

## An uncertainty principle on Sturm-Liouville hypergroups

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**Abstract:** As an analogue of the classical uncertainty inequality on the Euclidean space, we shall obtain a generalization on the Sturm-Liouville hypergroups  $(\mathbf{R}_+, *(A))$ . Especially, we shall obtain a condition on  $A$  under which the discrete part of the Plancherel formula vanishes.

**Key words:** uncertainty principle; Sturm-Liouville; hypergroup.

**1. Sturm-Liouville hypergroups.** Sturm-Liouville hypergroups are a class of one-dimensional hypergroups on  $\mathbf{R}_+ = [0, \infty)$  with the convolution structure related to the second order differential operators

$$L = \frac{d^2}{dx^2} + \frac{A'(x)}{A(x)} \frac{d}{dx},$$

where  $A$  satisfies the following conditions (see [1,2]):

- (1)  $A > 0$  on  $\mathbf{R}_+^* = (0, \infty)$ , and is in  $C^2(\mathbf{R}_+^*)$ ,
- (2) on a neighborhood of 0,

$$\frac{A'(x)}{A(x)} = \frac{2\alpha + 1}{x} + B(x),$$

where  $\alpha \geq -\frac{1}{2}$  and (a) if  $\alpha > 0$ ,  $B$  and  $B'$  are integrable, (b) if  $\alpha = 0$ ,  $\log xB$  and  $x \log xB'$  are integrable, (c) if  $-\frac{1}{2} < \alpha < 0$ ,  $x^{2\alpha}B$  and  $x^{2\alpha+1}B'$  are integrable, (d) if  $\alpha = -\frac{1}{2}$ ,  $B'$  is integrable,

- (3)  $\frac{A'}{A} \geq 0$  on  $\mathbf{R}_+^*$  and  $\lim_{x \rightarrow \infty} \frac{A'(x)}{A(x)} = 2\rho$ ,
- (4)  $\frac{1}{2} \left(\frac{A'}{A}\right)' + \frac{1}{4} \left(\frac{A'}{A}\right)^2 - \rho^2$  is integrable at  $\infty$ .

Since  $A'/A = (\log A)'$ , (3) implies that  $A$  is increasing, and thus,  $A(0) < \infty$ . Under the conditions (1) to (3), the second order differential equation:  $Lu + (\lambda^2 + \rho^2)u = 0$ ,  $\lambda \in \mathbf{C}$ , has a unique solution satisfying  $u(0) = 1$ ,  $u'(0) = 0$ , which we denote by  $\phi_\lambda$ . Furthermore, under (4), if  $\Im \lambda \geq 0$ , then there exists another solution  $\psi_\lambda(x)$ , which behaves as

$\sqrt{\pi/2} \sqrt{\lambda x} H_\alpha^{(1)}$  at  $\infty$ , where  $H_\alpha^{(1)}$  is the Hankel function. Similarly, we have  $\psi_\lambda^-(x)$  for  $\Im \lambda \leq 0$ , and for  $\lambda \in \mathbf{R}_+^*$ , there exists  $C(\lambda) \in \mathbf{C}$  such that  $\phi_\lambda(x) = C(\lambda)\psi_\lambda(x) + \overline{C(\lambda)}\psi_\lambda^-(x)$ .

Let  $C_{c,e}^\infty(\mathbf{R})$  denote the set of  $C^\infty$  even functions  $f$  on  $\mathbf{R}$ . For  $f \in C_{c,e}^\infty(\mathbf{R})$  the Fourier transform  $\hat{f}$  is defined by

$$\hat{f}(\lambda) = \int_0^\infty f(x)\phi_\lambda(x)A(x)dx.$$

Then the inverse transform is given as

$$f(x) = \sum_{\Lambda \in D} \pi_\Lambda \hat{f}(\Lambda) \phi_\Lambda(x) + \frac{1}{2\pi} \int_0^\infty \hat{f}(\lambda) \phi_\lambda(x) \frac{d\lambda}{|C(\lambda)|^2},$$

where  $D$  is a finite set in the interval  $i(0, \rho)$  and  $\pi_\Lambda = \|\phi_\Lambda\|_{L^2(\mathbf{R}_+, A dx)}^{-2}$ . We denote this decomposition as

$$f = {}^\circ f + f_P$$

and we call  $f_P$  and  ${}^\circ f$  the principal part and the discrete part of  $f$  respectively. We denote by  $\mathbf{F}(\nu) = (F(\lambda), \{a_\Lambda\})$  a function on  $\mathbf{R}_+ \cup D$  defined by

$$\mathbf{F}(\nu) = \begin{cases} F(\lambda) & \text{if } \nu = \lambda \in \mathbf{R}_+ \\ a_\Lambda & \text{if } \nu = \Lambda \in D. \end{cases}$$

We put  $\overline{\mathbf{F}}(\nu) = (\overline{F(\lambda)}, \{\overline{a_\Lambda}\})$  and define the product of  $\mathbf{F}(\nu) = (F(\lambda), \{a_\Lambda\})$  and  $\mathbf{G}(\nu) = (G(\lambda), \{b_\Lambda\})$  as

$$(\mathbf{F}\mathbf{G})(\nu) = (F(\lambda)G(\lambda), \{a_\Lambda b_\Lambda\}).$$

Let  $d\nu$  denote the measure on  $\mathbf{R}_+ \cup D$  defined by

$$\int_{\mathbf{R}_+ \cup D} \mathbf{F}(\nu) d\nu = \sum_{\Lambda \in D} \pi_\Lambda a_\Lambda + \frac{1}{2\pi} \int_0^\infty F(\lambda) |C(\lambda)|^{-2} d\lambda.$$

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For  $f \in C_{c,e}^\infty(\mathbf{R})$ , we put

$$\hat{\mathbf{f}}(\nu) = (\hat{f}(\lambda), \{\hat{f}(\Lambda)\}).$$

Then the Parseval formula on  $C_{c,e}^\infty(\mathbf{R})$  can be stated as follows: For  $f, g \in C_{c,e}^\infty(\mathbf{R})$

$$(5) \quad \int_0^\infty f(x)\overline{g(x)}A(x)dx = \int_{\mathbf{R}_+ \cup D} \hat{\mathbf{f}}(\nu)\overline{\hat{\mathbf{g}}(\nu)}d\nu.$$

The map  $f \rightarrow \hat{\mathbf{f}}$ ,  $f \in C_{c,e}^\infty(\mathbf{R})$ , is extended to an isometry between  $L^2(A) = L^2(\mathbf{R}_+, A(x)dx)$  and  $L^2(\nu) = L^2(\mathbf{R}_+ \cup D, d\nu)$ . Actually, each function  $f$  in  $L^2(A)$  is of the form

$$\begin{aligned} f(x) &= \sum_{\Lambda \in D} \pi_\Lambda \hat{f}_\Lambda \phi_\Lambda(x) \\ &\quad + \frac{1}{2\pi} \int_0^\infty \hat{f}(\lambda) \phi_\lambda(x) |C(\lambda)|^{-2} d\lambda \\ &= {}^\circ f + f_P \end{aligned}$$

and their  $L^2$ -norms are given as

$$\int_0^\infty |{}^\circ f(x)|^2 A(x)dx = \sum_{\Lambda \in D} \pi_\Lambda |\hat{f}_\Lambda|^2,$$

$$\int_0^\infty |f_P(x)|^2 A(x) = \frac{1}{2\pi} \int_0^\infty |\hat{f}_P(\lambda)|^2 |C(\lambda)|^{-2} d\lambda.$$

Therefore, if we define  $\hat{\mathbf{f}}(\nu) = (\hat{f}(\lambda), \{\hat{f}_\Lambda\})$ , then  $\|f\|_{L^2(A)} = \|\hat{\mathbf{f}}\|_{L^2(\nu)}$  holds. In particular, if  $f \in C_{c,e}^\infty(\mathbf{R})$ , then  $\hat{f}_\Lambda = \hat{f}(\Lambda)$  for all  $\Lambda \in D$ .

**2. Uncertainty inequality.** We retain the notations in the previous sections. We put for  $x \in \mathbf{R}_+$ ,

$$(6) \quad a(x) = \int_0^x A(t)dt \text{ and } v(x) = \frac{a(x)}{A(x)}$$

and for  $\lambda \in \mathbf{C}$ ,

$$w(\lambda) = (\lambda^2 + \rho^2)^{1/2}.$$

**Theorem 2.1.** For all  $f \in L^1(A) \cap L^2(A)$ ,

$$(7) \quad \|fv\|_{L^2(A)}^2 \int_{\mathbf{R}_+ \cup D} |\hat{\mathbf{f}}(\nu)|^2 w(\nu)^2 d\nu \geq \frac{1}{4} \|f\|_{L^2(A)}^4,$$

where the equality holds if and only if  $f$  is of the form

$$f(x) = ce^{\gamma \int_0^x v(t)dt}$$

for some  $c, \gamma \in \mathbf{C}$  and  $\Re\gamma < 0$ .

*Proof.* Without loss of generality we may suppose that  $f \in C_{c,e}^\infty(\mathbf{R})$ . Since  $(-Lf)^\wedge(\lambda) = \hat{f}(\lambda)(\lambda^2 + \rho^2) = \hat{f}(\lambda)w(\lambda)^2$  and  $w(\lambda)$  is positive on  $\mathbf{R}_+ \cup D$ , the Parseval formula (5) yields that

$$\begin{aligned} \int_{\mathbf{R}_+ \cup D} |\hat{\mathbf{f}}(\nu)|^2 w(\nu)^2 d\nu &= \int_0^\infty (-Lf)(x)\overline{f(x)}A(x)dx \\ &= \int_0^\infty |f'(x)|^2 A(x)dx. \end{aligned}$$

Hence it follows that

$$\begin{aligned} \int_0^\infty |f(x)|^2 v(x)^2 A(x)dx &\int_{\mathbf{R}_+ \cup D} |\hat{\mathbf{f}}(\nu)|^2 w(\nu)^2 d\nu \\ &= \int_0^\infty |f(x)|^2 v(x)^2 A(x)dx \int_0^\infty |f'(x)|^2 A(x)dx \\ &\geq \left( \int_0^\infty \Re(f(x)f'(x))v(x)A(x)dx \right)^2 \\ &= \frac{1}{4} \left( \int_0^\infty (|f(x)|^2)' a(x)dx \right)^2 \\ &= \frac{1}{4} \left( \int_0^\infty |f(x)|^2 A(x)dx \right)^2. \end{aligned}$$

Here we used the fact that  $a' = A$  (see (6)). Clearly, the equality holds if and only if  $fv = cf'$  for some  $c \in \mathbf{C}$ , that is,  $f'/f = c^{-1}v$ . This means that  $\log(f) = c^{-1} \int_0^x v(t)dt + C$  and thus, the desired result follows.  $\square$

**Remark 2.2.** When  $(\mathbf{R}_+, *(A))$  is the Bessel-Kingman hypergroup, the equality holds for  $e^{\gamma x^2}$ ,  $\Re\gamma < 0$ . However, when it is the Jacobi hypergroup, each function satisfying the equality has an exponential decay  $e^{\gamma x}$ .

Since  $w^2(\lambda) = \lambda^2 + \rho^2$ , (7) can be rewritten as follows:

**Corollary 2.3.** Let  $f$  be the same as in Theorem 2.1.

$$(8) \quad \|fv\|_{L^2(A)}^2 \int_{\mathbf{R}_+ \cup D} |\hat{\mathbf{f}}(\nu)|^2 \nu^2 d\nu \geq \frac{1}{4} \|f\|_{L^2(A)}^2 \int_0^\infty |f(x)|^2 (1 - 4\rho^2 v(x)^2) A(x)dx.$$

**3. Vanishing condition of the discrete part.** We shall prove that under the assumption:

$$(9) \quad 0 \leq v(x) \leq \frac{1}{2\rho},$$

it follows that  $D = \emptyset$ . We suppose that  $D \neq \emptyset$  and we take  $f = \pi_\Lambda \phi_\Lambda$ ,  $\Lambda \in D$ . Then, since  $\hat{\mathbf{f}}(\nu) = 1$  if  $\nu = \Lambda$  and 0 otherwise, it follows from (8) that

$$\begin{aligned} \|fv\|_{L^2(A)}^2 \pi_\Lambda \Lambda^2 &\geq \frac{1}{4} \|f\|_{L^2(A)}^2 \int_0^\infty |f(x)|^2 (1 - 4\rho^2 v(x)^2) A(x)dx. \end{aligned}$$

Here we recall that  $\Lambda^2 < 0$ , because  $D \subset i(0, \rho)$  and  $1 - 4\rho^2 v(x)^2 \geq 0$  by (9). This is contradiction. Therefore, we obtain the following

**Theorem 3.1.** *If  $0 \leq v \leq \frac{1}{2\rho}$ , then  $D = \emptyset$ .*

For example, if  $A$  satisfies the inequality:

$$(10) \quad a(x)A'(x) = \int_0^x A(x)dx \cdot A'(x) \leq A^2(x),$$

then  $A$  satisfies (9). Actually, (10) implies

$$v'(x) = \frac{A^2(x) - a(x)A'(x)}{A^2(x)} \geq 0.$$

Hence  $v$  is increasing on  $\mathbf{R}_+$  and  $v(x) = a(x)/A(x) \leq A(x)/A'(x)$  because  $A/A' > 0$  by (3). Then it follows from (3) that  $A$  satisfies (9).

**Corollary 3.2.** *If  $A$  satisfies the inequality (10), then  $D = \emptyset$ .*

**Remark 3.3.** It is well-known that  $D = \emptyset$  for Chébli-Triméche hypergroups where  $A'/A$  is decreasing and (4) is not required (cf. [1]). This fact easily follows from our argument. Since  $A/A'$  is increasing and  $0 \leq A/A' \leq 1/2\rho$  by (3), we see that  $a \leq A/2\rho$  by integration and thus, (9) holds.

Hence  $D = \emptyset$  by Theorem 3.1.

**4. Uncertainty principle.** We suppose that  $D = \emptyset$ . Then (8) is of the form:

$$\begin{aligned} \|fv\|_{L^2(A)}^2 & \frac{1}{2\pi} \int_0^\infty |\hat{f}(\lambda)|^2 \lambda^2 \frac{d\lambda}{|C(\lambda)|^2} \\ & \geq \frac{1}{4} \|f\|_{L^2(A)}^2 \int_0^\infty |f(x)|^2 (1 - 4\rho^2 v(x)^2) A(x) dx. \end{aligned}$$

Since  $v$  is increasing,  $v(0) = 0$ , and  $1 - 4\rho^2 v(x)^2 \geq 0$  by (9), it follows that  $f$  and  $\hat{f}$  both cannot be concentrated around the origin.

In general, if  $D \neq \emptyset$ , then we must pay attention to the discrete part of  $f$  to consider uncertainty principles. We refer to [3] for the Jacobi hypergroups.

## References

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