

On the Power Series Coefficients of the Riemann Zeta Function

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§1. Introduction and the main result.

The Laurent expansion of the Riemann zeta function $\zeta(s)$ about the pole can be written in the form, in [2],

$$(1) \quad \zeta(s) = \frac{1}{s-1} + \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \gamma_n (s-1)^n$$

with

$$\gamma_n = \lim_{N \rightarrow \infty} \left(\sum_{k=1}^N \frac{\log^n k}{k} - \frac{\log^{n+1} N}{n+1} \right).$$

Here $\log^0 k$ mean 1 for all k including $k=1$. γ_0 is the well known Euler constant, and, for $n \geq 1$, γ_n , sometimes called generalized Euler constants, have been studied by many authors ([1], Entry 13; or [3], p. 51). In this paper we shall give an asymptotic expansion of γ_n for arbitrary large n , which yields some interesting results on γ_n . They can be found in [4].

We begin by defining some notations. Let N be a nonnegative integer, and let n be a positive integer. In order to write our theorem, we need two functions $a=a(n)$ and $b=b(n)$ which are given by the following lemma.

LEMMA 1. *If $n > c_1$, where c_1 is a sufficiently large constant, then the system of the equations*

$$(2) \quad -(n+1) \frac{y}{x^2+y^2} + \frac{1}{2} \pi - \operatorname{Im} \psi(x+iy) = 0,$$

$$(3) \quad -(n+1) \frac{x}{x^2+y^2} - \log 2\pi + \operatorname{Re} \psi(x+iy) = 0,$$

with unknown x and y , satisfying $0 < y < x$ and $n^{1/2} < x < n$, has a unique

solution $x=a$, $y=b$, where $\psi(y)$ is the logarithmic derivative of the gamma function, i.e., $\psi(z)=\Gamma'(z)/\Gamma(z)$, and $\operatorname{Re} z$ and $\operatorname{Im} z$ mean the real and imaginary parts of z , respectively.

For given n with $n > c_1$, the pair $x=a$, $y=b$ is uniquely determined by the lemma. Hence these a and b can be considered as the functions of n , so that we denote these new functions by $a=a(n)$ and $b=b(n)$. We next define the functions $\phi(z)$, $g(y)$, $f(y)$ by

$$\begin{aligned}\phi(z) &= -(n+1)\log z - z \log(2\pi i) + \log \Gamma(z), \\ g(y) &= \operatorname{Re} \phi(a+iy) \quad \text{and} \quad f(y) = \operatorname{Im} \phi(a+iy)\end{aligned}$$

with a real variable y , namely,

$$\begin{aligned}g(y) &= -\frac{1}{2}(n+1)\log(a^2+y^2) - a \log 2\pi + \frac{1}{2}\pi y + \operatorname{Re} \log \Gamma(a+iy), \\ f(y) &= -(n+1)\arctan\left(\frac{y}{a}\right) - y \log 2\pi - \frac{1}{2}\pi a + \operatorname{Im} \log \Gamma(a+iy).\end{aligned}$$

Moreover, define the sequences h_k , u_k , v_k ($k=0, 1, 2, 3, \dots$) by

$$(4) \quad \sum_{k=0}^{\infty} h_k (y-b)^k = \exp\left\{\phi(a+iy) - \phi(a+ib) + \frac{1}{2}\phi''(a+ib)(y-b)^2\right\},$$

$$u_k = \operatorname{Re} h_k \quad \text{and} \quad v_k = \operatorname{Im} h_k,$$

then we have the following theorem which states the asymptotic expansion of γ_n .

THEOREM. *Let N be a nonnegative integer. If $n > c_2 \Gamma(N/3 + 17/6)$, where c_2 is a sufficiently large constant, then*

$$\begin{aligned}\gamma_n &= \frac{1}{\pi} n! e^{g(b)} \sum_{k=0}^N |h_{2k}| 2^{k+1/2} \Gamma\left(k + \frac{1}{2}\right) \{g''(b)^2 + f''(b)^2\}^{-k/2-1/4} \\ &\quad \times \cos\left\{f(b) - \left(k + \frac{1}{2}\right)\arctan\left(\frac{f''(b)}{g''(b)}\right) + \arctan\left(\frac{v_{2k}}{u_{2k}}\right)\right\} \\ &\quad + O\left\{\Gamma\left(N + \frac{3}{2}\right) 2^N n! e^{g(b)} n^{-N/3+1/6} \log^{2N/3-1/3} n\right\}.\end{aligned}$$

§2. Proof of Lemma 1 and some other lemmas.

PROOF OF LEMMA 1. The equations (2) and (3) are the imaginary and real parts of the equation

$$\phi'(x+iy) = -\frac{n+1}{x+iy} - \log(2\pi i) + \psi(x+iy) = 0,$$

respectively. Hence we will show that the equation $z\phi'(z)=0$ with $z=x+iy$ has a unique solution under the assumption of the lemma. Now, we put

$$(5) \quad h(y) = \text{Im } z\phi'(z) = -y \log 2\pi - \frac{1}{2}\pi x + \text{Im}(x+iy)\psi(x+iy).$$

Then, if we use the asymptotic expansion of $\psi(z)$ which is derived from the asymptotic expansion of $\log \Gamma(z)$ ([6], p. 251), we can prove that, for sufficiently large x , $h(y)$ is steadily increasing in $0 \leq y \leq x$, having the values $h(0) < 0$ and $h(x) > 0$. Therefore the equation $h(y)=0$ has a unique solution, say, y_x , in $0 < y < x$. We further set $z_x = x + iy_x$ and define

$$(6) \quad u(x) = \text{Re } z_x\phi'(z_x) = -(n+1) + x \log 2\pi + \frac{1}{2}\pi y_x + \text{Re}(x+iy_x)\psi(x+iy_x) \\ = -n-1 + \left(x + \frac{y_x}{x}\right)(-\log 2\pi + \text{Re } \psi(x+iy_x)).$$

Then, using the asymptotic expansion of $\psi(z)$, we can prove that, for sufficiently large n , $u(x)$ is steadily increasing in $n^{1/2} \leq x \leq n$, having the values $u(n^{1/2}) < 0$ and $u(n) > 0$. Hence the equation $u(x)=0$ has a unique solution, say a , in $n^{1/2} < x < n$. It follows that $z=a+iy_a$ is the unique solution of $z\phi'(z)=0$. This completes the proof.

LEMMA 2. For sufficiently large n ,

$$(7) \quad n \log^{-1} n < a < n \log^{-1} n + 2n \log^{-2} n \log \log n,$$

and

$$(8) \quad b = \frac{1}{2}\pi n \log^{-2} n + O(n \log^{-3} n \log \log n).$$

PROOF. We have from (5)

$$(9) \quad -b \log 2\pi - \frac{1}{2}\pi a + \text{Im}(a+ib)\psi(a+ib) = 0.$$

Using (9), and $\psi(z) = \log z - z^{-1}/2 + O(|z|^{-2})$, we have $b/a < 2 \log^{-1} a$, and further

$$(10) \quad \frac{b}{a} = \frac{1}{2}\pi \log^{-1} a + O(\log^{-2} a).$$

It follows from (6) that

$$(11) \quad \begin{aligned} n &= -1 + \left(a + \frac{b}{a}\right)(-\log 2\pi + \operatorname{Re} \psi(a + ib)) \\ &= a \log a - a \log 2\pi + O(a \log^{-1} a). \end{aligned}$$

Here we suppose that $a \leq n \log^{-1} n$. Then we have from (11) $n < a \log a \leq n \log^{-1} n \log(n \log^{-1} n) < n$ for sufficiently large n . This is a contradiction, and hence $a > n \log^{-1} n$. We next suppose that $a \geq n \log^{-1} n + 2a \log^{-2} n \log \log n$. Then, by (11), we have $n > a \log a - 2a \geq n \log^{-1} n(1 + 2 \log^{-1} n \log \log n) \log\{n \log^{-1} n(1 + 2 \log^{-1} n \log \log n)\} - 2n \log^{-1} n(1 + 2 \log^{-1} n \times \log \log n) > n$. This is a contradiction, and hence $a < n \log^{-1} n + 2 \log^{-2} n \times \log \log n$, which completes (7). Moreover, we can obtain (8) from (7) and (10).

LEMMA 3. $g(y)$ is steadily increasing in $0 \leq y \leq b$, has the maximum at $y=b$, and is steadily decreasing in $b \leq y \leq a$, having the properties $g(b) - g(b - \Delta) > (1/3) \log^3 a$, $g(b) - g(b + \Delta) > (1/3) \log^3 a$, where $\Delta = a^{1/2} \log a$; further, $g''(b) = -a^{-1} \log a + (\log 2\pi - 1)a^{-1} + O(a^{-1} \log^{-1} a)$, and $f''(b) = a^{-1} \pi + O(a^{-1} \log^{-1} a)$.

PROOF. Using the asymptotic expansion of $\log \Gamma(z)$, we have, for $0 \leq y \leq a$,

$$g''(y) = -\left(n + \frac{3}{2}\right) \frac{a^2 - y^2}{(a^2 + y^2)^2} - \frac{a}{a^2 + y^2} + O(a^{-3}) < 0.$$

It follows that $g'(y)$ is steadily decreasing in $0 \leq y \leq a$. Moreover, $g'(0) > 0$ and $g'(a) < 0$ for sufficiently large n , which leads to that $g(y)$ is steadily increasing in $0 \leq y \leq b$, has the maximum at $y=b$, and is steadily decreasing in $b \leq y \leq a$, since $g'(b) = 0$.

Let $\delta = (\pi/2) \log^{-1} a$. Then $a + ib = a(1 + i\delta)(1 + O(\delta))$. It follows that

$$\begin{aligned} \phi''(a + ib) &= \left(n + \frac{3}{2}\right) a^{-2} (1 + i\delta)^{-2} + a^{-1} (1 + i\delta)^{-1} + O(a^{-1} \delta) \\ &= a^{-1} (\log a - \log 2\pi) (1 - 2i\delta) + a^{-1} + O(a^{-1} \delta) \end{aligned}$$

Hence, we obtain $g''(b) = -a^{-1} \log a + (\log 2\pi - 1)a^{-1} + O(a^{-1} \delta)$ and $f''(b) = a^{-1} \pi + O(a^{-1} \delta)$.

Now, if we use Taylor's theorem ([6], p. 96), then we have

$$(12) \quad g(b + \Delta) = g(b) + \frac{1}{2} g''(b) \Delta^2 + \frac{1}{6} g'''(\xi) \Delta^3$$

with $b < \xi < b + \Delta$. Using (12), and

$$g'''(\xi) = -\left(n + \frac{3}{2}\right) \frac{2\xi^5 - 4a^2\xi^3 - 6a^4\xi}{(a^2 + \xi^2)^4} + \frac{2a\xi}{(a^2 + \xi^2)^2} + O(a^{-4}),$$

we have

$$g'''(\xi) < 12\left(n + \frac{3}{2}\right)a^{-3} + 2a^{-2} + O(a^{-4}) < 12a^{-2} \log a.$$

It follows that

$$g(b + \Delta) - g(b) < -\frac{3}{8}\Delta^2 a^{-1} \log a + 2\Delta^3 a^{-2} \log a < -\frac{1}{3} \log^3 a,$$

since $g''(b) < -(3/4)a^{-1} \log a$. In case $g(b) - g(b - \Delta)$, we have the same result as well, which completes the proof.

LEMMA 4. For $j \geq 0$, $|h_j| \leq (a^{-1} \log a)^{2j/3}$.

PROOF. We first prove that, if p_j ($j=0, 1, 2, 3, \dots$) is the power series coefficients defined by $\phi(a + iy) = \sum_{j=0}^{\infty} p_j (y - b)^j$, then, for $j \geq 2$,

$$(13) \quad |p_j| \leq \frac{1}{j} (a^{-1} \log a)^{j-1}.$$

In case $j \geq 3$, differentiating $\phi(a + iy)$ two times by y , we have

$$j(j-1)p_j = -\frac{1}{2\pi i} \int_C \phi''(a + iw)(w - b)^{-j+1} dw,$$

where C is the circle with center b and radius $a/2$ in the positive direction. We have, on C , $\min |a + iw| = a/2 + O(a \log^{-2} a)$, and $|\phi''(a + iw)| < 5a^{-1} \log a$. Therefore, for $j \geq 3$,

$$|p_j| < \frac{1}{j} (a^{-1} \log a)^{j-1} \frac{5}{j-1} (2 \log^{-1} a)^{j-2} < \frac{1}{j} (a^{-1} \log a)^{j-1}.$$

In case $j=2$, we have, by Lemma 3, $|p_2| = (1/2)|g''(b) + if'''(b)| < (1/2)a^{-1} \log a$, which implies (13) for all $j \geq 2$.

We next prove the inequality in the lemma. If we differentiate (4), write the power series expansion with the variable $y - b$, and equate the terms of order 0, 1, 2, 3, \dots of both sides, then we have the system of the following infinite set of equations;

$$\begin{aligned} h_j &= 0 & (j=1, 2), \\ jh_j &= \sum_{k=2}^{j-1} (k+1)p_{k+1}h_{j-1-k} & (j=3, 4, 5, 6, \dots). \end{aligned}$$

In order to prove the lemma, we use the mathematical induction with respect to j . For $j=0, 1, 2$, the inequality holds, since $h_0=1$, $h_1=0$, and $h_2=0$. We thus suppose that the inequality holds for $j=0, 1, 2, 3, \dots, J-1$. Then we have from (13)

$$|Jh_J| = \left| \sum_{k=2}^{J-1} (k+1)p_{k+1}h_{J-1-k} \right| \leq \sum_{k=2}^{J-1} (a^{-1} \log a)^k (a^{-1} \log a)^{2(J-1-k)/3} \\ < \{1 - (a^{-1} \log a)^{1/3}\}^{-1} (a^{-1} \log a)^{2J/3} < J(a^{-1} \log a)^{2J/3},$$

since $\{1 - (a^{-1} \log a)^{1/3}\}^{-1} < 3$ for sufficiently large a , which implies that the inequality holds for $j=J$. This completes the proof.

§ 3. Proof of Theorem.

In this section we use Vinogradov's symbol \ll ([3], p. XVI). A key idea in our proof is the fact that we can apply the saddle point method to our integral expression of γ_n . Here, we shall prove that $\gamma_n = I_1 + I_2 + I_3 + I_4 + I_5 + I_6 + I_7 + I_8 + J_9$, I 's are the error terms, and J_9 is the main term.

We first show

$$(14) \quad \gamma_n = \frac{1}{\pi} n! \operatorname{Re} \int_0^\infty (a+iy)^{-n-1} (2\pi)^{-a-iy} e^{-\pi y/2 + \pi i a/2} \Gamma(a+iy) \zeta(a+iy) dy \\ + \frac{1}{\pi} n! \operatorname{Re} \int_0^\infty (a+iy)^{-n-1} (2\pi)^{-a-iy} e^{\pi y/2 - \pi i a/2} \Gamma(a+iy) \zeta(a+iy) dy \\ = I_1 + J_2,$$

say. We have from (1)

$$\gamma_n = \frac{n!}{2\pi i} \int_C z^{-n-1} \zeta(1-z) dz,$$

where C is a contour enclosing 0. Using the order of $\zeta(s)$ ([5], pp. 81-82), we can deform C into the line from $a-i\infty$ to $a+i\infty$. Moreover, using the functional equation of $\zeta(s)$ ([5], p. 13), and changing the variable z to y with $z=a+iy$, we have

$$\gamma_n = \frac{n!}{2\pi} \int_{-\infty}^\infty (a+iy)^{-n-1} 2(2\pi)^{-a-iy} \cos\left\{\frac{1}{2}\pi(a+iy)\right\} \Gamma(a+iy) \zeta(a+iy) dy.$$

We now divide the integral into two parts; $\int_{-\infty}^0 + \int_0^\infty$. If we notice that the first integral is the complex conjugate of the second, we obtain (14), using $2 \cos\{\pi(a+iy)/2\} = \exp(-\pi y/2 + \pi i a/2) + \exp(\pi y/2 - \pi i a/2)$.

We next prove

$$(15) \quad I_1 \ll n! e^{\sigma(b)} n^{-N/3-1/3} \log^{N/3+1/3} n .$$

Using asymptotic expansion of $\log \Gamma(a+iy)$, we have

$$\begin{aligned} I_1 &\ll n! \int_0^\infty (a^2+y^2)^{-(n-a+3/2)/2} (2\pi e)^{-a} e^{-\pi y/2-y \arctan(y/a)} dy \\ &\ll n! \exp\{-(n-a+1)\log a - a(\log 2\pi + 1)\} \ll n! \exp\{g(0)\} , \end{aligned}$$

since $g(0) = -(n-a+3/2)\log a - a(\log 2\pi + 1) + O(1)$. It follows from Lemma 3 that

$$I_1 \ll n! \exp\{g(b-\Delta)\} \ll n! \exp\left\{g(b) - \frac{1}{3} \log^3 a\right\} ,$$

where $\Delta = a^{1/2} \log a$. By Lemma 1 and Lemma 2, we can prove $\log^2 a < N+4$, and hence we obtain (15).

We next divide J_2 into two parts;

$$J_2 = \frac{1}{\pi} n! \operatorname{Re} \left(\int_0^a + \int_a^\infty \right) = J_3 + I_2 ,$$

say. We have by Lemma 2 and Lemma 3

$$(16) \quad \begin{aligned} I_2 &\ll n! (2\pi e)^{-a} \int_a^\infty (a^2+y^2)^{-(n-a+3/2)/2} e^{\pi y/2-y \arctan(y/a)} dy \\ &\ll n! e^{\sigma(b)} n^{-N/3-4/3} \log^{N/3+4/3} n . \end{aligned}$$

We next divide J_3 into two parts;

$$J_3 = \frac{1}{\pi} n! \operatorname{Re} \left\{ \int_0^a e^{\phi(a+iy)} dy + \int_0^a e^{\phi(a+iy)} (\zeta(a+iy) - 1) dy \right\} = J_4 + I_3 ,$$

say. If we use the estimate $\zeta(a+iy) - 1 = O(2^{-a})$, then we obtain

$$(17) \quad I_3 \ll n! e^{\sigma(b)} n \log^{-1} n \exp(-n \log^{-1} n \log 2) .$$

We further divide J_4 into three parts;

$$J_4 = \frac{1}{\pi} n! \operatorname{Re} \left(\int_0^{b-\Delta} + \int_{b-\Delta}^{b+\Delta} + \int_{b+\Delta}^a \right) = I_4 + J_5 + I_5 ,$$

say. We then have by Lemma 3

$$\int_0^{b-\Delta} e^{\phi(a+iy)} dy \leq \int_0^{b-\Delta} e^{\sigma(y)} dy < a e^{\sigma(b-\Delta)} ,$$

so that

$$(18) \quad I_4 \ll n! e^{\sigma(b)} n^{-N/3-1/3} \log^{N/3-1/3} n .$$

We have also

$$(19) \quad I_5 \ll n! e^{\sigma(b)} n^{-N/3-1/3} \log^{N/3-1/3} n .$$

We next consider J_6 . We can write

$$J_6 = \frac{1}{\pi} n! \operatorname{Re} \int_{b-\Delta}^{b+\Delta} \exp \left\{ \phi(a+ib) - \frac{1}{2} \phi''(a+ib)(y-b)^2 \right\} \\ \times \exp \left\{ \phi(a+iy) - \phi(a+ib) + \frac{1}{2} \phi''(a+ib)(y-b)^2 \right\} dy .$$

We know that $\phi(a+iy)$ has the nearest singularity at $y=ia$, so that the power series $\sum_{k=0}^{\infty} h_k(y-b)^k$ defined by (4) converges in $b-\Delta \leq y \leq b+\Delta$. We now divide the power series into two parts;

$$\sum_{k=0}^{\infty} h_k(y-b)^k = \sum_{k=0}^{2N+1} h_k(y-b)^k + U_N$$

with

$$U_N = \sum_{k=2N+2}^{\infty} h_k(y-b)^k .$$

We thus have

$$J_6 = \frac{1}{\pi} n! \operatorname{Re} \int_{b-\Delta}^{b+\Delta} \exp \left\{ \phi(a+ib) - \frac{1}{2} \phi''(a+ib)(y-b)^2 \right\} \sum_{k=0}^{2N+1} h_k(y-b)^k dy \\ + \frac{1}{\pi} n! \operatorname{Re} \int_{b-\Delta}^{b+\Delta} \exp \left\{ \phi(a+ib) - \frac{1}{2} \phi''(a+ib)(y-b)^2 \right\} U_N dy = J_7 + I_6 ,$$

say. We are going to estimate I_6 . We have by Lemma 4

$$|U_N| \leq \sum_{k=2N+2}^{\infty} |h_k| |y-b|^k < \sum_{k=2N+2}^{\infty} \{(a^{-1} \log a)^{2/3} \Delta\}^k < 2 \{(a^{-1} \log a)^{2/3} \Delta\}^{2N+2} ,$$

since $(a^{-1} \log a)^{2/3} \Delta < 1/2$. We have further, for $b-\Delta \leq y \leq b+\Delta$, $(1/2) \phi''(b) \times (y-b)^2 < 0$, so that

$$I_6 \ll n! \Delta e^{\sigma(b)} \{(a^{-1} \log a)^{2/3} \Delta\}^{2N+2} .$$

Hence we get

$$(20) \quad I_6 \ll n! e^{\sigma(b)} n^{-N/3+1/6} \log^{11N/3+25/6} n .$$

We next divide J_7 into three parts;

$$J_7 = \frac{1}{\pi} n! \operatorname{Re} \sum_{k=0}^{2N+1} h_k \exp\{\phi(a+ib)\} \left(\int_{-\infty}^{\infty} - \int_{b+\Delta}^{\infty} - \int_{-\infty}^{b-\Delta} \right) = J_9 + I_7 + I_8,$$

say. We are going to estimate I_7 . For all k , the integrals in I_7 are estimated by

$$\ll \Delta^{k+1} \int_1^{\infty} \exp\left(-\frac{1}{3} \log^3 ay^2\right) y^k dy,$$

since $g''(b) < -(1/3)a^{-1} \log a$. We can now write the integrand as

$$\exp\left(-\frac{1}{6} \log^3 ay^2\right) \times \exp\left(-\frac{1}{6} \log^3 ay^2\right) y^k.$$

We see that the second term has the maximum at $y = (3k \log^{-1} a)^{1/2} < 1$. Hence the integrals are estimated by

$$\begin{aligned} &\ll \Delta^{k+1} \exp\left(-\frac{1}{6} \log^3 a\right) \int_1^{\infty} \exp\left(-\frac{1}{6} \log^3 ay^2\right) dy \\ &< \Delta^{k+1} \exp\left(-\frac{1}{6} \log^3 a\right) \int_1^{\infty} y \exp\left(-\frac{1}{6} \log^3 ay^2\right) dy \\ &= 3\Delta^{k+1} \exp\left(-\frac{1}{6} \log^3 a\right) \log^{-3} a. \end{aligned}$$

It follows that

$$(21) \quad I_7 \ll n! e^{g(b)} n^{-N/3-5/6} \log^{N/3-7/6} n.$$

We have also

$$(22) \quad I_8 \ll n! e^{g(b)} n^{-N/3-5/6} \log^{N/3-7/6} n.$$

We finally consider J_9 . If we notice that the integrals in J_9 are zeros for all odd k , we get

$$J_9 = \frac{1}{\pi} n! \operatorname{Re} \sum_{k=0}^N h_{2k} e^{\phi(a+ib)} \int_{-\infty}^{\infty} \exp\left\{-\frac{1}{2} \phi''(a+ib)(y-b)^2\right\} (y-b)^{2k} dy.$$

Here we change the variable y to x with $x = (y-b)^2$ in the integral, and recall the Euler's expression of $\Gamma(z)$ ([6], p. 241). Then we see that the integral is expressed by

$$\Gamma\left(k + \frac{1}{2}\right) 2^{k+1/2} \{g''(b)^2 + f''(b)^2\}^{-k/2-1/4} \exp\left\{-i\left(k + \frac{1}{2}\right) \arctan\left(\frac{f''(b)}{g''(b)}\right)\right\}.$$

Hence we finally obtain

$$(23) \quad J_9 = \frac{1}{\pi} n! e^{g(b)} \sum_{k=0}^N |h_{2k}| 2^{k+1/2} \Gamma\left(k + \frac{1}{2}\right) \{g''(b)^2 + f''(b)^2\}^{-k/2-1/4} \\ \times \cos\left\{f(b) - \left(k + \frac{1}{2}\right) \arctan\left(\frac{f''(b)}{g''(b)}\right) + \arctan\left(\frac{v_{2k}}{u_{2k}}\right)\right\}.$$

If we take account of the fact that $\gamma_n = I_1 + I_2 + I_3 + I_4 + I_5 + I_6 + I_7 + I_8 + J_9$, we get the formula from (15), (16), (17), (18), (19), (20), (21), (22), and (23). Although the O -term in the theorem is sharper than what we get, we can obtain the theorem, taking $N+1$ instead of N , and estimating the last term in the sum, and the O -term. This completes the proof.

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