

Formula for the Casimir Operator in Iwasawa Coördinates

Floyd L. WILLIAMS

University of Massachusetts at Amherst and Sophia University

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Introduction

Let G be a connected non-compact semisimple Lie group with an Iwasawa decomposition $G=KAN$, where the Lie algebra of K is \mathfrak{k} in the Cartan decomposition $\mathfrak{g}=\mathfrak{k}+\mathfrak{p}$ of the Lie algebra \mathfrak{g} of G , the Lie algebra \mathfrak{a} of A is a maximal abelian subspace of \mathfrak{p} , and the Lie algebra \mathfrak{n} of N is a sum of root spaces corresponding to a choice of positive *restricted* roots. Let Ω be the Casimir operator of G . We give a formula for Ω in terms of first order differential operators $\delta_z, \delta_H, \delta_x$ defined for left-invariant vector fields $(z, H, x) \in \mathfrak{k} \times \mathfrak{a} \times \mathfrak{n}$. Theorem 1.10 is the main result. Such a formula was first proposed in [1]. The arguments given there, however, are incomplete. Corrections are made in the present paper. In particular our formula reduces to Takahashi's formula [3] when $G=SL(2, R)$.

§1. Statement of the result.

Let \mathfrak{g} be a non-compact real semisimple Lie algebra with a Cartan decomposition $\mathfrak{g}=\mathfrak{k}+\mathfrak{p}$, and Cartan involution θ . The Killing form B is negative definite on \mathfrak{k} and positive definite on \mathfrak{p} . The formula

$$(1.1) \quad \langle x, y \rangle = -B(x, \theta y) \quad x, y \in \mathfrak{g}$$

defines a real positive definite inner product \langle, \rangle on \mathfrak{g} . Let $\mathfrak{a} \subset \mathfrak{p}$ be a maximal abelian subspace of \mathfrak{p} . For $\alpha \in \mathfrak{a}^*$ (the real dual space of \mathfrak{a}) let $\mathfrak{g}_\alpha = \{x \in \mathfrak{g} \mid [H, x] = \alpha(H)x \text{ for every } H \text{ in } \mathfrak{a}\}$. α is a *restricted* root (relative to \mathfrak{a}) if both α and the root space \mathfrak{g}_α are non-zero. Let $\Sigma \subset \mathfrak{a}^*$ denote the set of restricted roots, let $\Sigma^+ \subset \Sigma$ denote a choice of a system of positive restricted roots, let \mathfrak{n} denote the sum of the positive root spaces, and let (for $\alpha \in \Sigma$)

$$(1.2) \quad m_\alpha = \dim \mathfrak{g}_\alpha \quad \rho = \frac{1}{2} \sum_{\alpha \in \Sigma^+} m_\alpha \alpha .$$

Then \mathfrak{g} has an Iwasawa decomposition $\mathfrak{g} = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$ and a root space decomposition $\mathfrak{g} = \mathfrak{m} + \mathfrak{a} + \mathfrak{n} + \theta\mathfrak{n}$, where \mathfrak{m} is the *centralizer* of \mathfrak{a} in \mathfrak{k} . Let G be a connected Lie group with Lie algebra \mathfrak{g} . We regard elements of \mathfrak{g} as left-invariant vector fields on G . Let $K \subset G$ be the Lie subgroup with Lie algebra \mathfrak{k} . Then G has an Iwasawa decomposition $G = KAN$ where $A = \exp \mathfrak{a}$, $N = \exp \mathfrak{n}$. The subgroups K, A, N are closed; K is compact when and only when the center of G is finite. Given $(z, H, x) \in \mathfrak{k} \times \mathfrak{a} \times \mathfrak{n}$ we define differential operators $\delta_z, \delta_H, \delta_x: C^\infty(G) \rightarrow C^\infty(G)$ on G by

$$(1.3) \quad \begin{aligned} (\delta_z f)(kan) &= \frac{d}{dt} f(k(\exp tz)an) \Big|_{t=0} \\ (\delta_H f)(kan) &= \frac{d}{dt} f(k(\exp tH)an) \Big|_{t=0} \\ (\delta_x f)(kan) &= \frac{d}{dt} f(ka(\exp tx)n) \Big|_{t=0} \end{aligned}$$

where $f \in C^\infty(G)$ and $kan \in KAN$ is the Iwasawa decomposition of an element of G . For $y \in \mathfrak{g}$ arbitrary (a left-invariant operator) we also define \tilde{y} by

$$(1.4) \quad (\tilde{y}f)(b) = \frac{d}{dt} f((\exp ty)b) \Big|_{t=0}$$

for $f \in C^\infty(G)$, $b \in G$. For $a = \exp H \in A$, $H \in \mathfrak{a}$, we write $H = \log a$ and we set

$$(1.5) \quad e^\alpha(kan) = e^{\alpha(\log a)}$$

for $\alpha \in \Sigma$, $kan \in KAN$. Thus $e^\alpha \in C^\infty(G)$. Let Ω be the Casimir operator of G . If $\{E_i\}, \{F_j\}$ are orthonormal bases of $\mathfrak{k}, \mathfrak{p}$, relative to \langle, \rangle given in (1.1), then $\Omega = -\sum_i E_i^2 + \sum_j F_j^2 \in U\mathfrak{g}$, the universal enveloping algebra of \mathfrak{g} . Let $\tilde{\Omega} = -\sum_i \tilde{E}_i^2 + \sum_j \tilde{F}_j^2$; see (1.4). Then one knows that

$$(1.6) \quad \Omega = \tilde{\Omega} \quad \text{on the domain } C^\infty(G) .$$

(1.6) can be proved using the Cartan decomposition $G = K \exp(\mathfrak{p})$ of G . Now choose orthonormal bases $\{H_i\}_{i=1}^r$ of \mathfrak{a} , $\{u_i\}_{i=1}^s$ of \mathfrak{m} , and $\{x_i\}_{i=1}^t$ of \mathfrak{n} . Set

$$(1.7) \quad z_i = \frac{x_i + \theta x_i}{\sqrt{2}} \in \mathfrak{k}, \quad y_i = \frac{x_i - \theta x_i}{\sqrt{2}} \in \mathfrak{p}$$

for $1 \leq i \leq t$. Let \mathfrak{m}^\perp denote the orthogonal complement of \mathfrak{m} in \mathfrak{k} .

PROPOSITION 1.8. $\{z_i\}_{i=1}^t$ is an orthonormal basis of \mathfrak{m}^\perp . Hence $\{u_i\}_{i=1}^s \cup \{z_i\}_{i=1}^t$ is an orthonormal basis of \mathfrak{k} . Also $\{H_i\}_{i=1}^r \cup \{y_i\}_{i=1}^t$ is an orthonormal basis of \mathfrak{p} . Hence $\Omega = -\sum_{i=1}^s \tilde{u}_i^2 - \sum_{i=1}^t \tilde{z}_i^2 + \sum_{i=1}^r \tilde{H}_i^2 + \sum_{i=1}^t \tilde{y}_i^2$ by (1.6). We make the choice of the orthonormal basis $\{x_i\}_{i=1}^t$ of \mathfrak{n} as follows. Write $\Sigma^+ = \{\alpha_1, \dots, \alpha_q\}$, and for $1 \leq j \leq q$ let $\{x_{i(j)}\}_{i=1}^{m_{\alpha_j}}$ be an orthonormal basis of \mathfrak{g}_{α_j} . Then

$$(1.9) \quad \{x_i\}_{i=1}^t = \bigcup_{1 \leq j \leq q} \{x_{i(j)}\}_{i=1}^{m_{\alpha_j}}$$

is an orthonormal basis of \mathfrak{n} ; $\sum_{j=1}^q m_{\alpha_j} = t$ and $2\rho = \sum_{j=1}^q m_{\alpha_j} \alpha_j$; see (1.2).

Using the expression for Ω in Proposition 1.8 we will derive the following main

THEOREM 1.10. As above let $r = \dim \mathfrak{a}$ (= the rank of G), $s = \dim \mathfrak{m}$, $t = \dim \mathfrak{n}$, and $q = |\Sigma^+|$, the cardinality of Σ^+ . Relative to the inner product defined in (1.1) choose orthonormal bases $\{H_i\}_{i=1}^r$ of \mathfrak{a} , $\{u_i\}_{i=1}^s$ of \mathfrak{m} , and $\{x_i\}_{i=1}^t = \bigcup_{1 \leq j \leq q} \{x_{i(j)}\}_{i=1}^{m_{\alpha_j}}$ of \mathfrak{n} (as in (1.9)). Let $z_i = (x_i + \theta x_i)/\sqrt{2}$ as in (1.7). Then

$$(1.11) \quad \Omega = \sum_{i=1}^r (\delta_{H_i}^2 + 2\rho(H_i)\delta_{H_i}) + \sum_{j=1}^q \sum_{i=1}^{m_{\alpha_j}} (2e^{-2\alpha_j} \delta_{x_{i(j)}}^2 - 2\sqrt{2} e^{-\alpha_j} \delta_{x_{i(j)}} \delta_{x_{i(j)}}) \\ + \sum_{i=1}^t \delta_{z_i}^2 - \sum_{i=1}^t \tilde{z}_i^2 - \sum_{i=1}^s \tilde{u}_i^2 \quad \text{on } C^\infty(G);$$

see (1.2), (1.3), (1.4), and (1.5). Moreover we have

$$\left(\sum_{i=1}^t \delta_{z_i}^2 - \sum_{i=1}^t \tilde{z}_i^2 - \sum_{i=1}^s \tilde{u}_i^2 \right) f = \left(- \sum_{i=1}^s u_i^2 \right) f \quad \text{on } KA \text{ for } f \in C^\infty(G).$$

Because of the latter equation and the definitions (1.3), (1.4), we immediately obtain

COROLLARY 1.12. If $f \in C^\infty(G)$ satisfies $f(bn) = f(b)$ for $(b, n) \in G \times N$ (i.e. $f \in C^\infty(G/N)$) then

$$\Omega f = \left[\sum_{i=1}^r (\delta_{H_i}^2 + 2\rho(H_i)\delta_{H_i}) - \sum_{i=1}^s u_i^2 \right] f.$$

Also since δ_x, δ_x commute we have for $f \in C^\infty(G)$ satisfying $f(kb) = f(b)$ for $(k, b) \in K \times G$ (i.e. $f \in C^\infty(K \backslash G)$)

$$\Omega f = \left[\sum_{i=1}^r (\delta_{H_i}^2 + 2\rho(H_i)\delta_{H_i}) + 2 \sum_{j=1}^q \sum_{i=1}^{m_{\alpha_j}} e^{-2\alpha_j} \delta_{x_{i(j)}}^2 \right] f.$$

§2. Proof of Theorem 1.10.

Set $w_i = y_i$, $1 \leq i \leq t$, $w_i = H_{i-t}$, $t+1 \leq i \leq r+t$, $k_{ij}(kan) = \langle \text{Ad}(k)w_j, w_i \rangle$; $k_{ij} \in C^\infty(G)$. By Proposition 1.8 $\{w_i\}_{i=1}^{r+t}$ is an orthonormal basis of \mathfrak{p} and $\Omega \stackrel{(a.)}{=} -\sum_{i=1}^t \tilde{w}_i^2 - \sum_{i=1}^t \tilde{z}_i^2 + \sum_{i=1}^{r+t} \tilde{w}_i^2$. The main point is to compute $\sum_{j=1}^{r+t} \tilde{w}_j^2$.

LEMMA 2.1. $\tilde{w}_j = \sum_{i=1}^r k_{j,t+i} \delta_{H_i} + \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha_\mu}} k_{j,\nu(\mu)} \{-\delta_{x_{\nu(\mu)}} + \sqrt{2} e^{-\alpha_\mu} \delta_{x_{\nu(\mu)}}\}$; here $k_{j,\nu(\mu)} = \langle \text{Ad}(\cdot)y_{\nu(\mu)}, w_j \rangle$, $1 \leq j \leq r+t$.

PROOF. For $(a, b, x) \in G \times G \times \mathfrak{g}$ define $L_{axb}: C^\infty(G) \xrightarrow{\text{linear}} \mathcal{C}$ by

$$(2.2) \quad L_{axb}f = \frac{d}{dt} f(a(\exp tx)b) |_{t=0}$$

for $f \in C^\infty(G)$. For $kan \in G = KAN$ arbitrary $(\tilde{w}_j, f)(kan) = (d/dt)f((\exp tw_j)kan) |_{t=0}$ (by (1.4)) $= (d/dt)f(k \exp t \text{Ad}(k^{-1})w_j, an) |_{t=0} = L_{k \Delta \text{Ad}(k^{-1})w_j an} f \stackrel{(i)}{=} \sum_{i=1}^{r+t} k_{ij}(k^{-1}) L_{k w_i an} f$. For $1 \leq i \leq t$, $w_i = y_i = \sqrt{2} x_i - z_i$ (by (1.7)). $x_i = x_{\nu(\mu)} \in \mathfrak{g}_{\alpha_\mu}$, $1 \leq \nu \leq m_{\alpha_\mu}$, for some ν, μ (by (1.9)) so that $L_{k w_i an} \stackrel{(ii)}{=} \sqrt{2} L_{k x_{\nu(\mu)} an} - L_{k z_i an}$ where for $a = \exp H$, $H \in \mathfrak{a}$, $L_{k x_{\nu(\mu)} an} f = (d/dt)f(k \exp t x_{\nu(\mu)} an) |_{t=0} = (d/dt)f(ka \exp t \text{Ad}(a^{-1})x_{\nu(\mu)} n) |_{t=0} = (d/dt)f(ka \exp t e^{-\alpha_\mu(H)} x_{\nu(\mu)} n) |_{t=0} = L_{ka e^{-\alpha_\mu(H)} x_{\nu(\mu)} n} f = e^{-\alpha_\mu(H)} (\delta_{x_{\nu(\mu)}} f)(kan)$. Also by definition $L_{k z_i an} f = (\delta_{z_i} f)(kan)$. Thus (ii) becomes $L_{k w_i an} f \stackrel{(iii)}{=} \sqrt{2} e^{-\alpha_\mu(H)} (\delta_{x_{\nu(\mu)}} f)(kan) - (\delta_{z_i} f)(kan)$. On the other hand for $t+1 \leq i \leq r+t$, $w_i = H_{i-t}$ so that $\sum_{i=t+1}^{r+t} k_{ij}(k^{-1}) L_{k H_{i-t} an} f \stackrel{(iv)}{=} \sum_{i=1}^r k_{i+t,j}(k^{-1}) (\delta_{H_i} f)(kan)$. Since $k_{ij}(k^{-1}) = k_{ji}(k)$, Lemma 2.1 follows from (i), (iii), (iv), and (1.5).

LEMMA 2.2. $\sum_{j=1}^{r+t} \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha_\mu}} k_{j,\nu(\mu)} \delta_{x_{\nu(\mu)}} \sum_{i=1}^r k_{k,t+i} \delta_{H_i} = -\sum_{i=1}^r 2\rho(H_i) \delta_{H_i}$.

PROOF. For $(z, H, f) \in \mathfrak{k} \times \mathfrak{a} \times C^\infty(G)$, $kan \in G = KAN$ $(\delta_z k_{ij} \delta_H f)(kan) = (d/dt)k_{ij}(k \exp tz) |_{t=0} (\delta_H f)(kan) + k_{ij}(kan) (\delta_z \delta_H f)(kan)$. Now $k_{ij}(k \exp tz) = \langle \text{Ad}(\exp tz)w_j, \text{Ad}(k^{-1})w_i \rangle \Rightarrow (d/dt)k_{ij}(k \exp tz) |_{t=0} = \langle [z, w_j], \text{Ad}(k^{-1})w_i \rangle = \sum_{i=1}^{r+t} k_{ii}(kan) \langle [z, w_j], w_i \rangle$; i.e.

$$(2.3) \quad \delta_z k_{ij} \delta_H = \sum_{i=1}^{r+t} \langle [z, w_j], w_i \rangle k_{ii} \delta_H + k_{ij} \delta_z \delta_H$$

for $(z, H) \in \mathfrak{k} \times \mathfrak{a}$. Hence for $1 \leq j \leq r+t$

$$(2.4) \quad \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha_\mu}} k_{j,\nu(\mu)} \delta_{x_{\nu(\mu)}} \sum_{i=1}^r k_{j,t+i} \delta_{H_i} = \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha_\mu}} \sum_{i=1}^r \sum_{l=1}^{r+t} \langle [z_{\nu(\mu)}, w_{t+i}], w_l \rangle k_{j,\nu(\mu)} k_{jl} \delta_{H_i} \\ + \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha_\mu}} \sum_{i=1}^r k_{j,\nu(\mu)} k_{j,t+i} \delta_{x_{\nu(\mu)}} \delta_{H_i}.$$

Now

$$(2.5) \quad \sum_{j=1}^{r+t} k_{ji} k_{jk} = \delta_{ik} ,$$

and for $1 \leq i \leq r$, $\delta_{\nu(\mu), i+t} = 0$, $w_{i+t} = H_i$. Also from (1.7), $[z_i, y_i] \stackrel{(v)}{=} [\theta x_i, x_i]$ and since $\{x_{\nu(\mu)}\}_{\nu=1}^{m_{\alpha\mu}}$ is an orthonormal basis of $\mathfrak{g}_{\alpha\mu}$ we get $\langle [z_{\nu(\mu)}, H_i], y_{\nu(\mu)} \rangle = -B(H_i, [z_{\nu(\mu)}, y_{\nu(\mu)}]) = -B(H_i, [\theta x_{\nu(\mu)}, x_{\nu(\mu)}])$ (by (v)) $= B(\theta x_{\nu(\mu)}, [H_i, x_{\nu(\mu)}]) = -\alpha_{\mu}(H_i) \langle x_{\nu(\mu)}, x_{\nu(\mu)} \rangle = -\alpha_{\mu}(H_i)$. Thus summing over j in (2.4) gives $\sum_{j=1}^{r+t} \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} \delta_{x_{\nu(\mu)}} \sum_{i=1}^r k_{j,i+t} \delta_{H_i} = \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} \sum_{i=1}^r \langle [z_{\nu(\mu)}, H_i], y_{\nu(\mu)} \rangle \delta_{H_i} = -\sum_{\mu=1}^q (\sum_{\nu=1}^{m_{\alpha\mu}} \alpha_{\mu})(H_i) \delta_{H_i} = -\sum_{i=1}^r 2\rho(H_i) \delta_{H_i}$ (by (1.9)), which proves Lemma 2.2.

LEMMA 2.6. $\sum_{j=1}^{r+t} [\sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} (-\delta_{x_{\nu(\mu)}} + \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}})]^2 = \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} \{\delta_{x_{\nu(\mu)}}^2 + 2e^{-2\alpha_{\mu}} \delta_{x_{\nu(\mu)}}^2 - 2\sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}} \delta_{x_{\nu(\mu)}}\}$.

PROOF. Similar to (2.3) we have $k_{j,\nu(\mu)} \delta_{x_{\nu(\mu)}} k_{j,\lambda(r)} \delta_{x_{\lambda(r)}} = k_{j,\nu(\mu)} \sum_{i=1}^{r+t} \langle [z_{\nu(\mu)}, w_{\lambda(r)}], w_i \rangle k_{ji} \delta_{x_{\lambda(r)}} + k_{j,\nu(\mu)} k_{j,\lambda(r)} \delta_{x_{\nu(\mu)}} \delta_{x_{\lambda(r)}}$, and $-k_{j,\nu(\mu)} \delta_{x_{\nu(\mu)}} k_{j,\lambda(r)} \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\lambda(r)}} = -\sqrt{2} e^{-\alpha_{\mu}} k_{j,\nu(\mu)} \sum_{i=1}^{r+t} \langle [z_{\nu(\mu)}, w_{\lambda(r)}], w_i \rangle k_{ji} \delta_{x_{\lambda(r)}} - \sqrt{2} e^{-\alpha_{\mu}} k_{j,\nu(\mu)} k_{j,\lambda(r)} \delta_{x_{\nu(\mu)}} \delta_{x_{\lambda(r)}}$. Then using (2.5) again, the identity $\langle [z_{\nu(\mu)}, w_{\lambda(r)}], w_{\nu(\mu)} \rangle = \langle [z_{\nu(\mu)}, y_{\lambda(r)}], y_{\nu(\mu)} \rangle = -B([z_{\nu(\mu)}, y_{\nu(\mu)}], y_{\lambda(r)}) = -B([\theta x_{\nu(\mu)}, x_{\nu(\mu)}], y_{\lambda(r)})$ (by (v) above) $= 0$ (since $[\theta x, x] \in \mathfrak{a}$ for $x \in \mathfrak{g}_{\alpha}$, $\alpha \in \Sigma$), and the fact that δ_z, δ_x commute for $(z, x) \in \mathfrak{k} \times \mathfrak{n}$, the statement of Lemma 2.6 follows.

We are now in a position to compute $\sum_{j=1}^{r+t} \tilde{w}_j^2$. By Lemma 2.1 $\sum_{j=1}^{r+t} \tilde{w}_j^2 = \sum_{j=1}^{r+t} [(\sum_{i=1}^r k_{j,i+t} \delta_{H_i})^2 + \sum_{i=1}^r k_{j,i+t} \delta_{H_i} \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} (-\delta_{x_{\nu(\mu)}} + \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}}) + \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} (-\delta_{x_{\nu(\mu)}} + \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}}) \sum_{i=1}^r k_{j,i+t} \delta_{H_i} + \{\sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} (-\delta_{x_{\nu(\mu)}} + \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}})\}^2]$. Now $\sum_{j=1}^{r+t} (\sum_{i=1}^r k_{j,i+t} \delta_{H_i})^2 = \sum_{i=1}^r \sum_{j=1}^{r+t} k_{j,i+t} k_{j,i+t} \delta_{H_i} \delta_{H_i}$ (since $\delta_{H_i} k_{j,i+t} \delta_{H_i} = k_{j,i+t} \delta_{H_i} \delta_{H_i} = \sum_{i=1}^r \delta_{H_i}^2$ (by (2.5))). Also the sum over j of the second term in the above bracket is zero by (2.5) since $\delta_{i+t, \nu(\mu)} = 0$ for $1 \leq i \leq r$. Similarly $\sum_{j=1}^{r+t} \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} k_{j,\nu(\mu)} \sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}} \sum_{i=1}^r k_{j,i+t} \delta_{H_i} = 0$ by (2.5). Lemma 2.2 and Lemma 2.6 therefore imply

PROPOSITION 2.7. $\sum_{j=1}^{r+t} \tilde{w}_j^2 = \sum_{i=1}^r (\delta_{H_i}^2 + 2\rho(H_i) \delta_{H_i}) + \sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} \{\delta_{x_{\nu(\mu)}}^2 + 2e^{-2\alpha_{\mu}} \delta_{x_{\nu(\mu)}}^2 - 2\sqrt{2} e^{-\alpha_{\mu}} \delta_{x_{\nu(\mu)}} \delta_{x_{\nu(\mu)}}\}$.

Since $\sum_{\mu=1}^q \sum_{\nu=1}^{m_{\alpha\mu}} \delta_{x_{\nu(\mu)}}^2 = \sum_{i=1}^t \delta_{z_i}^2$ Proposition 2.7 and equation (a.) preceding Lemma 2.1 clearly imply formula (1.11) of Theorem 1.10. Given $f \in C^\infty(G)$, $ka \in KA$ (1.3) implies for $z \in \mathfrak{k}$, $(\delta_z f)(ka) = (d/dt) f(ka \exp t \text{Ad}(a^{-1})z)|_{t=0} = (\text{Ad}(a^{-1})z f)(ka)$ so that

$$(2.8) \quad \left(\left(\sum_{i=1}^t \delta_{z_i} \right) f \right) (ka) = \sum_{i=1}^t (\text{Ad}(a^{-1})z_i \text{Ad}(a^{-1})z_i f)(ka) .$$

On the other hand we have

LEMMA 2.9. *Let $\{E_i\}$ be an orthonormal basis of \mathfrak{k} . Then*

$$\left(\left(\sum_i \tilde{E}_i^2 \right) f \right) (k \exp y) = \sum_i (\text{Ad}(\exp -y) E_i)^2 f (k \exp y) \quad \text{for } (k, y) \in K \times \mathfrak{p}.$$

Lemma 2.9 follows essentially from the fact that $\text{Ad}(k^{-1})$ is an orthogonal transformation of \mathfrak{p} . Taking $\{E_i\} = \{u_i\} \cup \{z_i\}$ (see Proposition 2.8) and $y \in \mathfrak{a}$ in Lemma 2.9 we get for $a = \exp y \in A$

$$(2.10) \quad \left(\left(\sum_{i=1}^s \tilde{u}_i^2 + \sum_{i=1}^t \tilde{z}_i^2 \right) f \right) (ka) = \left(\left(\sum_{i=1}^s (\text{Ad}(a^{-1})u_i)^2 + \sum_{i=1}^t (\text{Ad}(a^{-1})z_i)^2 \right) f \right) (ka) \\ = \left(\left(\sum_{i=1}^s u_i^2 + \sum_{i=1}^t (\text{Ad}(a^{-1})z_i)^2 \right) f \right) (ka)$$

since $u_i \in \mathfrak{m}$, the centralizer of \mathfrak{a} in \mathfrak{k} . Thus (2.8) and (2.10) imply $(\sum_{i=1}^t \delta_{z_i}^2) f = (\sum_{i=1}^s (\tilde{u}_i^2 - u_i^2) + \sum_{i=1}^t \tilde{z}_i^2) f$ on KA , which completes the proof of Theorem 1.10.

In the very special case when $G = SL(2, \mathbf{R})$, Theorem 1.10 reduces to a result of Takahashi. Namely we have the following data. If \mathbf{Z} is the ring of integers, an Iwasawa decomposition of G is given by the diffeomorphism $\Phi: \mathbf{R}/4\pi\mathbf{Z} \times \mathbf{R} \times \mathbf{R} \rightarrow G$ defined by

$$(2.11) \quad \Phi(\theta, t, \xi) = k_\theta a_t n_\xi \stackrel{\text{def.}}{=} \begin{bmatrix} \cos \theta/2 & \sin \theta/2 \\ -\sin \theta/2 & \cos \theta/2 \end{bmatrix} \begin{bmatrix} e^{t/2} & 0 \\ 0 & e^{-t/2} \end{bmatrix} \begin{bmatrix} 1 & \xi \\ 0 & 1 \end{bmatrix} \\ \in KAN \stackrel{\text{def.}}{=} SO(2) \exp \mathfrak{a} \exp \mathfrak{n}, \quad \text{where}$$

$$(2.12) \quad \mathfrak{a} = \left\{ \begin{bmatrix} r & 0 \\ 0 & -r \end{bmatrix} \mid r \in \mathbf{R} \right\}, \quad \mathfrak{n} = \mathfrak{g}_\alpha = \left\{ \begin{bmatrix} 0 & r \\ 0 & 0 \end{bmatrix} \mid r \in \mathbf{R} \right\}, \\ \alpha \begin{bmatrix} r & 0 \\ 0 & -r \end{bmatrix} = 2r, \quad \Sigma^+ = \{\alpha\}; \mathfrak{m} = 0.$$

The vectors $H_1 = \begin{bmatrix} 1/2\sqrt{2} & 0 \\ 0 & -1/2\sqrt{2} \end{bmatrix}$, $x_1 = \begin{bmatrix} 0 & 1/2 \\ 0 & 0 \end{bmatrix}$ are an orthonormal basis of \mathfrak{a} , \mathfrak{n} respectively. Here $B(x, y) = 4 \text{ trace } xy$ and $\theta x = -x^\theta$ for $x, y \in \mathfrak{g} = \mathfrak{sl}(2, \mathbf{R})$. Hence, by (1.7), $Z_1 = \begin{bmatrix} 0 & 1/2\sqrt{2} \\ -1/2\sqrt{2} & 0 \end{bmatrix}$ and since K is abelian δ_{z_1} and \tilde{Z}_1 (in (1.3), (1.4)) coincide. Theorem 1.10 therefore gives

$$(2.13) \quad \Omega = \delta_{H_1}^2 + \frac{1}{\sqrt{2}} \delta_{H_1} + 2e^{-2\alpha} \delta_{x_1}^2 - 2\sqrt{2} e^{-\alpha} \delta_{z_1} \delta_{x_1}.$$

For $s, t \in \mathbf{R}$, $\exp sH_1 a_t = \exp \begin{bmatrix} (s/\sqrt{2} + t)/2 & 0 \\ 0 & -(s/\sqrt{2} + t)/2 \end{bmatrix} = a_{s/\sqrt{2} + t}$. For $\psi \in C^\infty(\mathbf{R}/4\pi\mathbf{Z} \times \mathbf{R} \times \mathbf{R})$ it follows that $(\delta_{H_1}(\psi \cdot \Phi^{-1})) \cdot \Phi(\theta, t, \xi) = \delta_{H_1}(\Psi \cdot \Phi^{-1})$

$(k_\theta a_t h_\xi) = (d/ds)\psi \cdot \Phi^{-1}(k_\theta(\exp sH_1)a_t n_\xi)|_{s=0} = (d/ds)\psi(\theta, s/\sqrt{2} + t, \xi)|_{s=0} = (1/\sqrt{2})$
 $(\partial\psi/\partial t)(\theta, t, \xi)$. Similarly $k_\theta \exp sZ_1 = \exp \begin{bmatrix} 0 & \theta/2 \\ -\theta/2 & 0 \end{bmatrix} \exp \begin{bmatrix} 0 & s/2\sqrt{2} \\ -s/2\sqrt{2} & 0 \end{bmatrix} =$
 $\exp \begin{bmatrix} 0 & (\theta + s/\sqrt{2})/2 \\ (\theta + s/\sqrt{2})/2 & 0 \end{bmatrix} = k_{\theta+s/\sqrt{2}} \Rightarrow (\delta_{Z_1}(\psi \cdot \Phi^{-1})) \cdot \Phi = (1/\sqrt{2})(\partial\psi/\partial\theta)$,
 and $(\delta_{X_1}(\psi \cdot \Phi^{-1})) \cdot \Phi = (1/2)(\partial\psi/\partial\xi)$ since $\exp sX_1 n_\xi = n_{s/2+\xi}$. Using $e^\alpha(k_\theta a_t n_\xi) = e^t$ (cf. (1.5)) we therefore obtain from (2.13)

COROLLARY 2.14 (R. Takahashi, Lemma 3 of [3]). *For $G = SL(2, \mathbf{R})$, the expression of Ω in terms of the Iwasawa coordinates (θ, t, ξ) in (3.11) is given by $(\Omega(\psi \cdot \Phi^{-1})) \cdot \Phi = (1/2)(\partial^2\psi/\partial t^2) + (1/2)(\partial\psi/\partial t) + (1/2)e^{-2t}(\partial^2\psi/\partial\xi^2) - e^{-t}(\partial^2\psi/\partial\theta\partial\xi)$ for $\psi \in C^\infty(\mathbf{R}/4\pi\mathbf{Z} \times \mathbf{R} \times \mathbf{R})$.*

The reader should note that the Casimir operator defined in [3] is not the same as our Ω , but rather coincides with -2Ω .

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Present Address:

DEPARTMENT OF MATHEMATICS AND STATISTICS,
 UNIVERSITY OF MASSACHUSETTS
 Amherst, Massachusetts 01003
 U.S.A.