### 266. Hermann Weyl: On Hodge's theory of harmonic integrals.

Hodge's fundamental existence theorem for harmonic integrals on Riemannian manifolds of any dimensionality is proved by the parametrix method. (The proof incorporated in Hodge's recent book on *Harmonic Integrals*, Cambridge, 1941, is wrong.) (Received July 1, 1942.)

267. Hassler Whitney: Differentiability of the remainder term in Taylor's formula.

If f(x) is of class  $C^m$ , and  $1 \le n \le m$ , then  $f(x) = \sum_{i=0}^{n-1} f^{(i)}(0) x^i/i! + x^n f_n(x)/n!$ . It is shown that  $f_n(x)$  is of class  $C^{m-n}$ , but not necessarily of higher class, and  $\lim_{x\to 0} x^k f^{(m-n+k)}(x) = 0$   $(k=1, \cdots, n)$ . A converse is true. A similar theorem holds in more dimensions. (Received July 28, 1942.)

268. Hassler Whitney: Note on differentiable even functions.

It is shown that an even function f(x) of class  $C^{2s}$  (or class  $C^{\infty}$ , or analytic) may be written as  $g(x^2)$ , with g of class  $C^s$  (or class  $C^{\infty}$ , or analytic). (Received July 28, 1942.)

269. Hassler Whitney: The general type of singularity of a set of 2n-1 smooth functions of n variables.

Let f be a mapping of class  $C^1$  of an n-manifold  $M^n$  into an  $M^{2n-1}$ . Then arbitrarily near f is a mapping f', regular except at isolated singular points; at each of these, a certain condition (C) holds. (C) involves first and second derivatives, but is independent of the coordinate system employed. If (C) holds at p, and the mapping is of class  $C^{4r+8}$  (or class  $C^{\infty}$ , or analytic), then coordinate systems about p and f(p), of class  $C^r$  (or class  $C^{\infty}$ , or analytic), exist such that the mapping is exactly  $y_1 = x_1^2$ ,  $y_i = x_i$ ,  $y_{n+i-1} = x_1x_i$  ( $i = 2, \cdots, n$ ). (Received July 28, 1942.)

#### APPLIED MATHEMATICS

270. Stefan Bergman: Operators in the theory of differential equations and their application. I.

By introducing  $u=x\cos\theta+y\sin\theta$ ,  $v=-x\sin\theta+y\cos\theta$  and  $\xi=(\sigma/2k)+\theta$ ,  $\eta = (\sigma/2k) - \theta$ , where  $\sigma_x = \sigma + k \sin 2\theta$ ,  $\sigma_y = \sigma - k \sin 2\theta$ ,  $\tau_{xy} = -k \cos 2\theta$  the equations of the theory of plasticity can be written in the form  $(\partial^2 u/\partial \xi \partial \eta) - u/4 = 0$ ,  $(\partial^2 v/\partial \xi \partial \eta) - v/4$ =0 (see Geiringer and Prager, Ergebnisse der exakten Naturwissenschaften, vol. 13. p. 350). Here  $\sigma_x$ ,  $\sigma_y$ ,  $\tau_{xy}$  are stresses, x, y, cartesian coordinates. Particular solutions of these equations can be written in the form  $u(\xi, \eta) = \int_{-1}^{1} \exp(t(\xi \eta)^{1/2}) \{f[\xi(1-t^2)/2]\}$  $+g[n(1-t^2)/2]$   $\{(1-t^2)^{1/2}dt$  where f and g are arbitrary twice continuously differentiable functions of one variable. (See Duke Mathematical Journal, vol. 6 (1940), pp. 538 and 557.) This class of functions possesses a base  $\{u_{\nu}(\xi, \eta)\}$  such that each  $u_{\nu}$  satisfies two (simple) ordinary linear differential equations of second order with rational coefficients. Entire solutions  $u_{\nu}$  of the above partial differential equation are such that every u defined in a convex domain can be approximated by sums of the form  $\sum_{\nu=1}^{n} a_{\nu}^{(n)} u_{\nu}$ . The author indicates an approximation procedure of a function u given by its boundary values. These functions u possess singularities which can be characterized in a way analogous to that in Comptes Rendus de l'Académie des Sciences, vol. 205 (1937), pp. 1360-1362. (Received June 3, 1942.)

## 271. Stefan Bergman: Operators in the theory of partial differential equations and their application. II.

Let  $v(x, y)e^{i\theta(x,y)}$  denote the velocity vector of an irrotational steady flow of compressible fluid. Let  $\zeta = \Lambda(v) + i\theta$ ,  $\bar{\zeta} = \Lambda(v) - i\theta$ , where  $d\Lambda(v)/dv = [1-M^2]^{1/2}/v$ , and  $M = v/[d_0^2 - (1/2)(k-1)v^2]^{1/2}$ ,  $d_0$  and k being constants. Finally: let  $E^* = 1 + t\zeta^{1/2}Q(\zeta, \bar{\zeta}, t\zeta^{1/2})$  where  $Q(\zeta, \bar{\zeta}, p)$  is an (arbitrary) solution of  $Q_p\bar{\zeta} + 2p(Q_\zeta\bar{\zeta} + PQ) + 2F = 0$ , and Q is supposed to be an odd function of p. Then  $\psi(v, \theta) = Re\{\int_{-1}^1 T(\zeta+\bar{\zeta})E^*(\zeta, \bar{\zeta}, t)f[(1/2)\zeta(1-t^2)]dt/(1-t^2)^{1/2}\}$  where f is an arbitrary analytic function of one complex variable is the stream function of a suitable subsonic flow, and the stream function of every flow can be represented in the above form. f and f are suitable functions of f and f are suitable functions of f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete. The author indicates a method of determining the constants f and f are complete.

#### 272. Vladimir Morkovin: On the deflection of anisotropic thin plates.

Deflections w of an anisotropic plate (with one plane of elastic symmetry) bounded by an analytic curve  $C_0$  are considered. The general solution of the differential equation for w is known to be expressible in terms of two analytic functions  $f_1(z_1)$  and  $f_2(z_2)$ , where the complex variables  $z_1$  and  $z_2$  are related to the variable  $z_0$  of the original plane by  $z_k = p_k z_0 + \bar{q}_k \bar{z}_0$ , the constants  $p_k$  and  $q_k$  depending on the material of the plate. (See S. N. Lechnitzky, Journal of Applied Mathematics and Mechanics, (n. s.), vol. 2 (1939), pp. 181–210.) Transformations  $z_k = \omega_k(\zeta_k)$  are found which make any point on  $C_0$  correspond to points  $e^{i\theta}$  on the circumferences  $\gamma_k$  of unit radii in new  $\zeta_1$  and  $\zeta_2$  planes having the same polar angle  $\theta$ , and which are conformal in some neighborhoods of  $\gamma_k$ . Then the functions  $\phi_k(\zeta_k) \equiv f_k(z_k)$  can be determined from the two given boundary conditions if these are expressed in terms of  $e^{i\theta}$ . A detailed solution illustrating this general procedure is carried out in the case of a clamped elliptic plate with polynomial loading. (Received July 31, 1942.)

#### GEOMETRY

## 273. H. S. M. Coxeter: A geometrical background for the description of de Sitter's world.

This paper begins with an elementary treatment of the process by which an elliptic or hyperbolic metric in the plane at infinity of affine space induces a Euclidean or Minkowskian metric in the whole space. The various kinds of sphere are defined, and are seen to provide models for non-Euclidean planes, including the "exterior-hyperbolic" plane which is a two-dimensional de Sitter's world. (See Eddington, *The Mathematical Theory of Relativity*, 1924, p. 165.) Then comes a simple proof of Study's theorem to the effect that one side of a triangle is greater than the sum of the other two, and finally a discussion of some cosmological paradoxes. (Received July 31, 1942.)

# 274. J.J.DeCicco: New proofs of the theorems of Beltrami and Kasner on linear families.

Here new proofs of the theorems of Beltrami and Kasner on linear families of