

## THEORY OF THE SPHERICALLY SYMMETRIC SPACE-TIMES. IV. CONFORMAL TRANSFORMATIONS<sup>1)</sup>

By

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(Received March, 20, 1952)

### § 1. Spherically symmetric conformal transformation.

In the previous papers we have obtained some properties of s. s. space-times.<sup>1)</sup> In the present paper, further, we shall study conformal transformations which transform a given s. s. space-time into s. s. one again.

A conformal transformation of a Riemannian space is given by

$$C(v): \quad g_{ij}^* = e^{2v} g_{ij}, \quad (v = v(x^k)). \quad (1.1)$$

If  $g_{ij}$  gives a s. s. space-time  $S_0$ , the space-time defined by  $g_{ij}^*$  is not necessarily s. s., but it is evident that when  $v$  is s. s. with respect to a c. s. ( $K$ ) of  $S_0$  i. e. when

$$v_i = -(\alpha^s v_s) \alpha_i + (\beta^s v_s) \beta_i, \quad (v_i \equiv \nabla_i v), \quad (1.2)$$

holds for ( $K$ ), the resulting space-time becomes s. s. again. We shall call such a transformation a *spherically symmetric conformal transformation* (or s. s. c. t. in short) with respect to the  $S_0$ . If (1.1) is s. s. with respect to  $S_0$  its inverse conformal transformation is also s. s. with respect to  $S_0^*$  defined by  $g_{ij}^*$ . If a scalar  $v$  is s. s. with respect to a c. s. it is also s. s. with respect to any other c. s. to within  $m$ -transformation at most.<sup>2)</sup> Hence we know that any  $v$  obtained from  $v$  of a s. s. c. t. by any  $m$ -transformation defines a s. s. c. t. again. In this paper we shall obtain the transformation equations of the c. s. corresponding to a s. s. c. t. to within  $\varepsilon$ -transformation and then study the general form of  $C(v)$  which preserves the spherical symmetry of  $S_0$ .

### § 2. Preliminary formulae.

Transformation equations for Christoffel symbols and curvature tensor corresponding to (1.1) are given by<sup>3)</sup>

$$\{i_j\}^* = \{i_j\} + \delta_i^k v_j + \delta_j^k v_i - g_{ij} g^{km} v_m, \quad (2.1)$$

$$e^{-2v} K_{hi,jk}^* = K_{hi,jk} + 4g_{(h(k} v_{i)j)} + 2p g_{(h(k} g_{i)j)}, \quad (2.2)$$

$$\text{where} \quad v_{ij} = \nabla_i \nabla_j v - v_i v_j, \quad p = g^{ij} v_i v_j. \quad (2.3)$$

Let  $C(v)$  be an arbitrary s. s. c. t. If we take any s. s. coordinate system of  $S_0$ , we have

$$ds^2 = -A(r, t) dr^2 - B(r, t) (d\theta^2 + \sin^2 \theta d\phi^2) + C(r, t) dt^2, \quad (2.4)$$

and  $v$  becomes a function of  $r$  and  $t$  by a suitable  $m$ -transformation at most. Then from (2.1), (2.2) and (2.3), we have

$$v_1 = v', \quad v_4 = \dot{v}, \quad v_2 = v_3 = 0, \quad (2.5)$$

$$v_{11} = v'' - (A'/2A)v' - (\dot{A}/2C)\dot{v} - v'^2, \quad \text{etc.}, \quad (2.6)$$

$$\{11\}^* = A'/2A + v', \quad \{11\}^4 = \dot{A}/2C + (A/C)\dot{v}, \quad \text{etc.}, \quad (2.7)$$

$$\begin{cases} \alpha^* = e^{-2v}(\alpha + v_{11}/A + v_{22}/B - p), & \beta = e^{-2v}(\beta + v_{22}/B - v_{44}/C - p), \\ \gamma^* = e^{-2v}(\gamma + v_{14}/C), & \xi^* = e^{-2v}(\xi + v_{11}/A - v_{44}/C - p), \\ \eta^* = e^{-2v}(\eta + 2v_{22}/B - p); & p = -v'^2/A + \dot{v}^2/C. \end{cases} \quad (2.8)$$

Using the formulae concerning c. s. given in (I) and the relation

$$\begin{aligned} (g^{st} - h^{st})v_{st} &= -2v_{22}/B, & g^{st}v_{st} &= \square v - p, \\ (h_{ij} &\equiv -\alpha_i \alpha_j + \beta_i \beta_j, & \square &\equiv g^{st} \nabla_s \nabla_t), \end{aligned} \quad (2.9)$$

we can prove:

**Theorem [2.1]** *Let  $v$  be s. s. with respect to a  $(K)$  of an  $S_0$ , then for the s. s. c. t.  $C(v)$  we have*

$$\rho^1{}^* = e^{-2v} \rho^1, \quad F^* = F - v, \quad (2.10)$$

$$\rho^4{}^* = e^{-2v} \{\rho^4 - 2(g^{st} - h^{st})v_{st} - 2p\} = e^{-2v}(\rho^4 - 2\square v + 2h^{st}v_{st}), \quad (2.11)$$

as the transformation formulae for  $\rho^1$ ,  $F$  and  $\rho^4$ . When  $S_0$  admits  $\omega$ -transformation (2.11) holds for all c. s. obtained from  $(K)$  by any  $\omega$ -transformation. Furthermore the formulae holds for all c. s. provided that  $v$  undergoes the same  $m$ -transformation as  $(K)$ .

We shall call the property stated at the end of the theorem the *invariancy under  $m$ -transformation* for brevity's sake. It is to be noted that the tensor  $h_{ij}$  is invariant under  $\omega$ -transformation.

§ 3.  $C(v)$  which transforms  $S_a$  into  $S_a^*$ .

Theorem [3.1] *Let  $v$  be s. s. with respect to a c. s. ( $K$ ) of an  $S_a$ , then a necessary and sufficient condition that a s. s. c. t.  $C(v)$  transform  $S_a$  into an  $S_a^*$  again is given by*

$$P \equiv \alpha^s \beta^t v_{st} = 0, \quad Q \equiv (\alpha^s \alpha^t + \beta^s \beta^t) v_{st} = 0. \tag{3.1}$$

*And this condition is invariant under  $m$ - and  $\omega$ -transformations.*

Proof: Taking the standard coordinate system for ( $K$ ) and using the theorems concerning c. s. obtained in (I) we can prove that the condition that  $C(v)$  transform  $S_a$  into  $S_a^*$  is given by ( $v_{14}=0, v_{11}/A+v_{44}/C=0$ ) which is equivalent to (3.1) by virtue of (2.6). Further if  $(\alpha_i', \beta_i')$  is any pair of c. v. obtained from  $(\alpha_i, \beta_i)$  by an  $\omega$ -transformation, we have

$$\begin{cases} P' \equiv \alpha'^s \beta'^t v_{st} = P \cosh 2\omega + \frac{1}{2} Q \sinh 2\omega, \\ Q' \equiv (\alpha'^s \alpha'^t + \beta'^s \beta'^t) v_{st} = 2P \sinh 2\omega + Q \cosh 2\omega, \end{cases} \tag{3.2}$$

from which the latter part of the theorem is obvious.

We shall denote such a s. s. c. t. as stated above by  $C_{aa}(v)$ . Then we easily obtain:

Theorem [3.2] *In  $C_{aa}(v)$ , the transformation equation for  $\overset{2}{\rho}(=\overset{3}{\rho})$  is given by*

$$\overset{2}{\rho}^* = \overset{3}{\rho}^* = e^{-2v} \{ \overset{2}{\rho} + 2(g^{st} - 2h^{st}) v_{st} \}, \tag{3.3}$$

where  $(\alpha_i, \beta_i)$  is any pair of c. v. of the  $S_a$ .

(3.3) coincides with the transformation equation for  $M$  given in § 5. Theorem [3.3] *In  $C_{aa}(v)$ , we can establish one to one correspondence between  $(\alpha_i, \beta_i; \sigma, \bar{\sigma}; \kappa, \bar{\kappa})$  of  $S_a$  and  $(\alpha_i^*, \beta_i^*; \sigma^*, \dots)$  of  $S_a^*$  by the following relations:*

$$\begin{cases} \alpha_i^* = e^v \alpha_i, \quad \beta_i^* = e^v \beta_i, \quad i. e. \quad \alpha^{*t} = e^{-v} \alpha^t, \quad \beta^{*t} = e^{-v} \beta^t; \\ \sigma^* = e^{-v} (\sigma - \beta^s v_s), \quad \bar{\sigma}^* = e^{-v} (\bar{\sigma} + \alpha v_s); \\ \kappa^* = e^{-v} (\kappa + \alpha^s v_s), \quad \bar{\kappa}^* = e^{-v} (\bar{\kappa} + \beta^s v_s), \end{cases} \tag{3.4}$$

*to within an  $m$ -transformation.*

Both  $S_a$  and  $S_a^*$  admit  $\omega$ -transformation of c. s. and if we use the formulae concerning this transformation we can easily prove the theorem. When both  $S_a$  and  $S_a^*$  are neither [A] nor [B], their c. s. are invariant under  $m$ -transformation,<sup>2)</sup> hence we obtain one to one correspondence between both c. s. Of course another way of correspondence may exist.

§ 4.  $C(v)$  which transforms  $S_a$  into  $S_b^*$ .

From [3.1] it follows that a necessary and sufficient condition that the  $S_b^*$  obtained by a s. s. c. t. from an  $S_a$  be  $S_b^*$  is that (3.1) be not satisfied. We shall denote such a  $C(v)$  by  $C_{ab}(v)$ . In this section we shall give transformation equations for c. s. under  $C_{ab}(v)$ . Here it is to be noticed that  $S_a$  admits  $\omega$ -transformation of c. s. though it is not the case for  $S_b^*$ .

Theorem [4.1] *In  $C_{ab}(v)$ , it holds that*

$$\begin{cases} \rho^{2*} = e^{-2v} \{ \rho^2 + 2(g^{st} - 2h^{st})v_{st} + 2e\sqrt{(P^2 - 4Q^2)} \}, \\ \rho^{3*} = e^{-2v} \{ \rho^2 + 2( \quad , \quad )v_{st} - 2e\sqrt{( \quad , \quad )} \}, \end{cases} \quad (4.1)$$

$$\alpha_i^* = e^v (\alpha_i \cosh \Omega + \beta_i \sinh \Omega), \quad \beta_i^* = e^v (\alpha_i \sinh \Omega + \beta_i \cosh \Omega), \quad (4.2)$$

where  $\tanh 2\Omega = -2Q/P$  and  $e = \pm 1$  is chosen so as to satisfy  $eP \geq 0$ . Furthermore if  $(\alpha'_i, \beta'_i)$  is obtained from  $(\alpha_i, \beta_i)$  by an  $\omega$ -transformation, it holds that

$$\begin{cases} \rho^{2*} = e^{-2v} \{ \rho^2 + 2(g^{st} - h'^{st})v_{st} + 2e'\sqrt{(P'^2 - 4Q'^2)} \}, \\ \rho^{3*} = e^{-2v} \{ \rho^2 + 2( \quad , \quad )v_{st} - 2e'\sqrt{( \quad , \quad )} \}, \end{cases} \quad (4.1')$$

$$\begin{cases} \alpha_i^* = e^v \{ \alpha'_i \cosh (\Omega' - \omega) + \beta'_i \sinh (\Omega' - \omega) \}, \\ \beta_i^* = e^v \{ \alpha'_i \sinh (\Omega' - \omega) + \beta'_i \cosh (\Omega' - \omega) \}, \end{cases} \quad (4.2')$$

where  $\omega$  is the parameter of the  $\omega$ -transformation,  $e'$  is chosen so as to satisfy  $e'P' \geq 0$ , and

$$\tanh 2\Omega' = (P' \sinh 2\omega - 2Q' \cosh 2\omega) / (P' \cosh 2\omega - 2Q' \sinh 2\omega). \quad (4.3)$$

Of course (4.1), (4.2), (4.1') and (4.2') are invariant under  $m$ -transformation.

Proof: Using the relation<sup>4)</sup>

$$\alpha_i^* = e^v (\alpha_i \cosh \zeta^* + \beta_i \sinh \zeta^*), \quad \beta_i^* = e^v (\alpha_i \sinh \zeta^* + \beta_i \cosh \zeta^*), \quad (4.4)$$

$$\tanh 2\zeta^* = 2\sqrt{C/A} \gamma^* / (\alpha^* - \beta^*) = -2Q/P, \quad (4.5)$$

( $P = v_{11}/A + v_{44}/C$ ,  $Q = -v_{14}/\sqrt{AC}$ ), and putting  $\zeta = \Omega$ , we can obtain (4.2) and (4.2'). Then from (2.8), we have

$$M^* = e^{-2v} \{ M + 2(g^{st} - 2h^{st})v_{st} \}, \quad (M = \rho^2 = \rho^3); \quad N^* = 2e^{-2v} \sqrt{(P^2 - 4Q^2)}. \quad (4.6)$$

But  $h^{st}$  and  $P^2 - 4Q^2$  are invariant under  $\omega$ -transformation, so we have (4.1) and (4.1').

Theorem [4.2] In  $C_{ab}(v)$ , it holds that

$$\begin{cases} \sigma^* = e^{-v} \{ (\sigma - \alpha^s \Omega_s - \beta^s v_s) \cosh \Omega - (\bar{\sigma} + \beta^s \Omega_s + \alpha^s v_s) \sinh \Omega \}, \\ \bar{\sigma}^* = e^{-v} \{ -(\quad, \quad) \sinh \Omega + (\quad, \quad) \cosh \Omega \}; \end{cases} \quad (4.7)$$

$$\begin{cases} \kappa^* = e^{-v} \{ (\kappa + \alpha^s v_s) \cosh \Omega + (\bar{\kappa} + \beta^s v_s) \sinh \Omega \}, \\ \bar{\kappa}^* = e^{-v} \{ (\quad, \quad) \sinh \Omega + (\quad, \quad) \cosh \Omega \}, \end{cases} \quad (4.8)$$

$$\begin{cases} \sigma^* = e^{-v} [ \{ \sigma' - \alpha'^s (\Omega'_s - \omega_s) - \beta'^s v_s \} \cosh (\Omega' - \omega) - \{ \bar{\sigma} + \beta'^s (\Omega'_s - \omega_s) + \alpha'^s v_s \} \\ \qquad \qquad \qquad \sinh (\Omega' - \omega) ], \\ \bar{\sigma}^* = e^{-v} [ - \{ \quad, \quad \} \sinh (\Omega' - \omega) + \{ \quad, \quad \} \cosh (\Omega' - \omega) ]; \end{cases} \quad (4.7')$$

$$\begin{cases} \kappa^* = e^{-v} \{ (\kappa' + \alpha'^s v_s) \cosh (\Omega' - \omega) + (\bar{\kappa}' + \beta'^s v_s) \sinh (\Omega' - \omega) \}, \\ \bar{\kappa}^* = e^{-v} \{ (\quad, \quad) \sinh (\Omega' - \omega) + (\quad, \quad) \cosh (\Omega' - \omega) \}, \end{cases} \quad (4.8')$$

where  $\Omega_s = \partial_s \Omega$  and  $\omega_s = \partial_s \omega$ , corresponding to  $(\alpha_i, \beta_i)$  and  $(\alpha'_i, \beta'_i)$  respectively as in [4.1].

Proof: Using the relations<sup>4)</sup>

$$\sigma^* = \bar{M}^* \cosh \zeta - \bar{N}^* \sinh \zeta, \quad \bar{\sigma}^* = -\bar{M}^* \sinh \zeta + \bar{N}^* \cosh \zeta, \quad (4.9)$$

$$\begin{aligned} \bar{M}^* &= e^{-v} (\sigma - \alpha^s \zeta_s^* - \beta^s v_s) \\ &= e^{-v} \{ (\sigma' + \alpha'^s \omega_s - \alpha'^s \zeta_s^* - \beta'^s v_s) \cosh \omega + \dots \}, \quad \text{etc.}, \end{aligned} \quad (4.10)$$

and putting  $\zeta^* = \Omega$ , we can prove (4.7) and (4.7'). Similarly we have (4.8) and (4.8'). Of course the right hand sides of (4.7') and (4.8') are independent of  $\omega$ , and (4.7), ... (4.8') are invariant under  $m$ -transformation.

### § 5. $C(v)$ which transforms $S_b$ into $S_b^*$ .

As in the preceding sections we can easily prove:

Theorem [5.1] A necessary and sufficient condition that a s. s. c. t. be  $C_{ba}(v)$ , (i. e. a s. s. c. t. which transforms  $S_b$  into  $S_a^*$ ), is given by

$$(\rho^2 - \rho^3) + 4P = 0, \quad (\text{i. e. } N + 2P = 0), \quad \text{and} \quad Q = 0. \quad (5.1)$$

When this condition is satisfied for a c. s. ( $K$ ) of  $S_b$  it holds for all c. s. of the  $S_b$  with a proviso that  $v$  undergoes the same  $m$ -transformation as ( $K$ ).

By this theorem we know that an  $S_b$  is transformed into  $S_b^*$  by a s. s. c. t. when and only when (5.1) is not satisfied. In the following we shall consider the transformation law of the c. s. under such a  $C_{ba}(v)$ .

Using the formulae concerning c. s. of  $S_b$  given in (I) we can easily obtain:

Lemma [5.2] *In  $C_{bb}(v)$ , the following relation concerning the transformation law for  $\zeta$ , the quantity which appears in the formulae for c. s. of  $S_b$ , holds*

$$\zeta^* = \zeta + \Omega, \quad (5.2)$$

where 
$$\tanh 2\Omega = -2Q/\{(\rho^2 - \rho^3)/4 + P\}. \quad (5.3)$$

By using this lemma, we have

Theorem [5.3] *In  $C_{bb}(v)$ , we have the transformation equations (4.2), (4.7) and (4.8) for  $(\alpha_i, \beta_i)$ ,  $(\sigma, \bar{\sigma})$  and  $(\kappa, \bar{\kappa})$  respectively, with a proviso that  $\Omega$  is given by (5.3), and for  $\rho^2$  and  $\rho^3$ , it holds that*

$$\left\{ \begin{array}{l} \rho^{2*} = e^{-2v} \{2(g^{st} - h^{st})v_{st} + (\rho^2 + 4\alpha^s \alpha^t v_{st}) \cosh^2 \Omega - (\rho^3 - 4\beta^s \beta^t v_{st}) \sinh^2 \Omega + 4Q \sinh 2\Omega\}, \\ \rho^{3*} = e^{-2v} \{2( \quad \quad )v_{st} - ( \quad \quad ) \sinh^2 \Omega + ( \quad \quad ) \cosh^2 \Omega - 4Q \sinh 2\Omega\}, \end{array} \right. \quad (5.4)$$

where  $\Omega$  is given by (5.3) again. These equations are invariant under  $m$ -transformation of  $S_b$ .

Proof: We can prove the former part by using the formulae concerning c. s. of an  $S_b$  given in (I), and then by substituting this result into the identities (3.9) of (I) we can obtain (5.4). From this theorem we have:

Corollary [5.4] *In  $C_{bb}(v)$ , it holds that*

$$\begin{aligned} M^* &= e^{-2v} \{M + 2(g^{st} - 2h^{st})v_{st}\}, \quad N^* = e^{-2v} \{(N + 2P) \cosh 2\Omega + 4Q \sinh 2\Omega\}, \\ (N^*)^2 &= e^{-4v} \{(N + 2P)^2 - 16Q^2\}. \end{aligned} \quad (5.5)$$

## § 6. $C(v)$ which transforms $S_b$ into $S_a^*$ .

The condition to be satisfied by  $C_{ba}(v)$  is given by [5.1]. If this condition is satisfied  $\Omega$  defined by (5.3) is indeterminate and we can not determine  $\alpha_i^*$ ,  $\beta_i^*$ ,  $\sigma^*$ , ... from the formulae given in the last section. This fact is natural since the c. s. of an  $S_a^*$  admits  $\omega$ -transformation.

Theorem [6.1] *In  $C_{ba}(v)$ , it holds that*

$$\rho^{2*} = \rho^{3*} = e^{-2v} \{(\rho^2 + \rho^3)/2 + 2(g^{st} - 2h^{st})v_{st}\},$$

and this holds for all c. s. of  $S_b$ .  $(\alpha_i^*, \beta_i^*; \sigma^*, \dots)$  is given by

$$\left\{ \begin{array}{l} \alpha_i^* = e^v \alpha_i, \quad \beta_i^* = e^v \beta_i; \quad \sigma^* = e^{-v} (\sigma - \beta^s v_s), \quad \bar{\sigma}^* = e^{-v} (\bar{\sigma} + \alpha^s v_s), \\ \kappa^* = e^{-v} (\kappa + \alpha^s v_s), \quad \bar{\kappa}^* = e^{-v} (\bar{\kappa} + \beta^s v_s), \end{array} \right. \quad (6.2)$$

and those obtained from this by  $m$ - and  $\omega$ -transformations.

Proof: To prove the latter part of the theorem we have only to show that  $(\alpha_i^*, \dots)$  given by (6.2) gives a c. s. of  $S_0^*$ . Since the proof is easy but tedious, we shall omit it here.

By the discussions given in §§ 3, 4, 5 and 6, we have succeeded in expressing the transformation formulae for c. s. of an  $S_0$  under any s. s. c. t. in invariant form in terms of the c. s. of given  $S_0$  and the conformal factor  $v$ .

### § 7. Conformal transformation which preserves spherical symmetry of an $S_0$ .

Let  $S^*$  be the space-time obtained from an  $S_0$  by (1.1) whose  $v$  is any function of  $x^i$ . Obviously  $S^*$  is not necessarily s. s. but in some special cases both  $S_0$  and  $S^*$  can be s. s. notwithstanding that (1.1) is not s. s. In this section we shall give some theorems concerning this problem.

We shall denote by  $S_{20}$  the s. s. space-time whose line element can be brought into the form

$$ds^2 = -A(r, t) dr^2 + C(r, t) dt^2 - B(d\theta^2 + \sin^2 \theta d\phi^2), \quad (7.1)$$

where  $B = \text{const.}$  and  $\xi = \text{const.}$  Hence  $S_{20}$  is an  $S_{II}$  whose  $(r, t)$ -space is of constant curvature (i. e.  $\rho = \text{const.}$ ). Specially when  $S_{20}$  is conformally flat we have  $\xi = -\eta = 1/B$ . Therefore in the coordinate system in which  $A=1$ , we have  $\xi = (\sqrt{C})'/\sqrt{C} = -1/B$  and the general form of  $C$  is given by  $(ae^{mr} + be^{-mr})^2$  where  $a$  and  $b$  are arbitrary functions of  $t$  and  $m = 1/\sqrt{B}$ . Theorem [7.1] *A necessary and sufficient condition that an  $S_0$  be conformal to an  $S_{20}$  is given by  $\rho e^{-2F} = \text{const.}$*  (7.2)

Proof: (i) (7.2) is necessary: The condition for  $S_0$  to be conformally flat is  $\rho = 0$  i. e. a special form of (7.2), hence we have only to deal with  $S_0$  whose  $\rho \neq 0$ . Then we have

$$ds^2 = e^{2v} ds^{*2}, \quad ds^{*2} = -A^* dr^2 + C^* dt^2 - B^*(d\theta^2 + \sin^2 \theta d\phi^2), \quad (7.3)$$

where  $ds^{*2}$  gives an  $S_{20}$  and  $v = v(r, \theta, \phi, t)$ . By the same method as the one used in the next theorem we know that a c. s. of  $S_0$  must be either of the following two forms: (a)  $P_{14} \neq 0$  and other  $P_{ij} = 0$ , (b)  $P_{23} \neq 0$  and other  $P_{ij} = 0$  where  $P_{ij} = \alpha_{(i} \beta_{j)}$ . Then we can easily obtain (7.2). (ii) By taking standard coordinate system for c. s. of  $S_0$  and using (2.10) we can prove that  $ds^{*2} = ds^2/B(r, t)$  gives an  $S_{20}$ .

**Theorem [7.2]** *Let  $S_0$  be not conformal to any  $S_{20}$ . A necessary and sufficient condition that the  $S^*$  obtained by  $g_{ij}^* = e^{2v}g_{ij}$  from the  $S_0$  be s. s. again is that  $v$  be  $\frac{1}{2}$ s. s. with respect to a c. s. of the  $S_0$ . (i. e. the  $C(v)$  be s. s. c. t.).*

**Proof:** The sufficiency is evident. Next, let  $S^*$  be s. s. and take any standard coordinate system for  $g_{ij}$ . Using the first fundamental equation for c. s. (F<sub>1</sub>) and (2.2), we can prove that only either of the following two cases is possible<sup>5)</sup>: (i) When  $P_{14} \neq 0$  and other  $P_{ij} = 0$ . From this we have  $\alpha_2 = \alpha_3 = \beta_2 = \beta_3 = 0$  and further by using (F<sub>2</sub>), (F<sub>3</sub>)<sup>6)</sup> and (2.1) we can prove  $v_2 = v_3 = 0$  i. e. the  $C(v)$  is s. s. (ii) When  $P_{23} \neq 0$  and other  $P_{ij} = 0$ . In the same way we have  $\alpha_1 = \alpha_4 = \beta_1 = \beta_4 = 0$  and from (F<sub>2</sub>) we have  $e^{2v} = e^{2\bar{v}}/B(r, t)$  where  $\bar{v} = \bar{v}(\theta, \phi)$ . Then using (F<sub>1</sub>) and (2.2) we can prove that  $d\bar{s}^2$  defined by  $\bar{g}_{ij} = e^{-2\bar{v}}g_{ij}^*$  gives an  $S_{20}$  contrary to our assumption. Thus the theorem is proved.

From the proof of this theorem we know that when (7.2) holds there exist some  $C(v)$ 's which are not s. s. and yet transform  $S_0$  into s. s.  $S^*$  again. (The constant in (7.2) is invariant under s. s. c. t. but not necessarily so under other ones). Concerning this we have the theorem:

**Theorem [7.3]** *The general form of the conformal transformation  $g_{ij}^* = e^{2v}g_{ij}$ , where both  $g_{ij}^*$  and  $g_{ij}$  define [B] (i. e. Minkowski type  $S_0$ ) is given by the following two types: (a) The one whose  $v_{ij} = -pg_{ij}/2 \neq 0$ . In this case the  $C(v)$  is s. s. (b) The one whose  $v_{ij} = 0$ . In this case the  $C(v)$  is not s. s.*

**Proof:** Using the coordinate system in which  $ds^2 = -dx^2 - dy^2 - dz^2 + dt^2$ , we have  $e^{-v} = -mx^t x_i + n_i x^t + s$  where  $x^t = x_t$ , and  $m, n_i$ , and  $s$  are constants satisfying  $n_i n^i = -4ms, (n_i = n^i)$ . Then we can prove that when  $m \neq 0$  or  $m = 0$  we have (a) or (b) respectively. (In the case of (a) we must use  $m$ -transformation).

In connection with the above we shall add the following theorem.

**Theorem [7.4]** *The fundamental form of any conformally flat  $S_0$  is reducible to*

$$ds^2 = A\{-(dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2) + \lambda dt^2\}, \tag{7.4}$$

where  $A = A(r, t)$  and  $\lambda = \{\alpha(t) + r^2 b(t)\}^2$ . Namely any conformally flat  $S_0$  is given by a suitable choice of  $A(r, t)$ ,  $\alpha(t)$  and  $b(t)$ .

**Proof:** (2.4) is transformable into isotropic form

$$ds^2 = -A(r, t)(dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2) + C(r, t) dt^2, \tag{7.5}$$

using which we can easily obtain (7.4).



From this theorem it must be that by a suitable  $C(v)$  (7.4) is transformable into [B]. If we assume that  $v=v(r, t)$ , however, this is not necessarily the case. This will be seen if we look for the condition to be satisfied by  $\lambda$  in order that (7.4) may give [B].

When  $b=0$ , by taking  $A=\text{const.}$  or  $=(t^2-r^2)^2$  taking  $a=1$  by a transformation of  $t$ , we have [B]<sup>7)</sup>. Hence we assume that  $b\neq 0$  and by a transformation of  $t$  we take  $b=1$ . Calculating  $K_{ij}{}^m=0$  for (7.4) we can show that (7.4) can be [B] when and only when  $a=0$  or  $\pm 1$  or  $(2t+q)^{-2}$  where  $q$  is any const., and we can obtain the corresponding  $A(r, t)$ . Examples of the line element of [B] in these cases are given by

$$\text{When } a=0: \quad ds^2 = -r^4 dm^2 + dt^2, \quad (dm^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2). \quad (7.6)$$

$$\text{When } a=1: \quad ds^2 = -e^{4t} (1+r^2)^{-2} dm^2 + e^{4t} dt^2, \quad (7.7)$$

$$\text{or putting } e^{2t} = 2\bar{t}, \quad ds^2 = 4\bar{t}^2 (1+r^2)^{-2} dm^2 + d\bar{t}^2. \quad (7.8)$$

$$\text{When } a=-1: \quad ds^2 = 4t^2 (1-r^2)^{-2} dm^2 + dt^2.$$

$$\text{When } a=(2t+q)^{-2}: \quad ds^2 = (2t+q)^{-2} (r^2-a)^{-2} \{-dm^2 + ((r^2+a)^2 dt^2)\}. \quad (7.10)$$

(7.8) was obtained by the writer from another point of view.<sup>7)</sup>

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### Notes and References

1) This paper is a continuation of Journ. Math. Soc. Japan **3** (1952), 317; this Journal **16** (1952), 67; this Journal **16** (1952), 291. These are cited as (I), (II) and (III) respectively and the same notations as in these papers are used throughout the present paper.

2) (I) and (II).

3) L. P. Eisenhart, *Riemannian Geometry*, Princeton (1926), 89.

4) § 5 of (I).

5) Here  $(\alpha_t, \beta_t)$  denotes any pair of c. v. of  $S^*$ .

6) § 1 of (I).

7) H. Takeno, This Journal **11** (1941), 224.