

**THE RIESZ INTEGRAL AND AN  $L^p - L^q$   
ESTIMATE FOR THE CAUCHY PROBLEM  
OF THE WAVE OPERATOR**

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ABSTRACT. In 1949, M. Riesz [3] generalized the Riemann-Liouville integral of one-variable to high dimensional Euclidean spaces and obtained a powerful method now known as the Riesz integral for studying wave operators. In this paper we apply the Riesz integral to get the global space-time estimate

$$\|u\|_q \leq C\{\|w\|_p + t^{(1-n)/(n+1)}(\|g\|_p + \|\nabla f\|_p)\}$$

where  $1/q = 1/p - 2/(n+1)$ ,  $1/p + 1/q = 1$ , and  $u$  is the solution of the Cauchy problem  $\square u(x, t) = w(x, t)$  in  $\mathbf{R}_+^{n+1}$ ,  $u(x, 0) = f(x)$ , and  $\partial_t u(x, 0) = g(x)$ .

**1. The Riesz distribution.** For  $x = (x_1, x_2, \dots, x_n) \in \mathbf{R}^n$ , we denote  $|x| = \sqrt{x_1^2 + x_2^2 + \dots + x_n^2}$ . Let  $\mathbf{R}^{n+1} = \{(x, t) : x \in \mathbf{R}^n, t \in \mathbf{R}\}$ , and define

$$\rho^\lambda = \begin{cases} (t^2 - |x|^2)^{\lambda/2} & \text{if } t \geq |x| \\ 0 & \text{otherwise.} \end{cases}$$

For  $\text{Re } \lambda > -2$ ,  $\rho^\lambda$  is a locally integrable function on  $\mathbf{R}^{n+1}$  and so defines a distribution

$$\langle \rho^\lambda, \phi \rangle = \int_{\mathbf{R}^{n+1}} \rho^\lambda \phi(x, t) dx dt$$

for  $\phi \in \mathcal{D}(\mathbf{R}^{n+1})$ . In spherical coordinates, the above integral can be written as

$$\langle \rho^\lambda, \phi \rangle = \int_0^\infty \int_0^t (t^2 - r^2)^{\lambda/2} r^{n-1} \bar{\phi}(r, t) dr dt$$

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Received by the editors on June 22, 1992.

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where  $\bar{\phi}(r, t) = \int_{S^{n-1}} \phi(r\omega, t) d\sigma(\omega)$  is a  $C^\infty$ -function in  $(r^2, t)$  with compact support. By the change of variable  $r = \sqrt{st}$

$$(1.1) \quad \langle \rho^\lambda, \phi \rangle = \int_0^\infty t^{\lambda+n} \Phi_\lambda(t) dt$$

where

$$\Phi_\lambda(t) = \frac{1}{2} \int_0^1 (1-s)^{\lambda/2} s^{(n-2)/2} \bar{\phi}(\sqrt{st}, t) ds$$

is a holomorphic function in  $\lambda$  for  $\operatorname{Re} \lambda > -2$  and can be extended by analytic continuation to a meromorphic function in  $\lambda \in \mathbf{C}$  with poles at  $\lambda = -2, -4, -6, \dots$ . For  $\lambda \neq -2, -4, -6, \dots$ ,  $\Phi_\lambda(t)$  is a  $C^\infty$ -function in  $t$  with compact support.

Now let  $n$  be odd,  $\lambda = -n - 2k$ . Then

$$\begin{aligned} \Phi_{-n-2k}^{(2k-1)}(0) &= \frac{1}{2} \int_0^1 (1-s)^{(-n-2k)/2} s^{(n-2)/2} \left[ \frac{\partial^{2k-1}}{\partial t^{2k-1}} \bar{\phi}(\sqrt{st}, t) \right]_{t=0} ds \\ &= \frac{1}{2} \int_0^1 (1-s)^{(-n-2k)/2} s^{(n-2)/2} \\ &\quad \cdot \left\{ \sum_{j=0}^{2k-1} \binom{2k-1}{j} s^{j/2} \partial_1^j \bar{\phi}(0, 0) \partial_2^{2k-1-j} \bar{\phi}(0, 0) \right\} ds. \end{aligned}$$

For  $j$  even,

$$\int_0^1 (1-s)^{(-n-2k)/2} s^{(n+j-2)/2} ds = \frac{\Gamma\left(\frac{-n-2k+2}{2}\right) \Gamma\left(\frac{n+j}{2}\right)}{\Gamma\left(\frac{-2k+j+2}{2}\right)} = 0.$$

But in another case, for  $j$  odd,  $\partial_1^j \bar{\phi}(0, 0) = 0$ . In other words,  $\Phi_{-n-2k}(t)$  has zero derivative of order  $2k - 1$  at  $t = 0$ , and hence

$$\operatorname{Res}_{\lambda=-n-2k} \langle \rho^\lambda, \phi \rangle = \frac{\Phi_{-n-2k}^{(2k-1)}(0)}{(2k-1)!} = 0$$

by (1.1) and  $\operatorname{Res}_{a=-j} t_+^a = ((-1)^{j-1}/(j-1!)) \cdot \delta^{(j-1)}$  [2, p. 68]. Therefore,  $\rho^\lambda$  is a holomorphic function in  $\lambda$  for  $\operatorname{Re} \lambda > -2$  and can be extended by

analytic continuation to a meromorphic function in  $\lambda \in \mathbf{C}$  with poles at

$$(1.2) \quad \begin{cases} \text{(i)} & \lambda = -2, -4, -6, \dots \\ \text{(ii)} & \lambda = -n - 1, -n - 3, -n - 5, \dots \end{cases}$$

For  $n$  even,  $\rho^\lambda$  has simple poles at these points; for  $n$  odd, it has simple poles at  $\lambda = -2, -4, \dots, -n + 1$  and double poles at  $\lambda = -n - 1, -n - 3, \dots$

Let the wave operator be denoted by

$$\square = \frac{\partial^2}{\partial t^2} - \sum_{j=1}^n \frac{\partial^2}{\partial x_j^2}.$$

For  $\text{Re } \lambda > -2$ ,  $\square \rho^{\lambda+2} = (\lambda + 2)(\lambda + n + 1)\rho^\lambda$ , and so by iteration for  $k = 1, 2, 3, \dots$ ,

$$\begin{aligned} \square^k \rho^{\lambda+2k} &= (\lambda+2)(\lambda+4) \cdots (\lambda+2k)(\lambda+n+1) \\ &\quad \cdot (\lambda+n+3) \cdots (\lambda+n+2k-1)\rho^\lambda; \end{aligned}$$

that is,

$$(1.3) \quad \begin{aligned} \langle \rho^\lambda, \phi \rangle &= \frac{\langle \square^k \rho^{\lambda+2k}, \phi \rangle}{(\lambda+2)(\lambda+4) \cdots (\lambda+2k)(\lambda+n+1)(\lambda+n+3) \cdots (\lambda+n+2k-1)}. \end{aligned}$$

By analytic continuation (1.3) holds also for  $\lambda \in \mathbf{C}$  except at the singularities of (1.2).

The distribution  $\rho^\lambda$  can be normalized. Its construction was first given by M. Riesz [3].

**Definition.** The *Riesz distribution* is defined by

$$Z_\alpha = \frac{\rho^{\alpha-n-1}}{2^{\alpha-1} \pi^{(n-1)/2} \Gamma(\alpha/2) \Gamma((\alpha-n+1)/2)}.$$

The constant  $H(\alpha, n) = 2^{\alpha-1} \pi^{(n-1)/2} \Gamma(\alpha/2) \Gamma((\alpha-n+1)/2)$  is so determined that

$$(1.4) \quad \langle Z_\alpha, e^{-t} \rangle = 1.$$

We note that  $\Gamma(\alpha/2)$  has simple poles at  $\alpha = 0, -2, -4, \dots$ , (i.e.,  $\alpha - n - 1 = -n - 1, -n - 3, -n - 5, \dots$ ) and  $\Gamma((\alpha - n + 1)/2)$  has simple poles at  $\alpha - n + 1 = 0, -2, -4, \dots$ , (i.e.,  $\alpha - n - 1 = -2, -4, -6, \dots$ ). Hence  $Z_\alpha$  is an entire function of  $\alpha \in \mathbf{C}$  and satisfies

$$(1.5) \quad \begin{aligned} \square Z_\alpha &= Z_{\alpha-2} \\ \square^k Z_\alpha &= Z_{\alpha-2k}, \quad k = 0, 1, 2, \dots \end{aligned}$$

By calculating the residues of  $\langle \rho^\lambda, \phi \rangle$ , we have

$$(1.6) \quad \begin{aligned} Z_0 &= \delta \\ Z_{-2k} &= \square^k \delta, \quad k = 0, 1, 2, \dots \end{aligned}$$

Moreover, the support of  $Z_\alpha$ , for all complex  $\alpha \in \mathbf{C}$ , is contained in the forward cone  $C = \{(x, t) \in \mathbf{R}^{n+1} : t \geq |x|\}$ .

Combining (1.5) and (1.6), we obtain easily

$$(1.7) \quad \square^k Z_{2k} = \delta, \quad k = 0, 1, 2, \dots$$

In particular,  $\square Z_2 = \delta$ , so  $Z_2$  is a fundamental solution of wave operator (cf. [2, Section 6.2]).

The convolution property of  $Z_\alpha$  is given by

**Theorem 1.8.**

$$Z_\alpha * Z_\beta = Z_{\alpha+\beta}.$$

*Note.* Since  $\text{Supp}(Z_\alpha)$  and  $\text{Supp}(Z_\beta)$  are contained in the cone  $C = \{(x, t) \in \mathbf{R}^{n+1} : t \geq |x|\}$ ,  $\text{Supp}(Z_\alpha * Z_\beta)$  is concentrated in the compact set  $C \cap G$ , where  $G$  is the reflection of  $C$  and translated by some vector on  $\mathbf{R}^{n+1}$ . That implies the convolution  $Z_\alpha * Z_\beta$  exists.

*Proof.* It suffices to verify for  $\text{Re } \alpha, \text{Re } \beta$  large enough. For  $(x, t) \in \mathbf{R}^{n+1}$ , let

$$T = \int_D (\tau^2 - \xi_1^2 \dots - \xi_n^2)^{(\alpha-n-1)/2} \cdot ((t-\tau)^2 - (x_1-\xi_1)^2 \dots - (x_n-\xi_n)^2)^{(\beta-n-1)/2} d\xi d\tau$$

where  $D = \{(\xi, \tau) \in \mathbf{R}^{n+1} : \tau \geq |\xi|, t - \tau \geq |x - \xi|\}$ . By a rotation of the space axes followed by a two-dimensional Lorentz transformation,

$$T = \int_{D_1} (\tau^2 - \xi_1^2 - \dots - \xi_n^2)^{(\alpha-n-1)/2} ((s-\tau)^2 - \xi_1^2 - \dots - \xi_n^2)^{(\beta-n-1)/2} d\xi d\tau$$

where  $s = \sqrt{t^2 - |x|^2}$ ,  $D_1 = \{(\xi, \tau) \in \mathbf{R}^{n+1} : \tau \geq |\xi|, s - \tau \geq |\xi|\}$ . Thus, we have

$$T = \Omega_n \int_{D_2} (\tau^2 - \eta^2)^{(\alpha-n-1)/2} ((s-\tau)^2 - \eta^2)^{(\beta-n-1)/2} \eta^{n-1} d\eta d\tau$$

where  $D_2 = \{(\eta, \tau) \in \mathbf{R}^2 : 0 \leq \tau - \eta \leq s, 0 \leq \tau + \eta \leq s\}$  and  $\Omega_n$  is the hypersurface area of the unit sphere in  $\mathbf{R}^n$ . By the transformation  $\tau + \eta = u$  and  $\tau - \eta = v$ ,

$$\begin{aligned} T &= \frac{\Omega_n}{2} \int_0^s \int_0^s u^{(\alpha-n-1)/2} (s-u)^{(\beta-n-1)/2} v^{(\alpha-n-1)/2} \\ &\quad \cdot (s-v)^{(\beta-n-1)/2} \left(\frac{u-v}{2}\right)^{n-1} du dv \\ &= s^{\alpha+\beta-n-1} B_n(\alpha, \beta) \end{aligned}$$

where

$$\begin{aligned} B_n(\alpha, \beta) &= \frac{\Omega_n}{2^n} \int_0^1 \int_0^1 u^{(\alpha-n-1)/2} (1-u)^{(\beta-n-1)/2} v^{(\alpha-n-1)/2} \\ &\quad \cdot (1-v)^{(\beta-n-1)/2} (u-v)^{n-1} du dv \end{aligned}$$

depends only on  $\alpha, \beta$ , and  $n$ . We put  $e^{-t} = e^{-(t-\tau)-\tau}$  into (1.4), then

$$H(\alpha, n)H(\beta, n) = \langle T, e^{-t} \rangle = B_n(\alpha, \beta)H(\alpha + \beta, n)$$

and so  $B_n(\alpha, \beta) = H(\alpha, n)H(\beta, n)/H(\alpha + \beta, n)$ . For  $\phi \in \mathcal{D}(\mathbf{R}^{n+1})$ ,

$$\begin{aligned}
 \langle Z_\alpha * Z_\beta, \phi \rangle &= \frac{1}{H(\alpha, n)H(\beta, n)} \int_{s \geq |y|} (s^2 - |y|^2)^{(\alpha-n-1)/2} dy ds \\
 &\quad \cdot \int_{t \geq |x|} (t^2 - |x|^2)^{(\beta-n-1)/2} \phi(x+y, t+s) dx dt \\
 &= \frac{1}{H(\alpha, n)H(\beta, n)} \int_{s \geq |y|} (s^2 - |y|^2)^{(\alpha-n-1)/2} dy ds \\
 &\quad \cdot \int_{t-s \geq |x-y|} ((t-s)^2 - |x-y|^2)^{(\beta-n-1)/2} \phi(x, t) dx dt \\
 &= \frac{1}{H(\alpha, n)H(\beta, n)} \int_{t \geq |x|} \phi(x, t) dx dt \\
 &\quad \cdot \int_{(y,s) \in \mathcal{D}} (s^2 - |y|^2)^{(\alpha-n-1)/2} \\
 &\quad \quad \cdot ((t-s)^2 - |x-y|^2)^{(\beta-n-1)/2} dy ds \\
 &= \frac{1}{H(\alpha+\beta, n)} \int_{t \geq |x|} \phi(x, t) (t^2 - |x|^2)^{(\alpha+\beta-n-1)/2} dx dt \\
 &= \langle Z_{\alpha+\beta}, \phi \rangle.
 \end{aligned}$$

## 2. The solution of the Cauchy problem and $L^p - L^q$ estimate.

As mentioned in the previous section, a fundamental solution of the wave equation is given by

$$E(x, t) = Z_2(x, t) = \frac{\rho^{1-n}}{2\pi^{(n-1)/2}\Gamma((3-n)/2)}$$

which can be used to solve the Cauchy problem

$$(2.1) \quad \begin{cases} \square u(x, t) = w(x, t) & \text{for } t \geq 0 \\ u(x, 0) = f(x) \\ \partial_t u(x, 0) = g(x) \end{cases}$$

where  $w$  is a function on  $\mathbf{R}^{n+1}$  that vanishes for  $t < 0$ ,  $f$  and  $g$  are functions on  $\mathbf{R}^n$ , all of those are assumed to be sufficiently differentiable.

The Riesz distribution can be interpreted as an operator. For any  $\alpha \in \mathbf{C}$ , any  $\phi \in \mathcal{D}(\mathbf{R}^{n+1})$  or  $\mathcal{S}(\mathbf{R}^{n+1})$  with support contained in a translation of the upper half space  $t \geq 0$ , we define

$$I^\alpha \phi = Z_\alpha * \phi$$

which is known as the *Riesz integral*. Corresponding to (1.5), (1.6), (1.7), and Theorem 1.8, the Riesz integral  $I^\alpha$  has the following properties:

$$\begin{aligned} \square I^\alpha &= I^{\alpha-2}, & \square^k I^\alpha &= I^{\alpha-2k} \\ I^0 &= \text{identity}, & I^{-2k} &= \square^k \\ I^2 \square &= \square I^2 = \text{identity}, & I^{2k} \square^k &= \square^k I^{2k} = \text{identity} \\ I^\alpha I^\beta &= I^{\alpha+\beta} \end{aligned}$$

for  $k = 0, 1, 2, \dots$ .

For  $(x, t) \in \mathbf{R}^{n+1}$ , let  $\Omega = \{(\xi, \tau) \in \mathbf{R}^{n+1} : t - \tau \geq |x - \xi|, 0 \leq \tau \leq t\}$ , then  $\partial\Omega = \{(\xi, \tau) \in \mathbf{R}^{n+1} : t - \tau = |x - \xi|, 0 \leq \tau \leq t\} \cup \{(\xi, 0) \in \mathbf{R}^{n+1} : |x - \xi| \leq t\} \equiv B_1 \cup B_2$ . We use Green's formula

$$\int_{\Omega} (u \square v - v \square u) dV = \int_{\partial\Omega} \left( u \frac{\partial v}{\partial N} - v \frac{\partial u}{\partial N} \right) dS.$$

Let  $v = v(\xi, \tau) = (1/H(\alpha + 2, n))((t - \tau)^2 - |x - \xi|^2)^{(\alpha-n+1)/2}$ . Then  $\square_{(\xi, \tau)} v(\xi, \tau) = (1/H(\alpha, n))((t - \tau)^2 - |x - \xi|^2)^{(\alpha-n-1)/2}$ . Now  $v = 0$ ,  $\partial v / \partial N = 0$  on  $B_1$ . Hence

$$\begin{aligned} (2.2) \quad & \frac{1}{H(\alpha, n)} \int_{t-\tau \geq |x-\xi| \text{ and } \tau \geq 0} u(\xi, \tau) ((t-\tau)^2 - |x-\xi|^2)^{(\alpha-n-1)/2} d\xi d\tau \\ & - \frac{1}{H(\alpha+2, n)} \int_{t-\tau \geq |x-\xi| \text{ and } \tau \geq 0} \square u(\xi, \tau) ((t-\tau)^2 \\ & \quad \cdot - |x-\xi|^2)^{(\alpha-n+1)/2} d\xi d\tau \\ & = \frac{1}{H(\alpha+2, n)} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(\alpha-n+1)/2} \partial_\tau u(\xi, 0) d\xi \\ & \quad + \frac{\alpha-n+1}{H(\alpha+2, n)} t \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(\alpha-n-1)/2} u(\xi, 0) d\xi. \end{aligned}$$

Let  $u$  be the desired solution of (2.1) that vanishes for  $t < 0$ , taking  $\alpha = 0$  in (2.2),

$$\begin{aligned}
 (2.3) \quad u(x, t) &= I^2 w(x, t) + \frac{1}{H(2, n)} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} g(\xi) d\xi \\
 &\quad + \frac{1-n}{H(2, n)} t \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(-1-n)/2} f(\xi) d\xi \\
 &= I^2 w(x, t) + \frac{1}{H(2, n)} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} g(\xi) d\xi \\
 &\quad + \frac{1}{H(2, n)} \frac{d}{dt} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} f(\xi) d\xi \\
 &= (E *_{(x,t)} w)(x, t) + (E *_{(x)} g)(x, t) + \left( \frac{\partial E}{\partial t} *_{(x)} f \right)(x, t)
 \end{aligned}$$

where  $*_{(x,t)}$  and  $*_{(x)}$  denote the convolution with respect to variable  $(x, t)$  or  $x$  only.

Moreover, if  $n$  is even,  $g \in C^{(n+2)/2}(\mathbf{R}^n)$ , and  $f \in C^{(n+4)/2}(\mathbf{R}^n)$ , we can write

$$\begin{aligned}
 &\frac{1}{H(2, n)} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} g(\xi) d\xi \\
 &= (2\pi)^{-n/2} \left( \frac{1}{t} \frac{\partial}{\partial t} \right)^{(n-2)/2} t^{n-1} \int_{|\xi| \leq 1} \frac{g(x+t\xi)}{\sqrt{1-|\xi|^2}} d\xi
 \end{aligned}$$

and

$$\begin{aligned}
 &\frac{1}{H(2, n)} \frac{\partial}{\partial t} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} f(\xi) d\xi \\
 &= (2\pi)^{-n/2} \frac{\partial}{\partial t} \left( \frac{1}{t} \frac{\partial}{\partial t} \right)^{(n-2)/2} t^{n-1} \int_{|\xi| \leq 1} \frac{f(x+t\xi)}{\sqrt{1-|\xi|^2}} d\xi.
 \end{aligned}$$

If  $n$  is odd and  $n \geq 3$ ,  $g \in C^{(n+1)/2}(\mathbf{R}^n)$  and  $f \in C^{(n+3)/2}(\mathbf{R}^n)$ , then we have

$$\begin{aligned}
 &\frac{1}{H(2, n)} \int_{|x-\xi| \leq t} (t^2 - |x-\xi|^2)^{(1-n)/2} g(\xi) d\xi \\
 &= \frac{1}{2} (2\pi)^{(1-n)/2} \left( \frac{1}{t} \frac{\partial}{\partial t} \right)^{(n-3)/2} t^{n-2} \int_{|\xi|=1} g(x+t\xi) d\sigma(\xi)
 \end{aligned}$$



and

$$\begin{aligned} & \frac{1}{H(2, n)} \frac{\partial}{\partial t} \int_{|x-\xi| \leq t} (t^2 - |x - \xi|^2)^{(1-n)/2} f(\xi) d\xi \\ &= \frac{1}{2} (2\pi)^{(1-n)/2} \frac{\partial}{\partial t} \left( \frac{1}{t} \frac{\partial}{\partial t} \right)^{(n-3)/2} t^{n-1} \int_{|\xi|=1} f(x + t\xi) d\sigma(\xi). \end{aligned}$$

If also  $w(x, t)$  belongs to  $C^{[n/2]+1}(\mathbf{R}^{n+1})$  that vanishes for  $t < 0$ , then (2.3) is the classical solution of the Cauchy problem (2.1).

For the case  $n = 0$ , the Riesz integral has the form

$$I_1^\alpha f(x) = \frac{1}{\Gamma(\alpha)} \int_a^x (x - t)^{\alpha-1} f(t) dt$$

which is also called *Riemann-Liouville integral*. Hardy-Littlewood [1] has proved that  $I_1^\alpha$  is of type  $(p, q)$  with

$$1/q = 1/p - \alpha, \quad 0 < \alpha < 1/p, \quad p > 1.$$

We expect the Riesz integral with  $\alpha = 2$

$$\begin{aligned} I^2 \phi(x, t) &= Z_2 * \phi(x, t) \\ &= \frac{1}{H(2, n)} \int_{t-\tau \geq |x-\xi|} ((t-\tau)^2 - |x-\xi|^2)^{(1-n)/2} \phi(\xi, \tau) d\xi d\tau \end{aligned}$$

has similar property for  $1/q = 1/p - 2/(n + 1)$  and  $1/p + 1/q = 1$ .

By Stein-Weiss [4, p. 171, Theorem 4.15],

$$\mathcal{F}((1 - |x|^2)_+^\lambda)(\xi) = (2\pi)^{n/2} 2^\lambda \Gamma(\lambda + 1) |\xi|^{(-2\lambda-n)/2} J_{(2\lambda+n)/2}(|\xi|)$$

and so

$$\begin{aligned} \mathcal{F}_x((t^2 - |x|^2)_+^{\lambda/2})(\xi) &= t^\lambda \mathcal{F}_x(\delta_{t^{-1}}(1 - |x|^2)_+^{\lambda/2})(\xi) \\ &= t^{\lambda+n} \delta_t \mathcal{F}_x((1 - |x|^2)_+^{\lambda/2})(\xi) \\ &= t^{\lambda+n} \delta_t \left( (2\pi)^{n/2} 2^{\lambda/2} \Gamma(\lambda/2 + 1) \right. \\ &\quad \left. \cdot |\xi|^{(-\lambda-n)/2} J_{(\lambda+n)/2}(|\xi|) \right) \\ &= (2\pi)^{n/2} 2^{\lambda/2} \Gamma(\lambda/2 + 1) \left| \frac{\xi}{t} \right|^{(-\lambda-n)/2} J_{(\lambda+n)/2}(|t\xi|) \end{aligned}$$

where  $\mathcal{F}f(\xi) = \hat{f}(\xi) = \int_{\mathbf{R}^n} e^{-ix\xi} f(x) dx$  denotes the Fourier transform,  $\mathcal{F}_x \phi(\xi, t) = \tilde{\phi}(\xi, t)$  the Fourier transform with respect to  $x$ ,  $\mathcal{F}_\xi^{-1}$  the inverse Fourier transform with respect to  $\xi$ , and  $\delta_t f(x) = f(tx)$ . Hence

$$\tilde{E}(\xi, t) = \frac{1}{2\pi^{(n-1)/2} \Gamma((3-n)/2)} \mathcal{F}_x((t^2 - |x|^2)_+^{(1-n)/2})(\xi) = \frac{\sin |\xi| t}{|\xi|}$$

and so by (2.3)

$$\begin{aligned} \tilde{u}(\xi, t) &= \int_{-\infty}^{\infty} \tilde{E}(\xi, t-s) \tilde{w}(\xi, s) ds + \frac{\sin |\xi| t}{|\xi|} \hat{g}(\xi) + \cos |\xi| t \cdot \hat{f}(\xi) \\ &= \int_0^t \frac{\sin |\xi|(t-s)}{|\xi|} \tilde{w}(\xi, s) ds + \frac{\sin |\xi| t}{|\xi|} \hat{g}(\xi) + \cos |\xi| t \cdot \hat{f}(\xi). \end{aligned}$$

Strichartz [5] has proved that

$$\begin{aligned} \|\mathcal{F}_\xi^{-1} \left( \frac{\sin |\xi|}{|\xi|} \hat{g}(\xi) \right)\|_q &\leq C_1 \|g\|_p \\ \|\mathcal{F}_\xi^{-1} (\cos |\xi| \cdot \hat{f}(\xi))\|_q &\leq C_2 \|\nabla f\|_p \end{aligned}$$

for  $1/p + 1/q = 1$  and  $1/q = 1/p - 2/(n+1)$ ,  $n \geq 2$ . Hence the solution  $u(x, t)$  of problem (2.1) with  $w = 0$  satisfies that

$$\begin{aligned} \|u\|_q &\leq \|\mathcal{F}_\xi^{-1} \left( \frac{\sin |\xi| t}{|\xi|} \hat{g}(\xi) \right)\|_q + \|\mathcal{F}_\xi^{-1} (\cos |\xi| t \cdot \hat{f}(\xi))\|_q \\ &= \|t^n \mathcal{F}_\xi^{-1} \left( \frac{\sin |\xi| t}{|\xi|} t^{-n} \hat{g}(\xi) \right)\|_q \\ &\quad + \|t^n \mathcal{F}_\xi^{-1} (\cos |\xi| t \cdot t^{-n} \hat{f}(\xi))\|_q \\ &= \|t^{n+1} \mathcal{F}_\xi^{-1} \delta_t \left( \frac{\sin |\xi|}{|\xi|} \cdot t^{-n} \delta_{t^{-1}} \hat{g}(\xi) \right)\|_q \\ (2.4) \quad &\quad + \|t^n \mathcal{F}_\xi^{-1} \delta_t (\cos |\xi| \cdot t^{-n} \delta_{t^{-1}} \hat{f}(\xi))\|_q \\ &= \|t \delta_{t^{-1}} \mathcal{F}_\xi^{-1} \left( \frac{\sin |\xi|}{|\xi|} \widehat{\delta_t g}(\xi) \right)\|_q \\ &\quad + \|\delta_{t^{-1}} \mathcal{F}_\xi^{-1} (\cos |\xi| \cdot \widehat{\delta_t f}(\xi))\|_q \\ &= t^{1+(n/q)} \|\mathcal{F}_\xi^{-1} \left( \frac{\sin |\xi|}{|\xi|} \widehat{\delta_t g}(\xi) \right)\|_q \end{aligned}$$

$$\begin{aligned} & + t^{n/q} \|\mathcal{F}_\xi^{-1}(\cos|\xi| \cdot \widehat{\delta_t f}(\xi))\|_q \\ & \leq C_1 t^{1+(n/q)} \|\delta_t g\|_p + C_2 t^{n/q} \|\nabla(\delta_t f)\|_p \\ & = C_1 t^{1+(n/q)-(n/p)} \|g\|_p + C_2 t^{1+(n/q)-(n/p)} \|\nabla f\|_p. \end{aligned}$$

On the other hand, let  $f = g = 0$  in problem (2.1), by Jensen's inequality and (2.4)

$$\begin{aligned} \|u(\cdot, t)\|_q & = \|\mathcal{F}_\xi^{-1} \int_0^t \frac{\sin|\xi|(t-s)}{|\xi|} \tilde{w}(\xi, s) ds\|_q \\ & \leq \int_0^t \|\mathcal{F}_\xi^{-1} \left( \frac{\sin|\xi|(t-s)}{|\xi|} \tilde{w}(\xi, s) \right)\|_q ds \\ & \leq C_3 \int_0^t (t-s)^{1+(n/q)-(n/p)} \|w(\cdot, s)\|_p dx. \end{aligned}$$

The last integral is exactly the Riemann-Liouville integral  $I_1^{2/(n+1)}(\|w(\cdot, t)\|_p)$  when  $1 + n/q - n/p = 1 - (2n)/(n + 1) = (1 - n)/(n + 1)$ . Now we take  $L^q$ -norm with respect to the time variable  $t$ , then

$$(2.5) \quad \|u\|_q \leq C_3 \|I_1^{2/(n+1)}(\|w(\cdot, t)\|_p)\|_q \leq C_4 \|w\|_p$$

since  $I_1^{2/(n+1)}$  is of type  $(p, q)$  with  $1/q = 1/p - 2/(n + 1)$ .

Combining (2.4) and (2.5) we obtain the global space-time estimate for the Cauchy problem (2.1)

$$\|u\|_q \leq C (\|w\|_p + t^{(1-n)/(n+1)} (\|g\|_p + \|\nabla f\|_p))$$

where  $1/q = 1/p - 2/(n + 1)$  and  $1/p + 1/q = 1$ .

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