# POON'S SELF-DUAL METRICS AND KÄHLER GEOMETRY

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### Abstract

It is shown that the self-dual conformal metrics on connected sums of  $\mathbf{CP}_2$ 's recently produced by Y. S. Poon arise from zero scalar curvature Kähler metrics on blow-ups of  $\mathbf{C}^2$  by adding a point at infinity and reversing the orientation.

As noted by many authors ([4], [5], [6]), a complex surface with Kähler metric has anti-self-dual Weyl curvature iff the scalar curvature vanishes. On what would initially appear to be a completely unrelated front, Poon ([8], [9]) has produced positive scalar curvature self-dual metrics on connected sums of two and three complex projective planes. In fact, however, these phenomena are closely related:

**Theorem.** Let  $M = m\mathbf{CP}_2$ ,  $0 \le m \le 3$ , be equipped with a self-dual metric g of positive scalar curvature. There exists at least one point  $p \in M$  such that  $(M - \{p\}, g)$  is conformally isometric to  $\mathbf{C}^2$  with m points blown up equipped with an asymptotically flat Kähler metric of zero scalar curvature.

(**Remark.** The conformal isometry, of course, reverses orientation.)

**Proof.** Let  $\pi: Z \to M$  be the canonical projection from the twistor space Z onto M; recall [1] that Z consists of all orthogonal almost-complex structure tensors on M inducing the reverse orientation. There exists ([8], [9]) a complex surface  $\Sigma \subset Z$  isomorphic to  $\mathbb{CP}_2$  blown up at m points such that  $\pi|_{\Sigma}: \Sigma \to M$  is a diffeomorphism away from a projective line  $L \subset \Sigma$  sent to a point  $p \in M$ ; e.g. when m = 0,  $M = S^4$ ,  $Z = \mathbb{CP}_3$ , and  $\Sigma$  is a hyperplane. By construction,  $(\pi|_{\Sigma})^*g$  is a Hermitian metric on  $\Sigma - L$  but degenerates at L. Identifying  $\Sigma - L$  with  $\mathbb{C}^2$  blow up at m points, let

$$\hat{g} = (1 + r^2)^2 (\pi |_{\Sigma})^* g,$$

where r is the Euclidean distance from the origin in  $\mathbb{C}^2$ ; this is not only Hermitian, but asymptotically flat, differing from the standard metric only by terms of order  $1/r^2$  because the projection  $\Sigma \to M$  is standard on the

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second neighborhood of  $L \subset \Sigma$ . Notice that the Weyl curvature of  $\hat{g}$  is antiself-dual with respect to the complex orientation.

(The second infinitesimal neighborhood of any twistor line is isomorphic to the second neighborhood of the zero section in  $\mathscr{O}(1) \oplus \mathscr{O}(1) \to \mathbf{CP}_1$ ; this amounts to the fact that any conformal metric is flat to first order, but may be seen more directly from the obstruction theory of [7] via the vanishing of  $H^1(\mathbf{CP}_1, T \otimes N^*)$  and  $H^1(\mathbf{CP}_1, \hat{T} \otimes \bigcirc^2 N^*)$ , where  $T \cong \mathscr{O}(2)$  is the tangent bundle,  $N \cong \mathscr{O}(1) \oplus \mathscr{O}(1)$  is the normal bundle, and  $\hat{T} \cong T \oplus N$  is the extended tangent bundle. Since  $\Sigma$  is a blow-up of  $\mathbf{CP}_2$ , it contains a complex 2-parameter family of projective lines, one of which is L. This implies that the second neighborhood of  $L \subset \Sigma$  corresponds to one of the  $\mathscr{O}(1)$  factors of  $\mathscr{O}(1) \oplus \mathscr{O}(1) \to \mathbf{CP}_1$ .)

We now proceed by modifying an argument given for compact surfaces by Boyer [2, p. 522]. Let  $\omega$  be the (1, 1) form associated with  $\hat{g}$ , and let  $\beta$  be the 1-form defined by

$$-d\omega = \beta \wedge \omega.$$

Because  $\hat{g}$  is anti-self-dual,  $d\beta$  is an anti-self-dual 2-form. Letting  $B_r \subset (\Sigma - L)$  be the blow up of the ball of radius r, and letting dV denote the metric volume form of  $(\Sigma - L, \hat{g})$ , we have

$$-\int_{B_r} \|d\beta\|^2 \, dV = -\int_{B_r} d\beta \wedge {}^*d\beta = \int_{B_r} d\beta \wedge d\beta$$
$$= \int_{B_r} d(\beta \wedge d\beta) = \int_{\partial B_r} \beta \wedge d\beta.$$

But  $\beta$  and  $d\beta$  are of order  $1/r^3$  and  $1/r^4$ , respectively, so

$$-\int_{\Sigma-L} \|d\beta\|^2 \, dV = \lim_{r \to \infty} \int_{\partial B_r} \beta \wedge d\beta = 0.$$

Hence  $d\beta = 0$ . But  $\Sigma - L$  is simply connected, so  $\beta = df$ . Hence  $d(e^f \omega) = 0$ , so  $h = e^f \hat{g}$  is Kähler and anti-self-dual, and thus has scalar curvature zero. Since f differs from a constant by terms of order  $1/r^2$ , h is asymptotically flat. q.e.d.

These Kähler metrics may be written down explicitly for m = 0, 1. For m = 0, h is the standard flat metric on  $\mathbb{C}^2$ . For m = 1, h is the metric with Kähler potential on  $\mathbb{C}^2 - \{0\}$  given by  $\|\overrightarrow{z}\|^2 + \log \|\overrightarrow{z}\|^2$ ; this metric, first pointed out in this context by Burns [3], is the restriction of the standard product metric on  $\mathbb{CP}_1 \times \mathbb{C}^2$  to the blow-up

$$\tilde{\mathbf{C}}^2 = \{ [\overrightarrow{w}], \overrightarrow{z}) \in \mathbf{CP}_1 \times \mathbf{C}^2 \mid \overrightarrow{w} \wedge \overrightarrow{z} = 0 \}.$$

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Its remarkable conformal isometry with the Fubini-Study metric is obtained by extending the map  $\Phi: \mathbf{C}^2 - \{0\} \to \mathbf{C}^2 - \{0\}: \overrightarrow{z} \to \overrightarrow{z}/||\overrightarrow{z}||^2$  to a diffeomorphism between  $\widetilde{\mathbf{C}}^2$  and  $\mathbf{CP}_2 - \{\text{point}\}$ .

The explicit form of the potentials for m = 2, 3 remains a topic for further investigation. One may hope to produce an ansatz for the necessary potentials for arbitrary values of m, thereby producing self-dual metrics on all connected sums of  $\mathbb{CP}_2$ 's; but while such an ansatz may give the general solution for  $m \leq 3$ , one should only expect to produce special solutions in this manner for larger values of m. The reason is that the Kähler form gives rise to a solution of the twistor equation  $\nabla_{A'}{}^{(A}\omega^{BC)} = 0$  and thus to an element of  $H^0(Z, K^{-1/2})$ . The Riemann-Roch formula predicts that sections of  $K^{-1/2}$ must exist for small values of m, but gives no information when m is large; similarly our complex surface  $\Sigma$  need not exist. Nonetheless, there seems to be much to be learned here.

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