# TORIC FIBRATIONS AND MIRROR SYMMETRY 

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#### Abstract

The relation between the quantum $\mathcal{D}$-modules of a smooth variety $X$ and a toric bundle is studied here. We describe the relation completely when $X$ is a semi-ample complete intersection in a toric variety. In this case, we obtain all the relations in the small quantum cohomology ring of the bundle.


## 1. Introduction and Goals

For a smooth, projective variety $Y$ we denote by $Y_{k, \beta}$ the moduli stack of rational stable maps of class $\beta \in H_{2}(Y, \mathbb{Z})$ with $k$-markings (Fulton et al [8]) and [ $Y_{k, \beta}$ ] its virtual fundamental class (Behrend et al [3], Li et al [13]). Genus zero GromovWitten invariants are defined as appropriate integrals over $\left[Y_{k, \beta}\right]$. We let $e: Y_{1, \beta} \rightarrow$ $Y$ be the evaluation map, $\psi$ - the first chern class of the cotangent line bundle on $Y_{1, \beta}$ and $\mathrm{ft}: Y_{1, \beta} \rightarrow Y_{0, \beta}$ - the forgetful morphism.
The formal completion of an arbitrary ring $\mathcal{R}$ along the semigroup $M Y$ of the rational curves of $Y$ is defined to be

$$
\begin{equation*}
\mathcal{R}\left[\left[q^{\beta}\right]\right]:=\left\{\sum_{\beta \in \mathrm{MY}} a_{\beta} q^{\beta}, \quad a_{\beta} \in \mathcal{A}, \quad \beta-\text { effective }\right\} \tag{1}
\end{equation*}
$$

where $\beta \in H_{2}(Y, \mathbb{Z})$ is effective if it is a positive linear combination of rational curves. For each $\beta$, the set of $\alpha$ such that $\alpha$ and $\beta-\alpha$ are both effective is finite, hence $\mathcal{R}\left[\left[q^{\beta}\right]\right]$ behaves like a power series. Alternatively, we may define

$$
q^{\beta}:=q_{1}^{d_{1}} \cdot \ldots \cdot q_{k}^{d_{k}}=\exp \left(t_{1} d_{1}+\ldots+t_{k} d_{k}\right)
$$

where $\left\{d_{1}, d_{2}, \ldots, d_{k}\right\}$ are the coordinates of $\beta$ relative to the dual of a nef basis $\left\{p_{1}, \ldots, p_{k}\right\}$ of $H^{2}(Y, \mathbb{Q})$..
Let $*$ denote the small quantum product of $Y$. The small quantum cohomology ring

$$
\left(Q H_{s}^{*} Y, *\right)
$$

is a deformation of the cohomology ring $\left(H^{*}\left(Y, \mathbb{Q}\left[q^{\beta}\right]\right), \cup\right)$. Its structural constants are three point Gromov-Witten invariants of genus zero. Let $\hbar$ be a formal variable and

$$
J_{\beta}(Y):=e_{*}\left(\frac{\left[Y_{1, \beta}\right]}{\hbar(\hbar-\psi)}\right)=\sum_{k=0}^{\infty} \frac{1}{\hbar^{2+k}} e_{*}\left(\psi^{k} \cap\left[Y_{1, \beta}\right]\right)
$$

[^0]The sum is finite for dimension reasons. For $t=\left(t_{0}, t_{1}, \ldots, t_{k}\right)$, let

$$
t p:=t_{0}+\sum_{i=1}^{k} t_{i} p_{i}
$$

The $\mathcal{D}$-module for the quantum differential equation of $Y$

$$
1 \leq i \leq k, \hbar \partial / \partial t_{i}=p_{i} *
$$

is generated by (Givental [10])

$$
J(Y)=\exp \left(\frac{t p}{\hbar}\right) \sum_{\beta \in H_{2}(Y, \mathbb{Z})} q^{\beta} J_{\beta}(Y)
$$

where we use the convention $J_{0}=1$. The generator $J(Y)$ encodes all of the genus zero, one marking Gromov-Witten invariants and gravitational descendants of $Y$. The generator $J(Y)$ is an element of the completion $H^{*}(Y, \mathbb{Q})[t]\left[\left[q^{\beta}\right]\right]$ that may be used to produce relations in $Q H_{s}^{*} Y$ in the following way: let

$$
\mathcal{P}\left(\hbar, \hbar \partial / \partial t_{i}, q_{i}\right)
$$

be a polynomial differential operator where $q_{i}$ and $\hbar$ act via multiplication and $q_{i}=e^{t_{i}}$ are on the left of derivatives. If

$$
\mathcal{P}\left(\hbar, \hbar \partial / \partial t_{i}, q_{i}\right) J(Y)=0
$$

then

$$
\mathcal{P}\left(0, p_{i}, q_{i}\right)=0
$$

is a relation in the small quantum cohomology ring $Q H_{s}^{*} Y$.
If $Y$ is a complete intersection in a toric variety, $J(Y)$ is related to an explicit hypergeometric series $I(Y)$ via a change of variables (Givental [8], Lian et al [12],[13]). Furthermore, if $Y$ is Fano then the change of variables is trivial, i.e.

$$
J(Y)=I(Y)
$$

Since $I(Y)$ is known explicitly, this yields two immediate benefits.
(1) The one point Gromov-Witten invariants and gravitational descendants of $Y$ are determined completely.
(2) Differential operators that annihilate $I(Y)$ are easy to find, hence producing relations in the small quantum cohomology ring of $Y$.

In this paper we seek to relativize these results for Fano toric bundles, hence extending the results of the papers Elezi [6],[7]

## 2. Toric Bundles and Mirror Theorems

Toric varieties and bundles. We follow the approach and the terminology of Oda [15]. Let $\mathbb{M} \simeq \mathbb{Z}^{m}$ be a free abelian group of $\operatorname{rank} m, \mathbb{N}=\operatorname{Hom}(\mathbb{M}, \mathbb{Z})$ its dual, and $<,>: \mathbb{M} \times \mathbb{N} \mapsto \mathbb{Z}$ the pairing between them. Let $Y$ be an m-dimensional smooth, toric variety determined by a fan $\Sigma \subset \mathbb{N} \otimes \mathbb{R}$. Denote by

$$
\Sigma(1)=\left\{\rho_{1}, \ldots, \rho_{m}, \rho_{m+1}, \ldots, \rho_{r=m+k}\right\}
$$

the one dimensional cones of $\Sigma$ and $D_{1}, \ldots, D_{r}$ the corresponding toric divisors. Let $v_{i}=\left(v_{i 1}, \ldots, v_{i m}\right)$ be the first lattice point along the ray $\rho_{i}$. Let

$$
\left\{a_{1}, a_{2}, \ldots, a_{k}\right\}
$$

with $a_{j}:=\left(a_{1 j}, a_{2 j}, \ldots, a_{m j}, a_{m+1 j}, \ldots, a_{r j}\right)$ be a basis of the lattice of relations $\Lambda$ between $v_{1}, \ldots, v_{r}$. There is a short exact sequence

$$
\begin{equation*}
0 \rightarrow \Lambda \rightarrow \mathbb{Z}^{\Sigma(1)} \xrightarrow{h} \mathbb{N} \rightarrow 0 \tag{2}
\end{equation*}
$$

where $h\left(c_{1}, c_{2}, \ldots c_{r}\right)=c_{1} v_{1}+\ldots+c_{r} v_{r}$. The lattice $\Lambda$ may be identified with $\operatorname{Hom}\left(A_{m-1}(Y), \mathbb{Z}\right) \simeq H_{2}(Y, \mathbb{Z})$. Under this isomorphism, $a_{i j}$ is the intersection of $a_{j}$, when interpreted as a two dimensional cycle, with the toric divisor $D_{i}$. We choose $a_{j}$ so that $\left\{a_{1}, \ldots, a_{k}\right\}$ is a generating set for the Mori cone of classes of effective curves. Then $a_{i 1}, \ldots, a_{i k}$ are the coordinates of $D_{i}$ with respect to the nef basis $\left\{p_{1}, \ldots, p_{k}\right\}$ dual to $\left\{a_{1}, \ldots, a_{k}\right\}$.
Assume that $\rho_{1}, \ldots, \rho_{m}$ generate a maximal dimensional cone in $\Sigma$. Since $Y$ is smooth, $\left\{v_{1}, v_{2}, \ldots, v_{m}\right\}$ forms a $\mathbb{Z}$-basis of $\mathbb{N}$ and the absolute value of the matrix

$$
\left(a_{i j}\right) ; i=m+1, \ldots, r ; j=1,2, \ldots, k
$$

is 1 .
The cohomology ring $H^{*}(Y, \mathbb{Z})$ is generated by the divisors $D_{1}, \ldots, D_{r}$ subject to the following two types of relations:
Type One: Whenever $\left\{\rho_{j_{1}}, \ldots, \rho_{j_{s}}\right\}$ do not generate a cone in $\Sigma$, the intersection

$$
\begin{equation*}
D_{j_{1}} \cdot \ldots \cdot D_{j_{s}}=0 \tag{3}
\end{equation*}
$$

Type Two: For each $1 \leq i \leq m$,

$$
\begin{equation*}
D_{i}=\sum_{j=1}^{k} a_{i j} p_{j} \tag{4}
\end{equation*}
$$

From the short exact sequence (2) we obtain

$$
\begin{equation*}
0 \rightarrow \mathbb{T}^{k} \xrightarrow{\alpha} \mathbb{T}^{r} \xrightarrow{\beta} \mathbb{T}^{m} \rightarrow 0 \tag{5}
\end{equation*}
$$

where the maps are defined as follows:

$$
\alpha\left(t_{1}, t_{2}, \ldots t_{k}\right)=\left(\prod_{i=1}^{k} t_{i}^{a_{1 i}}, \ldots, \prod_{i=1}^{k} t_{i}^{a_{r i}}\right), \beta\left(t_{1}, \ldots t_{r}\right)=\left(\prod_{i=1}^{r} t_{i}^{v_{i 1}}, \ldots, \prod_{i=1}^{r} t_{i}^{v_{i m}}\right)
$$

Let $Z(\Sigma) \subset \mathbb{C}^{r}$ be the variety whose ideal is generated by the products of those variables which do not generate a cone in $\Sigma$. The toric variety $Y$ is the geometric quotient (Cox [5])

$$
\mathbb{C}^{r}-Z(\Sigma) / / \mathbb{T}^{k}
$$

where the torus acts as follows

$$
\begin{equation*}
t \cdot x=\left(\prod_{i=1}^{k} t_{i}^{a_{1 i}} x_{1}, \ldots, \prod_{i=1}^{k} t_{i}^{a_{r i}} x_{r}\right) \tag{6}
\end{equation*}
$$

The short exact sequence (5) yields an action of the quotient $\mathbb{T}:=\mathbb{T}^{m}$ on $Y$.
The first chern class of the tangent bundle to $Y$ is equal to

$$
\sum_{i=1}^{r} D_{i}=\sum_{i=1}^{k} n_{i} p_{i}
$$

The toric variety $Y$ is Fano iff $n_{i}>0$ for all $i$.

We relativize the previous construction as follows. Consider the principal $\mathbb{T}$-bundle

$$
\mathbb{E}:=\oplus_{i=1}^{m}\left(L_{i}-\{0\}\right) \rightarrow X
$$

where $L_{i}$ are line bundles over a smooth, projective variety $X$. Let $\mathbb{T}$ act fibrewisely on $\mathbb{E}$ and the diagonally on the first $m$-homogeneous coordinates of $Y$. The quotient space

$$
Y(\mathbb{E}):=\mathbb{E} \times_{\mathbb{T}} Y
$$

is a toric bundles over $X$ with fiber isomorphic to $Y$. The bundle $Y(\mathbb{E})$ inherits a T-action.
There is a projection map $\pi: Y(\mathbb{E}) \rightarrow Y$. The maximal cone generated by $\left\{\rho_{1}, \rho_{2}, \ldots, \rho_{m}\right\}$ determines a $\mathbb{T}$ fixed point $q$ in $Y$ whose homogeneous coordinates are $(0,0, \ldots, 0,1,1, \ldots 1)$. In the relativized setting, the $\mathbb{T}$-equivariant inclusion

$$
q \hookrightarrow Y
$$

yields a map

$$
q(\mathbb{E}) \simeq X \stackrel{s}{\hookrightarrow} Y(\mathbb{E})
$$

which is a section of $\pi$. This is also a fixed point component for the action of $\mathbb{T}$ on $Y(\mathbb{E})$. The other $\mathbb{T}$-fixed points of $Y$ yield sections of $\pi$ and these are all the fixed point components.
Toric divisors lift to divisors in $Y(\mathbb{E})$; these liftings will be denoted by the same letter in this paper. It was shown in Sankaran and Uma [17] that the two types of relations (3) and (4) lift in a natural way in $H^{*}(Y(\mathbb{E}), \mathbb{Z})$; namely

$$
D_{j_{1}} \cdot \ldots \cdot D_{j_{s}}=0
$$

whenever $\left\{\rho_{j_{1}}, \ldots, \rho_{j_{s}}\right\}$ do not generate a cone in $\Sigma$, and

$$
D_{i}=\sum_{j=1}^{k} a_{i j} p_{j}+c_{1}\left(L_{i}\right)
$$

for each $1 \leq i \leq m$, where as in the case of $H^{*}(Y, \mathbb{Z})$ the divisors

$$
p_{1}, \ldots, p_{k}
$$

generate freely $H^{*}(Y(\mathbb{E}), \mathbb{Z})$. In fact, there is a simple relation between the $\mathbb{T}$ equivariant cohomology of $Y$ and the cohomology of $Y(\mathbb{E})$ which will be used throughout this paper. Recall, that the rational cohomology of the classifying space $B \mathbb{T}$ is $\mathbb{Q}\left[\lambda_{1}, \ldots \lambda_{m}\right]$ where $\lambda_{i}$ is the first chern class of the equivariant line bundle corresponding to the character

$$
\nu_{i}: \mathbb{T} \rightarrow \mathbb{C}^{*} \nu_{i}\left(t_{1}, \ldots, t_{m}\right)=t_{i}
$$

A relation in the equivariant cohomology ring of $Y$ becomes a relation in $H^{*}(Y(\mathbb{E}))$ after substituting $c_{1}\left(L_{i}\right)$ for $\lambda_{i}$.

We may assume that $L_{i}=\mathcal{O}_{X}, i>m$ without loss of generality. This is due to the fact that $\rho_{1}, \ldots, \rho_{m}$ generate a maximal cone in $\Sigma$.

The quantum $\mathcal{D}$-module structure of a toric bundle. The generator $J$ of a quantum $\mathcal{D}$-structure is weighted by the lattice points of the Mori cone. Hence we first study the relation between the Mori cones of $Y$ and $Y(\mathbb{E})$.

Lemma 1. If $L_{i}^{*}$ are generated by global sections, then the liftings of the nef divisors $p_{1}, \ldots, p_{k}$ in $Y(\mathbb{E})$ are also nef. Furthermore, the Mori cone of $Y(\mathbb{E})$ is a direct sum of the Mori cone of $X$, embedded via the section $s$, and the Mori cone of the fiber $Y$.

Proof. In toric varieties, every nef divisor $p$ is generated by global sections (Oda [14]). Let $x_{1}, x_{2}, \ldots, x_{r}$ be homogeneous coordinates in $Y$. The vector space of global sections $H^{0}(\mathcal{O}(p))$ has a monomial basis

$$
\prod_{i=1}^{r} x_{i}^{m_{i}}
$$

Let $\left\{\phi_{i j}\right\}$ be a collection of generating sections for the line bundles $L_{i}^{*}$. The "monomials"

$$
\prod_{i=1}^{r}\left(x_{i} \phi_{i j}\right)^{m_{i}}
$$

are generating sections the line bundle

$$
\prod_{i=1}^{r}\left(\mathcal{O}\left(D_{i}\right) \otimes\left(L_{i}^{*}\right)\right)^{m_{i}}
$$

which is isomorphic to $\mathcal{O}(p)$ in $Y(\mathbb{E})$. Thus $p$ lifts to a nef divisor in $Y(\mathbb{E})$.
This shows that the addition of $p_{1}, \ldots, p_{k}$ to a nef basis $\left\{p_{k+1}, \ldots, p_{l}\right\}$ of $X$ yields a nef basis

$$
\left\{p_{1}, \ldots, p_{l}\right\}
$$

of $Y(\mathbb{E})$. Now for a curve $C \subset Y(\mathbb{E})$ we have

$$
\pi_{*}\left([C]-s_{*}\left(\pi_{*}([C])\right)\right)=0
$$

Notice that the restrictions of the divisors $p_{1}, p_{2}, \ldots, p_{k}$ in the section $q(\mathbb{E})$ are all zero since they may be written as $\mathbb{Z}$-linear combinations of $D_{m+1}, \ldots, D_{m+k}$. Hence $\forall i=1,2, \ldots, k, p_{i} \cdot\left([C]-s_{*}\left(\pi_{*}([C])\right)\right) \geq 0$ and we have a unique decomposition

$$
[C]=s_{*}\left(\pi_{*}([C])\right)+\left[C^{\prime}\right]
$$

where $\left[C^{\prime}\right]$ and $\pi_{*}([C])$ are curve classes respectively in the fiber of $\pi$ and $X$.
We introduce a "mixed" $I(Y(\mathbb{E}))$ that admits contributions from both $J(X)$ and an $\mathbb{E}$-twisted $J(Y)$. Let $(\nu, d)$ denote a curve class in the Mori cone of $Y(\mathbb{E})$, with $\nu$ a curve class in the fiber of $\pi$ and $d$ a curve class in $X$.
Define

$$
I(Y(\mathbb{E})):=\exp \left(\frac{t p}{\hbar}\right) \sum_{(d, \nu)} q_{1}{ }^{\nu} q_{2}{ }^{d} \prod_{i=1}^{m} \frac{\prod_{m=0}^{\infty}\left(D_{i}+m \hbar\right)}{\prod_{m=0}^{D_{i}(\nu, d)}\left(D_{i}+m \hbar\right)} \pi^{*}\left(J_{d}(X)\right)
$$

If $X$ is a point then $Y(\mathbb{E})=Y$. Furthermore, as mentioned in the introduction $J(Y)=I(Y)$ if $Y$ is a Fano toric variety. In this paper we show that the same holds for the relativized $Y(\mathbb{E})$.

Proposition 1. If $X$ is a semi-ample complete intersection in a toric variety, and both $Y$ and $Y(\mathbb{E})$ are Fano, then $J(Y(\mathbb{E}))=I(Y(\mathbb{E}))$.

Proposition 1 will follow as a corollary of another statement which we now formulate and prove.
Let $Z$ be a toric variety, $\tilde{L}_{i}, i=0,1, \ldots, n$ toric line bundles over $Z$ and $\tilde{\mathbb{E}}=\oplus_{i=0}^{n} \tilde{L}_{i}$. The bundle

$$
\pi: Y(\tilde{\mathbb{E}}) \rightarrow Z
$$

is also a toric variety (Oda [15]). The edges of the fan for $Y(\tilde{\mathbb{E}})$ corresponds to the liftings $B_{1}, \ldots, B_{r}$ to $Y(\mathbb{E})$ of the toric base divisors $b_{1}, \ldots, b_{r}$ and the divisors $D_{i}$ from $Y$.
Let $\mathcal{L}_{a}: a=1,2, \ldots, l$ be globally generated line bundles over $Z$ and $X$ the zero locus of a generic section s of

$$
V=\oplus_{a=1}^{l} \mathcal{L}_{a}
$$

Such an $X$ will be called a semi-ample complete intersection. Denote by $L_{i}$ and $\mathbb{E}$ the restrictions of $\tilde{L}_{i}$ and $\tilde{\mathbb{E}}$ to $X$. The total space of $Y(\mathbb{E})$ is easily seen to be the zero locus of the section $\pi^{*}(s)$ of the pull back bundle $\pi^{*}(V)$.

Assume that the line bundles $\tilde{L}_{i}^{*}$ are globally generated and $-K_{Z}-\sum_{a=1}^{l} c_{1}\left(\mathcal{L}_{a}\right)+$ $\sum_{i=0}^{n} c_{1}\left(\tilde{L}_{i}\right)$ is ample. (This will ensure that the conditions of Proposition 1 for the bundle $Y(\mathbb{E})$ over $X$ are satisfied.)

Let $V_{d}$ be the bundle on $Z_{1, d}$ whose fiber over the moduli point $\left(C, x_{1}, f\right)$ is $\oplus_{a} H^{0}\left(f^{*}\left(\mathcal{L}_{a}\right)\right)$. Denote by $s_{V}$ its canonical section induced by $s$, i.e.

$$
s_{V}\left(\left(C, x_{1}, f\right)\right)=f^{*}(s)
$$

The stack theoretic zero section of $s_{V}$ is the disjoint union

$$
\begin{equation*}
Z\left(s_{V}\right)=\coprod_{i_{*}(\beta)=d} X_{1, \beta} \tag{7}
\end{equation*}
$$

The map $i_{*}: H_{2} X \rightarrow H_{2} Z$ is not injective in general, hence the zero locus $Z\left(s_{V}\right)$ may have more then one connected component. An example is the quadric surface in $\mathbb{P}^{3}$. The sum of the virtual fundamental classes $\left[X_{1, \beta}\right]$ is the refined top Chern class of $V_{d}$ with respect to $s_{V}$.
Let $\tilde{V}_{\nu, d}$ and $\tilde{s}_{V}$ be the pull backs of $V_{d}$ and $s_{V}$ via the stack morphism

$$
Y(\tilde{\mathbb{E}})_{1,(\nu, d)} \rightarrow Z_{1, d}
$$

The zero section of $\tilde{s}_{V}$ is the disjoint union

$$
z\left(\tilde{s}_{V}\right)=\coprod_{i_{*}(\beta)=d} Y(\mathbb{E})_{1,(\nu, \beta)}
$$

It follows that

$$
\sum_{i_{*}(\beta)=d}\left[Y(\mathbb{E})_{1,(\nu, \beta)}\right]=c_{\text {top }}\left(\tilde{V}_{\nu, d}\right) \cap\left[Y(\tilde{\mathbb{E}})_{1,(\nu, d)}\right]
$$

Recall that the nef basis $\left\{p_{1}, p_{2}, \ldots, p_{k}, p_{k+1}, \ldots p_{l}\right\}$ of $Y(\mathbb{E})$ is obtained by completing a nef basis $\left\{p_{k+1}, \ldots, p_{l}\right\}$ of $X$. We will use $t p$ to denote both $\sum_{i=1}^{l} t_{i} p_{i}$ and $\sum_{i=k+1}^{l} t_{i} p_{i}$. The difference will be clear from the context.
Consider the following generating functions

$$
J^{V}(Y(\tilde{\mathbb{E}}))=\exp \left(\frac{t p}{\hbar}\right) \sum_{(\nu, d)} q_{1}^{\nu} q_{2}^{d} e_{*}\left(\frac{c_{\mathrm{top}}\left(\tilde{V}_{\nu, d}\right) \cap\left[Y(\tilde{\mathbb{E}})_{1,(\nu, d)}\right]}{\hbar(\hbar-c)}\right)
$$

and

$$
\tilde{I}^{V}(Y(\tilde{\mathbb{E}}))=\exp \left(\frac{t p}{\hbar}\right) \sum_{(\nu, d)} q_{1}^{\nu} q_{2}^{d} \Omega_{\nu, d} \pi^{*} e_{*}\left(\frac{c_{\mathrm{top}}\left(V_{d}\right) \cap\left[Z_{1, d}\right]}{\hbar(\hbar-c)}\right)
$$

where

$$
\Omega_{\nu, d}=\prod_{i=1}^{m} \frac{\prod_{m=0}^{\infty}\left(D_{i}+m \hbar\right)}{\prod_{m=0}^{D_{i}(\nu, d)}\left(D_{i}+m \hbar\right)}
$$

Proposition 2. If $-K_{Y}-\sum_{a=1}^{l} c_{1}\left(\mathcal{L}_{a}\right)-\sum_{i=0}^{n} c_{1}\left(\tilde{L}_{i}\right)$ is ample then

$$
J^{V}((\tilde{\mathbb{E}}))=\tilde{I}^{V}(Y(\tilde{\mathbb{E}}))
$$

Proof. Let

$$
I_{d}^{V}(Z)=\prod_{a} \frac{\prod_{m=-\infty}^{\mathcal{L}_{a}(d)}\left(\mathcal{L}_{a}+m \hbar\right)}{\prod_{m=-\infty}^{0}\left(\mathcal{L}_{a}+m \hbar\right)} \prod_{i} \frac{\prod_{m=-\infty}^{0}\left(B_{i}+m \hbar\right)}{\prod_{m=-\infty}^{B_{i}(d)}\left(B_{i}+m \hbar\right)}
$$

From Givental [9], Lian et al [12], Lian et al [13] we know that $J^{V}(Y(\tilde{E}))$ is related via a mirror transformation to

$$
I^{V}(Y(\tilde{\mathbb{E}}))=\exp \left(\frac{t p}{\hbar}\right) \cdot \sum q_{1}^{\nu} q_{2}^{d} \Omega_{\nu, d} I_{d}^{V}(Z)
$$

Likewise

$$
J^{V}(Z)=\exp \left(\frac{t p}{\hbar}\right) \sum q_{2}^{d} e_{*}\left(\frac{c_{\mathrm{top}}\left(V_{d}\right) \cap\left[Z_{1, d}\right]}{\hbar(\hbar-c)}\right)
$$

is related to

$$
I^{V}(Z)=\exp \left(\frac{t p}{\hbar}\right) \sum q_{2}^{d} I_{d}^{V}(Z)
$$

Since $-K_{Y(\tilde{E})}-\sum_{a} c_{1}\left(\mathcal{L}_{a}\right)$ and $-K_{Z}-\sum_{a} c_{1}\left(\mathcal{L}_{a}\right)$ are ample, the mirror transformations are particularly simple. Indeed, both series can be written as power series of $\hbar^{-1}$ as follows:

$$
I^{V}(Y(\tilde{E}))=1+\frac{P_{1}\left(q_{1}, q_{2}\right)}{\hbar}+o\left(\hbar^{-1}\right), I^{V}(Z)=1+\frac{P_{2}\left(q_{2}\right)}{\hbar}+o\left(\hbar^{-1}\right)
$$

where $P_{1}\left(q_{1}, q_{2}\right), P_{2}\left(q_{2}\right)$ are both polynomials supported respectively in

$$
\Lambda_{1}:=\left\{(\nu, d) \mid\left(-K_{Y(\tilde{E})}-\sum c_{1}\left(\mathcal{L}_{a}\right)\right)=1 ; D_{j} \geq 0, \forall j ; \quad B_{i} \geq 0, \forall i\right\}
$$

and

$$
\Lambda_{2}:=\left\{d \mid\left(-K_{Z}-\sum c_{1}\left(\mathcal{L}_{a}\right)\right)=1 ; B_{i} \geq 0 \forall i\right\}
$$

Then

$$
J^{V}(Y(\tilde{E}))=\exp \left(\frac{-P_{1}\left(q_{1}, q_{2}\right)}{\hbar}\right) I^{V}(Y(\tilde{E}))
$$

and

$$
J^{V}(Z)=\exp \left(\frac{-P_{2}\left(q_{2}\right)}{\hbar}\right) I^{V}(Z)
$$

Simple algebraic manipulations show that

- $c_{1}\left(\tilde{L}_{j}\right) \cdot d=0, \forall d \in \Lambda_{2}, \forall j=1,2, \ldots, n$
- $\Lambda_{1}=\left\{(0, d) \mid d \in \Lambda_{2}\right\}$.

It follows that $\Omega_{0, d}=1, \forall d \in \Lambda_{2}$ hence $P_{1}\left(q_{1}, q_{2}\right)=P_{2}\left(q_{2}\right)$. Notice also that if we expand

$$
\exp \left(\frac{-P_{2}\left(q_{2}\right)}{\hbar}\right)=\sum_{\alpha} c_{\alpha} q_{2}^{\alpha}
$$

then

$$
c_{1}\left(\tilde{L}_{j}\right) \cdot \alpha=0, \forall j=1,2, \ldots, n
$$

Hence for each $(\nu, d) \in M \mathbb{P}(\tilde{V})$ we have $\Omega_{\nu, d}=\Omega_{\nu, d+\alpha}$. Now the proposition follows easily.

Proof. of Proposition 1. We know return to the proof of Proposition 1. Recall that the map

$$
\begin{equation*}
i_{*}: H_{2}(X) \rightarrow H_{2}(Z) \tag{8}
\end{equation*}
$$

is not necessarily injective in general. If it is, then

$$
\left[X_{1, \beta}\right]=c_{\mathrm{top}}\left(V_{i_{*}(\beta)}\right) \cap\left[Y_{1, i_{*}(\beta)}\right]
$$

and

$$
\left[Y(\mathbb{E})_{1,(\nu, \beta)}\right]=c_{\text {top }}\left(\tilde{V}_{\nu, i_{*}(\beta)}\right) \cap\left[Y(\tilde{\mathbb{E}})_{1,\left(\nu, i_{*}(\beta)\right)}\right]
$$

In this case one can easily show that

$$
i_{*}\left(J_{\nu, \beta}(Y(\mathbb{E}))\right)=J_{\nu, i_{*}(\beta)}^{V}(Y(\tilde{\mathbb{E}}))
$$

and

$$
i_{*}\left(I_{\nu, \beta}(Y(\mathbb{E}))\right)=\tilde{I}_{\nu, i_{*}(\beta)}^{V}(Y(\tilde{\mathbb{E}}))
$$

Proposition 2 shows that Proposition 1 holds for complete intersection in toric varieties for which the map (8) is injective.

## 3. Lifting the Quantum Cohomology Structure

In this section we use Proposition 1 to study small quantum cohomology ring of $Y(\mathbb{E})$. As explained in the introduction, some of the relations in the small quantum cohomology ring come from differential operators.

Proposition 3. Whenever Proposition 1 holds, quantum differential operators of $X$ may be lifted in $Y(\mathbb{E})$, while the quantum differential operators of the fiber $Y$ may be extended to $Y(\mathbb{E})$. Both types of operators produce relations in the quantum cohomology $Q H_{s}^{*} Y(\mathbb{E})$.

Proof. Recall that $D_{i}=\sum a_{i j} p_{j}$. Let

$$
c_{1}\left(L_{i}\right)=\sum_{j=k+1}^{l} c_{i j} p_{j}, i=0,1, \ldots, n .
$$

Recall that the nef basis $\left\{p_{1}, p_{2}, \ldots, p_{k}, p_{k+1}, \ldots p_{l}\right\}$ of $H^{2}(Y(\mathbb{E}), \mathbb{Z})$ is obtained by completing a nef basis $\left\{p_{k+1}, \ldots, p_{l}\right\}$ of $X$. Let

$$
\mathcal{P}\left(\hbar, \hbar \partial / \partial t_{k+1}, \ldots, \hbar \partial / \partial t_{l}, q_{2}\right)=\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \mathcal{P}_{\alpha}
$$

be a polynomial differential operator with $\Lambda$ a finite subset of the Mori cone of $X$. Suppose that

$$
\begin{gathered}
0=\mathcal{P} J(X)=\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \sum_{\beta} \mathcal{P}_{\alpha}\left(\exp \left(\frac{p t}{\hbar}\right) q_{2}^{\beta}\right) J_{\beta}(X) \\
=\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \sum_{\beta} c_{\alpha, \beta} \exp \left(\frac{p t}{\hbar}\right) q_{2}^{\beta} J_{\beta}(X)=\exp \left(\frac{p t}{\hbar}\right) \sum_{\alpha \in \Lambda, \beta} q_{2}^{\alpha+\beta} c_{\alpha, \beta} J_{\beta}(X) .
\end{gathered}
$$

Let

$$
\delta_{\alpha}=\prod_{i=1}^{n} \prod_{r_{i}=0}^{-L_{i} \cdot \alpha-1}\left(\sum_{j=1}^{k} a_{i j} \hbar \frac{\partial}{\partial t_{j}}+\sum_{j=k+1}^{l} c_{i j} \hbar \frac{\partial}{\partial t_{j}}-r_{i} \hbar\right), \tilde{\mathcal{P}}=\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \delta_{\alpha} \mathcal{P}_{\alpha}
$$

with the convention that if

$$
L_{i}(\alpha)=0
$$

the factors of $\delta_{\alpha}$ corresponding to $L_{i}$ are missing. Notice that

$$
L_{n+1}(\alpha)=\ldots=L_{m}(\alpha)=0
$$

since we have chosen $L_{i}$ to be trivial for $i>n$. We compute

$$
\begin{gathered}
\tilde{\mathcal{P}} J(Y(\mathbb{E}))=\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \delta_{\alpha} \sum_{\nu, \beta} \mathcal{P}_{\alpha}\left(q_{2}^{\beta} \exp \left(\frac{p t}{\hbar}\right)\right) q_{1}^{\nu} \Omega_{\nu, \beta} J_{\beta}= \\
\sum_{\alpha \in \Lambda} q_{2}^{\alpha} \delta_{\alpha} \sum_{\nu, \beta} c_{\alpha, \beta} \exp \left(\frac{p t}{\hbar}\right) q_{1}^{\nu} q_{2}^{\beta} \Omega_{\nu, \beta} J_{\beta}
\end{gathered}
$$

One can easily show that

$$
\delta_{\alpha}\left(\exp \left(\frac{p t}{\hbar}\right) q_{1}^{\nu} q_{2}^{\beta} \Omega_{\nu, \beta}\right)=\exp \left(\frac{p t}{\hbar}\right) q_{1}^{\nu} q_{2}^{\beta} \Omega_{\nu, \alpha+\beta}
$$

It follows that

$$
\tilde{\mathcal{P}} J(Y(\mathbb{E}))=\exp \left(\frac{p t}{\hbar}\right) \sum_{\nu} q_{1}^{\nu} \sum_{\alpha \in \Lambda, \beta} c_{\alpha, \beta} q_{2}^{\alpha+\beta} \Omega_{\nu, \alpha+\beta} J_{\beta}(X)=0
$$

Hence the relation $\mathcal{P}\left(0, p_{k+1}, \ldots, p_{l}, q_{2}\right)=0$ in $Q H_{s}^{*} X$ lifts into the relation

$$
\mathcal{P}\left(0, p_{k+1}, \ldots, p_{l}, q_{2} \prod_{i=1}^{n} D_{i}\right)=0
$$

in $Q H_{s}^{*} Y(\mathbb{E})$, where

$$
\left(\prod_{i=1}^{n} D_{i}\right)^{\alpha}:=\prod_{i=1}^{n} D_{i}^{-L_{i}(\alpha)}, \forall \alpha \in M X .
$$

For a curve class $\nu$ in the fiber of $\pi$, consider the following differential operator

$$
\begin{gathered}
\Delta_{\nu}\left(\hbar \frac{\partial}{\partial t_{1}}, \ldots, \hbar \frac{\partial}{\partial t_{l}}, q_{j}\right):=\prod_{i: D_{i}(\nu)>0} \prod_{m=0}^{D_{i}(\nu)-1}\left(\sum_{j=1}^{k} a_{i j} \hbar \frac{\partial}{\partial t_{j}}-\sum_{j=k+1}^{l} c_{i j} \hbar \frac{\partial}{\partial t_{j}}+m \hbar\right) \\
-q^{\nu} \prod_{i: D_{i}(\nu)<0} \prod_{m=0}^{-D_{i}(\nu)-1}\left(\sum_{j=1}^{k} a_{i j} \hbar \frac{\partial}{\partial t_{j}}-\sum_{j=k+1}^{l} c_{i j} \hbar \frac{\partial}{\partial t_{j}}+m \hbar\right) .
\end{gathered}
$$

It is easy to show that it satisfies

$$
\Delta_{\nu} J(Y(\mathbb{E}))=0
$$

It follows that

$$
\Delta_{\nu}\left(p_{1}, \ldots, p_{l}, q_{j}\right)=0
$$

in $Q H_{s}^{*} Y(\mathbb{E})$, i.e.

$$
\prod_{i=1}^{r} D_{i}^{D_{i}(\nu)}=q^{\nu}
$$

These are precisely the extensions to $Y(\mathbb{E})$ of the small quantum cohomolgy relations of the fiber $Y$.

Sometimes all the relations in $Q H_{s}^{*} X$ come from quantum differential operators, hence $Q H_{s}^{*} X$ pulls back to $Q H_{s}^{*} Y(\mathbb{E})$. This is the case when $X$ is a Fano toric variety. The results of this section yield a complete description of $Q H^{*} Y(E)$ which generalizes previous results of Costa et al [4] and Qin et al [15] and Givental [9].

## 4. The General (Nontoric) Case

We believe that Proposition 1 holds for any $X$. On one end, the equality of the $d=0$ terms in $J(Y(\mathbb{E}))=I(Y(\mathbb{E}))$ is easy to establish. Indeed, the relative GromovWitten theory of the $Y$-bundle over $B \mathbb{T}$ associated with the universal bundle $E \mathbb{T} \mapsto$ $B \mathbb{T}$ is precisely the $\mathbb{T}$-equivariant GW theory of $Y$ (Astashkevich and Sadov [1]). The latter pulls back under the classifying map $X \mapsto B \mathbb{T}$ to the relative GW theory of $Y(\mathbb{E})$ over $X$. It follows that the restriction of $J(Y(\mathbb{E}))$ to $\nu=0$ is obtained by substituting $c_{1}\left(L_{i}\right)$ for $\lambda_{i}$ in $J^{\mathbb{T}}(Y)$. Since $Y$ is assumed to be Fano, the generator $J^{\mathbb{T}}(Y)$ is known (see for example [8]) and the substitution $c_{1}\left(L_{i}\right) \mapsto \lambda_{i}$ is easily seen to yield the desired equality. At the other end, the $\nu=0$ equality follows as an application of the equivariant quantum Lefshetz principle for the action of a torus on the fibers of $Y(\mathbb{E})$. The fixed point component relevant for the equivariant and localization considerations ([12]) consists of the maps that land in the section $s(X)$. The top chern class of the virtual normal bundle for this component is that of the $\mathbb{H}^{1}$-bundle for $\oplus_{i=1}^{m} L_{i}$. Calculations are easy to carry out (see for example Elezi [7]).

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