_{l_j}] \\ &\geq 1 - \sum_j l_j^{-p} \\ &\geq 1 - \text{const} l_0^{-p} \geq 1 - K'_p N^{-p}, \end{aligned}$$

and we have used (2.19) and (2.20) in the last inequality. This proves Theorem 1.2_A under the assumption that $|x| \geq N/2$.

If $|x| < N/2$ we set $l_j = N \cdot 4^{j+1}$, $j=0, 1, 2, \dots$. Let $\|\cdot\|$ denote the operator norm. By Lemma 2.4 we know that the inequalities

$$\begin{aligned} |G_0(E + ie, v; 0, x)| &\leq \|G_0\| \leq e^{mN/2} \\ &\leq e^{(N-|x|)}, \quad |x| \leq \frac{N}{2}, \end{aligned} \tag{2.24}$$

hold with probability

$$1 - 2e^{-mN/4} |A_0| \bar{q}(E, \delta) \geq 1 - 2(8N)^v e^{-mN/4} \bar{q}(E, \delta). \tag{2.25}$$

In (2.24) and (2.25) it is *not* required that A_0 be admissible, since in Lemma 2.4 it is not assumed that A is admissible. (In fact, admissibility of a set A is irrelevant, whenever one needs only a norm estimate on G_A .) We denote by $\mathcal{E}^{(0)}$ the event that inequality (2.24) holds. Thus by (2.25)

$$P(\mathcal{E}^{(0)}) \geq 1 - 2(8N)^v e^{-mN/4} \bar{q}(E, \delta), \tag{2.26}$$

with $m = m(E, \delta)$ as in Theorem 2.1.

It is still assumed that the regions A_j , $j=1, 2, 3, \dots$, are admissible. The resolvent expansion on the right hand side of (2.21) can now be estimated by using (2.24) and Theorem 2.3, provided the potential v belongs to the subset

$$\mathcal{E}^{(0)} \cap \left[\bigcap_{j=1}^{\infty} \mathcal{E}_{l_j} \right].$$

Furthermore

$$\begin{aligned} P \left(\mathcal{E}^{(0)} \cap \left[\bigcap_{j=1}^{\infty} \mathcal{E}_{l_j} \right] \right) &\geq 1 - P(\sim \mathcal{E}^{(0)}) - \sum_{j=1}^{\infty} P(\sim \mathcal{E}_{l_j}) \\ &\geq 1 - \frac{K''_p}{N^p}, \end{aligned} \tag{2.27}$$

if $|E - 2v|$ or δ is large.

This completes the proof of Theorem 1.2_A.

3. Decay Estimates on Green's Functions

In this section we prove inductively a family of inequalities for $G_A(E + i\varepsilon, v; x, y)$ indexed by distance scales $|x - y| \approx d_k \equiv \exp[K_0(\frac{\delta}{4})^k]$, for admissible A such that $A \cap S_k = \emptyset$. Let \mathcal{S}_k indicate the following inequality:

$$(\mathcal{S}_k) \quad |G_A(E + i\varepsilon, v; x, y)| \leq e^{-m_k|x-y|}, \tag{3.1}$$

for $|x - y| \geq \frac{1}{5}d_k \equiv \frac{1}{5}r_{k-1}$, assuming that A is $(k-1)$ -admissible, and

$$A \cap S_k = \emptyset.$$

The exponential decay rate ("mass") m_k is given by

$$m_k = \ln\left(\frac{N(E, \delta)}{2v}\right) - \sum_{i=0}^k C d_i^{-(1/4)}, \tag{3.2}$$

for some finite constant C , with $N(E, \delta)$ as in (2.2).

We note that the sequence $(m_k)_{k=0}^\infty$ is decreasing and that if we choose $d_0 = e^{K_0}$ large enough

$$m \equiv \lim_{k \rightarrow \infty} m_k \geq \frac{1}{2} \ln(N(E, \delta)) > 0, \tag{3.3}$$

provided $|E| + \sqrt{\delta}$ is chosen sufficiently large.

It is quite easy to prove \mathcal{S}_0 , see Appendix A. We shall prove \mathcal{S}_k by induction. These inequalities together with (3.3) imply Theorem 2.1.

We now assume \mathcal{S}_i , $i = 0, 1, \dots, k$ and propose to establish \mathcal{S}_{k+1} .

Since A is chosen in such a way that

$$A \cap S_{k+1} = A \cap (S_k \sim S_k^g) = \emptyset, \tag{3.4}$$

we see that \mathcal{S}_k holds for admissible subsets of

$$A \sim (S_k^g \cap A). \tag{3.5}$$

We now consider those components C_k^α of S_k^g which belong to A . By (3.4) and (3.5), these sets represent the only "new singularities" of the potential v which might obstruct the validity of \mathcal{S}_{k+1} . [If $A \cap S_k^g = \emptyset$ then $A \cap S_k = \emptyset$, by (3.4), and \mathcal{S}_{k+1} follows trivially from \mathcal{S}_k .] We shall prove \mathcal{S}_{k+1} by using \mathcal{S}_k on $G_{A \sim (S_k^g \cap A)}(E + i\varepsilon; x, y)$ and putting back the couplings to $S_k^g \cap A$ perturbatively. The convergence of our perturbative expansions can be controlled with the help of Condition A(k), Sect. 2, and our induction hypotheses.

We now start discussing the details of the proof. Let R_k^α be a k -admissible region containing a component C_k^α of $S_k^g(E; v)$ and with the property that

$$\left. \begin{aligned} \frac{1}{5}d_{k+1} \leq \text{diam } R_k^\alpha \leq \frac{3}{2}d_{k+1}, \\ \text{dist}(\partial R_k^\alpha, \overline{C_k^\alpha}) > 0, \quad \text{with } \overline{C_k^\alpha} \text{ as in (2.7)}. \end{aligned} \right\} \tag{3.7}$$

Since, by Condition A(k), (b),

$$\text{dist}(C_k^\alpha, S_k \sim C_k^\alpha) \geq 2r_k \equiv 2d_{k+1}, \tag{3.8}$$

[see (2.5)], it follows that

$$R_k^\alpha \cap (S_k \sim C_k^\alpha) = \emptyset,$$

in particular,

$$R_k^\alpha \cap S_{k+1} = \emptyset,$$

and R_k^α is an admissible set. The existence of regions R_k^α with the above properties is checked in Appendix D. Our first technical result in this section is

Lemma 3.1. *Let $R \equiv R_k^\alpha$ be as above, and choose d_0 and $|E| + \sqrt{\delta}$ sufficiently large, independently of k . If*

$$|x - y| \geq \frac{1}{5}d_{k+1},$$

then

$$|G_R(E + i\varepsilon, v; x, y)| \leq \exp[-(m_k - \mu_k)|x - y|], \tag{3.9}$$

where

$$\mu_k = 45d_k^{-1/4}m_0.$$

Proof. Let $C \equiv C_k^\alpha$, $\bar{C} \equiv \bar{C}_k^\alpha$ [see (2.7)], and let B be a $(k - 1)$ -admissible set such that $R \sim B$ is $(k - 1)$ -admissible, as well, and

$$\left. \begin{aligned} \bar{C} \supset B \supset C, \\ \text{dist}(\partial B, \{x, y\}) \geq d_k \\ \text{dist}(\sim B, C) \leq 3d_k. \end{aligned} \right\} \tag{3.10}$$

Since $|x - y| \geq \frac{1}{5}d_{k+1}$, at most one site, x or y may be contained in any one of the sets C, B, \bar{C} , while the other one is at a distance $\geq d_k$ from \bar{C} . The existence of a set B with the above properties is verified in Appendix D. We set

$$\bar{\gamma} = \partial \bar{C}, \quad \gamma = \partial B, \quad \text{and} \quad Q = R \sim B.$$

The geometrical situation envisaged here is depicted in Fig. 1.

We define $\bar{\Gamma}$ and Γ to be the operators corresponding to $\bar{\gamma}$ and γ , respectively [see (2.9)]. To prove (3.9) we now distinguish the following two cases:

- (1) $x \in Q, y \in Q$; and
- (2) $x \in B, y \in Q$,

[the case $y \in B, x \in Q$ is clearly analogous to (2)]. We first study (1): By using

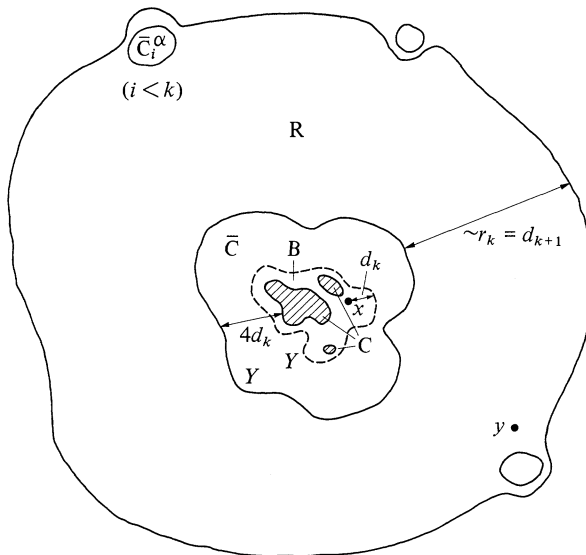


Fig. 1

the resolvent identities (2.12) we obtain

$$G_R(E + i\varepsilon; x, y) = [G_Q + G_Q \Gamma G_R \Gamma G_Q](E + i\varepsilon; x, y), \tag{3.11}$$

since, for x and y in Q , the term

$$[G_Q \Gamma G_{Q/R}](E + i\varepsilon; x, y)$$

vanishes (Γ couples Q to B !). Since Q is $(k-1)$ -admissible and $Q \cap S_k = \emptyset$, \mathcal{I}_k implies

$$|G_Q(E + i\varepsilon; u, w)| \leq \exp[-m_k |u - w|], \tag{3.12}$$

provided $|u - w| \geq \frac{1}{5}d_k$. By (3.10), inequality (3.12) can be applied to estimate G_Q in both terms on the right side of (3.11). To complete our estimate of $|[G_Q \Gamma G_R \Gamma G_Q](E + i\varepsilon; x, y)|$ we appeal to the following

Sub-Lemma 3.2. *Let u and w belong to B . Then*

$$|G_R(E + i\varepsilon; u, w)| \leq 2 \exp \sqrt{d_k}. \quad \square \tag{3.13}$$

Assuming (3.13), and using (3.12), we see that Eq. (3.11) yields the bound

$$\begin{aligned} |G_R(E + i\varepsilon; x, y)| &\leq e^{-m_k |x - y|} + \sum_{\substack{(u, u') \in \gamma \\ (w, w') \in \gamma}} e^{-m_k(|x - u| + |w - y|)} 2e^{\sqrt{d_k}} \\ &\leq e^{-m_k |x - y|} \{1 + 2|\gamma|^2 e^{7m_k d_k + \sqrt{d_k}}\}, \end{aligned} \tag{3.14}$$

and we have used the fact that

$$|u' - w'| \leq \text{diam } B \leq 7d_k,$$

by (3.10) and Condition A(k), (a). In (3.14) (u, u') and (w, w') are bonds (nearest neighbor pairs) belonging to $\gamma = \partial B$, with $u, w \in Q$, and $|\gamma|$ is the area of γ . If we set

$$\mu_k = \frac{9d_k m_0}{(1/5)d_{k+1}} = 45d_k^{-1/4} m_0,$$

we see that, for $|x - y| \geq \frac{1}{5}d_{k+1}$ and $d_0, |E| + \sqrt{\delta}$ so large that $\sqrt{d_k} \leq m_k d_k$,

$$|G_R(E + i\varepsilon; x, y)| \leq \exp[-(m_k - \mu_k)|x - y|].$$

We now come to the discussion of case (2), i.e., $x \in B, y \in Q$: We use the resolvent identity

$$G_R = G_{Q/R} + G_R \Gamma G_{Q/R},$$

which yields

$$G_R(E + i\varepsilon; x, y) = [G_R \Gamma G_Q](E + i\varepsilon; x, y),$$

because [see (2.11)]

$$G_{Q/R}(E + i\varepsilon; x, y) = 0, \quad \text{for } x \in B \subset \sim Q, \quad y \in Q,$$

and

$$G_{Q/R}(E + i\varepsilon; w, y) = G_Q(E + i\varepsilon; w, y), \quad \text{for } y \in Q.$$

Hence, using Sub-Lemma 3.2 to bound $|G_R(E + i\varepsilon; x, u)|$, x, u in B , and \mathcal{J}_k to bound $|G_Q(E + i\varepsilon; u', y)|$ (which is possible, since

$$|y - u'| \geq |y - x| - \text{diam } B \geq \frac{1}{5}d_{k+1} - 7d_k > \frac{1}{5}d_k,$$

for d_0 large enough), we obtain

$$\begin{aligned} |G_R(E + i\varepsilon; x, y)| &\leq \sum_{(u, u') \in \gamma} |G_R(E + i\varepsilon; x, u)| |G_Q(E + i\varepsilon; u', y)| \\ &\leq 2|\gamma| e^{V\sqrt{d_k}} e^{-m_k(|x-y| - 7d_k)} \\ &\leq \exp[-(m_k - \mu_k)|x - y|]. \end{aligned}$$

This completes the proof of Lemma 3.1.

We now turn to our

Proof of Sub-Lemma 3.2. By alternated application of the two resolvent identities

$$\begin{aligned} G_R &= G_{\bar{C}/R} + G_{\bar{C}/R} \bar{\Gamma} G_R \\ &= G_{B/R} + G_{B/R} \Gamma G_R \end{aligned}$$

we obtain the expansion

$$G_R(E + i\varepsilon; u, w) = [G_{\bar{C}} + G_{\bar{C}} \bar{\Gamma} G_{B/R} + G_{\bar{C}} \bar{\Gamma} G_{B/R} \Gamma G_{\bar{C}} + \dots](E + i\varepsilon; u, w). \quad (3.15)$$

Note that since Γ and u, w are all contained in \bar{C} , $G_{\bar{C}/R} = G_{\bar{C}}$, in all terms contributing to the right side of (3.15). Factors $G_{\bar{C}}$ on the right side of (3.15) are estimated by using the norm bound

$$\|G_{\bar{C}}(E + i\varepsilon; u, w)\| \leq \|G_{\bar{C}}(E + i\varepsilon; \cdot, \cdot)\| \leq e^{V\sqrt{d_k}}, \quad (3.16)$$

for u and w in \bar{C} . This bound follows from Condition A(k), (c), see (2.6), i.e.,

$$\text{dist}(\text{spec } H_{\bar{C}}, E) \geq e^{-V\sqrt{d_k}}.$$

Next, we observe that, in all non-vanishing terms on the right side of (3.15), we have

$$G_{B/R} = G_{R \sim B} \equiv G_Q,$$

since $\bar{\Gamma} \subset Q$. These factors can be estimated by using the induction hypothesis \mathcal{J}_k , since

$$\text{dist}(\bar{\gamma}, \gamma) \geq d_k > \frac{1}{5}d_k, \quad (3.17)$$

and $(R \sim B) \cap S_k = \emptyset$. [To show (3.17) we recall that, by (2.7) and (3.10),

$$\text{dist}(\bar{\gamma}, \gamma) \geq 4d_k - \text{dist}(\sim B, C) \geq d_k.]$$

Thus, applying \mathcal{J}_k to G_Q we find

$$|(\bar{\Gamma} G_Q \Gamma)(z, z')| \leq e^{-m_k|z - z'|}. \quad (3.18)$$

By (3.16) and (3.18), the series on the right side of (3.15) can be bounded by

$$e^{V\sqrt{d_k}} (1 - |\bar{\gamma}| |\gamma| e^{-m_k d_k + V\sqrt{d_k}})^{-1} \leq e^{V\sqrt{d_k}} \left(1 - e^{-\frac{m}{2} d_k + V\sqrt{d_k}}\right)^{-1} \leq 2e^{V\sqrt{d_k}},$$

if $|E| + \sqrt{\delta}$ is large enough and d_0 is chosen sufficiently large. This completes the proof of Sub-Lemma 3.2, and hence of Lemma 3.1.

We now use Lemma 3.1 to complete the *induction step*,

$$\mathcal{I}_k \text{ implies } \mathcal{I}_{k+1} \tag{3.19}$$

Thus, let A be some k -admissible region, with

$$A \cap S_{k+1} = \emptyset.$$

We may also assume that

$$\text{diam } A \geq \frac{1}{5}d_{k+1};$$

otherwise \mathcal{I}_{k+1} is empty [see (3.1)]. Furthermore we may assume that

$$A \cap S_k^g \neq \emptyset,$$

since otherwise \mathcal{I}_{k+1} follows trivially from \mathcal{I}_k , as already remarked. If

$$\text{diam } A \leq \frac{3}{2}d_{k+1},$$

\mathcal{I}_{k+1} follows directly from Lemma 3.1, by setting $R = A$ in Lemma 3.1 [see (3.7)]. Thus we may assume henceforth that

$$\text{diam } A > \frac{3}{2}d_{k+1}.$$

To establish \mathcal{I}_{k+1} , we shall combine Lemma 3.1 with a method closely related to one used in [13] to prove exponential clustering of spin-spin correlations in statistical mechanics. We need the following easy geometrical result which is proven in Appendix D.

Lemma 3.3. *If d_0 is chosen sufficiently large then, given any site $p \in A$, there is a k -admissible region $R_p \subset A$ such that*

$$\text{dist}(p, \partial R_p \sim \partial A) \geq \frac{1}{2}d_{k+1}, \tag{3.20}$$

$$\text{diam } R_p \leq \frac{3}{2}d_{k+1}, \tag{3.21}$$

and

$$\text{dist}(\{x, y\}, \partial R_p) \geq \frac{1}{5}d_{k+1}, \tag{3.22}$$

where x and y are two sites in A , with $|x - y| \geq \frac{1}{5}d_{k+1}$.

Let $\gamma_p = \partial R_p$ and let Γ_p be the corresponding operator (coupling R_p and $\sim R_p$). We shall iterate the resolvent identity

$$\begin{aligned} G_A(E + i\varepsilon; p, y) &= [G_{R_p} + G_{R_p} \Gamma_p G_A](E + i\varepsilon; p, y) \\ &= G_{R_p}(E + i\varepsilon; p, y) + \sum_{(u, u') \in \gamma_p} G_{R_p}(E + i\varepsilon; p, u) \\ &\quad \cdot G_A(E + i\varepsilon; u', y). \end{aligned} \tag{3.23}$$

(See Fig. 2 for a graphical representation of a typical term.)

In order to prove \mathcal{I}_{k+1} [see (3.1) and (3.2)], we begin by setting $p = x$. The factor $G_A(E + i\varepsilon; u', y)$ on the right side of (3.23) is expanded by iterating (3.23) with $p = u'$, etc. This process is repeated infinitely often. The resulting series expansion is

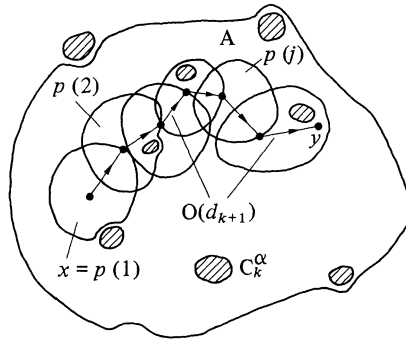


Fig. 2

rapidly convergent, and the remainder term, after n iterations, tends to 0, as $n \rightarrow \infty$, whenever $\varepsilon \neq 0$, because the matrix elements of all Greens' functions $G_{R_{p(j)}}(E + i\varepsilon; p(j), p(j + 1))$, $j = 1, 2, \dots$, appearing in the iteration of (3.23), are evaluated at sites which are separated by a distance of at least $\frac{1}{5}d_{k+1}$. If $R_{p(j)} \cap S_k = \emptyset$ then the bound

$$|G_{R_{p(j)}}(E + i\varepsilon; p(j), p(j + 1))| \leq e^{-m'_k |p(j) - p(j + 1)|}, \tag{3.24}$$

with $m'_k = m_k$, follows directly from \mathcal{S}_k . If $R_{p(j)}$ contains a component C_k^α of S_k^g then (3.24) follows from Lemma 3.1 with $R = R_{p(j)}$, and $m'_k = m_k - \mu_k$ (where μ_k is as in Lemma 3.1).

In order to estimate the sum of products of these Green's functions, let

$$m''_k \equiv m_k - \mu_k.$$

Using (3.24), we bound the n^{th} term in our expansion by

$$\sum_{\{p(2), \dots, p(n)\}} \prod_{j=1}^n e^{-m''_k |p(j) - p(j + 1)|} \tag{3.25}$$

with $p(1) = x, p(n + 1) = y$. Setting

$$l_j \equiv |p(j) - p(j + 1)|,$$

l_1, \dots, l_n are constrained by the following conditions:

$$\sum_{j=1}^n l_j \geq |x - y| \geq \frac{1}{5}d_{k+1}, \tag{3.26}$$

and

$$\frac{3}{2}d_{k+1} \geq l_j \geq \frac{1}{2}d_{k+1}, \quad \text{for } j = 1, 2, \dots, n - 1. \tag{3.27}$$

These conditions follow from (3.23), and from (3.21) and (3.22). We note that, by (3.26) and (3.27),

$$n \leq (2l/d_{k+1}) + 1 \leq 7l/d_{k+1}. \tag{3.28}$$

By summing (3.25) over all values of n and taking (3.26)–(3.28) into account, we obtain the following upper bound for the expansion resulting from our iterations of (3.23):

$$\begin{aligned}
 |G_A(E + i\varepsilon; x, y)| &\leq \sum_{l \geq |x-y|} e^{-m_k^l} (\text{const } d_{k+1})^{7vl/d_{k+1}} \\
 &\leq \exp[-m_{k+1}|x-y|],
 \end{aligned}
 \tag{3.29}$$

with

$$\begin{aligned}
 m_{k+1} &\equiv m_k'' - 7v \ln(\text{const } d_{k+1})/d_{k+1} \\
 &\geq m_k - 50m_0 d_{k+1}^{-1/4},
 \end{aligned}$$

provided d_0 is chosen sufficiently large. In (3.29) it is assumed, implicitly, that the remainder term, after n iterations of (3.23), tends to 0, as $n \rightarrow \infty$. That remainder term can be bounded by

$$\left[\sum_{\{p(2), \dots, p(n)\}} \prod_{j=1}^n e^{-m_k^j |p(j) - p(j+1)|} \right] \|G_A\| \leq (\text{const } d_{k+1} e^{-m_k^j d_{k+1}^{1/2}})^n e^{-1}$$

which tends to 0, as $n \rightarrow \infty$, for arbitrary $\varepsilon \neq 0$, provided d_0 is chosen sufficiently large.

To summarize, we have proven that for every k -admissible region A with $A \cap S_{k+1} = \emptyset$ and arbitrary sites x, y in A with the property that $|x-y| \geq \frac{1}{5}d_{k+1}$

with

$$\left. \begin{aligned}
 |G_A(E + i\varepsilon; x, y)| &\leq e^{-m_{k+1}|x-y|} \\
 m_{k+1} &\geq m_k - 50m_0 d_{k+1}^{-1/4}.
 \end{aligned} \right\}
 \tag{3.30}$$

Thus \mathcal{I}_{k+1} holds [see (3.1) and (3.2)], and the induction step is finished. The proof of Theorem 2.1 is therefore complete.

Remark. The condition that A and R_p be admissible regions is important, because ∂R_p and ∂A are required not to intersect S_k^q for Lemma 3.1 to apply.

In the remaining sections we estimate the probability of the event that

$$A \cap S_k = \emptyset,$$

where $S_k = S_k(E; v)$ depends on the potential v . This will yield Theorems 2.2 and 2.3.

4. Entropy – Combinatorial Bounds

In this section we establish some general combinatorial results which provide estimates on the number of subsets $D \subset \mathbb{Z}^v$ which contain the origin and have prescribed “volumes,” in a sense that we shall make precise. These results play an important role in the proof of our basic probability estimates stated in Theorem 2.3 (see Sect. 6). Let $\mathbb{Z}^v(n) \equiv 2^n \mathbb{Z}^v$, $n \geq 0$. With each site $k \in \mathbb{Z}^v(n-1)$ we associate a cube $c_n(k)$ with sides of length 2^n parallel to the lattice axes and centered at k . Each such cube is called an n -cube. [For $n=0$, $\mathbb{Z}^v(-1) \equiv \mathbb{Z}^v$, and a 0-cube is a site in \mathbb{Z}^v .] We note that two adjacent n cubes overlap in a rectangular region of width 2^{n-1} , for $n \geq 1$. Using this fact it is easy to show

Proposition 4.1. *Every finite subset $D \subset \mathbb{Z}^v$ is entirely contained in a single n -cube, c_n , provided that $2^n \geq 2 \text{diam} D$.*

By \mathcal{C}_n we denote an arbitrary collection of n -cubes and define $|\mathcal{C}_n|$ to be the cardinality of \mathcal{C}_n . Given a subset $D \subset \mathbb{Z}^v$ we define $\mathcal{C}_n(D)$ to be a minimal family of n -cubes [i.e. $|\mathcal{C}_n(D)|$ is minimal] which cover D . We define the “volume” of D on scale 2^n to be

$$V_n(D) \equiv |\mathcal{C}_n(D)|, \tag{4.1}$$

and the total volume

$$V(D) = \sum_{n=0}^{n_0(D)} V_n(D), \tag{4.2}$$

where $n_0(D)$ is the smallest integer such that $2^{n_0(D)} \geq 2 \text{diam} D$.

Theorem 4.2. *If $N(V)$ is the number of subsets $D \subset \mathbb{Z}^v$ such that $V(D) = V$ and $0 \in D$ then*

$$N(V) \leq 3^v 2^V (2 \cdot 3^v)^V \leq e^{K_v V}, \tag{4.3}$$

for some constant K_v .

Remark. Note $\ln N(V)$ is related to what in physics is called *entropy*.

To prove this theorem, let us first consider the family \mathcal{F} of sets D containing the origin with specified n -volumes: $V_n(D) = V_n$. Let $N(\{V_n\})$ be the cardinality of \mathcal{F} . Now consider a collection \mathcal{C}_n of n -cubes such that $|\mathcal{C}_n| = V_n$. We think of \mathcal{C}_n as specifying our set D “on scale 2^n .” Let $N(\mathcal{C}_n, V_{n-1})$ be the number of covers \mathcal{C}_{n-1}^α subordinate to \mathcal{C}_n (i.e. each cube of \mathcal{C}_{n-1}^α is covered by some cube in \mathcal{C}_n) such that $|\mathcal{C}_{n-1}^\alpha| = V_{n-1}$, $\alpha = 1, 2, \dots, N(\mathcal{C}_n, V_{n-1})$.

Proposition 4.3.

$$N(\mathcal{C}_n, V_{n-1}) \leq (2 \cdot 3^v)^{V_{n-1}}. \tag{4.4}$$

This bound clearly implies that

$$N(\{V_n\}) \leq 3^v (2 \cdot 3^v)^V, \tag{4.5}$$

where the factor 3^v arises from counting the number of covers \mathcal{C}_{n_0} , ($n_0 = n_0(D)$), consisting of a single n_0 -cube which encloses the origin. Since there are less than 2^V choices of sequences (V_n) , $n = 0, 1, 2, \dots$ for which $\sum V_n = V$, Theorem 4.1 follows.

Proof of Proposition 4.3. Given a cube $c \in \mathcal{C}_n$ and a collection \mathcal{C}_{n-1}^α as specified above, let N_c^α be the number of $(n-1)$ -cubes in \mathcal{C}_{n-1}^α contained in c . Clearly

$$\sum_{c \in \mathcal{C}_n} N_c^\alpha = V_{n-1}.$$

We now ask how many collections \mathcal{C}_{n-1}^α are there for which $N_c^\alpha = N_c$ with $\sum N_c = V_{n-1}$. Note that there are 3^v possible positions of an $(n-1)$ -cube inside some arbitrary n -cube. Thus the answer to the question posed above is that the number of collections with specified $N_c^\alpha = N_c$ is bounded by

$$\prod_{c \in \mathcal{C}_n} \binom{3^v}{N_c} \leq 3^{v V_{n-1}}.$$

Now summing over the $2^{V_{n-1}}$ choices of $\{N_c\}_{c \in \mathcal{C}_n}$ we see that (4.4) follows. \square

Remark. These bounds are very crude and can be improved with more effort. This might be important in other applications, e.g. to expansions in statistical mechanics.

In order to obtain probability estimates for the components $D \equiv C_i^\alpha$ of sets of singular sites, introduced in Sects. 2 and 3, it is convenient to define a family $\mathcal{C}'_n(D)$ of “isolated” cubes. [Cubes in $\mathcal{C}'_n(D)$ will turn out to cover subsets of D on which our eigenvalue condition (2.6) is violated.] More precisely, let $1 \leq \alpha < 2$, and $M \geq 2$ be specified constants. We define

$$\begin{aligned} \mathcal{C}'_n &\equiv \mathcal{C}'_n(\alpha, M) \\ &\equiv \{c : c \in \mathcal{C}_n, \text{dist}(c, c') \geq M2^{\alpha n}, \text{ for all } c' \in \mathcal{C}_n, c' \neq c\}, \\ V'_n &\equiv |\mathcal{C}'_n|, \quad V' \equiv \sum_{n=1}^{n_0} V'_n. \end{aligned} \tag{4.6}$$

Theorem 4.4. *For $\alpha < 2, M < \infty$ there are finite constants $K_0(\alpha, M)$ and $K'(\alpha, M)$ such that, for arbitrary collections $\mathcal{C}_n, n=0, 1, 2, \dots$, with the properties described above,*

$$V \leq K_0(\alpha, M)V_0 + K'(\alpha, M)V'.$$

Moreover, $K'(1, M) = 2$.

Remarks. For $v=2$ a variant of this result has been proven in Sect. 3 of [8]. The arguments given there combined with Proposition 4.1 prove Theorem 4.4 for $v=2$. The extension to arbitrary dimension v is immediate. (See also [7] for a discussion of this and related results when $v=1$.)

Sketch of Proof of Theorem 4.4. Let $\gamma(n)$ be defined to be the integer part of $\frac{1}{\alpha}(n-1 - \log_2[2M+3\sqrt{v}])$. The basic observation required in the proof is the following inequality: For $n > 1 + \log_2[2M+3\sqrt{v}]$,

$$V_n \leq \frac{1}{2}V_{\gamma(n)} + V'_{\gamma(n)} \quad (\text{provided } V_{\gamma(n)} \geq 2). \tag{4.7}$$

To prove (4.7) we define $\mathcal{C}''_m = \mathcal{C}_m \sim \mathcal{C}'_m$, for all m , and note that, given an arbitrary cube $c \in \mathcal{C}''_{\gamma(n)}$, one can find a cube $c' \in \mathcal{C}''_{\gamma(n)}, c' \neq c$ such that

$$\text{dist}(c, c') < M2^{\alpha\gamma(n)}.$$

By the definition of $\gamma(n)$ and Proposition 4.1, c and c' can be covered by a single cube in \mathcal{C}_n . More generally, if c_1, \dots, c_l are l cubes belonging to $\mathcal{C}''_{\gamma(n)}$ such that for each $j \in \{1, \dots, l\}$ there is a $\sigma(j) \neq j$ with the property that

$$\text{dist}(c_j, c_{\sigma(j)}) < M2^{\alpha\gamma(n)},$$

then one shows quite easily, using Proposition 4.1, that $[l/2]$ cubes in \mathcal{C}_n suffice to cover $\{c_1, \dots, c_l\}$, where $[l/2]$ = integer part of $l/2$. However, it will require $|\mathcal{C}'_{\gamma(n)}|$ different cubes in \mathcal{C}_n to cover all cubes in $\mathcal{C}'_{\gamma(n)}$. Hence

$$V_n \equiv |\mathcal{C}_n| \leq \frac{1}{2}|\mathcal{C}''_{\gamma(n)}| + |\mathcal{C}'_{\gamma(n)}| \leq \frac{1}{2}V_{\gamma(n)} + V'_{\gamma(n)},$$

which is inequality (4.7). Given n , we can iterate (4.7) $N(n)$ times, where $N(n)$ is the largest integer for which $\gamma^{N(n)}(n) \geq 0$. (Here γ^m denotes the m -fold composition of the

function γ with itself.) Using the trivial fact that $V_m \leq V_0$, for all $m > 0$, we obtain by iterating (4.7)

$$V_n \leq \left(\frac{1}{2}\right)^{N(n)} V_0 + \sum_{m=1}^{N(n)} \left(\frac{1}{2}\right)^{m-1} V_{\gamma^m(n)}. \tag{4.8}$$

By summing inequality (4.8) over all values of $n \leq n_0$ and using simple properties of γ and $N(n)$ which are discussed in detail in Sect. 3 of [8], Theorem 4.4 follows. [In particular, since $N(n) \sim \left(\frac{\log_2 n}{\log_2 \alpha}\right) + \text{const}$, $\sum_{n=1}^{\infty} \left(\frac{1}{2}\right)^{N(n)} \leq \text{const} \left(\sum_{n=1}^{\infty} \left(\frac{1}{n}\right)^{1/\log_2 \alpha}\right)$ converges, provided $\alpha < 2$.] \square

In certain applications to high- and low-temperature expansions in statistical mechanics it would be useful to have better constants in the bounds proven in Theorems 4.2 and 4.4.

5. Probability of Occurrence of Singular Sets

5.1. Statement of the Main Theorem

Let D be some finite subset of the lattice \mathbb{Z}^v . With D we associate a subset $V_{D,j} \subset \Omega$, of potentials, v , as follows:

$$V_{D,j} = \{v : D \text{ is a component of } S_j^g(E; v)\}, \tag{5.1}$$

where $S_j^g(E; v)$ is the set of ‘‘gentle singular points of strength j ,’’ as constructed in Sect. 2.

The purpose of this section is to establish an upper bound on

$$P_{D,j} \equiv \int_{V_{D,j}} dP(v) \equiv \text{Prob}(V_{D,j}). \tag{5.2}$$

Thus, let $v \in V_{D,j}$ and suppose, first, that

$$\text{diam } D \leq d_{j-1}. \tag{5.3}$$

Since D is a component of S_j^g and $S_j \subset S_{j-1}$, D clearly satisfies Condition A(j-1), (a) and (b). Since the sets, S_i^g , $i=0, 1, 2, \dots$, of singular points of strength i are always chosen to be maximal, D would be a component of S_{j-1}^g if Condition A(j-1), (c) were satisfied. Thus, since $D \notin S_{j-1}^g$, it follows that

$$\text{dist}(\text{spec } H_{\bar{D}}(v), E) < e^{-V \bar{d}_{j-1}} \tag{5.4}$$

for all $v \in V_{D,j}$, with D satisfying (5.3). Here $\bar{D} = \{x \in \mathbb{Z}^v : \text{dist}(x, D) \leq 4d_{j-1}\}$; see (2.7).

Second, if

$$\text{diam } D > d_j, \tag{5.5}$$

then $V_{D,j}$ is clearly empty, and hence

$$P_{D,j} = 0. \tag{5.6}$$

Let $\chi_{D,j}^<$ be the characteristic function of the set of potentials

$$\{v : \text{dist}(\text{spec } H_D(v), E) < e^{-\sqrt{d_j-1}}, \text{diam } D \leq d_{j-1}\},$$

and let χ_D be the characteristic function of the set, V_D , of potentials defined by

$$\{v : \exists i \text{ such that } D \text{ is a component of } S_i^g(E, v)\}. \tag{5.7}$$

Clearly $V_{D,j} \subset V_D$. Our discussion can now be summarized by the following inequality:

$$P_{D,j} \leq \begin{cases} \int \chi_{D,j}^<(v) \chi_D(v) dP(v), & \text{if } \text{diam } D \leq d_{j-1} \\ P_D \equiv \int \chi_D(v) dP(v), & \text{if } \text{diam } D > d_{j-1} \\ 0, & \text{if } \text{diam } D > d_j. \end{cases} \tag{5.8}$$

If $\text{diam } D \leq d_{j-1}$, $P_{D,j}$ can be further estimated by

$$P_{D,j} \leq (P_{D,j}^<)^{1/2} P_D^{1/2}, \tag{5.9}$$

where

$$P_{D,j}^< \equiv \int \chi_{D,j}^<(v) dP(v). \tag{5.10}$$

Inequality (5.9) follows from the first inequality in (5.8), by the Schwarz inequality. We have thus reduced the problem of estimating $P_{D,j}$ to estimating P_D and $P_{D,j}^<$. By Lemma 2.4, if $|E - 2v|$ or δ is sufficiently large then

$$P_{D,j}^< \leq 2\bar{q}(E, \delta)(9d_{j-1})^y \exp(-\sqrt{d_{j-1}}/2), \tag{5.11}$$

where $\bar{q}(E, \delta) \rightarrow 0$ as $|E| \rightarrow \infty$ or $\delta \rightarrow \infty$. Thus it remains to estimate P_D .

Theorem 5.1. *If $|E - 2v|$ or δ is sufficiently large*

$$P_D \leq \exp[-k_0(E, \delta)V_0(D) - k'(E, \delta)V'(D)], \tag{5.12}$$

where $V_0(D)$ and $V'(D)$ are as defined in Sect. 4, and k_0, k' are constants which tend to $+\infty$, as $|E| \rightarrow \infty$ or $\delta \rightarrow \infty$.

In Sect. 6 we shall combine the inequality in Theorem 5.1 with the entropy estimates of Sect. 4 to complete our probability estimates for the decay of the Green's function, $G(E + i\epsilon, v; x, y)$, announced in Theorems 1.2 and 1.2_A of the introduction.

5.2. Proof of Theorem 5.1

Let us first recall that $\mathcal{C}'_n(D) \subset \mathcal{C}_n(D)$ is the set of isolated n -cubes c with the property that $c \in \mathcal{C}'_n(D)$ iff

$$\text{dist}(c, \tilde{c}) \geq 2 \cdot 2^{(5/3)^n},$$

for all $\tilde{c} \in \mathcal{C}_n(D)$ with $\tilde{c} \neq c$. (We have set $M = 2, \alpha = 5/3$; see Sect. 4 for definitions.) The aim in Theorem 5.1 is to associate a small factor e^{-k} with each isolated cube c in $\mathcal{C}'_n, n \geq 1$. Roughly speaking, isolated subsets of D are extremely unlikely, because on all such subsets the eigenvalue condition (2.6) must be violated. Lemma 2.4 can be used to show that this occurs with low probability, since $\kappa \simeq e^{-\sqrt{d_n}}$.

(Note that if the eigenvalue condition were satisfied then $c \cap D$ would itself be a component belonging to S_m^g , $m < j$, because of our maximality condition.)

In order to make the above ideas more precise we shall first bound the characteristic function of V_D by a product of characteristic functions. To define these characteristic functions let $j(n)$ be the smallest integer such that

$$d_{j(n)} \equiv \exp[K_0(\frac{\delta}{4})^{j(n)}] \geq 2^n. \tag{5.13}$$

If $c \in \mathcal{C}'_n(D)$ we set

$$\overline{c \cap D} = \{x : \text{dist}(x, c \cap D) \leq 4d_{j(n)}\}, \tag{5.14}$$

and define $\chi_{n,c}(v)$ to be the characteristic function of

$$\{v | \text{dist}(\text{spec } H_{\overline{c \cap D}}(v), E) \leq \exp - \sqrt{d_{j(n)}}\}. \tag{5.15}$$

When $n=0$ we set $\mathcal{C}'_0(D) = D$ and identify $c \in \mathcal{C}'$ with a site in D . We define $\chi_{0,c}$ to be the characteristic function of the set

$$\{v : |v(c) - (E - 2v)| \leq N(E, \delta)\}, \tag{5.16}$$

where $N(E, \delta)$ is defined after (2.2). Let

$$I(K_0) \equiv \{0\} \cup \{n \in \mathbb{Z} : n \geq K_0 / \ln 2\}, \text{ with } K_0 \text{ as in (2.4)}. \tag{5.17}$$

Note that $j(n) \geq 1$ for all $0 \neq n \in I(K_0)$. We adopt the convention that an empty product is defined to be 1 and that $\chi_{n,c}(v) = 1$ for $n \notin I(K_0)$.

In Sect. 5.3 we shall prove the following lemma.

Lemma 5.2. *For $\chi_{n,c}$ defined as above*

$$\chi_{V_D}(v) \leq \prod_{n=0}^{\infty} \prod_{c \in \mathcal{C}'_n(D)} \chi_{n,c}(v). \quad \square \tag{5.18}$$

Assuming this lemma we shall prove Theorem 5.1 by using Hölder's inequality together with the statistical independence of $\chi_{n,c}$, for disjoint c 's. In fact, if c_1 and $c_2 \in \mathcal{C}'_n(D)$ then

$$\overline{c_1 \cap D} \cap \overline{c_2 \cap D} = \emptyset, \text{ if } c_1 \neq c_2,$$

because

$$\begin{aligned} \text{dist}(\overline{c_1 \cap D}, \overline{c_2 \cap D}) &\geq 2 \cdot 2^{(5/3)n} - 9d_{j(n)} \\ &\geq 2 \cdot 2^{(5/3)n} - 9 \cdot 2^{(5/4)n} \\ &\geq 0, \text{ if } K_0 \text{ is large enough.} \end{aligned}$$

This implies the independence of $\chi_{n,c}$, $c \in \mathcal{C}'_n(D)$. Hence

$$\int \prod_{c \in \mathcal{C}'_n(D)} \chi_{n,c}(v) dP(v) = \prod_{c \in \mathcal{C}'_n(D)} \int \chi_{n,c}(v) dP(v). \tag{5.19}$$

Next we combine Lemma 5.2 and (5.19) with the Hölder inequality

$$|\int f g dP| \leq (\int |f|^{1/r} dP)^r (\int |g|^{1/(1-r)} dP)^{1-r}, \tag{5.20}$$

$0 < r < 1$, and use the obvious fact that $\chi^q = \chi$, $q > 0$. This yields

$$\begin{aligned} \int \chi_{V_D}(v) dP(v) &\leq \int \left\{ \prod_{n=0}^{\infty} \prod_{c \in \mathcal{C}'_n(D)} \chi_{n,c}(v) \right\} dP(v) \\ &\leq \left[\prod_{c \in \mathcal{C}'_0(D)} \int \chi_{0,c} dP \right]^{1-r} \cdot \left[\int \left\{ \prod_{n=1}^{\infty} \prod_{c \in \mathcal{C}'_n(D)} \chi_{n,c} \right\} dP \right]^r \\ &\leq \dots \\ &\leq \prod_{n=0}^{\infty} \prod_{c \in \mathcal{C}'_n(D)} [\int \chi_{n,c} dP]^{r^n(1-r)} \end{aligned} \tag{5.21}$$

by successive applications of (5.19) and (5.20).

Proof of Theorem 5.1. By the definition of $\chi_{n,c}$, (5.15), and the fact that $d_{j(n)} \geq 2^n$, we see that Lemma 2.4, with $\kappa = \exp(-\sqrt{2^n})$, implies

$$[\int \chi_{n,c} dP]^{r^n(1-r)} \leq \bar{q}(E, \delta)^{r^n(1-r)} |\overline{c \cap D}|^{r^n(1-r)} 2 \exp[(r-1)(r\sqrt{2})^n/2]. \tag{5.22}$$

We now choose $r = 0.8 > 1/\sqrt{2}$. Note that the right side of (5.22) is small when n is large. When n is small we use the fact that $\bar{q}(E, \delta)$ is small for $|E|$ or δ large. Therefore each factor in the product on the right side of (5.21) yields a small factor, for each cube in \mathcal{C}'_n , and Theorem 5.1 follows.

5.3. Proof of Lemma 5.2

Let $j(n)$ be given by (5.13), and $\overline{c \cap D}$ by (5.14). For each $c \in \mathcal{C}'_n(D)$ we propose to show that

$$\text{dist}(\text{spec } H_{\overline{c \cap D}}(v), E) \leq \exp(-\sqrt{d_{j(n)}}), \tag{5.23}$$

for all $v \in V_D$, with V_D as in (5.7). From this Lemma 5.2 clearly follows. Since $c \in \mathcal{C}'_n(D)$, we have

$$\begin{aligned} \text{diam}(D) &\geq \text{dist}(c, \mathcal{C}'_n(D) \sim c) \geq 2 \cdot 2^{(5/3)n} \\ &\geq 2 \cdot 2^{(5/4)2n} \geq 2 \cdot d_{j(n)}^{5/4} \end{aligned} \tag{5.24}$$

and we have used (5.13).

Let $i(n)$ be the largest integer such that

$$d_{i(n)} \equiv \exp[K_0(5/4)^{i(n)}] \leq 2^{(5/3)n}. \tag{5.25}$$

Then, by (5.24) and (2.4), we see that D can be a component of $S_k^q(E; v)$ only for some value of k satisfying

$$k > i(n).$$

If $j \leq k$, and C_j^α is a component of $S_j^q(E; v)$ different from D then, by (2.5),

$$\text{dist}(c \cap D, C_j^\alpha) \geq \text{dist}(D, C_j^\alpha) \geq 2d_j^{(5/4)}, \tag{5.26}$$

and if $j > k$

$$\begin{aligned} \text{dist}(c \cap D, C_j^\alpha) &\geq \text{dist}(D, S_k \sim D) \\ &\geq 2d_{k+1} > 2d_{i(n)+1} > 2 \cdot 2^{(5/3)n} \geq 2d_{j(n)}^{5/4}. \end{aligned} \tag{5.27}$$

Inequalities (5.24), (5.26), and (5.27) imply that $c \cap D$ satisfies Condition A(j(n)), (a) and (b). Now recall that the sets $S_j^g(E; v)$ are required to be maximal. Thus if $c \cap D$ satisfied Condition A(j(n)), (c) it would be a component of $S_{j(n)}^g(E; v)$. Since $j(n) < i(n) < k$, (5.23) now follows.

It should be pointed out that, for each $c \in \mathcal{C}'_n(D)$, the integer $j(n)$ and the validity of (5.23) do not depend on the potential v , as long as $v \in V_D$.

6. Proofs of Theorems 2.2 and 2.3, Conclusions

In this section we use the results of Sects. 4 and 5 and Theorem 2.1, proven in Sect. 3, to complete the proofs of Theorems 2.2 and 2.3. We recall that it has been shown in Sect. 2 that our main technical result, Theorem 1.2_A of Sect. 1, follows from Theorem 2.3 and Lemma 2.4. A proof of Lemma 2.4 and related matters are contained in Appendices A and C.

The main points which remain to be discussed are the following:

(1) Let $\sigma \in (0, 1)$ and $l \gg 1$ be given. Let k be the largest integer for which

$$\frac{1}{5}d_k \leq l^\sigma. \tag{6.1}$$

We must estimate the probability, $P(l, \sigma)$, that there exists a $(k - 1)$ -admissible region A containing 0 such that

$$l/2 \leq \min_{b \in \partial A} \text{dist}(0, b) \leq \max_{b \in \partial A} \text{dist}(0, b) \leq l, \tag{6.2}$$

and with the property that

$$A \cap S_k(E; v) = \emptyset. \tag{6.3}$$

For such regions it follows from Theorem 2.1 that

$$|G_A(E + i\varepsilon, v; x, y)| \leq e^{-m|x-y|}, \tag{6.4}$$

with

$$m = m(E, \delta) \geq \frac{1}{2} \ln(N(E, \delta)),$$

provided

$$|x - y| \geq l^\sigma \geq \frac{1}{5}d_k. \tag{6.5}$$

Theorem 2.3 follows if we can show that given any $p < \infty$, there is a finite constant $c_{p, \sigma}$ such that, for $|E| + \sqrt{\delta} \geq c_{p, \sigma}$,

$$P(l, \sigma) \geq 1 - l^p. \tag{6.6}$$

Our proof of this estimate proceeds as follows:

First we notice that, with probability 1, there exists a $(k - 1)$ -admissible region A satisfying (6.2). This is shown in Appendix D.

Thus it is enough to estimate the conditional probability, P_k , that $A \cap S_k(E; v) = \emptyset$, given that A is a $(k - 1)$ -admissible set satisfying (6.2). But P_k can be bounded below by the probability $P_k(l)$ that any set A which satisfies (6.2) does not intersect $S_k(E; v)$. Clearly

$$P_k(l) \geq 1 - \text{const } l^\nu \max_{|x| \leq l} (1 - P_{\lambda, k}) \tag{6.7}$$

where $P_{x,k}$ is the probability that $x \notin S_k(E; v)$. Since the distribution $dP(v)$ of potentials v is translation-invariant,

$$\min_{|\Lambda| \leq l} P_{\Lambda,k} = P_{0,k}. \tag{6.8}$$

Next, we bound $P_{0,k}$ from below by

$$P_{0,k} \geq 1 - \sum_{j=k}^{\infty} p_j - p_{\infty}, \tag{6.9}$$

where $p_j \equiv \text{Prob}(V_{0,j})$ is the probability that $0 \in S_j^q(E; v)$. Here

$$V_{x,j} \equiv \{v : S_j^q(E; v) \ni x\},$$

see (2.15), and p_{∞} is the probability that

$$S_{\infty} \equiv \bigcap_{k=0}^{\infty} S_k = S_0 \sim \left(\bigcup_{k=0}^{\infty} S_k^q \right)$$

is non-empty.

Thus, we must show

(2) for arbitrary finite regions A

$$\text{Prob}\{v : S_{\infty} \equiv S_{\infty}(E; v|_A) = \emptyset\} = 1,$$

i.e. $p_{\infty} = 0$; and

(3) given an arbitrary $q < \infty$, there is a finite constant c_q such that

$$p_j \leq d_j^{-q} \tag{6.10}$$

if $|E| + \sqrt{\bar{\delta}} \geq c_q$.

Remarks. 1) We note that (3) clearly yields Theorem 2.2.

2) From (6.9), (2) and (3) it follows that

$$P_{0,k} \geq 1 - \text{const} d_k^{-q}, \quad q > 0,$$

(because $p_j \leq d_j^{-q} = \exp[-qK_0(5/4)^j]$). Now, recall that k is the largest integer with the property that

$$\frac{1}{5}d_k \leq l^{\sigma}, \quad \text{i.e.,} \quad d_{k+1} = d_k^{5/4} \geq 5l^{\sigma}.$$

Thus

$$P_{0,k} \geq 1 - \text{const} l^{-\frac{4}{5}\sigma q}. \tag{6.11}$$

By the definition of $P(l, \sigma)$, (6.7), (6.8), and (6.11), it now follows that

$$\begin{aligned} P(l, \sigma) &\geq 1 - \text{const} l^{v - \frac{4}{5}\sigma q} \\ &\geq 1 - l^{-p}, \end{aligned}$$

for any finite p , if $|E| + \sqrt{\bar{\delta}} \geq c_q$, with $q > \frac{5}{4\sigma}(p + v)$. This proves our basic lower bound on $P(l, \sigma)$, see (6.6). Thus the proof of Theorem 2.3 is complete.

We now turn to our proofs of (2) and (3); see (6.10). Point (2) is settled easily: Since A is assumed to be a finite region, Conditions A(i), (a) and (b) of Sect. 2 will

be satisfied for i large enough. Thus for S_∞ to be non-empty it is necessary that E is an eigenvalue of $H_A = -\Delta_A + v\chi_A$. Since $\text{spec} H_A$ is discrete, this event has probability 0.

To settle (3) we use the results in Sects. 4 and 5: To begin, note that

$$p_j \equiv \text{Prob}(V_{0,j}) \leq \sum_{D \geq 0} P_{D,j}, \tag{6.12}$$

where $P_{D,j}$ is the probability that the set $D \subset \mathbb{Z}^v$ is actually a component of $S_j^q(E; v)$. Thus we need an estimate on $P_{D,j}$. Such an estimate has been established in Sect. 5. Let $V(D)$ and $V'(D)$ be defined as in Sect. 4; see (4.2) and (4.6). Then, by (6.12) and inequality (5.8),

$$\begin{aligned}
 p_j &\leq \sum_{V=1}^{\infty} \sum_{\substack{D \geq 0 \\ V(D)=V}} P_{D,j} \\
 &\leq \sum_{V=1}^{\infty} \left\{ \sum_{\substack{D \geq 0 \\ V(D)=V \\ \text{diam} D \leq d_{j-1}}} (P_{D,j}^<)^{1/2} P_D^{1/2} + \sum_{\substack{D \geq 0 \\ V(D)=V \\ \text{diam} D \geq d_{j-1}}} P_D \right\}, \tag{6.13}
 \end{aligned}$$

where $P_{D,j}^<$ has been defined in (5.10) and P_D in (5.8). The second inequality in (6.13) follows from (5.8). By (5.11),

$$P_{D,j}^< \leq 2\bar{q}(E, \delta)(9d_{j-1})^v e^{-V\bar{d}_{j-1}/2} \equiv \Phi_j(E, \delta),$$

and by Theorem 5.1,

$$P_D \leq \exp[-k_0(E, \delta)V_0(D) - k'(E, \delta)V'(D)],$$

$$\tag{6.14}$$

if $|E| + \sqrt{\delta}$ is large enough. [Here $V_0(D) \equiv |D|$ is the number of sites in D .]

Next, we use Theorem 4.4, with $\alpha = \frac{5}{3}$ and $M = 2$, to conclude that

$$V(D) \leq K_0 V_0(D) + K' V'(D),$$

so that, by (6.14),

$$P_D \leq \exp[-K(E, \delta)V(D)], \tag{6.15}$$

for some constant $K(E, \delta)$ which tends to $+\infty$, as $(|E| + \delta) \rightarrow \infty$. It is an immediate consequence of the definition of V (see Sect. 4) that

$$V(D) \geq \text{const} \ln \text{diam} D. \tag{6.16}$$

Thus, by (6.13)–(6.16),

$$p_j \leq \sum_V \Phi_j(E, \delta)^{1/2} e^{-K(E, \delta)V/2} N(V) + \sum_{V \geq \text{const} \ln d_j} e^{-K(E, \delta)V} N(V),$$

where

$$N(V) \equiv \# \{D : D \geq 0, V(D) = V\}.$$

By Theorem 4.2,

$$N(V) \leq e^{K_v V}.$$

Thus

$$p_j \leq \Phi_j(E, \delta)^{1/2} \sum_V e^{\left(K_V - \frac{1}{2}K(E, \delta)\right)V} + \sum_{V \geq \text{const} \ln d_j} e^{(K_V - K(E, \delta))V}. \tag{6.17}$$

By (6.14) and (6.17),

$$p_j \leq (1/d_j)^{q(E, \delta)},$$

with $q(E, \delta) \rightarrow \infty$, as $|E| + \delta \rightarrow \infty$. This proves (3).

The proofs of the results announced in Sects. 1 (Theorems 1.1 and 1.2) and 2 (Theorem 2.3) are now complete.

To conclude, we wish to mention some interesting open problems.

1) It appears to be rather straightforward to extend our analysis to systems with off-diagonal disorder, or to the problem of wave propagation in a random environment on \mathbb{Z}^v , provided the disorder is large (or the energy is constrained to a suitable region).

2) It would be interesting to improve our methods in such a way that not only absence of diffusion (for large disorder or at low energies), but localization, in the strong sense of the word, is obtained, i.e.

$$\text{spec} H \cap I = \sigma_{pp}(H) \cap I \neq \emptyset$$

if I is chosen to be a union of intervals located at sufficiently large values of $|E|$ (and if the disorder is large, I is the whole real line). A proof of this requires mastering more serious small divisor problems. While we have envisaged a tentative strategy to get at such results we do not know of any complete analysis.

3) As the reader has noticed, estimates on the density of states are an important ingredient of our analysis. It would be interesting to improve the present estimates on the density of states (see [9, 10]) and, for example, establish bounds on and smoothness properties of $\varrho(E)$, for E in the region where extended states are expected.

4) It would be interesting to prove general upper bounds on the diffusion constant $D(E)$, for arbitrary E . This will probably require new ideas.

5) It may be possible to apply our methods to almost periodic Schrödinger operators on \mathbb{Z}^v . In particular one might consider a limit periodic potential with random coefficients. A primary problem in an analysis of almost periodic operators will be control over the density of states.

Appendix A: The Path Expansion for G

In this appendix we expand the Green's function in the off diagonal matrix elements of $-\Delta$ to obtain a formal series for $G_A(E)$. The terms of the series are indexed in a natural way by paths ω which represent a sequence of nearest neighbor pairs of lattice sites in \mathbb{Z}^v . If D is the diagonal operator

$$D_{ij} = (2v + v_j - E - i\varepsilon)^{-1} \delta_{ij},$$

and J denotes the off diagonal elements of Δ_A , i.e.

$$J_{ij} = 1 \quad \text{if} \quad |i - j| = 1, \quad i, j \in A,$$

and 0 otherwise. Then formally, see [1, 14]

$$\begin{aligned} G_A(E + i\varepsilon, v; x, y) &= [D^{-1} - J]^{-1}(x, y) \\ &= \sum_{n=0}^{\infty} [D(JD)^n](x, y) \\ &= \sum_{\omega: x \rightarrow y} \prod_{j \in A} D_{jj}^{n_j(\omega)}, \end{aligned} \quad (\text{A.1})$$

where ω ranges over nearest neighbor paths starting at x and ending at y which are constrained to lie in A . Also, $n_j(\omega)$ denotes the number of times ω visits the lattice site $j \in A$. Note that if $S_0 \cap A = \emptyset$ then $|D_j| \leq \frac{1}{N}$, $N \equiv N(E, \delta)$ [see (2.2) for the definition of $N(E, \delta)$] and

$$\begin{aligned} |G_A(E; x, y)| &\leq \sum_{\omega: x \rightarrow y} \prod_{j \in A} \left(\frac{1}{2v + \varepsilon} \right)^{n_j(\omega)} \left(\frac{2v + \varepsilon}{N} \right)^{n_j(\omega)} \\ &\leq \left| \frac{2v + \varepsilon}{N} \right|^{|x-y|} \sum_{\omega: x \rightarrow y} \prod_{j \in A} \left(\frac{1}{2v + \varepsilon} \right)^{n_j(\omega)} \\ &= \left| \frac{2v + \varepsilon}{N} \right|^{|x-y|} (-\Delta + \varepsilon)^{-1}(x, y) \end{aligned} \quad (\text{A.2})$$

with ε chosen such that $\left| \frac{2v + \varepsilon}{N} \right| \leq 1$. [This is possible if an appropriate definition of $N(E, \delta)$ is made.]

In general, it is clear that the factors D_{jj} may become arbitrarily large, and (A.1) will not be absolutely convergent. This is the problem of small divisors. Our strategy is to selectively expand the off diagonal matrix elements of Δ , depending upon v . In this way we obtain a convergent series.

Representation (A.1) is also useful for estimates on the density of states, see (1.11),

$$\begin{aligned} \varrho(E) &= \lim_{\varepsilon \rightarrow 0} \mathcal{I} m \int dP(v) G(E + i\varepsilon, v; 0, 0) \\ &= \lim_{\varepsilon \rightarrow 0} \mathcal{I} m \sum_{\omega: 0 \rightarrow 0} \prod_j \int D_{jj}^{n_j(\omega)} d\lambda(v_j). \end{aligned}$$

If $d\lambda(v_j)$ is, for example, the Gaussian given by (1.14), then by deforming the contour of integration so as to avoid the zero of D_j^{-1} , it is easy to show that if $|E - 2v|$ is large or γ is large then

$$|\int D_{jj}^{n_j(\omega)} d\lambda(v_j)| \leq a^{n_j(\omega)},$$

where $a = a(E, \gamma) \rightarrow 0$ as either $|E - 2v|$ or γ become large. Thus if $|a| < (2v)^{-1}$ the integrated series converges and $\varrho(E)$ is bounded and analytic in E . Similarly one can prove that the *average* Green's function decays in $|x - y|$ for large $|E|$. In order to make this argument rigorous it is important to first take ε large so that the series (A.1) is manifestly convergent. We then let $\varepsilon \rightarrow 0$ after integration. Moreover one can obtain with some extra effort good bounds on $\varrho(E)$ as $|E - 2v| \rightarrow \infty$ (see [9, 14] for details).

We remark that the above analysis breaks down if one tries to estimate the average of the *absolute value* of the Green's function because deformation of our contour is no longer possible.

Appendix B: Proof of (1.28)

Let $U(a)$ be the multiplication operator on $l_2(\mathbb{Z}^v)$ given by $e^{ia \cdot j}$. By direct calculation we see that

$$U(-a)H_A(v)U(a) = H_A(v) + Q_A(a),$$

where $Q_A(a)$ is a bounded operator independent of v which has an analytic continuation in a satisfying

$$\|Q_A(a)\| \leq C|a| \tag{B.1}$$

for $a \in \mathbb{C}^v$ and $|a| \leq 1$. If we choose $|a| \leq \frac{\varepsilon}{2C}$ then

$$U(-a)G_A(E + i\varepsilon)U(a) = [H_A - E + i\varepsilon Q_A(a)]^{-1} \tag{B.2}$$

is a bounded operator on $l_2(\mathbb{Z}^v)$ with norm less than $2\varepsilon^{-1}$. Now observe that the x, y matrix element of (B.2) which equals

$$e^{ia \cdot (x-y)} G_A(E + i\varepsilon; x, y)$$

must be bounded in absolute value by $2\varepsilon^{-1}$. This completes our proof of (1.28).

Remarks. A similar result holds if there is a gap in the spectrum of H , with ε playing the role of the width of the gap.

In the above we have relied on the so-called Combes-Thomas argument (see [15] and references therein).

Appendix C: Bounds on Density of States

Proof of Lemma 2.4. Let $\chi_A(\kappa, E)$ be the characteristic function of the set

$$\{v | \text{dist}(\text{spec } H_A(v), E) \leq \kappa\}, \tag{C.1}$$

and let $N_A(E, v)$ be as in (1.10). Following Wegner [10],

$$\begin{aligned} \int \chi_A(\kappa, E) dP &\leq \int [N_A(E + \kappa, v) - N_A(E - \kappa, v)] dP \\ &= \int_{|E' - E| \leq \kappa} \frac{d}{dE'} N_A(E', v) dE' dP \\ &= \sum_j - \int_{|E' - E| \leq \kappa} \frac{d}{dv_j} N_A(E', v) dE' dP \\ &\leq \sum_j \sup_{v_j} \left(\frac{d\lambda(v_j)}{dv_j} \right) \int \delta_j N_A(E', v) dE' \prod_{k \neq j} d\lambda(v_k), \end{aligned} \tag{C.2}$$

where

$$\delta_j N_A(E', v) \equiv N_A(E', v^-) - N_A(E', v^+) \tag{C.3}$$

and $v_k^\pm = v_k$ for $k \neq j$, and $v_j^+ > v_j^-$ denote the endpoints of the support of $d\lambda$ (possibly $\pm \infty$). Note that the integrands above are all nonnegative and that we have used the relation $N(E, v) = N(0, v - E)$. Since v_j is a rank one perturbation of $H_A(v)$ it follows from the minimax principle that the eigenvalues of $H_A(v)$ are interlaced $\lambda_k^- \leq \lambda_k^+ \leq \lambda_{k+1}^-$, hence $0 \leq \delta_j N \leq 1$, a.e. in v (since the λ_k are distinct with probability 1). Thus

$$\int \chi_A(\kappa, E) \leq 2\kappa \delta^{-1} |A|. \tag{C.4}$$

This establishes Lemma 2.4 if δ is large.

If $|E|$ is large, $N_A(E, v) \geq 1$ only if for at least one j , $|v_j - E + 2v| \leq 2v$ which occurs with probability less than

$$|A| \int_{|v - E + 2v| \leq 2v} d\lambda(v). \tag{C.5}$$

When $d\lambda(v)$ is Gaussian, then the integral is small for large E . Lemma 2.4 follows by taking the geometric mean of (C.4) and (C.5).

Remark. We have used only the fact that $\frac{d\lambda(v)}{dv}$ is bounded and that for large $|E|$ (C.5) is small. If we assume that $\lambda(v_j)$ is Gaussian then the factor $\sqrt{2\kappa}$ which appears in Lemma 2.4 may be replaced by 2κ . This follows from the method of Appendix A (see also [9]).

Appendix D: Admissible Sets

Lemma. *Let $D_2 \supset D_1$ be rectangular regions such that*

$$\text{dist}(D_1, \sim D_2) \geq 10d_k. \tag{D.1}$$

If d_0 is sufficiently large, then there is a k -admissible set A such that $D_2 \supset A \supset D_1$.

Remarks. The admissible sets needed in Sects. 3 and 6 are given by the above lemma after a suitable choice of D_1 and D_2 .

The proof of this lemma is nearly the same as the argument which appears in Appendix E of [8].

Proof. Let

$$\tilde{C}_j^\alpha = \{x \in \mathbb{Z}^v : \text{dist}(x, \overline{C}_j^\alpha) \leq 2\}, \tag{D.2}$$

and let

$$U = \bigcup_{\alpha, j \leq k} \tilde{C}_j^\alpha.$$

Now consider the components of U (defined in the usual nearest neighbor sense) which meet ∂D_1 – call them U^β . We claim that

$$\text{diam}(U^\beta) \leq 10d_r, \tag{D.3}$$

where r is the largest integer such that $U^\beta \supset \tilde{C}_r^\alpha$ for some α . Let A be defined as the union of D_1 and $\bigcup_{\beta} U^\beta$. Clearly (D.3) implies that $A \subset D_2$.

We shall establish (D.3) by induction on r . Clearly (D.3) holds if $r \leq 1$ using (2.4) and (2.5). Now let U^β and $r > 1$ be given and let $\tilde{C}_r^\alpha, i = 1, 2, \dots, m$ denote those \tilde{C}_r^α

contained in U^β . We define Q^γ to be the components of

$$U^\beta \sim \bigcup_i^m \tilde{C}_r^{\alpha_i}.$$

By induction,

$$\text{diam } Q^\gamma \leq 10d_{r-1}$$

which is much less than

$$\text{dist}(\tilde{C}_r^{\alpha_i}, \tilde{C}_r^{\alpha_j}) \geq 2d_{r+1} - 8d_r - 4$$

for d_0 large. In the last inequality we have used (2.5), (2.7), and (D.2). Therefore each Q^γ can meet at most one $\tilde{C}_r^{\alpha_i}$. Since U^β is connected, it follows that U^β contains exactly one $C_r^{\alpha_1}$ (i.e. $m=1$). Therefore if d_0 is sufficiently large,

$$\begin{aligned} \text{diam } U^\beta &\leq \text{diam } \tilde{C}_r^{\alpha_1} + 20d_{r-1} \\ &\leq 9d_r + 4 + 20d_{r-1} \\ &\leq 10d_r. \quad \square \end{aligned}$$

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