

On the a priori estimate for solutions of the Cauchy problem for some non-linear wave equations

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For the global Cauchy problem of wave equation, the existence of an a priori estimate of the solution is very useful as we have shown recently in another report [1] [2] for one special type of the non-linear wave equation :

$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2} + f(u) \frac{\partial u}{\partial t} + g(u)$$

under relatively weak conditions.

Here, we note that a priori estimate is also obtained for wave equation of a little different type with more than one space dimension which is identical to the equation treated by Konrad Jorgens [3] in the case of 3 dimension and without damping term.

At first we shall treat the case in which the space dimension is 2. Our Cauchy problem is the following : Find the solution of the equation

$$(1) \quad \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial t^2} + f\left(\frac{\partial u}{\partial t}\right) + g(u)$$

satisfying the following initial conditions.

$$(2) \quad \begin{cases} u(x, y, 0) = u_0(x, y) \\ u_t(x, y, 0) = u_1(x, y) \end{cases}$$

where $u_0(x, y)$ belongs to C^3 and $u_1(x, y)$ belongs to C^2 .

Here we do not solve this problem, but we obtain an a priori estimate for the solution of this problem assuming $u_0(x, y)$ and

$u_i(x, y)$ have compact carriers¹⁾ under conditions for functions f and g :

Conditions

- i) $f(u)$ and $g(u)$ are continuously differentiable in $-\infty < u < +\infty$
- ii) $\operatorname{sgn} uf(u) \geq 0$ and $G(u) = \int_0^u g(u) du \geq 0$ for $|u| > M$.
- iii) $f'(u) \geq -k$ (k is one positive constant)
 $|g'(u)| \leq \text{Polynomial of } |u|$

Before we proceed to write our results, we define two generalized energies $E_0(t)$ and $E_1(t)$ for our solutions $u(x, y, t)$ of (1) and (2).

$$(4) \quad E_0(t) = \iint \left[G(u) + \frac{1}{2} \left(\frac{\partial u}{\partial t} \right)^2 + \frac{1}{2} \left(\frac{\partial u}{\partial x} \right)^2 + \frac{1}{2} \left(\frac{\partial u}{\partial y} \right)^2 \right] dx dy$$

$$(5) \quad E_1(t) = \frac{1}{2} \iint \left[\left(\frac{\partial^2 u}{\partial t \partial x} \right)^2 + \left(\frac{\partial^2 u}{\partial t \partial y} \right)^2 + \left(\frac{\partial^2 u}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 u}{\partial y \partial x} \right) + \left(\frac{\partial^2 u}{\partial y^2} \right)^2 \right] dx dy$$

and we see

$$(6) \quad E_0(0) = \iint \left[G(u_0) + \frac{1}{2} (u_1)^2 + \frac{1}{2} \left(\frac{\partial u_0}{\partial x} \right)^2 + \frac{1}{2} \left(\frac{\partial u_0}{\partial y} \right)^2 \right] dx dy$$

$$(7) \quad E_1(0) = \iint \frac{1}{2} \left[\left(\frac{\partial u_1}{\partial x} \right)^2 + \left(\frac{\partial u_1}{\partial y} \right)^2 + \left(\frac{\partial^2 u_0}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 u_0}{\partial x \partial y} \right)^2 + \left(\frac{\partial^2 u_0}{\partial y^2} \right)^2 \right] dx dy$$

where always integrals are taken in whole x, y plane which is possible, because, $u_1(x, y)$ and $u_0(x, y)$ have compact carriers and $u(x, y, t)$ also.

Now we estimate the energy $E_0(t)$ of the solution by the initial energy. First we transform (1) and (2) into a system of equations.

$$(8) \quad \begin{cases} p = \frac{\partial u}{\partial x}, & q = \frac{\partial u}{\partial y} \\ \frac{\partial u}{\partial t} = v \\ \frac{\partial v}{\partial t} = \frac{\partial p}{\partial x} + \frac{\partial q}{\partial y} - f(v) - g(u) \\ \frac{\partial p}{\partial t} = \frac{\partial v}{\partial x} \\ \frac{\partial q}{\partial t} = \frac{\partial v}{\partial y} \end{cases}$$

$$(9) \quad \begin{cases} v(x, y, 0) = u_1(x, y) \\ u(x, y, 0) = u_0(x, y) \\ p(x, y, 0) = \frac{\partial u_0}{\partial x}(x, y) \\ q(x, y, 0) = \frac{\partial u_0}{\partial y}(x, y) \end{cases}$$

and

$$(10) \quad E_0(t) = \iint \left[G(u) + \frac{v^2}{2} + \frac{p^2}{2} + \frac{q^2}{2} \right] dx dy$$

Differentiating (10) with respect to t and Considering (8), we have

$$(11) \quad \begin{aligned} \frac{dE_0}{dt} &= \iint \left[g(u)v + v \frac{\partial v}{\partial t} + p \frac{\partial p}{\partial t} + q \frac{\partial q}{\partial t} \right] dx dy \\ &= \iint \left[g(u)v + v \frac{\partial p}{\partial x} + v \frac{\partial q}{\partial y} + p \frac{\partial u}{\partial x} + q \frac{\partial v}{\partial y} - f(v)v - g(u)v \right] dx dy \\ &= \iint [-vf(v)] dx dy \leq \iint_{|v| \leq M} [vf(v)] dx dy \\ &\leq L \iint \left[\frac{v^2}{2} + \frac{p^2}{2} + \frac{q^2}{2} + G(u) \right] dx dy + L_0 l, \end{aligned}$$

$$(12) \quad \frac{dE_0(t)}{dt} \leq LE_0(t) + L_0 l \quad 0 \leq t \leq h,$$

where l is the area of the carrier of $u_0(x, y)$ multiplied by $2h$.

$$(13) \quad E_0(t) \leq e^{Lh}E_0(0) + e^{Lh}L_0lh = e^{Lh}(E_0(0) + L_0lh)$$

Next, we proceed to estimate $E_1(t)$. We write

$$E_1(t) = \frac{1}{2} \iint [v_x^2 + v_y^2 + p_x^2 + p_y^2 + q_x^2 + q_y^2] dx dy.$$

Differentiating (8) by x and y , we have:

$$(14) \quad \begin{cases} \frac{\partial v_x}{\partial t} = \frac{\partial p_x}{\partial x} + \frac{\partial q_x}{\partial y} - f'(v)v_x - g'(u)p \\ \frac{\partial p_x}{\partial t} = \frac{\partial v_x}{\partial x} \\ \frac{\partial q_x}{\partial t} = \frac{\partial v_x}{\partial y} \end{cases}$$

$$\begin{cases} \frac{\partial v_y}{\partial t} = \frac{\partial p_y}{\partial x} + \frac{\partial q_y}{\partial y} - f'(v)v_y - g'(u)q \\ \frac{\partial p_y}{\partial t} = \frac{\partial v_y}{\partial x} \\ \frac{\partial q_y}{\partial t} = \frac{\partial v_y}{\partial y} \end{cases}$$

And we obtain,

$$\begin{aligned} (15) \quad \frac{dE_1}{dt} &= \iint \left[v_x \frac{\partial v_x}{\partial t} + v_y \frac{\partial v_y}{\partial t} + p_x \frac{\partial p_x}{\partial t} + q_x \frac{\partial q_x}{\partial t} + p_y \frac{\partial p_y}{\partial t} + q_y \frac{\partial q_y}{\partial t} \right] dx dy \\ &= \iint \left[-f'(v)v_x^2 - g'(u)pv_x - f'(v)v_y^2 - g'(u)qv_y \right] dx dy \\ &= \iint \left[k(v_x^2 + v_y^2) - g'(u)pv_x - g'(u)qv_y \right] dx dy. \end{aligned}$$

We consider the integral :

$$I_1 = \iint |g'(u)pv_x| dx dy.$$

By the condition (iii)

$$\begin{aligned} I_1 &\leq \iint |u|^\alpha |p| |v_x| dx dy \leq \sqrt{\iint |u|^{2\alpha} p^2 dx dy} \sqrt{\iint v_x^2 dx dy} \\ &\leq \left(\iint |u|^{4\alpha} dx dy \right)^{1/4} \cdot \left(\iint p^4 dx dy \right)^{1/4} \left(\iint v_x^2 dx dy \right)^{1/2} \\ &\leq c \left[\iint (p^2 + q^2) dx dy \right]^{\alpha/2} \cdot E_1(t) \\ &\leq cE_0(t)^{\alpha/2} E_1(t), \end{aligned}$$

$$[\log E_1(t)]_0^t \leq c \int_0^t E_0(\tau)^{\alpha/2} d\tau + k,$$

$$E_1(t) \leq E_1(0) e^{c \int_0^t E_0(\tau)^{\alpha/2} d\tau} \leq E_1(0) e^{c \int_0^h E_0(\tau)^{\alpha/2} d\tau + k}.$$

Then we obtain by the Sobolev's lemma

$$\begin{aligned} |u(x, y, t)| &\leq C |E_0(t) + E_1(t)| \\ &\leq C \{ (E_0(0) + L_0 l h) e^{Lh} + E_1(0) e^{c(E_0(0) + L_0 l h)^{\alpha/2} e^{L\alpha h} h + k} \}. \end{aligned}$$

This is our desired results.

We proceed to show that similar results can be obtained for the

case of 3 space dimension under more stringent condition. We replace condition iii) by,

$$\text{iii) } f'(u) \geq -k, \quad |g'(u)| \leq c|u|^2$$

Our equation is the following :

$$(16) \quad \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = \frac{\partial^2 u}{\partial t^2} + f\left(\frac{\partial u}{\partial t}\right) + g(u)$$

We can estimate the energy $E_0(t)$ by the same argument of the preceding case, where $E_0(t)$ is defined.

$$(17) \quad E_0(t) = \iiint \left\{ G(u) + \frac{1}{2}(v^2 + p_1^2 + p_2^2 + p_3^2) \right\} dy$$

where $v = \frac{\partial u}{\partial t}$, $p = \frac{\partial u}{\partial x}$, $p_2 = \frac{\partial u}{\partial y}$, $p_3 = \frac{\partial u}{\partial z}$. Then we have

$$(18) \quad E_0(t) \leq e^{Lh} \{E_0(0) + M_0 l\} \quad 0 \leq t \leq h.$$

$E_1(t)$ is the integral :

$$\frac{1}{2} \iiint \sum_{x,y,z} \{v_x^2 + p_{1x}^2 + p_{2x}^2 + p_{3x}^2\} dV.$$

We obtain by the condition iii).

$$\begin{aligned} \frac{dE_1(t)}{dt} &\leq - \sum \iiint f'(v)v_x^2 + \sum \iiint 3u^2 p_1 v_x dV. \\ &\leq k \sum \iiint v_x^2 + \sum \iiint 3u^2 |p_1 v_x| dV. \end{aligned}$$

We treat the last term by the similar inequality as we have used in (15) :

$$\begin{aligned} & \left| \iiint u^2 p_1 v_x dV. \right| \\ & \leq \left[\iiint u^4 p_1^2 dV \right]^{1/2} \left[\iiint v_x^2 dV \right]^{1/2} \\ & \leq \left[\left(\iiint u^6 dV \right)^{2/3} \left(\iiint p_1^6 dV \right)^{1/3} \right]^{1/2} \left[\iiint v_x^2 dV \right]^{1/2} \\ & \leq \left(\iiint u^6 dV \right)^{1/3} \left(\iiint p_1^6 dV \right)^{1/6} \left[\iiint v_x^2 dV \right]^{1/2} \\ & \leq \left[\iiint (p_1^2 + p_2^2 + p_3^2) dV \right] \left[\iiint (p_{1x}^2 + p_{1y}^2 + p_{1z}^2) dV \right]^{1/2} \left[\iiint v_x^2 dV \right]^{1/2} \\ & \leq E_0(t) E_1(t) \end{aligned}$$

Just similarly we can estimate other terms. It is easy to see that the maximum norm of $U(x, y, z, t)$ for $0 \leq t \leq h$ is majorized by $E_0(0)$ and $E_1(0)^2$.

NOTES

- 1) We assume also that the solution $u(x, y, t)$ and its derivatives of 3rd order with respect to x, y and t are square integrable in xy space for all t . By the Sobolev's work [4], we can find always this solution for sufficiently small t , for our Cauchy data.
- 2) We could not find the bound for $\frac{\partial u}{\partial t}$ by $E_0(0)$ and $E_1(0)$, therefore the existence of a global solution of the Cauchy problem is not proved for the equation (1) and (16). But if $f\left(\frac{\partial u}{\partial t}\right)$ is linear for $\frac{\partial u}{\partial t}$, we can easily prove the global existence of the solution of the Cauchy problem.

BIBLIOGRAPHY

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