

## Hilbert Spaces of Analytic Functions and the Gegenbauer Polynomials

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### Introduction.

Let  $F$  be the Fock type Hilbert space of analytic functions  $f(z)$  of  $n$  complex variables  $z = (z_1, z_2, \dots, z_n) \in C^n$ , with the scalar product

$$(f, g) = \pi^{-n} \int_{C^n} \overline{f(z)} g(z) \exp(-|z_1|^2 - \dots - |z_n|^2) dz_1 \cdots dz_n,$$

with

$$dz_1 \cdots dz_n = dx_1 \cdots dx_n dy_1 \cdots dy_n, \quad z_j = x_j + iy_j,$$

and let  $H$  be the usual Hilbert space  $L^2(\mathbf{R}^n)$ . Bargmann constructed in [1] a unitary mapping  $A$  from  $H$  to  $F$  given by an integral operator whose kernel is related in some definite sense to the Hermite polynomials. More precisely,  $f = A\phi$  for  $\phi \in H$  is defined by

$$f(z) = \int_{\mathbf{R}^n} A(z, q) \phi(q) d^n q$$

with

$$A(z, q) = \pi^{-n/4} \prod_{j=1}^n \exp\left\{-\frac{1}{2}(z_j^2 + q_j^2) + 2^{1/2} z_j q_j\right\}.$$

The purpose of the present paper is to show that similar constructions are possible for some other classical orthogonal polynomials.

### §1. The arguments for the Gegenbauer polynomials.

Let  $\lambda$  be a positive real number. The Gegenbauer polynomials  $C_m^\lambda$ ,  $m = 0, 1, 2, \dots$ , are defined as the coefficients in the expansion

$$(1 - 2zq + z^2)^{-\lambda} = \sum_{m=0}^{\infty} C_m^\lambda(q) z^m, \quad (-1 < q < 1, |z| < 1)$$

and have the following orthogonality relation:

$$\int_{-1}^1 C_m^\lambda(q) C_k^\lambda(q) (1 - q^2)^{\lambda - 1/2} dq = \begin{cases} 0 & (m \neq k) \\ \frac{\pi \Gamma(m + 2\lambda)}{2^{2\lambda - 1} (m + \lambda) m! [\Gamma(\lambda)]^2} & (m = k) \end{cases}.$$

Let  $\phi_m^\lambda$  stand for the normalization of  $C_m^\lambda$  with respect to  $K_\lambda = L^2((-1, 1), (1 - q^2)^{\lambda - 1/2})$ , with the scalar product

$$(\phi, \psi)_\lambda = \int_{-1}^1 \overline{\phi(q)} \psi(q) (1 - q^2)^{\lambda - 1/2} dq.$$

$\{\phi_m^\lambda \mid m = 0, 1, 2, \dots\}$  is a complete orthonormal system in  $K_\lambda$  since  $C_m^\lambda$  is a polynomial of degree  $m$ .

The Hilbert space  $F_\lambda$  consists of analytic functions  $f$  of one complex variable on  $B$ , the unit disk,  $|z| < 1$ . The inner product is given by

$$\langle f, g \rangle_\lambda = \int_B \overline{f(z)} g(z) \rho_\lambda(|z|^2) dx dy \quad (z = x + iy)$$

where

$$\rho_\lambda(t) = \begin{cases} \frac{1}{\Gamma(2\lambda - 1)} t^{\lambda - 1} \int_t^1 s^{-\lambda} (1 - s)^{2\lambda - 2} ds & (\lambda > 1/2) \\ t^{\lambda - 1} \left\{ \frac{\Gamma(1 - \lambda)}{\Gamma(\lambda)} - \frac{1}{\Gamma(2\lambda - 1)} \int_0^t s^{-\lambda} (1 - s)^{2\lambda - 2} ds \right\} & (0 < \lambda \leq 1/2) \end{cases}.$$

The next equality is easily computed for  $\lambda > 1/2$ .

$$\int_B |z|^{2m} \rho_\lambda(|z|^2) dx dy = \frac{\pi m!}{(m + \lambda) \Gamma(m + 2\lambda)}. \tag{1}$$

From the analyticity with respect to  $\lambda$ , this equality holds also for  $0 < \lambda \leq 1/2$ .

**PROPOSITION 1.** *Suppose that  $f$  is an element of  $F_\lambda$  with the power series expansion*

$$f(z) = \sum_{m=0}^{\infty} \alpha_m z^m.$$

*Then*

$$\|f\|_{\lambda}^2 = \langle f, f \rangle_{\lambda} = \sum_{m=0}^{\infty} \frac{\pi m!}{(m+\lambda)\Gamma(m+2\lambda)} |\alpha_m|^2 .$$

PROOF. For  $0 < \sigma < 1$ ,

$$\begin{aligned} & \int_{|z| \leq \sigma} |f(z)|^2 \rho_{\lambda}(|z|^2) dx dy \\ &= \sum_{m=0}^{\infty} |\alpha_m|^2 \int_{|z| \leq \sigma} |z|^{2m} \rho_{\lambda}(|z|^2) dx dy \leq \|f\|_{\lambda}^2 , \end{aligned}$$

which implies our assertion immediately.

q.e.d.

COROLLARY. If  $f \in F_{\lambda}$ , then

$$|f(z)| \leq \|f\|_{\lambda} (h_{\lambda}(|z|^2))^{1/2} ,$$

where

$$h_{\lambda}(\xi) = \sum_{m=0}^{\infty} \frac{(m+\lambda)\Gamma(m+2\lambda)}{\pi m!} \xi^m .$$

The preceding corollary asserts that the strong convergence in  $F_{\lambda}$  implies the uniform convergence on every compact subset of  $B$ . Hence  $F_{\lambda}$  is a Hilbert space and from Proposition 1 we see that a complete orthonormal system in  $F_{\lambda}$  is given by the functions

$$u_m^{\lambda}(z) = \left( \frac{(m+\lambda)\Gamma(m+2\lambda)}{\pi m!} \right)^{1/2} z^m .$$

The main theorem is

THEOREM 1. A unitary operator,  $f = A_{\lambda}\phi$ , of  $K_{\lambda}$  onto  $F_{\lambda}$  is defined by

$$f(z) = \int_{-1}^1 A_{\lambda}(z, q)\phi(q)(1-q^2)^{\lambda-1/2} dq ,$$

where

$$A_{\lambda}(z, q) = \frac{2^{\lambda-1/2}\Gamma(\lambda)\lambda}{\pi} \frac{1-z^2}{(1-2zq+z^2)^{\lambda+1}} .$$

PROOF. It is easy to see that

$$A_{\lambda}(z, q) = \frac{2^{\lambda-1/2}\Gamma(\lambda)}{\pi} \sum_{m=0}^{\infty} (m+\lambda)C_m^{\lambda}(q)z^m .$$

(i.e.,  $A_{\lambda}(z, q)$  can be regarded as a generating function for the Gegenbauer polynomials.) Therefore

$$A_\lambda(z, q) = \sum_{m=0}^{\infty} \phi_m^\lambda(q) u_m^\lambda(z). \quad (2)$$

We can consider that the right-hand side is the Fourier expansion for  $A_\lambda(z, q)$  as a function of  $q$ , because  $\sum_{m=0}^{\infty} |u_m^\lambda(z)|^2 < \infty$  for  $z \in B$ .

Let  $\phi \in K_\lambda$ , then

$$\begin{aligned} (A_\lambda \phi)(z) &= \left( \sum_{m=0}^{\infty} \overline{u_m^\lambda(z)} \phi_m^\lambda, \phi \right)_\lambda \\ &= \sum_{m=0}^{\infty} (\phi_m^\lambda, \phi)_\lambda u_m^\lambda(z). \end{aligned}$$

Hence,  $\|A_\lambda \phi\|_\lambda^2 = \sum_{m=0}^{\infty} |(\phi_m^\lambda, \phi)_\lambda|^2$ , i.e.,

$$\|A_\lambda \phi\|_\lambda = \|\phi\|_\lambda. \quad (3)$$

Substituting  $\phi_m^\lambda$  for  $\phi$ , we obtain

$$u_m^\lambda = A_\lambda \phi_m^\lambda. \quad (4)$$

It follows from (3), (4) that  $A_\lambda$  is a unitary operator of  $K_\lambda$  onto  $F_\lambda$ .

q.e.d.

The inverse operator  $A_\lambda^{-1}$ , which exists by Theorem 1, cannot be expressed as an integral operator like  $A_\lambda$ . But we have

**PROPOSITION 2.** *If  $f \in F_\lambda$ , then*

$$(A_\lambda^{-1} f)(q) = \text{l.i.m.}_{\sigma \rightarrow 1} \int_{|z| \leq \sigma} \overline{A_\lambda(z, q)} f(z) \rho_\lambda(|z|^2) dx dy,$$

or, more precisely

$$\lim_{\sigma \rightarrow 1} \int_{-1}^1 \left| (A_\lambda^{-1} f)(q) - \int_{|z| \leq \sigma} \overline{A_\lambda(z, q)} f(z) \rho_\lambda(|z|^2) dx dy \right|^2 (1 - q^2)^{\lambda - 1/2} dq = 0.$$

**PROOF.** Let  $f$  be an element of  $F_\lambda$  with the power series expansion

$$f(z) = \sum_{m=0}^{\infty} \alpha_m u_m^\lambda(z).$$

Using (2), for  $0 < \sigma < 1$ , we obtain

$$\begin{aligned} & \int_{|z| \leq \sigma} \overline{A_\lambda(z, q)} f(z) \rho_\lambda(|z|^2) dx dy \\ &= \sum_{m=0}^{\infty} \phi_m^\lambda(q) \alpha_m \int_{|z| \leq \sigma} |u_m^\lambda(z)|^2 \rho_\lambda(|z|^2) dx dy. \end{aligned}$$

The right-hand side can be regarded as the Fourier expansion for the left-hand side. Hence,

$$\begin{aligned} & \int_{-1}^1 \left| (A_\lambda^{-1}f)(q) - \int_{|z| \leq \sigma} \overline{A_\lambda(z, q)} f(z) \rho_\lambda(|z|^2) dx dy \right|^2 (1-q^2)^{\lambda-1/2} dq \\ &= \sum_{m=0}^\infty |\alpha_m - \alpha_m M_m^\lambda(\sigma)|^2 \\ &= \sum_{m=0}^\infty |1 - M_m^\lambda(\sigma)|^2 |\alpha_m|^2, \end{aligned}$$

where

$$M_m^\lambda(\sigma) = \int_{|z| \leq \sigma} |u_m^\lambda(z)|^2 \rho_\lambda(|z|^2) dx dy .$$

Therefore, we obtain the assertion.

q.e.d.

REMARK. The Gegenbauer polynomials for  $\lambda=1/2$  coincide with the Legendre polynomials  $P_m$ ,  $m=0, 1, 2, \dots$ . Hence if we put  $\lambda=1/2$  in Theorem 1, we obtain the desired result for the Legendre polynomials. In particular, we remark that in this case  $\rho_{1/2}(t) = t^{-1/2}$ .

§ 2. Some application to spherical harmonics.

We consider some application of Theorem 1 to the Funk-Hecke formula.

For a fixed  $n \geq 3$ , let  $S^{n-1}$  denote the surface of the unit sphere in  $R^n$ ,  $d\sigma$  the element of surface area on  $S^{n-1}$  and  $\omega_n$  the total surface area. Let finally  $H_k$  be the space of all spherical harmonics of order  $k$  on  $S^{n-1}$ .

Now the Funk-Hecke formula is given in the following theorem.

THEOREM 2 (Funk-Hecke). Suppose that  $\phi$  is an element of  $L^1((-1, 1), (1-q^2)^{(n-3)/2})$ . Then, for  $S_k \in H_k$  and  $\omega \in S^{n-1}$ ,

$$\begin{aligned} & \int_{S^{n-1}} \phi((\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ &= \omega_{n-1} S_k(\omega) \int_{-1}^1 \phi(q) P(k, q) (1-q^2)^{(n-3)/2} dq, \end{aligned}$$

where

$$P(k, q) = \frac{k! \Gamma(n-2)}{\Gamma(k+n-2)} C_k^{(n-2)/2}(q) .$$

First of all, we let  $r$  be an element of the open interval  $(-1, 1)$  and set  $\phi(q) = A_{(n-2)/2}(r, q)$  in the Funk-Hecke formula. Using the formula (4)  $u_m^\lambda = A_\lambda \phi_m^\lambda$  with  $\lambda = (n-2)/2$ , we then obtain the following:

$$\begin{aligned} & \int_{S^{n-1}} A_{(n-2)/2}(r, (\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ &= 2^{-(n-3)/2} \omega_{n-1} \frac{\Gamma(n-2)}{\Gamma((n-2)/2)} r^k S_k(\omega). \end{aligned} \quad (5)$$

On the other hand, the integral kernel  $A_{(n-2)/2}(r, q)$  and the Poisson kernel

$$p_r(q) = \frac{1-r^2}{(1-2rq+r^2)^{n/2}}$$

differ only by a constant. So we obtain

$$r^k S_k(\omega) = \omega_n^{-1} \int_{S^{n-1}} p_r((\omega, \tau)) S_k(\tau) d\sigma(\tau). \quad (6)$$

This equation is well-known as the Poisson integral. Thus we conclude that the equation (6) is only a special case of the Funk-Hecke formula.

Conversely, it is possible to obtain the Funk-Hecke formula from the equation (6). Since the element of surface area  $d\sigma$  is invariant under the orthogonal group  $O(n)$ , we obtain that

$$\begin{aligned} & \int_{S^{n-1}} \phi_m^\lambda((\omega, \tau)) \phi_l^\lambda((\omega, \tau)) d\sigma(\tau) \\ &= \omega_{n-1} \int_{-1}^1 \phi_m^\lambda(q) \phi_l^\lambda(q) (1-q^2)^{\lambda-1/2} dq \\ &= \omega_{n-1} \delta_{ml} \end{aligned}$$

with  $\lambda = (n-2)/2$  and  $\omega \in S^{n-1}$ . Therefore, for  $S_k \in H_k$ ,

$$\begin{aligned} & \sum_{m=0}^{\infty} u_m^\lambda(r) \int_{S^{n-1}} \phi_m^\lambda((\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ &= \int_{S^{n-1}} \sum_{m=0}^{\infty} u_m^\lambda(r) \phi_m^\lambda((\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ &= \int_{S^{n-1}} A_\lambda(r, (\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ &= 2^{-(\lambda-1/2)} \omega_{n-1} \frac{\Gamma(2\lambda)}{\Gamma(\lambda)} r^k S_k(\omega) \end{aligned}$$

with  $\lambda = (n-2)/2$  and  $-1 < r < 1$ . This implies

$$\begin{aligned} u_m^\lambda(r) \int_{S^{n-1}} \phi_m^\lambda((\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ = \delta_{mk} 2^{-(\lambda-1/2)} \omega_{n-1} \frac{\Gamma(2\lambda)}{\Gamma(\lambda)} r^k S_k(\omega), \end{aligned}$$

which is equivalent to the following:

$$\begin{aligned} \int_{S^{n-1}} \phi_m^\lambda((\omega, \tau)) S_k(\tau) d\sigma(\tau) \\ = \omega_{n-1} S_k(\omega) \int_{-1}^1 \phi_m^\lambda(q) P(k, q) (1-q^2)^{\lambda-1/2} dq. \end{aligned}$$

Hence, we obtain the Funk-Hecke formula for  $\phi_m^\lambda$  with  $\lambda = (n-2)/2$ . It follows from this result that the Funk-Hecke formula is true for any  $\phi$ .

### References

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