

THE BEST CONSTANT IN A BMO-INEQUALITY FOR THE BEURLING-AHLFORS TRANSFORM

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Introduction. The Beurling–Ahlfors operator Tf for $f \in L^2(\mathbf{C})$ is defined by the following relation between Fourier transforms:

$$(Tf)^\wedge(\xi) = m(\xi)\hat{f}(\xi),$$

where $m(\xi) = (\xi/|\xi|)^2$ for every $\xi \in \mathbf{C} - \{0\}$.

Clearly T is a unitary operator on $L^2(\mathbf{C})$ commuting with translations and dilations. Another definition of T is the convolution formula

$$(Tf)(z) = -\text{p.v.} \frac{1}{\pi} \int \frac{f(t) d\sigma(t)}{(z-t)^2},$$

where p.v. means the principal value and $d\sigma(t)$ the Lebesgue measure in \mathbf{C} . This operator can be regarded as an analogue of the Hilbert transform in the complex plane with an even kernel.

The importance of the Beurling–Ahlfors operator to the elliptic equations ([3], [4], [13]) as well as to quasiconformal mappings in the plane lies in the fact that it changes the complex derivative $\partial_{\bar{z}}$ into ∂_z : in symbols,

$$(1) \quad T\left(\frac{\partial w}{\partial \bar{z}}\right) = \frac{\partial w}{\partial z}$$

for every w in the Sobolev space $\mathcal{W}_2^1(\mathbf{C})$. We shall appeal to this formula to evaluate Tf for some particular functions f .

As an operator of Calderon–Zygmund type, the Beurling–Ahlfors transform is bounded in $L^p(\mathbf{C})$ for all $1 < p < \infty$. This breaks down for $p = \infty$. However, in this limiting case T extends to a bounded operator from $L^\infty(\mathbf{C})$ into BMO-spaces [12].

Fix $1 \leq p < \infty$. A function $f \in L^p_{\text{loc}}(\mathbf{R}^n)$ is said to be of bounded mean oscillation (briefly, BMO_p) if

$$\|f\|_{\text{BMO}_p} = \sup_B \left(\int_B \left| f(x) - \int_B f(y) dy \right|^p dx \right)^{1/p} < \infty,$$

where the supremum is taken over all balls B in \mathbf{R}^n and

$$\int_B f(y) dy = \frac{1}{|B|} \int_B f(y) dy = f_B$$

is the average of f on B .

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All BMO_p spaces for $1 \leq p < \infty$ are essentially equivalent; namely,

$$\|f\|_{BMO_p} \leq C(p, q, n) \|f\|_{BMO_q}$$

for $1 \leq p, q < \infty$. Actually, by Hölder's inequality,

$$(2) \quad \|f\|_{BMO_p} \leq \|f\|_{BMO_q}$$

whenever $1 \leq p \leq q < \infty$.

The purpose of this paper is to identify the $L^\infty \rightarrow BMO_2$ norm of the Beurling-Ahlfors operator.

THEOREM 1. *We have*

$$(3) \quad \left(\int_B |Tf(z) - (Tf)_B|^2 d\sigma(z) \right)^{1/2} \leq 3$$

for every ball $B \subset \mathbb{C}$ and every function $f \in L^2(\mathbb{C})$ such that $|f(z)| \leq 1$, a.e. This inequality is sharp.

Interest in the BMO-inequalities is motivated by their relations with quasiconformal mappings. In this connection we mention the beautiful results of Reimann [11], who discovered that quasiconformal mappings are invariant transformations of independent variables for BMO-functions. We also refer to several deep results due to Jones [9] and Astala and Gehring [1], who characterized domains having the BMO-extension property.

1. Preliminaries. It clearly suffices to prove (3) when \mathbf{B} is the unit disk in \mathbb{C} . From now on we shall assume that

$$\mathbf{B} = \{z \in \mathbb{C}; |z| < 1\} \quad \text{and} \quad \Omega = \{t \in \mathbb{C}; |t| \geq 1\}.$$

The following formulas are worth recording. For every integer $k \geq 0$ set

$$(4) \quad \rho_k(z) = \bar{z}^k \chi_{\mathbf{B}}(z).$$

Then

$$(5) \quad T\rho_k(z) = -z^{-k-2} \chi_{\Omega}(z).$$

In order to see these we apply (1) to the functions

$$w_k(z) = \frac{1}{k+1} [\bar{z}^{k+1} \chi_{\mathbf{B}}(z) + z^{-k-1} \chi_{\Omega}(z)].$$

Based on formula (1) we also derive that, for

$$(6) \quad \rho(z) = \left(\frac{z}{|z|} \right)^2 \chi_{\mathbf{B}}(z),$$

$$(7) \quad T\rho(z) = (1 + \log|z|^2) \chi_{\mathbf{B}}(z).$$

In this case (1) applies to $w(z) = (z \log|z|^2) \chi_{\mathbf{B}}(z)$. Formulas (4) and (5) immediately imply the following.

COROLLARY 1. Let $\mathcal{H}^2(\mathbf{B})$ denote the space of functions square integrable and analytic in \mathbf{B} . Then, for every $h \in \mathcal{H}^2(\mathbf{B})$,

$$(8) \quad T(\chi_{\mathbf{B}} \bar{h})(z) = 0$$

for $z \in \mathbf{B}$.

2. **Certain orthogonality properties of T .** The proof of Theorem 1 depends on the orthogonality properties of T . We begin with the following fundamental identities:

$$(9) \quad \int f(z) Tg(z) d\sigma(z) = \int g(z) Tf(z) d\sigma(z),$$

$$(10) \quad \int Tf(z) \overline{Tg(z)} d\sigma(z) = \int f(z) \overline{g(z)} d\sigma(z),$$

for every $f, g \in L^2(\mathbf{C})$.

The sets \mathbf{B} and Ω induce naturally an orthogonal decomposition of $L^2(\mathbf{C})$,

$$L^2(\mathbf{C}) = L^2(\mathbf{B}) \oplus L^2(\Omega),$$

obtained by decomposing any function f into the sum $f = \chi_{\mathbf{B}} f + \chi_{\Omega} f$. This decomposition is not invariant under the Beurling-Ahlfors transform. However, we have the following lemma.

LEMMA 1. Let $b \in L^2(\mathbf{B})$ and $\omega \in L^2(\Omega)$. Then the functions Tb and $T\omega$ are orthogonal in $L^2(\mathbf{B})$ and in $L^2(\Omega)$, that is,

$$(11) \quad \int_{\mathbf{B}} Tb(z) \overline{T\omega(z)} d\sigma(z) = 0;$$

also,

$$\int_{\Omega} Tb(z) \overline{T\omega(z)} d\sigma(z) = 0.$$

Moreover,

$$(12) \quad \int_{\mathbf{B}} Tb(z) d\sigma(z) = 0.$$

Proof. In view of identity (9), the integral in (11) takes the form

$$\int (Tb) \overline{(\chi_{\mathbf{B}} T\omega)} = \int b T(\chi_{\mathbf{B}} \overline{T\omega}) = \int_{\mathbf{B}} b T(\chi_{\mathbf{B}} \overline{T\omega}).$$

Observe that $T\omega \in \mathcal{H}^2(\mathbf{B})$ and by Corollary 1 the function $T(\chi_{\mathbf{B}} \overline{T\omega})$ vanishes on \mathbf{B} . Hence the equality (11) follows. The second equality follows from (11) together with the fact that T is a unitary operator on $L^2(\mathbf{C})$ and the supports of b and ω are disjoint.

Now replace $T\omega$ in (11) by $\chi_{\mathbf{B}}$, which is obviously in $\mathcal{H}^2(\mathbf{B})$, and repeat the same arguments to obtain (12). □

COROLLARY 2. *Let $f \in L^2(\mathbf{C})$. Then*

$$(13) \quad \int_{\mathbf{B}} |Tf - (Tf)_{\mathbf{B}}|^2 = \int_{\mathbf{B}} |Tb|^2 + \int_{\mathbf{B}} |T\omega - (T\omega)_{\mathbf{B}}|^2,$$

where b and ω are $L^2(\mathbf{B})$ and $L^2(\Omega)$ components of f , respectively.

Notice that $|b| \leq \chi_{\mathbf{B}}$ and $|\omega| \leq \chi_{\Omega}$, whenever $|f| \leq 1$.

Proof. Identity (13) is a simple consequence of (11) and (12). □

Corollary 2 leads naturally to two independent extremal problems. We begin by considering the simpler one.

3. The estimate in $L^2(\mathbf{B}) \cap L^\infty(\mathbf{C})$.

LEMMA 2. *Let b be measurable and such that $|b(z)| \leq \chi_{\mathbf{B}}(z)$, a.e. Then*

$$(14) \quad \int_{\mathbf{B}} |Tb(z)|^2 d\sigma(z) \leq 1.$$

The inequality is sharp.

Proof. Inequality (14) is an immediate consequence of the fact that T is an isometry in $L^2(\mathbf{C})$:

$$\int_{\mathbf{B}} |Tb|^2 \leq \frac{1}{|\mathbf{B}|} \int |Tb|^2 = \frac{1}{|\mathbf{B}|} \int |b|^2 = \int_{\mathbf{B}} |b|^2 \leq 1.$$

This argument also shows how to achieve equality in (14). For this we require (i) $\chi_{\Omega}(z)Tb(z) = 0$ and (ii) $|b(z)| = \chi_{\mathbf{B}}(z)$. Certainly the function ρ defined by (6) satisfies these two requirements. This ends the proof of Lemma 2. □

4. The estimate in $L^2(\Omega) \cap L^\infty(\mathbf{C})$.

LEMMA 3. *Let $\omega \in L^2(\mathbf{C})$ be such that $|\omega(z)| \leq \chi_{\Omega}(z)$, a.e. Then*

$$(15) \quad \int_{\mathbf{B}} |T\omega(z) - (T\omega)_{\mathbf{B}}|^2 d\sigma(z) \leq 8.$$

The inequality is sharp.

For the extremal function see the remark after the proof.

Proof. Clearly $\chi_{\mathbf{B}}T\omega \in \mathcal{H}^2(\mathbf{B})$, and for $z \in \mathbf{B}$ we have the following Taylor expansion:

$$\begin{aligned} T\omega(z) &= -\frac{1}{\pi} \int_{\Omega} \frac{\omega(t) d\sigma(t)}{(z-t)^2} = -\frac{1}{\pi} \sum_{k=0}^{\infty} (k+1)z^k \int t^{-k-2} \omega(t) d\sigma(t) \\ &= \sum_{k=0}^{\infty} (k+1)c_k z^k, \end{aligned}$$

where

$$(16) \quad c_k = -\frac{1}{\pi} \int t^{-k-2} \omega(t) d\sigma(t)$$

for $k = 0, 1, 2, \dots$. Since $c_0 = T\omega(0) = (T\omega)_{\mathbf{B}}$,

$$(17) \quad T\omega(z) - (T\omega)_{\mathbf{B}} = \sum_{k=1}^{\infty} (k+1)c_k z^k$$

for $z \in \mathbf{B}$. This last series is also well defined when we drop the assumption $\omega \in L^2(\mathbf{C})$, because of the following uniform estimate:

$$|c_k| \leq \frac{1}{\pi} \int_{\Omega} |t|^{-k-2} d\sigma(t) \leq \frac{1}{\pi} \int_{\Omega} |t|^{-3} d\sigma(t) = 2$$

for $k = 1, 2, \dots$. Here we have

$$(18) \quad \int_{\Omega} |t|^{-3} d\sigma(t) = 2\pi \int_1^{\infty} r^{-2} dr = 2\pi.$$

Notice that

$$\int_{\mathbf{B}} |z|^{2k} d\sigma(z) = 2 \int_0^1 r^{2k+1} dr = \frac{1}{k+1}$$

for $k = 1, 2, 3, \dots$. This, together with the mutual orthogonality of the functions $z^k \chi_{\mathbf{B}}(z)$ and (17), gives

$$(19) \quad \int_{\mathbf{B}} |T\omega(z) - (T\omega)_{\mathbf{B}}|^2 d\sigma(z) = \sum_{k=1}^{\infty} (k+1)|c_k|^2,$$

where c_k are defined by (16).

Our goal now is to maximize the above series subject to the conditions $|\omega(z)| \leq \chi_{\Omega}(z)$, $\omega \in L^2(\mathbf{C})$. Variational arguments suggest consideration of the following auxilliary function, analytic in Ω :

$$(20) \quad \mathcal{Q}(t) = -\frac{1}{\pi} \sum_{k=1}^{\infty} (k+1)\bar{c}_k t^{-k-2}.$$

This function appears in the corresponding Lagrange-Euler equation. It turns out that the local maxima must take the form

$$\omega(t) = \chi_{\Omega}(t) \exp[-i \arg \mathcal{Q}(t)]$$

(cf. [5]). But, in order to avoid a delicate question on existence of the extremals, we do not exploit this extra information.

From the definition of c_k it follows that

$$(21) \quad \begin{aligned} \sum_{k=1}^{\infty} (k+1)|c_k|^2 &= \sum_{k=1}^{\infty} (k+1)\bar{c}_k \left[-\frac{1}{\pi} \int t^{-k-2} \omega(t) d\sigma(t) \right] \\ &= \int \omega(t) \mathcal{Q}(t) d\sigma(t) \leq \int_{\Omega} |\mathcal{Q}(t)| d\sigma(t). \end{aligned}$$

Hence, by Hölder's inequality,

$$(22) \quad \sum_{k=1}^{\infty} (k+1)|c_k|^2 \leq \left[\int_{\Omega} |t|^3 |\mathcal{Q}(t)|^2 d\sigma(t) \right]^{1/2} \left[\int_{\Omega} |t|^{-3} d\sigma(t) \right]^{1/2}.$$

We appeal now to the fact that the functions $t^{-k-2}|t|^{3/2}$, $k = 1, 2, \dots$, are mutually orthogonal in $L^2(\Omega)$ and that

$$\int_{\Omega} |t|^{-2k-1} d\sigma(t) = 2\pi \int_1^{\infty} \frac{dr}{r^{2k}} = \frac{2\pi}{2k-1}$$

for $k = 1, 2, 3, \dots$

This gives

$$\begin{aligned} \int_{\Omega} |t|^3 |\mathcal{Q}(t)|^2 d\sigma(t) &= \frac{1}{\pi^2} \int_{\Omega} \left| \sum_{k=1}^{\infty} (k+1) \overline{c_k} t^{-k-2} |t|^{3/2} \right|^2 d\sigma(t) \\ (23) \qquad &= \frac{1}{\pi^2} \sum_{k=1}^{\infty} (k+1)^2 |c_k|^2 \int_{\Omega} |t|^{-2k-1} d\sigma(t) \\ &= \frac{2}{\pi} \sum_{k=1}^{\infty} \frac{(k+1)^2}{2k-1} |c_k|^2 \leq \frac{4}{\pi} \sum_{k=1}^{\infty} (k+1) |c_k|^2. \end{aligned}$$

This, together with (22) and (18), implies

$$\sum_{k=1}^{\infty} (k+1) |c_k|^2 \leq \sqrt{\frac{4}{\pi}} \sqrt{2\pi} \left(\sum_{k=1}^{\infty} (k+1) |c_k|^2 \right)^{1/2}.$$

Hence

$$(24) \qquad \sum_{k=1}^{\infty} (k+1) |c_k|^2 \leq 8.$$

Finally, in view of (19) the proof of inequality (15) is complete.

To achieve equality in (24) ω must have equality in (21), (22), and (23), which can occur (respectively) under the following conditions:

- (i) $\omega(t) = \overline{\mathcal{Q}(t)} / |\mathcal{Q}(t)|$, a.e. in Ω ;
- (ii) $|t|^3 |\mathcal{Q}(t)|^2 = \lambda |t|^{-3}$, a.e. in Ω with $\lambda > 0$;
- (iii) $c_k = 0$ for $k = 2, 3, 4, \dots$

Since $t^3 \mathcal{Q}(t)$ is analytic in Ω , condition (ii) constrains that $\mathcal{Q}(t) = at^{-3}$, $a \in \mathbf{C}$ for $t \in \Omega$. Then (i) becomes

$$(25) \qquad \omega(t) = e^{i\theta} \left(\frac{t}{|t|} \right)^3 \chi_{\Omega}(t), \quad 0 \leq \theta < 2\pi.$$

Finally, we require (iii). Here luck is with us, since for ω defined by (25) the coefficients c_k with $k \geq 2$ vanish. The proof of Lemma 3 is complete. \square

REMARK. It is a classical result that the norm of a linear bounded operator in a Banach space is attained (if at all) on the extremal points of the unit ball. For $L^{\infty}(\Omega)$ the only extremal points are the unimodular functions. Obviously they are not in $L^2(\mathbf{C})$. That is why we could not expect the extremals for inequality (15) to be in the class $L^2(\Omega) \cap L^{\infty}(\mathbf{C})$. On the other hand, inequality (15), which was proved originally for $\omega \in L^2(\Omega) \cap L^{\infty}(\mathbf{C})$, extends now to all $\omega \in L^{\infty}(\mathbf{C})$. The equality in (15) is attained, in this extension, only for the function (25).

Theorem 1 follows readily from Corollary 2 and Lemmas 2 and 3.

EPILOGUE. By (2) we also have

$$(26) \quad \|Tf\|_{\text{BMO}_1} \leq 3\|f\|_\infty.$$

However, this inequality is not sharp. The John–Nirenberg Lemma [8] implies the following.

There exists $\mu > 0$ such that

$$(27) \quad \int_B \exp[\mu|Tb - (Tb)_B|] < \infty$$

for every b , with $|b(z)| \leq \chi_B(z)$, a.e.

In order to estimate μ consider the following example:

$$\rho(z) = \left(\frac{z}{|z|}\right)^2 \chi_B(z).$$

By (7), $T\rho(z) = (1 + \log|z|^2)\chi_B(z)$. A computation shows that

$$\int_B \exp[\mu|T\rho - (T\rho)_B|] = \frac{2\mu}{(1-\mu^2)e} + \frac{e^\mu}{1+\mu} < \infty$$

for $|\mu| < 1$. We believe that (27) holds for every b and $|\mu| < 1$. If this happens to be true, then several interesting L^p -estimates for T and consequently for the derivatives of a quasiconformal mapping would follow ([2], [6], [7], [10]).

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