

$$\{ [f(x + (D - 1)\theta) - f(x - \theta D)]F(y) \}_{y=0} = \theta \frac{df(x)}{dx}.$$

Blissard's remark, "An equation which has a representative quantity is not susceptible to any algebraic operation by which the indices would be affected," becomes

$$(Df)^2 \neq D^2f.$$

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ON FOURTH ORDER SELF-ADJOINT DIFFERENCE SYSTEMS*

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A linear difference expression for which the differential transform is self-adjoint (anti-self-adjoint) we shall call self-adjoint (anti-self-adjoint).† We choose two fourth order difference equations

$$(1) \quad \begin{aligned} L^+(u) &\equiv p(x)[u(x+2) + u(x-2)] \\ &+ \lambda[u(x+1) + u(x-1)] + R(x)u(x) = 0, \end{aligned}$$

$$(2) \quad \begin{aligned} L^-(u) &\equiv p(x)[u(x+2) - u(x-2)] \\ &+ \lambda[u(x+1) - u(x-1)] = 0, \end{aligned}$$

where $L^+(u)$ is self-adjoint and $L^-(u)$ anti-self-adjoint for the range $(x = a, a+1, \dots, b-1; b-a \geq 4)$. $R(x)$ and $p(x)$ are both real, $p(x)$ being a non-vanishing periodic function of period two; λ is a parameter.

Let the functions (y_1, y_2, y_3, y_4) constitute a fundamental set of solutions for either (1) or (2), and (w_1, w_2, w_3, w_4) the set adjoint to it. The two sets are related by the equations

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† J. Kaucky, *Sur les équations aux différences finies qui sont identiques à leurs adjointes*, Publications of the Faculty of Sciences, University of Masaryk, No. 22 (1922). For a discussion of adjoint differential expressions of infinite order, see H. T. Davis, *The Theory of Linear Operators*, 1936, pp. 474-475.

$$\begin{aligned}
 w_i(x)p(x) &= \frac{A_{4j}(x+1)}{A(x+1)}, \\
 (3) \quad A(x) &= \begin{vmatrix} y_1(x-2) \cdots y_4(x-2) \\ \cdots \cdots \cdots \cdots \cdots \\ y_1(x+1) \cdots y_4(x+1) \end{vmatrix},
 \end{aligned}$$

in which $A(x)$ denotes the Casorati determinant and $A_{ij}(x)$ the cofactors of its elements. Since w_i satisfies the equation $L(u) = 0$, one can write

$$(4) \quad w_i = c_{i1}y_1(x) + \cdots + c_{i4}y_4(x), \quad (i = 1, 2, 3, 4).$$

The c_{ij} have properties stated in the following theorem.*

THEOREM 1. *The matrix C^+ of the substitution (4) is skew-symmetric and C^- symmetric. Further, there is a set of relations involving c_{ij}^+ and the second order minors of $A^+(x)$ from which the c_{ij}^+ may be calculated explicitly:*

$$\begin{aligned}
 \sum_{i>j=1}^4 c_{ij}^+ \begin{pmatrix} y_i & y_j \\ -2 & -1 \end{pmatrix}^- = 0, \quad \sum c_{ij}^+ \begin{pmatrix} y_i & y_j \\ -2 & 0 \end{pmatrix}^- = -\frac{1}{p(x)}, \\
 \sum c_{ij}^+ \begin{pmatrix} y_i & y_j \\ -2 & 1 \end{pmatrix}^- = \frac{\lambda}{p(x)p(x+1)},
 \end{aligned}$$

$$\begin{aligned}
 \sum_{j>i=1}^4 c_{ij}^+ \begin{pmatrix} y_i & y_j \\ -1 & 0 \end{pmatrix}^- = 0, \quad \sum c_{ij}^+ \begin{pmatrix} y_i & y_j \\ -1 & 1 \end{pmatrix}^- = -\frac{1}{p(x+1)}, \\
 \sum c_{ij}^+ \begin{pmatrix} y_i & y_j \\ 0 & 1 \end{pmatrix}^- = 0.
 \end{aligned}$$

We resolve (4) to get

$$(5) \quad c_{ij}^\pm = [w_i(x-2)A_{1j}(x) + \cdots + w_i(x+1)A_{4j}(x)] \frac{1}{A(x)}.$$

We use (1, 2, 3) to furnish the following relations

* We introduce a (+) and (-) convention to distinguish between quantities associated with (1) and (2) respectively. Also we set

$$u(x+r)v(x+s) \pm v(x+r)u(x+s) = \begin{pmatrix} u & v \\ r & s \end{pmatrix}^\pm.$$

$$\begin{aligned}
 w_j(x-1)p(x-1) &= \frac{A_{4j}(x)}{A(x)}, \\
 w_j(x-2)p(x) + \lambda w_j(x-1) &= \frac{A_{3j}(x)}{A(x)}, \\
 \pm w_j^\pm(x)p(x) &= \frac{A_{1j}^\pm(x)}{A^\pm(x)}, \\
 \mp [\lambda w_j^\pm(x) + p(x+1)w_j^\pm(x+1)] &= \frac{A_{2j}^\pm(x)}{A^\pm(x)}.
 \end{aligned}
 \tag{6}$$

Combining (5) and (6) we find

$$\begin{aligned}
 c_{ij}^\pm &= p(x+1) \begin{pmatrix} w_i^\pm & w_j^\pm \\ 1 & -1 \end{pmatrix}^\mp + p(x) \begin{pmatrix} w_i^\pm & w_j^\pm \\ 0 & -2 \end{pmatrix}^\mp \\
 &\quad + \lambda \begin{pmatrix} w_i^\pm & w_j^\pm \\ 0 & -1 \end{pmatrix}^\mp.
 \end{aligned}
 \tag{7}$$

Obviously we have $c_{ij}^+ = -c_{ji}^+$ and $c_{ij}^- = c_{ji}^-$. Combine (3) and (4) to eliminate the w 's. These equations together with $L^+(y_i^\pm) = 0$ yield the set of equations involving the minors. The sixth order determinant composed of the two-rowed minors is non-vanishing ($= [A^+(x)]^3$).*

In the development of adjoint difference systems the formula

$$\sum f(x+1) = f(x) + \sum f(x)
 \tag{8}$$

is used to provide a Lagrange relation

$$\sum_{x=a}^{b-1} [vL(u) - uL(v)] = \Pi(u, v) = U_1U_8 + \dots + U_8V_1,
 \tag{9}$$

in which the U_i are eight linearly independent forms arbitrarily chosen:

$$\begin{aligned}
 U_i &= a_{i1}u(a-2) + \dots + a_{i4}u(a+1) \\
 &\quad + b_{i1}u(b-2) + \dots + b_{i4}u(b+1).
 \end{aligned}
 \tag{10}$$

If the two systems

$$\begin{aligned}
 L(u) &= 0, & U_1 &= U_2 = U_3 = U_4 = 0, \\
 L(v) &= 0, & V_1 &= V_2 = V_3 = V_4 = 0,
 \end{aligned}
 \tag{11}$$

* Turnbull, *The Theory of Determinants, Matrices, and Invariants*, 1929, p. 87.

are equivalent, we shall call them self-adjoint difference systems.

THEOREM 2. *Given the fourth order difference systems composed of the equations*

$$L^\pm(u) = 0, \quad U_i = 0, \quad (i = 1, 2, 3, 4),$$

defined in (1), (2), (10), let $u(x)$ and $v(x)$ be any pair of functions satisfying $U_i = 0$; then $\Pi(u, v) \equiv 0$ is a necessary and sufficient condition that the given system be self-adjoint.

That the above condition is necessary needs no proof. Let (u_1, u_2, u_3, u_4) be four linearly independent functions satisfying $U_i = 0$ and u any linear combination of them. Through substitution the identity $\Pi(u, v) \equiv 0$ gives

$$(12) \quad U_5(u_i)V_4(u) + \cdots + U_8(u_i)V_1(u) = 0, \quad (i = 1, 2, 3, 4).$$

Since the set (U_1, U_2, \cdots, U_8) is linearly independent, it follows that the four systems of constants comprising the coefficients of the V 's are linearly independent and we have $V_1(u) = V_2(u) = V_3(u) = V_4(u) = 0$. A similar argument shows that any function satisfying the given boundary conditions will also satisfy the adjoint boundary conditions.

We record some examples which fulfill the condition for self-adjointness.

$$(13) \quad \begin{aligned} U_1^\pm &= a_{11}^\pm u(a-2) + a_{12}u(a-1) + p(a)u(a) \\ &\quad + b_{13}u(b) + b_{14}u(b+1) = 0, \\ U_2^\pm &= \pm a_{12}u(a-2) + a_{22}^\pm u(a-1) + \lambda u(a) \\ &\quad + p(a+1)u(a+1) + b_{24}u(b+1) = 0, \\ U_3^\pm &= \pm b_{13}u(a-2) \pm p(b)u(b-2) \pm \lambda u(b-1) \\ &\quad + b_{33}^\pm u(b) + b_{34}u(b+1) = 0, \\ U_4^\pm &= \pm b_{14}u(a-2) \pm b_{24}u(a-1) \pm p(b+1)u(b-1) \\ &\quad \pm b_{34}u(b) + b_{44}^\pm u(b+1) = 0, \end{aligned}$$

with the agreement that $a_{11}^- = a_{22}^- = b_{33}^- = b_{44}^- = 0$. For self-adjoint Sturmian boundary conditions we make the added restrictions, $b_{13} = b_{14} = b_{24} = 0$.

We now introduce a function $G^\pm(x, t)$ defined for $(a-1 < x < b)$,

($a \leqq t < b$), satisfying the given boundary conditions and for which $L^\pm [G^\pm(i, j)] = \delta_{ij}$.

THEOREM 3. *Let $\lambda_1^\pm, \lambda_2^\pm, \dots$ be sets of characteristic values for the systems $L^\pm(u) = 0, U_i = 0$ defined in (13). There exists a λ_k^\pm in the interval $(\lambda_i^- \leqq \lambda_k^\pm < \lambda_{i+1}^-)$ provided*

$$G^-(b, b - 1) - G^+(b, b - 1) - \sum_{x=a}^{b-1} [G^-(x - 1, x) + G^+(x - 1, x)] \neq 0, \tag{13}$$

$(\lambda_i^- \leqq \lambda \leqq \lambda_{i+1}^-).$

If $(b - a)$ is an odd integer, then every value of λ^- is a characteristic value.

Let $D^\pm(\lambda) = 0$ be the characteristic equations for (13). By writing $D^\pm(\lambda)$ in determinant form one finds

$$(14) \quad \frac{dD^\pm}{d\lambda} = 2D^\pm \left[G^\pm(b, b - 1) \pm \sum_{x=a}^{b-1} G^\pm(x - 1, x) \right].$$

This relation enables us to write

$$(15) \quad \frac{d}{d\lambda} \left(\frac{D^-}{D^+} \right) = \frac{2D^-}{D^+} \left[G^-(b, b - 1) - G^+(b, b - 1) - \sum_{x=a}^{b-1} \{ G^-(x - 1, x) + G^+(x - 1, x) \} \right].$$

Between the two real consecutive zeros λ_i^- and λ_{i+1}^- of $D^-(\lambda)$ either the bracketed expression or $D^+(\lambda)$ must vanish. By assumption the bracketed expression does not vanish.

The proof of the final statement in the theorem consists merely in noticing that $D^-(\lambda)$ for this case is a skew-symmetric determinant of odd order.