

## BERNOULLI PERCOLATION ABOVE THRESHOLD: AN INVASION PERCOLATION ANALYSIS

BY J. T. CHAYES,<sup>1</sup> L. CHAYES<sup>1</sup> AND C. M. NEWMAN<sup>2</sup>

*Cornell University, Cornell University and University of Arizona*

Using the invasion percolation process, we prove the following for Bernoulli percolation on  $\mathbb{Z}^d$  ( $d > 2$ ): (1) exponential decay of the truncated connectivity,  $\tau'_{xy} \equiv P(x \text{ and } y \text{ belong to the same finite cluster}) \leq \exp(-m\|x - y\|)$ ; (2) infinite differentiability of  $P_\infty(p)$ , the infinite cluster density, and of  $\chi'(p)$ , the expected size of finite clusters, as functions of  $p$ , the density of occupied bonds; and (3) upper bounds on the cluster size distribution tail,  $P_n \equiv P(\text{the cluster of the origin contains exactly } n \text{ bonds}) \leq \exp(-[c/\log n]n^{(d-1)/d})$ . Such results (without the  $\log n$  denominator in (3)) were previously known for  $d = 2$  and  $p > p_c$ , the usual percolation threshold, or for  $d > 2$  and  $p$  close to 1. We establish these results for all  $d > 2$  when  $p$  is above a limit of "slab thresholds," conjectured to coincide with  $p_c$ .

**1. Introduction and results.** Many of the properties of Bernoulli percolation can be characterized in terms of the behavior of large finite clusters. This behavior is well understood below the threshold at which the expected cluster size diverges. In particular, both the (truncated) connectivity function and the cluster size distribution tail are known to decay exponentially ([10], [13], [3]). On the other hand, previous results about these quantities above the percolation threshold ([15], [2], [14]) are quite incomplete. This note represents a moderately successful attempt to remedy this situation by exploiting the relation between standard percolation and a recently invented dynamic growth model known as invasion percolation ([16], [5], [21], [8]).

In this paper, we will treat standard Bernoulli bond percolation on  $\mathbb{Z}^d$  ([4], [14]), in which the nearest neighbor bonds are independently "occupied" or "vacant" with probability  $p$  or  $1 - p$ . All our results can easily be extended to Bernoulli site percolation and to lattices other than  $\mathbb{Z}^d$ . It is natural to think of the hypercubic bond percolation model as a random graph whose vertex set is  $\mathbb{Z}^d$  and whose edges are the occupied nearest neighbor bonds. The cluster  $C(x)$  of  $x$  in  $\mathbb{Z}^d$  is then defined as that connected component of this random graph which contains the vertex  $x$ . Its size or volume  $|C(x)|$  is the number of occupied bonds in  $C(x)$ . For  $x, y$  in  $\mathbb{Z}^d$ , we then define the connectivity,

$$(1.1) \quad \tau_{xy} = P(y \in C(x)) = P(x \in C(y)),$$

---

Received August 1985.

<sup>1</sup>National Science Foundation Postdoctoral Research Fellow. Research supported in part by NSF Grant No. PHY-82-03669.

<sup>2</sup>John S. Guggenheim Memorial Fellow. Research supported in part by NSF Grant No. MCS-80-19384.

AMS 1980 subject classifications. Primary 60K35; secondary 60D05.

Key words and phrases. Bernoulli percolation, invasion percolation, truncated connectivity function, cluster size distribution.

and the truncated connectivity,

$$(1.2) \quad \tau'_{xy} = P(y \in C(x) \text{ and } |C(x)| < \infty).$$

Other basic quantities of interest are the cluster size distribution,

$$(1.3) \quad P_n = P(|C(0)| = n); \quad n = 0, 1, 2, \dots,$$

the infinite cluster density,

$$(1.4) \quad P_\infty = P(|C(0)| = \infty) = 1 - \sum_{n=0}^{\infty} P_n,$$

the expected cluster size,

$$(1.5) \quad \chi = E(|C(0)|),$$

and the expected size of finite clusters,

$$(1.6) \quad \chi' = E(|C(0)| \mathbf{1}_{|C(0)| < \infty}) = \sum_{n=0}^{\infty} n P_n.$$

Our main results are stated in terms of the above definitions with  $|C(0)|$  as the number of bonds in  $C(0)$ . They also hold (with suitable changes of constants) when  $|C(0)|$  is replaced by the number of sites in  $C(0)$ .

There are many possible definitions of a threshold or critical point for  $p$ . These are nontrivial (i.e., not equal to 1 or 0) for every  $d \geq 2$ , but are rigorously known to coincide only for  $d = 2$  ([12], [14]). Two of these thresholds are

$$(1.7) \quad \pi_c = \sup\{p: \chi(p) < \infty\}$$

and

$$(1.8) \quad p_c = \sup\{p: P_\infty(p) = 0\}.$$

It is of course clear that  $\pi_c \leq p_c$ .

For  $p$  in  $(0, \pi_c)$ , there exist constants (depending on  $p$ )  $m_1$  and  $m_2$  in  $(0, \infty)$ , such that

$$(1.9) \quad \exp(-m_1 \|y - x\|) \leq \tau_{xy} = \tau'_{xy} \leq \exp(-m_2 \|y - x\|) \quad \text{for } x, y \text{ in } \mathbb{Z}^d,$$

where  $\|x\|$  denotes (e.g.) the  $l_1$  norm of  $x$  in  $\mathbb{Z}^d$ , and constants  $c_1$  and  $c_2$  in  $(0, \infty)$  such that

$$(1.10) \quad \exp(-c_1 n) \leq P_n \leq \exp(-c_2 n) \quad \text{for } n = 1, 2, \dots$$

The lower bounds of (1.9)–(1.10) are elementary. The upper bounds were originally obtained in [10] and then rederived in [13], [3].

A lower bound for  $\tau'$  as in (1.9) is easily seen to be valid for any  $p$  in  $(0, 1)$  and any  $d$ . For example, by considering a path from  $x$  to  $y$  containing exactly  $\|y - x\|$  occupied bonds and surrounded by at most  $(2d - 2)(\|y - x\| + 1) + 2$  vacant bonds, one obtains

$$(1.11) \quad \tau'_{xy} \geq (1 - p)^{2d} [p(1 - p)^{(2d-2)}]^{\|y-x\|}.$$

A lower bound for  $P_n$  above threshold is less elementary. It is known ([2]) that

for  $p$  in  $(p_c, 1)$ , there exists  $c_3 = c_3(p) < \infty$  such that

$$(1.12) \quad P_n \geq \exp(-c_3 n^{(d-1)/d}) \quad \text{for } n = 1, 2, \dots,$$

a result which was first derived for  $p$  close to 1 in [15]. We remark that this explicitly shows that an upper bound as in (1.10), i.e., an exponential tail for the cluster size distribution, cannot be valid above  $p_c$ . Furthermore, the lower bound (1.12) is valid for a large class of nonindependent percolation models ([2]).

The focus of this paper is on upper bounds complementary to (1.11)–(1.12) above threshold. Previously such results were known only for  $d = 2$  or for  $p$  close to 1. The upper bound for  $d = 2$  and  $p > \frac{1}{2}$  ( $= p_c = \pi_c$ ),

$$(1.13) \quad P_n \leq \exp(-c_4 n^{1/2}),$$

was first obtained by Kesten ([14]). For general  $d$ , it is known ([15]) that

$$(1.14) \quad P_n \leq \exp(-c_4 n^{(d-1)/d})$$

for  $p$  close to 1. In fact, it is not hard to show ([8]) that (1.14) holds for  $p > 1 - \pi_c^*$ , where  $\pi_c^*$  is the analogue of  $\pi_c$  for the dual percolation model (bonds for  $d = 2$ , plaquettes for  $d = 3$ , etc.). An argument similar to that of [8], but for  $\tau'$  rather than  $P_n$ , yields an exponential upper bound for  $\tau'$ , as in (1.9) for  $p > 1 - \pi_c^*$ . Since  $1 - \pi_c^* = 1 - \pi_c = \frac{1}{2} = \pi_c = p_c$  for  $d = 2$  ([18], [19], [12]), the two-dimensional noncritical behavior of  $\tau$ ,  $\tau'$ , and  $P_n$  are all well characterized.

In this paper, we treat dimension  $d \geq 3$  and derive upper bounds complementary to (1.11)–(1.12) and some related results above one or the other of two thresholds introduced in [1],  $p_{c,d-1}^\infty$  and  $\hat{p}_{c,2}^\infty$ , which are conjectured to coincide with  $p_c$  (and  $\pi_c$ ). To define  $p_{c,d-1}^\infty$ , first denote by  $p_{c,d-1}^K$  for  $K = 0, 1, 2, \dots$ , the percolation threshold defined as in (1.8) but with  $\mathbb{Z}^d$  replaced by  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$ ;  $p_{c,d-1}^\infty$  is then the limit of  $p_{c,d-1}^K$  as  $K \rightarrow \infty$ . The threshold  $\hat{p}_{c,2}^\infty$  is defined similarly as the limit of  $\hat{p}_{c,2}^K$ , the percolation threshold in  $(\mathbb{Z}^+)^2 \times \{0, \dots, K\}^{d-2}$ , where  $(\mathbb{Z}^+)^2$  is the positive quadrant of  $\mathbb{Z}^2$ . Our main results are as follows.

**THEOREM 1.** *For  $d > 2$  and  $p > p_{c,d-1}^\infty$ , there exists  $m = m(p) > 0$  such that*

$$(1.15) \quad \tau'_{xy} \leq \exp(-m\|y - x\|) \quad \text{for all } x, y \text{ in } \mathbb{Z}^d.$$

Our proof of Theorem 1 leads to an inequality relating the critical behavior of  $m(p)$ ,  $P_\infty(p)$ , and a quantity characterizing the approach of  $p_{c,d-1}^K$  to  $p_{c,d-1}^\infty$ . See the remark following the proof of Theorem 5 in Section 2.

As a fairly direct consequence of Theorem 1, we obtain the following result concerning the smoothness of  $P_\infty$  and  $\chi'$ .

**THEOREM 2.** *For  $d > 2$ , both  $P_\infty$  and  $\chi'$  are infinitely differentiable functions of  $p$  on  $(p_{c,d-1}^\infty, 1)$ .*

We note that infinite differentiability of  $P_\infty(p)$  on  $(p_c, 1)$  was known previously only for  $d = 2$  ([18]), although continuity of  $P_\infty(p)$  on  $(\hat{p}_{c,2}^\infty, 1)$  for  $d > 2$  follows from uniqueness of the infinite cluster ([1]) and the results of [20].

Finally, we obtain upper bounds on  $P_n$ .

**THEOREM 3.** *For  $d > 2$  and  $p > \hat{p}_{c,2}^\infty$ , there exists  $c = c(p) > 0$  such that*

$$(1.16) \quad P_n \leq \exp(-[c/(\log n)]n^{(d-1)/d}) \quad \text{for } n = 2, 3, \dots$$

The extra ingredients needed beyond Theorems 1 and 3 to complete (together with (1.11)–(1.12)) the characterization of the asymptotic behavior of  $\tau'$  and  $P_n$  above threshold would be a proof that  $\pi_c = p_c = p_{c,d-1}^\infty = \hat{p}_{c,2}^\infty$  and a replacement of  $c/(\log n)$  by a somewhat larger “constant.”

The proofs of Theorems 1 and 2 are presented in Section 2 of this paper; that of Theorem 3 in Section 3. Theorem 2 is a consequence of Theorem 1 and conventional percolation theory arguments. On the other hand, our proofs of Theorems 1 and 3 are based on the use of invasion percolation as a tool for analyzing standard Bernoulli percolation. We complete this section with some equivalent definitions of invasion percolation and a proposition which gives two simple but crucial inequalities relating the invasion and standard percolation models. For a more detailed discussion of invasion percolation, see [8] or [6].

Invasion percolation was introduced and studied numerically as a model of transport in a random medium in [16], [5], [21]. There are at least three equivalent definitions of the model ([8]), all of which are useful in obtaining results about invasion percolation itself and in relating it to standard percolation. These definitions are as follows.

(i) *Deterministic invasion.* Denote by  $\mathbb{B}_d$  the set of nearest neighbor bonds on  $\mathbb{Z}^d$ . Let  $\{W_b; b \in \mathbb{B}_d\}$  be i.i.d. random variables uniformly distributed on  $[0, 1]$ . The invasion percolation process is specified by a sequence  $\emptyset = C_0 \subset C_1 \subset C_2 \cdots$  of random subsets of  $\mathbb{B}_d$ , defined as follows.  $C_1$  is that bond  $b$  touching the origin with the minimum value of  $W_b$ . For  $n \geq 1$ ,  $C_{n+1}$  is the union of  $C_n$  with the single additional bond  $b$  in  $\partial C_n$  with the minimum value of  $W_b$ , where

$$(1.17) \quad \partial C_n \equiv \{b: b \text{ shares an endpoint with some } b' \text{ in } C_n\} \setminus C_n.$$

Denote the sequence of bonds absorbed into the growing invaded region by  $b_n$ , so that  $C_n = \{b_1, \dots, b_n\}$  for each  $n$ . Denote the sequence of  $W$  values of these absorbed bonds by

$$(1.18) \quad X_n = W_{b_n}.$$

(ii) *Dynamic growth.* This is an algorithmic procedure in which bond values are assigned to bonds in  $\mathbb{B}_d$  as needed while the invaded region is growing. Choose some deterministic ordering on  $\mathbb{B}_d$ . Let  $Y_1, Y_2, \dots$  be i.i.d. and uniformly distributed on  $[0, 1]$ . For each of the  $2d$  bonds touching the origin, assign one of  $Y_1, \dots, Y_{2d}$  according to the chosen bond order. The  $Y_i$  assigned to the bond  $b$  is then defined to be  $W_b$ . Define  $C_1$  as above. Note that there are now  $2d - 1$

bonds in  $\partial C_1$  which do not touch the origin and therefore have not been assigned values; to each of these bonds  $b$ , assign in order one of  $Y_{2d+1}, \dots, Y_{4d-1}$  and define that to be  $W_b$ . For  $n \geq 2$  define  $C_n$  as above and then define  $W_b$  for each  $b \in \partial C_n \setminus \partial C_{n-1}$  by assigning in order one of the next unassigned  $Y_i$ 's. Define  $b_n$  and  $X_n$  as above.

(iii) *Percolation cluster method.* This is an intermediate scheme which is dynamic in the sense of method (ii) but provides partial advance global information in the sense of method (i). Let  $y$  be in  $(0,1)$  and choose a random configuration  $\omega$  of occupied bonds from the standard Bernoulli bond percolation model with density  $y$ . Then, as in method (ii), values are assigned dynamically to bonds as they join the boundary of the invaded region. Here, however, if  $b$  is occupied (in  $\omega$ ), its value is assigned uniformly from  $[0, y]$ , while unoccupied bonds are assigned values uniformly from  $[y, 1]$ .

Within any of the above formulations, we define for  $0 \neq x \in \mathbb{Z}^d$ , the time  $T_x$  at which  $x$  is invaded as

$$(1.19) \quad T_x = \begin{cases} \min\{n: C_n \text{ touches } x\}, & \text{if } C_\infty \equiv \bigcup_{n=1}^\infty C_n \text{ touches } x, \\ +\infty, & \text{otherwise.} \end{cases}$$

Here,  $C$  touches  $x$  means that  $x$  is an endpoint of some  $b \in C$ . Thus  $T_x < \infty$  if and only if  $x$  is eventually invaded.  $T_0$  is defined to be 0.

The following proposition states two simple inequalities which explain why an analysis of invasion percolation should be useful in obtaining Theorems 1 and 3. In the proposition, we specifically display the  $p$  dependence of  $\tau'$  and  $P_n$ .

PROPOSITION 4. For  $d \geq 1$  and any  $p$  in  $(0, 1)$ ,

$$(1.20) \quad \tau'_{0x}(p) \leq P(T_x < \infty \text{ and } X_n > p \text{ for some } n > T_x) \quad \text{for all } x \text{ in } \mathbb{Z}^d,$$

$$(1.21) \quad P_n(p) \leq B_{n+1}(p) \equiv P(X_{n+1} > p) \quad \text{for } n = 0, 1, 2, \dots$$

PROOF. These inequalities are essentially immediate in the context of the percolation cluster method (method (iii) above) with  $y = p$ . Indeed, to prove (1.20) it suffices to show that if the invasion process is constructed on any configuration  $\omega$  which contributes to  $\tau'_{0x}(p)$ , then  $T_x < \infty$  and  $X_n > p$  for some  $n > T_x$ . To this end, suppose that, in the configuration  $\omega$ ,  $C(0)$  is finite and contains  $x$ . Then the invasion process on  $\omega$  will have  $\{b_1, \dots, b_{|C(0)|}\}$  coinciding with the bonds of  $C(0)$ , so that  $T_x \leq |C(0)| < \infty$ . Furthermore,  $b_{|C(0)|+1}$  must be a vacant bond of  $\omega$  and hence  $X_{|C(0)|+1} > p$ . Inequality (1.21) follows in the same fashion by considering Bernoulli configurations at density  $p$  in which  $|C(0)| = n$ . □

**2. Decay of the truncated connectivity.**

**PROOF OF THEOREM 1.** Let  $z \in \mathbb{Z}^d$  and assume, without loss of generality, that  $z_d \geq |z_1|, \dots, |z_{d-1}|$ . Defining  $T_N$  to be the time at which the invaded region first reaches the hyperplane  $x_d = N$ , i.e.,

$$(2.1) \quad T_N = \min\{T_x : x_d = N\},$$

we may easily bound the right-hand side of (1.20) by

$$(2.2) \quad \tau'_{0z}(p) \leq P(T_{z_d} < \infty \text{ and } X_n > p \text{ for some } n > T_{z_d}).$$

Since

$$(2.3) \quad z_d = \max_i |z_i| \geq \|z\|/d,$$

Theorem 1 is an immediate consequence of the following:

**THEOREM 5.** Suppose  $d > 2$  and  $p > p_{c,d-1}^K$  for some  $K \geq 0$ . Then for  $N \geq 0$ ,

$$(2.4) \quad P(T_N < \infty \text{ and } X_n > p \text{ for some } n > T_N) < (1 - \rho(K, p))^{N/(K+1)},$$

where  $\rho(K, p) > 0$  is the probability that  $|C(0)| = \infty$  (i.e.,  $P_\infty$ ) for standard Bernoulli bond percolation on  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$ .

**PROOF.** This argument uses a combination of the dynamic growth method and a variation of the percolation cluster method fitted to percolation in slabs. Let  $L_j$ , for  $j \in \mathbb{Z}$ , denote the slab which is a translate by  $(j - 1)(K + 1)$  of  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$ ,

$$(2.5) \quad L_j = \mathbb{Z}^{d-1} \times \{(j - 1)(K + 1), (j - 1)(K + 1) + 1, \dots, j(K + 1) - 1\}.$$

Let us denote by  $\psi_j$  the time at which  $L_j$  is first invaded and by  $S_j$  the site at which it is first invaded;  $\psi_j = \inf\{T_x : x \in L_j\}$  and  $S_j$  is the  $y$  such that  $\psi_j = T_y$  (defined only if  $\psi_j < \infty$ ). Invasion percolation may be constructed by using the percolation cluster method (with  $y = p$ ) within each slab  $L_j$  by choosing a configuration  $\omega_j$  (for Bernoulli percolation on  $L_j$ ) at the time  $\psi_j$  when it is first needed. The bond variables for bonds between  $L_j$  and  $L_{j+1}$  may be assigned according to the dynamic growth method, as they are needed. Furthermore,  $\omega_j$  may be chosen by first choosing a configuration  $\omega'_j$  from a Bernoulli percolation model on  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$  which is independent of all previously chosen  $\omega'_j$ 's and all previously assigned bond variables, and then translating in  $\mathbb{Z}^d$  so that the origin in  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$  is shifted to  $S_j$ .

Define  $D_j(p)$  as the event that in  $\omega_j$ , the cluster  $C(S_j)$  is infinite. If  $T_N < \infty$  for some  $N \geq 0$ , then  $\psi_{N'} < \infty$  for every integer  $N' \geq 0$  such that  $(N' - 1) \times (K + 1) \leq N$ . The largest such integer  $\bar{N}$  satisfies

$$(2.6) \quad \bar{N} > N/(K + 1).$$

If  $\psi_{\bar{N}} < \infty$  and  $D_j(p)$  occurs for some  $j = 1, \dots, \bar{N}$  (say  $j'$ ), then clearly  $X_n \leq p$

for every  $n \geq \psi_j$ , and hence

$$(2.7) \quad \begin{aligned} &P(T_N < \infty \text{ and } X_n > p \text{ for some } n > T_N) \\ &\leq P\left(\psi_{\bar{N}} < \infty \text{ and } \bigcap_{j=1}^{\bar{N}} (D_j(p))^c\right). \end{aligned}$$

By the above construction, conditional on  $\psi_{\bar{N}} < \infty$ ,  $D_{\bar{N}}(p)$  occurs with probability  $\rho(K, p)$  independently of the occurrence of  $D_1(p), \dots, D_{\bar{N}-1}(p)$ . Thus

$$(2.8) \quad \begin{aligned} &P\left(\psi_{\bar{N}} < \infty \text{ and } \bigcap_{j=1}^{\bar{N}} (D_j(p))^c\right) \\ &= (1 - \rho(K, p))P\left(\psi_{\bar{N}} < \infty \text{ and } \bigcap_{j=1}^{\bar{N}-1} (D_j(p))^c\right) \\ &\leq (1 - \rho(K, p))P\left(\psi_{\bar{N}-1} < \infty \text{ and } \bigcap_{j=1}^{\bar{N}-1} (D_j(p))^c\right). \end{aligned}$$

Since

$$(2.9) \quad P(\psi_1 < \infty \text{ and } (D_1(p))^c) = P((D_1(p))^c) = 1 - \rho(K, p),$$

(2.4) follows inductively from (2.6), (2.7), and (2.8). This completes the proof of Theorem 5 and hence of Theorem 1.  $\square$

**REMARK.** (2.2), (2.3), and (2.4) combine to give the explicit bound

$$(2.10) \quad \tau'_{0x} \leq (1 - \rho(K, p))^{\|x\|/(dK+d)} \leq \exp\left(-\frac{\rho(K, p)}{d(K+1)}\|x\|\right) \text{ for all } x \text{ in } \mathbb{Z}^d.$$

Let us define the inverse correlation length  $(\xi(p))^{-1}$  as the supremum of those  $m \geq 0$  such that

$$(2.11) \quad \tau'_{0x} \leq \exp(-m\|x\|) \text{ for all sufficiently large } x.$$

Then (2.11) implies that

$$(2.12) \quad \xi(p) \leq \frac{d(K+1)}{\rho(K, p)}.$$

Let us next define a length scale  $\bar{K}(p)$ , associated with the convergence of  $p_{c,d-1}^K$  to  $p_{c,d-1}^\infty$ , by

$$(2.13) \quad \bar{K}(p) = \min\{K: \rho(K, p) \geq (\frac{1}{2})\bar{P}_\infty(p)\},$$

where  $\bar{P}_\infty(p)$  is the percolation density in the half-space,  $\mathbb{Z}^{d-1} \times \{0, 1, 2, \dots\}$ . Then (2.12) yields

$$(2.14) \quad \xi(p) \leq \frac{d(\bar{K}(p) + 1)}{2\bar{P}_\infty(p)}.$$

The inequality (2.14) is nonvacuous only if the length scale  $\bar{K}(p)$  is finite. This can be verified whenever the infinite cluster in the half-space is unique and

contains an infinite path within some finite slab,  $\mathbb{Z}^{d-1} \times \{0, \dots, K_0\}$ , a condition which is known to hold whenever  $p > \hat{p}_{d,2}^\infty$  ([1]). Indeed, when this is the case, the probability that the origin is connected to an infinite path within the slab  $\mathbb{Z}^{d-1} \times \{0, \dots, K\}$  converges, as  $K \rightarrow \infty$ , to  $\bar{P}_\infty$ .

Assuming that  $p_c = p_{c,d-1}^\infty = \hat{p}_{c,2}^\infty$  and that as  $p$  approaches  $p_c$  from above,  $\xi$ ,  $\bar{P}_\infty$ , and  $\bar{K}$  exhibit power law behavior in some appropriate sense,

$$(2.15) \quad \begin{aligned} \xi(p) &\sim (p - p_c)^{-\nu'}, & \bar{P}_\infty(p) &\sim (p - p_c)^{\bar{\beta}}, \quad \text{and} \\ \bar{K}(p) &\sim (p - p_c)^{-\kappa}, \end{aligned}$$

one may obtain from (2.14) the critical exponent inequality

$$(2.16) \quad \nu' \leq \bar{\beta} + \kappa.$$

**PROOF OF THEOREM 2.** We first obtain an estimate on the tail of the finite cluster distribution as a corollary of Theorem 1. There is a positive constant  $h$ , depending only on  $d$ , such that  $|C(0)| = n$  implies that some  $x$  in  $\mathbb{Z}^d$  with  $\|x\| \geq hn^{1/d}$  belongs to  $C(0)$ . Thus, by elementary arguments and (1.15),

$$(2.17) \quad \begin{aligned} P_n &\leq P(|C(0)| < \infty \text{ and } x \in C(0) \text{ for some } x \text{ with } \|x\| \geq hn^{1/d}) \\ &\leq \sum_{\|x\| \geq hn^{1/d}} \tau'_{0x} \leq \sum_{\|x\| \geq hn^{1/d}} \exp(-m\|x\|) \\ &\leq \exp(-m_1 n^{1/d}) \quad \text{for } n = 0, 1, 2, \dots, \end{aligned}$$

where  $m_1(p) > 0$  whenever  $m(p) > 0$ . As an immediate consequence of (2.17) one has that all moments of the finite cluster size distribution are finite:

$$E(|C(0)|^k \mathbf{1}_{|C(0)| < \infty}) = \sum_{n=0}^\infty n^k P_n < \infty \quad \text{for } k = 1, 2, \dots.$$

Furthermore, it follows from (2.10), (2.17) and the monotonicity in  $p$  of  $\rho(K, p)$  that  $\sum n^k P_n$  is uniformly convergent away from  $p_{c,d-1}^\infty$ ; i.e., that for any  $\epsilon > 0$ ,

$$(2.18) \quad \lim_{N \rightarrow \infty} \sup_{p \geq p_{c,d-1}^\infty + \epsilon} \left[ \sum_{n=N}^\infty n^k P_n(p) \right] = 0, \quad \text{for } k = 1, 2, \dots.$$

The proof will be completed by using standard arguments ([18]) to show that (2.18) implies infinite differentiability of  $P_\infty$  and  $\chi$ , in fact of all moments of the finite cluster size distribution. Since

$$(2.19) \quad P_\infty(p) = 1 - \sum_{n=0}^\infty P_n(p) \quad \text{and} \quad \chi(p) = \sum_{n=0}^\infty n P_n(p),$$

it suffices by (2.18) to show that  $P_n(p)$  is infinitely differentiable and that for each  $k$ , there is some  $H_k$  and  $k'$  so that

$$(2.20) \quad \left| \left( \frac{d}{dp} \right)^k P_n(p) \right| \leq H_k n^{k'} P_n(p) \quad \text{for } n = 0, 1, \dots,$$



where  $k'$  is independent of  $p$ , and  $H_k$  is uniformly bounded on compact subsets of  $(0, 1)$ . Now if  $A$  denotes a (bond) lattice animal, i.e., a possible configuration of a finite  $C(0)$ ,  $|A|$  denotes the number of bonds in  $A$  and  $|\partial A|$  denotes the number of boundary bonds of  $A$ , then

$$(2.21) \quad P_n(p) = \sum_{A: |A|=n} p^{|A|}(1-p)^{|\partial A|},$$

and hence  $P_n$  is smooth. It is straightforward to show that

$$(2.22) \quad \left| \left( \frac{d}{dp} \right)^k p^m(1-p)^n \right| \leq H'_k(p) [\max(m, n)]^k p^m(1-p)^n,$$

where  $H'_k$  is a polynomial in  $p^{-1}$  and  $(1-p)^{-1}$  independent of  $m$  and  $n$ . Since for any  $|A|$ ,  $|\partial A| \leq 2d|A|$ , it follows that (2.20) is valid with  $H_k = (2d)^k H'_k$  and  $k' = k$ . The proof is now complete.  $\square$

**3. Upper bounds on the finite cluster size distribution.** Note that as a consequence of the truncated connectivity bound (1.15) of Theorem 1, we have already obtained in (2.17) an upper bound on the cluster size distribution. This upper bound, in view of the lower bound (1.12) and the high density upper bound (1.14), has the wrong power of  $n$  (for  $d > 2$ ). Our object in this section is to obtain the upper bound (1.16) of Theorem 3 which has the correct power of  $n$  [but also an unfortunate logarithmic factor which, as explained in Section 1, is known not to be needed for  $d = 2$  ([14]) or for  $p$  close to 1 ([15], [8])].

The proof of Theorem 3 uses an extension of the invasion percolation technique introduced above in the proof of Theorem 1 together with renormalization methods such as those of [1] and [7]. It also exploits an elegant inequality of Loomis and Whitney [17] concerning the volumes of  $(d - 1)$ -dimensional projections of subsets of  $\mathbb{Z}^d$ . Since this inequality is perfectly suited for use in percolation and other models of aggregation, we state it as a separate theorem in order to bring it to the attention of other researchers in the field. We omit the proof which involves nothing more complicated than Hölder's inequality and applies equally well to subsets of  $\mathbb{R}^d$ .

**THEOREM 6 ([17]).** *Suppose  $\Lambda \subset \mathbb{Z}^d$ ,  $d \geq 2$ . Denote by  $\Lambda_i$  for  $i = 1, \dots, d$ , the  $(d - 1)$ -dimensional projection of  $\Lambda$  perpendicular to coordinate  $i$ :*

$$(3.1) \quad \Lambda_i = \left\{ (x_1, \dots, x_{d-1}) \in \mathbb{Z}^{d-1}: (x_1, \dots, x_{i-1}, x_0, x_i, \dots, x_{d-1}) \in \Lambda \right. \\ \left. \text{for some } x_0 \in \mathbb{Z} \right\}.$$

Denote by  $|\Lambda'|$  the number of sites in  $\Lambda'$ . Then

$$(3.2) \quad \prod_{i=1}^d |\Lambda_i| \geq |\Lambda|^{d-1}.$$

We remark that the proof of Theorem 6 shows that equality holds between the two sides of (3.2) (for finite  $|\Lambda|$ ) if and only if  $\Lambda$  is a generalized rectangle,

i.e., if  $\Lambda$  is the product of its  $d$  one-dimensional projections. We further note that (3.2) implies that

$$(3.3) \quad \max_i |\Lambda_i| \geq |\Lambda|^{(d-1)/d},$$

which is the only consequence of Theorem 6 that we will use below.

**PROOF OF THEOREM 3.** By (1.21), it suffices to bound  $P(X_{n+1} > p)$  by the right-hand side of (1.16) for  $p > \hat{p}_{c,2}^\infty$ . By adjusting the constant  $c$ , it clearly suffices to obtain the bound only for large  $n$ . At time  $n$ , the invaded region  $C_n$  contains exactly  $n$  bonds and hence touches at least  $n/d$  sites in  $\mathbb{Z}^d$ . Let  $C_{n,i}$  denote the  $(d - 1)$ -dimensional projection perpendicular to coordinate  $i$  of the set of sites touched by  $C_n$ . By (3.3), at least one of these projections contains  $(n/d)^{(d-1)/d}$  or more sites. Thus

$$(3.4) \quad \begin{aligned} P(X_{n+1} > p) &= P\left(X_{n+1} > p \text{ and } \bigcup_{i=1}^d \{|C_{n,i}| \geq (n/d)^{(d-1)/d}\}\right) \\ &\leq dP(X_{n+1} > p \text{ and } |C_{n,d}| \geq (n/d)^{(d-1)/d}). \end{aligned}$$

The remainder of the proof parallels that of Theorem 5 but with two major differences. First we replace the slabs,  $L_j$  ( $j \in \mathbb{Z}$ ) of size  $(\infty)^{d-1} \times K$  used previously, by pillars,  $U_j$  ( $j \in \mathbb{Z}^{d-1}$ ) of size  $(K_n)^{d-1} \times \infty$ :

$$(3.5) \quad U_j = U_{n,j} = (\{0, \dots, K_n\}^{d-1} \times \mathbb{Z}) + j(K_n + 1, \dots, K_n + 1, 0).$$

Second, the  $K_n$ 's, unlike the previous  $K$ , will be chosen to increase with  $n$ . This will be done in such a way as to insure that the following quantity is bounded away from zero as  $n \rightarrow \infty$ :

$$(3.6) \quad \rho_n(p) \equiv \inf_{x \in U_{n,0}} P(\text{the cluster } C(x) \text{ in standard Bernoulli bond percolation on } U_{n,0} \text{ contains more than } n \text{ bonds}).$$

As in the proof of Theorem 5, it is most convenient to consider the invasion percolation process as constructed by using the percolation cluster method within each pillar, with the Bernoulli configuration assigned at the time the pillar is first invaded, and using the dynamic growth method for the bonds not within a pillar. For each  $n$ , we define  $\theta_j$  ( $= \theta_{n,j}$ ), for  $j = 1, 2, \dots$ , as the time that the  $j$ th pillar is invaded (the slabs are ordered (randomly) in the order of invasion) and  $\sigma_j$  ( $= \sigma_{n,j}$ ) as the "beachhead" point (on the surface of that  $j$ th pillar) which is invaded at time  $\theta_j$ . Thus  $\theta_1 = 0$ ,  $\sigma_1 = (0, \dots, 0)$ , and the event  $\theta_j \leq n$  is identical to the event that  $C_n$  touches at least  $j$  distinct pillars. If fewer than  $j$  pillars are eventually invaded, then  $\theta_j = \infty$  and  $\sigma_j$  is undefined.

Denote by  $C(\sigma_j)$  the Bernoulli cluster of  $\sigma_j$  within the  $j$ th pillar. A (lengthy) moment's thought yields the following inequality, analogous to (2.7):

$$(3.7) \quad P(X_{n+1} > p \text{ and } \theta_M \leq n) \leq P\left(\theta_M \leq n \text{ and } \bigcap_{j=1}^M (|C(\sigma_j)| \leq n)\right).$$

We may use (3.4) together with (3.7) to bound  $P(X_{n+1} > p)$ , since

$|C_{n,d}| \geq (n/d)^{(d-1)}$  implies that  $C_n$  touches at least  $(n/d)^{(d-1)/d}/(K_n + 1)^{d-1}$  pillars. Let us define

$$(3.8) \quad M_n = \left[ \text{greatest integer} \leq (n/d)^{(d-1)/d}/(K_n + 1)^{d-1} \right].$$

Then

$$(3.9) \quad \begin{aligned} P(X_{n+1} > p) &\leq dP\left(\theta_{M_n} \leq n \text{ and } \bigcap_{j=1}^{M_n} (|C(\sigma_j)| \leq n)\right) \\ &\leq dP\left(\theta_{M_n} < \infty \text{ and } \bigcap_{j=1}^{M_n} (|C(\sigma_j)| \leq n)\right). \end{aligned}$$

The conditional probability that  $|C(\sigma_M)| \leq n$  is *not* independent of the history of the percolation process up to time  $\theta_M$  (as was the analogous slab probability which led to the first equality of (2.8)) since it depends on the position of the beachhead point  $\sigma_M$ , but this conditional probability *is* bounded above by  $1 - \rho_n(p)$ . Thus,

$$(3.10) \quad \begin{aligned} &P\left(\theta_M < \infty \text{ and } \bigcap_{j=1}^M (|C(\sigma_j)| \leq n)\right) \\ &\leq (1 - \rho_n(p))P\left(\theta_M < \infty \text{ and } \bigcap_{j=1}^{M-1} (|C(\sigma_j)| \leq n)\right) \\ &\leq (1 - \rho_n(p))P\left(\theta_{M-1} < \infty \text{ and } \bigcap_{j=1}^{M-1} (|C(\sigma_j)| \leq n)\right) \cdots \\ &\leq (1 - \rho_n(p))^M. \end{aligned}$$

Hence by (1.21) and (3.9), we have the inequality

$$(3.11) \quad P_n(p) \leq B_{n+1}(p) \equiv P(X_{n+1} > p) \leq d(1 - \rho_n(p))^{M_n},$$

where  $\rho_n(p)$  is defined in (3.5)–(3.6) and  $M_n$  is defined in (3.8). If  $K_n$  is chosen so that for some  $\bar{K} < \infty$ ,

$$(3.12) \quad K_n \leq \bar{K}(\log n)^{1/(d-1)} \text{ for large } n,$$

then  $M_n$  will be eventually larger than  $c'n^{(d-1)/d}/(\log n)$  for some  $c' > 0$  and consequently (3.11) will imply (1.16) for some  $c > 0$ , providing that

$$(3.13) \quad \liminf_{n \rightarrow \infty} \rho_n(p) > 0$$

for such a choice of  $K_n$ . The following theorem implies that  $K_n$  can be chosen so that (3.12) and (3.13) will be simultaneously satisfied, providing  $p > \hat{p}_{c,2}^\infty$ , and hence completes the proof of Theorem 3.  $\square$

**THEOREM 7.** *Define*

$$(3.14) \quad U_n(K) = \{0, \dots, K\}^{d-1} \times \{0, \dots, n + 1\}$$

and  $\rho_{n,p}(K, x)$  to be the probability, in standard Bernoulli bond percolation on

$U_n(K)$  with bond density  $p$ , that  $|C(x)| > n$ . If  $d \geq 3$  and  $p > \hat{p}_{c,2}^\infty$ , then there exists a sequence of positive integers,  $K_n$ , such that

$$(3.15) \quad \limsup_{n \rightarrow \infty} K_n / (\log n)^{1/(d-1)} < \infty$$

and

$$(3.16) \quad \inf_{n,x} \{ \rho_{n,p}(K_n, x) : n \geq 1, x \in \partial_d U_n(K_n) \} > 0,$$

where  $\partial_d U_n(K_n)$  denotes the  $x_d = 0$  portion of the boundary of  $U_n(K_n)$  perpendicular to the  $x_d$  axis.

**PROOF.** If  $p > \hat{p}_{c,2}^\infty$ , then there is some  $N < \infty$  such that  $p > \hat{p}_{c,2}^N$ , where  $\hat{p}_{c,2}^N$  is the percolation threshold in a quadrant layer  $\hat{L}_N = (\mathbb{Z}^+)^2 \times \{0, \dots, N\}^{d-2}$ . Define  $\Delta = \frac{1}{2}(p - \hat{p}_{c,2}^N)$  so that, at bond density  $p - \Delta$  there is also percolation in the quadrant layer  $\hat{L}_N$ .

First we will show that there exist constants  $c_1 < \infty$  and  $c_2 > 0$  (independent of  $n$ ) such that if  $K_n \geq c_1(\log n)^{1/(d-1)}$ , then for  $n \geq 1$ ,

$$(3.17) \quad \max_x \{ \rho_{n,p-\Delta}(K_n, x) : x \in \partial_d U_n(K_n) \} \geq c_2.$$

In other words, we will show that, at density  $p - \Delta$ , some point  $x$  in  $\partial_d U_n(K_n)$  has  $|C(x)| \geq n$  with uniformly positive probability (in contrast to (3.16) which says that, at density  $p$ , every point  $x$  in  $\partial_d U_n(K_n)$  has  $|C(x)| \geq n$  with uniformly positive probability). By adjusting  $c_2$ , it clearly suffices to prove (3.17) for  $n$  large.

To establish (3.17), we will rely on the construction of [1] which shows that whenever the bond density exceeds  $\hat{p}_{c,2}^N$ , it is possible to construct renormalized “block” bond events  $B_{(i,j)}^{N,L,J}$  between nearest neighbor “sites”  $i, j$  of a rescaled lattice  $V^{N,L,J}$ . Here  $V^{N,L,J}$  is a rotation of  $\hat{L}_N$  through  $45^\circ$  parallel to the  $x_1, x_2$  plane, with squares of side  $J$  as “sites” and a nearest neighbor Euclidean distance  $|i - j| = \sqrt{2}L$ . (See [1] for precise definitions.) The relevant properties of the renormalized bond events are as follows: (1) The density of block bonds may be made arbitrarily large in the sense that  $P(B_{(i,j)}^{N,L,J}) \rightarrow 1$  as  $L$  and  $J$  tend to infinity in an appropriate manner. (2) The bond events are transitive in the sense that if both  $B_{(i,j)}^{N,L,J}$  and  $B_{(j,k)}^{N,L,J}$  occur, then the sites  $i$  and  $k$  are connected by a path of occupied bonds in  $\hat{L}_N$ . (3) Although the bond events are not entirely independent, the only dependence is due to bond events that share a vertex. The relevance of this construction for our purposes is that whenever the bond density exceeds  $\hat{p}_{c,2}^N$ , it is possible to estimate the probability of certain events by choosing the scales  $L$  and  $J$  of the lattice  $V^{N,L,J}$  so that  $P(B_{(i,j)}^{N,L,J})$  is in a Peierls regime and doing estimates on (say)  $2d$  independent sublattices.

Let  $\delta > 0$  and consider the system at density  $p - \Delta$ . Choose  $L$  and  $J$  so that

$$(3.18) \quad 1 - P(B_{(i,j)}^{N,L,J}) < e^{-2d(\lambda_d + \delta)},$$

where  $\lambda_d$  is the Peierls constant for the numbers of contours composed of

$(d - 1)$  cells passing through a given  $(d - 2)$  cell. Take  $n$  much larger than  $L$ . Let  $\mathcal{C}_x$  denote the event of a crossing of the box  $U_n(K_n)$  by occupied bonds from the point  $x \in \partial_d U_n(K_n)$  to some point on the opposite (i.e.,  $x_d = n + 1$ ) boundary, so that  $\cap_x \mathcal{C}_x^c$  is the event that there is no left-right crossing of  $U_n(K_n)$ . A straightforward counting argument shows that

$$(3.19) \quad P\left(\bigcap_x \mathcal{C}_x^c\right) \leq (\text{const})[2(d - 1)nK_n^{d-2}]e^{-\delta K_n^{d-1}}.$$

(See [7], Theorem 4.1b and use the additional factor of  $2d$  in the exponent in (3.18) to guarantee an independent sublattice estimate.) Thus if

$$(3.20) \quad n \leq e^{\delta K_n^{d-1}/2},$$

then  $P(\cap_x \mathcal{C}_x^c) \leq e^{-\delta' K_n^{d-1}}$  for some  $\delta' > 0$ . However, by the Harris-FKG inequality ([11], [9]),  $P(\cap_x \mathcal{C}_x^c) \geq \prod_x P(\mathcal{C}_x^c)$ , which implies that

$$(3.21) \quad \max_x P_{p-\Delta}(\mathcal{C}_x) \geq 1 - e^{-\delta'} > 0.$$

Since the event  $\mathcal{C}_x$  insures that  $|C(x)| \geq n$ , this proves the claim (3.17).

Next, we use (3.17) to obtain a lower bound on  $\min_x P_p(\mathcal{C}_x)$ . Denote by  $\bar{x}_\Delta$  the maximizing point of (3.17) (or (3.21)) and consider any  $x \in \partial_d U_n(K_n)$ . Let  $G_{x, \bar{x}_\Delta}$  denote the event that  $x$  and  $\bar{x}_\Delta$  are connected by a path of occupied bonds within  $U_n(K_n)$ . Clearly, if the event  $G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta}$  occurs, then  $\mathcal{C}_x$  occurs; hence, it suffices to bound  $P_p(G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta})$  below. In order to do this, we will bound the probability that  $G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta}$  “almost occurs” at density  $p - \Delta$ , and then use a lemma of [1] which allows us to relate this to the probability that  $G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta}$  actually occurs at density  $p$ . The lemma is as follows:

Let  $A$  be an event which is nondecreasing in the sense of FKG, and let  $E^M(A) \supset A$  denote that set of configurations in which the event  $A$  occurs if the configuration is altered on no more than  $M$  bonds. Then

$$(3.22) \quad P_p(A) \geq \varepsilon^M P_{p-\varepsilon}(E^M(A)).$$

We will also use the easily verified fact that for any nondecreasing events,  $A$  and  $B$ ,

$$(3.23) \quad E^{M_1+M_2}(A \cap B) \subset E^{M_1}(A) \cap E^{M_2}(B).$$

We have

$$(3.24) \quad \begin{aligned} P_p(\mathcal{C}_x) &\geq P_p(G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta}) \\ &\geq \Delta^{2(d-2)^2} N P_{p-\Delta} \left( E^{2(d-2)^2 N} (G_{x, \bar{x}_\Delta} \cap \mathcal{C}_{\bar{x}_\Delta}) \right) \\ &\geq \Delta^{2(d-2)^2} N P_{p-\Delta} \left( E^{2(d-2)^2 N} (G_{x, \bar{x}_\Delta}) \right) \max_x P_{p-\Delta}(\mathcal{C}_x). \end{aligned}$$

Hence it suffices to bound

$$P_{p-\Delta} \left( E^{2(d-2)^2 N} (G_{x, \bar{x}_\Delta}) \right).$$

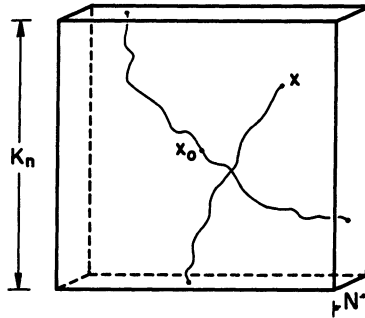


FIG. 1. Connection of  $x$  to  $x_0$  within  $T_N(K_n)$ .

To do this, we first note that if each of the points  $x$  and  $\bar{x}_\Delta$  is connected to the center  $x_0$  of  $\partial_d U_n(K_n)$  within  $U_n(K_n)$ , then  $G_{x, \bar{x}_\Delta}$  occurs. Thus

$$(3.25) \quad \begin{aligned} P_{p-\Delta} \left( E^{2(d-2)^2 N} (G_{x, \bar{x}_\Delta}) \right) &\geq P_{p-\Delta} \left( E^{2(d-2)^2 N} (G_{x, x_0} \cap G_{x_0, \bar{x}_\Delta}) \right) \\ &\geq \left[ \min_x P_{p-\Delta} \left( E^{(d-2)^2 N} (G_{x, x_0}) \right) \right]^2, \end{aligned}$$

where  $x \in \partial_d U_n(K_n)$ . Thus we must show that, with uniformly positive probability, any point in  $\partial_d U_n(K_n)$  can be connected to  $x_0$  within  $U_n(K_n)$  (by altering at most  $(d - 2)^2 N$  bonds). In fact, we will show that this occurs in the “leftmost”  $N$  layer (along the  $x_d$  axis) of  $U_n(K_n)$ :

$$(3.26) \quad T_n(K_n) = \{0, \dots, K_n\}^{d-1} \times \{0, \dots, N\}.$$

For simplicity, let us first consider  $d = 3$ , for which  $T_N(K_n)$  is simply a  $K_n \times K_n \times N$  rectangle, the  $x_d = 0$  plane of which is the square boundary  $\partial_d U_n(K_n)$  of size  $K_n \times K_n$ . Take  $x \in \partial_d U_n(K_n)$  and assume, without loss of generality, that  $x$  lies in the “upper right” quadrant of  $U_n(K_n)$  (see Figure 1).

For the remainder of this argument, it is convenient to consider  $T_N(K_n)$  as being embedded in the  $N$  layer  $L_N = \mathbb{Z}^2 \times \{0, \dots, N\}^{d-2}$ . Let  $D^+$  ( $D^-$ ) denote the event that  $x_0$  is connected to infinity within the upper left (lower right) quadrant of the  $N$  layer centered at  $x_0$ . Let  $D_x$  denote the event that  $x$  is connected to infinity within the lower left quadrant of an  $N$  layer with  $x_0$  translated to  $x$ . If  $D^+ \cap D^- \cap D_x$  occurs, the two-dimensional projection of the path from  $x$  must intersect the projection of one of the paths from  $x_0$ , so that the actual paths must be within a distance  $N$  of each other. Thus, by the Harris-FKG inequality

$$(3.27) \quad P_{p-\Delta} \left( E^N (G_{x, x_0}) \right) \geq P_{p-\Delta} (D^+ \cap D^- \cap D_x) \geq \left[ \hat{P}_\infty^{N,3}(p - \Delta) \right]^3 \quad \text{in } d = 3,$$

where  $\hat{P}_\infty^{N,d}(p - \Delta)$  is the percolation probability for the quadrant  $N$  layer  $\hat{L}_N$  in dimension  $d$ .

The construction is somewhat more intricate for  $d > 3$ . Now we can regard  $T_N(K_n)$  as consisting of  $N + 1$   $(d - 1)$ -hypercubes, each of size  $K_n^{d-1}$ , the

“leftmost” of which (along the  $x_d$  axis) is the boundary  $\partial_d U_n(K_n)$ . The object is to show that any point within  $\partial_d U_n(K_n)$  misses being connected to the center point  $x_0$  by at most  $(d - 2)^2 N$  bonds. Thus consider the point  $\tilde{x} = (\tilde{x}_1, \tilde{x}_2, \dots, \tilde{x}_{d-1}, 0) \in \partial_d U_n(K_n)$ . Let us take a “slice” of  $T_N(K_n)$  of size  $K_n^2 \times N^{d-2}$ , which contains the point  $\tilde{x}$ ,

$$(3.28) \quad S_N^{1,2}(K_n) = \{x: 0 \leq x_1; x_2 \leq K_n; \tilde{x}_3 - N \leq x_3 \leq \tilde{x}_3; \dots; \\ \tilde{x}_{d-1} - N \leq x_{d-1} \leq \tilde{x}_{d-1}; 0 \leq x_d \leq N\}.$$

In this expression, we have assumed, without loss of generality, that  $\tilde{x}_k \geq N$  for  $k = 3, \dots, d - 1$ . We can repeat the previous layer construction (see Figure 1) in an attempt to connect  $\tilde{x}$  to the central point of  $\partial_d S_N^{1,2}(K_n)$ . It again costs no more than  $[\hat{P}_\infty^{N,d}(p - \Delta)]^3$  to insure that the two-dimensional projections of the paths intersect. In this case, however, the actual paths must be within a distance  $(d - 2)N$  (= maximum distance between two points in a cube of size  $N^{d-2}$ ).

In order to “connect” the central point of  $\partial_d S_N^{1,2}(K_n)$  to  $x_0$ , we must repeat this argument on  $d - 3$  additional slices  $S_N^{i,j}(K_n)$ , orthogonal to the slice  $\partial_d S_N^{1,2}(K_n)$ , each time “connecting” the central point of the boundary of one slice to that of the next slice. Each such “connection” costs  $[\hat{P}_\infty^{N,d}(p - \Delta)]^3$  and is guaranteed to miss by no more than  $(d - 2)N$  bonds. Thus, at a total cost of no more than  $[\hat{P}_\infty^{N,d}(p - \Delta)]^{3(d-2)}$ , we can insure a connection between  $x_0$  and  $\tilde{x}$  which is missing a maximum of  $(d - 2)[(d - 2)N]$  bonds. We have

$$(3.29) \quad \min_x P_{p-\Delta} \left( E^{(d-2)^2 N}(G_{x,x_0}) \right) \geq [\hat{P}_\infty^{N,d}(p - \Delta)]^{3(d-2)},$$

which, along with (3.21), (3.24), and (3.25), provides a uniform lower bound on  $\min_x P_p(\mathcal{C}_x)$  and thus completes the theorem.  $\square$

**Acknowledgments.** The authors wish to thank H. Groemer for bringing to our attention the fact that Theorem 6 had already been obtained in [17]. We also all thank the Institut des Hautes Études Scientifiques and N. H. Kuiper, O. Lanford and D. Ruelle for their support and hospitality. Two of us (J.T.C. and L.C.) thank the Department of Mathematics and the Program in Applied Mathematics at the University of Arizona for their hospitality. J.T.C. and L.C. also thank the National Science Foundation for its support and the opportunities it has provided. One of us (C.M.N.) thanks the John S. Guggenheim Memorial Foundation and the National Science Foundation for their support.

**Note added in proof.** Although it does not directly affect our results, in this paper we distinguished between the critical points  $p_c$  and  $\pi_c$  [cf. (1.7)–(1.8)]. It has now been proved that  $p_c = \pi_c$  in all dimensions [see M. V. Menshikov, *Dokl. Akad. Nauk. SSSR* **288** 1308–1311 (1986) (in Russian); M. V. Menshikov, S. A. Molchanov and A. F. Sidorenko, *Itogi Nauki i Tekhniki (Series of Probability Theory, Mathematical Statistics, Theoretical Cybernetics)* **24** 53–110 (1986) (in Russian); M. Aizenman and D. J. Barsky, *Comm. Math. Phys.* **108** 489–526 (1987)]. Of more direct relevance is a recent proof that the quantity  $\tau'_{0,n}$  [cf. (1.2)]

has a well-defined rate of exponential decay:  $-1/\xi''(p) = \lim_{n \rightarrow \infty} n^{-1} \log \tau'_{0n}(p)$ , for which our  $m(p)$  in Theorem 1 provides a nontrivial bound (see J. T. Chayes, L. Chayes, G. R. Grimmett, H. Kesten and R. H. Schonmann, in preparation).

## REFERENCES

- [1] AIZENMAN, M., CHAYES, J. T., CHAYES, L., FRÖHLICH, J. and RUSSO, L. (1983). On a sharp transition from area law to perimeter law in a system of random surfaces. *Comm. Math. Phys.* **92** 19–69.
- [2] AIZENMAN, M., DELYON, F. and SOUILLARD, B. (1980). Lower bounds on the cluster size distribution. *J. Statist. Phys.* **23** 267–280.
- [3] AIZENMAN, M. and NEWMAN, C. M. (1984). Tree graph inequalities and critical behavior in percolation models. *J. Statist. Phys.* **36** 107–143.
- [4] BROADBENT, S. R. and HAMMERSLEY, J. M. (1957). Percolation processes. I. Crystals and mazes. *Proc. Cambridge Philos. Soc.* **53** 629–641.
- [5] CHANDLER, R., KOPLIK, J., LERMAN, K. and WILLEMSSEN, J. F. (1982). Capillary displacement and percolation in porous media. *J. Fluid Mech.* **119** 249–267.
- [6] CHAYES, J. T. and CHAYES, L. (1986). Percolation and random media. In *Critical Phenomena, Random Systems and Gauge Theories, Les Houches Session XLIII, 1984* (K. Osterwalder and R. Stora, eds.) 1000–1142. North-Holland, Amsterdam.
- [7] CHAYES, J. T. and CHAYES, L. (1986). Critical points and intermediate phases on wedges of  $\mathbb{Z}^d$ . *J. Phys. A* **19** 3033–3048.
- [8] CHAYES, J. T., CHAYES, L. and NEWMAN, C. M. (1985). The stochastic geometry of invasion percolation. *Comm. Math. Phys.* **101** 383–407.
- [9] FORTUIN, C., KASTELEYN, P. and GINIBRE, J. (1971). Correlation inequalities on some partially ordered sets. *Comm. Math. Phys.* **22** 89–103.
- [10] HAMMERSLEY, J. M. (1957). Percolation processes. Lower bounds for the critical probability. *Ann. Math. Statist.* **28** 790–795.
- [11] HARRIS, T. E. (1960). A lower bound for the critical probability in a certain percolation process. *Proc. Cambridge Philos. Soc.* **56** 13–20.
- [12] KESTEN, H. (1980). The critical probability of bond percolation on the square lattice equals  $\frac{1}{2}$ . *Comm. Math. Phys.* **74** 41–59.
- [13] KESTEN, H. (1981). Analyticity properties and power law estimates of functions in percolation theory. *J. Statist. Phys.* **25** 717–756.
- [14] KESTEN, H. (1982). *Percolation Theory for Mathematicians*. Birkhäuser, Boston.
- [15] KUNZ, H. and SOUILLARD, B. (1978). Essential singularity in percolation problems and asymptotic behavior of cluster size distribution. *J. Statist. Phys.* **19** 77–106.
- [16] LENORMAND, R. and BORIES, S. (1980). Description d'un mécanisme de connexion de liaison destiné à l'étude du drainage avec piégeage en milieu poreux. *C. R. Acad. Sci. Paris Sér. B* **291** 279–282.
- [17] LOOMIS, L. H. and WHITNEY, H. (1949). An inequality related to the isoperimetric inequality. *Bull. Amer. Math. Soc.* **55** 961–962.
- [18] RUSSO, L. (1978). A note on percolation. *Z. Wahrsch. verw. Gebiete* **43** 39–48.
- [19] SEYMOUR, P. D. and WELSH, D. J. A. (1978). Percolation probabilities on the square lattice. *Ann. Discrete Math.* **3** 227–245.
- [20] VAN DEN BERG, J. and KEANE, M. (1984). On the continuity of the percolation probability function. *Contemp. Math.* **26** 61–65.
- [21] WILKINSON, D. and WILLEMSSEN, J. F. (1983). Invasion percolation: A new form of percolation theory. *J. Phys. A* **16** 3365–3376.

J. T. CHAYES AND L. CHAYES  
DEPARTMENT OF MATHEMATICS  
UNIVERSITY OF CALIFORNIA  
LOS ANGELES, CALIFORNIA 90024

C. M. NEWMAN  
DEPARTMENT OF MATHEMATICS  
UNIVERSITY OF ARIZONA  
TUCSON, ARIZONA 85721