

## SCALAR CURVATURE RIGIDITY OF GEODESIC BALLS IN $S^n$

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### Abstract

In this paper, we prove a scalar curvature rigidity result for geodesic balls in  $S^n$ . This result contrasts sharply with the counterexamples to Min-Oo's conjecture constructed in [7].

### 1. Introduction

This paper is concerned with rigidity phenomena involving the scalar curvature. These questions are motivated to a large extent by the positive mass theorem in general relativity, which was proved by Schoen and Yau [18] and Witten [20]. An important corollary of this theorem is that any Riemannian metric on  $\mathbb{R}^n$  which has nonnegative scalar curvature and agrees with the Euclidean metric outside a compact set is necessarily flat.

It was observed by Miao [16] that the positive mass theorem implies the following rigidity result for metrics on the unit ball:

**Theorem 1.** *Suppose that  $g$  is a Riemannian metric on the unit ball  $B^n \subset \mathbb{R}^n$  with the following properties:*

- *The scalar curvature of  $g$  is nonnegative.*
- *The induced metric on the boundary  $\partial B^n$  agrees with the standard metric on  $\partial B^n$ .*
- *The mean curvature of  $\partial B^n$  with respect to  $g$  is at least  $n - 1$ .*

*Then  $g$  is isometric to the standard metric on  $B^n$ .*

An important generalization of Theorem 1 was proved by Shi and Tam [19].

Motivated by the positive mass theorem, Min-Oo [17] posed the following question:

**Min-Oo's Conjecture.** *Suppose that  $g$  is a Riemannian metric on the hemisphere  $S_+^n$  with the following properties:*

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- The scalar curvature of  $g$  is at least  $n(n-1)$ .
- The induced metric on the boundary  $\partial S_+^n$  agrees with the standard metric on  $\partial S_+^n$ .
- The boundary  $\partial S_+^n$  is totally geodesic with respect to  $g$ .

Then  $g$  is isometric to the standard metric on  $S_+^n$ .

Min-Oo's conjecture has been verified in many special cases (see e.g. [11], [13], [14]). A related rigidity result for real projective space  $\mathbb{RP}^3$  was established in [3] (see also [4]). A survey of these and other related results can be found in [6].

Very recently, counterexamples to Min-Oo's conjecture were constructed in [7].

**Theorem 2** (S. Brendle, F.C. Marques, A. Neves [7]). *Given any integer  $n \geq 3$ , there exists a smooth Riemannian metric  $\hat{g}$  on the hemisphere  $S_+^n$  with the following properties:*

- The scalar curvature of  $\hat{g}$  is at least  $n(n-1)$  at each point on  $S_+^n$ .
- The scalar curvature of  $\hat{g}$  is strictly greater than  $n(n-1)$  at some point on  $S_+^n$ .
- The metric  $\hat{g}$  agrees with the standard metric in a neighborhood of  $\partial S_+^n$ .

The proof of Theorem 2 relies on a perturbation analysis.

In this paper, we study the analogous rigidity question for geodesic balls in  $S^n$  of radius less than  $\frac{\pi}{2}$ . To fix notation, let  $\bar{g}$  be the standard metric on  $S^n$  and let  $f : S^n \rightarrow \mathbb{R}$  denotes the restriction of the coordinate function  $x_{n+1}$  to  $S^n$ . We will consider a domain of the form  $\Omega = \{f \geq c\}$ . If  $c \geq \frac{2}{\sqrt{n+3}}$ , we have the following rigidity result:

**Theorem 3.** *Consider the domain  $\Omega = \{f \geq c\}$ , where  $c \geq \frac{2}{\sqrt{n+3}}$ . Let  $g$  be a Riemannian metric on  $\Omega$  with the following properties:*

- $R_g \geq n(n-1)$  at each point in  $\Omega$ .
- The metrics  $g$  and  $\bar{g}$  induce the same metric on  $\partial\Omega$ .
- $H_g \geq H_{\bar{g}}$  at each point on  $\partial\Omega$ .

*If  $g - \bar{g}$  is sufficiently small in the  $C^2$ -norm, then  $\varphi^*(g) = \bar{g}$  for some diffeomorphism  $\varphi : \Omega \rightarrow \Omega$  with  $\varphi|_{\partial\Omega} = \text{id}$ .*

We remark that the conclusion of Theorem 3 holds under the weaker assumption that  $g$  is close to  $\bar{g}$  in  $W^{2,p}$ -norm for  $p > n$ .

Note that Theorem 3 is false for the hemisphere  $\{f \geq 0\}$ : by Theorem 4 in [7], there exist Riemannian metrics on the hemisphere which satisfy the assumptions of Theorem 3 and are arbitrary close to the standard metric  $\bar{g}$  in the  $C^\infty$ -topology, but which are not isometric to  $\bar{g}$ .

The proof of Theorem 3 relies on a perturbation analysis which is similar in spirit to Bartnik's work on the positive mass theorem (cf. [1], Section 5). Similar techniques have been employed in the study

of the total scalar curvature functional (see e.g. [2], Section 4G) and the Yamabe flow (cf. [5]). Dai, Wang, and Wei [8],[9] have obtained interesting stability results for manifolds with parallel spinors, as well as for Kähler-Einstein manifolds.

**2. The scalar curvature and boundary mean curvature of a perturbed metric**

In this section, we consider a smooth manifold  $\Omega$  with boundary  $\partial\Omega$ . Let  $\bar{g}$  be a fixed Riemannian metric on  $\Omega$ . Moreover, we consider another Riemannian metric  $g = \bar{g} + h$ , where  $|h|_{\bar{g}} \leq \frac{1}{2}$  at each point in  $\Omega$ . For abbreviation, we write  $(h^2)_{ik} = \bar{g}^{jl} h_{ij} h_{kl}$ .

**Proposition 4.** *The scalar curvature of  $g$  satisfies the pointwise estimate*

$$\begin{aligned} & \left| R_g - R_{\bar{g}} + \langle \text{Ric}_{\bar{g}}, h \rangle - \langle \text{Ric}_{\bar{g}}, h^2 \rangle \right. \\ & + \frac{1}{4} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_j h_{lq} - \frac{1}{2} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} \\ & \left. + \frac{1}{4} \bar{g}^{pq} \partial_p(\text{tr}_{\bar{g}}(h)) \partial_q(\text{tr}_{\bar{g}}(h)) + \bar{D}_i \left( g^{ik} g^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \right) \right| \\ & \leq C |h| |\bar{D}h|^2 + C |h|^3. \end{aligned}$$

Here,  $\bar{D}$  denotes the Levi-Civita connection with respect to  $\bar{g}$ , and  $C$  is a positive constant which depends only on  $n$ .

*Proof.* The Levi-Civita connection with respect to  $g$  is given by

$$D_X Y = \bar{D}_X Y + \Gamma(X, Y),$$

where  $\Gamma$  is defined by

$$g(\Gamma(X, Y), Z) = \frac{1}{2} ((\bar{D}_X h)(Y, Z) + (\bar{D}_Y h)(X, Z) - (\bar{D}_Z h)(X, Y)).$$

In local coordinates, the tensor  $\Gamma$  can be written in the form

$$\Gamma_{jk}^m = \frac{1}{2} g^{lm} (\bar{D}_j h_{kl} + \bar{D}_k h_{jl} - \bar{D}_l h_{jk}).$$

With this understood, the scalar curvature of  $g$  is given by

$$\begin{aligned} R_g &= g^{ik} (\text{Ric}_{\bar{g}})_{ik} + g^{ik} g^{jl} g_{pq} \Gamma_{il}^q \Gamma_{jk}^p - g^{ik} g^{jl} g_{pq} \Gamma_{jl}^q \Gamma_{ik}^p \\ &\quad - g^{ik} g^{jl} (\bar{D}_{i,k}^2 h_{jl} - \bar{D}_{i,l}^2 h_{jk}) \end{aligned}$$

(cf. [7], Proposition 16). This implies

$$\begin{aligned} & \left| R_g - R_{\bar{g}} + \langle \text{Ric}_{\bar{g}}, h \rangle - \langle \text{Ric}_{\bar{g}}, h^2 \rangle \right. \\ & - \frac{3}{4} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_j h_{lq} + \frac{1}{2} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} \\ & + \frac{1}{4} \bar{g}^{pq} \partial_p(\text{tr}_{\bar{g}}(h)) \partial_q(\text{tr}_{\bar{g}}(h)) - \bar{g}^{ij} \bar{g}^{pq} \bar{D}_i h_{jp} \partial_q(\text{tr}_{\bar{g}}(h)) \\ & \left. + \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{jp} \bar{D}_k h_{lq} + g^{ik} g^{jl} (\bar{D}_{i,k}^2 h_{jl} - \bar{D}_{i,l}^2 h_{jk}) \right| \\ & \leq C |h| |\bar{D}h|^2 + C |h|^3. \end{aligned}$$

Hence, we obtain

$$\begin{aligned} & \left| R_g - R_{\bar{g}} + \langle \text{Ric}_{\bar{g}}, h \rangle - \langle \text{Ric}_{\bar{g}}, h^2 \rangle \right. \\ & + \frac{1}{4} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_j h_{lq} - \frac{1}{2} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} \\ & \left. + \frac{1}{4} \bar{g}^{pq} \partial_p(\text{tr}_{\bar{g}}(h)) \partial_q(\text{tr}_{\bar{g}}(h)) + \bar{D}_i (g^{ik} g^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk})) \right| \\ & \leq C |h| |\bar{D}h|^2 + C |h|^3, \end{aligned}$$

as claimed.

q.e.d.

In the next step, we estimate the mean curvature of the boundary  $\partial\Omega$  with respect to the metric  $\bar{g}$ . To that end, we assume that  $g$  and  $\bar{g}$  induce the same metric on the boundary  $\partial\Omega$ ; in other words, we assume that  $h(X, Y) = 0$  whenever  $X$  and  $Y$  are tangent vectors to  $\partial\Omega$ .

**Proposition 5.** *Assume that  $g$  and  $\bar{g}$  induce the same metric on the boundary  $\partial\Omega$ . Then the mean curvature of  $\partial\Omega$  with respect to  $g$  satisfies*

$$\begin{aligned} & \left| 2(H_g - H_{\bar{g}}) - \left( h(\bar{\nu}, \bar{\nu}) - \frac{1}{4} h(\bar{\nu}, \bar{\nu})^2 + \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 \right) H_{\bar{g}} \right. \\ & \left. + \left( 1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu}) \right) \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) \right| \\ & \leq C |h|^2 |\bar{D}h| + C |h|^3. \end{aligned}$$

Here,  $\{e_a : 1 \leq a \leq n-1\}$  is a local orthonormal frame on  $\partial\Omega$ , and  $C$  is a positive constant that depends only on  $n$ .

*Proof.* Using the identity

$$H_g \nu - H_{\bar{g}} \bar{\nu} = - \sum_{a=1}^{n-1} (D_{e_a} e_a - \bar{D}_{e_a} e_a) = - \sum_{a=1}^{n-1} \Gamma(e_a, e_a),$$

we obtain

$$\begin{aligned} & 2(H_g g(\nu, \bar{\nu}) - H_{\bar{g}} g(\bar{\nu}, \bar{\nu})) \\ &= -2 \sum_{a=1}^{n-1} g(\Gamma(e_a, e_a), \bar{\nu}) = - \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)). \end{aligned}$$

Clearly,  $g(\bar{\nu}, \bar{\nu}) = 1 + h(\bar{\nu}, \bar{\nu})$ . Moreover, it is easy to see that the vector  $\bar{\nu} - \sum_{a=1}^{n-1} h(e_a, \bar{\nu}) e_a$  is orthogonal to  $\partial\Omega$  with respect to  $g$ . From this, we deduce that

$$\bar{\nu} - \sum_{a=1}^{n-1} h(e_a, \bar{\nu}) e_a = \left(1 + h(\bar{\nu}, \bar{\nu}) - \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2\right)^{\frac{1}{2}} \nu,$$

hence

$$g(\nu, \bar{\nu}) = \left(1 + h(\bar{\nu}, \bar{\nu}) - \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2\right)^{\frac{1}{2}}.$$

Substituting these identities into the previous formula for  $H_g$ , the assertion follows. q.e.d.

### 3. Perturbations of the standard metric on $S^n$

We now consider perturbations of the standard metric  $\bar{g}$  on  $S^n$ . To fix notation, let  $f : S^n \rightarrow \mathbb{R}$  denote the restriction of the coordinate function  $x_{n+1}$  to  $S^n$ , and let  $\Omega = \{f \geq c\}$  be a geodesic ball centered at the north pole. Here,  $c$  is a positive real number which will be specified later.

Let  $g$  be a Riemannian metric on  $\Omega$ . We will assume throughout that  $g$  and  $\bar{g}$  induce the same metric on the boundary  $\partial\Omega$ . Moreover, we assume that  $g = \bar{g} + h$ , where  $|h|_{\bar{g}} \leq \frac{1}{2}$  at each point in  $\Omega$ .

Our goal in this section is to estimate the integral

$$\int_{\Omega} (R_g - n(n-1)) f \, d\text{vol}_{\bar{g}}$$

(see also [12]).

**Proposition 6.** *We have*

$$\begin{aligned}
& \left| \int_{\Omega} (R_g - n(n-1) - (n-1)|h|_{\bar{g}}^2) f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} \right. \\
& + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\
& + \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& + \int_{\Omega} \bar{g}^{ip} \bar{g}^{kq} \bar{g}^{jl} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& + \int_{\partial\Omega} \bar{g}^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k f \, d\sigma_{\bar{g}} \\
& - \int_{\partial\Omega} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k f \, d\sigma_{\bar{g}} \\
& \left. - \int_{\partial\Omega} \bar{g}^{kq} \bar{g}^{jl} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^p f \, d\sigma_{\bar{g}} \right| \\
& \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}} + C \int_{\partial\Omega} |h|^2 |\bar{D}h| \, d\sigma_{\bar{g}},
\end{aligned}$$

where  $C$  is a positive constant that depends only on  $n$  and  $c$ .

*Proof.* Using Proposition 4 and the divergence theorem, we obtain

$$\begin{aligned}
& \left| \int_{\Omega} (R_g - n(n-1) + (n-1) \text{tr}_{\bar{g}}(h) - (n-1)|h|_{\bar{g}}^2) f \, d\text{vol}_{\bar{g}} \right. \\
& + \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\
& + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} - \int_{\Omega} g^{ik} g^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& + \int_{\partial\Omega} g^{ik} g^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{g}_{im} \bar{\nu}^m f \, d\sigma_{\bar{g}} \left. \right| \\
& \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}}.
\end{aligned}$$

Here,  $\bar{\nu}$  denotes the outward-pointing unit normal vector to  $\partial\Omega$  with respect to the metric  $\bar{g}$ . Using the identity  $\bar{D}_{i,k}^2 f = -f \bar{g}_{ik}$ , we obtain

$$\begin{aligned}
& \int_{\Omega} (n-1) \text{tr}_{\bar{g}}(h) f \, d\text{vol}_{\bar{g}} - \int_{\Omega} \bar{g}^{ik} \bar{g}^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& = - \int_{\partial\Omega} (\text{tr}_{\bar{g}}(h) \partial_{\bar{\nu}} f - h(\bar{\nu}, \bar{\nabla} f)) \, d\sigma_{\bar{g}} = 0.
\end{aligned}$$

Thus, we conclude that

$$\begin{aligned} & \left| \int_{\Omega} (R_g - n(n-1) - (n-1)|h|_{\bar{g}}^2) f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} \right. \\ & + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\ & - \int_{\Omega} (g^{ik} g^{jl} - \bar{g}^{ik} \bar{g}^{jl}) (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\ & + \left. \int_{\partial\Omega} g^{ik} g^{jl} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{g}_{im} \bar{\nu}^m f \, d\sigma_{\bar{g}} \right| \\ & \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}}. \end{aligned}$$

From this, the assertion follows easily. q.e.d.

In the remainder of this section, we will assume that  $h$  is divergence-free in the sense that  $\bar{g}^{ik} \bar{D}_i h_{kl} = 0$ .

**Proposition 7.** *Assume that  $h$  is divergence-free. Then*

$$\begin{aligned} & \left| \int_{\Omega} (R_g - n(n-1)) f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} \right. \\ & + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} \\ & + \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\partial\Omega} (|h|_{\bar{g}}^2 + 3h(\bar{\nu}, \bar{\nu})^2) \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} \\ & + \int_{\partial\Omega} \bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k f \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}} \\ & - \int_{\partial\Omega} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k f \, d\sigma_{\bar{g}} \\ & - \left. \int_{\partial\Omega} \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p f \, d\sigma_{\bar{g}} \right| \\ & \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}} + C \int_{\partial\Omega} |h|^2 |\bar{D}h| \, d\sigma_{\bar{g}}, \end{aligned}$$

where  $C$  is a positive constant that depends only on  $n$  and  $c$ .

*Proof.* Since  $\bar{g}$  has constant sectional curvature 1, we have

$$\bar{D}_{i,l}^2 h_{jq} = \bar{D}_{i,i}^2 h_{jq} + h_{lq} \bar{g}_{ij} - h_{iq} \bar{g}_{jl} + h_{jl} \bar{g}_{iq} - h_{ij} \bar{g}_{lq}.$$

Since  $h$  is divergence-free, it follows that

$$\bar{g}^{ij} \bar{D}_{i,l}^2 h_{jq} = n h_{lq} - \text{tr}_{\bar{g}}(h) \bar{g}_{lq}.$$

This implies

$$\begin{aligned}
& - \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \partial_i f \, d\text{vol}_{\bar{g}} - \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\
& = \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_{i,l}^2 h_{jq} f \, d\text{vol}_{\bar{g}} - \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}} \\
& = n \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} - \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} - \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}}.
\end{aligned}$$

From this, we deduce that

$$\begin{aligned}
& \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& - \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\
& = \frac{1}{2} \int_{\Omega} \bar{g}^{ik} \partial_k (|h|_{\bar{g}}^2) \partial_i f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} \bar{D}_l h_{jk} \partial_i f \, d\text{vol}_{\bar{g}} \\
& + \frac{n}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\
& - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}}.
\end{aligned}$$

Integration by parts gives

$$\begin{aligned}
& \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\
& - \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\
& = -\frac{1}{2} \int_{\Omega} |h|_{\bar{g}}^2 \Delta_{\bar{g}} f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} h_{jk} \bar{D}_{i,l}^2 f \, d\text{vol}_{\bar{g}} \\
& + \frac{1}{2} \int_{\partial\Omega} |h|_{\bar{g}}^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{ik} \bar{g}^{jp} h_{pq} h_{jk} \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}} \\
& + \frac{n}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\
& - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}} \\
& = \frac{2n-1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} - \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\
& + \frac{1}{2} \int_{\partial\Omega} |h|_{\bar{g}}^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{ik} \bar{g}^{jp} h_{pq} h_{jk} \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}} \\
& - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}}.
\end{aligned}$$



Moreover, we have

$$\begin{aligned} & \int_{\Omega} \bar{g}^{ip} \bar{g}^{kq} \bar{g}^{jl} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\ &= \int_{\Omega} \bar{g}^{ip} \bar{g}^{kq} h_{pq} \partial_k (\text{tr}_{\bar{g}}(h)) \partial_i f \, d\text{vol}_{\bar{g}} \\ &= - \int_{\Omega} \bar{g}^{ip} \bar{g}^{kq} h_{pq} \text{tr}_{\bar{g}}(h) \bar{D}_{i,k}^2 f \, d\text{vol}_{\bar{g}} + \int_{\partial\Omega} \bar{g}^{ip} h_{pq} \text{tr}_{\bar{g}}(h) \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}} \\ &= \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} + \int_{\partial\Omega} \bar{g}^{ip} h_{pq} \text{tr}_{\bar{g}}(h) \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}}. \end{aligned}$$

Putting these facts together, we obtain

$$\begin{aligned} & \int_{\Omega} \bar{g}^{ip} \bar{g}^{kq} \bar{g}^{jl} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\ &+ \int_{\Omega} \bar{g}^{ik} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \partial_i f \, d\text{vol}_{\bar{g}} \\ &- \frac{1}{2} \int_{\Omega} \bar{g}^{ij} \bar{g}^{kl} \bar{g}^{pq} \bar{D}_i h_{kp} \bar{D}_l h_{jq} f \, d\text{vol}_{\bar{g}} \\ &= \frac{2n-1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\ &+ \frac{1}{2} \int_{\partial\Omega} |h|_{\bar{g}}^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{ik} \bar{g}^{jp} h_{pq} h_{jk} \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}} \\ &+ \int_{\partial\Omega} \bar{g}^{ip} h_{pq} \text{tr}_{\bar{g}}(h) \partial_i f \bar{\nu}^q \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}} \\ &= \frac{2n-1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\ &+ \frac{1}{4} \int_{\partial\Omega} (|h|_{\bar{g}}^2 + 3h(\bar{\nu}, \bar{\nu})^2) \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j f \, d\sigma_{\bar{g}}. \end{aligned}$$

Hence, the assertion follows from Proposition 6. q.e.d.

#### 4. Analysis of the boundary terms

In this section, we analyze the boundary terms in Proposition 7. As in the previous section, we assume that  $\bar{g}$  is the standard metric on  $S^n$ , and  $\Omega = \{f \geq c\}$  centered at the north pole. Moreover, we consider a Riemannian metric on  $\Omega$  of the form  $g = \bar{g} + h$ , where  $|h|_{\bar{g}} \leq \frac{1}{2}$  at each point in  $\Omega$ .

**Proposition 8.** *Assume that  $h$  is divergence-free. Then*

$$\begin{aligned}
& \int_{\partial\Omega} \bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j d\sigma_{\bar{g}} \\
& - \int_{\partial\Omega} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k d\sigma_{\bar{g}} \\
& - \int_{\partial\Omega} \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p d\sigma_{\bar{g}} \\
& = - \int_{\partial\Omega} (1 - h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) d\sigma_{\bar{g}} \\
& + \int_{\partial\Omega} (1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})) h(\bar{\nu}, \bar{\nu}) H_{\bar{g}} d\sigma_{\bar{g}} \\
& + \frac{3n-2}{2(n-1)} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} d\sigma_{\bar{g}}.
\end{aligned}$$

Here,  $\{e_a : 1 \leq a \leq n-1\}$  is a local orthonormal frame on  $\partial\Omega$ , and  $C$  is a positive constant that depends only on  $n$  and  $c$ .

*Proof.* Let  $\{e_a : 1 \leq a \leq n-1\}$  be a local orthonormal frame on  $\partial\Omega$ . Since  $h$  is divergence-free, we have

$$\begin{aligned}
& \bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k - \frac{1}{2} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j \\
& - \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k - \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p \\
& = -(1 - h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) \\
& + (1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (\bar{D}_{e_a} h)(e_a, \bar{\nu}) - \frac{1}{2} \sum_{a=1}^{n-1} (\bar{D}_{e_a} h)(\bar{\nu}, \bar{\nu}) h(e_a, \bar{\nu}) \\
& + \frac{3}{2} \sum_{a,b=1}^{n-1} h(e_a, \bar{\nu}) (\bar{D}_{e_b} h)(e_a, e_b) - \sum_{a,b=1}^{n-1} h(e_a, \bar{\nu}) (\bar{D}_{e_a} h)(e_b, e_b).
\end{aligned}$$

At this point, we define a one-form  $\omega$  on  $\partial\Omega$  by  $\omega(e_a) = (1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})) h(e_a, \bar{\nu})$ . Since  $\partial\Omega$  is umbilic with respect to  $\bar{g}$ , we have

$$\bar{D}_{e_a} \bar{\nu} = \frac{1}{n-1} H_{\bar{g}} e_a,$$

where  $H_{\bar{g}}$  denotes the mean curvature of  $\partial\Omega$  with respect to the metric  $\bar{g}$ . Using this relation, we obtain the following formula for the divergence

of  $\omega$ :

$$\begin{aligned} \operatorname{div}_{\partial\Omega}(\omega) &= \left(1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})\right) \sum_{a=1}^{n-1} (\bar{D}_{e_a} h)(e_a, \bar{\nu}) - \frac{1}{2} \sum_{a=1}^{n-1} (\bar{D}_{e_a} h)(\bar{\nu}, \bar{\nu}) h(e_a, \bar{\nu}) \\ &\quad - \left(1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})\right) h(\bar{\nu}, \bar{\nu}) H_{\bar{g}} - \frac{1}{n-1} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}}. \end{aligned}$$

Moreover, we have the pointwise identities

$$\sum_{b=1}^{n-1} (\bar{D}_{e_b} h)(e_a, e_b) = \frac{n}{n-1} h(e_a, \bar{\nu}) H_{\bar{g}}$$

and

$$\sum_{b=1}^{n-1} (\bar{D}_{e_a} h)(e_b, e_b) = \frac{2}{n-1} h(e_a, \bar{\nu}) H_{\bar{g}}.$$

Putting these facts together, we obtain

$$\begin{aligned} &\bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k - \frac{1}{2} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j \\ &\quad - \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k - \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p \\ &= -(1 - h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (2 (\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) \\ &\quad + \left(1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu})\right) h(\bar{\nu}, \bar{\nu}) H_{\bar{g}} + \frac{3n-2}{2(n-1)} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} + \operatorname{div}_{\partial\Omega}(\omega). \end{aligned}$$

Therefore, the assertion follows from the divergence theorem. q.e.d.

Combining Proposition 8 and Proposition 5, we can draw the following conclusion:

**Corollary 9.** *If  $h$  is divergence-free, then we have*

$$\begin{aligned} &\left| \int_{\partial\Omega} (2 - h(\bar{\nu}, \bar{\nu})) (H_g - H_{\bar{g}}) d\sigma_{\bar{g}} \right. \\ &\quad - \int_{\partial\Omega} \bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k d\sigma_{\bar{g}} + \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j d\sigma_{\bar{g}} \\ &\quad + \int_{\partial\Omega} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k d\sigma_{\bar{g}} \\ &\quad + \int_{\partial\Omega} \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p d\sigma_{\bar{g}} \\ &\quad \left. + \frac{1}{4} \int_{\partial\Omega} h(\bar{\nu}, \bar{\nu})^2 H_{\bar{g}} d\sigma_{\bar{g}} + \frac{n}{2(n-1)} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} d\sigma_{\bar{g}} \right| \\ &\leq C \int_{\partial\Omega} |h|^2 |\bar{D}h| d\sigma_{\bar{g}} + C \int_{\partial\Omega} |h|^3 d\sigma_{\bar{g}}, \end{aligned}$$

where  $C$  is a positive constant that depends only on  $n$  and  $c$ .

*Proof.* It follows from Proposition 5 that

$$\begin{aligned} & \left| \int_{\partial\Omega} (2 - h(\bar{\nu}, \bar{\nu})) (H_g - H_{\bar{g}}) d\sigma_{\bar{g}} \right. \\ & - \int_{\partial\Omega} \left( h(\bar{\nu}, \bar{\nu}) - \frac{3}{4} h(\bar{\nu}, \bar{\nu})^2 + \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 \right) H_{\bar{g}} d\sigma_{\bar{g}} \\ & + \int_{\partial\Omega} (1 - h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) d\sigma_{\bar{g}} \left. \right| \\ & \leq C \int_{\partial\Omega} |h|^2 |\bar{D}h| d\sigma_{\bar{g}} + C \int_{\partial\Omega} |h|^3 d\sigma_{\bar{g}}. \end{aligned}$$

Moreover, we have

$$\begin{aligned} & \int_{\partial\Omega} \bar{g}^{jl} \bar{D}_k h_{jl} \bar{\nu}^k d\sigma_{\bar{g}} - \frac{1}{2} \int_{\partial\Omega} \bar{g}^{kl} \bar{g}^{pq} h_{kp} \bar{D}_l h_{jq} \bar{\nu}^j d\sigma_{\bar{g}} \\ & - \int_{\partial\Omega} \bar{g}^{jp} \bar{g}^{lq} h_{pq} (\bar{D}_k h_{jl} - \bar{D}_l h_{jk}) \bar{\nu}^k d\sigma_{\bar{g}} \\ & - \int_{\partial\Omega} \bar{g}^{kq} \bar{g}^{jl} h_{pq} \bar{D}_k h_{jl} \bar{\nu}^p d\sigma_{\bar{g}} \\ & = - \int_{\partial\Omega} (1 - h(\bar{\nu}, \bar{\nu})) \sum_{a=1}^{n-1} (2(\bar{D}_{e_a} h)(e_a, \bar{\nu}) - (\bar{D}_{\bar{\nu}} h)(e_a, e_a)) d\sigma_{\bar{g}} \\ & + \int_{\partial\Omega} \left( 1 - \frac{1}{2} h(\bar{\nu}, \bar{\nu}) \right) h(\bar{\nu}, \bar{\nu}) H_{\bar{g}} d\sigma_{\bar{g}} \\ & + \frac{3n-2}{2(n-1)} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} d\sigma_{\bar{g}} \end{aligned}$$

by Proposition 8. Putting these facts together, the assertion follows.

q.e.d.

**Theorem 10.** *Assume that  $h$  is divergence-free. Then*

$$\begin{aligned}
 & \left| \int_{\Omega} (R_g - n(n-1)) f \, d\text{vol}_{\bar{g}} + \int_{\partial\Omega} (2 - h(\bar{\nu}, \bar{\nu})) (H_g - H_{\bar{g}}) f \, d\sigma \right. \\
 & + \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} \\
 & + \frac{1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\
 & + \int_{\partial\Omega} h(\bar{\nu}, \bar{\nu})^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} + \frac{1}{2} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} \\
 & + \frac{1}{4} \int_{\partial\Omega} h(\bar{\nu}, \bar{\nu})^2 H_{\bar{g}} f \, d\sigma_{\bar{g}} + \frac{n}{2(n-1)} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} f \, d\sigma_{\bar{g}} \left. \right| \\
 & \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}} \\
 & + C \int_{\partial\Omega} |h|^2 |\bar{D}h| \, d\sigma_{\bar{g}} + C \int_{\partial\Omega} |h|^3 \, d\sigma_{\bar{g}}.
 \end{aligned}$$

Here,  $C$  is a positive constant that depends only on  $n$  and  $c$ .

*Proof.* Recall that  $f$  is constant along the boundary  $\partial\Omega$ . Hence, the assertion is a consequence of Proposition 7 and Corollary 9. q.e.d.

### 5. Proof of Theorem 3

To prove Theorem 3, we need an analogue of Ebin’s slice theorem for manifolds with boundary [10] (see also [12]). The proof is standard, and works on any compact manifold with boundary.

**Proposition 11.** *Fix a real number  $p > n$ . If  $\|g - \bar{g}\|_{W^{2,p}(\Omega, \bar{g})}$  is sufficiently small, we can find a diffeomorphism  $\varphi : \Omega \rightarrow \Omega$  such that  $\varphi|_{\partial\Omega} = \text{id}$  and  $h = \varphi^*(g) - \bar{g}$  is divergence-free. Moreover,*

$$\|h\|_{W^{2,p}(\Omega, \bar{g})} \leq N \|g - \bar{g}\|_{W^{2,p}(\Omega, \bar{g})},$$

where  $N$  is a positive constant that depends only on  $\Omega$ .

*Proof.* Let  $\mathcal{S}$  denote the space of symmetric two-tensors on  $\Omega$  of class  $W^{2,p}$ , and let  $\mathcal{M}$  denote the space of Riemannian metrics on  $\Omega$  of class  $W^{2,p}$ . Moreover, let  $\mathcal{X}$  denote the space of vector fields of class  $W^{3,p}$  that vanish along the boundary  $\partial\Omega$ , and let  $\mathcal{D}$  denote the space of all diffeomorphisms  $\varphi : \Omega \rightarrow \Omega$  of class  $W^{3,p}$  satisfying  $\varphi|_{\partial\Omega} = \text{id}$ . Clearly, the tangent space to  $\mathcal{M}$  at  $\bar{g}$  can be identified with  $\mathcal{S}$ ; similarly, the tangent space to  $\mathcal{D}$  at the identity can be identified with  $\mathcal{X}$ .

There is a natural action

$$A : \mathcal{D} \times \mathcal{M} \rightarrow \mathcal{M}, \quad (\varphi, g) \rightarrow \varphi^*(g).$$

Let us consider the linearization of  $A$  around the point  $(\text{id}, \bar{g})$ . This gives a map  $L : T_{\text{id}}\mathcal{D} \rightarrow T_{\bar{g}}\mathcal{M}$ . The map  $L$  sends a vector field  $\xi \in \mathcal{X}$  to the Lie derivative  $\mathcal{L}_\xi(\bar{g}) \in \mathcal{S}$ . Standard elliptic regularity theory implies that

$$\mathcal{S} = \{\mathcal{L}_\xi(\bar{g}) : \xi \in \mathcal{X}\} \oplus \{h \in \mathcal{S} : h \text{ is divergence-free}\}$$

(compare [12], p. 523). Hence, the assertion follows from the implicit function theorem. q.e.d.

We now complete the proof of Theorem 3. Let  $g$  be a Riemannian metric on the domain  $\Omega = \{f \geq c\}$  with the following properties:

- $R_g \geq n(n - 1)$  at each point in  $\Omega$ .
- The metrics  $g$  and  $\bar{g}$  induce the same metric on  $\partial\Omega$ .
- $H_g \geq H_{\bar{g}}$  at each point on  $\partial\Omega$ .

If  $\|g - \bar{g}\|_{W^{2,p}(\Omega, \bar{g})}$  is sufficiently small, Proposition 11 implies the existence of a diffeomorphism  $\varphi : \Omega \rightarrow \Omega$  such that  $\varphi|_{\partial\Omega} = \text{id}$  and  $h = \varphi^*(g) - \bar{g}$  is divergence-free.

Note that  $R_{\varphi^*(g)} \geq n(n - 1)$  at each point in  $\Omega$  and  $H_{\varphi^*(g)} \geq H_{\bar{g}}$  at each point on  $\partial\Omega$ . Applying Theorem 10 to the metric  $\varphi^*(g) = \bar{g} + h$ , we obtain

$$\begin{aligned} & \frac{1}{4} \int_{\Omega} |\bar{D}h|^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{4} \int_{\Omega} |\bar{\nabla}(\text{tr}_{\bar{g}}(h))|^2 f \, d\text{vol}_{\bar{g}} \\ & + \frac{1}{2} \int_{\Omega} |h|_{\bar{g}}^2 f \, d\text{vol}_{\bar{g}} + \frac{1}{2} \int_{\Omega} \text{tr}_{\bar{g}}(h)^2 f \, d\text{vol}_{\bar{g}} \\ & + \int_{\partial\Omega} h(\bar{\nu}, \bar{\nu})^2 \partial_{\bar{\nu}} f \, d\sigma_{\bar{g}} + \frac{1}{2} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 \partial_{\nu} f \, d\sigma_{\bar{g}} \\ & + \frac{1}{4} \int_{\partial\Omega} h(\bar{\nu}, \bar{\nu})^2 H_{\bar{g}} f \, d\sigma_{\bar{g}} + \frac{n}{2(n-1)} \int_{\partial\Omega} \sum_{a=1}^{n-1} h(e_a, \bar{\nu})^2 H_{\bar{g}} f \, d\sigma_{\bar{g}} \\ & \leq C \int_{\Omega} |h| |\bar{D}h|^2 \, d\text{vol}_{\bar{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\bar{g}} \\ & + C \int_{\partial\Omega} |h|^2 |\bar{D}h| \, d\sigma_{\bar{g}} + C \int_{\partial\Omega} |h|^3 \, d\sigma_{\bar{g}}. \end{aligned}$$

If we choose  $c \geq \frac{2}{\sqrt{n+3}}$ , then

$$\frac{1}{4} H_{\bar{g}} f + \partial_{\bar{\nu}} f = \frac{n-1}{4} \frac{f^2}{|\bar{\nabla} f|} - |\bar{\nabla} f| = \frac{n-1}{4} \frac{c^2}{\sqrt{1-c^2}} - \sqrt{1-c^2} \geq 0$$

at each point on  $\partial\Omega$ . This implies

$$\begin{aligned} & \frac{1}{4} \int_{\Omega} |\overline{D}h|^2 f \, d\text{vol}_{\overline{g}} + \frac{1}{4} \int_{\Omega} |\overline{\nabla}(\text{tr}_{\overline{g}}(h))|^2 f \, d\text{vol}_{\overline{g}} \\ & + \frac{1}{2} \int_{\Omega} |h|_{\overline{g}}^2 f \, d\text{vol}_{\overline{g}} + \frac{1}{2} \int_{\Omega} \text{tr}_{\overline{g}}(h)^2 f \, d\text{vol}_{\overline{g}} \\ & \leq C \int_{\Omega} |h| |\overline{D}h|^2 \, d\text{vol}_{\overline{g}} + C \int_{\Omega} |h|^3 \, d\text{vol}_{\overline{g}} \\ & + C \int_{\partial\Omega} |h|^2 |\overline{D}h| \, d\sigma_{\overline{g}} + C \int_{\partial\Omega} |h|^3 \, d\sigma_{\overline{g}}. \end{aligned}$$

By the trace theorem, the error terms on the right hand side are bounded from above by  $C \|h\|_{C^1(\Omega, \overline{g})} \|h\|_{W^{1,2}(\Omega, \overline{g})}^2$ . Hence, if  $\|h\|_{C^1(\Omega, \overline{g})}$  is sufficiently small, then  $h$  vanishes identically, and therefore  $\varphi^*(g) = \overline{g}$ . This completes the proof of Theorem 3.

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