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Sample path MDP for the current and the tagged particle in the SSEP*

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Abstract

We prove sample path moderate deviation principles (MDP) for the current and the tagged particle in the symmetric simple exclusion process, which extends the results in [39], where the MDP was only proved at any fixed time.

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1 Introduction

The exclusion process is informally defined as a family of indistinguishable particles performing random walks on some graph subjected to the exclusion rule, *i.e.*, there is at most one particle at each site. It has been a long-standing problem to investigate the behaviors of a typical particle in the exclusion process, which is called the tagged particle. Due to the interactions with the other particles, the replacement of the tagged particle itself is not a Markov process. Despite that, much progress has been made since Spitzer introduced the exclusion process in [36]. We refer the readers to [24, 21] and references therein for an excellent survey on the tagged particle process.

The model is usually called the symmetric simple exclusion process (SSEP) if the underlying graph is the one-dimensional infinite lattice \mathbb{Z} and a particle jumps to its left and right neighbors at equal rates 1/2. Since the total number of particles is conserved by the dynamics, the Bernoulli product measure with constant density $\rho \in [0, 1]$, denoted by ν_{ρ} , is invariant for the SSEP, see [23] for example. Assume the initial distribution of the SSEP is the Bernoulli product measure ν_{ρ} conditioned on having a particle at

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the origin, and let X(t) be the position of this tagged particle at time t. Since particles cannot jump over each other in the SSEP, the tagged particle turns out to be sub-diffusive. Precisely speaking, Arratia [1] and De Masi and Ferrari [5] showed that

$$\frac{X_t}{t^{1/4}} \Rightarrow \mathcal{N}(0, \sigma_X^2), \quad t \to +\infty,$$

where $\mathcal{N}(0, \sigma_X^2)$ is the normal distribution with mean zero and variance $\sigma_X^2 := \sqrt{2/\pi}(1-\rho)/\rho$. The above central limit theorem was then extended to the invariance principle by Peligrad and Sethuraman [27],

$$\left\{\frac{X_{tN^2}}{N^{1/2}}: 0 \le t \le T\right\} \Rightarrow \left\{B_t^{1/4}: 0 \le t \le T\right\}, \quad N \to +\infty,$$

where $B^{1/4}_{\cdot}$ is the fractional Brownian motion with Hurst index 1/4 and T > 0 is fixed. Jara and Landim [15] also proved central limit theorems for the tagged particle when the initial distribution of the SSEP is non-equilibrium. Recently, Conroy and Sethuraman [4] showed that the scaled position of the tagged particle converges to a Gumble limit law when the process starts from the step configuration.

Large and moderate deviation behaviors of the tagged particle in the SSEP have also been studied. The large deviation behaviors were investigated in [33, 12]. When the initial distribution of the SSEP is the Bernoulli product measure ν_{ρ} conditioned on having a particle at the origin, the same authors [39] proved that for any t > 0, the sequence $\{X(tN^2)/a_N\}_{N\geq 1}$ satisfies moderate deviation principles with decay rate a_N^2/N and with rate function $\sqrt{2\pi\rho\alpha^2}/[4(1-\rho)\sqrt{t}]$, where $\sqrt{N} \ll a_N \ll N$.

We comment that the above deviation results only considered the behaviors of the tagged particle at any fixed time t. The main aim of this paper is to investigate sample path moderate deviations for the tagged particle. Roughly speaking, we proved that when the SSEP starts from the same initial distribution as in [39], the sequence of processes $\{X(tN^2)/a_N : 0 \le t \le T\}_{N\ge 1}$ satisfies moderate deviation principles with decay rate a_N^2/N , where $\sqrt{N \log N} \ll a_N \ll N$ and T > 0 is fixed. Moreover, the moderate deviation rate function is proportional to the large deviation rate function of the fractional Brownian motion with Hurst index 1/4. See Theorem 2.2 for details.

Since the relative ordering of the particles is conserved in the SSEP, closely related to the position of the tagged particle is the current of the process. This current-tagged particle relationship has been used in the above literature [5, 15, 33, 39]. To prove Theorem 2.2, we also proved sample path moderate deviation principles for the current, which is the content of Theorem 2.1 and has its own interest.

The main idea of the proof is as follows. As in [39], by using moderate deviation principles from hydrodynamic limits of the SSEP [10], we first prove finite-dimensional moderate deviation principles for the current and the tagged particle, where the rate function is given by a variational formula. In [39], this variational problem was solved by constructing a minimizer of the problem directly and cannot be extended to the finite-dimensional case straightforwardly. Instead, we use the Fourier approach and then obtain an explicit formula for the finite-dimensional deviation rate function. Then, we need to prove the exponential tightness of the sample path of the current and the tagged particle, which requires more delicate analysis compared with [39] since all the estimates should be uniform in time. Finally, using standard results from large deviations theory, we obtain another variational formula for the rate function of the sample path, which could be solved easily from known results on fractional Brownian motion.

Very recently, the first named author in [38] proved non-equilibrium moderate deviations from hydrodynamic limits of the SSEP. It remains a challenge to extend the results in this paper to the non-equilibrium case by using the above results. Another possible direction is to consider the sample path moderate deviation principles for the tagged particle in the SSEP in random environments or in the asymmetric exclusion process. We leave those as future works.

1.1 Related literature

For the tagged particle in exclusion processes, Saada [29] and Rezakhanlou [28] proved the law of large numbers. When the process starts from its equilibrium measure, except the literature mentioned above, central limit theorems and invariance principles were proved by Kipnis and Varadhan [20] in the symmetric case in all dimensions except the one-dimensional nearest neighbor case, by Varadhan [37] in the asymmetric meanzero case, by Kipnis [18] in the one-dimensional nearest neighbor asymmetric case and by Sethuraman, Varadhan and Yau [35] in the asymmetric case in three or higher dimensions. In dimensions $d \leq 2$ when the underlying random walk has a drift except for the one-dimensional nearest-neighbor case, a full CLT or invariance principle remains open. However, Sethuraman [31] proved that the tagged particle is diffusive in this case.

Jara [13] and the second named author [40] proved CLT for the tagged particle when particles perform random walks with long jumps. The tagged particle in the SSEP with bond disorder was considered in [16]. In [22, 25, 26], the authors considered the tagged particle process in order to check the validity of the Einstein relation. The heat kernel bound for the tagged particle in the SSEP was also obtained in [11]. In [30, 34], the authors investigated large deviations of the tagged particle in the asymmetric exclusion process. Besides on the integer lattice \mathbb{Z}^d , the tagged particle was also considered on regular trees [2] and on Galton–Watson trees [9].

The behavior of the tagged particle has also been investigated in other interacting particle processes, such as in zero range processes [14, 17, 32] or in stirring-exclusion processes [3].

1.2 Notation

For integers m, n > 0 and $U, V \subset \mathbb{R}$, let $\mathcal{C}^{m,n}(U \times V)$ be the space of functions on $U \times V$ which are continuously m (resp. n) times differentiable on the first (resp. second) variable, and let $\mathcal{C}_c^{m,n}(U \times V)$ be those functions in $\mathcal{C}^{m,n}(U \times V)$ with compact support. Let $\mathcal{S}(\mathbb{R})$ be the space of Schwartz functions on \mathbb{R} and let $\mathcal{S}'(\mathbb{R})$ be its dual space, *i.e.*, the space of tempered distributions. For a metric space $(\Sigma, d(\cdot, \cdot))$, let $\mathcal{D}([0, \infty), \Sigma)$ be the space of càdlàg functions equipped with the Skorohod topology.

For any column vector $\boldsymbol{\alpha} = (\alpha_1, \alpha_2, \dots, \alpha_n)^T \in \mathbb{R}^n$, denote by $\|\boldsymbol{\alpha}\|_{\infty} := \sup_{1 \leq j \leq n} |\alpha_j|$ the uniform norm of $\boldsymbol{\alpha}$. For $a \in \mathbb{R}$, let [a] be the largest integer smaller than or equal to a. For two positive sequences $\{a_N\}$ and $\{b_N\}$, we write $a_N \ll b_N$ if $\lim_{N \to \infty} a_N/b_N = 0$.

1.3 Outline of the paper

In Section 2, we introduce the model and state the main results of the paper. We recall moderate deviations from hydrodynamic limits of the SSEP and prove several basic properties of the rate function from hydrodynamic limits in Section 3. Section 4 is devoted to identifying the rate function of the tagged particle or the current with a variational formula. We prove the exponential tightness of the tagged particle and the current in Section 5. Finally, the proof of Theorems 2.1 and 2.2 is presented in Section 6. To make the paper easier to follow, we put the tedious calculations in Section A.

2 Model and results

2.1 Model

The state space of the one-dimensional symmetric simple exclusion process (SSEP) is $\Omega = \{0, 1\}^{\mathbb{Z}}$. For a configuration $\eta \in \Omega$, $\eta(x) = 1$ if and only if there is one particle at site x. The infinitesimal generator of the process acts on local functions $f : \Omega \to \mathbb{R}$ as

$$\mathscr{L}f(\eta) = \frac{1}{2} \sum_{x \in \mathbb{Z}} \{ f(\eta^{x,x+1}) - f(\eta) \}.$$
 (2.1)

Above, we call f a local function if it depends on η only through a finite number of sites, and $\eta^{x,y}$ is the configuration obtained from η by exchanging the values of $\eta(x)$ and $\eta(y)$, that is,

$$\eta^{x,y}(z) = \begin{cases} \eta(z) & \text{if } z \neq x, y; \\ \eta(y) & \text{if } z = x; \\ \eta(x) & \text{if } z = y. \end{cases}$$

For each $\rho \in [0,1]$, let ν_{ρ} be the Bernoulli product measure on Ω with marginals given by

$$\nu_{\rho}\{\eta:\eta(x)=1\}=\rho, \quad \forall x\in\mathbb{Z}.$$

It is well known that ν_{ρ} is reversible and ergodic for the SSEP, *cf.* [23] for example.

Denote by $\{\eta_t\}_{t\geq 0}$ the Markov process with generator \mathscr{L} . For a probability measure μ on Ω , denote by \mathbb{P}_{μ} the measure on the path space $\mathcal{D}([0,\infty),\Omega)$ of the process η_t with initial measure μ , and by \mathbb{E}_{μ} the corresponding expectation.

2.2 MDP for the current

We first recall basic facts of the fractional Brownian motion. Let $\{B_t^{1/4}\}_{t\geq 0}$ be the fractional Brownian motion with Hurst parameter 1/4, which is a mean-zero continuous-time Gaussian process with covariance given by

$$\operatorname{Cov}\left(B_t^{1/4}, B_s^{1/4}\right) = \frac{1}{2}\left(t^{1/2} + s^{1/2} - |t - s|^{1/2}\right) =: a(t, s)$$

for any $s,t \ge 0$. It is shown in [6] that the fractional Brownian motion has the following representation,

$$B_t^{1/4} = \int_0^t \mathcal{K}(t,s) dB_s, \quad \forall t \ge 0,$$
(2.2)

where $\{B_t\}_{t\geq 0}$ is the standard one-dimensional Brownian motion starting from the origin, and the kernel function is explicitly given by

$$\mathcal{K}(t,s) = \frac{(t-s)^{-1/4}}{\sqrt{V}\Gamma(3/4)} F(1/4, -1/4, 3/4, 1-\frac{t}{s})$$

for any $0 \le s < t$, where

$$V = \frac{8\Gamma(3/2)\cos(\pi/4)}{\pi} \quad \text{and} \quad F(\alpha, \beta, \gamma, z) = \sum_{k=0}^{+\infty} \frac{(\alpha)_k(\beta)_k}{(\gamma)_k k!} z^k$$

with $(a)_k$ defined as $\Gamma(a+k)/\Gamma(a)$. Moreover, it is easy to see that for any $s, t \ge 0$,

$$\int_0^{t\wedge s} \mathcal{K}(t,\tau) \mathcal{K}(s,\tau) d\tau = a(t,s).$$

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Fix a time horizon T > 0. Let \mathcal{H} be the set of càdlàg functions $f : [0,T] \to \mathbb{R}$ such that there exists a function $h_f \in L^2[0,T]$ satisfying

$$f(t) = \int_0^t \mathcal{K}(t,s) h_f(s) ds, \quad \forall 0 \le t \le T.$$

For any $f \in \mathcal{D}([0,T],\mathbb{R})$, define

$$\mathcal{I}_{path}(f) = \begin{cases} \frac{1}{2} \int_0^T (h_f(s))^2 ds & \text{if } f \in \mathcal{H}, \\ +\infty & \text{otherwise.} \end{cases}$$

By (2.2) and [7, Theorem 3.4.12], $\mathcal{I}_{path}(f)$ is the large deviation rate function of the sequence of the processes $\{\frac{1}{\sqrt{N}}B_t^{1/4}: 0 \le t \le T\}_{N \ge 1}$.

For each $x \in \mathbb{Z}$, let $J_{x,x+1}(t)$ be the net number of particles across the bond (x, x+1) up to time t, *i.e.*, the number of particles jumping from x to x + 1 minus the number of particles jumping from x + 1 to x during the time interval [0, t].

The first result concerns moderate deviations for the path of the current when the process starts from the initial measure ν_{ρ} for some $\rho \in (0, 1)$. Throughout the article, we assume $\{a_N\}_{N\geq 1}$ to be a sequence of positive numbers such that

$$\lim_{N \to \infty} \frac{a_N}{N} = \lim_{N \to \infty} \frac{N \log N}{a_N^2} = 0.$$

Theorem 2.1. Fix $\rho \in (0,1)$. For $f \in \mathcal{D}([0,T],\mathbb{R})$, define

$$\mathcal{I}_{cur}(f) = \frac{1}{\sigma_J^2} \mathcal{I}_{path}(f)$$

where $\sigma_J^2 := \sigma_J^2(\rho) := \sqrt{\frac{2}{\pi}} \rho(1-\rho)$. Then, the sequence of the processes

$$\left\{\frac{1}{a_N}J_{-1,0}(tN^2): 0 \le t \le T\right\}_{N \ge 1}$$

satisfies moderate deviation principles with rate a_N^2/N and with rate function \mathcal{I}_{cur} . Precisely speaking, for any closed set $\mathcal{C} \subseteq \mathcal{D}([0,T],\mathbb{R})$,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \frac{1}{a_N} J_{-1,0}(tN^2) : 0 \le t \le T \right\} \in \mathcal{C} \right) \le -\inf_{f \in \mathcal{C}} \mathcal{I}_{cur}(f),$$

and for any open set $\mathcal{O} \subseteq \mathcal{D}([0,T],\mathbb{R})$,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \frac{1}{a_N} J_{-1,0}(tN^2) : 0 \le t \le T \right\} \in \mathcal{O} \right) \ge -\inf_{f \in \mathcal{O}} \mathcal{I}_{cur}(f).$$

2.3 MDP for the tagged particle

Note that the particles in the SSEP are indistinguishable. In this subsection, we distinguish one particular particle, put it initially at the origin, and call it the *tagged* particle. Denote by X(t) the position of the tagged particle at time t. By convention, X(0) = 0. Since particles cannot jump over each other, the relative ordering of the particles is conserved.

Observe that the tagged particle process $\{X(t)\}_{t\geq 0}$ itself is not Markovian. However, both the coupled process $\{(\eta_t, X(t))\}_{t\geq 0}$ and the environment process $\{\xi_t\}_{t\geq 0}$ are Markovian, where $\xi_t(x) = \eta_t(X(t) + x)$ for $x \in \mathbb{Z}$ and $t \geq 0$.

For $\rho \in (0,1)$, let ν_{ρ}^{*} be the measure obtained from the Bernoulli product measure ν_{ρ} conditioned on having a particle at the origin, *i.e.*,

$$\nu_{\rho}^{*}(\cdot) = \nu_{\rho}(\cdot | \eta(0) = 1).$$

It is also well known that ν_{ρ}^* is invariant and ergodic for the environment process $\{\xi_t\}_{t\geq 0}$.

The second result of the article concerns moderate deviations for the sample path of the tagged particle when the process starts from the initial measure ν_{ρ}^{*} for some $\rho \in (0, 1)$.

Theorem 2.2. Fix $\rho \in (0,1)$. For $f \in \mathcal{D}([0,T],\mathbb{R})$, define

$$\mathcal{I}_{tag}(f) = \frac{1}{\sigma_X^2} \mathcal{I}_{path}(f)$$

where $\sigma_X^2 = \sqrt{\frac{2}{\pi}} \frac{1-\rho}{\rho}$. Then, the sequence of the tagged particle processes

$$\left\{\frac{1}{a_N}X(tN^2): 0 \le t \le T\right\}_{N \ge 0}$$

satisfies moderate deviation principles with rate a_N^2/N and with rate function \mathcal{I}_{tag} . Precisely speaking, for any closed set $\mathcal{C} \subseteq \mathcal{D}([0,T],\mathbb{R})$,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(\left\{ \frac{1}{a_N} X(tN^2) : 0 \le t \le T \right\} \in \mathcal{C} \right) \le -\inf_{f \in \mathcal{C}} \mathcal{I}_{tag}(f),$$

and for any open set $\mathcal{O} \subseteq \mathcal{D}([0,T],\mathbb{R})$,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \Big(\left\{ \frac{1}{a_N} X(tN^2) : 0 \le t \le T \right\} \in \mathcal{O} \Big) \ge -\inf_{f \in \mathcal{O}} \mathcal{I}_{tag}(f).$$

3 Preliminary estimates

3.1 Moderate deviations from hydrodynamic limits

We first recall moderate deviation principles for the empirical measure of the SSEP when the process starts from its invariant measure ν_{ρ} for some $\rho \in (0,1)$, see [10] for details. For each $N \ge 1$, the centered empirical measure $\mu_t^N \in \mathcal{S}'(\mathbb{R})$ of the SSEP, which acts on Schwartz functions $G \in \mathcal{S}(\mathbb{R})$, is defined as

$$\langle \mu_t^N, G \rangle = \frac{1}{a_N} \sum_{x \in \mathbb{Z}} (\eta_{tN^2}(x) - \rho) G(x/N).$$

In this subsection, we assume the sequence of positive numbers $\{a_N\}_{N\geq 1}$ satisfies

$$\lim_{N \to \infty} \frac{a_N}{N} = \lim_{N \to \infty} \frac{N}{a_N^2} = 0.$$

Let $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ be the set of càdlàg functions from [0,T] to $\mathcal{S}'(\mathbb{R})$ endowed with the topology under which $\lim_{n\to+\infty} \theta^n = \theta$ in $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ if and only if

$$\{\theta_t^n(G)\}_{0 \le t \le T} \to \{\theta_t(G)\}_{0 \le t \le T}$$

in $\mathcal{D}([0,T],\mathbb{R})$ for any $G \in \mathcal{S}(\mathbb{R})$. For $G \in \mathcal{C}_c^{1,\infty}([0,T] \times \mathbb{R})$ and $\mu \in \mathcal{D}([0,T],\mathcal{S}'(\mathbb{R}))$, define the linear functional $l(\cdot, \cdot)$ as

$$l(\mu, G) = \langle \mu_T, G_T \rangle - \langle \mu_0, G_0 \rangle - \int_0^T \langle \mu_s, \left(\partial_s + (1/2) \partial_u^2 \right) G_s \rangle ds$$

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Above, $G_s(u) = G(s, u)$ and $\langle \cdot, \cdot \rangle$ is the bilinear form on $\mathcal{S}'(\mathbb{R}) \times \mathcal{S}(\mathbb{R})$ such that $\langle \mu, g \rangle = \mu(g)$ is the output of μ acting on g for $\mu \in \mathcal{S}'(\mathbb{R})$ and $g \in \mathcal{S}(\mathbb{R})$. The rate function $\mathcal{Q}_T(\mu) := \mathcal{Q}_{T,dyn}(\mu) + \mathcal{Q}_0(\mu_0)$ is defined as

$$\begin{aligned} \mathcal{Q}_{T,dyn}(\mu) &= \sup_{G \in \mathcal{C}_c^{1,\infty}([0,T] \times \mathbb{R})} \left\{ l(\mu,G) - \frac{\chi(\rho)}{2} \int_0^T \int_{\mathbb{R}} (\partial_u G)^2(s,u) \, du \, ds \right\}, \\ \mathcal{Q}_0\left(\mu_0\right) &= \sup_{\phi \in \mathcal{C}_c^{\infty}(\mathbb{R})} \left\{ \langle \mu_0, \phi \rangle - \frac{\chi(\rho)}{2} \int_{\mathbb{R}} \phi^2(u) du \right\}, \end{aligned}$$

where $\chi(\rho) = \rho(1-\rho)$ is the compressibility of the SSEP.

Theorem 3.1 ([10]). Let $\rho \in (0,1)$. The sequence of processes $\{\mu_t^N : 0 \le t \le T\}_{N\ge 1}$ satisfies the MDP with decay rate a_N^2/N and with rate function $\mathcal{Q}_T(\mu)$. More precisely, for any closed set $\mathcal{C} \subseteq \mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$,

$$\limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\{ \mu_t^N : 0 \le t \le T \} \in \mathcal{C} \right) \le - \inf_{\mu \in \mathcal{C}} \mathcal{Q}_T(\mu),$$

and for any open set $\mathcal{O} \subseteq \mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$,

$$\liminf_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\{ \mu_t^N : 0 \le t \le T \} \in \mathcal{O} \right) \ge - \inf_{\mu \in \mathcal{O}} \mathcal{Q}_T(\mu).$$

Remark 3.2. Let η_{\cdot} and η_{\cdot}^* be two SSEPs with initial distributions ν_{ρ} and ν_{ρ}^* respectively. By the basic coupling (see [23] for example), one could couple η_{\cdot} and η_{\cdot}^* together such that $\eta_t(x) = \eta_t^*(x)$ for all but one x. As a consequence, the associated empirical measures of the two processes satisfy

$$|\langle \mu_t^N - \mu_t^{N,*}, G \rangle| \le \frac{\|G\|_{\infty}}{a_N}, \quad \forall t \ge 0, \; \forall G \in \mathcal{S}(\mathbb{R}).$$

Thus, the above theorem also holds under $\mathbb{P}_{\nu_{\alpha}^*}$.

3.2 Properties

In this subsection, we establish several properties of the measure μ when it satisfies $Q_T(\mu) < \infty$. Before that, we first introduce some notation. For $H, G \in \mathcal{C}_c^{1,2}([0,T] \times \mathbb{R})$, define

$$[H,G] := \int_0^T \int_{\mathbb{R}} \partial_u H(t,u) \partial_u G(t,u) \, du \, dt.$$

In order to make $[\cdot, \cdot]$ an inner product, we say $H \sim G$ if [H - G, H - G] = 0. Let \mathcal{H}^1 be the Hilbert space obtained as the completion of $\mathcal{C}_c^{1,2}([0,T] \times \mathbb{R}))/\sim$ with respect to the inner product $[\cdot, \cdot]$.

Lemma 3.3. Let $\mu \in S'(\mathbb{R})$. If $Q_T(\mu) < \infty$, then there exist functions $\psi \in L^2(\mathbb{R})$ and $H \in \mathcal{H}^1$ such that, for any $\phi \in S(\mathbb{R})$,

$$\left\langle \mu_{0},\phi\right\rangle =\int_{\mathbb{R}}\psi(u)\phi(u)du,$$

and that μ is the unique weak solution of the following PDE

$$\begin{cases} \partial_t \mu(t, u) = (1/2) \partial_u^2 \mu(t, u) - \chi(\rho) \partial_u^2 H(t, u), & t > 0, \ u \in \mathbb{R}, \\ \mu(0, u) = \psi(u), & u \in \mathbb{R}. \end{cases}$$
(3.1)

In particular, $\mu = \mu(t, u)$ is a function. Moreover,

$$Q_0(\mu_0) = \frac{\|\psi\|_{L^2(\mathbb{R})}^2}{2\chi(\rho)}, \quad Q_{T,dyn}(\mu) = \frac{\chi(\rho)}{2}[H,H]$$

The proof of the above lemma follows directly from Riesz's representation theorem. We refer the readers to Appendix A.1 for its proof.

In the above lemma, we say μ is a weak solution to (3.1) if for any $G \in \mathcal{C}_c^{1,2}([0,T] \times \mathbb{R})$ and for any $0 < t \leq T$,

$$\int_{\mathbb{R}} \mu(t, u) G(t, u) du = \int_{\mathbb{R}} \psi(u) G(0, u) du + \int_{0}^{t} \int_{\mathbb{R}} \mu(s, u) (\partial_{s} + (1/2)\partial_{u}^{2}) G(s, u) du ds + \chi(\rho) \int_{0}^{t} \int_{\mathbb{R}} \partial_{u} G(s, u) \partial_{u} H(s, u) du ds.$$

Actually, μ has the following explicit expression: for any $t \ge 0$ and for any $u \in \mathbb{R}$,

$$\mu(t,u) = \int_{\mathbb{R}} p_t(u-v)\psi(v)dv - \chi(\rho) \int_0^t \int_{\mathbb{R}} p'_{t-s}(u-v)\partial_v H(s,v)\,dv\,ds,$$
 (3.2)

where $p_t(u)$ is the heat kernel

$$p_t(u) = \frac{1}{\sqrt{2\pi t}} e^{-u^2/(2t)}$$

and $p'_t(u) := \frac{\partial}{\partial u} p_t(u)$.

Using the explicit formula of μ , we have the following lemma, which states that the macroscopic current is well-defined and can be explicitly characterized if the MDP rate function of the fluctuation fields is finite. Below, we shall approximate μ by a sequence $\{\mu_n\}$ in \mathcal{A} (introduced in Subsection 3.3), and the following lemma allows us to calculate the limit of the current associated to the sequence $\{\mu^n\}$, see the proof of Proposition 4.2 and Lemma 6.3.

Lemma 3.4. If $Q_T(\mu) < \infty$, then the limit

$$\int_0^\infty [\mu(t, u) - \mu(0, u)] du := \lim_{M \to \infty} \int_0^M [\mu(t, u) - \mu(0, u)] du$$

is well-defined. Moreover,

$$\int_0^\infty [\mu(t,u) - \mu(0,u)] du = \int_0^t \int_{\mathbb{R}} \left[-\frac{1}{2} p_s'(v) \psi(v) + \chi(\rho) p_{t-s}(v) \partial_v H(s,v) \right] dv \, ds.$$

Proof. By (3.2),

$$\int_0^M [\mu(t, u) - \mu(0, u)] du = \mathrm{I} - \mathrm{II},$$

.

where

$$I = \int_0^M \int_{\mathbb{R}} \left[p_t(u-v) - p_0(u-v) \right] \psi(v) \, dv \, du,$$

$$II = \chi(\rho) \int_0^t \int_0^M \int_{\mathbb{R}} p'_{t-s}(u-v) \partial_v H(s,v) \, dv \, du \, ds$$

For the first term, we write it as

$$\begin{split} \mathbf{I} &= \int_0^M \int_{\mathbb{R}} \int_0^t \dot{p}_s(u-v)\psi(v) \, ds \, dv \, du = \frac{1}{2} \int_0^M \int_{\mathbb{R}} \int_0^t p_s''(u-v)\psi(v) \, ds \, dv \, du \\ &= \frac{1}{2} \int_{\mathbb{R}} \int_0^t \left[p_s'(M-v) - p_s'(v) \right] \psi(v) \, ds \, dv. \end{split}$$

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Above, $\dot{p}_s(u) := \frac{\partial}{\partial s} p_s(u)$ and we used the identity $\dot{p}_s(u) = (1/2) p_s''(u)$. For the second term,

$$II = \chi(\rho) \int_0^t \int_{\mathbb{R}} \left[p_{t-s}(M-v) - p_{t-s}(v) \right] \partial_v H(s,v) \, dv \, ds.$$

Thus, we conclude the proof once we show that

$$\lim_{M \to \infty} \int_{\mathbb{R}} \left(\int_0^t p'_s(M-v) \, ds \right) \psi(v) \, dv = 0, \tag{3.3}$$

$$\lim_{M \to \infty} \int_0^t \int_{\mathbb{R}} p_{t-s}(M-v) \partial_v H(s,v) \, dv \, ds = 0.$$
(3.4)

One could check directly that, for any t > 0, $p_s(u) \in L^2([0, t] \times \mathbb{R})$ and

$$\int_0^t p'_s(\cdot) \, ds \in L^2(\mathbb{R}). \tag{3.5}$$

We refer the readers to Appendix A.2 for the proof of (3.5). Finally, (3.3) and (3.4) follow from the following claim: if $f, g \in L^2(\mathbb{R})$, then

$$\lim_{M \to \infty} \int_{\mathbb{R}} f(M-u)g(u)du = 0.$$

The similar result holds when $f, g \in L^2([0, t] \times \mathbb{R})$. Indeed, by splitting

$$\int_{\mathbb{R}} f(M-u)g(u)du = \int_{|u| > M/2} f(M-u)g(u)du + \int_{|u| \le M/2} f(M-u)g(u)du,$$

and by Cauchy-Schwarz inequality, we bound

$$\Big|\int_{\mathbb{R}} f(M-u)g(u)du\Big| \le \|f\|_{L^{2}(\mathbb{R})} \Big(\int_{|u|>M/2} g(u)^{2} du\Big)^{1/2} + \|g\|_{L^{2}(\mathbb{R})} \Big(\int_{u>M/2} f(u)^{2} du\Big)^{1/2},$$

which converges to zero as $M \to \infty$. This concludes the proof.

Remark 3.5. By repeating the above arguments, one could also show that if $Q_T(\mu) < \infty$, then for any $0 < t \leq T$,

$$\lim_{M \to \infty} \frac{1}{M} \int_0^M u[\mu(t, u) - \mu(0, u)] \, du = 0.$$

Remark 3.6. Using (3.2) and Lemma 3.3, it is not difficult to show that the moderate deviation rate function $Q_T(\cdot)$ is good. Here we give an outline of the check of this property. The property that $Q_T(\cdot)$ is lower semi-continuous follows directly from the fact that $Q_T(\cdot)$ is the supremum of a class of continuous functions from $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ to \mathbb{R} . Consequently,

$$\{\mu: \mathcal{Q}_T(\mu) \leq c\} := \mathbb{L}_c$$

is closed for any c > 0. Hence, we only need to show that \mathbb{L}_c is compact. Using (3.2), Lemma 3.3, the formula of integration by parts, Cauchy-Schwarz inequality and $\partial_t p(t, u) = \frac{1}{2} \partial_u^2 p(t, u)$, we have

$$\begin{aligned} \left| \left\langle \mu_{t}, f \right\rangle \right| &\leq \|\psi\|_{L^{2}(\mathbb{R})} \|f\|_{L^{2}(\mathbb{R})} + \chi(\rho)\sqrt{T} \|f'\|_{L^{2}(\mathbb{R})}\sqrt{[H,H]} \\ &= \sqrt{2\chi(\rho)\mathcal{Q}_{0}(\mu_{0})} \|f\|_{L^{2}(\mathbb{R})} + \sqrt{T} \|f'\|_{L^{2}(\mathbb{R})}\sqrt{2\chi(\rho)\mathcal{Q}_{T,dyn}(\mu)} \\ &\leq 2\sqrt{\chi(\rho)\mathcal{Q}_{T}(\mu)} \left(\|f\|_{L^{2}(\mathbb{R})} + \sqrt{T} \|f'\|_{L^{2}(\mathbb{R})} \right) \end{aligned}$$

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and

$$\begin{aligned} \left| \left\langle \mu_{t}, f \right\rangle - \left\langle \mu_{s}, f \right\rangle \right| \\ &\leq \frac{t-s}{2} \|f''\|_{L^{2}(\mathbb{R})} \|\psi\|_{L^{2}(\mathbb{R})} + \chi(\rho)\sqrt{t-s}\|f'\|_{L^{2}(\mathbb{R})}\sqrt{[H,H]} \\ &\quad + \frac{\chi(\rho)}{2}(t-s)\sqrt{T}\|f'''\|_{L^{2}(\mathbb{R})}\sqrt{[H,H]} \\ &\leq \sqrt{2\mathcal{Q}_{T}(\mu)\chi(\rho)} \left(\frac{t-s}{2}\|f''\|_{L^{2}(\mathbb{R})} + \sqrt{t-s}\|f'\|_{L^{2}(\mathbb{R})} + \frac{t-s}{2}\sqrt{T}\|f'''\|_{L^{2}(\mathbb{R})} \right) \end{aligned}$$

for any μ such that $\mathcal{Q}_T(\mu) < +\infty$, any $f \in \mathcal{S}(\mathbb{R})$ and any $0 \le s \le t \le T$. Therefore, for any $f \in \mathcal{S}(\mathbb{R})$, $\{\mu_t(f): 0 \le t \le T\}_{\mu \in \mathbb{L}_c}$ are uniformly bounded and equicontinuous. As a result, the compactness of \mathbb{L}_c follows from Ascoli-Arzela theorem.

3.3 Alternative formulas for Q_T

Introduce

$$\mathcal{A} = \big\{ \mu : \text{there exist } H \in \mathcal{C}_c^{1,2}([0,T] \times \mathbb{R})) \text{ and } \psi \in \mathcal{C}_c^2(\mathbb{R}) \text{ such that } \mu \text{ satisfies (3.1)} \big\}.$$

One could check directly that if $\mu \in A$, then for any t > 0, $\mu(t, \cdot)$ is integrable on \mathbb{R} and vanishes at infinity. For $\mu \in A$, define J the macroscopic instantaneous current as

$$J = -\frac{1}{2}\partial_u \mu + \chi(\rho)\partial_u H.$$
(3.6)

Obviously,

$$\partial_t \mu + \partial_u J = 0.$$

Define K(t, u) the macroscopic current across the macroscopic point u during time [0, t] as

$$K(t,u) := \int_0^t J(s,u) ds.$$

Then, for $u \in \mathbb{R}$ and t > 0,

$$\partial_u K(t, u) = \mu(0, u) - \mu(t, u).$$
 (3.7)

By direct calculations, for $\mu \in \mathcal{A}$,

$$\chi(\rho)\mathcal{Q}_{T}(\mu) = \chi(\rho)\mathcal{Q}_{T}(\mu_{0}, K) := \frac{1}{2} \int_{0}^{T} \int_{\mathbb{R}} \left(\partial_{t}K\right)^{2} (t, u) du dt + \frac{1}{4} \int_{\mathbb{R}} \left(\partial_{u}K\right)^{2} (T, u) du + \frac{1}{8} \int_{0}^{T} \int_{\mathbb{R}} \left(\partial_{u}\mu_{0}(u) - \partial_{u}^{2}K(t, u)\right)^{2} du dt + \frac{1}{2} \int_{\mathbb{R}} \left[\mu_{0}^{2}(u) - \mu_{0}(u)\partial_{u}K(T, u)\right] du.$$
(3.8)

The above formula is exactly [39, Eqn. (3.8)], whose proof is presented in Appendix A.3.

4 Variational formula for the rate function

In this section, we present a variational formula for the rate function \mathcal{I}_{path} . For any $n \geq 1$, any column vector $\boldsymbol{\alpha} \in \mathbb{R}^n$ and any $0 \leq t_1 < t_2 < \ldots < t_n$, let

$$\mathcal{I}_{\{t_j\}_{j=1}^n}(\alpha) = \frac{1}{2} \alpha^T A_{\{t_j\}_{j=1}^n}^{-1} \alpha$$
(4.1)

where $A_{\{t_j\}_{j=1}^n}$ is the $n \times n$ matrix with (k, ℓ) -component equal to $a(t_k, t_\ell)$ for $1 \le k, \ell \le n$. Recall $a(\cdot, \cdot)$ is the covariance function of the fractional Brownian motion with Hurst parameter 1/4.

The following variational formula relates \mathcal{I}_{path} the rate function of the sample path to the corresponding finite-dimensional rate functions, which is a direct consequence of [8, Theorem 4.28].

Proposition 4.1. For any $f \in \mathcal{D}([0,T],\mathbb{R})$,

$$\begin{split} \mathcal{I}_{path}(f) \\ &= \sup \Big\{ \mathcal{I}_{_{\{t_j\}_{j=1}^n}} \left(f_{_{\{t_j\}_{j=1}^n}} \right): \ n \geq 1, 0 \leq t_1 < t_2 < \ldots < t_n \leq T, t_j \in \Delta_f^c \text{ for all } 1 \leq j \leq n \Big\}, \\ \text{where } f_{_{\{t_j\}_{j=1}^n}} = (f(t_1), f(t_2), \ldots, f(t_n))^T \text{ and } \Delta_f^c \text{ is the set of continuities of } f. \end{split}$$

Proof. It is well known that for any $0 \le t_1 < t_2 < \ldots < t_n$, $\mathcal{I}_{\{t_j\}_{j=1}^n}(\cdot)$ is the large deviation rate function of the sequence of Gaussian vectors $\left\{\frac{1}{\sqrt{N}}\left(B_{t_1}^{1/4}, B_{t_2}^{1/4}, \ldots, B_{t_n}^{1/4}\right)\right\}_{N\ge 1}$, and $\mathcal{I}_{path}(\cdot)$ is the rate function of the sequence $\left\{\frac{1}{\sqrt{N}}B_t^{1/4}: 0 \le t \le T\right\}_{N\ge 1}$, see [7] for example. Then, by [8, Theorem 4.28], the result follows immediately. \Box

The following lemma gives a variational formula for the finite-dimensional rate function appearing in the last proposition, which is the main result of this section.

Proposition 4.2. For any positive integer n, any column vector $\boldsymbol{\alpha} = (\alpha_1, \dots, \alpha_n)^T \in \mathbb{R}^n$, and any $0 < t_1 < t_2 < \dots < t_n \leq T$, we have

$$\inf\left\{\mathcal{Q}_T(\mu): \int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j, \ 1 \le j \le n\right\} = \frac{\sqrt{2\pi}}{4\chi(\rho)} \boldsymbol{\alpha}^T A^{-1} \boldsymbol{\alpha}, \tag{4.2}$$

where $A = A_{\{t_i\}_{i=1}^n}$ for short.

Note that

$$\frac{\sqrt{2\pi}}{4\chi(\rho)}\boldsymbol{\alpha}^T A^{-1}\boldsymbol{\alpha} = \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}).$$

In the rest of this section, we prove Proposition 4.2.

4.1 Restricted to $\mu \in A$

We first show that we could restrict the infimum in Proposition 4.2 to $\mu \in A$, which is the content of the next result.

Proposition 4.3 (The case $n \ge 1$). For any positive integer n, any column vector $\alpha \in \mathbb{R}^n$, and any $0 < t_1 < t_2 < \ldots < t_n \le T$, we have

$$\inf\left\{\mathcal{Q}_T(\mu): \mu \in \mathcal{A}, \ \int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j, \ 1 \le j \le n\right\} = \frac{\sqrt{2\pi}}{4\chi(\rho)} \boldsymbol{\alpha}^T A^{-1} \boldsymbol{\alpha}, \ (4.3)$$

where $A = A_{\{t_j\}_{j=1}^n}$.

We first prove Proposition 4.2 from the last result, and then prove the last proposition in the next subsection.

Proof of Proposition 4.2. By Proposition 4.3,

$$\inf\left\{\mathcal{Q}_T(\mu): \int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j, \ 1 \le j \le n\right\} \le \frac{\sqrt{2\pi}}{4\chi(\rho)} \boldsymbol{\alpha}^T A^{-1} \boldsymbol{\alpha},$$

and it remains to prove the opposite direction of the last inequality. Take μ such that $Q_T(\mu) < \infty$ and

$$\int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j, \ \forall 1 \le j \le n.$$

Then, by Lemma 3.3, there exist $\psi \in L^2(\mathbb{R})$ and $H \in \mathcal{H}^1$ such that μ is the unique weak solution to the PDE (3.1). For $\varepsilon > 0$, let $\psi^{\varepsilon} \in \mathcal{C}^2_c(\mathbb{R})$ and $H^{\varepsilon} \in \mathcal{C}^{1,2}_c([0,T] \times \mathbb{R})$ such that

$$\psi^{\varepsilon} \to \psi \quad \text{in } L^2(\mathbb{R}), \quad \text{and} \quad H^{\varepsilon} \to H \quad \text{in } \mathcal{H}^1, \quad \varepsilon \to 0.$$

Denote by μ^{ε} the solution of the PDE (3.1) associated with ψ^{ε} and H^{ε} . Finally, let

$$\alpha_j^{\varepsilon} := \int_0^\infty [\mu^{\varepsilon}(t_j, u) - \mu^{\varepsilon}(0, u)] du, \quad \forall 1 \le j \le n.$$

By Lemmas 3.3 and 3.4,

$$\lim_{\varepsilon \to 0} \mathcal{Q}_T(\mu^{\varepsilon}) = \mathcal{Q}_T(\mu), \quad \lim_{\varepsilon \to 0} \alpha_j^{\varepsilon} = \alpha_j, \quad \forall 1 \le j \le n.$$

Thus,

$$\begin{aligned} \mathcal{Q}_T(\mu) &= \lim_{\varepsilon \to 0} \mathcal{Q}_T(\mu^\varepsilon) \\ &\geq \limsup_{\varepsilon \to 0} \inf \left\{ \mathcal{Q}_T(\mu) : \mu \in \mathcal{A}, \ \int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j^\varepsilon, \ 1 \le j \le n \right\} \\ &= \limsup_{\varepsilon \to 0} \frac{\sqrt{2\pi}}{4\chi(\rho)} \alpha_\varepsilon^T A^{-1} \alpha_\varepsilon = \frac{\sqrt{2\pi}}{4\chi(\rho)} \alpha^T A^{-1} \alpha, \end{aligned}$$

where $\alpha_{\varepsilon} = (\alpha_1^{\varepsilon}, \alpha_2^{\varepsilon}, \dots, \alpha_n^{\varepsilon})^T$. We conclude the proof by taking the infimum over μ on the left-hand side of the last inequality.

4.2 **Proof of Proposition 4.3**

The proof of Proposition 4.3 is based on Fourier approach and is divided into several lemmas. Observe that if $\mu \in A$, then by (3.7) and the fact that μ vanishes at infinity,

$$K(t,0) = \int_0^\infty [\mu(t,u) - \mu(0,u)] du, \quad t > 0,$$

which allows us to reduce the constraints inside the infimum in (4.3) to $K(t_j, 0) = \alpha_j$ for $1 \le j \le n$. Obviously, we also have K(0, u) = 0 for any $u \in \mathbb{R}$ by the definition of K(t, u).

First, we consider the case n = 1 and $K(T, 0) = \alpha$ for some $\alpha \in \mathbb{R}$. The following result was proved in [39] by constructing the minimizer of the variational problem directly. Here, we present a different approach, which allows us to generalize it to the case n > 1.

Lemma 4.4. Fix $\alpha \in \mathbb{R}$. Then,

$$\inf\{\mathcal{Q}_T(\mu): \mu \in \mathcal{A}, \ K(0, \cdot) = 0, \ K(T, 0) = \alpha\} = \frac{\sqrt{2\pi}\alpha^2}{4\chi(\rho)\sqrt{T}}.$$
(4.4)

Proof. Recall that we sometimes write $Q_T(\mu)$ as $Q_T(\mu_0, K)$ to indicate the dependence of Q_T on μ_0 and K, see (3.8). The proof is divided into several steps:

Step 1. We first minimize $\mathcal{Q}_T(\mu_0, K)$ subject to $K(0, \cdot) = 0$ and $K(T, \cdot)$ being a known function. We claim that the infimum is attained at the point $(K^{T,\alpha}, \mu_0^{T,\alpha})$ such that

$$-\partial_t^2 K^{T,\alpha}(t,u) + \frac{1}{4} \partial_u^4 K^{T,\alpha}(t,u) - \frac{1}{4} \partial_u^3 \mu_0^{T,\alpha}(u) = 0,$$
(4.5)

$$-\frac{T}{4}\partial_u^2\mu_0^{T,\alpha}(u) + \frac{1}{4}\int_0^T \partial_u^3 K^{T,\alpha}(t,u)dt + \mu_0^{T,\alpha}(u) - \frac{1}{2}\partial_u K^{T,\alpha}(T,u) = 0.$$
(4.6)

Indeed, the above equations are obtained by solving, for any μ_0 and K such that $K(0,\cdot) = K(T,\cdot) = 0$,

$$\frac{d}{d\varepsilon}\mathcal{Q}_T(\mu_0^{T,\alpha} + \varepsilon\mu_0, K^{T,\alpha})|_{\varepsilon=0} = \frac{d}{d\varepsilon}\mathcal{Q}_T(\mu_0^{T,\alpha}, K^{T,\alpha} + \varepsilon K)|_{\varepsilon=0} = 0.$$

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For any function $f \in \mathcal{S}(\mathbb{R})$, define the Fourier transform of f as

$$\mathcal{F}f(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{iu\xi} f(u) du, \quad \xi \in \mathbb{R}.$$

Then, $(\mathcal{F}f')(\xi) = -i\xi\mathcal{F}f(\xi), \ \xi \in \mathbb{R}$. This permits us to rewrite (4.5) and (4.6) as

$$-\partial_t^2 \mathcal{F} K^{T,\alpha}(t,\xi) + \frac{\xi^4}{4} \mathcal{F} K^{T,\alpha}(t,\xi) - \frac{i\xi^3}{4} \mathcal{F} \mu_0^{T,\alpha}(\xi) = 0,$$
(4.7)

$$\left\{1 + \frac{T\xi^2}{4}\right\} \mathcal{F}\mu_0^{T,\alpha}(\xi) = -i\left\{\frac{\xi}{2}\mathcal{F}K^{T,\alpha}(T,\xi) + \frac{\xi^3}{4}\int_0^T \mathcal{F}K^{T,\alpha}(t,\xi)dt\right\}.$$
(4.8)

Solving (4.7), we have

$$\mathcal{F}K^{T,\alpha}(t,\xi) = C_1(\xi)e^{t\xi^2/2} + C_2(\xi)e^{-t\xi^2/2} + \frac{i\mathcal{F}\mu_0^{T,\alpha}(\xi)}{\xi}$$
(4.9)

for two functions $C_1(\xi)$ and $C_2(\xi)$. Since we assumed $K^{T,\alpha}(0,\cdot) = 0$ and $K^{T,\alpha}(T,\cdot)$ is known, the coefficients are determined by

$$\begin{cases} C_1(\xi) + C_2(\xi) + \frac{i\mathcal{F}\mu_0^{T,\alpha}(\xi)}{\xi} = 0, \\ C_1(\xi)e^{T\xi^2/2} + C_2(\xi)e^{-T\xi^2/2} + \frac{i\mathcal{F}\mu_0^{T,\alpha}(\xi)}{\xi} = \mathcal{F}K^{T,\alpha}(T,\xi). \end{cases}$$
(4.10)

Inserting (4.9) into (4.8), we have

$$\mathcal{F}\mu_0^{T,\alpha}(\xi) = -\frac{i\xi}{2} \Big\{ \mathcal{F}K^{T,\alpha}(T,\xi) + C_1(\xi) \big[e^{T\xi^2/2} - 1 \big] + C_2(\xi) \big[1 - e^{-T\xi^2/2} \big] \Big\}.$$
(4.11)

Finally, using (4.11), and solving (4.10), we have

$$\begin{cases} C_1(\xi) = \frac{1}{2(e^{T\xi^2/2} - 1)} \mathcal{F} K^{T,\alpha}(T,\xi), \\ C_2(\xi) = -\frac{e^{T\xi^2/2}}{2(e^{T\xi^2/2} - 1)} \mathcal{F} K^{T,\alpha}(T,\xi). \end{cases}$$

In conclusion, we have shown that, for $0 \le t \le T$,

$$\mathcal{F}\mu_0^{T,\alpha}(\xi) = -\frac{i\xi}{2}\mathcal{F}K^{T,\alpha}(T,\xi),\tag{4.12}$$

$$\mathcal{F}K^{T,\alpha}(t,\xi) = \frac{1}{2} \Big[\frac{e^{t\xi^2/2} - e^{(T-t)\xi^2/2}}{e^{T\xi^2/2} - 1} + 1 \Big] \mathcal{F}K^{T,\alpha}(T,\xi).$$
(4.13)

Step 2. By Parseval's identity and (3.8),

$$\chi(\rho)\mathcal{Q}_{T}(\mu_{0},K) = \int_{0}^{T} \int_{\mathbb{R}} \left\{ \frac{1}{2} \left| \partial_{t}\mathcal{F}K(t,\xi) \right|^{2} + \frac{1}{8} \left| -i\xi\mathcal{F}\mu_{0}(\xi) + \xi^{2}\mathcal{F}K(t,\xi) \right|^{2} \right\} d\xi \, dt \\ + \frac{1}{4} \int_{\mathbb{R}} \left\{ \left| \mathcal{F}\mu_{0}(\xi) \right|^{2} + \left| \mathcal{F}\mu_{0}(\xi) + i\xi\mathcal{F}K(T,\xi) \right|^{2} \right\} d\xi.$$
(4.14)

For example, to obtain the last line, just note that

$$\begin{split} &\frac{1}{4} \int_{\mathbb{R}} \left(\partial_u K\right)^2 (T, u) du + \frac{1}{2} \int_{\mathbb{R}} \left[\mu_0^2(u) - \mu_0(u) \partial_u K(T, u)\right] du \\ &= \frac{1}{4} \int_{\mathbb{R}} \left|\xi \mathcal{F}K(T, \xi)\right|^2 d\xi + \frac{1}{2} \int_{\mathbb{R}} \left[|\mathcal{F}\mu_0(\xi)|^2 + i\xi \overline{\mathcal{F}\mu_0(\xi)} \mathcal{F}K(T, \xi)\right] d\xi \\ &= \frac{1}{4} \int_{\mathbb{R}} \left\{\left|\mathcal{F}\mu_0(\xi)\right|^2 + \left|\mathcal{F}\mu_0(\xi) + i\xi \mathcal{F}K(T, \xi)\right|^2\right\} d\xi. \end{split}$$

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The first term on the right hand side could be calculated similarly. Inserting $(K^{T,\alpha}, \mu_0^{T,\alpha})$ into (4.14), and by (4.12) and (4.13),

$$\chi(\rho)\mathcal{Q}_{T}(\mu_{0}^{T,\alpha},K^{T,\alpha}) = \int_{0}^{T}\int_{\mathbb{R}}\frac{\xi^{4}}{32[e^{T\xi^{2}/2}-1]^{2}}|\mathcal{F}K^{T,\alpha}(T,\xi)|^{2}\left[\left(e^{t\xi^{2}/2}+e^{(T-t)\xi^{2}/2}\right)^{2}+\left(e^{t\xi^{2}/2}-e^{(T-t)\xi^{2}/2}\right)^{2}\right]d\xi\,dt + \int_{\mathbb{R}}\frac{\xi^{2}}{8}|\mathcal{F}K^{T,\alpha}(T,\xi)|^{2}\,d\xi = \int_{\mathbb{R}}\frac{\xi^{2}e^{T\xi^{2}/2}}{4[e^{T\xi^{2}/2}-1]}|\mathcal{F}K^{T,\alpha}(T,\xi)|^{2}\,d\xi.$$
(4.15)

Step 3. Finally, we need to optimize $\mathcal{Q}_T(\mu_0^{T,\alpha}, K^{T,\alpha})$ over $\mathcal{F}K^{T,\alpha}(T,\xi)$. The boundary condition $K^{T,\alpha}(T,0) = \alpha$ is equivalent to

$$\frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \mathcal{F} K^{T,\alpha}(T,\xi) d\xi = \alpha.$$
(4.16)

By Cauchy-Schwarz inequality,

$$\Big(\int_{\mathbb{R}} \mathcal{F}K^{T,\alpha}(T,\xi)d\xi\Big)^2 \le \Big(\int_{\mathbb{R}} \frac{\xi^2 e^{T\xi^2/2}}{4[e^{T\xi^2/2}-1]} |\mathcal{F}K^{T,\alpha}(T,\xi)|^2 d\xi\Big) \Big(\int_{\mathbb{R}} \frac{4[e^{T\xi^2/2}-1]}{\xi^2 e^{T\xi^2/2}} d\xi\Big).$$

Thus, the infimum of (4.15) is attained at the point

$$\mathcal{F}K^{T,\alpha}(T,\xi) = \frac{\alpha(1-e^{-T\xi^2/2})}{\sqrt{T\xi^2}}$$

and equals

$$(\sqrt{2\pi\alpha})^2 \Big(\int_{\mathbb{R}} \frac{4(e^{T\xi^2/2} - 1)}{\xi^2 e^{T\xi^2/2}} \, d\xi \Big)^{-1} = (\sqrt{2\pi\alpha})^2 \Big(\int_{\mathbb{R}} \frac{4(1 - e^{-T\xi^2/2})}{\xi^2} \, d\xi \Big)^{-1} = \frac{\sqrt{2\pi\alpha^2}}{4\sqrt{T}}.$$

In the last identity, we used Equation (A.1). This concludes the proof.

The next lemma extends the last result to the case n = 1 and $K(t, 0) = \alpha$ for $0 < t \le T$ and $\alpha \in \mathbb{R}$.

Lemma 4.5. Fix $0 < t \leq T$ and $\alpha \in \mathbb{R}$. Then,

$$\inf\{\mathcal{Q}_T(\mu): \mu \in \mathcal{A}, \ K(0, \cdot) = 0, \ K(t, 0) = \alpha\} = \frac{\sqrt{2\pi}\alpha^2}{4\chi(\rho)\sqrt{t}}.$$
(4.17)

Proof. For any t > 0, let $\mu_T^{t,\alpha}$ be the minimizer of (4.17), which does exist by the following construction. Denote by $H_T^{t,\alpha}$ the function H in (3.1), respectively $K_T^{t,\alpha}$ the function K in (3.7), associated with the function $\mu_T^{t,\alpha}$. We call $\mu_T^{t,\alpha}$, or correspondingly $(K_T^{t,\alpha}, \mu_{T,0}^{t,\alpha})$, the minimizer of the rate function $\mathcal{Q}_T(\mu)$ with current α at time t. In particular, when T = t,

$$\mu^{T,\alpha}=\mu^{T,\alpha}_T,\quad H^{T,\alpha}=H^{T,\alpha}_T,\quad K^{T,\alpha}=K^{T,\alpha}_T,$$

the precise expressions of which are given through the proof of Lemma 4.4.

We claim the infimum in (4.17) is attained at the point $\mu_T^{t, \alpha}$ such that

(i) for $0 \leq s \leq t,$ the function $\mu_T^{t,\alpha}$ is the unique solution to

$$\begin{cases} \partial_s \mu_T^{t,\alpha}(s,u) = \frac{1}{2} \partial_u^2 \mu_T^{t,\alpha}(s,u) - \chi(\rho) \partial_u^2 H^{t,\alpha}(s,u), & u \in \mathbb{R}, \\ \mu_T^{t,\alpha}(0,u) = \mu_0^{t,\alpha}(u), & u \in \mathbb{R}, \end{cases}$$
(4.18)

where $H^{t,\alpha}$ is the function H associated with the minimizer of the rate function $Q_t(\mu)$ with current α at time t;

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(ii) for $t \leq s \leq T$, the function $\mu_T^{t,\alpha}$ is the unique solution to

$$\begin{cases} \partial_s \mu_T^{t,\alpha}(s,u) = \frac{1}{2} \partial_u^2 \mu_T^{t,\alpha}(s,u), & u \in \mathbb{R}, \\ \mu_T^{t,\alpha}(t,u) = \mu^{t,\alpha}(t,u), & u \in \mathbb{R}. \end{cases}$$
(4.19)

Roughly speaking, the minimizer $\mu_T^{t,\alpha}$ evolves as the minimizer of the rate function $\mathcal{Q}_t(\mu)$ with current α at time t during time interval [0,t], and then evolves according to the hydrodynamic equation of the SSEP after time t. This claim is obvious since

$$\mathcal{Q}_{T,dyn}(\mu) = rac{\chi(\rho)}{2}[H,H],$$

and thus

$$\inf\{\mathcal{Q}_{T}(\mu): \mu \in \mathcal{A}, \ K(0, \cdot)=0, \ K(t, 0)=\alpha\} \ge \inf\{\mathcal{Q}_{t}(\mu): \mu \in \mathcal{A}, \ K(0, \cdot)=0, \ K(t, 0)=\alpha\}$$

Moreover, by Lemma 4.4, the left-hand side of (4.17) equals

$$\mathcal{Q}_T(\mu_T^{t,\alpha}) = \mathcal{Q}_t(\mu^{t,\alpha}) + \int_t^T 0 ds = \frac{\sqrt{2\pi\alpha^2}}{4\chi(\rho)\sqrt{t}}$$

which concludes the proof.

One could check directly that the functions $\mu_T^{t,\alpha}$ and $(K_T^{t,\alpha}, \mu_{T,0}^{t,\alpha})$ constructed in the last proof satisfy the following properties:

(P1) for $0 \le s \le t$,

$$\mathcal{F}\mu_{T,0}^{t,\alpha}(\xi) = -\frac{i\xi}{2}\mathcal{F}K^{t,\alpha}(t,\xi), \qquad (4.20)$$

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{1}{2} \Big[\frac{e^{s\xi^2/2} - e^{(t-s)\xi^2/2}}{e^{t\xi^2/2} - 1} + 1 \Big] \mathcal{F}K^{t,\alpha}(t,\xi),$$
(4.21)

where

$$\mathcal{F}K^{t,\alpha}(t,\xi) = \frac{\alpha(1-e^{-t\xi^2/2})}{\sqrt{t\xi^2}};$$

(P2) for $t \leq s \leq T$,

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{1}{2} \left[1 + e^{-(s-t)\xi^2/2} \right] \mathcal{F}K^{t,\alpha}(t,\xi).$$
(4.22)

Moreover, by direct calculations (see Section A.4), for $0 \leq s \leq T$,

$$K_T^{t,\alpha}(s,0) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \mathcal{F} K_T^{t,\alpha}(s,\xi) d\xi = \frac{\alpha a(t,s)}{\sqrt{t}},$$
(4.23)

where $a(t,s) = \frac{1}{2}(\sqrt{t} + \sqrt{s} - \sqrt{|t-s|})$ is the covariance function of the fractional Brownian motion with Hurst parameter $\frac{1}{4}$.

(P1) is exactly (4.12) and (4.13). (P2) follows from (4.19) and (3.7). Indeed, by (4.19),

$$\mathcal{F}\mu_T^{t,\alpha}(s,\xi) = \mathcal{F}\mu^{t,\alpha}(t,\xi)e^{-(s-t)\xi^2/2}, \quad t \le s \le T,$$

and by (3.7),

$$-i\xi \mathcal{F}K_T^{t,\alpha}(s,\xi) = \mathcal{F}\mu_{T,0}^{t,\alpha}(\xi) - \mathcal{F}\mu_T^{t,\alpha}(s,\xi), \quad s \ge 0.$$

$$(4.24)$$

Taking s = t in the last expression, and then using (P1), we obtain

$$\mathcal{F}\mu_T^{t,\alpha}(t,\xi) = \frac{i\xi}{2}\mathcal{F}K^{t,\alpha}(t,\xi).$$

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Therefore, by (4.24), for $s \ge t$,

$$-i\xi\mathcal{F}K_T^{t,\alpha}(s,\xi) = -\frac{i\xi}{2}\mathcal{F}K^{t,\alpha}(t,\xi) - \mathcal{F}\mu^{t,\alpha}(t,\xi)e^{-(s-t)\xi^2/2} = -\frac{i\xi}{2} \left[1 + e^{-(s-t)\xi^2/2}\right]\mathcal{F}K^{t,\alpha}(t,\xi),$$

as claimed.

Finally, we are ready to prove Proposition 4.3.

Proof of Proposition 4.3. Let

$$\mu_T^n := \mu_T^{n, \{t_j\}_{j=1}^n, \{\alpha_j\}_{j=1}^n}, \quad (K_T^n, \mu_{T,0}^n) := (K_T^{n, \{t_j\}_{j=1}^n, \{\alpha_j\}_{j=1}^n}, \mu_{T,0}^{n, \{t_j\}_{j=1}^n, \{\alpha_j\}_{j=1}^n})$$

be the minimizer of (4.2). We claim that

$$K_T^n = \sum_{j=1}^n \beta_j K_T^{t_j,1}, \quad \mu_{T,0}^n = \sum_{j=1}^n \beta_j \mu_{T,0}^{t_j,1}$$
(4.25)

for some coefficients $\beta_j \in \mathbb{R}$, $1 \leq j \leq n$. Recall that for $1 \leq j \leq n$, $\left(K_T^{t_j,1}, \mu_{T,0}^{t_j,1}\right)$ is the minimizer of the rate function $\mathcal{Q}_T(\mu)$ with current one at time t_j , and is constructed in (4.18) and (4.19). We prove claim (4.25) in Section A.5. The coefficients $\{\beta_j\}$ satisfy the following constraints

$$\sum_{j=1}^{n} \beta_j K_T^{t_j,1}(t_k,0) = \alpha_k, \quad 1 \le k \le n.$$
(4.26)

By (4.23), the last line is equivalent to

$$\sum_{j=1}^{n} \beta_j \frac{a(t_j, t_k)}{\sqrt{t_j}} = \alpha_k, \quad 1 \le k \le n.$$

Let $D = \text{diag}(d_j)_{1 \le j \le n}$ be the diagonal matrix with $d_j = 1/\sqrt{t_j}$. Let $\beta = (\beta_1, \ldots, \beta_n)^T$ be the column vector. Then, we have shown that

$$AD\beta = \alpha.$$

Note that AD is invertible, which ensures the existence and uniqueness of β . Inserting (4.25) into (4.14), we have

$$\chi(\rho)\mathcal{Q}_T(\mu_{T,0}^n, K_T^n) = \boldsymbol{\beta}^T Q \boldsymbol{\beta},\tag{4.27}$$

where $Q = (q_{jk})_{1 \le j,k \le n}$ is the matrix with q_{jk} being the real part of

$$\begin{split} &\int_{0}^{T} \int_{\mathbb{R}} \Big\{ \frac{1}{2} \partial_{t} \mathcal{F} K_{T}^{t_{j},1}(t,\xi) \overline{\partial_{t} \mathcal{F} K_{T}^{t_{k},1}(t,\xi)} \\ &\quad + \frac{1}{8} \Big[-i\xi \mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) + \xi^{2} \mathcal{F} K_{T}^{t_{j},1}(t,\xi) \Big] \overline{\big[-i\xi \mathcal{F} \mu_{T,0}^{t_{k},1}(\xi) + \xi^{2} \mathcal{F} K_{T}^{t_{k},1}(t,\xi) \big]} \Big\} d\xi \, dt \\ &\quad + \frac{1}{4} \int_{\mathbb{R}} \Big\{ \mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) \overline{\mathcal{F} \mu_{T,0}^{t_{k},1}(\xi)} \\ &\quad + \big[\mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) + i\xi \mathcal{F} K_{T}^{t_{j},1}(T,\xi) \big] \overline{\big[\mathcal{F} \mu_{T,0}^{t_{k},1}(\xi) + i\xi \mathcal{F} K_{T}^{t_{k},1}(T,\xi) \big]} \Big\} d\xi. \end{split}$$

After a long computation (see Section A.6), we have

$$q_{jk} = \frac{\sqrt{2\pi}}{4\sqrt{t_j t_k}} a_{jk}.$$
(4.28)

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Equivalently,

$$Q = \frac{\sqrt{2\pi}}{4} DAD.$$

Therefore,

$$\chi(\rho)\mathcal{Q}_T(\mu_{T,0}^n, K_T^n) = \frac{\sqrt{2\pi}}{4} \boldsymbol{\alpha}^T A^{-1} \boldsymbol{\alpha},$$

which concludes the proof.

5 Exponential tightness

In this section, we prove the sequences of the two processes

$$\{J_{-1,0}(tN^2)/a_N: 0 \le t \le T\}_{N\ge 1}$$
 and $\{X(tN^2)/a_N: 0 \le t \le T\}_{N\ge 1}$

are exponentially tight.

5.1 A uniform super-exponential estimate

We first state a super-exponential estimate for the space average of the current. The following result was proved by the same authors in [39] without the supremum over time inside the probability.

Lemma 5.1. For any $\delta > 0$,

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}\Big(\sup_{0 \le t \le T} \Big| \frac{1}{nNa_N} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2) \Big| > \delta\Big) = -\infty$$
(5.1)

under both $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$ and $\mathbb{P} = \mathbb{P}_{\nu_{\rho}^*}$.

Proof. By Remark 3.2, we only need to discuss the case where $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$. By Markov's inequality, for any K > 0, the expression in (5.1) is bounded by

$$-\delta K + \frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \sup_{0 \le t \le T} \Big| \frac{a_N K}{nN^2} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2) \Big| \Big\} \Big].$$
(5.2)

Since *K* could be taken arbitrarily large, we only need to prove that, for any fixed K > 0,

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \sup_{0 \le t \le T} \Big| \frac{a_N K}{n N^2} \sum_{x=0}^{n N-1} J_{x,x+1}(t N^2) \Big| \Big\} \Big] = 0.$$

By using the inequality

$$\exp\{\sup_{0 \le t \le T} |x_t|\} \le \exp\{\sup_{0 \le t \le T} x_t\} + \exp\{\sup_{0 \le t \le T} (-x_t)\}$$

for any trajectory $\{x_t\}$, and $\log(a+b) \leq \log 2 + \max\{\log a, \log b\}$ for any a, b > 0, without loss of generality, it suffices to prove that, for any fixed K > 0,

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_\rho} \Big[\exp \Big\{ \sup_{0 \le t \le T} \frac{a_N K}{nN^2} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2) \Big\} \Big] = 0.$$
(5.3)

Since for each $x \in \mathbb{Z}$, $\{J_{x,x+1}(t)\}$ is a compound Poisson process with intensity

$$\frac{1}{2}\int_0^t (\eta_s(x) - \eta_s(x+1))ds,$$

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and there are no common jumps between those compound Poisson processes,

$$M_t^{N,1} = \exp\left\{\frac{a_N K}{nN^2} \sum_{x=0}^{N-1} J_{x,x+1}(tN^2) - \Gamma_{n,K}^{N,1}(t)\right\}$$

is a mean-one exponential martingale, where

$$\Gamma_{n,K}^{N,1}(t) = \frac{1}{2} \sum_{x=0}^{nN-1} \left(e^{a_N K/(nN^2)} - 1 \right) \int_0^{tN^2} (\eta_s(x) - \eta_s(x+1)) ds.$$

Then, by Cauchy-Schwarz inequality, the expectation in (5.3) is bounded by

$$\mathbb{E}_{\nu_{\rho}} \left[\sup_{0 \le t \le T} \{ M_{t}^{N,1} \} \exp \left\{ \sup_{0 \le t \le T} \Gamma_{n,K}^{N,1}(t) \right\} \right] \\ \le \mathbb{E}_{\nu_{\rho}} \left[\sup_{0 \le t \le T} (M_{t}^{N,1})^{2} \right]^{1/2} \mathbb{E}_{\nu_{\rho}} \left[\exp \left\{ \sup_{0 \le t \le T} 2\Gamma_{n,K}^{N,1}(t) \right\} \right]^{1/2} \right]^{1/2}$$

As a consequence, the expression in (5.3) is bounded by

$$\frac{N}{2a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\sup_{0 \le t \le T} (M_t^{N,1})^2 \Big] + \frac{N}{2a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \sup_{0 \le t \le T} 2\Gamma_{n,K}^{N,1}(t) \Big\} \Big].$$

For the martingale term in the last line, we have

$$\begin{split} (M_t^{N,1})^2 &= \exp\Big\{\frac{1}{2}\Big[\frac{4a_NK}{nN^2}\sum_{x=0}^{nN-1}J_{x,x+1}(tN^2) - 2\Gamma_{n,K}^{N,2}(t)\Big] + \Gamma_{n,K}^{N,2}(t) - 2\Gamma_{n,K}^{N,1}(t)\Big\} \\ &=: (M_t^{N,2})^{1/2}\exp\Big\{\Gamma_{n,K}^{N,2}(t) - 2\Gamma_{n,K}^{N,1}(t)\Big\}, \end{split}$$

where

$$2\Gamma_{n,K}^{N,2}(t) = \frac{1}{2} \sum_{x=0}^{nN-1} \left(e^{4a_N K/(nN^2)} - 1 \right) \int_0^{tN^2} (\eta_s(x) - \eta_s(x+1)) ds,$$

and

$$M_t^{N,2} = \exp\left\{\frac{4a_NK}{nN^2}\sum_{x=0}^{nN-1}J_{x,x+1}(tN^2) - 2\Gamma_{n,K}^{N,2}(t)\right\}$$

is also a martingale. By Doob's inequality and Cauchy-Schwarz inequality,

$$\frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \left[\sup_{0 \le t \le T} (M_t^{N,1})^2 \right] \le \frac{N \log 4}{a_N^2} + \frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \left[(M_T^{N,1})^2 \right] \\
\le \frac{N \log 4}{a_N^2} + \frac{N}{2a_N^2} \log \mathbb{E}_{\nu_{\rho}} \left[M_T^{N,2} \right] + \frac{N}{2a_N^2} \log \mathbb{E}_{\nu_{\rho}} \left[\exp \left\{ 2\Gamma_{n,K}^{N,2}(T) - 4\Gamma_{n,K}^{N,1}(T) \right\} \right].$$
(5.4)

Note that the second term on the right-hand side of (5.4) is zero. Since $a_N \gg \sqrt{N}$, the first term on the right-hand side of (5.4) converges to zero as $N \to \infty$. Therefore, to prove (5.3), we only need to prove

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{2a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \sup_{0 \le t \le T} 2\Gamma_{n,K}^{N,1}(t) \Big\} \Big] = 0,$$
(5.5)

since the third term on the right-hand side of (5.4) could be treated in the same way. Since $|e^x - 1 - x| \le (x^2/2)e^{|x|}$,

$$\Gamma_{n,K}^{N,1}(t) \le \frac{a_N K}{2nN^2} \int_0^{tN^2} (\eta_s(0) - \eta_s(nN)) ds + \frac{C_{N,n} t a_N^2 K^2}{n^2 N^2},$$

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where $C_{N,n}$ is uniformly bounded in N and n. Thus, we are left to show

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_\rho} \Big[\exp \Big\{ \sup_{0 \le t \le T} \frac{a_N K}{nN^2} \int_0^{tN^2} (\eta_s(0) - \eta_s(nN)) ds \Big\} \Big] = 0.$$
(5.6)

Since

$$\sup_{0 \le t \le T} \frac{a_N K}{nN^2} \int_0^{tN^2} (\eta_s(0) - \eta_s(nN)) ds \le \sup_{0 \le k \le TN} \frac{a_N K}{nN^2} \int_0^{kN} (\eta_s(0) - \eta_s(nN)) ds + \frac{a_N K}{nN},$$

where the supremum on the right-hand side is over integers k, the left-hand side in (5.6) is bounded by

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \sup_{0 \le k \le TN} \frac{a_N K}{nN^2} \int_0^{kN} (\eta_s(0) - \eta_s(nN)) ds \Big\} \Big].$$

Using the inequality

$$\log \mathbb{E}[\exp\{\sup_{0 \le k \le TN} X_k\}] \le \log \mathbb{E}\Big[\sum_{k=0}^{TN} e^{X_k}\Big] \le \log(TN+1) + \max_{0 \le k \le TN} \log \mathbb{E}[e^{X_k}],$$

and since $a_N \gg \sqrt{N \log N}$, we only need to prove

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \max_{0 \le k \le TN} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_\rho} \Big[\exp \Big\{ \frac{a_N K}{nN^2} \int_0^{kN} (\eta_s(0) - \eta_s(nN)) ds \Big\} \Big] = 0, \quad (5.7)$$

which has been shown in the proof of [39, Lemma 4.1]. We also refer the readers to Section A.7 for details. This concludes the proof. $\hfill \Box$

5.2 Exponential tightness of the current

To show the sequence of processes $\{J_{-1,0}(tN^2)/a_N : 0 \le t \le T\}_{N\ge 1}$ is exponentially tight, we only need to prove the following estimates for the current.

Lemma 5.2. We have the following estimates for the current under both $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$ and $\mathbb{P} = \mathbb{P}_{\nu_{\rho}^*}$:

1. for any fixed T > 0,

$$\limsup_{M \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}\left(\sup_{0 \le t \le T} \frac{1}{a_N} \left| J_{-1,0}(tN^2) \right| > M\right) = -\infty;$$
(5.8)

2. for any $\varepsilon > 0$,

$$\limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \mathbb{P}\Big(\sup_{0 < t \le \delta} \frac{1}{a_N} \big| J_{-1,0}((t+\tau)N^2) - J_{-1,0}(\tau N^2) \big| > \varepsilon\Big) = -\infty,$$
(5.9)

where \mathcal{T}_T is the set of all stopping times bounded by T.

Proof. As in the last lemma, we only need to deal with the case where $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$. For any n > 0, introduce the function $G_n : \mathbb{R} \to \mathbb{R}$ as

$$G_n(u) = (1 - u/n)^+ \mathbf{1} \{ u \ge 0 \}, \quad u \in \mathbb{R}.$$

Then, one could check directly that, for any n > 0,

$$\frac{1}{a_N}J_{-1,0}(tN^2) = \left\langle \mu_t^N, G_n \right\rangle - \left\langle \mu_0^N, G_n \right\rangle + \frac{1}{nNa_N} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2).$$
(5.10)

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Thus, we only need to show that the estimates in the lemma hold respectively for the terms on the right-hand side of the last line, *i.e.*,

$$\limsup_{n \to \infty} \limsup_{M \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\sup_{0 \le t \le T} \left| \frac{1}{n N a_N} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2) \right| > M \right) = -\infty,$$
(5.11)

 $\limsup_{n \to \infty} \limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T}$

$$\frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le \delta} \Big| \frac{1}{nNa_N} \sum_{x=0}^{nN-1} [J_{x,x+1}((t+\tau)N^2) - J_{x,x+1}(\tau N^2)] \Big| > \varepsilon \Big) = -\infty, \quad (5.12)$$

and

$$\limsup_{n \to \infty} \limsup_{M \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le T} \left| \left\langle \mu_t^N, G_n \right\rangle \right| > M \Big) = -\infty, \quad (5.13)$$

$$\limsup_{n \to \infty} \limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le \delta} \left| \left\langle \mu_{t+\tau}^N - \mu_{\tau}^N, G_n \right\rangle \right| > \varepsilon \Big) = -\infty.$$
(5.14)

Note that the left hand side of (5.11) is bounded by

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le T} \Big| \frac{1}{n N a_N} \sum_{x=0}^{nN-1} J_{x,x+1}(tN^2) \Big| > 1 \Big).$$

Thus, (5.11) is true by Lemma 5.1. Since the probability in (5.12) is bounded by

$$2\mathbb{P}_{\nu_{\rho}}\Big(\sup_{0\leq t\leq T}\Big|\frac{1}{nNa_{N}}\sum_{x=0}^{nN-1}J_{x,x+1}(tN^{2})\Big|>\varepsilon/2\Big),$$

the estimate in (5.12) also follows directly from Lemma 5.1.

The estimates in (5.13) and (5.14) have been proved in [10, Lemma 3.2] for Schwartz test functions. Thus, we only need to find $\tilde{G}_n \in \mathcal{S}(\mathbb{R})$ such that, for any $\varepsilon > 0$,

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le T} \left| \left\langle \mu_t^N, G_n - \widetilde{G}_n \right\rangle \right| > \varepsilon \Big) = -\infty.$$
(5.15)

We underline that this is not obvious since we only have the priori bound

$$\left|\left\langle \mu_t^N, G_n - \widetilde{G}_n\right\rangle\right| \le C(n) \frac{N}{a_N}$$

and note that $a_N \ll N$. To this end, we choose $\tilde{G}_n \in \mathcal{S}(\mathbb{R})$ such that

$$|\delta_n(u)| := |G_n(u) - \tilde{G}_n(u)| \le Cn^{-1}, \quad \forall u \in \mathbb{R},$$

and that the support of $\delta_n(\cdot)$ is contained in [-2n, 2n]. For $0 \le i \le N^3$, let $t_i = iT/N^3$. Then, the probability in (5.15) is bounded by

$$\mathbb{P}_{\nu_{\rho}}\left(\sup_{0\leq i\leq N^{3}}\left|\left\langle\mu_{t_{i}}^{N},\delta_{n}\right\rangle\right|>\varepsilon/2\right)+\mathbb{P}_{\nu_{\rho}}\left(\sup_{0\leq t\leq T}\left|\left\langle\mu_{t}^{N},\delta_{n}\right\rangle\right|-\sup_{0\leq i\leq N^{3}}\left|\left\langle\mu_{t_{i}}^{N},\delta_{n}\right\rangle\right|>\varepsilon/2\right).$$
 (5.16)

For the first term in the last line, since the process is stationary,

$$\mathbb{P}_{\nu_{\rho}}\left(\sup_{0\leq i\leq N^{3}}\left|\left\langle \mu_{t_{i}}^{N},\delta_{n}\right\rangle\right|>\varepsilon/2\right)\leq (N^{3}+1)\mathbb{P}_{\nu_{\rho}}\left(\left|\left\langle \mu_{0}^{N},\delta_{n}\right\rangle\right|>\varepsilon/2\right).$$

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Then, using $a_N \gg \sqrt{N \log N}$, for any A > 0,

$$\begin{split} &\lim_{n \to \infty} \sup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le i \le N^3} \left| \left\langle \mu_{t_i}^N, \delta_n \right\rangle \right| > \varepsilon/2 \Big) \\ &= \limsup_{n \to \infty} \sup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\left| \left\langle \mu_0^N, \delta_n \right\rangle \right| > \varepsilon/2 \Big) \\ &\le -\frac{A\varepsilon}{2} + \limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \frac{a_N A}{N} \Big| \sum_{x \in \mathbb{Z}} \bar{\eta}_0(x) \delta_n(x/N) \Big| \Big\} \Big]. \end{split}$$

As before, we could remove the absolute value inside the exponential in the last line. Since ν_{ρ} is a product measure, and by Taylor's expansion up to second order, we bound the last line by

$$\begin{aligned} -\frac{A\varepsilon}{2} + \limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \sum_{x \in \mathbb{Z}} \log \mathbb{E}_{\nu_{\rho}} \Big[\exp \Big\{ \frac{a_N A}{N} \bar{\eta}_0(x) \delta_n(x/N) \Big\} \Big] \\ \leq -\frac{A\varepsilon}{2} + \limsup_{n \to \infty} \limsup_{N \to \infty} \frac{CA^2}{N} \sum_{x \in \mathbb{Z}} \delta_n^2(x/N) = -\frac{A\varepsilon}{2}. \end{aligned}$$

Since A could be arbitrarily large,

$$\limsup_{n \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le i \le N^3} \left| \left\langle \mu_{t_i}^N, \delta_n \right\rangle \right| > \varepsilon/2 \Big) = -\infty.$$

It remains to bound the second term in (5.16). Before that, we first recall the stirring representation for the symmetric exclusion process. Initially, we put a particle at each site of \mathbb{Z} . Those particles are distinguishable. For each bond (x, x + 1), we have a Poisson process with parameter $\frac{1}{2}$, and assume those Poisson processes are independent. When the Poisson clock associated with the bond (x, x + 1) rings, we exchange the particles at sites x and x + 1. Let ξ_t^x be the position at time t of the particle initially at site x. For $\eta \in \Omega$, let

$$\eta_t(x) = 1$$
 if and only if $\xi_t^y = x$ for some $y \in \mathbb{Z}$ such that $\eta(y) = 1$.

Then, $\{\eta_t\}$ is a version of the SSEP with initial configuration η . Note that for each x, $\{\xi_t^x\}$ is a continuous-time simple random walk with jump rate one. Using this representation, for any $s, t \ge 0$, we have

$$\left\langle \mu_t^N - \mu_s^N, \delta_n \right\rangle = \frac{1}{a_N} \sum_{y \in \mathbb{Z}} \left[\delta_n \left(\frac{\xi_{tN^2}^y}{N} \right) - \delta_n \left(\frac{\xi_{sN^2}^y}{N} \right) \right] \eta_0(y).$$

Thus, we bound the second term in (5.16) by

$$\begin{split} \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le T} \left| \left\langle \mu_{t}^{N}, \delta_{n} \right\rangle \right| &- \sup_{0 \le i \le N^{3}} \left| \left\langle \mu_{t_{i}}^{N}, \delta_{n} \right\rangle \right| > \varepsilon/2 \Big) \\ & \le \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le i \le N^{3}} \sup_{t_{i} \le t \le t_{i+1}} \left| \left\langle \mu_{t}^{N} - \mu_{t_{i}}^{N}, \delta_{n} \right\rangle \right| > \varepsilon/2 \Big) \le N^{3} \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le t_{1}} \left| \left\langle \mu_{t}^{N} - \mu_{0}^{N}, \delta_{n} \right\rangle \right| > \varepsilon/2 \Big) \\ & = N^{3} \mathbb{P}_{\nu_{\rho}} \Big(\sup_{0 \le t \le t_{1}} \left| \frac{1}{a_{N}} \sum_{y \in \mathbb{Z}} \left[\delta_{n} \Big(\frac{\xi_{tN^{2}}^{y}}{N} \Big) - \delta_{n} \Big(\frac{y}{N} \Big) \right] \eta_{0}(y) \Big| > \varepsilon/2 \Big). \end{split}$$

To conclude the proof, it is sufficient to show that, for any n > 0,

$$\limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\sup_{0 \le t \le t_1} \left| \frac{1}{a_N} \sum_{y \in \mathbb{Z}} \left[\delta_n \left(\frac{\xi_{tN^2}^y}{N} \right) - \delta_n \left(\frac{y}{N} \right) \right] \eta_0(y) \right| > \varepsilon/2 \right) = -\infty.$$
(5.17)

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For the above sum over $y \in \mathbb{Z}$, we consider the two cases |y| > 3Nn and $|y| \le 3Nn$ respectively.

For the case |y| > 3Nn, $\delta_n(y/N) = 0$ since the support of δ_n is contained in [-2n, 2n]. Then, we bound

$$\begin{split} & \mathbb{P}_{\nu_{\rho}}\Big(\sup_{0\leq t\leq t_{1}}\Big|\frac{1}{a_{N}}\sum_{|y|>3Nn}\delta_{n}\Big(\frac{\xi_{tN^{2}}^{y}}{N}\Big)\eta_{0}(y)\Big|>\varepsilon/2\Big)\\ & \leq \mathbb{P}_{\nu_{\rho}}\Big(\sup_{0\leq t\leq t_{1}}|\xi_{tN^{2}}^{y}|\leq 2Nn \quad \text{for some } |y|>3Nn\Big)\\ & \leq \sum_{|y|>3Nn}\mathbb{P}_{\nu_{\rho}}\Big(\sup_{0\leq t\leq t_{1}}|\xi_{tN^{2}}^{y}-y|\geq |y|-2Nn\Big). \end{split}$$

Since for each y, $|\xi_{tN^2}^y - y|$ is stochastically bounded by a Poisson process with parameter tN^2 , and recall $t_1 = T/N^3$, the last line is bound by

$$\sum_{|y|>3Nn} e^{-c(|y|-2Nn)} \le e^{-cNn}$$

for some constant c > 0. Therefore,

$$\limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho} \Big(\sup_{0 \le t \le t_1} \Big| \frac{1}{a_N} \sum_{|y| > 3Nn} \delta_n \Big(\frac{\xi_{tN^2}^y}{N} \Big) \eta_0(y) \Big| > \varepsilon/2 \Big) = -\infty.$$

It remains to bound

$$\mathbb{P}_{\nu_{\rho}}\left(\sup_{0\leq t\leq t_{1}}\left|\frac{1}{a_{N}}\sum_{|y|\leq 3Nn}\left[\delta_{n}\left(\frac{\xi_{tN^{2}}^{y}}{N}\right)-\delta_{n}\left(\frac{y}{N}\right)\right]\eta_{0}(y)\right|>\varepsilon/2\right)\\ \leq \mathbb{P}_{\nu_{\rho}}(A_{N}^{c})+\mathbb{P}_{\nu_{\rho}}\left(\left\{\sup_{0\leq t\leq t_{1}}\left|\frac{1}{a_{N}}\sum_{|y|\leq 3Nn}\left[\delta_{n}\left(\frac{\xi_{tN^{2}}^{y}}{N}\right)-\delta_{n}\left(\frac{y}{N}\right)\right]\eta_{0}(y)\right|>\varepsilon/2\right\}\cap A_{N}\right),$$
(5.18)

where

$$A_N = \bigg\{ \sup_{|y| \le 3Nn} \sup_{0 \le t \le t_1} |\xi_{tN^2}^y - y| \le a_N^{3/2} N^{-1/2} \bigg\}.$$

Standard large deviation estimates yield

$$\mathbb{P}_{\nu_{\rho}}(A_{N}^{c}) \leq 7Nne^{-ca_{N}^{3/2}N^{-1/2}}$$

for some constant c > 0, and thus

$$\limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho}(A_N^c) = -\infty.$$

We claim that the event inside the second probability on the right-hand side of (5.18) is empty for N large enough, which is sufficient to conclude the proof. Indeed, first note that δ_n is only discontinuous at the origin. Moreover, on the event A_N , for each $0 \le t \le t_1$, the cardinality of

$$B_{N,t} := \{ |y| \le 3Nn : y \ge 0, \xi_{tN^2}^y < 0 \text{ or } y < 0, \xi_{tN^2}^y \ge 0 \}$$

is bounded by $2a_N^{3/2}N^{-1/2}$. Thus, on the event A_N ,

$$\sup_{0 \le t \le t_1} \left| \frac{1}{a_N} \sum_{y \in B_{N,t}} \left[\delta_n \left(\frac{\xi_{tN^2}^y}{N} \right) - \delta_n \left(\frac{y}{N} \right) \right] \eta_0(y) \right| \le \frac{4}{n} \sqrt{\frac{a_N}{N}},$$

which goes to zero as $N \to \infty$. By the piecewise smoothness of δ_n ,

$$\sup_{0 \le t \le t_1} \left| \frac{1}{a_N} \sum_{|y| \le 3Nn, y \notin B_{N,t}} \left[\delta_n \left(\frac{\xi_{tN^2}^y}{N} \right) - \delta_n \left(\frac{y}{N} \right) \right] \eta_0(y) \right| \le C(n) \sqrt{\frac{a_N}{N}}$$

which also vanishes as $N \to \infty$, thus concluding the proof.

5.3 Exponential tightness of the tagged particle

As in the last subsection, if follows from the following estimates that the sequence of processes $\{X(tN^2)/a_N: 0 \le t \le T\}_{N \ge 1}$ is exponentially tight.

Lemma 5.3. We have the following estimates for the replacement of the tagged particle:

1. for any fixed T > 0,

$$\limsup_{M \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(\sup_{0 \le t \le T} \frac{1}{a_N} |X(tN^2)| > M \right) = -\infty;$$
(5.19)

2. for any $\varepsilon > 0$,

$$\limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \Big(\sup_{0 < t \le \delta} \frac{1}{a_N} \big| X((t+\tau)N^2) - X(\tau N^2) \big| > \varepsilon \Big) = -\infty,$$
(5.20)

where \mathcal{T}_T is the set of all stopping times bounded by T.

Proof. We first prove (5.19). Since the process is symmetric with respect to the origin, we could remove the absolute value inside the probability in (5.19). If $X(tN^2) > a_NM$ for some $0 \le t \le T$, then $J_{-1,0}(tN^2) \ge \sum_{x=0}^{[a_NM]} \eta_{tN^2}(x)$ for some $0 \le t \le T$. Thus, the probability in (5.19) is bounded by

$$\mathbb{P}_{\nu_{\rho}^{*}}\Big(\sup_{0 \le t \le T} J_{-1,0}(tN^{2}) \ge \inf_{0 \le t \le T} \sum_{x=0}^{[a_{N}M]} \eta_{tN^{2}}(x)\Big) \le \mathbb{P}_{\nu_{\rho}^{*}}\Big(\sup_{0 \le t \le T} J_{-1,0}(tN^{2}) \ge \rho a_{N}M/2\Big) + \mathbb{P}_{\nu_{\rho}^{*}}\Big(\inf_{0 \le t \le T} \sum_{x=0}^{[a_{N}M]} \eta_{tN^{2}}(x) \le \rho a_{N}M/2\Big).$$

To finish the proof, it remains to show that

$$\limsup_{M \to \infty} \limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \Big(\sup_{0 \le t \le T} J_{-1,0}(tN^2) \ge \rho a_N M/2 \Big) = -\infty,$$

which follows directly from (5.8), and that, for any fixed M > 0,

$$\limsup_{N \to \infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \Big(\inf_{0 \le t \le T} \sum_{x=0}^{[a_N M]} \eta_{tN^2}(x) \le \rho a_N M/2 \Big) = -\infty.$$
(5.21)

Now, we prove (5.21). For $0 \le i \le N^2$, define $t_i = iT/N^2$. Then, for $t_i \le t \le t_{i+1}$,

$$\sum_{x=0}^{[a_NM]} \eta_{tN^2}(x) = \sum_{x=0}^{[a_NM]} \eta_{t_iN^2}(x) + J_{-1,0}(t_iN^2, tN^2) - J_{[a_NM], [a_NM]+1}(t_iN^2, tN^2),$$

which implies that

$$\begin{split} &\inf_{0\leq t\leq T}\sum_{x=0}^{[a_NM]}\eta_{tN^2}(x)\geq \inf_{0\leq i\leq N^2}\sum_{x=0}^{[a_NM]}\eta_{t_iN^2}(x) \\ &-\sup_{0\leq i\leq N^2}\sup_{t_i\leq t\leq t_{i+1}}\left|J_{-1,0}(t_iN^2,tN^2)\right|-\sup_{0\leq i\leq N^2}\sup_{t_i\leq t\leq t_{i+1}}\left|J_{[a_NM],[a_NM]+1}(t_iN^2,tN^2)\right|. \end{split}$$

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This permits us to bound the probability in (5.21) by

$$\mathbb{P}_{\nu_{\rho}^{*}}\left(\inf_{0\leq i\leq N^{2}}\sum_{x=0}^{[a_{N}M]}\eta_{t_{i}N^{2}}(x)\leq 3\rho a_{N}M/4\right) \\
+\mathbb{P}_{\nu_{\rho}^{*}}\left(\sup_{0\leq i\leq N^{2}}\sup_{t_{i}\leq t\leq t_{i+1}}\left|J_{-1,0}(t_{i}N^{2},tN^{2})\right|\geq \rho a_{N}M/8\right) \\
+\mathbb{P}_{\nu_{\rho}^{*}}\left(\sup_{0\leq i\leq N^{2}}\sup_{t_{i}\leq t\leq t_{i+1}}\left|J_{[a_{N}M],[a_{N}M]+1}(t_{i}N^{2},tN^{2})\right|\geq \rho a_{N}M/8\right)$$

By Remark 3.2, we could replace the above $\mathbb{P}_{\nu_{\rho}^*}$ with $\mathbb{P}_{\nu_{\rho}}$. Since ν_{ρ} is invariant for the SSEP, using the large deviation principle of the sum of i.i.d. random variables, there exists C > 0 independent of N and $1 \le i \le N^2$ such that

$$\limsup_{N \to +\infty} \frac{1}{a_N} \log \mathbb{P}_{\nu_{\rho}} \Big(\sum_{x=0}^{[a_N M]} \eta_{t_i N^2}(x) \le 3\rho a_N M/4 \Big) = -C$$

for all $1 \leq i \leq N^2$. Then, since $a_N \gg \sqrt{N \log N}$ and

$$\mathbb{P}_{\nu_{\rho}}\Big(\inf_{0 \le i \le N^2} \sum_{x=0}^{[a_N M]} \eta_{t_i N^2}(x) \le 3\rho a_N M/4\Big) \le (N^2 + 1) \mathbb{P}_{\nu_{\rho}}\Big(\sum_{x=0}^{[a_N M]} \eta_0(x) \le 3\rho a_N M/4\Big),$$

we have

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \Big(\inf_{0 \le i \le N^2} \sum_{x=0}^{[a_N M]} \eta_{t_i N^2}(x) \le 3\rho a_N M/4 \Big) = -\infty.$$
(5.22)

Since a particle crosses an edge at rate at most 1, $\{|J_{-1,0}(t_iN^2, tN^2)|\}_{t \ge t_i}$ is stochastically dominated from above by a Poisson process with rate 1. Thus,

$$\begin{split} \mathbb{P}_{\nu_{\rho}}(\sup_{0 \le i \le N^{2}} \sup_{t_{i} \le t \le t_{i+1}} \left| J_{-1,0}(t_{i}N^{2}, tN^{2}) \right| \ge \rho a_{N}M/8) \\ \le N^{2} \mathbb{P}_{\nu_{\rho}}(\sup_{t_{i} \le t \le t_{i+1}} \left| J_{-1,0}(t_{i}N^{2}, tN^{2}) \right| \ge \rho a_{N}M/8) \le CN^{2} \exp\{-Ca_{N}M\} \end{split}$$

for some constant C > 0, which implies

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\sup_{0 \le i \le N^2} \sup_{t_i \le t \le t_{i+1}} \left| J_{-1,0}(t_i N^2, t N^2) \right| \ge \rho a_N M/8 \right) = -\infty.$$
(5.23)

By translation invariance, (5.23) still holds when $J_{-1,0}(t_iN^2, tN^2)$ is replaced by

$$J_{[a_N M],[a_N M]+1}(t_i N^2, tN^2)$$

thus concluding the proof of (5.21).

Now we prove (5.20). We first remove the absolute value inside the probability in (5.20) as before. For any M>0, we define

$$B_{M,N,\delta} = \Big\{ \sup_{0 \le t \le T+\delta} \frac{1}{a_N} |X(tN^2)| \le M \Big\}.$$

By (5.19),

$$\limsup_{M \to +\infty} \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*}(B_{M,N,\delta}^c) = -\infty.$$

Hence, to prove (5.20), we only need to show that

$$\limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\sup_{0 < t \le \delta} \frac{1}{a_N} \left(X((t+\tau)N^2) - X(\tau N^2) \right) > \varepsilon, B_{N,M,\delta} \right) = -\infty$$
(5.24)

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for any M > 0. We further define

$$C_{N,M,\varepsilon,\delta} = \Big\{ \inf_{-a_N M \le k \le a_N M} \inf_{0 \le t \le T+\delta} \sum_{x=0}^{\lfloor a_N \varepsilon \rfloor} \eta_{tN^2}(x+k) \ge \frac{\rho a_N \varepsilon}{2} \Big\},$$

where the infimum is over integers k. Then, by (5.21),

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*}(C_{N,M,\varepsilon,\delta}^c) = -\infty.$$

Therefore, to prove (5.24), we only need to show that, for any M > 0,

$$\limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \Big(\sup_{0 < t \le \delta} \frac{1}{a_N} \left(X((t+\tau)N^2) - X(\tau N^2) \right) > \varepsilon, \\ B_{N,M,\delta}, C_{N,M,\varepsilon,\delta} \Big) = -\infty.$$
(5.25)

Conditioned on the event $B_{N,M,\delta} \bigcap C_{N,M,\varepsilon,\delta}$, we have $X(\tau N^2) = k$ for some $-a_N M \le k \le a_N M$. Moreover, if $\frac{1}{a_N} \left(X((t+\tau)N^2) - X(\tau N^2) \right) > \varepsilon$ for some $0 < t \le \delta$, then

$$J_{k-1,k}((t+\tau)N^2) - J_{k-1,k}(\tau N^2) = \sum_{\substack{x=k\\x=k}}^{X((t+\tau)N^2)} \eta_{(t+\tau)N^2}(x)$$
$$\geq \sum_{\substack{x=k\\x=k}}^{k+[a_N\varepsilon]} \eta_{(t+\tau)N^2}(x) = \sum_{\substack{x=0\\x=0}}^{[a_N\varepsilon]} \eta_{(t+\tau)N^2}(x+k) \ge \frac{\rho a_N\varepsilon}{2}$$

for some $-a_N M \le k \le a_N M$. Thus, to prove (5.25), we only need to show that

$$\limsup_{\delta \to 0} \limsup_{N \to \infty} \sup_{\tau \in \mathcal{T}_T} \frac{N}{a_N^2} \log \times \sum_{k:|k| \le a_N M} \mathbb{P}_{\nu_{\rho}^*} \left(\sup_{0 < t \le \delta} \frac{1}{a_N} \left(J_{k-1,k}((t+\tau)N^2) - J_{k-1,k}(\tau N^2) \right) \ge \frac{\rho \varepsilon}{2} \right) = -\infty, \quad (5.26)$$

which follows immediately from (5.9) since $a_N \gg \sqrt{N \log N}$. This concludes the proof. \Box

6 Proof of Theorems 2.1 and 2.2

In this section, we prove Theorems 2.1 and 2.2. We first state two lemmas concerning finite-dimensional moderate deviation principles for the current and the tagged particle. **Lemma 6.1.** For any $n \ge 1$ and for any $0 \le t_1 < t_2 < \ldots < t_n$, the sequence

$$\left\{\frac{1}{a_N}\left(J_{-1,0}(t_1N^2),\ldots,J_{-1,0}(t_nN^2)\right)\right\}_{N>1}$$

satisfies moderate deviation principles with decay rate a_N^2/N and with rate function $\sigma_J^{-2}\mathcal{I}_{\{t_j\}_{j=1}^n}$ under both $\mathbb{P}_{\nu_{\rho}}$ and $\mathbb{P}_{\nu_{\rho}^*}$. Recall the definition of $\mathcal{I}_{\{t_j\}_{j=1}^n}$ from (4.1). More precisely, for any closed set $\mathcal{C} \subseteq \mathbb{R}^n$,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}\left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2)\right) \in \mathcal{C}\right) \le -\inf_{\alpha \in \mathcal{C}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\alpha),$$
(6.1)

and for any open set $\mathcal{O} \subseteq \mathbb{R}^n$,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}\left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2)\right) \in \mathcal{O}\right) \ge -\inf_{\boldsymbol{\alpha} \in \mathcal{O}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}),$$
(6.2)

where $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$ or $\mathbb{P} = \mathbb{P}_{\nu_{\rho}^*}$.

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Lemma 6.2. For any $n \ge 1$ and for any $0 \le t_1 < t_2 < \ldots < t_n$, the sequence

$$\left\{\frac{1}{a_N}\left(X(t_1N^2),\ldots,X(t_nN^2)\right)\right\}_{N\geq 1}$$

satisfies moderate deviation principles with decay rate a_N^2/N and with rate function $\sigma_X^{-2}\mathcal{I}_{\{\iota_i\}_{i=1}^n}$. More precisely, for any closed set $\mathcal{C} \subseteq \mathbb{R}^n$,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{C} \right) \le -\inf_{\boldsymbol{\alpha} \in \mathcal{C}} \frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}), \quad (6.3)$$

and for any open set $\mathcal{O} \subseteq \mathbb{R}^n$,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{O} \right) \ge -\inf_{\boldsymbol{\alpha} \in \mathcal{O}} \frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}).$$
(6.4)

We shall prove Lemmas 6.1 and 6.2 in later subsections. Now we utilize these two lemmas to prove Theorems 2.1 and 2.2.

Proof of Theorems 2.1 and 2.2. We first prove Theorem 2.1. In Lemma 5.2, we have shown that the sequence $\{\frac{1}{a_N}J_{-1,0}(tN^2): 0 \le t \le T\}_{N\ge 1}$ is exponentially tight. Using [8, Theorem 4.28] and Lemma 6.1, the sequence $\{\frac{1}{a_N}J_{-1,0}(tN^2): 0 \le t \le T\}_{N\ge 1}$ satisfies moderate deviation principles with decay rate a_N^2/N and with rate function given by

$$\sup\left\{\frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}\left(f_{\{t_j\}_{j=1}^n}\right): \ n \ge 1, 0 \le t_1 < t_2 < \ldots < t_n \le T, t_j \in \Delta_f^c \text{ for all } 1 \le j \le n\right\}$$

for all $f \in \mathcal{D}([0,T],\mathbb{R})$, where $f_{\{t_j\}_{j=1}^n}$ and Δ_f^c are defined as in Proposition 4.1. Using Proposition 4.1 again, the last supremum equals

$$\frac{1}{\sigma_J^2} \mathcal{I}_{path}(f) = \mathcal{I}_{cur}(f),$$

which concludes the proof of Theorem 2.1. For Theorem 2.2, we use Lemmas 5.3 and 6.2 instead. $\hfill \Box$

6.1 Proof of Lemma 6.1

For $\boldsymbol{\alpha} \in \mathbb{R}^n$ and r > 0, let

$$B(\boldsymbol{\alpha}, r) = \{ \boldsymbol{\beta} \in \mathbb{R}^n : \| \boldsymbol{\beta} - \boldsymbol{\alpha} \|_{\infty} < r \}$$

be the cube of radius r centered at α . For any M > 0, shorten

$$\overline{B}_M = \overline{B(\mathbf{0}, M)} = \{ \boldsymbol{\alpha} \in \mathbb{R}^n : \| \boldsymbol{\alpha} \|_{\infty} \le M \}.$$

We first prove the lower bound (6.2).

Proof of (6.2). For any open set $\mathcal{O} \subseteq \mathbb{R}^n$ and for any point $\alpha \in \mathcal{O}$, using (5.1) and (5.10), for sufficiently small $\varepsilon > 0$ such that $B(\alpha, \varepsilon) \subseteq \mathcal{O}$,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{O} \right) \\
\geq \liminf_{m \to +\infty} \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \left\langle \mu_{t_j}^N, G_m \right\rangle - \left\langle \mu_0^N, G_m \right\rangle \right\}_{j=1}^n \in B(\boldsymbol{\alpha}, \varepsilon/2), \\
\left| \frac{1}{m N a_N} \sum_{x=0}^{m N-1} J_{x,x+1}(t_j N^2) \right| < \varepsilon/2 \text{ for all } 1 \le j \le n \right) \\
= \liminf_{m \to +\infty} \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \left\langle \mu_{t_j}^N, G_m \right\rangle - \left\langle \mu_0^N, G_m \right\rangle \right\}_{j=1}^n \in B(\boldsymbol{\alpha}, \varepsilon/2) \right). \quad (6.5)$$

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Note that $G_m \notin S(\mathbb{R})$. By (5.15) and Theorem 3.1, the last line is bounded from below by

$$\liminf_{m \to +\infty} \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \left\langle \mu_{t_j}^N, \tilde{G}_m \right\rangle - \left\langle \mu_0^N, \tilde{G}_m \right\rangle \right\}_{j=1}^n \in B(\boldsymbol{\alpha}, \varepsilon/4) \right) \\ \geq -\limsup_{m \to +\infty} \inf \left\{ \mathcal{Q}_{t_n}(\mu) : \left\{ \left\langle \mu_{t_j}, \tilde{G}_m \right\rangle - \left\langle \mu_0, \tilde{G}_m \right\rangle \right\}_{j=1}^n \in B(\boldsymbol{\alpha}, \varepsilon/4) \right\}.$$

Let $\mu_{t_n}^n := \mu_{t_n}^{n, \{t_j\}_{j=1}^n, \{\alpha_j\}_{j=1}^n}$ be the minimizer of the following variational problem

$$\inf\left\{\mathcal{Q}_{t_n}(\mu): \mu \in \mathcal{A}, \ \int_0^\infty [\mu(t_j, u) - \mu(0, u)] du = \alpha_j, \ 1 \le j \le n\right\}$$

which was introduced in the proof of Proposition 4.3. Since $\mu_{t_n}^n \in \mathcal{A}$, by the explicit expression of μ given in (3.2), $\mu_{t_n}^n(t, u) \in L^1(\mathbb{R})$ for any $t \in [0, t_n]$. Thus, by dominated convergence theorem, for all $1 \leq j \leq n$,

$$\lim_{n \to +\infty} \left\langle \mu_{t_n}^n(t_j, \cdot) - \mu_{t_n}^n(0, \cdot), \tilde{G}_m \right\rangle = \int_0^\infty \left[\mu_{t_n}^n(t_j, u) - \mu_{t_n}^n(0, u) \right] du = \alpha_j.$$
(6.6)

This implies that, for m large enough,

$$\left\{\left\langle \mu_{t_n}^n(t_j), \tilde{G}_m \right\rangle - \left\langle \mu_{t_n}^n(0), \tilde{G}_m \right\rangle \right\}_{j=1}^n \in B(\boldsymbol{\alpha}, \varepsilon/4).$$

Consequently,

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{O} \right)$$

$$\geq -\mathcal{Q}_{t_n}(\mu_{t_n}^n) = -\frac{\sqrt{2\pi}}{4\chi(\rho)} \boldsymbol{\alpha}^T A_{\{t_j\}_{j=1}^n}^{-1} \boldsymbol{\alpha} = -\frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}).$$

Since $oldsymbol{lpha} \in \mathcal{O}$ is arbitrary,

$$\begin{split} \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{O} \right) &\geq \sup_{\boldsymbol{\alpha} \in \mathcal{O}} -\frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}) \\ &= -\inf_{\boldsymbol{\alpha} \in \mathcal{O}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}). \end{split}$$

By Remark 3.2, the lower bound also holds under $\mathbb{P}_{\nu_a^*}$, thus concluding the proof. \Box

Now, we prove the upper bound (6.1). By exponential tightness of the current (see Lemma 5.2), we only need to prove the estimate for any compact subset in \mathbb{R}^n .

Proof of (6.1). As in the proof of the lower bound, we only need to deal with the case where $\mathbb{P} = \mathbb{P}_{\nu_{\rho}}$. Let $\mathcal{K} \subseteq \mathbb{R}^n$ be any compact subset. For any r > 0, we could find a sequence of points $\alpha^j \in \mathcal{K}$, $j = 1, 2, \ldots, \ell(r)$, such that

$$\mathcal{K} \subseteq \bigcup_{j=1}^{\ell} B(\boldsymbol{\alpha}^j, r).$$

By (5.10), for any m > 0 and any r > 0,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{K} \right) \\ \leq \max_{1 \le j \le \ell} \limsup_{N \to +\infty} \frac{N}{a_N^2} \\ \log \mathbb{P}_{\nu_{\rho}} \left(\left\{ \left\langle \mu_{t_i}^N, G_m \right\rangle - \left\langle \mu_0^N, G_m \right\rangle + \frac{1}{m N a_N} \sum_{x=0}^{mN-1} J_{x,x+1}(t_i N^2) \right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, r) \right),$$

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Using Lemma 5.1, the right-hand side of the last inequality is bounded from above by, for any r > 0,

$$\limsup_{m \to \infty} \max_{1 \le j \le \ell} \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho} \left(\left\{ \left\langle \mu_{t_i}^N, G_m \right\rangle - \left\langle \mu_0^N, G_m \right\rangle \right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, 2r) \right).$$

By (5.15), we could replace the above G_m with the Schwartz function \tilde{G}_m . Then, using moderate deviation principles from hydrodynamic limits (see Theorem 3.1), we finally bound the last expression from above by, for any r > 0,

$$\limsup_{m \to \infty} \max_{1 \le j \le \ell} -\inf \left\{ \mathcal{Q}_{t_n}(\mu) : \left\{ \left\langle \mu_{t_i}, \tilde{G}_m \right\rangle - \left\langle \mu_0, \tilde{G}_m \right\rangle \right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, 3r) \right\}.$$

We first show that there exists some constant C_0 independent of r, j, m such that, for m large enough,

$$\inf \left\{ \mathcal{Q}_{t_n}(\mu) : \left\{ \left\langle \mu_{t_i}, \tilde{G}_m \right\rangle - \left\langle \mu_0, \tilde{G}_m \right\rangle \right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, 3r) \right\} \le C_0.$$
(6.7)

Indeed, let $\mu_{t_n}^{n,j} := \mu_{t_n}^{n,\{t_i\}_{i=1}^n,\{\alpha_i^j\}_{i=1}^n}$ be as in the proof of (6.2), where α_i^j is the *i*-th component of α^j for $1 \le i \le n$. By (6.6), for *m* large enough,

$$\left\{\left\langle \mu_{t_n}^{n,j}(t_i) - \mu_{t_n}^{n,j}(0), \tilde{G}_m \right\rangle\right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, 3r).$$

Therefore, the infimum in (6.7) is bounded from above by

$$\mathcal{Q}_{t_n}(\mu_{t_n}^{n,j}) = \frac{\sqrt{2\pi}}{4\chi(\rho)} (\boldsymbol{\alpha}^j)^T A_{\{t_i\}_{i=1}^n}^{-1} \boldsymbol{\alpha}^j \le \sup_{\boldsymbol{\alpha} \in \mathcal{K}} \frac{\sqrt{2\pi}}{4\chi(\rho)} \boldsymbol{\alpha}^T A_{\{t_i\}_{i=1}^n}^{-1} \boldsymbol{\alpha} := C_0,$$

which proves (6.7).

For any r, j, m and any $\varepsilon > 0$, there exists $\mu_{\varepsilon}^{r,j,m}$ such that

$$\left\{\left\langle \mu_{\varepsilon}^{r,j,m}(t_{i},\cdot) - \mu_{\varepsilon}^{r,j,m}(0,\cdot), \tilde{G}_{m}\right\rangle\right\}_{i=1}^{n} \in B(\boldsymbol{\alpha}^{j},3r),\tag{6.8}$$

and

$$\mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r,j,m}) - \varepsilon \leq \inf \left\{ \mathcal{Q}_{t_n}(\mu) : \left\{ \left\langle \mu_{t_i}, \tilde{G}_m \right\rangle - \left\langle \mu_0, \tilde{G}_m \right\rangle \right\}_{i=1}^n \in B(\boldsymbol{\alpha}^j, 3r) \right\}.$$

Moreover, by (6.7),

$$\sup_{r,j,m} \mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r,j,m}) \le C_0 + \varepsilon.$$
(6.9)

So far, we have shown that, for any $\varepsilon > 0$,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{K} \right) \\
\leq -\liminf_{r \to 0} \liminf_{m \to +\infty} \min_{1 \le j \le \ell} \mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r,j,m}) + \varepsilon = -\liminf_{r \to 0} \liminf_{m \to +\infty} \mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r,\tilde{j},m}) + \varepsilon, \quad (6.10)$$

where $\tilde{j} := j(r, m, \varepsilon) := \arg \min_{1 \le j \le \ell} \mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r, j, m}).$

By (6.9), we could extract a subsequence along which the limit on the right-hand side of (6.10) is attained as $r \to 0, m \to +\infty$. According to Remark 3.6, $\{\mu_{\varepsilon}^{r,j,m} : r, j, m\}$ is a subset of the compact set $\mathbb{L}_{C_0+\varepsilon}$. Hence, by further extracting a subsequence, there exists μ_{ε}^* such that $\mu_{\varepsilon}^{r,\tilde{j},m}$ converges to μ_{ε}^* in $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$, and by the compactness of \mathcal{K} , there also exists some $\boldsymbol{\alpha} = (\alpha_1, \alpha_2, \ldots, \alpha_n) \in \mathcal{K}$ such that $\alpha^{\tilde{j}}$ converges to α along this subsequence. Then, by the lower-semicontinuity of \mathcal{Q}_{t_n} ,

$$\mathcal{Q}_{t_n}(\mu_{\varepsilon}^*) \leq \liminf_{r \to 0} \liminf_{m \to +\infty} \mathcal{Q}_{t_n}(\mu_{\varepsilon}^{r,j,m}),$$

and thus,

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{K} \right) \le -\mathcal{Q}_{t_n}(\mu_{\varepsilon}^*) + \varepsilon.$$
 (6.11)

Moreover, by Lemma 6.3 below,

$$\int_0^\infty \left[\mu_\varepsilon^*(t_i, u) - \mu_\varepsilon^*(0, u)\right] du = \alpha_i, \quad \forall 1 \le i \le n.$$
(6.12)

Therefore, by Proposition 4.2,

$$\begin{split} \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in \mathcal{K} \right) \\ &\leq - \inf_{\pmb{\alpha} \in \mathcal{K}} \inf_{\mu} \left\{ \mathcal{Q}_{t_n}(\mu) : \int_0^\infty \left[\mu(t_i, u) - \mu(0, u) \right] du = \alpha_i, \ \forall 1 \leq i \leq n \right\} + \varepsilon \\ &= - \inf_{\pmb{\alpha} \in \mathcal{K}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\pmb{\alpha}) + \varepsilon. \end{split}$$

We conclude the proof of (6.1) by letting $\varepsilon \to 0$.

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Lemma 6.3. If the sequence $\mu^m \in \mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$, $m \ge 1$, satisfies that

$$\lim_{m \to \infty} \mu^m = \mu \text{ in } \mathcal{D}([0,T], \mathcal{S}'(\mathbb{R})), \quad \sup_{m \ge 1} \mathcal{Q}_T(\mu_m) \le C_0$$

for some $\mu \in \mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ and for some finite constant C_0 , then for any $0 < t \leq T$,

$$\lim_{m\to\infty}\int_{\mathbb{R}}[\mu^m(t,u)-\mu^m(0,u)]\tilde{G}_m(u)\,du=\int_0^\infty[\mu(t,u)-\mu(0,u)]\,du.$$

Proof. For any $m \ge n$, we write

$$\int_{\mathbb{R}} [\mu^{m}(t,u) - \mu^{m}(0,u)] \tilde{G}_{m}(u) \, du$$

=
$$\int_{\mathbb{R}} [\mu^{m}(t,u) - \mu^{m}(0,u)] [\tilde{G}_{m}(u) - \tilde{G}_{n}(u)] \, du + \int_{\mathbb{R}} [\mu^{m}(t,u) - \mu^{m}(0,u)] \tilde{G}_{n}(u) \, du. \quad (6.13)$$

We first prove that the first term on the right-hand side in the last identity converges to zero as $m \to \infty, n \to \infty$. By Lemma 3.3, Eqn. (3.2) and the uniform boundedness of $Q_T(\mu_m)$, there exists some constant C_1 such that, for any t > 0,

$$\sup_{m\geq 1} \|\mu^m(t,\cdot)\|_{L^2(\mathbb{R})} \leq C_1.$$

Since $\|G_m - \tilde{G}_m\|_{L^2(\mathbb{R})}^2 \le Cm^{-1}$, by Cauchy-Schwarz inequality, we only need to show that

$$\lim_{n \to \infty} \lim_{m \to \infty} \int_{\mathbb{R}} [\mu^m(t, u) - \mu^m(0, u)] [G_m(u) - G_n(u)] \, du = 0.$$

We could also find $\tilde{\mu}^m \in \mathcal{A}$ such that $\|\tilde{\mu}^m(t,\cdot) - \mu^m(t,\cdot)\|_{L^2(\mathbb{R})}^2 \leq Cm^{-2}$. Since $\|G_m\|_{L^2(\mathbb{R})}^2 \leq Cm$, using Cauchy-Schwarz inequality again, it is sufficient to show that

$$\lim_{n \to \infty} \lim_{m \to \infty} \int_{\mathbb{R}} [\tilde{\mu}^m(t, u) - \tilde{\mu}^m(0, u)] [G_m(u) - G_n(u)] \, du = 0.$$
(6.14)

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Let \tilde{J}^m be the current associated with $\tilde{\mu}^m$, i.e., $\partial_t \tilde{\mu}^m + \partial_u \tilde{J}^m = 0$. Since

$$G_m(u) - G_n(u) = \begin{cases} 0 & \text{if } u \le 0 \text{ or } u \ge m; \\ \frac{u}{n} - \frac{u}{m} & \text{if } 0 \le u \le n; \\ 1 - \frac{u}{m} & \text{if } n \le u \le m, \end{cases}$$

using integration by parts formula, we rewrite the left-hand side in (6.14) as

$$\left(\frac{1}{n} - \frac{1}{m}\right) \int_0^t \int_0^n \tilde{J}^m(s, u) du ds - \frac{1}{m} \int_0^t \int_n^m \tilde{J}^m(s, u) du ds.$$

Note that by (3.6) and the uniform boundedness of $\mathcal{Q}_T(\mu_m)$, there exists some constant C_2 such that $\sup_{0 \le t \le T} \sup_{m \ge 1} \|\tilde{J}^m(t, \cdot)\|_{L^2(\mathbb{R})} \le C_2$. Then, by Cauchy-Schwarz inequality, the last expression is bounded by $C(n^{-1/2} + m^{-1/2})$, concluding the proof of (6.14).

For the second term in (6.13), letting $m \to \infty$, it converges to

$$\int_{\mathbb{R}} [\mu(t, u) - \mu(0, u)] \tilde{G}_n(u) du.$$

Note that $\mu(t, \cdot)$ may not be in $L^1(\mathbb{R})$, and thus we cannot use dominated convergence theorem directly. Instead, we first replace $\tilde{G}_n(u)$ with $G_n(u)$ as in the above argument, and then only need to deal with

$$\int_{\mathbb{R}} [\mu(t,u) - \mu(0,u)] G_n(u) du = \int_0^n [\mu(t,u) - \mu(0,u)] du - \frac{1}{n} \int_0^n u[\mu(t,u) - \mu(0,u)] du$$

by the definition of G_n . At last, letting $n \to \infty$ and using Lemma 3.4 and Remark 3.5, we conclude the proof.

6.2 Proof of Lemma 6.2

We start with the lower bound (6.4).

Proof of (6.4). Observe that if $J_{-1,0}(t) \ge 0$, then

$$J_{-1,0}(t) = \sum_{x=0}^{X(t)-1} \eta_t(x) = \sum_{x=0}^{X(t)-1} (\eta_t(x) - \rho) + \rho X(t),$$
(6.15)

and if $J_{-1,0}(t) < 0$, then

$$J_{-1,0}(t) = -\sum_{x=X(t)}^{-1} \eta_t(x) = -\sum_{x=X(t)}^{-1} (\eta_t(x) - \rho) + \rho X(t).$$
(6.16)

For any given $\boldsymbol{\alpha} = (\alpha_1, \dots, \alpha_n)^T \in \mathcal{O}$ and any $\varepsilon > 0$, by (6.2),

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \alpha, \varepsilon \rho/10) \right)$$

$$\geq - \inf_{\boldsymbol{\beta} \in B(\rho \alpha, \varepsilon \rho/10)} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\beta}) > -\infty.$$
(6.17)

For M > 0, let $D_{M,N}$ be the event that $|X(t_jN^2)| \le a_N M$ for all $1 \le j \le n$. Using (5.19), we have

$$\limsup_{M \to +\infty} \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(D_{M,N}^c \right) = -\infty.$$

Hence, there exists $M_1 > 0$ such that

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10), D_{M_1, N} \right) \\
= \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10) \right) \\
\geq - \inf_{\boldsymbol{\beta} \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10)} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\beta}). \quad (6.18)$$

Let F_N be the event that $\left|\frac{1}{a_N}\sum_{x=0}^k (\eta_{t_jN^2}(x) - \rho)\right| < \frac{\rho\varepsilon}{10}$ and $\left|\frac{1}{a_N}\sum_{x=-k}^{-1} (\eta_{t_jN^2}(x) - \rho)\right| < \frac{\rho\varepsilon}{10}$ for all $1 \le k \le M_1 a_N$ and all $1 \le j \le n$. By Remark 3.2 and standard large deviation results, we have

$$\limsup_{N \to +\infty} \frac{1}{a_N} \log \mathbb{P}_{\nu_{\rho}^*} \left(F_N^c \right) < 0$$

and thus

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(F_N^c \right) = -\infty.$$
(6.19)

Using (6.18) and (6.19), we have

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10), D_{M_1,N}, F_N \right)$$

$$\geq - \inf_{\boldsymbol{\beta} \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10)} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\beta}).$$
(6.20)

Conditioned on the event $D_{M_1,N} \bigcap F_N$, according to (6.15) and (6.16), if

$$\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \alpha, \varepsilon \rho/10),$$

then

$$\left|\frac{1}{a_N}X(t_jN^2) - \alpha_j\right| < \varepsilon/5 < \varepsilon, \quad \forall 1 \le j \le n.$$

Consequently, for sufficiently small $\varepsilon > 0$ such that $B(\alpha, \varepsilon) \subseteq \mathcal{O}$, we have

$$\begin{split} & \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{O} \right) \\ & \geq \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in B(\boldsymbol{\alpha}, \varepsilon) \right) \\ & \geq \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10), D_{M_1,N}, F_N \right) \\ & \geq - \inf_{\boldsymbol{\beta} \in B(\rho \boldsymbol{\alpha}, \varepsilon \rho/10)} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n} \left(\boldsymbol{\beta} \right) \end{split}$$

Since ε is arbitrary and $\mathcal{I}_{_{\{t_j\}_{j=1}^n}}$ is continuous, let $\varepsilon\to 0,$ we have

$$\liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{O} \right) \ge -\frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\rho \alpha)$$
$$= -\frac{1}{\sigma_Y^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\alpha).$$

Since $oldsymbol{lpha} \in \mathcal{O}$ is arbitrary,

$$\begin{split} \liminf_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{O} \right) &\geq \sup_{\boldsymbol{\alpha} \in \mathcal{O}} \left(-\frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}) \right) \\ &= -\inf_{\boldsymbol{\alpha} \in \mathcal{O}} \frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}) \end{split}$$

the proof is completed. \Box

and the proof is completed.

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At last, we prove the upper bound (6.3).

Proof of (6.3). According to Lemma 5.3, we only need to prove (6.3) for all compact set $\mathcal{K} \subseteq \mathbb{R}^n$. For any compact set $\mathcal{K} \subseteq \mathbb{R}^n$, there exists $M_2 > 0$ such that $\mathcal{K} \subseteq \overline{B}_{M_2}$. For given $0 < \varepsilon < 1$ and each $N \ge 1$, we denote by V_N the event that $\left|\frac{1}{a_N}\sum_{x=0}^k \left(\eta_{t_jN^2}(x) - \rho\right)\right| \le \varepsilon$ and $\left|\frac{1}{a_N}\sum_{x=-k}^{-1} \left(\eta_{t_jN^2}(x) - \rho\right)\right| \le \varepsilon$ for all $1 \le k \le (M_2 + 1)a_N$ and $1 \le j \le n$. Then, similar with (6.19), we have

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*}(V_N^c) = -\infty$$

and thus

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{K} \right)$$
$$= \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_\rho^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{K}, V_N \right). \quad (6.21)$$

Conditioned on V_N , if $\frac{1}{a_N} (X(t_1N^2), \dots, X(t_nN^2)) \in \mathcal{K}$, then, by (6.15) and (6.16),

$$\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in (\rho \mathcal{K})_{\varepsilon} \,,$$

where $\rho \mathcal{K} = \{\rho \beta : \beta \in \mathcal{K}\}$ and for some subset $\mathcal{A} \subset \mathbb{R}^n$ and some a > 0,

$$\mathcal{A}_a := \{ oldsymbol{eta} \in \mathbb{R}^n: \ \|oldsymbol{eta} - oldsymbol{\gamma}\|_\infty \leq a ext{ for some } oldsymbol{\gamma} \in \mathcal{A} \}.$$

Since \mathcal{K} is compact, $(\rho \mathcal{K})_{\varepsilon}$ is also compact for all $\varepsilon > 0$. As a result, by (6.1), for any $0 < \varepsilon < 1$,

$$\begin{split} \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{K}, V_N \right) \\ & \leq \limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(J_{-1,0}(t_1 N^2), \dots, J_{-1,0}(t_n N^2) \right) \in (\rho \mathcal{K})_{\varepsilon} \right) \\ & \leq -\inf_{\beta \in (\rho \mathcal{K})_{\varepsilon}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\beta). \end{split}$$

Since $\mathcal{I}_{\{t_i\}_{i=1}^n}$ is uniformly continuous on \bar{B}_{M_2+1} and $(\rho K)_{\varepsilon}$ is compact,

$$\lim_{\varepsilon \to 0} \inf_{\beta \in (\rho \mathcal{K})_{\varepsilon}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\beta) = \inf_{\beta \in \rho \mathcal{K}} \frac{1}{\sigma_J^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\beta) = \inf_{\alpha \in \mathcal{K}} \frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\alpha).$$

Therefore, by (6.21),

$$\limsup_{N \to +\infty} \frac{N}{a_N^2} \log \mathbb{P}_{\nu_{\rho}^*} \left(\frac{1}{a_N} \left(X(t_1 N^2), \dots, X(t_n N^2) \right) \in \mathcal{K} \right) \le - \inf_{\boldsymbol{\alpha} \in \mathcal{K}} \frac{1}{\sigma_X^2} \mathcal{I}_{\{t_j\}_{j=1}^n}(\boldsymbol{\alpha}),$$

and the proof is completed.

A Calculations

A.1 Proof of Lemma 3.3

We only prove the statement for $\mathcal{Q}_{T,dyn}$, and the remaining is much simpler. If $\mathcal{Q}_{T,dyn}(\mu) < \infty$, then there exists some constant A > 0 such that, for any $G \in \mathcal{C}_c^{1,\infty}([0,T] \times \mathbb{R})$,

$$l(\mu, G) - \frac{\chi(\rho)}{2}[G, G] \le A.$$

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Replacing *G* with aG for any $a \in \mathbb{R}$, we have that

$$al(\mu, G) - a^2 \frac{\chi(\rho)}{2} [G, G] \le A, \quad \forall a \in \mathbb{R}.$$

Optimizing over *a*,

$$l(\mu, G)^2 \le 2\chi(\rho)A[G, G], \quad \forall G \in \mathcal{C}^{1,\infty}_c([0, T] \times \mathbb{R}).$$

Since $\mathcal{C}_c^{1,\infty}([0,T]\times\mathbb{R})$ is dense in \mathcal{H}^1 , $l(\mu,\cdot)$ is a linear bounded functional on \mathcal{H}^1 . Thus, there exists some function $H\in\mathcal{H}^1$ such that

$$l(\mu, G) = \chi(\rho)[H, G], \quad \forall G \in \mathcal{H}^1.$$

In particular,

$$\mathcal{Q}_{T,dyn}(\mu) = \chi(\rho) \sup_{G \in \mathcal{C}_c^{1,\infty}([0,T] \times \mathbb{R})} \left\{ [H,G] - \frac{1}{2} [G,G] \right\} = \frac{\chi(\rho)}{2} [H,H],$$

which concludes the proof.

A.2 Proof of (3.5)

In this subsection, we prove that, for any t > 0,

$$\int_0^t p_s'(\cdot) \, ds \in L^2(\mathbb{R}),$$

where $p_s(\cdot)$ is the heat kernel. By direct calculations,

$$p_s'(u) = -\frac{u}{s}p_s(u).$$

Making the change of variables $u^2/2s \mapsto \tau$,

$$\int_0^t p_s'(\cdot) \, ds = \int_0^t -\frac{u}{\sqrt{2\pi}s^{3/2}} e^{-u^2/(2s)} ds = -\int_{u^2/2t}^{+\infty} \frac{1}{\sqrt{\pi\tau}} e^{-\tau} \, d\tau.$$

Thus,

$$\Big|\int_0^t p_s'(u)\,ds\Big|\leq C\min\Big\{\frac{\sqrt{t}}{|u|}e^{-u^2/(2t)},1\Big\}$$

for some finite constant C, which concludes the proof.

A.3 Proof of (3.8)

Recall in Lemma 3.3 we have shown that

$$\mathcal{Q}_{T,dyn}(\mu) = \frac{\chi(\rho)}{2} \int_0^T \int_{\mathbb{R}} \left(\partial_u H\right)^2 (t, u) du dt$$

Replacing $\partial_u H$ with $\chi(\rho)^{-1}[J + \frac{1}{2}\partial_u\mu]$, we have

$$\chi(\rho)\mathcal{Q}_{T,dyn}(\mu) = \int_0^T \int_{\mathbb{R}} \left\{ \frac{1}{2} J^2(t,u) + \frac{1}{8} \left(\partial_u \mu\right)^2(t,u) + \frac{1}{2} J(t,u) \partial_u \mu(t,u) \right\} du dt.$$

Integrating by parts and using $\partial_t \mu + \partial_u J = 0$,

$$\int_0^T \int_{\mathbb{R}} \frac{1}{2} J(t, u) \partial_u \mu(t, u) du dt = \frac{1}{4} \int_{\mathbb{R}} \left\{ \mu^2(T, u) - \mu_0^2(u) \right\} du.$$

Thus,

$$\chi(\rho)\mathcal{Q}_{T,dyn}(\mu) = \int_0^T \int_{\mathbb{R}} \left\{ \frac{1}{2} J^2(t,u) + \frac{1}{8} \left(\partial_u \mu \right)^2(t,u) \right\} du dt + \frac{1}{4} \int_{\mathbb{R}} \left\{ \mu^2(T,u) - \mu_0^2(u) \right\} du.$$

Finally, we obtain (3.8) by using (3.7) and the formula for Q_0 in Lemma 3.3.

A.4 Proof of (4.23)

Recall that for $0 \le s \le t$,

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}\xi^2} \Big[e^{-(t-s)\xi^2/2} - e^{-s\xi^2/2} + 1 - e^{-t\xi^2/2} \Big],$$

and for $t \leq s \leq T$,

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}\xi^2} \left[e^{-(s-t)\xi^2/2} - e^{-s\xi^2/2} + 1 - e^{-t\xi^2/2} \right].$$

Using the integeral formula

$$\int_{\mathbb{R}} \frac{1 - e^{-ax^2}}{x^2} dx = 2\sqrt{\pi a}, \quad a > 0,$$
(A.1)

one easily proves (4.23). To prove (A.1), the left-hand side of the identity equals

$$\int_{\mathbb{R}} \int_0^a e^{-yx^2} dy dx = \int_0^a \frac{\sqrt{\pi}}{\sqrt{y}} dy = 2\sqrt{\pi a},$$

as claimed.

A.5 Proof of (4.25)

We prove (4.25) by calculus of variations. For any $\mu_0, \hat{\mu}_0 \in \mathcal{C}^2_c(\mathbb{R})$ and $K, \hat{K} \in \mathcal{C}^{1,2}_c([0,T] \times \mathbb{R})$, we denote by $\Lambda\left((\mu_0, K), (\hat{\mu}_0, \hat{K})\right)$ the positive-definite quadratic form

$$\begin{split} &\int_0^T \int_{\mathbb{R}} \frac{1}{2} \partial_t K(t, u) \partial_t \hat{K}(t, u) du dt \\ &+ \int_0^T \int_{\mathbb{R}} \frac{1}{8} \left(\partial_u \mu_0(u) - \partial_u^2 K(t, u) \right) \left(\partial_u \hat{\mu}_0(u) - \partial_u^2 \hat{K}(t, u) \right) du dt \\ &+ \int_{\mathbb{R}} \frac{1}{2} \left(\mu_0(u) - \frac{1}{2} \partial_u K(T, u) \right) \left(\hat{\mu}_0(u) - \frac{1}{2} \partial_u \hat{K}(T, u) \right) du + \int_{\mathbb{R}} \frac{1}{8} \partial_u K(T, u) \partial_u \hat{K}(T, u) du. \end{split}$$

Then, using (3.8), we have

$$\chi(\rho)Q_T(\mu_0, K) = \Lambda((\mu_0, K), (\mu_0, K)).$$
(A.2)

For any given $\hat{\mu}_0 \in \mathcal{C}^2_c(\mathbb{R})$ and $\hat{K} \in \mathcal{C}^{1,2}_c([0,T] \times \mathbb{R})$ such that $\hat{K}(0, \cdot) = 0$ and $\hat{K}(t_j, 0) = 0$, we define

$$G_{j}(\varepsilon) = \Lambda\left((\mu_{T,0}^{t_{j},1}, K_{T}^{t_{j},1}) + \varepsilon(\hat{\mu}_{0}, \hat{K}), (\mu_{T,0}^{t_{j},1}, K_{T}^{t_{j},1}) + \varepsilon(\hat{\mu}_{0}, \hat{K})\right)$$

for all $\varepsilon \in \mathbb{R}$. By Lemma 4.5,

$$\frac{d}{d\varepsilon}G_j(\varepsilon)\Big|_{\varepsilon=0}=0$$

and hence

$$\Lambda((\mu_{T,0}^{t_j,1}, K_T^{t_j,1}), (\hat{\mu}_0, \hat{K})) = 0.$$
(A.3)

For any given $\mu_0 \in \mathcal{C}^2_c(\mathbb{R})$ and $K \in \mathcal{C}^{1,2}_c([0,T] \times \mathbb{R})$ such that $K(0, \cdot) = 0$ and $K(t_j, 0) = \alpha_j$ for all $1 \leq j \leq n$, let

$$\tilde{K} = K - \sum_{j=1}^{n} \beta_j K_T^{t_j,1}$$
 and $\tilde{\mu}_0 = \mu_0 - \sum_{j=1}^{n} \beta_j \mu_{T,0}^{t_j,1}$,

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where $\{\beta_j\}_{j\geq 1}$ satisfy (4.26), then

$$\tilde{K}(0, \cdot) = 0$$
 and $\tilde{K}(t_j, 0) = 0$

for all $1 \leq j \leq n$. Consequently, by (A.3),

$$\Lambda\left(\left(\sum_{j=1}^{n}\beta_{j}\mu_{T,0}^{t_{j},1},\sum_{j=1}^{n}\beta_{j}K_{T}^{t_{j},1}\right),(\tilde{\mu}_{0},\tilde{K})\right)=0.$$
(A.4)

For any $\varepsilon \in \mathbb{R}$, let

$$\tilde{G}(\varepsilon) = \Lambda \left(\left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right) + \varepsilon(\tilde{\mu}_{0}, \tilde{K}), \left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right) + \varepsilon(\tilde{\mu}_{0}, \tilde{K}) \right).$$
By (A.4), for all $\varepsilon \in \mathbb{R}$.

By (A.4), for all $\varepsilon \in \mathbb{R}$,

$$\begin{split} \tilde{G}(\varepsilon) &= \Lambda \left(\left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right), \left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right) \right) + \varepsilon^{2} \Lambda(\tilde{\mu}_{0}, \tilde{K}) \\ &\geq \Lambda \left(\left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right), \left(\sum_{j=1}^{n} \beta_{j} \mu_{T,0}^{t_{j},1}, \sum_{j=1}^{n} \beta_{j} K_{T}^{t_{j},1} \right) \right) \right). \end{split}$$

Therefore,

$$\Lambda\left((\mu_0, K), (\mu_0, K)\right) = \tilde{G}(1) \ge \Lambda\left(\left(\sum_{j=1}^n \beta_j \mu_{T,0}^{t_j,1}, \sum_{j=1}^n \beta_j K_T^{t_j,1}\right), \left(\sum_{j=1}^n \beta_j \mu_{T,0}^{t_j,1}, \sum_{j=1}^n \beta_j K_T^{t_j,1}\right)\right)$$
and (4.25) follows from (A.2)

and (4.25) follows from (A.2).

A.6 Proof of (4.28)

As in the proof of Proposition 4.3, q_{jk} is the real part of

$$\begin{split} \int_{0}^{T} \int_{\mathbb{R}} \Big\{ \frac{1}{2} \partial_{s} \mathcal{F} K_{T}^{t_{j},1}(s,\xi) \overline{\partial_{s} \mathcal{F} K_{T}^{t_{k},1}(s,\xi)} \\ &+ \frac{1}{8} \Big[-i\xi \mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) + \xi^{2} \mathcal{F} K_{T}^{t_{j},1}(s,\xi) \Big] \overline{\big[-i\xi \mathcal{F} \mu_{T,0}^{t_{k},1}(\xi) + \xi^{2} \mathcal{F} K_{T}^{t_{k},1}(s,\xi) \big]} \Big\} d\xi \, ds \\ &+ \frac{1}{4} \int_{\mathbb{R}} \Big\{ \mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) \overline{\mathcal{F} \mu_{T,0}^{t_{k},1}(\xi)} \\ &+ \big[\mathcal{F} \mu_{T,0}^{t_{j},1}(\xi) + i\xi \mathcal{F} K_{T}^{t_{j},1}(T,\xi) \big] \overline{\big[\mathcal{F} \mu_{T,0}^{t_{k},1}(\xi) + i\xi \mathcal{F} K_{T}^{t_{k},1}(T,\xi) \big]} \Big\} d\xi. \quad (A.5) \end{split}$$

Also recall from (P1) and (P2) after the proof of Lemma 4.5 that

$$\mathcal{F}\mu_{T,0}^{t,\alpha}(\xi) = (-i\xi)\frac{\alpha}{2\sqrt{t}\xi^2} [1 - e^{-t\xi^2/2}],$$

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}\xi^2} [1 - e^{-t\xi^2/2} + e^{-(t-s)\xi^2/2} - e^{-s\xi^2/2}], \quad 0 \le s \le t,$$

$$\mathcal{F}K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}\xi^2} [1 - e^{-t\xi^2/2} + e^{-(s-t)\xi^2/2} - e^{-s\xi^2/2}], \quad t \le s \le T.$$

Without loss of generality, assume $t_j < t_k.$ Since for $0 \leq s \leq t$,

$$\partial_s \mathcal{F} K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{4\sqrt{t}} [e^{-(t-s)\xi^2/2} + e^{-s\xi^2/2}],$$
$$-i\xi \mathcal{F} \mu_{T,0}^{t,\alpha}(\xi) + \xi^2 \mathcal{F} K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}} [e^{-(t-s)\xi^2/2} - e^{-s\xi^2/2}],$$

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and for $t \leq s \leq T$,

$$\partial_s \mathcal{F} K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{4\sqrt{t}} [-e^{-(s-t)\xi^2/2} + e^{-s\xi^2/2}],$$
$$-i\xi \mathcal{F} \mu_{T,0}^{t,\alpha}(\xi) + \xi^2 \mathcal{F} K_T^{t,\alpha}(s,\xi) = \frac{\alpha}{2\sqrt{t}} [e^{-(s-t)\xi^2/2} - e^{-s\xi^2/2}],$$

the first two lines in (A.5) equal $A(0,t_j) + A(t_j,t_k) + A(t_k,T)$, where

$$\begin{split} A(0,t_j) &= \frac{1}{32\sqrt{t_j t_k}} \int_0^{t_j} \int_{\mathbb{R}} \left\{ [e^{-(t_j-s)\xi^2/2} + e^{-s\xi^2/2}] [e^{-(t_k-s)\xi^2/2} + e^{-s\xi^2/2}] \\ &+ [e^{-(t_j-s)\xi^2/2} - e^{-s\xi^2/2}] [e^{-(t_k-s)\xi^2/2} - e^{-s\xi^2/2}] \right\} d\xi \, ds \\ &= \frac{1}{16\sqrt{t_j t_k}} \int_0^{t_j} \int_{\mathbb{R}} \left\{ e^{-(\frac{t_j+t_k}{2}-s)\xi^2} + e^{-s\xi^2} \right\} d\xi \, ds = \frac{\sqrt{\pi}}{16\sqrt{t_j t_k}} \int_0^{t_j} \left\{ \frac{1}{\sqrt{\frac{t_j+t_k}{2}-s}} + \frac{1}{\sqrt{s}} \right\} ds \\ &= \frac{\sqrt{\pi}}{8\sqrt{t_j t_k}} \left\{ \sqrt{\frac{t_j+t_k}{2}} - \sqrt{\frac{t_k-t_j}{2}} + \sqrt{t_j} \right\}, \end{split}$$

$$\begin{split} A(t_j, t_k) &= \frac{1}{32\sqrt{t_j t_k}} \int_{t_j}^{t_k} \int_{\mathbb{R}} \Big\{ [e^{-(t_k - s)\xi^2/2} + e^{-s\xi^2/2}] [-e^{-(s - t_j)\xi^2/2} + e^{-s\xi^2/2}] \\ &+ [e^{-(t_k - s)\xi^2/2} - e^{-s\xi^2/2}] [e^{-(s - t_j)\xi^2/2} - e^{-s\xi^2/2}] \Big\} d\xi \, ds \\ &= \frac{1}{16\sqrt{t_j t_k}} \int_{t_j}^{t_k} \int_{\mathbb{R}} \Big\{ e^{-s\xi^2} - e^{-(s - \frac{t_j}{2})\xi^2} \Big\} d\xi \, ds = \frac{\sqrt{\pi}}{8\sqrt{t_j t_k}} \Big\{ \sqrt{t_k} - \sqrt{t_j} - \sqrt{t_k - \frac{t_j}{2}} + \sqrt{\frac{t_j}{2}} \Big\}, \end{split}$$

 $\quad \text{and} \quad$

$$\begin{split} A(t_k,T) &= \frac{1}{32\sqrt{t_j t_k}} \int_{t_k}^T \int_{\mathbb{R}} \Big\{ [-e^{-(s-t_k)\xi^2/2} + e^{-s\xi^2/2}] [-e^{-(s-t_j)\xi^2/2} + e^{-s\xi^2/2}] \\ &+ [e^{-(s-t_k)\xi^2/2} - e^{-s\xi^2/2}] [e^{-(s-t_j)\xi^2/2} - e^{-s\xi^2/2}] \Big\} d\xi ds \\ &= \frac{1}{16\sqrt{t_j t_k}} \int_{t_k}^T \int_{\mathbb{R}} \Big\{ e^{-(s-\frac{t_k+t_j}{2})\xi^2} + e^{-s\xi^2} - e^{-(s-\frac{t_j}{2})\xi^2} - e^{-(s-\frac{t_k}{2})\xi^2} \Big\} d\xi ds \\ &= \frac{\sqrt{\pi}}{8\sqrt{t_j t_k}} \Big\{ \sqrt{T - \frac{t_k+t_j}{2}} - \sqrt{\frac{t_k-t_j}{2}} + \sqrt{T} - \sqrt{t_k} - \sqrt{T - \frac{t_j}{2}} + \sqrt{t_k - \frac{t_j}{2}} - \sqrt{T - \frac{t_k}{2}} + \sqrt{\frac{t_k}{2}} \Big\}. \end{split}$$

Thus, the first two lines in (A.5) equal

$$\frac{\sqrt{\pi}}{8\sqrt{t_j t_k}} \Big\{ \sqrt{\frac{t_j + t_k}{2}} - 2\sqrt{\frac{t_k - t_j}{2}} + \sqrt{T} + \sqrt{\frac{t_j}{2}} + \sqrt{\frac{t_k}{2}} + \sqrt{T - \frac{t_k + t_j}{2}} - \sqrt{T - \frac{t_j}{2}} - \sqrt{T - \frac{t_k}{2}} \Big\}.$$

Similarly, since

$$\mathcal{F}\mu_{T,0}^{t,\alpha}(\xi) + i\xi \mathcal{F}K_T^{t,\alpha}(T,\xi) = \frac{i\alpha}{2\xi\sqrt{t}} [e^{-(T-t)\xi^2/2} - e^{-T\xi^2/2}],$$

the last two lines in (A.5) equal

$$\frac{1}{16\sqrt{t_j t_k}} \int_{\mathbb{R}} \xi^{-2} \Big\{ [1 - e^{-t_j \xi^2/2}] [1 - e^{-t_k \xi^2/2}] + [e^{-(T - t_j) \xi^2/2} - e^{-T \xi^2/2}] \\ \times [e^{-(T - t_k) \xi^2/2} - e^{-T \xi^2/2}] \Big\} d\xi.$$

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Using the integeral formula

$$\int_{\mathbb{R}} \frac{1 - e^{-ax^2}}{x^2} dx = 2\sqrt{\pi a}, \quad a > 0,$$

the last two lines in (A.5) equal

$$\frac{\sqrt{\pi}}{8\sqrt{t_j t_k}} \Big\{ \sqrt{\frac{t_j}{2}} + \sqrt{\frac{t_k}{2}} - \sqrt{\frac{t_j + t_k}{2}} - \sqrt{T - \frac{t_j + t_k}{2}} - \sqrt{T + \sqrt{T - \frac{t_k}{2}}} + \sqrt{T - \frac{t_j}{2}} \Big\}.$$

Therefore,

$$q_{jk} = \frac{\sqrt{2\pi}}{8\sqrt{t_j t_k}} (\sqrt{t_k} + \sqrt{t_j} - \sqrt{t_k - t_j}),$$

as claimed.

A.7 Proof of (5.7)

By Feynman-Kac formula ([19, Lemma A.1.7.2]), we may bound the expression in (5.7) by

$$\frac{kN^2}{a_N^2} \sup_{f:\nu_{\rho}-\text{density}} \Big\{ \int \frac{a_N K}{nN^2} \sum_{x=0}^{nN-1} \big(\eta(x) - \eta(x+1)\big) f(\eta)\nu_{\rho}(d\eta) - \big\langle -\mathscr{L}\sqrt{f}, \sqrt{f} \big\rangle_{\rho} \Big\}.$$

Above, we call f is a ν_{ρ} - density if $f \ge 0$ and $\int f d\nu_{\rho} = 1$, and for two local functions $f, g: \Omega \to \mathbb{R}$,

$$\langle f,g \rangle_{\rho} = \int f(\eta)g(\eta) \,\nu_{\rho}(d\eta).$$

Since ν_{ρ} is invariant for the SSEP,

$$\left\langle -\mathscr{L}\sqrt{f}, \sqrt{f} \right\rangle_{\rho} = \frac{1}{4} \sum_{x \in \mathbb{Z}} \int \left(\sqrt{f}(\eta^{x,x+1}) - \sqrt{f}(\eta) \right)^2 \nu_{\rho}(d\eta).$$

Making the change of variables $\eta \mapsto \eta^{x,x+1}$ and using the Cauchy-Schwarz inequality, for any A > 0, we bound the first term inside the above brace by

$$\begin{split} \frac{a_N K}{2nN^2} \sum_{x=0}^{nN-1} \int \left(\eta(x) - \eta(x+1)\right) \left(f(\eta) - f(\eta^{x,x+1})\right) \nu_\rho(d\eta) \\ \leq \frac{a_N K}{4nN^2} \sum_{x=0}^{nN-1} \left\{ A \int \left(\sqrt{f}(\eta^{x,x+1}) - \sqrt{f}(\eta)\right)^2 \nu_\rho(d\eta) + \frac{1}{A} \int \left(\sqrt{f}(\eta^{x,x+1}) + \sqrt{f}(\eta)\right)^2 \nu_\rho(d\eta) \right\} \\ & \leq \frac{a_N K A}{nN^2} \left\langle -\mathscr{L}\sqrt{f}, \sqrt{f} \right\rangle_\rho + \frac{a_N K}{NA}. \end{split}$$

Taking $A = nN^2/(a_N K)$, we finally bound the expression in (5.7) by $kK^2/(Nn)$. Since $k \leq TN$, the proof is concluded.

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